

A microscopic approach to cosmic structure formation

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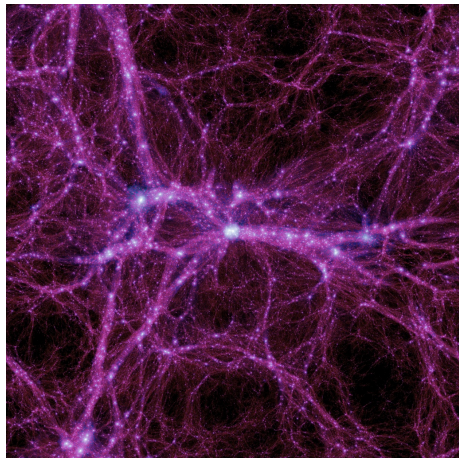
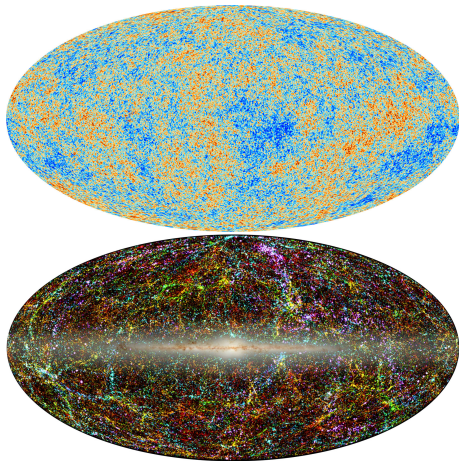
with

Felix Fabis, Daniel Berg, Elena Kozlikin, Robert Lilow,
Carsten Littek, Celia Viermann

