

# Slow-roll Inflation in the era of Euclid

**Sébastien Clesse**

S.C., C. Ringeval, *work in progress*

RWTH - Aachen University  
Institute of Theoretical Particle Physics and Cosmology (TKK)

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Gravity at the Largest Scales  
2015, October 26-28, Heidelberg

**RWTH**AACHEN  
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## Euclid Forecasts

Slow-roll  
single-field  
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Summary and  
perspectives $\Lambda$ -CDM

Extended models and parameterizations

Fisher Forecasts



MCMC forecasts

Planck mock data



Planck 2013, Planck 2015

In the context of inflation:

Scalar spectral index  $n_s$ , running  $\alpha_s$ , tensor-to-scalar ratio  $r$ Hubble-flow parameters  $\epsilon_1, \epsilon_2, \epsilon_3$ , slow-roll parameters, etc

Single field models - ASPIC library

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Inflaton field  $\phi$  evolves slowly along its potential  $V(\phi)$

### Hubble-flow functions (slow-roll parameters)

$$\epsilon_0 \equiv \frac{H_{\text{ini}}}{H} \quad \text{and} \quad \epsilon_{i+1} \equiv \frac{d \ln |\epsilon_i|}{dN} \quad (N, \text{ e-fold number})$$

$$\text{In slow-roll : } \epsilon_1 \simeq \frac{M_{\text{pl}}^2}{2} \left( \frac{V_{,\phi}}{V} \right)^2 \quad \epsilon_2 \simeq 2M_{\text{pl}}^2 \left[ \left( \frac{V_{,\phi}}{V} \right)^2 - \frac{V_{,\phi\phi}}{V} \right]$$

$$\epsilon_2 \epsilon_3 \simeq 2M_{\text{pl}}^4 \left[ \frac{V_{,\phi\phi\phi} V_{,\phi}}{V^2} - 3 \frac{V_{,\phi\phi}}{V} \left( \frac{V_{,\phi}}{V} \right)^2 + 2 \left( \frac{V_{,\phi}}{V} \right)^4 \right]$$

### Observable predictions

- Scalar power spectrum amplitude:  $A_s \equiv \mathcal{P}_\zeta(k_p) \simeq \frac{H^2}{8\pi^2 M_{\text{pl}}^2 \epsilon_1}$
- Scalar spectral index (*at second order*):  
 $n_s = 1 - 2\epsilon_1 - \epsilon_2 + (4C + 2)\epsilon_1^2 + (2C - 1)\epsilon_1\epsilon_2 + C\epsilon_2^2 - C\epsilon_2\epsilon_3$
- Running :  $\alpha_s = 4\epsilon_1^2 + 2\epsilon_1\epsilon_2 + \epsilon_2^2 - \epsilon_2\epsilon_3$
- Tensor to scalar ratio:  $r \simeq 16\epsilon_1$ , Tensor spectral index:  $n_t \simeq -2\epsilon_1$

evaluated at  $t_*$  when the pivot scale  $k_* = 0.05 \text{Mpc}^{-1}$  exits the Hubble radius ( $k_* = a(t_*)H(t_*)$ ) and  $C \approx -0.73$

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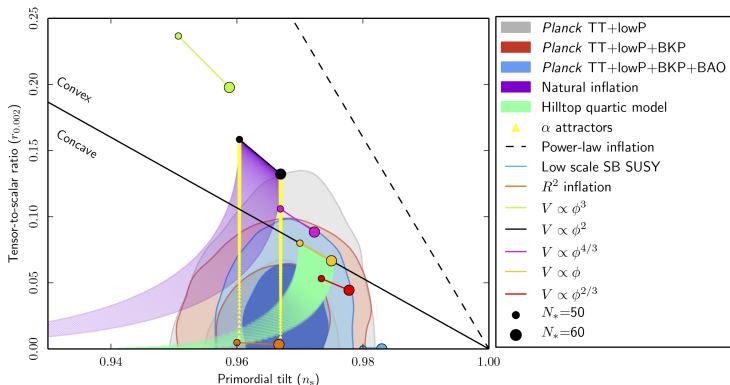
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## Constraints on $n_s, r, \alpha_s \dots$ (Planck 2015, 1502.02111)



