

# Ultraviolet Dynamics of Fermions and Gravity

Cold Quantum Coffee on June 26, 2018

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**Marc Schiffer**, Heidelberg University

with A. Eichhorn and S. Lippoldt arXiv 18xx.xxxxx

with A. Eichhorn, S. Lippoldt, J. M. Pawłowski and M. Reichert arXiv 1807.xxxxx



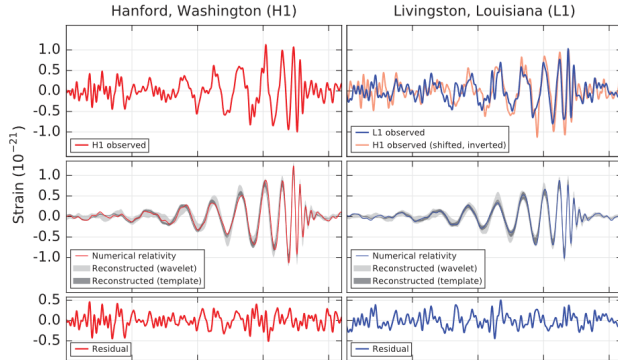
## Motivation for Quantum Gravity

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# Why quantum gravity?

## General Relativity:

- several high precision tests
- latest confirmation:  
Existence of black holes by LIGO  
collaboration



[LIGO Collaboration, 2016]

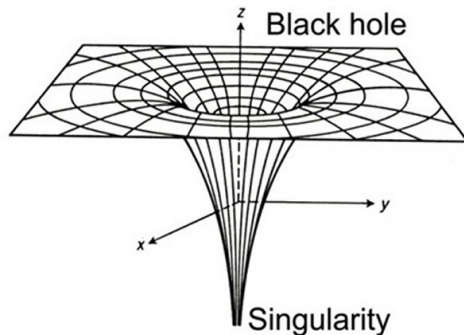
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Source: Northern Arizona University

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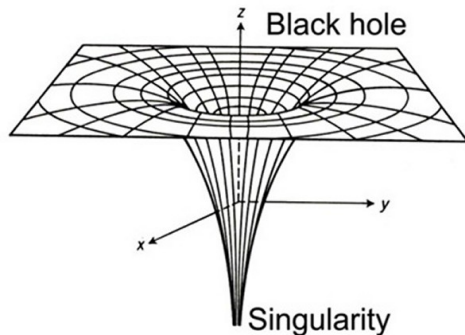
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- For  $E > M_{\text{PL}} = \sqrt{\frac{\hbar c}{G_{\text{N}}}}$ :  
QG effects are expected

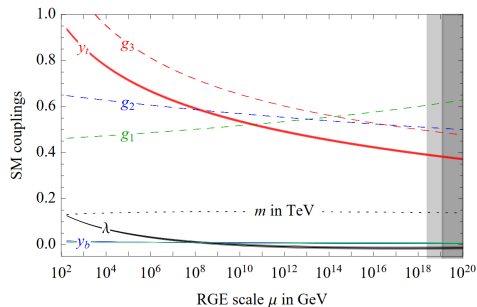


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# Why quantum gravity with matter?

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- well tested low energy model
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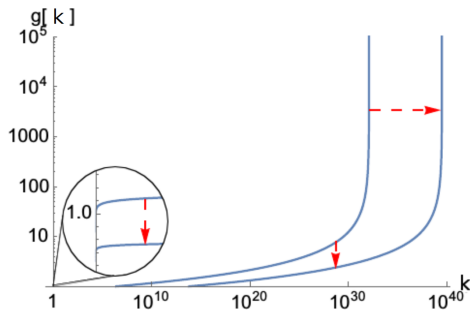


Image credits: Fleur Versteegen

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  - SM is only effective theory
  - breakdown beyond Planck scale
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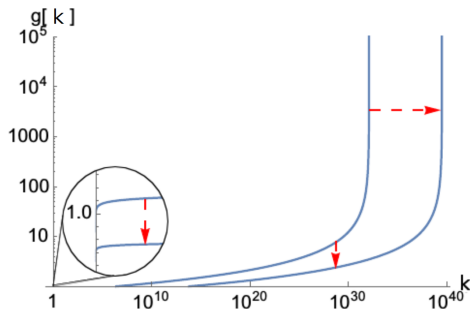


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compatibility with SM in IR provides test for quantum theory of gravity

