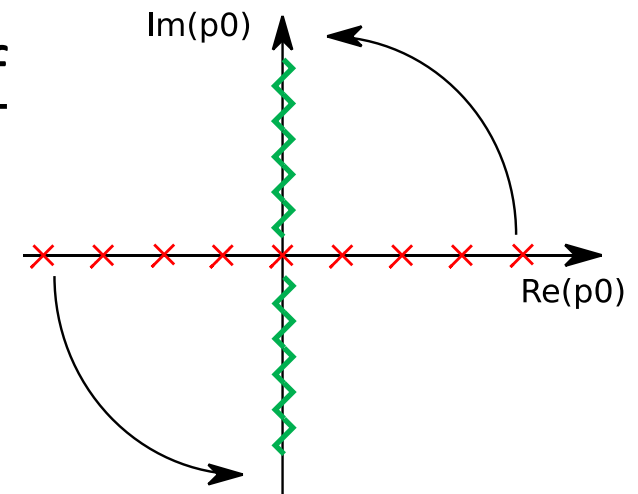


# Spectral functions from the FRG

666. WE-Heraeus-Seminar in Bad Honnef

Nicolas Wink



## Reconstructing the gluon

Anton K. Cyrol,<sup>1</sup> Jan M. Pawłowski,<sup>1,2</sup> Alexander Rothkopf,<sup>1,3</sup> and Nicolas Wink<sup>1</sup>

<sup>1</sup>*Institute for Theoretical Physics, Universität Heidelberg, Philosophenweg 12, D-69120 Germany*

<sup>2</sup>*ExtreMe Matter Institute EMMI, GSI, Planckstr. 1, D-64291 Darmstadt, Germany*

<sup>3</sup>*Faculty of Science and Technology, University of Stavanger, NO-4036 Stavanger, Norway*

We reconstruct the gluon spectral function in Landau gauge QCD from numerical data for the gluon propagator. The reconstruction relies on two novel ingredients: Firstly we derive analytically the low frequency asymptotics of the spectral function. Secondly we construct a functional basis from a careful consideration of the analytic properties of the gluon propagator in Landau gauge. This allows us to reliably capture the non-perturbative regime of the gluon spectrum. We also compare different reconstruction methods and discuss the respective systematic errors.

Uploaded today!

Base on:

Finite temperature calculation

Pawłowski, Strodthoff, NW arxiv:1711.07444

Gluon reconstruction

Cyrol, Rothkopf, Pawłowski, NW arxiv:1804.today

## Spectral functions in QCD

What is a spectral function?

Defined via the retarded propagator

$$\rho(\omega) = 2 \operatorname{Im} G_R(\omega)$$

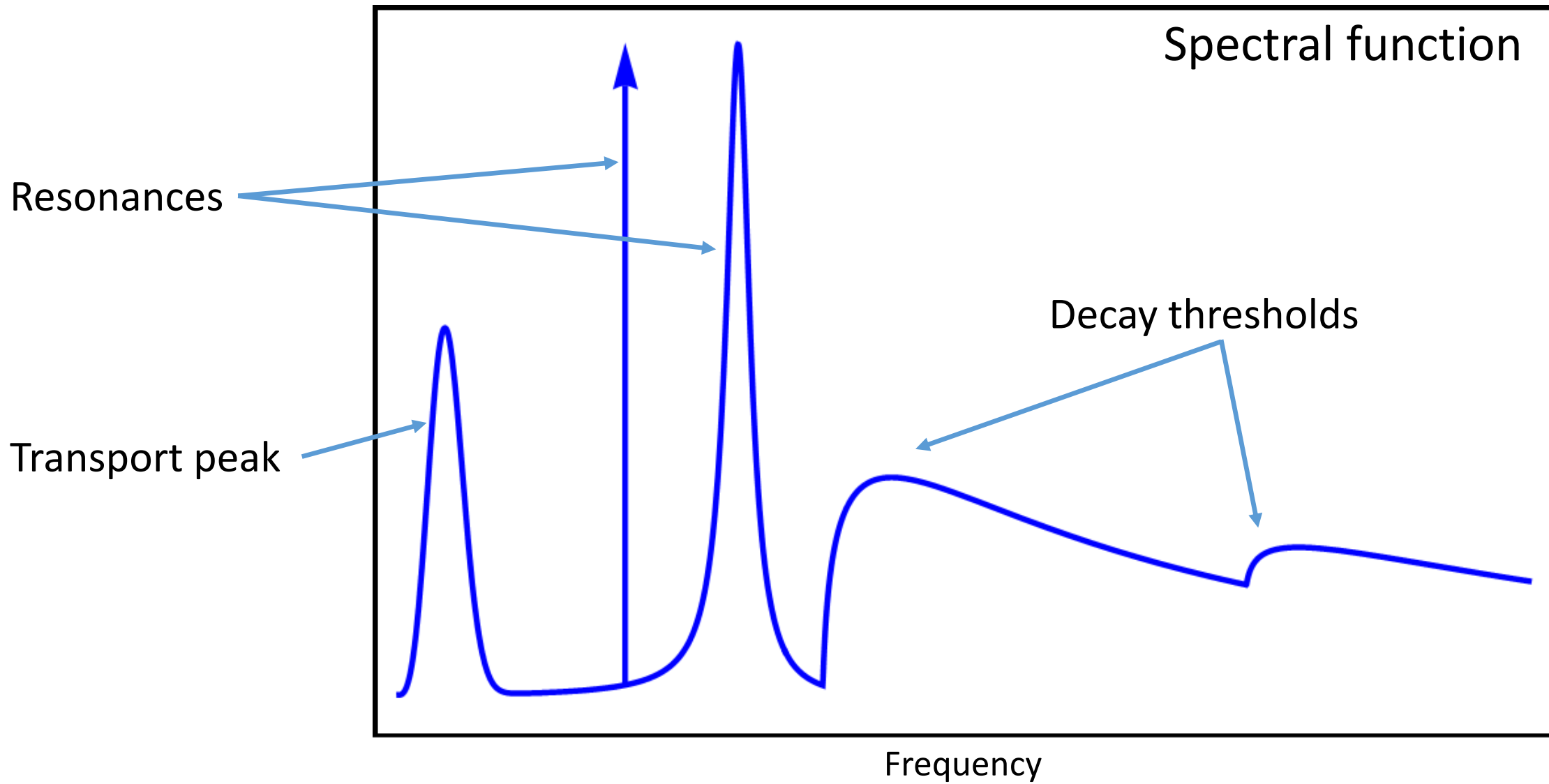
Encodes the spectrum of the theory

Why calculate spectral functions?

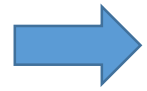
Encodes the spectrum of the theory

“Easy” to extract other observables

## Spectral functions in QCD

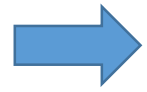


## Spectral functions in QCD

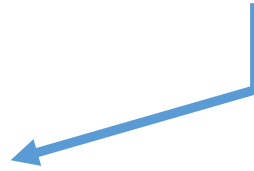


Requires non-perturbative correlation functions in Minkowski space-time

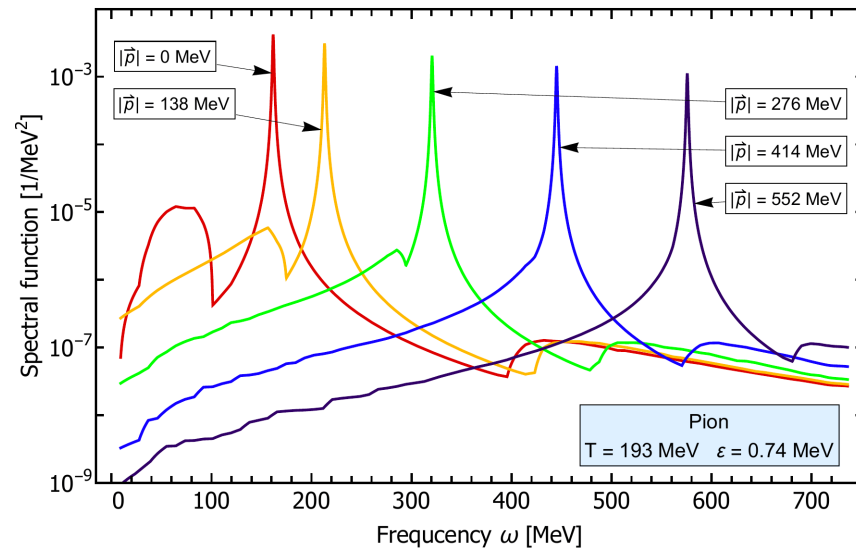
# Spectral functions in QCD



Requires non-perturbative correlation functions in Minkowski space-time

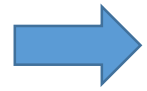


Direct calculation

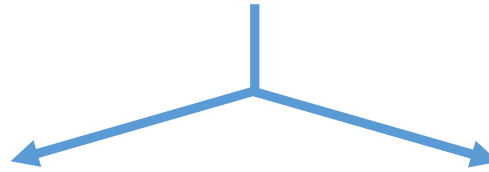


Pawlowski, Strodthoff, NW, arxiv:1711.07444

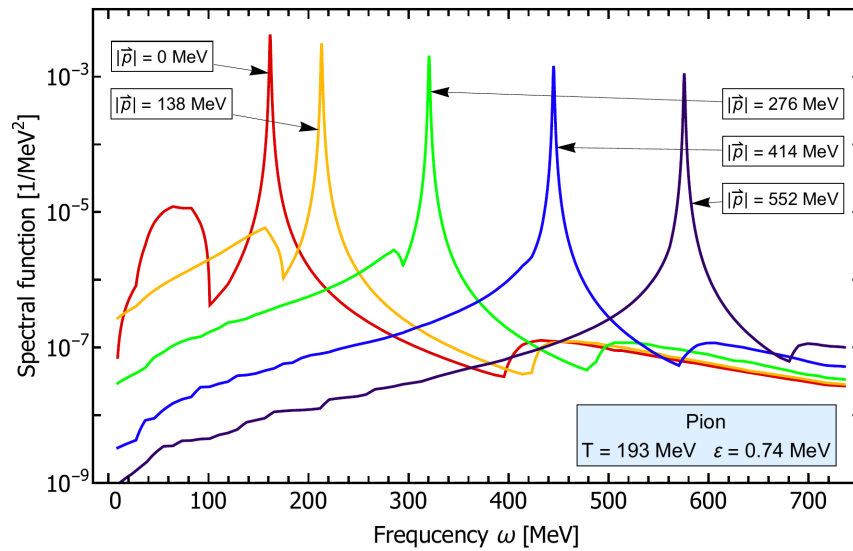
## Spectral functions in QCD



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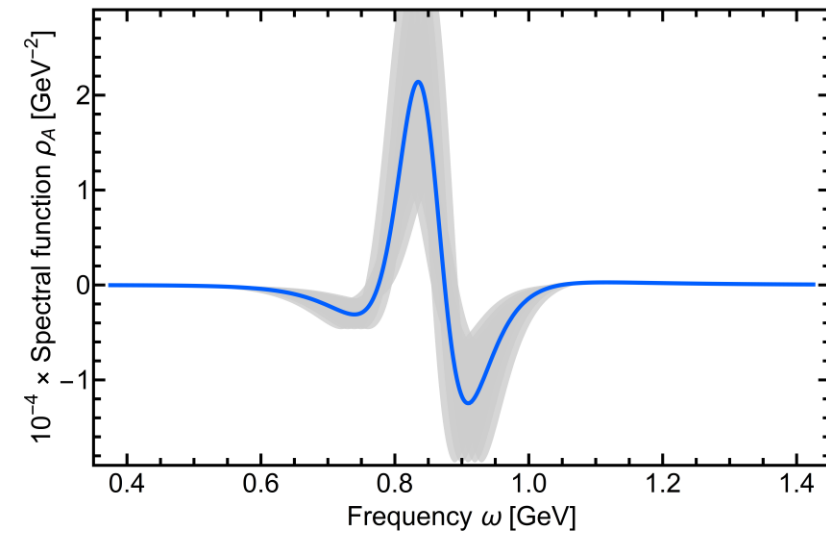


## Direct calculation



Pawlowski, Strodthoff, NW, arxiv:1711.07444

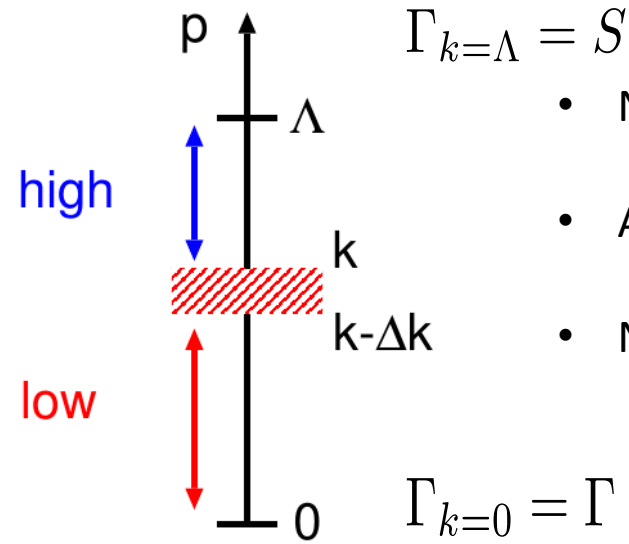
## Reconstruction



Cyrol, Pawlowski, Rothkopf, NW arxiv:1804.today

QCD from the FRG

# Functional Renormalization Group



- Non-perturbative first principle method
- Access to physical mechanisms
- No sign problem
  - Chemical potential
  - Real time

cf. poster by Anton Cyrol

cf. talk by Mario Mitter

cf. talk by Fabian Rennecke

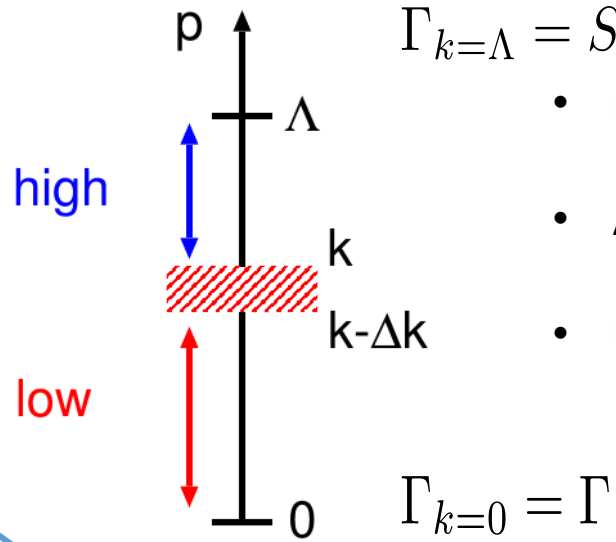
.....