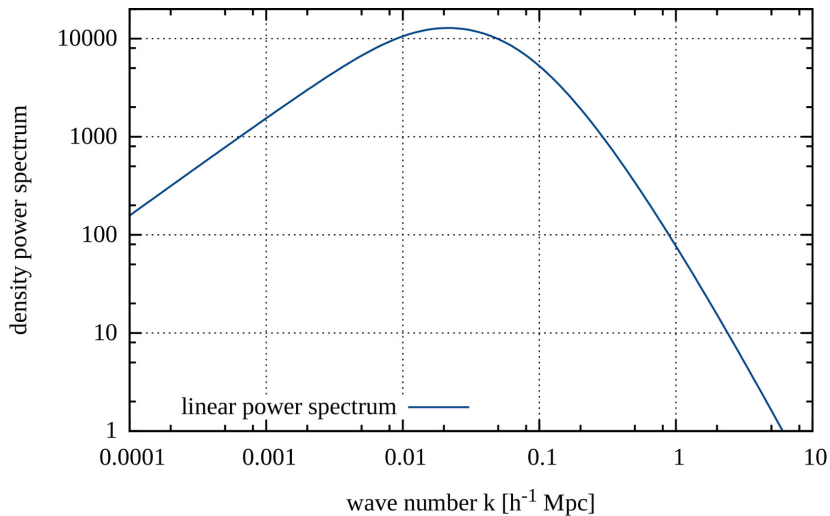
Heidelberg University, ZAH, ITA

Gravity on the Largest Scales, Heidelberg, October 26, 2015

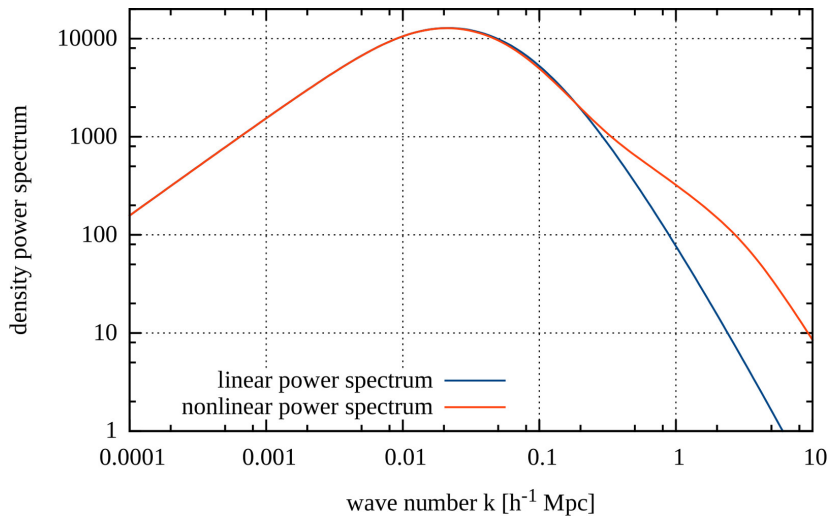
Problems in cosmic structure formation



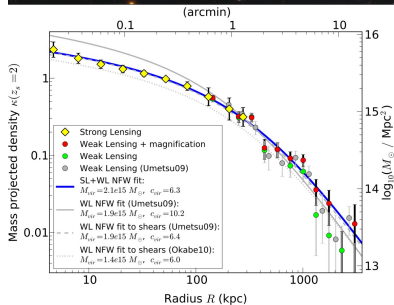
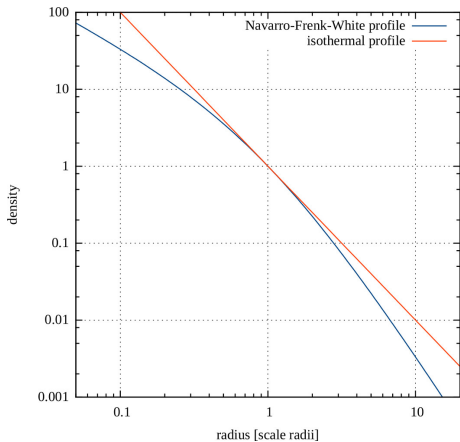
Problems in cosmic structure formation



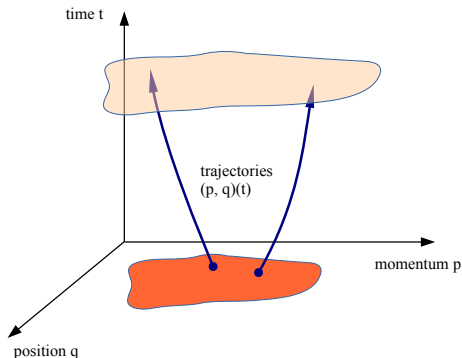
Problems in cosmic structure formation



Problems in cosmic structure formation



- Non-equilibrium statistics of N classical particle trajectories
- Describe particle ensemble by partition sum (generating functional) Z
- Derive statistical properties by functional derivatives



- Free functional:

$$Z_0[J, K] = \int d\Gamma^{(i)} \exp\left(i \int \langle J, \bar{x} \rangle dt\right)$$

with phase-space trajectories

$$\bar{x}(t) = \begin{pmatrix} \bar{q} \\ \bar{p} \end{pmatrix} = \underbrace{G_R(t, 0)x^{(i)}}_{\text{free motion}} - \underbrace{\int_0^t G_R(t, t')K(t')dt'}_{\text{interaction}}$$

- Initial phase-space distribution

$$d\Gamma^{(i)} = P(q, p)dqdp$$

- Include collective fields

$$Z_0[H, J, K] = \exp(iH\hat{\Phi}) \int d\Gamma^{(i)} \exp\left(i \int \langle J, \bar{x} \rangle dt\right)$$

with collective-field operator

$$\hat{\Phi} = \begin{pmatrix} \hat{\Phi}_\rho \\ \hat{\Phi}_B \end{pmatrix}, \quad \hat{\Phi}_\rho = \exp\left(-\vec{k} \cdot \frac{\delta}{\delta J}\right)$$

- Collective *response field* B quantifies the reaction of the ensemble to changes in particle positions,

$$\hat{\Phi}_B = \left(\vec{k} \cdot \frac{\delta}{\delta K}\right) \hat{\Phi}_\rho = \hat{b} \hat{\Phi}_\rho$$