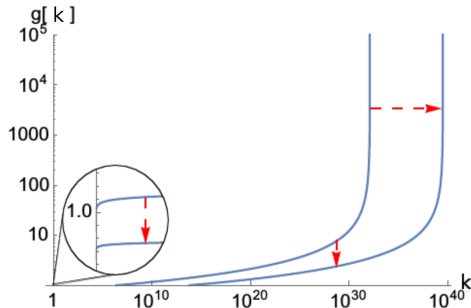


Image credits: Fleur Versteegen

Motivation for Quantum Gravity

Asymptotically Safe Quantum Gravity

Effective Universality for Gravity and Matter

Induced Couplings at UV Fixed Point

# Asymptotically Safe Quantum Gravity

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# How to quantize Gravity?

- $[G_N] = 2 - d$   
 $\Rightarrow$  GR is *perturbatively non-renormalizable* in  $d = 4$

[G. 't Hooft and M. J. G. Veltman, 1974]

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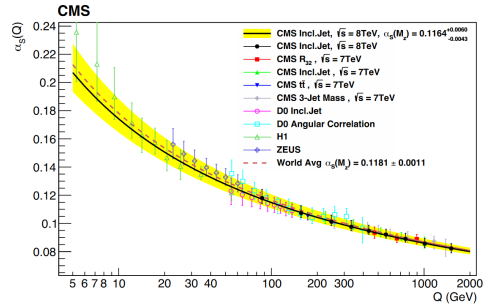
[M. H. Goroff and A. Sagnotti, 1986]

- Effective field theory approach: Loss of predictivity at  $M_{\text{PL}}$

[J. F. Donoghue and B. R. Holstein, 2015]

# Asymptotic safety

- Asymptotic freedom
  - ▶ all couplings vanish in the UV
  - ▶ perturbative renormalizability



CMS Collaboration, 2017

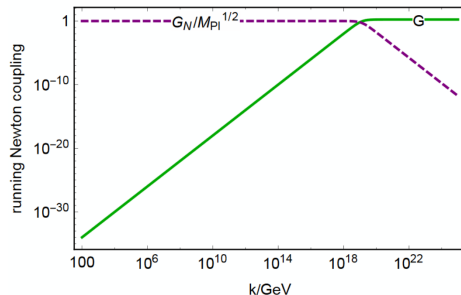
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- Asymptotic safety

[S. Weinberg, 1979]

- ▶ all dimensionless couplings enter a scale invariant regime
- ▶ interacting theory in the UV
- ▶ non-perturbative renormalizability



A. Eichhorn, 2017

# Search for Asymptotic safety

Study RG-flow of dimensionless couplings  $g_i = \bar{g}_i k^{-d_{\bar{g}_i}}$

$$\beta_{g_i}(\vec{g}) = k \partial_k g_i = \underbrace{-d_{\bar{g}_i} g_i}_{\text{dimensional}} + \underbrace{f_i(\vec{g})}_{\text{quantum}}$$

$\Rightarrow$  balancing of dimensional term with quantum correction can lead to AS

- Linearized  $\beta$ -functions

$$\beta_{g_i} = \beta_{g_i} \Big|_{\mathbf{g}=\mathbf{g}^*} + \sum_j \left( \frac{\partial \beta_{g_i}}{\partial g_j} \right) \Big|_{\mathbf{g}=\mathbf{g}^*} (g_j - g_j^*) + \mathcal{O}((g_j - g_j^*)^2)$$

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- Solution to linearized flow equations

$$g_i(k) = g_i^* + \sum_j c_j V_j^i \left( \frac{k}{k_0} \right)^{-\Theta_j} \quad \text{with} \quad -\text{eig}(M) = \Theta_i.$$

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$$\underline{\text{Re}(\Theta_i) < 0}$$

- irrelevant direction
- $g_i(k) \xrightarrow{k \rightarrow \infty} g_i^*$
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# Critical Exponents

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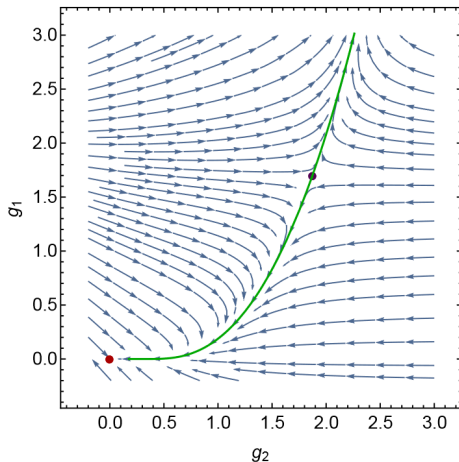
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Number of relevant directions determines **predictivity**