PRE-2020 experiments (Errard, Feeney, Peiris, Jaffe, 1509.0677):

- BICEP3  $\times$  Keck Array  $\times$  Planck  $\Rightarrow \sigma(r) = 0.013$
- PIPER  $\times$  Planck  $\Rightarrow \sigma(r) = 8.3 \times 10^{-3}$
- SPT-3G  $\times$  Planck  $\Rightarrow \sigma(r) = 4.8 \times 10^{-3}$

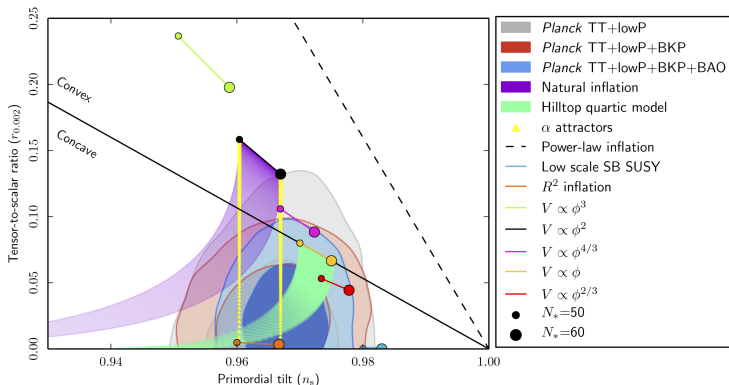
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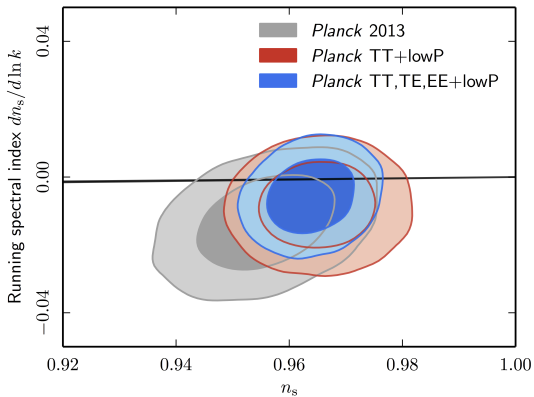
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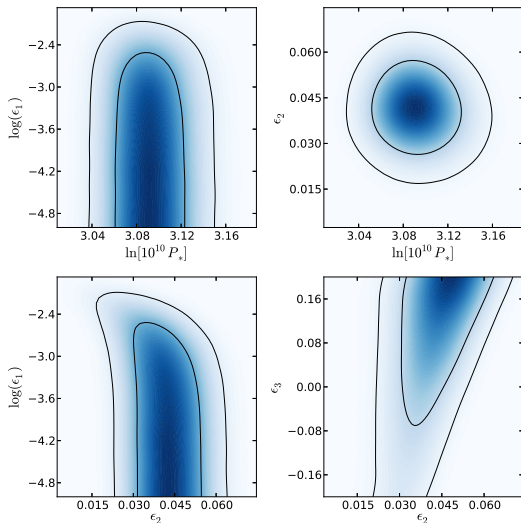
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perspectives...or on  $\epsilon_1, \epsilon_2, \epsilon_3$  (Martin, Ringeval, Trota, Vennin, 1312.3529)



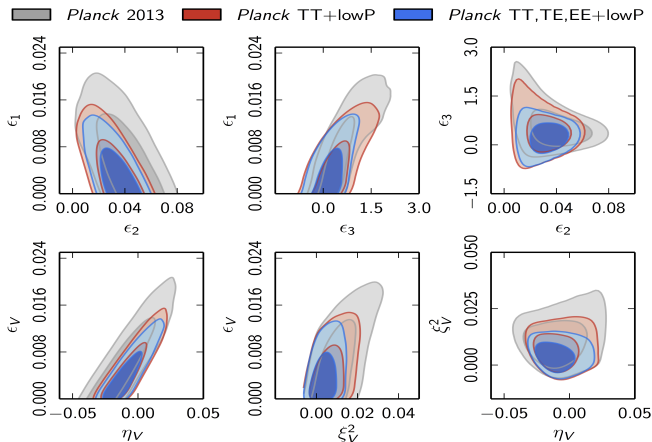
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...and on single field models (Martin, Ringeval, Trota, Vennin, 1312.3529)