QCD from the FRG

Functional Renormalization Group

Flow equation for QCD

$$\partial_t \Gamma_k[\Phi] = \frac{1}{2} \text{[diagram 1]} - \text{[diagram 2]} - \text{[diagram 3]} + \frac{1}{2} \text{[diagram 4]}$$



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Bound states efficiently taken into account via Dynamical Hadronization

cf. poster by Anton Cyrol  
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QCD from the FRG

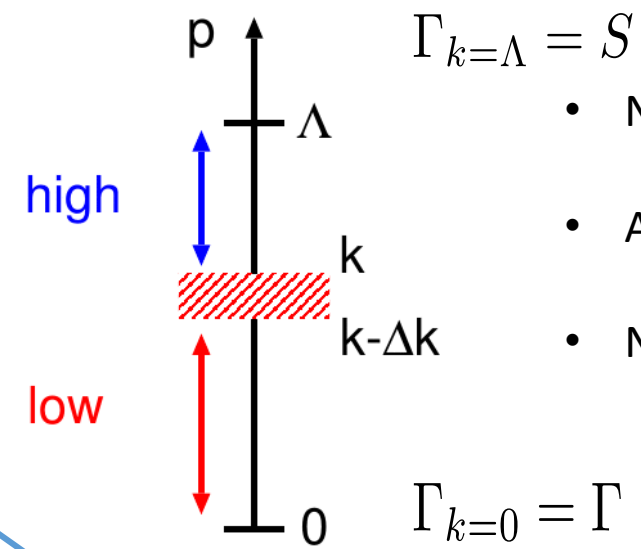
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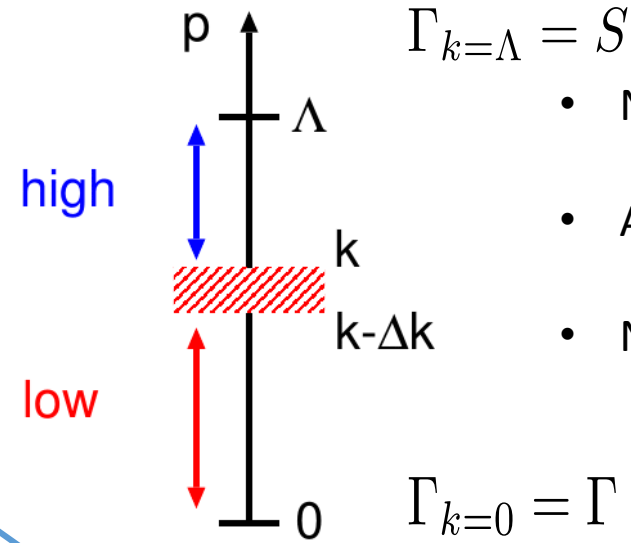
Functional Renormalization Group

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$$\partial_t \Gamma_k[\Phi] = \frac{1}{2} \text{[orange loop]} - \text{[dashed loop]} - \text{[solid loop]} + \frac{1}{2} \text{[blue loop]}$$

Reconstruction

Direct calculation



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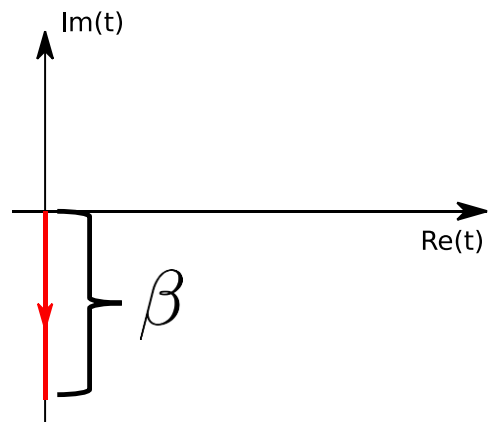
.....

Collaborative effort fQCD collaboration:

J. Braun, A. Cyrol, W.-j. Fu, M. Leonhardt, M. Mitter, J.M. Pawłowski, M. Pospiech, F. Rennecke, C. Schneider, NW

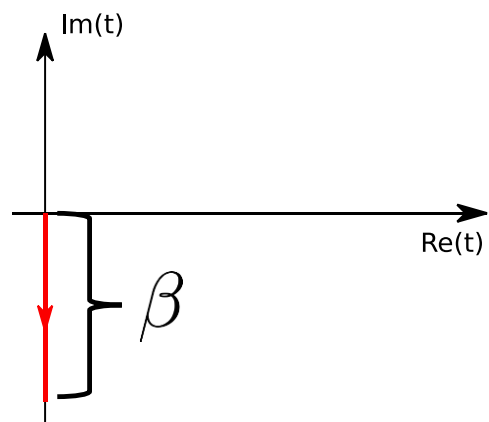
## Introduction

## From imaginary to real times

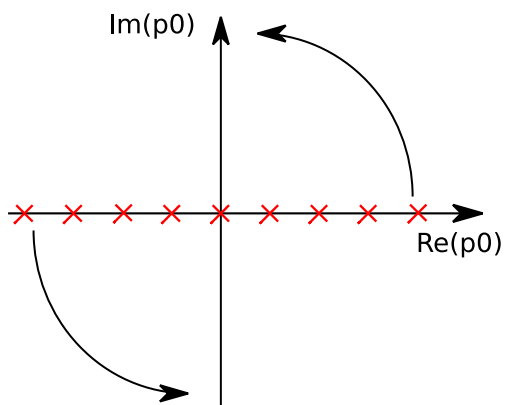


Matsubara contour

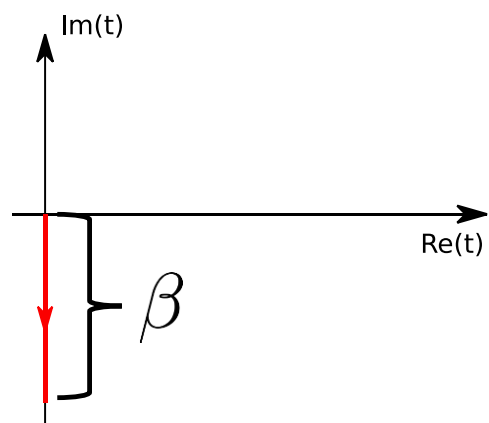
## From imaginary to real times



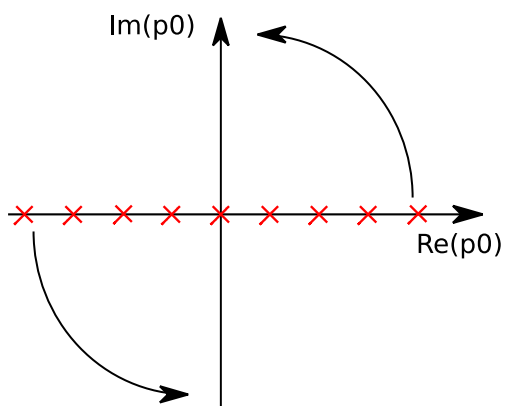
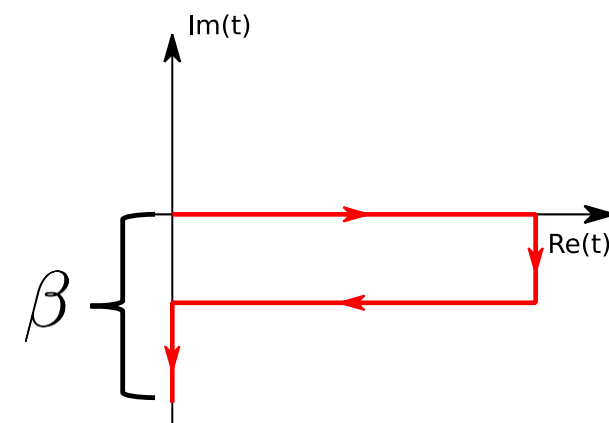
Matsubara contour

Continuation from  
Matsubara frequencies

## From imaginary to real times

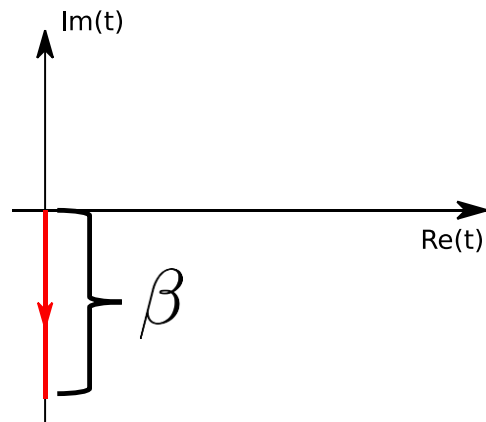


Matsubara contour

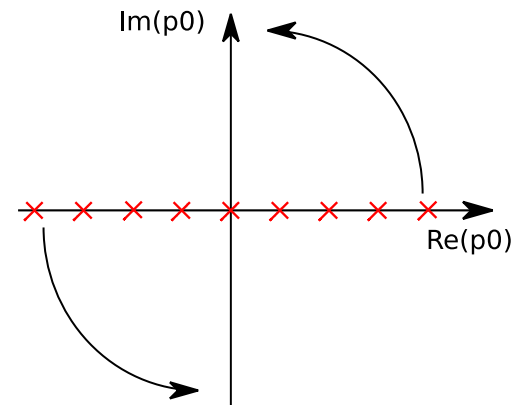
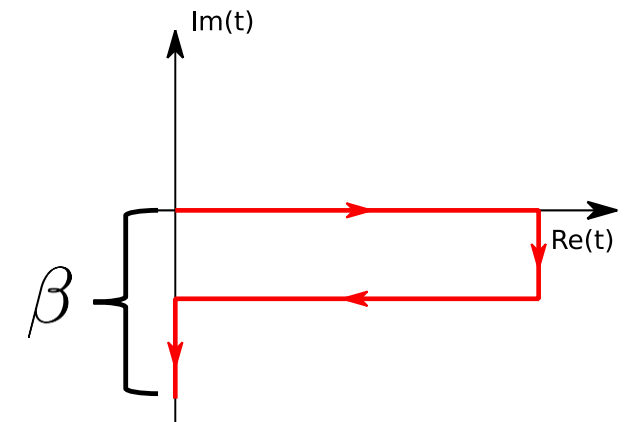
Continuation from  
Matsubara frequencies

Schwinger-Keldysh contour

## From imaginary to real times



Matsubara contour

Continuation from  
Matsubara frequencies

Schwinger-Keldysh contour

Use analyticity constraints and KMS condition to obtain real time correlation functions from Matsubara formalism

## From imaginary to real times

### Prerequisites :

Assume the existence of a spectral representation

$$G(p_0, \vec{p}) = \int_{\eta>0} 2\eta \frac{\rho(\eta, \vec{p})}{p_0^2 + \eta^2} + \sum_{j \in \{\text{poles}\}} \frac{R_j}{p_0^2 + M_j^2}$$



Possible to allow for additional complex conjugate poles



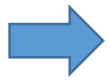
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Possible to allow for additional complex conjugate poles



Strong constraints on the analytic structure from the existence of a spectral representation

## From imaginary to real times

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Assume the existence of a spectral representation

$$G(p_0, \vec{p}) = \int_{\eta>0} 2\eta \frac{\rho(\eta, \vec{p})}{p_0^2 + \eta^2} + \sum_{j \in \{\text{poles}\}} \frac{R_j}{p_0^2 + M_j^2}$$

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➔ Strong constrains on the analytic structure from the existence of a spectral representation

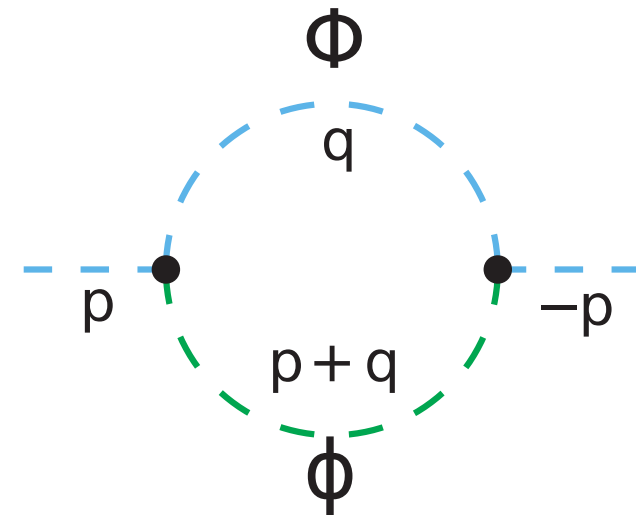
### Example :