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[A. Eichhorn, 2017]

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# Tool: Functional Renormalization Group

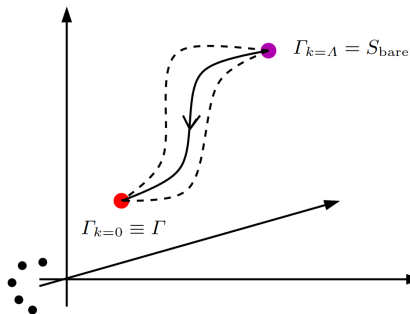
Non-Perturbative Renormalisation Group Equation [Wetterich, 1993], [Reuter, 1996]

$$k \partial_k \Gamma_k = \frac{1}{2} \text{STr} \left( \left( \Gamma_k^{(2)} + R_k \right)^{-1} k \partial_k R_k \right) = \frac{1}{2} \text{ (Feynman diagram: a circle with a cross inside)}$$

$\Gamma_k$  = scale dependent effective action

$R_k$  = IR regulator

- exact 1-loop equation
- extract  $\beta$ -functions via projection
- truncation needed  $\rightarrow$  not closed



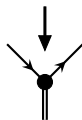
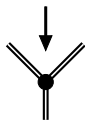
[H. Gies, 2006]

## Effective Universality for Gravity and Matter

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- Einstein-Hilbert gravity minimally coupled to fermions, i.e.

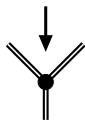
$$S = - \frac{1}{16 \pi G_{\text{N}}} \int d^4x \sqrt{g} (R - 2\Lambda) + \sum_{i=1}^{N_f} \int d^4x \sqrt{g} \bar{\psi}^i \not{\nabla} \psi^i$$



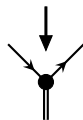
# Avatars of the Newton coupling

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$$\sim \sqrt{G_{3h}}$$



$$\sim \sqrt{G_{h\bar{\psi}\psi}}$$

- two different "avatars" of the Newton coupling

- Classically: Diffeomorphism invariance  
⇒ there should only be **one** Newton coupling

# Effective universality

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⇒ there should only be **one** Newton coupling
- On the quantum level
  - ▶  $[G_N] = -2 \Rightarrow$  2-loop universality is lost  
[\[Weinberg, 1995\]](#)
  - ▶ Gauge fixing, Regulator

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 $\Rightarrow$  there should only be **one** Newton coupling

- On the quantum level

- $[G_N] = -2 \Rightarrow$  2-loop universality is lost

[\[Weinberg, 1995\]](#)

- Gauge fixing, Regulator

$\Rightarrow$  set of identities (mSTI's) relates avatars

- Effective universality:**

Quantitative agreement of different avatars of the Newton coupling

[\[A. Eichhorn, P. Labus, J. M. Pawłowski and M. Reichert, 2018\]](#)

- Compare both avatars on the level of their  $\beta$ -functions at  $(\mu^*, \lambda_3^*, G_{3h} = G_{h\bar{\psi}\psi})$

$$\beta_{G_{3h}} = 2G - 3.4G^2 + 0.37G^3 + \mathcal{O}(G^4)$$

$$\beta_{G_{h\bar{\psi}\psi}} = 2G - 2.8G^2 + 0.42G^3 + \mathcal{O}(G^4)$$

Measure of effective universality

$$\varepsilon(G, \mu, \lambda_3) = \left| \frac{\Delta\beta_{G_i} - \Delta\beta_{G_j}}{\Delta\beta_{G_i} + \Delta\beta_{G_j}} \right|_{G_i=G_j}$$

with

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# Effective universality for fermions and gravity