Euclid Forecasts

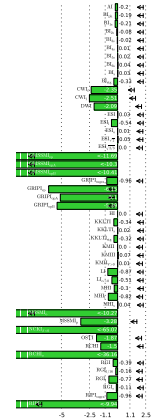
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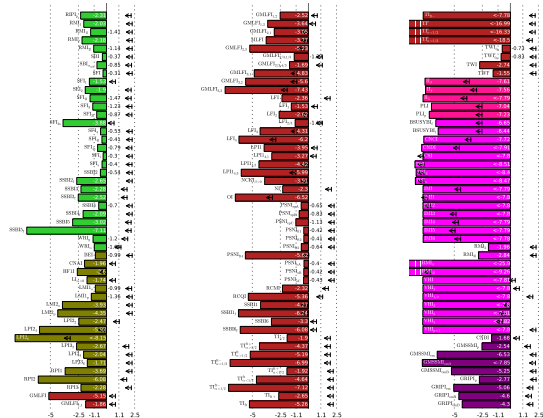
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Schwarzschild-Terrero-Escalante Classification:  
■ 0.0-0.2   ■ 0.2-0.5   ■ 0.5-1.0   ■ 1.0-1.5   ■ 1.5-2.0

Bayesian Evidences  $\ln(\mathcal{E}/\mathcal{E}_{\text{HI}})$  and  $\ln(\mathcal{L}_{\text{max}}/\mathcal{E}_{\text{HI}})$



J.Martin, C.Ringeval, R.Trota, V.Vennin  
 ASPIC project

Displayed Evidences: 193

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Observables:

- **3D Galaxy Power spectrum**  $P_g(k, \mu, z)$  Probes the matter power spectrum  $P_\delta(k)$  - biased
- **Galaxy Angular Power Spectrum** (cross-correlation between redshift bins  $i$  and  $j$ ) :  $C_{l,ij}^{(g)}$ .  
Probes the matter power spectrum  $P_\delta(k)$  - biased
- **Shear Angular Power Spectrum:**  $C_{l,ij}^{(s)}$   
Probes the metric perturbation  $\Psi$  power spectrum
- **Cross-correlation galaxy-shear spectra:**  $C_{l,ij}^{(g,s)}$
- **Cluster Mass Function** :  $N_{ij}$ , number of clusters in redshift bin  $i$  and mass bin  $j$   
Probes the matter power spectrum  $P_\delta(k)$

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- **Huang, Verde, Vernizzi, 1201.5955**

MCMC, mock Planck, 3D galaxy power spectrum

$$\Delta n_s = 0.002, \Delta \alpha_s = 0.003$$

- **Audren, Lesgourgues, Bird, Haehnelt, Viel, 1210.2194**

MCMC, mock Planck, galaxy P.S., shear angular P.S.

$$\text{Gal.: } \Delta n_s = 2.8/0.37/3.4 \times 10^{-3}$$

$$(k_{max} = 0.1/0.6, \text{ no uncorr. err. / incl. uncorr. err.})$$

$$\text{Shear: } \Delta n_s = 2.4/2.7 \times 10^{-3} \text{ (no uncorr. err. / incl. uncorr. err.)}$$

- **Euclid Theory Group, Euclid Living Rev., 1206.1225**

MCMC, mock Planck, Fisher matrix: Euclid only

- **Basse, Hamann, Hannetas, Wong, 1209.1043 and 1409.3469**

MCMC, mock Planck, galaxy and shear angular P.S. + clusters

Data	$\sigma(\log A_s)$	$\sigma(n_s)$	$\sigma(dn_s/d \ln k)$	$\sigma(r)$
c	0.011	0.0052	0.0074	0.028
cs	0.0091	0.0030	0.0030	0.027
cg	0.0046	0.0035	0.0048	0.027
ccl	0.0068	0.0034	0.0064	0.026
escl	0.0066	0.0028	0.0029	0.026
csgxcl	0.0032	0.0025	0.0017	0.026