One-loop perturbation theory

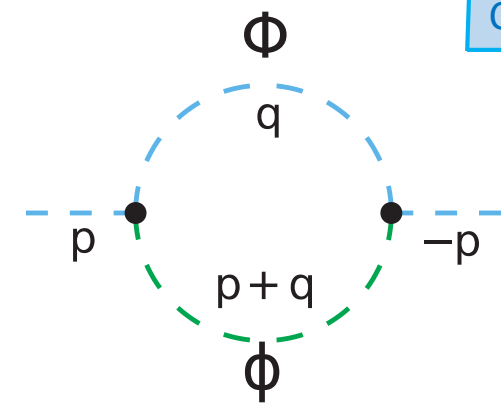
Two bosonic fields with interaction  $\sim \Phi\Phi\varphi$

Calculate  $\Gamma^{(2)}(p)$  for  $p^0 \in \mathbb{C}$

Calculate Matsubara sum  $\oint_{T,q} G_1(q+p)G_2(q)$



# Illustrative example



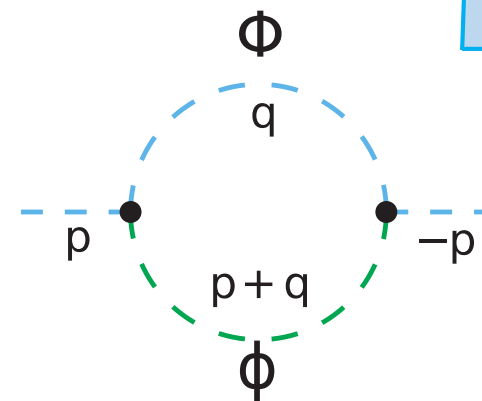
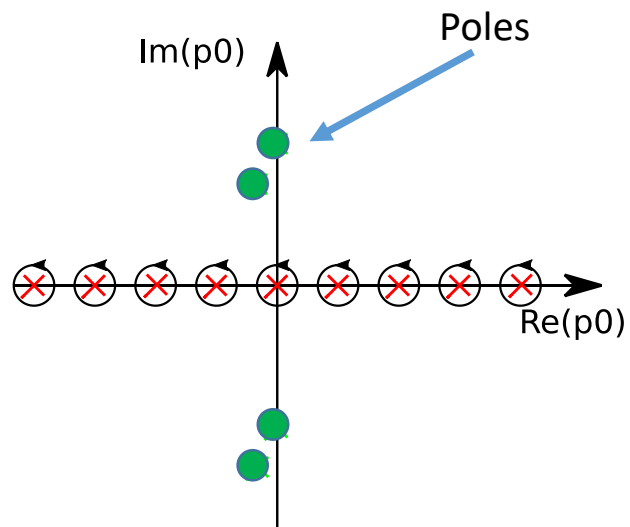
$$\sum_T \frac{1}{(q_0 + p_0)^2 + (\epsilon_{q+p}^1)^2} \frac{1}{(q_0)^2 + (\epsilon_q^2)^2}$$

Illustrative example

Replace sum by contour integral:

$$T \sum_n f(2\pi nT) = -\frac{1}{2} \int_C dz f(z) [1 + 2n_B(iz)]$$

Bosonic occupation number



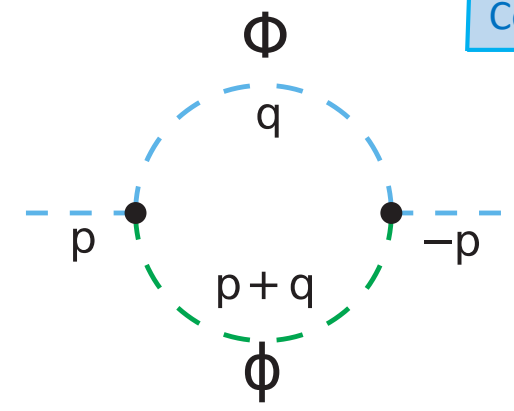
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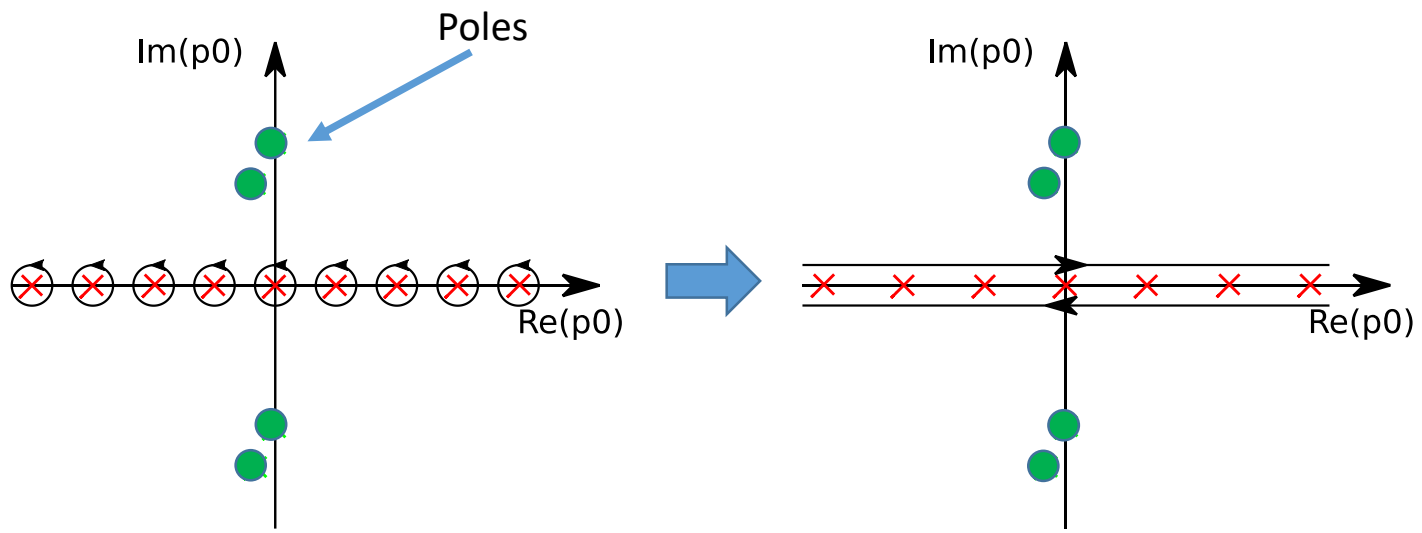
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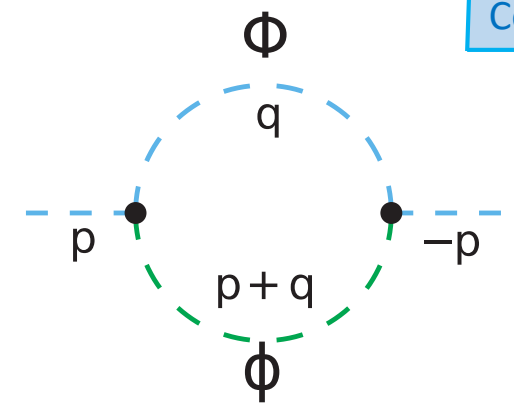


Illustrative example

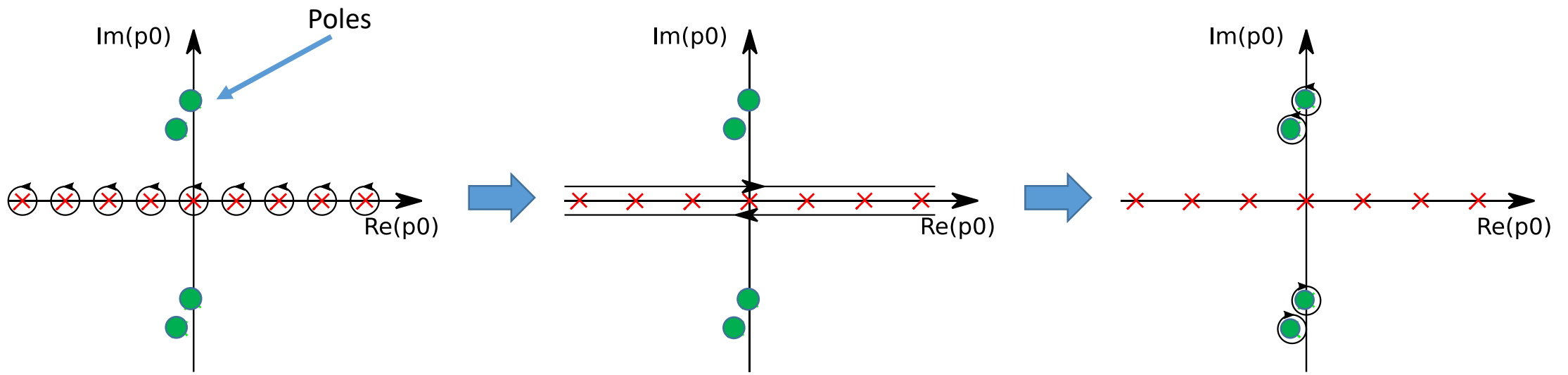
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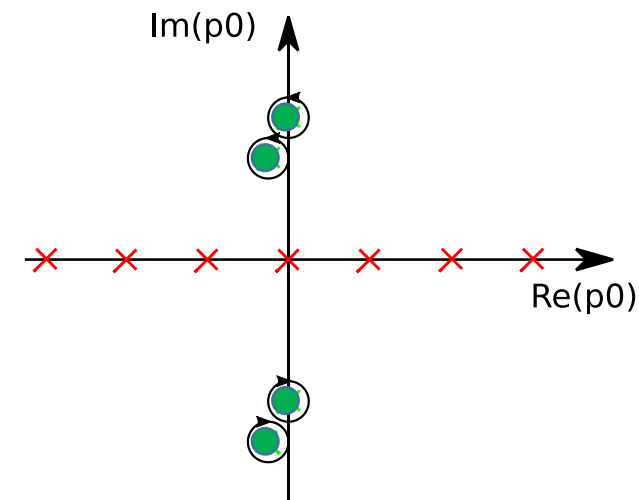


$$\sum_T \frac{1}{(q_0 + p_0)^2 + (\epsilon_{q+p}^1)^2} \frac{1}{(q_0)^2 + (\epsilon_q^2)^2}$$



# Illustrative example

$$\frac{1}{i} \sum_{\pm} (\text{Res}_1^{\pm} \cdot [1 + 2n_B(-ip_0 + \epsilon_{q+p}^1)] + \text{Res}_2^{\pm} \cdot [1 + 2n_B(\epsilon_q^2)])$$

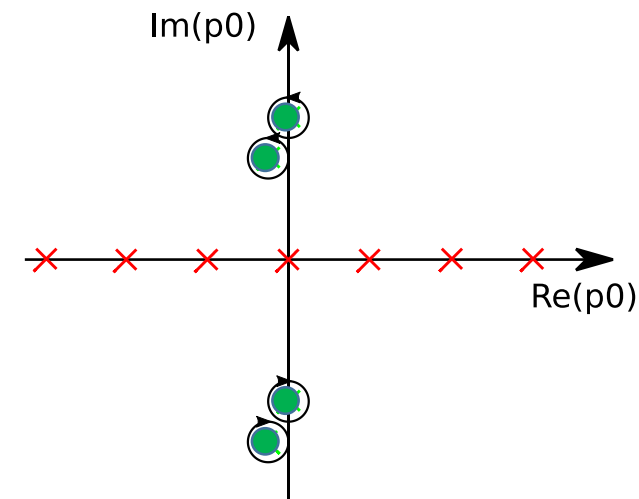


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$$p_0 = 2m\pi T \quad m \in \mathbb{Z}$$

Identify ambiguity of the analytic continuation





# Illustrative example

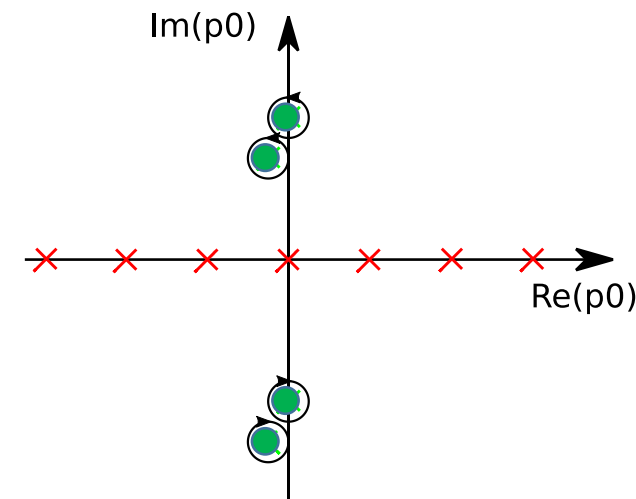
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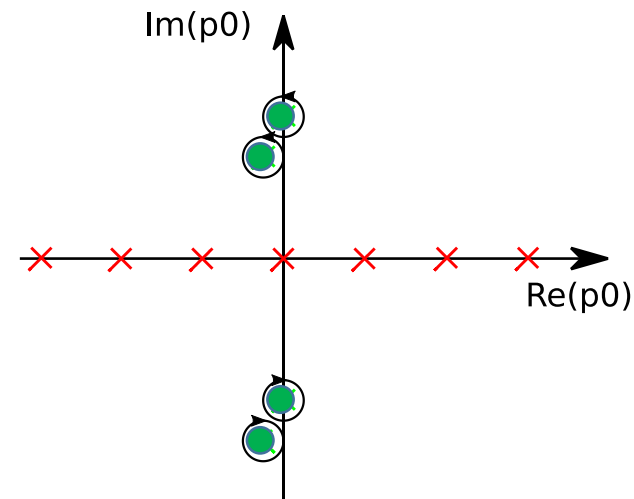
$$e^{ip_0} = 1$$

