

Evolution of linear perturbations in (Λ)LTB models

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Heidelberg, October 26, 2015

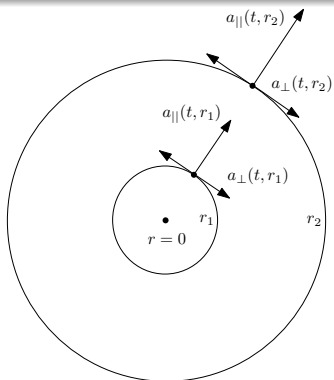
Outline

- 1 (Λ) LTB models in general
- 2 Linear perturbations
- 3 Evolution equations and features
- 4 Results

Simplest inhomogeneous generalisation of the standard FLRW metric based on an exact solution of Einstein's field equations

- 1 radially inhomogeneous
- 2 spherically symmetric about single "central" worldline
- 3 dust solution

$$T_{\mu\nu} = \rho(t, r)u_{\mu}u_{\nu}$$

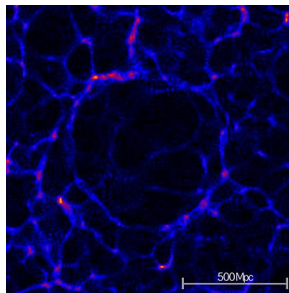


$$ds^2 = -dt^2 + \frac{a_{\parallel}^2(t, r)}{1 - \kappa(r)r^2}dr^2 + r^2 a_{\perp}^2(t, r)d\Omega^2$$

Motivation

Original idea:

- large Gpc scale spherical void (embedded into FLRW)
- accounts for DE effects on backward lightcone

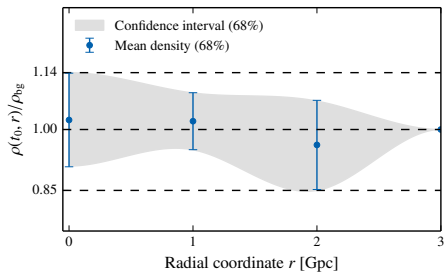


<http://www.thphys.uni-heidelberg.de/~cosmo/dokuwiki/doku.php/altmodels>

Problems:

- global $h_{\text{CMB}} < 0.4 \longleftrightarrow$ local $h_{\text{SNe}} = 0.738$ (Riess et al. (2011))
- no evidence for strong anisotropy away from our position (kSZ effect)

- cosmological constant Λ
- Λ + radial inhom. \rightarrow Λ LTB
- constrain deviations from homogeneity \Rightarrow test the Copernican Principle



from Redlich et al. (2014)

Provide **more observables** to

- 1 definitely distinguish and Λ LTB model from Λ CDM
- 2 confirm and strengthen findings on spherical voids on a broad scientific basis

investigate linear structure growth in (Λ)LTB models

Linear structure growth in Λ LTB: structures evolve **position-dependent** and **anisotropic**

spherically symmetry at background level:

expansion into **spherical harmonics** $Y^{(\ell m)}(\theta, \phi)$ and cov. derivatives
(Clarkson et al. (2010), Gundlach & Martín-García (2000), Gerlach & Sengupta (1978))

$$\phi = \sum_{(\ell m)} \phi^{(\ell m)} Y^{(\ell m)}$$

\Rightarrow Scalar-Vector-Tensor
expressions on \mathcal{S}^2

$$\phi_a = \sum_{(\ell m)} \phi^{(\ell m)} Y_a^{(\ell m)} + \bar{\phi}^{(\ell m)} \bar{Y}_a^{(\ell m)}$$

Sets of gauge invariant
quantities can be found

$$\phi_{ab} = \sum_{(\ell m)} \phi^{(\ell m)} Y_{ab}^{(\ell m)} + \bar{\phi}^{(\ell m)} \bar{Y}_{ab}^{(\ell m)}$$

different from FLRW gauge invariant quantities $\{\Psi, \Phi, V_i, h_{ij}\}$!

metric perturbations - Clarkson et al. (2010)