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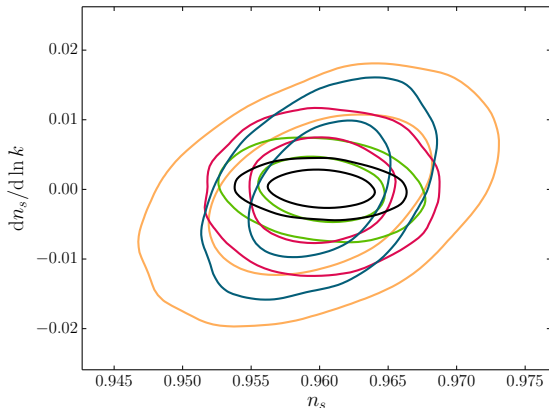
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Forecasts on  $n_s, \alpha_s, r$  (Basse, Hamann, Hannetas, Wong, 1409.3469)  
CMB, CMB+shear, CMB+Galaxy, CMB+Clusters,  
CMB+shear+galaxy+clusters+cross-correlation



Adding clusters increases the lever arm to constrain the running

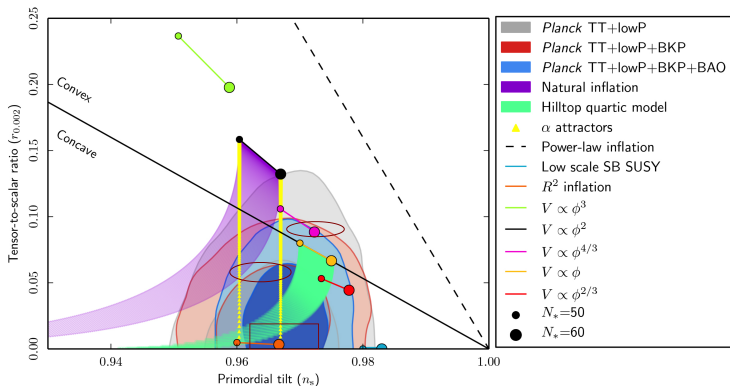
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Our forecasts:

- Modified version of the modified CAMB/COSMOMC version from **Basse, Hamann, Hannestas, Wong, 1409.3469**
- Modified version CLASS/MONTEPYTHON from **Audren, Lesgourgues, Bird, Haehnelt, Viel, 1210.2194**
- REAL Planck data instead of mock CMB data
- Forecasts for Hubble-flow slow-roll parameters

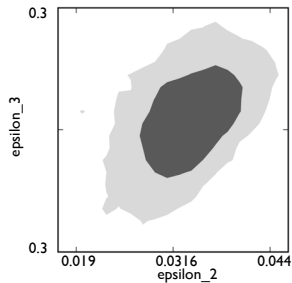
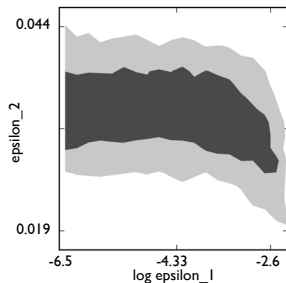
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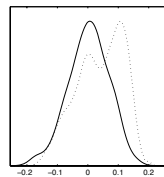
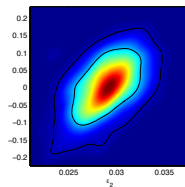
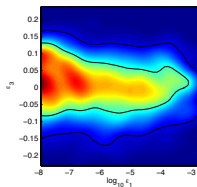
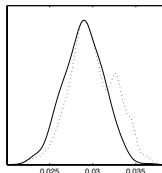
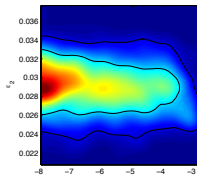
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Forecasts on  $\epsilon_1, \epsilon_2, \epsilon_3$  (S.C. Ringeval, *in preparation*)

Shear + galaxy + cross. corr. + clusters



4 times better reconstruction of  $\epsilon_2$  (!!! More statistics needed!!!)