Identify ambiguity of the analytic continuation



Illustrative example

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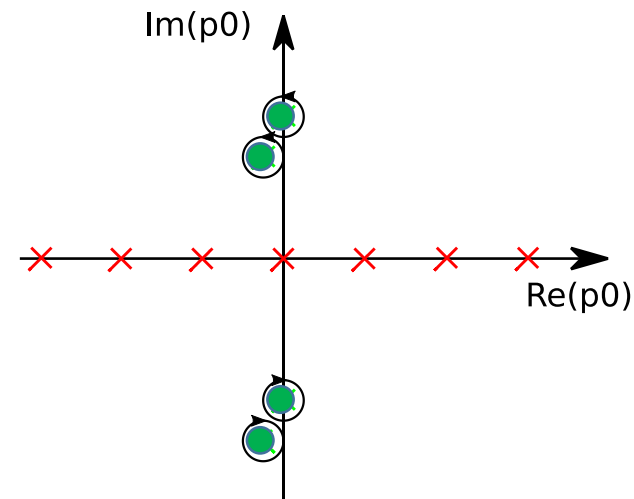
Mathematically rigorous  
 Baym and Mermin, *Journal of Mathematical Physics* 2, 232 (1961)  
 Evans, *Nucl.Phys.* B374 (1992)

Analyticity off the imaginary axis

Correct decay behaviour at infinity

Illustrative example

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 Baym and Mermin, *Journal of Mathematical Physics* 2, 232 (1961)  
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Analyticity off the imaginary axis

Correct decay behaviour at infinity

Unique physical analytic continuation identified by setting  $e^{ip_0} = \pm 1$  everywhere

Remarks

## Remarks

- Numerically accessible

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- Corresponds to a contour deformation at vanishing temperature

[Strodthoff, PRD 95 \(2017\) no.7, 076002](#)

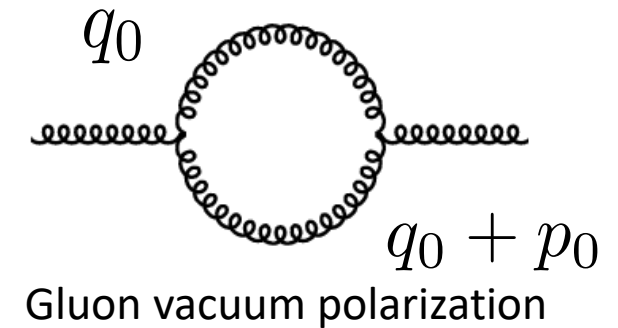
[Pawlowski, Strodthoff, NW, arxiv:1711.07444](#)

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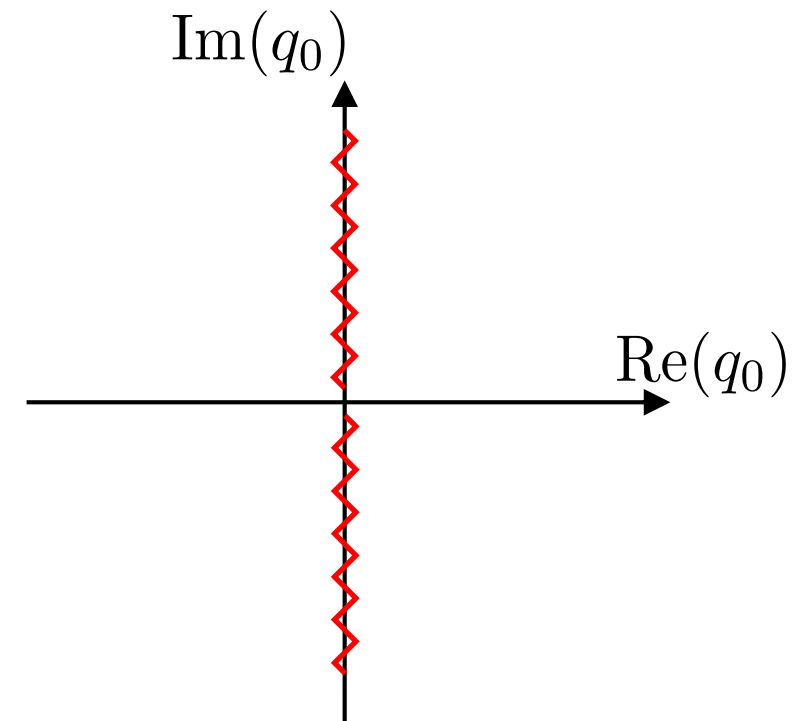
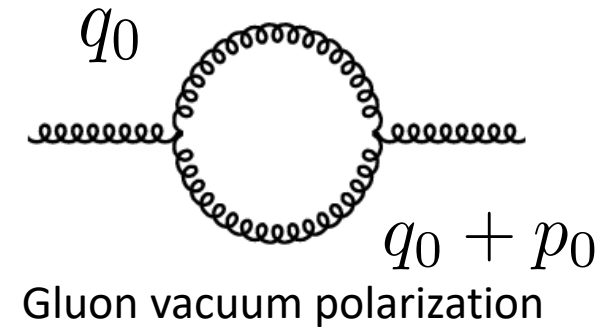


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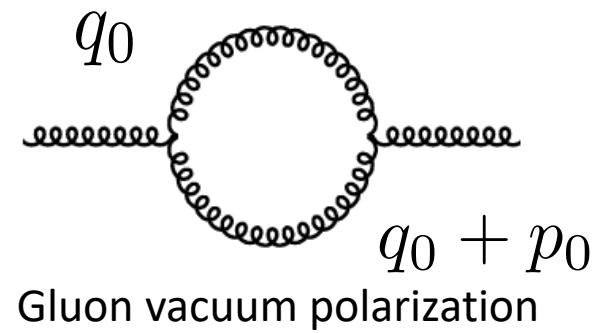
Analytic structure of the gluon



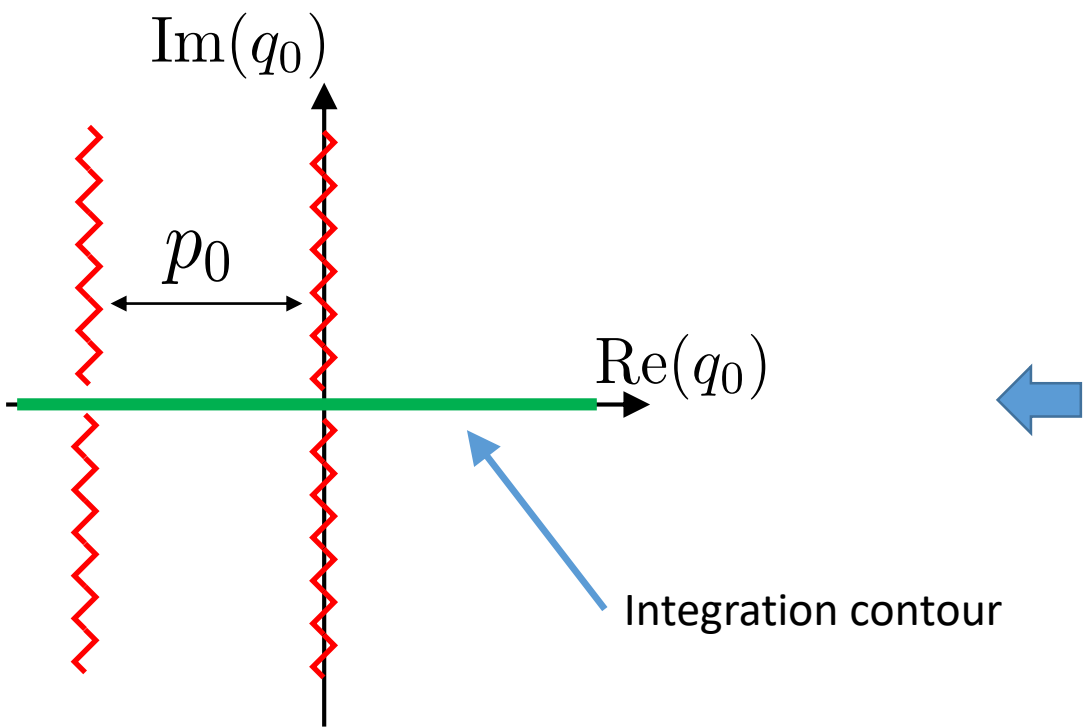
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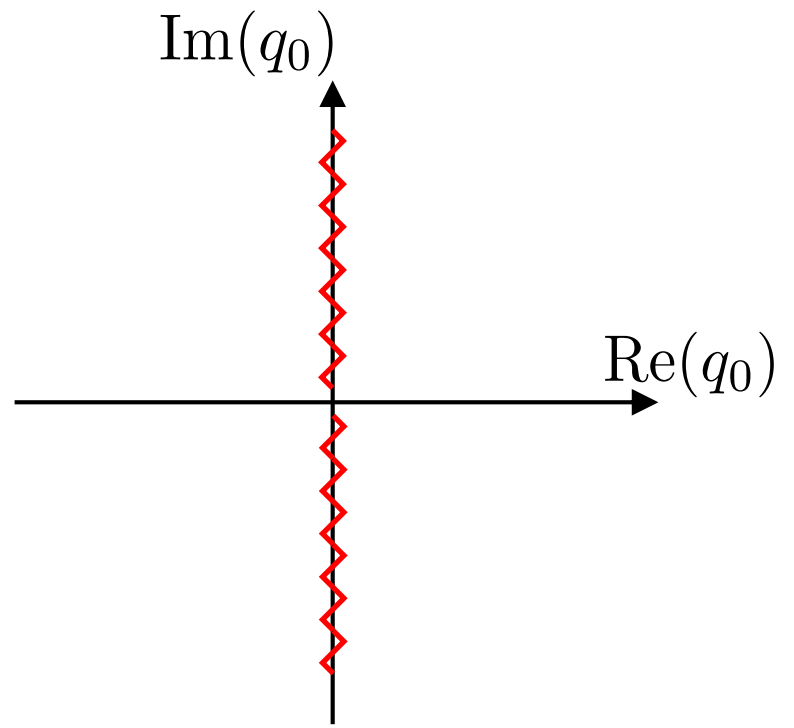
[Strodthoff, PRD 95 \(2017\) no.7, 076002](#)  
[Pawlowski, Strodthoff, NW, arxiv:1711.07444](#)