- Include interaction

$$Z[H, J, K] = \exp(i\hat{S}_I) Z_0[H, J, K]$$

with interaction operator

$$\hat{S}_I = - \int d1 \hat{\Phi}_B(-1) v(1) \hat{\Phi}_\rho(1)$$

where v is the Fourier-transformed potential

- Perturbation theory:

$$e^{i\hat{S}} = 1 + i\hat{S} - \frac{1}{2}\hat{S}^2 + \dots$$

- n -point density correlator

$$\begin{aligned} G_{\rho\rho\dots\rho}(12\dots n) &= \hat{\Phi}_\rho \hat{\Phi}_\rho \dots \hat{\Phi}_\rho Z_0[J, K] \Big|_{J, K=0} \\ &= Z_0[L, 0] \end{aligned} \tag{1}$$

with the shift L given by

$$L = - \sum_{i=1}^n \begin{pmatrix} \vec{k}_i \\ 0 \end{pmatrix} \delta_{\text{D}}(t - t_i)$$

- k -th order interaction terms are of the form

$$\delta G_{\rho\rho}(12) = \int d1' b(-1') v(1') Z_0[L]$$

- 1 Choose initial phase-space measure

$$d\Gamma^{(i)} = P(q, p) d^{3N}q d^{3N}p$$

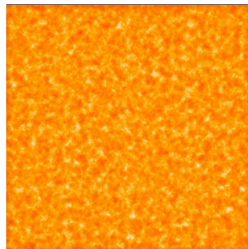
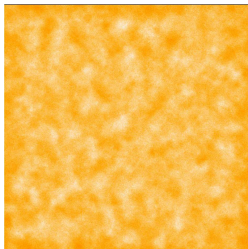
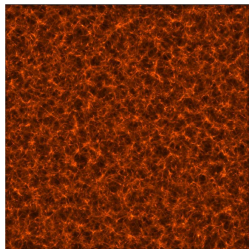
fully specified by initial power spectrum

- 2 Change time coordinate

$$t \rightarrow \tau = D_+(t) - D_+(t_i)$$

- 3 Adapt Green's function to expanding universe

Improved Zel'dovich trajectories



$$\vec{q}(\tau) = \vec{q}^{(i)} + \vec{p}^{(i)} \int_0^\tau \exp(h(\tau')) d\tau'$$

$$h(\tau) = g^{-1}(\tau) - 1, \quad g(\tau) = a^2 D_+ f E(a)$$

$$\hat{v}(k) = -\frac{3}{2k^2} \frac{a}{g^2} \propto a^{-2} \text{ (EdS)}$$

- In excellent approximation, the generating functional after application of arbitrary many density operators is

$$Z_0[L, 0] = V^{-N} \int dq \exp\left(-\frac{1}{2} L_p^\top C_{pp} L_p\right) \exp(i \langle L_q, q \rangle)$$

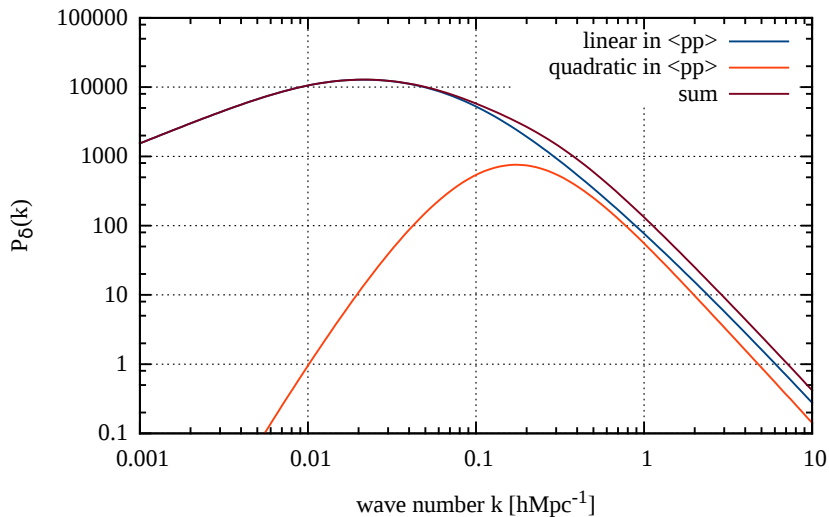
with the position and momentum shifts

$$L_q = - \sum_{i=1}^n \vec{k}_i, \quad L_p = - \sum_{i=1}^n g(t, t_i) \vec{k}_i$$