Euclid could distinguish between concave vs. convex potentials

i.e. large vs. small field models

Forecasts for PRISM-like experiment, Martin, Ringeval, Vennin, 1407.4034

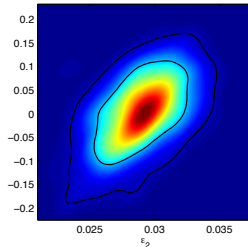
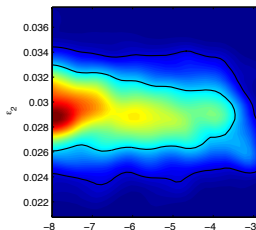
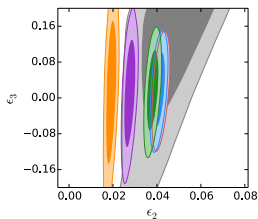
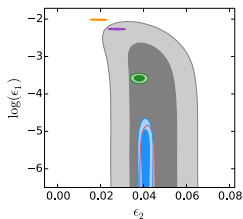
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## Euclid: implications for single-field inflation

- Constraints on  $\alpha_s$  optimally improved by a factor  $\sim 4$
- Constraints on  $n_s$  and  $\alpha_s$  optimally improved by a factor  $\sim 2$
- Reconstruction of  $\epsilon_2$  (linked to  $V_{,\phi\phi}$ )
- Distinction between concave and convex potentials at  $\gtrsim 2\sigma$

## Some perspectives:

- **More statistics needed!**
- **Code validation!**
- *"GI alignment too difficult to be taken into account"*
- Bias: shape and marginalization
- Forecasts for  $\sim 200$  models of the ASPIC library
- Detecting features in the scalar PS - non-gaussianity
- Joint forecasts with 21cm experiments (SKA-2020)
- Joint forecasts with future CMB experiments (CoRE+)
- Degeneracies with DE/modified gravity parameters

A lot of work  $\Rightarrow$  please join the Euclid-TH WP4 (initial conditions) !

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**Thank you for your attention**

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Specifications: as in Basse, Hamann, Hannestas, Wong, 1409.3469

- 2 redshift bins for cosmic shear
- 11 redshift bins from  $z = 0$  to  $z = 3$ , with  $\Delta z = 0.15$  for galaxy
- 10 redshift bins / 10 mass bins for clusters from  $z = 0$  to  $z = 2$
- Cluster mass range:  $[M_{\text{th}}, 10^{16} M_{\odot}]$
- Field of view:  $f_{\text{sky}} = 0.37$
- 30 galaxies per arcmin<sup>2</sup>
- Photometric uncertainty: simple Gaussian with  $\sigma(z) = 0.03(1 + z)$
- Known bias  $b = 1$  (optimistic!)
- Subdominant redshift space distortions
- Non-linearity cut-off:  $l_{\text{max}} = 2000$

As in Audren, Lesgourgues, Bird, Haehnelt, Viel, 1210.2194

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- **Galaxy Angular Power Spectrum** (cross-correlation between redshift bins  $i$  and  $j$ ) :

$$C_{l,ij}^{(g)} = 4\pi \int d \ln k S_{l,i}^{(g)}(k) S_{l,j}^{(g)}(k) \mathcal{P}_\zeta(k) \quad (1)$$

Source function :  $S_{l,i}^{(g)} = \int d\chi j_l(k\chi) \mathcal{W}_i^{(g)}(\chi) T_\delta(k, \eta_0 - \chi)$

Window function:  $\mathcal{W}_i^{(g)}(\chi) = b(k, \chi) \frac{H(z) dn/dz(z)}{n_i}$

with  $n_i \equiv \int_{\Delta z_i} dz' dn/dz'(z')$  and  $dn/dz \propto z^2 \exp(-(z/z_0)^\beta)$

Parameters:  $\beta = 1$ ,  $z_0 = 0.3$

Noise :  $C_{ij}^{\text{noise}} = \delta_{ij} n_i^{-2}$

- **Shear Angular Power Spectrum:**

$$C_{l,ij}^{(s)} = 4\pi \int d \ln k S_{l,i}^{(s)}(k) S_{l,j}^{(s)}(k) \mathcal{P}_\zeta(k) \quad (2)$$

Source function :

$$S_{l,i}^{(s)} = -2\sqrt{\frac{l(l^2-1)(l+2)}{4}} \int d\chi j_l(k\chi) \mathcal{W}_i^{(s)}(\chi) T_\Psi(k, \eta_0 - \chi)$$