Gluon vacuum polarization



Analytic structure gluon polarization diagram

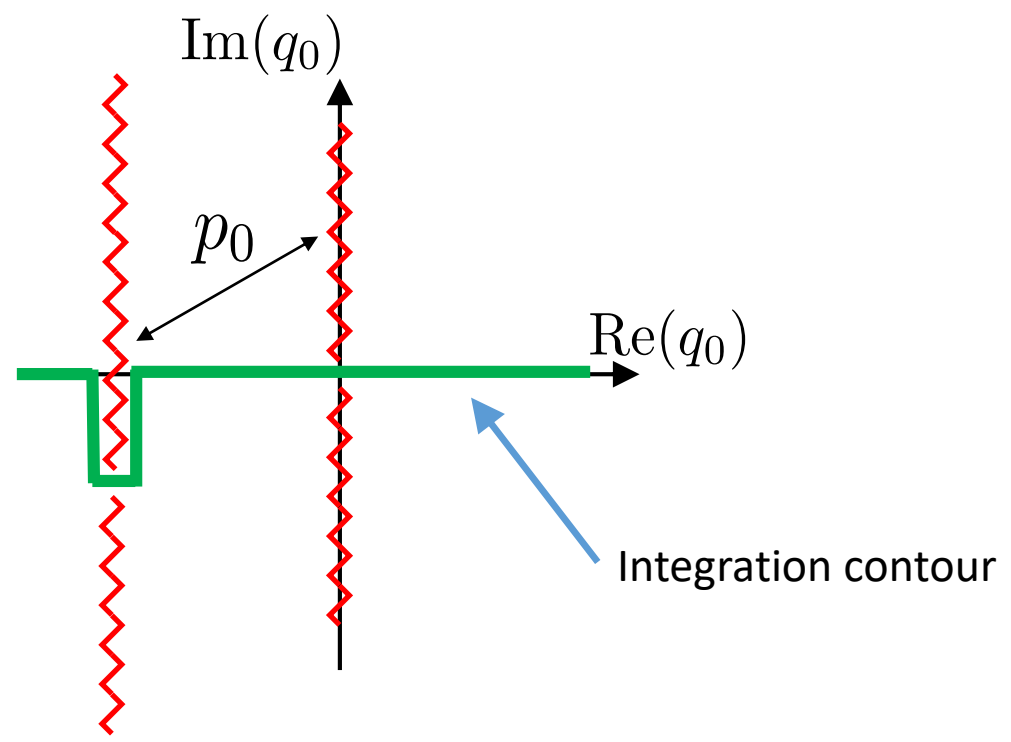
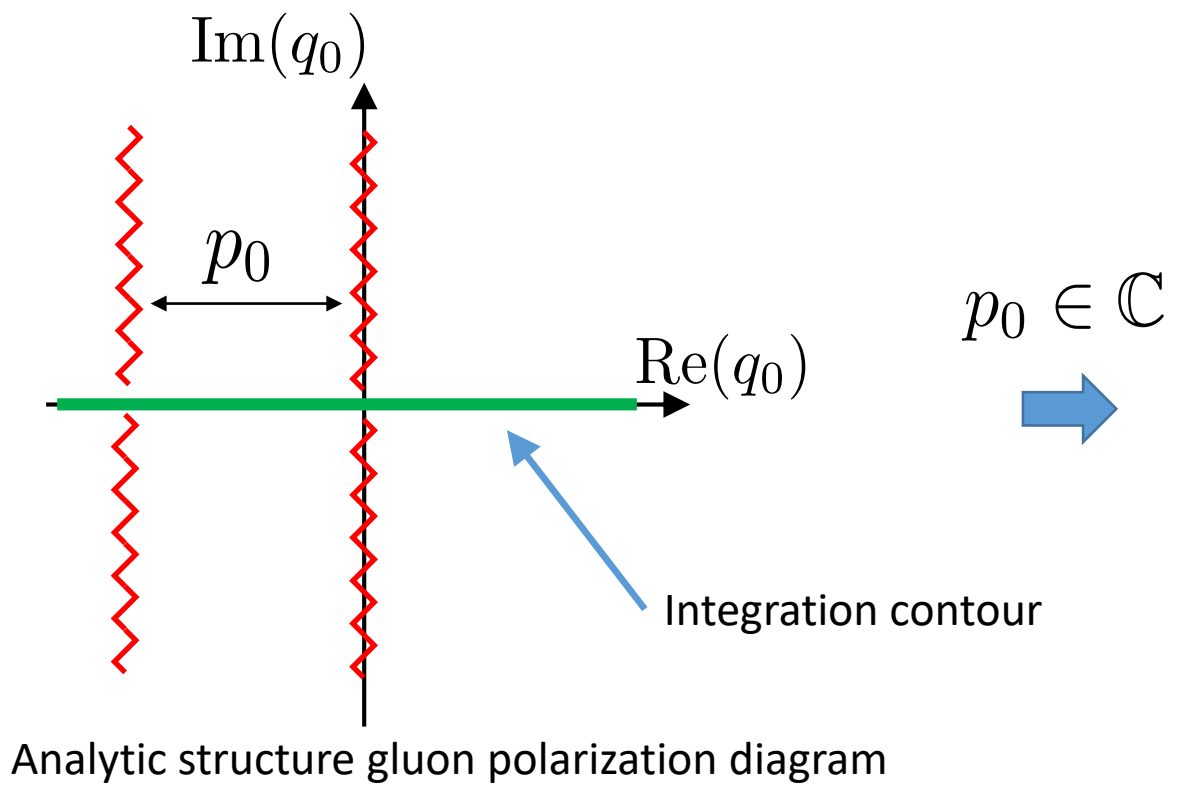
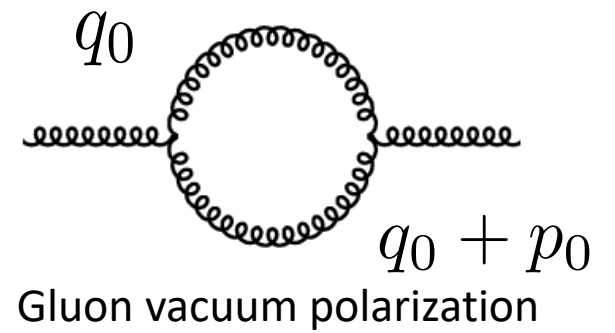


Analytic structure of the gluon

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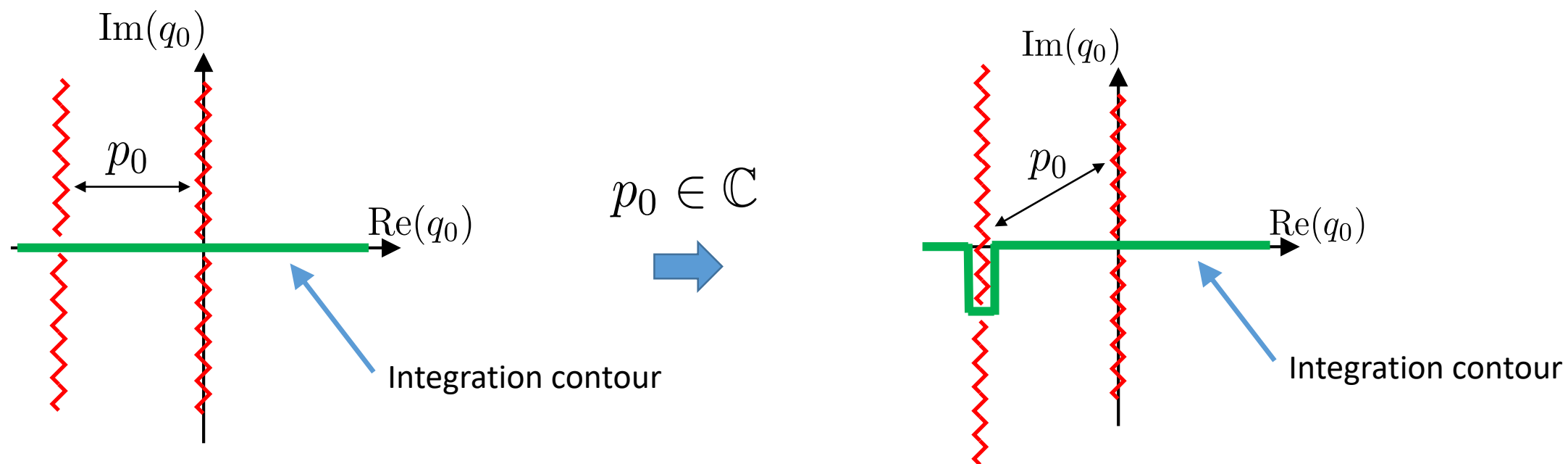
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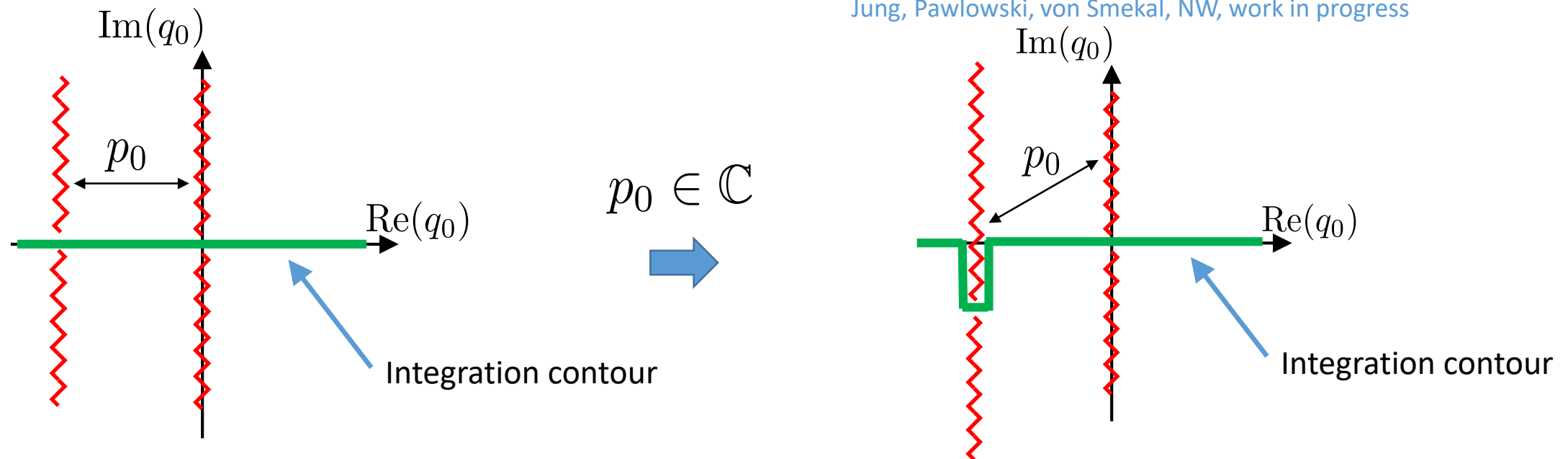


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$$G(p_0, \vec{p}) = \int_{\eta>0} 2\eta \frac{\rho(\eta, \vec{p})}{p_0^2 + \eta^2}$$

[Jung, Pawlowski, von Smekal, NW, work in progress](#)




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## Regulator poles

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## Regulator poles

 $R_k(\vec{q}^2)$ 


No changes

Kamikado, Strodthoff, von Smekal, Wambach, Eur.Phys.J. C74, 2806 (2014)  
Tripolt, Strodthoff, von Smekal, Wambach, Phys.Rev. D89, 034010 (2014)

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- Numerically accessible
- Corresponds to a contour deformation at vanishing temperature  
[Strodthoff, PRD 95 \(2017\) no.7, 076002](#)  
[Pawlowski, Strodthoff, NW, arxiv:1711.07444](#)
- Considering poles is sufficient
- Branch cuts can be mapped to poles via spectral/integral representations  
[Jung, Pawlowski, von Smekal, NW, work in progress](#)
- Generalization to the FRG  no new conceptual problems

$$G(p_0, \vec{p}) = \int_{\eta>0} 2\eta \frac{\rho(\eta, \vec{p})}{p_0^2 + \eta^2}$$

## Regulator poles

$$R_k(\vec{q}^2)$$

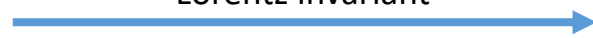


No changes

[Kamikado, Strodthoff, von Smekal, Wambach, Eur.Phys.J. C74, 2806 \(2014\)](#)
[Tripolt, Strodthoff, von Smekal, Wambach, Phys.Rev. D89, 034010 \(2014\)](#)

$$R_k(q^2)$$

Lorentz invariant



Additional poles

[Foerchinger, JHEP 1205 \(2012\) 021](#)
[Pawlowski, Strodthoff, Phys.Rev. D92 \(2015\)](#)
[Pawlowski, Strodthoff, NW arxiv:1711.07444](#)

Direct calculation



## Application to the $O(N)$ -Model

Spectral functions of the  $O(N)$  model

## Application to the O(N)-Model

Effective description of the lightest mesons

Spectral functions of the O(N) model

$$\rho(\omega, \vec{p}) = -2 \operatorname{Im} G_{\mathbf{R}}(\omega, \vec{p})$$

## Application to the O(N)-Model

Effective description of the lightest mesons

Truncation:

$$\Gamma_k = \sum_{T,q} \Delta\Gamma_{\sigma}^{(2)} + \Delta\Gamma_{\pi}^{(2)} + V(\sigma)$$

$$\text{Vacuum : } \Delta\Gamma_x^{(2)} = \Gamma_x^{(2)}(q^2) - \Gamma_x^{(2)}(0)$$

$$\text{Finite Temperature : } \Delta\Gamma_x^{(2)} = Z_x q^2$$

## Spectral functions of the O(N) model

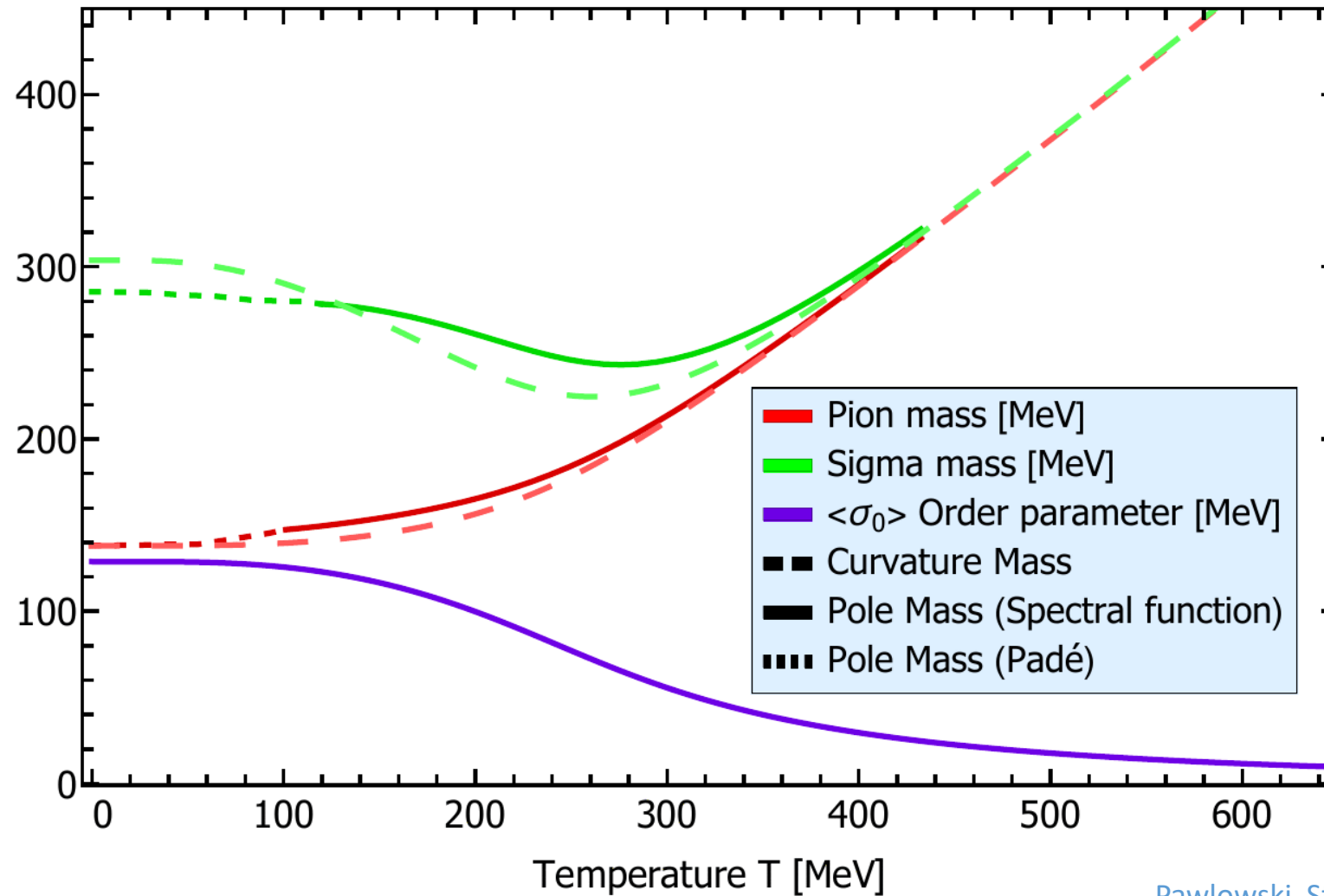
$$\rho(\omega, \vec{p}) = -2 \text{Im } G_R(\omega, \vec{p})$$