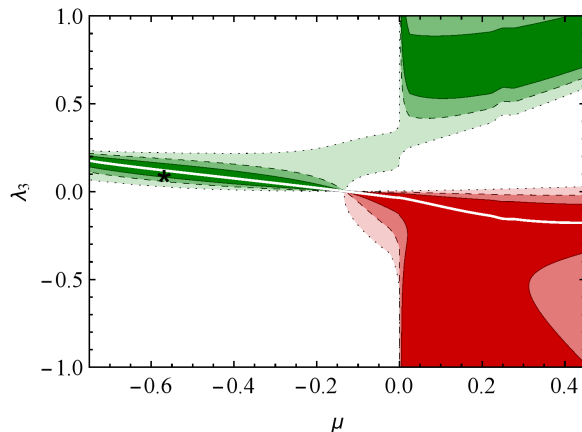
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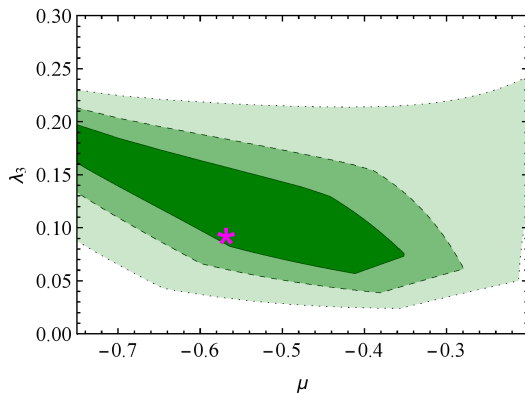


[A. Eichhorn, S. Lippoldt, MS, in preparation]

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- Compare fermions, scalars, vector fields and gravity

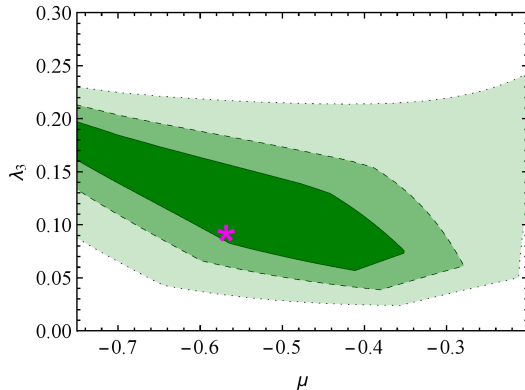
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Zoomed in plot region!

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## Implications

- Highly non-trivial cancellations
- *Not* expected for truncation artifact

→ strong hint for physical nature of asymptotically safe fixed point

## Induced Couplings at UV Fixed Point

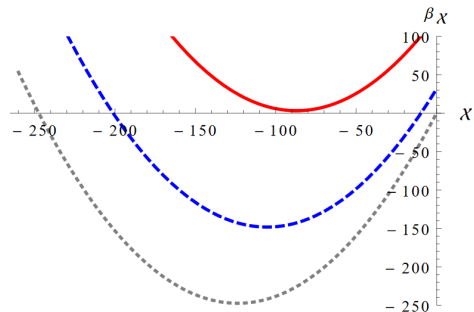
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[A. Eichhorn, A. Held and J. M. Pawłowski, 2016]

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- Explicit computations confirm expectation

[A. Eichhorn and H. Gies, 2011]

[A. Eichhorn and A. Held, 2017]

	scalars	fermions	gauge fields
scalars	GFP $\lambda_4 \phi^4$ sGFP $\rho (\partial_\mu \phi \partial^\mu \phi)^2$	GFP $y \phi \bar{\psi} \psi$ sGFP $\partial_\mu \phi \partial^\mu \phi \bar{\psi} \not{\nabla} \psi$	sGFP $\partial_\mu \phi \partial^\mu \phi F^2$
fermions		sGFP $(\bar{\psi} \gamma_\mu \psi)^2$	GFP $e \bar{\psi} \not{A} \psi$ sGFP $\bar{\psi} \not{\nabla} \psi F^2$
gauge fields			GFP <small>gluon self-interaction in <math>F^2</math></small> sGFP $w_2(F^2)^2 _{\text{Abelian}}$
complex scalars			GFP $e^2 A_\mu \phi^\dagger A^\mu \phi$ sGFP $\partial_\mu \phi^\dagger \partial^\mu \phi F^2$

[A. Eichhorn, A. Held and J. M. Pawłowski, 2016]

$$S = S_{\text{EH}} + S_{\text{Dirac}} \\ + \sigma \int d^4x \sqrt{g} R^{\mu\nu} \left( \bar{\psi}^i \gamma_\mu \overleftrightarrow{\nabla}_\nu \psi^i \right)$$