Window function:  $\mathcal{W}_i^{(s)}(\chi) = \int_\chi^\infty dx' \frac{x-x'}{xx'} \frac{H(z) dn/dz(z)}{n_i}$

Noise :  $C_{ij}^{\text{noise}} = \delta_{ij} \langle \gamma^2 \rangle^{1/2} n_i^{-2}$  with  $\langle \gamma^2 \rangle^{1/2} \sim 0.35$



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- **Galaxy Angular Power Spectrum** (cross-correlation between redshift bins  $i$  and  $j$ ) :

$$C_{l,ij}^{(g)} = 4\pi \int d \ln k S_{l,i}^{(g)}(k) S_{l,j}^{(g)}(k) \mathcal{P}_{\zeta}(k) \quad (1)$$

Source function :  $S_{l,i}^{(g)} = \int d\chi j_l(k\chi) \mathcal{W}_i^{(g)}(\chi) T_{\delta}(k, \eta_0 - \chi)$

Window function:  $\mathcal{W}_i^{(g)}(\chi) = b(k, \chi) \frac{H(z) dn/dz(z)}{n_i}$

with  $n_i \equiv \int_{\Delta z_i} dz' dn/dz'(z')$  and  $dn/dz \propto z^2 \exp(-(z/z_0)^\beta)$

Parameters:  $\beta = 1$ ,  $z_0 = 0.3$

Noise :  $C_{ij}^{\text{noise}} = \delta_{ij} n_i^{-2}$

- **Shear Angular Power Spectrum:**

$$C_{l,ij}^{(s)} = 4\pi \int d \ln k S_{l,i}^{(s)}(k) S_{l,j}^{(s)}(k) \mathcal{P}_{\zeta}(k) \quad (2)$$

Source function :

$$S_{l,i}^{(s)} = -2\sqrt{\frac{l(l^2-1)(l+2)}{4}} \int d\chi j_l(k\chi) \mathcal{W}_i^{(s)}(\chi) T_{\Psi}(k, \eta_0 - \chi)$$

Window function:  $\mathcal{W}_i^{(s)}(\chi) = \int_{\chi}^{\infty} d\chi' \frac{\chi - \chi'}{\chi \chi'} \frac{H(z) dn/dz(z)}{n_i}$

Noise :  $C_{ij}^{\text{noise}} = \delta_{ij} \langle \gamma^2 \rangle^{1/2} n_i^{-2}$  with  $\langle \gamma^2 \rangle^{1/2} \sim 0.35$

Euclid Forecasts

Slow-roll  
single-field  
inflation

After Planck...

Euclid  
Observables

Euclid forecasts

Summary and  
perspectives

- Cluster Mass Function :

$N_{ij}$ , number of clusters in redshift bin  $i$  and mass bin  $j$

$$N_{ij} = \Delta\Omega \int_0^\infty dz \int_0^\infty dM dV \mathcal{W}_{ij}^c(M, z) \frac{dn_{\text{ST}}}{dM}(M, z) \quad (3)$$

$\Delta\Omega$  is the survey solid angle,  $dV$  is a comoving volume element  
 $\mathcal{W}_{ij}^c$  is the window function for redshift and mass bins

Sheth-Tormen mass function :

$$\frac{dn_{\text{ST}}}{dM}(M, z) = \sqrt{\frac{1}{2\pi}} A \frac{\rho_m}{M^2} [1 + (a\nu)^{-p}] \sqrt{a\nu} \exp\left(-\frac{a\nu}{2}\right) \frac{d \ln \nu}{d \ln M} \quad (4)$$

with  $\nu(M, z) \equiv \delta_c^2(M, z)/\sigma_m^2$  and  
 $\sigma_m^2(M, z) \equiv \int_0^\infty d \ln k |W(k, R)|^2 \mathcal{P}_m(k)$

$\delta_c(M, z)$  is the linear threshold density obtained from the model of spherical collapse

$W(k, R)$  is the Fourier transform of the top-hat filter

Parameters:  $a = 0.707$ ,  $A = 0.32$ ,  $p = 0.3$

**Cosmology enters into the linear matter power spectrum  $\mathcal{P}_m(k)$**