$$\partial_k \Delta\Gamma_{\pi,k}^{(2)} = \text{Diagram 1} + \text{Diagram 2}$$

$$\partial_k \Delta\Gamma_{\sigma,k}^{(2)} = \text{Diagram 3} + \text{Diagram 4}$$

## Application to the O(N)-Model

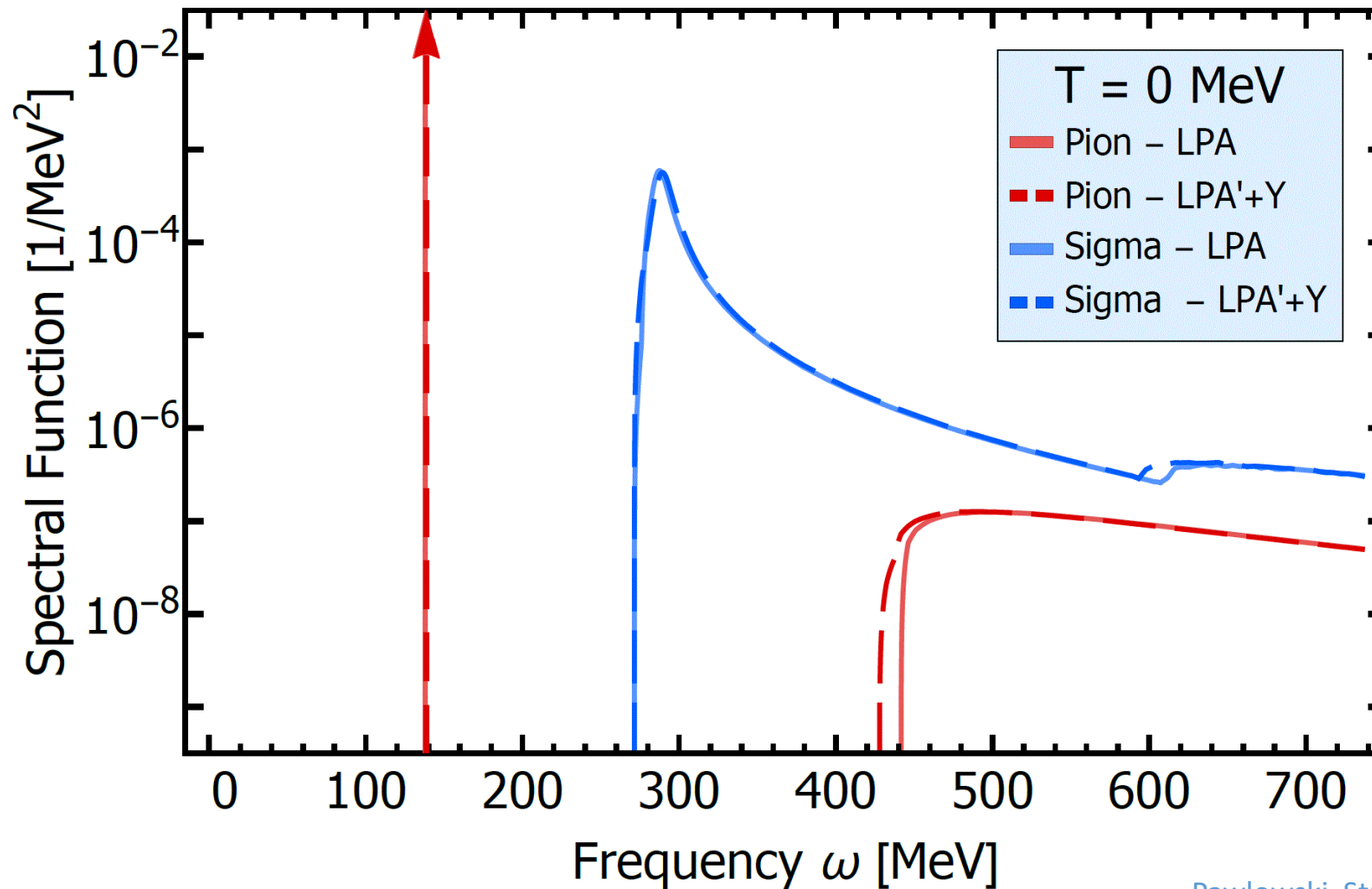
## Phase structure



Pawlowski, Strodthoff, NW, arxiv:1711.07444

## Application to the O(N)-Model

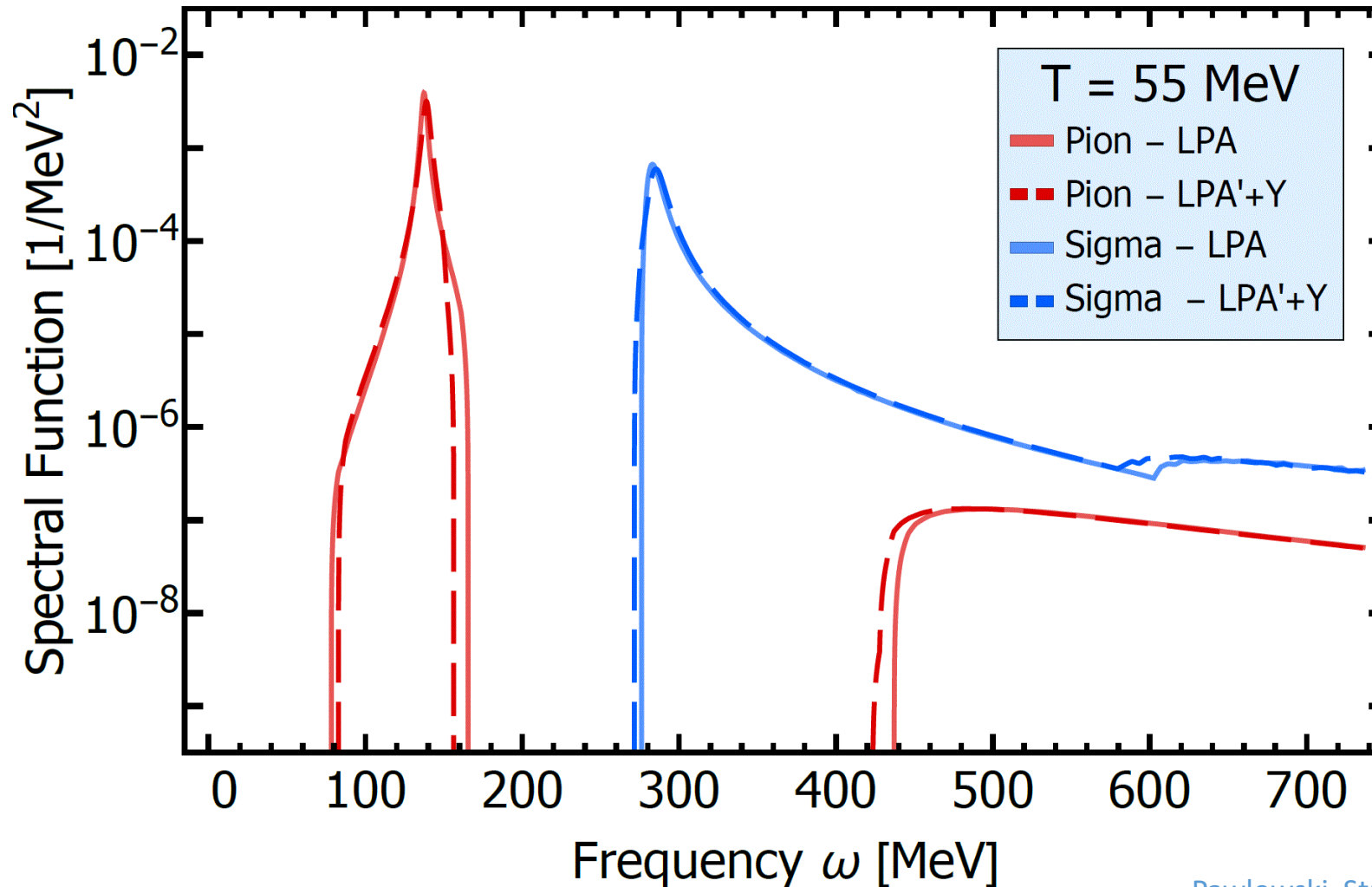
## Finite temperature spectral functions



Pawlowski, Strodthoff, NW, arxiv:1711.07444

## Application to the O(N)-Model

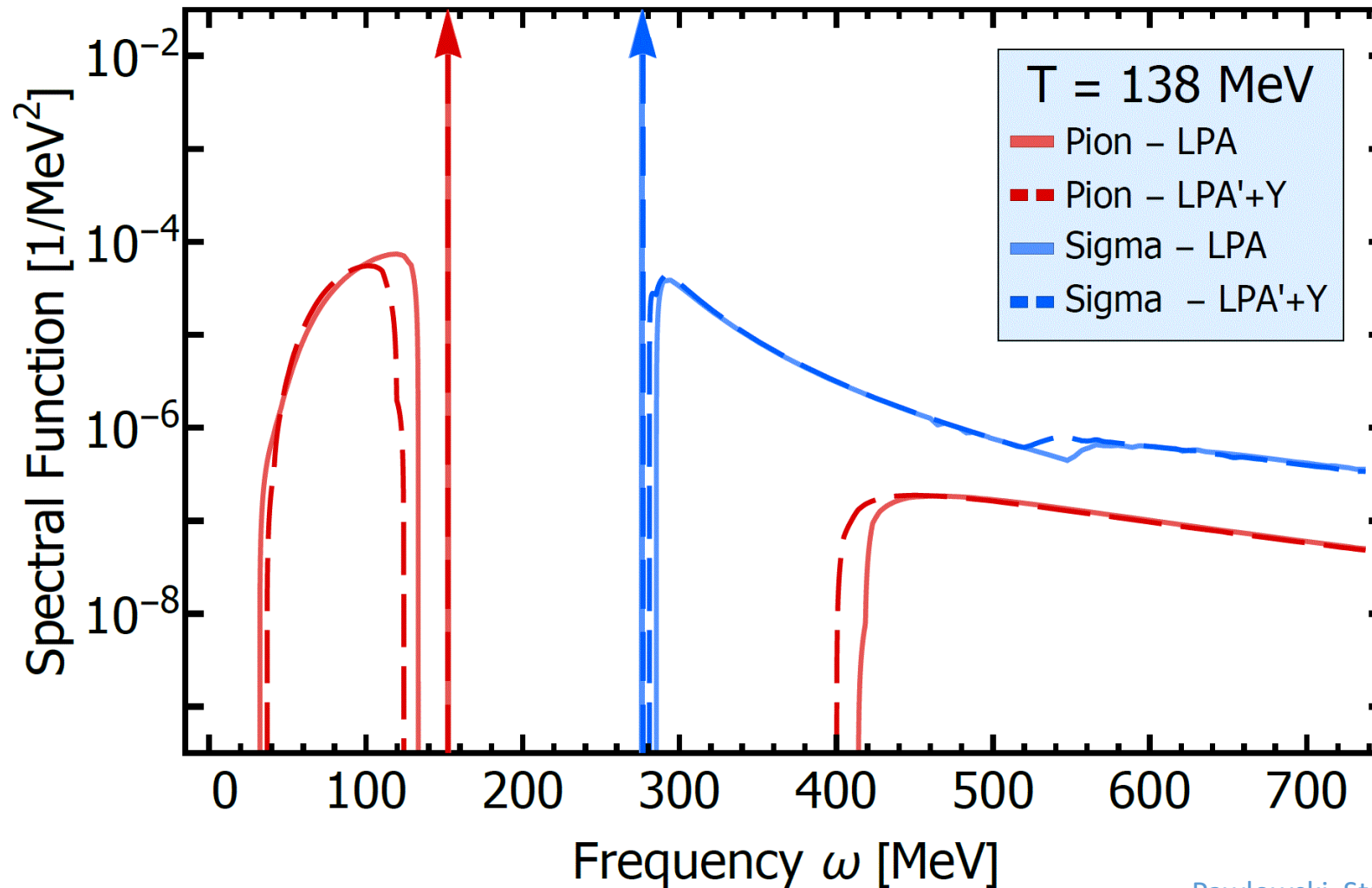
## Finite temperature spectral functions