Abstract set of gauge invariant metric perturbations:

$$\{\chi^{(\ell m)}, \varphi^{(\ell m)}, \varsigma^{(\ell m)}, \eta^{(\ell m)}\}$$

$$\ddot{\chi} - \frac{\chi'' + C\chi'}{Z^2} + 3H_{\parallel}\dot{\chi} - \left[A - \frac{(\ell-1)(\ell+2)}{r^2 a_{\perp}^2} \right] \chi = \mathcal{S}_{\chi}(\varsigma, \varsigma', \varphi, \dot{\varphi})$$

$$\ddot{\varphi} + 4H_{\perp}\dot{\varphi} - \left(\frac{2\kappa}{a_{\perp}^2} - \Lambda \right) \varphi = \mathcal{S}_{\varphi}(\varsigma, \chi, \dot{\chi}, \chi', \ell)$$

$$\dot{\varsigma} + 2H_{\parallel}\varsigma = \mathcal{S}_{\varsigma}(\chi')$$

$$\eta = 0$$

Similar for matter perturbations $\{\Delta^{(\ell m)}, w^{(\ell m)}, v^{(\ell m)}\}$

Dynamical coupling of gauge invariants !

main differences compared to FLRW:

- **dynamical coupling** during spacetime evolution
- **physical interpretation**: LTB gauge invariants \neq FLRW gauge invariants

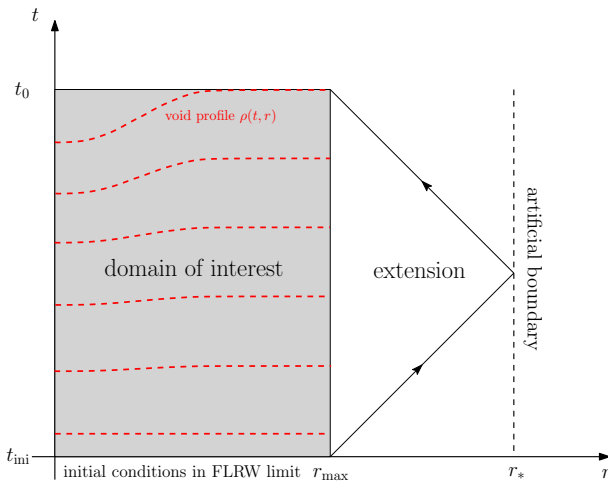


Numerical solution

- initial conditions
- boundary conditions
- discretization (Finite Elements)

extract observables

- light propagation:
- sensitive to combination of gauge-invariants \leftarrow Ricci- and Weyl focussing terms
- consequences: κ, γ



Dirichlet B.C.s

regularity at $r = 0$:

$$a^{(\ell m)}(t, 0) = 0$$

Gundlach et al. (2000)

artificial boundary

at $r = r_*$:

$$a^{(\ell m)}(t, r_*) = 0$$

(adapted from

February et al. (2014))

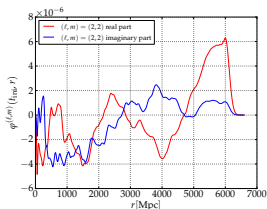
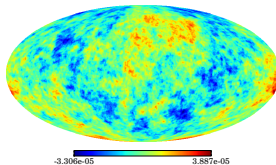
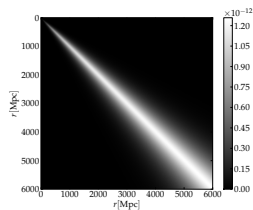
- initial universe at large redshift sufficiently homogeneous and isotropic
- start from initial scalar perturbations in spatially flat FLRW limit (Ψ)
- sample spherical harmonic coefficients from multivariate Gaussian distribution

covariance matrices:

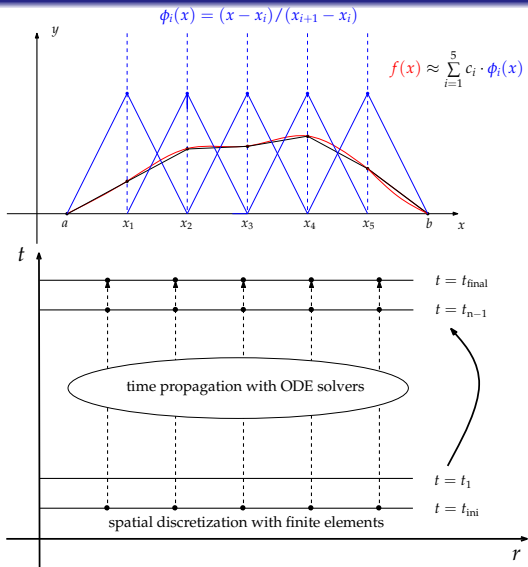
$$C_{\Psi}^{\ell}(r_i, r_j) = \frac{2}{\pi} \int_0^{\infty} dk k^2 P_{\Psi}(k) j_{\ell}(kr_i) j_{\ell}(kr_j)$$