- C_{pp} is the momentum correlation matrix

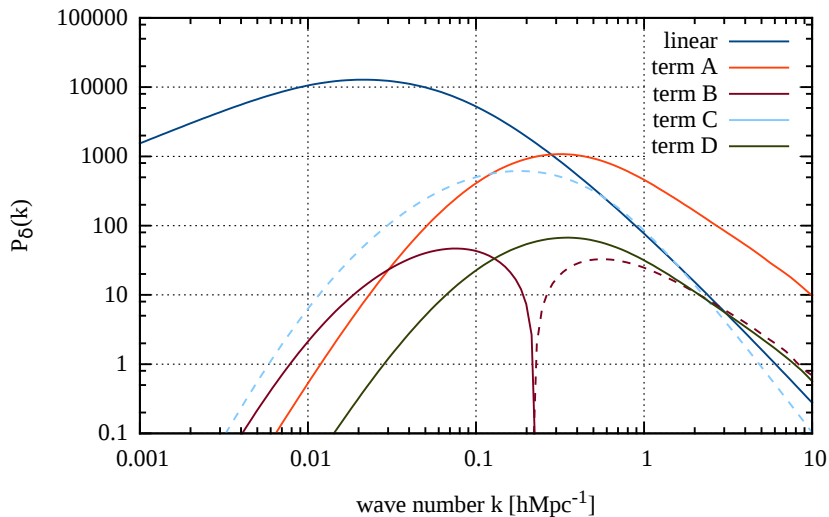
$$C_{pp} = \int \frac{d^3 k}{(2\pi)^3} \frac{\vec{k} \otimes \vec{k}}{k^4} P_\delta(k) e^{i\vec{k} \cdot \vec{x}}$$

Density power spectrum



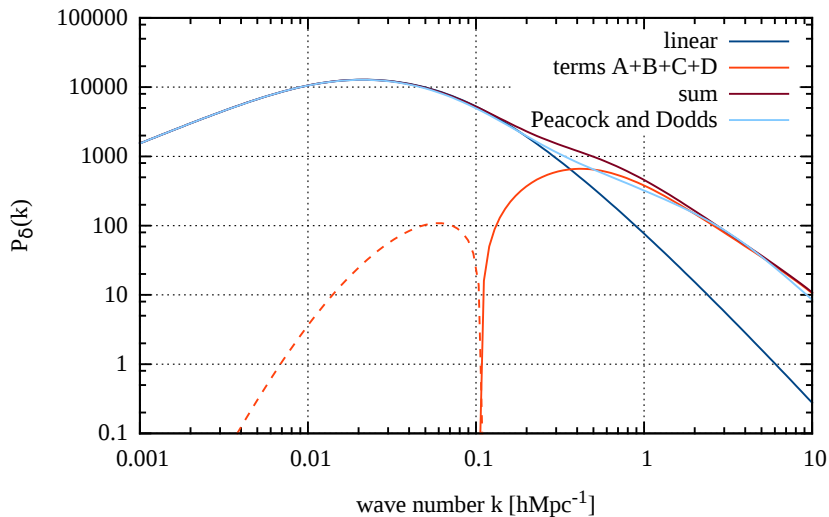
Improved Zel'dovich trajectories, no interaction

Density power spectrum



Improved Zel'dovich trajectories plus first-order interaction

Density power spectrum



Improved Zel'dovich trajectories plus first-order interaction

- 1 Non-equilibrium statistical field theory for dark-matter particles set up
- 2 Hamiltonian equations of motion, simple Green's function
- 3 Expansion parameter is deviation from unperturbed (improved Zel'dovich) trajectories
- 4 n -point statistics for collective fields obtained by functional derivatives
- 5 First-order perturbation theory reproduces numerical results already quite well

Further applications