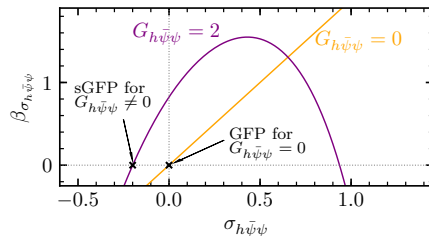
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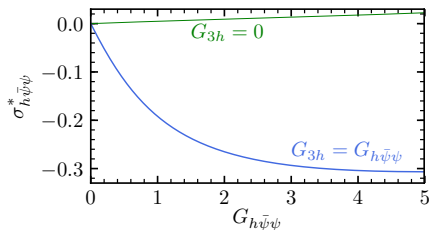
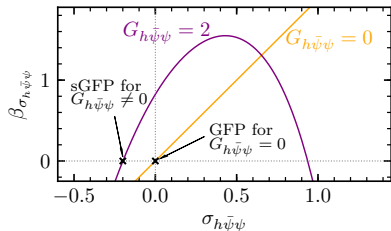
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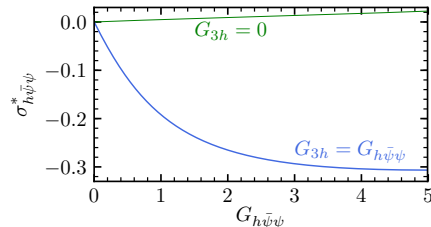
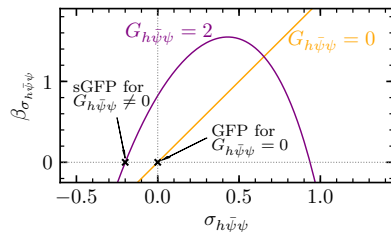
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Asymptotic safety passes non-minimal test for a UV complete theory of gravity and matter

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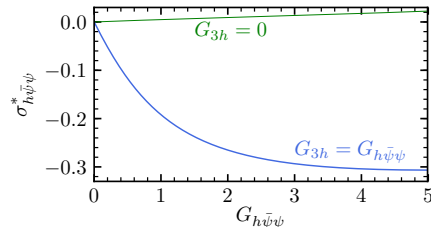
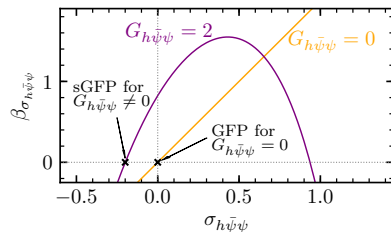
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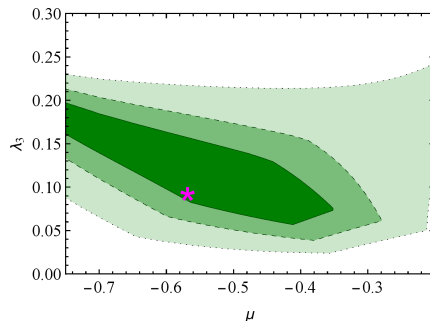
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Interacting nature of asymptotically safe fixed point percolates into chiral symmetry-protected matter sector

# Summary and Outlook

- Effective universality
  - ▶ hint for physical nature of asymptotically safe fixed point
  - ▶ guideline/ justification for future/past truncations
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- Explore restrictions of the gravitational parameter space under the inclusion of non-minimal couplings (weak gravity bound [\[A. Eichhorn and A. Held, 2017\]](#) )

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**Thank you for your attention!**

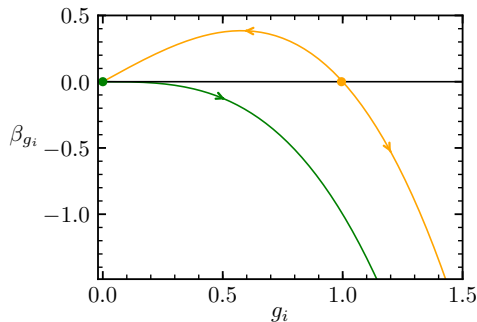
# Example for Asymptotic Safety

## Example: Yang-Mills Theory in $d = 4 + \epsilon$ at 1-loop

[M. E. Peskin, 1980], [M. Creutz, 1979], [H. Gies, 2003]

$$[\bar{g}] = -\frac{\epsilon}{2}$$

$$\begin{aligned}\beta_g &= \frac{\epsilon}{2} g - b_0 g^3 \\ &= g \left( \frac{\epsilon}{2} - b_0 g^2 \right).\end{aligned}$$



# Nontrivial cancellations