- obtain initial profiles

$$\varphi^{(\ell m)}(t_{\text{ini}}, r) = -2\Psi^{(\ell m)}(t_{\text{ini}}, r)$$



- discretize domain of interest in radius (**Finite Elements**)
- expand variables in basis polynomials
- time evolution of coefficients (**method of lines**)



Dune

Distributed and Unified Numerics Environment

evolve each spherical harmonic coefficients $(\ell, m) \leftarrow$ parallelised

Overview on algorithm & angular Powerspectra

background model

- void density profile
- asymptotically embedded into FLRW model



evolve backwards to initial time/redshift ($z_{\text{ini}} \sim 100$)



- Evolve linear PDEs for each (ℓ, m) -mode forward in time

$$C^\ell(z) = \sum_{m=-\ell}^{\ell} \frac{|a^{(\ell m)}(t(z), r(z))|^2}{2\ell + 1}$$



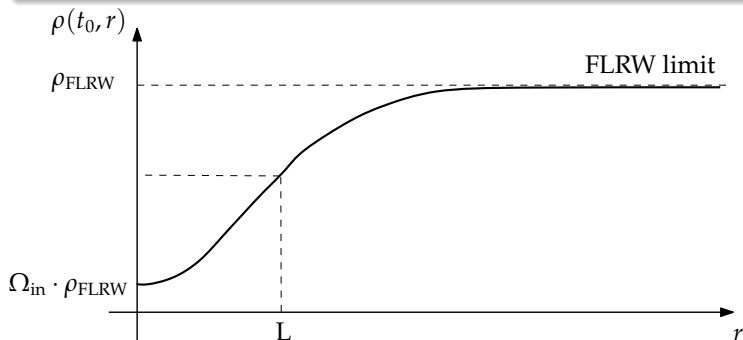
- **initial perturbations** (FLRW limit)

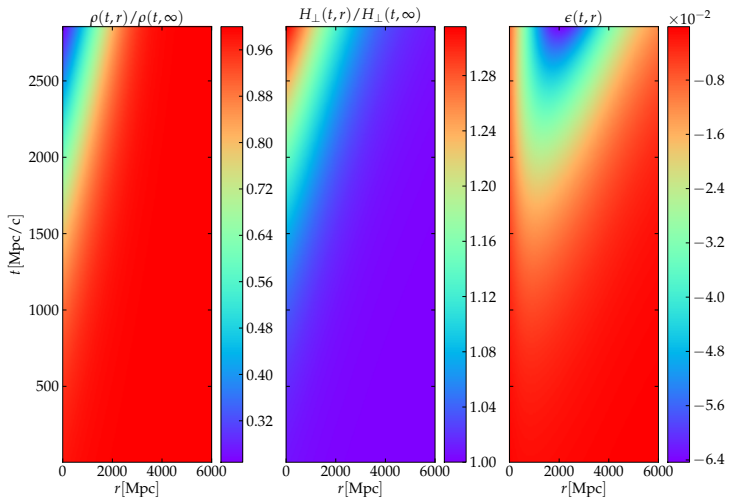
- $\varphi^{(\ell m)} = -2\Psi^{(\ell m)}$
- $\chi^{(\ell m)} = \zeta^{(\ell m)} = 0$

LTB model embedded into EdS

Gaussian void profile: $\rho(t_0, r) = \rho_{\text{bg}} \cdot f(r, \Omega_{\text{in}}, L)$

$$f(r, \Omega_{\text{in}}, L) = 1 + (\Omega_{\text{in}} - 1) \exp\left(-\frac{r^2}{2L^2}\right)$$

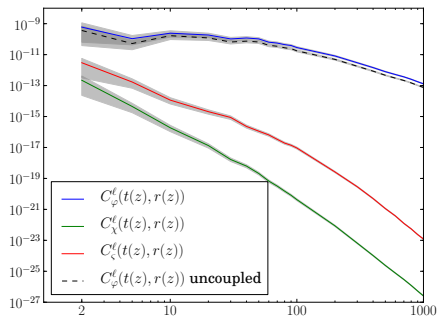




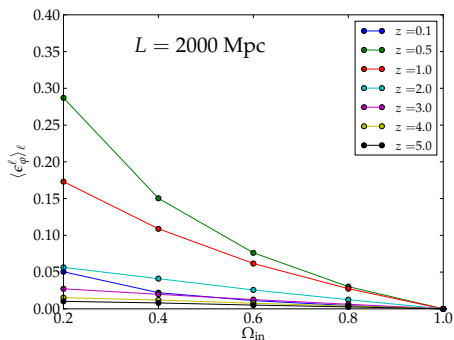
$$\epsilon(t, r) = \frac{H_{\parallel} - H_{\perp}}{H_{\parallel} + 2H_{\perp}}$$

average coupling strength

$$\epsilon_{\varphi}^{\ell}(z) = \left| \frac{\sqrt{C_{\varphi}^{\ell}(t(z), r(z))} - \sqrt{C_{\varphi, \text{uc}}^{\ell}(t(z), r(z))}}{\sqrt{C_{\varphi, \text{uc}}^{\ell}(t(z), r(z))}} \right|$$