Pawlowski, Strodthoff, NW, arxiv:1711.07444

## Application to the O(N)-Model

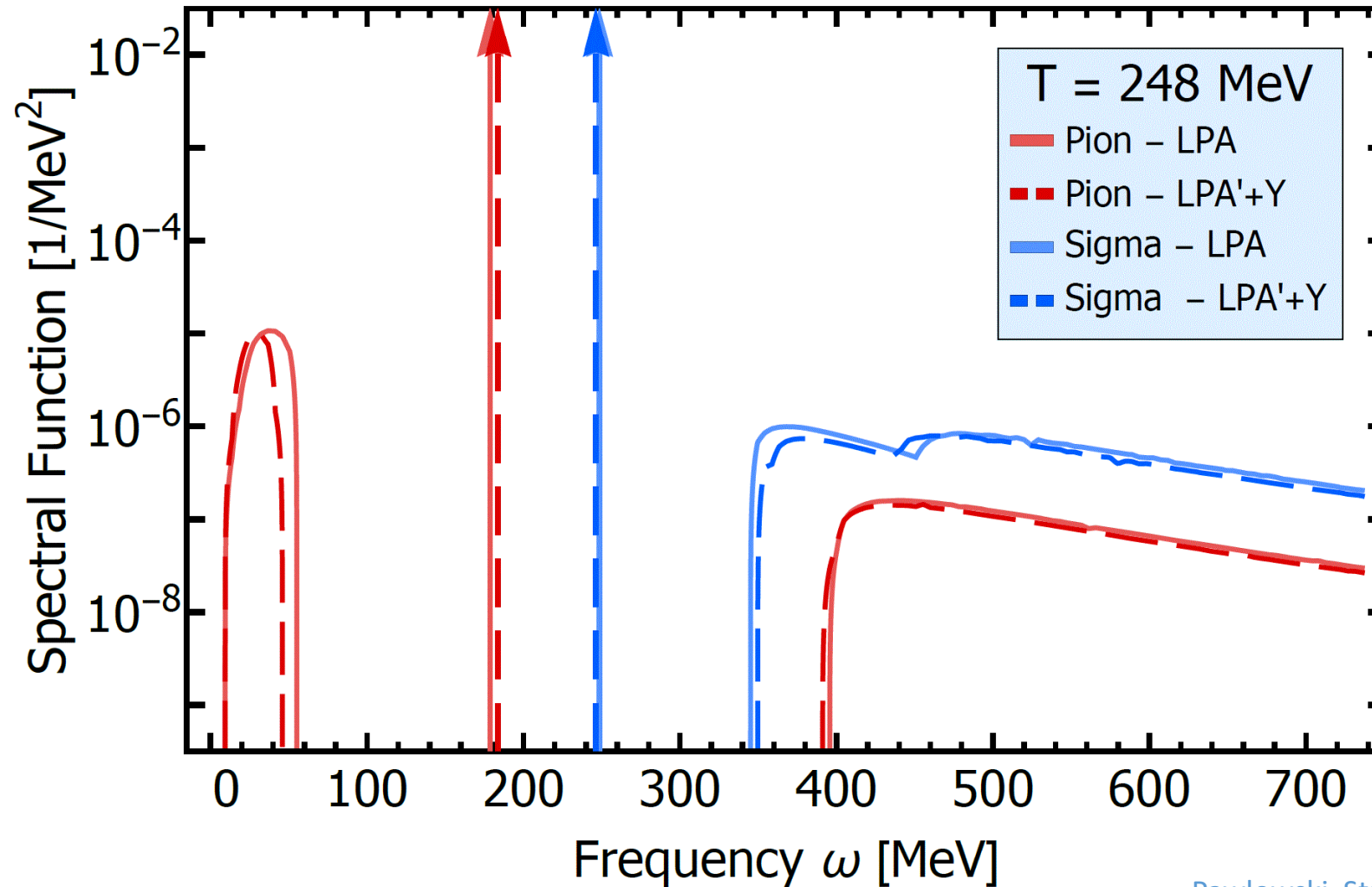
## Finite temperature spectral functions