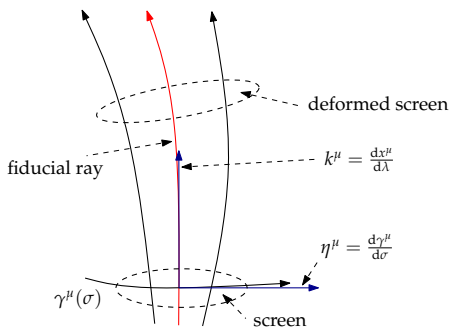


angular powerspectra at
 $z = 0.5$, $\Omega_{\text{in}} = 0.2$, $L = 2\text{Gpc}$



coupling strength

Application to best fit Λ LTB nearly done !



Geodesic deviation:

$$\nabla_k^2 \eta^\mu = R^\mu{}_{\nu\alpha\beta} k^\alpha k^\beta \eta^\nu$$

decomposition into Sachs basis, affine parameter :

$$\eta^\mu = \eta_1 n_1^\mu + \eta_2 n_2^\mu$$

$$\frac{d}{d\lambda} (\dots) = k^\mu \nabla_\mu (\dots)$$

$$\eta_a(\lambda) = D_{ab}(\lambda) \left. \frac{d\eta_b}{d\lambda} \right|_{\lambda=0}$$

$$\frac{d^2 D_{ab}}{d\lambda^2} = \mathcal{T}_{ac} D_{cb}$$

$$\text{with } \mathcal{T}_{ab} = \underbrace{-\frac{1}{2} R_{\alpha\beta} k^\alpha k^\beta \delta_{ab}}_{\text{Ricci focussing}} + \underbrace{C_{\alpha\beta\gamma\delta} n_a^\alpha k^\beta k^\gamma n_b^\delta}_{\text{Weyl focussing}}$$

Express Ricci and Weyl focussing in terms of LTB gauge invariants

Jacobi equation: $D_{ab} = D_A^{(\text{LTB})} \delta_{ab} + D_{ab}^{(1)}$ (central observer)

$$\frac{d^2 D_{ab}^{(1)}}{d\lambda^2} = -4\pi G\rho(t, r) (1+z)^2 D_{ab}^{(1)} + ra_{\perp}(t, r) \mathcal{T}_{ab}^{(1)} \left[\{\eta^{(\ell m)}, \chi^{(\ell m)}, \varsigma^{(\ell m)}, \varphi^{(\ell m)}\} \right]$$

Lensing amplification:

$$(A_{ab}) = \frac{(D_{ab})}{D_A^{(\text{LTB})}} = \frac{(D_{ab})}{ra_{\perp}} = \begin{pmatrix} 1 - \kappa & 0 \\ 0 & 1 - \kappa \end{pmatrix} - \begin{pmatrix} \gamma_1 & \gamma_2 \\ \gamma_2 & -\gamma_1 \end{pmatrix}.$$

Lensing quantities

$$\kappa = -\frac{1}{2} \frac{D_{11}^{(1)} + D_{22}^{(1)}}{ra_{\perp}}$$

$$\gamma_1 = -\frac{1}{2} \frac{D_{11}^{(1)} - D_{22}^{(1)}}{ra_{\perp}}$$

$$\gamma_2 = -\frac{1}{2} \frac{D_{12}^{(1)} + D_{21}^{(1)}}{ra_{\perp}}$$

- determine angular power spectra $C_{\kappa\kappa}^{\ell}$ and $C_{\gamma\gamma}^{\ell}$
- compare to observationally determined spectra

Conclusion

- **numerical solution** of evolution equations of gauge invariant perturbations in Λ LTB spacetimes (proof of concept)
- **coupling strength** significant in sufficiently deep and large voids (without Λ)
- application to best fit Λ LTB models nearly done
- observable predictions from **light propagation** (work in progress)