Pawlowski, Strodthoff, NW, arxiv:1711.07444

## Application to the O(N)-Model

## Finite temperature spectral functions

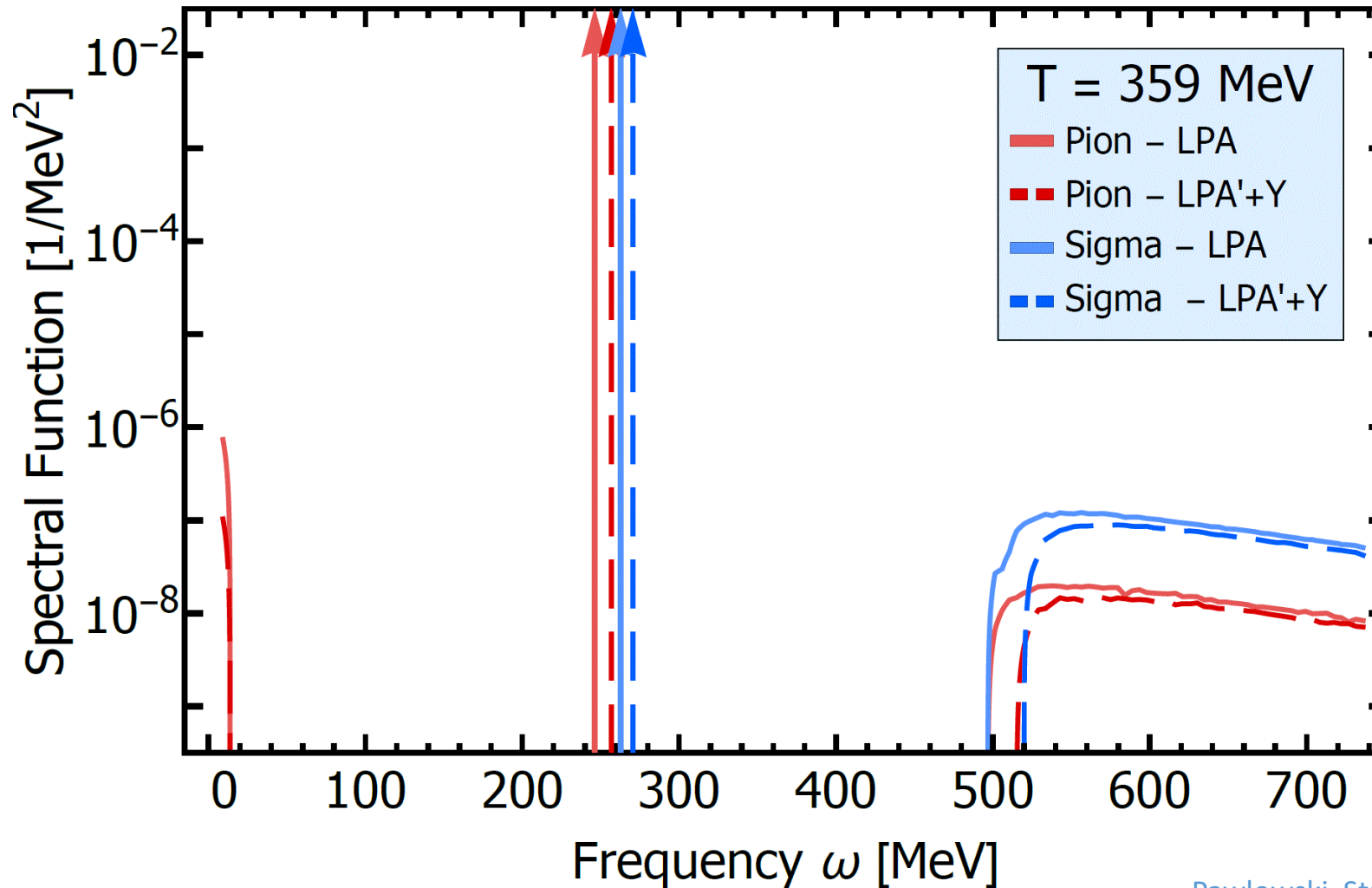


Pawlowski, Strodthoff, NW, arxiv:1711.07444



## Application to the O(N)-Model

## Finite temperature spectral functions

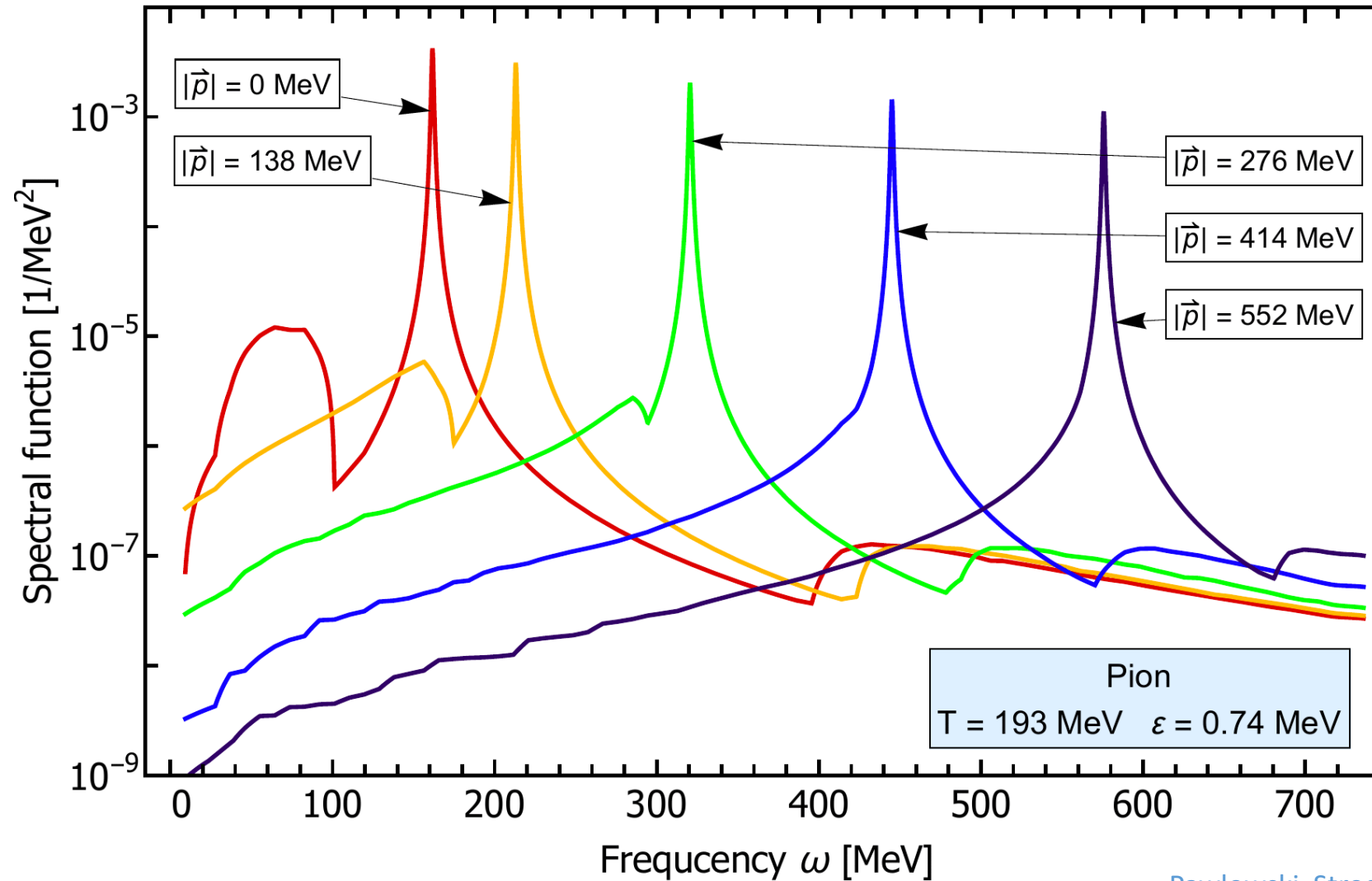


Pawlowski, Strodthoff, NW, arxiv:1711.07444

## Pion meson

## Application to the O(N)-Model

Finite temperature spectral function for various external momenta

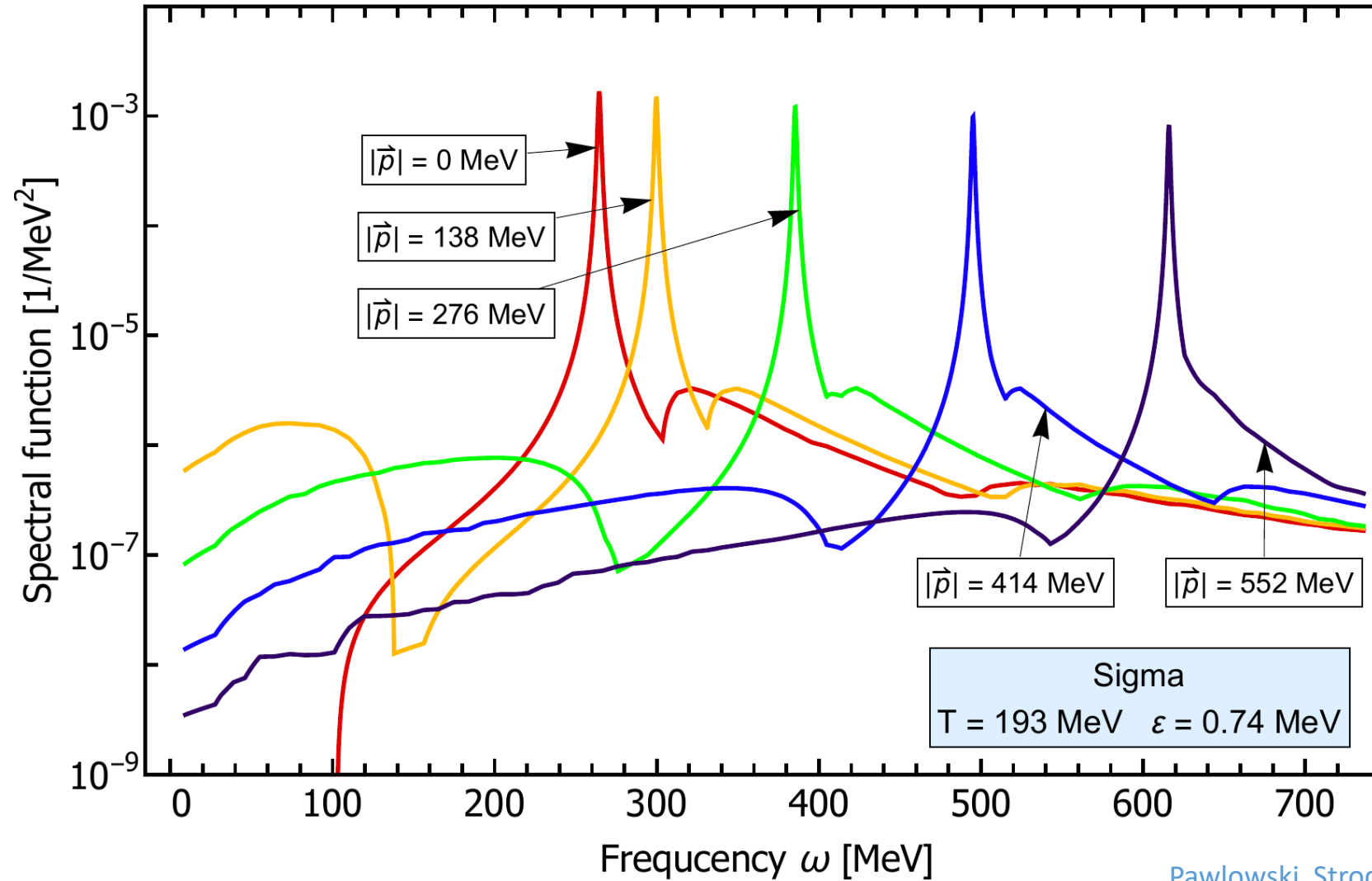


Pawlowski, Strodthoff, NW, arxiv:1711.07444

## Sigma meson

## Application to the O(N)-Model

Finite temperature spectral function for various external momenta



Pawlowski, Strodthoff, NW, arxiv:1711.07444

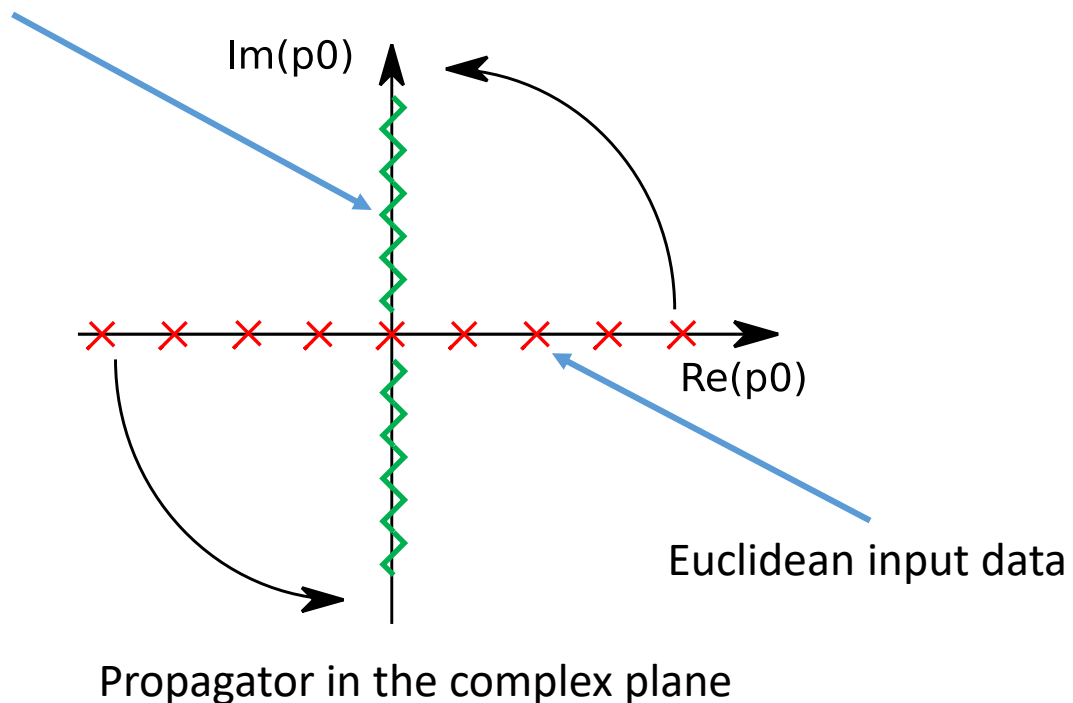
Reconstruction

# Reconstructing spectral functions

Spectral function defined as the discontinuity of the propagator

$$\rho(\omega, \vec{p}) = 2 \lim_{\varepsilon \rightarrow 0} \text{Im} \underbrace{G_E(-i(\omega + i\varepsilon), \vec{p})}_{\text{retarded propagator}}$$

Spectral function (discontinuity)



Propagator in the complex plane

# Reconstructing spectral functions

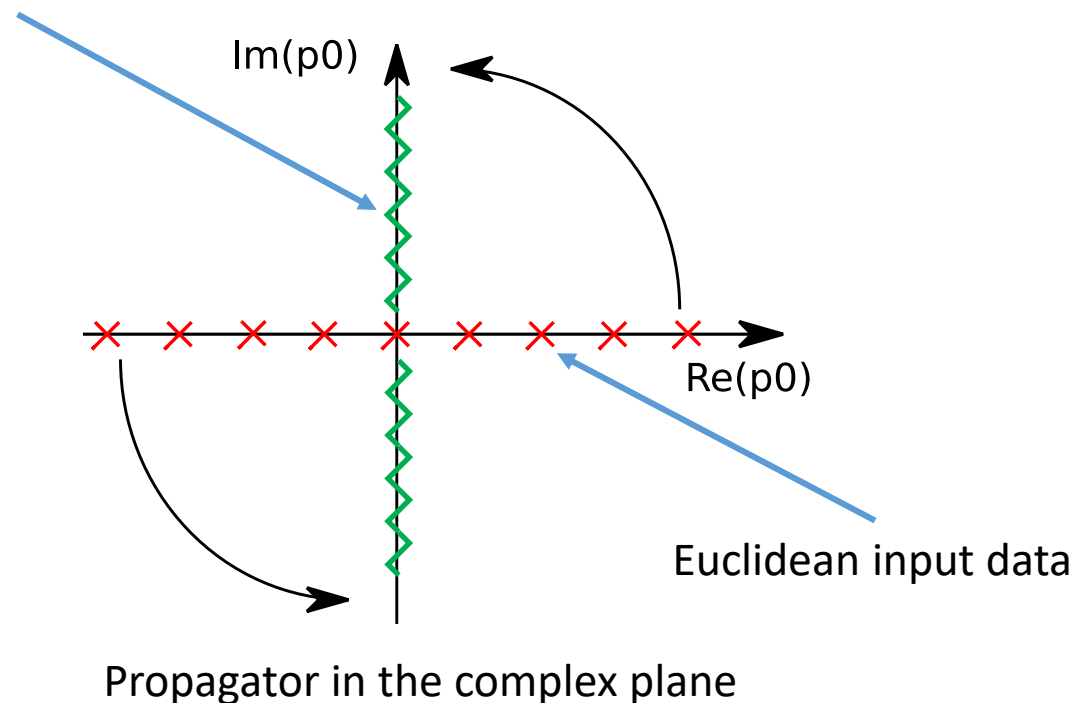
Spectral function defined as the discontinuity of the propagator

$$\rho(\omega, \vec{p}) = 2 \lim_{\varepsilon \rightarrow 0} \text{Im} \underbrace{G_E(-i(\omega + i\varepsilon), \vec{p})}_{\text{retarded propagator}}$$

Gluon violates reflection-positivity

➔ Spectral function positive and negative

Spectral function (discontinuity)



# Reconstructing spectral functions

Spectral function defined as the discontinuity of the propagator

$$\rho(\omega, \vec{p}) = 2 \lim_{\varepsilon \rightarrow 0} \text{Im} \underbrace{G_E(-i(\omega + i\varepsilon), \vec{p})}_{\text{retarded propagator}}$$

Gluon violates reflection-positivity

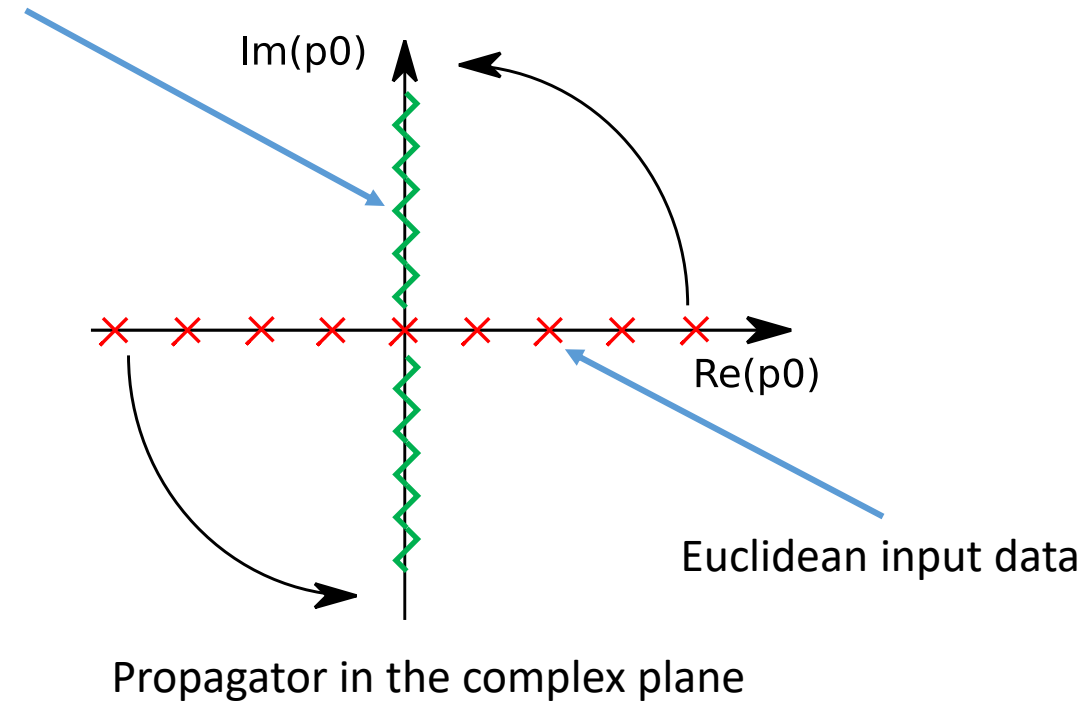
➔ Spectral function positive and negative

Use as much prior knowledge as possible!

- Analytic properties of the gluon spectral function
- Existence of a spectral representation has strong implications on the complex structure

➔ Construct suitable functional basis

Spectral function (discontinuity)



Propagator in the complex plane

YM from the FRG

System of coupled equations

$$\partial_t \text{---} \text{---}^{-1} = \text{---} \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} \text{---}$$

$$\partial_t \text{---} \text{---} \text{---}^{-1} = \text{---} \text{---} \text{---} \text{---} - 2 \text{---} \text{---} \text{---} \text{---} - \frac{1}{2} \text{---} \text{---} \text{---} \text{---}$$

$$\partial_t \text{---} \text{---} \text{---} = - \text{---} \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} \text{---} \text{---} = - \text{---} \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} \text{---} \text{---} \text{---} = + \text{---} \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} \text{---} - 2 \text{---} \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} \text{---} + \text{perm.}$$

Aiming at apparent convergence

- Only requires coupling at a perturbative scale
  - Absorbed during scale setting
- All quantities are fully dressed and momentum dependent
- Using Landau gauge

cf. poster by Anton Cyrol

cf. talk by Mario Mitter

Vacuum Yang-Mills

[Cyrol, Fister, Mitter, Pawłowski, Strodthoff, Phys.Rev. D94 \(2016\)](#)

Vacuum QCD

[Cyrol, Mitter, Pawłowski, Strodthoff, arXiv:1706.06326](#)

Finite temperature Yang-Mills

[Cyrol, Mitter, Pawłowski, Strodthoff, arXiv:1708.03482](#)

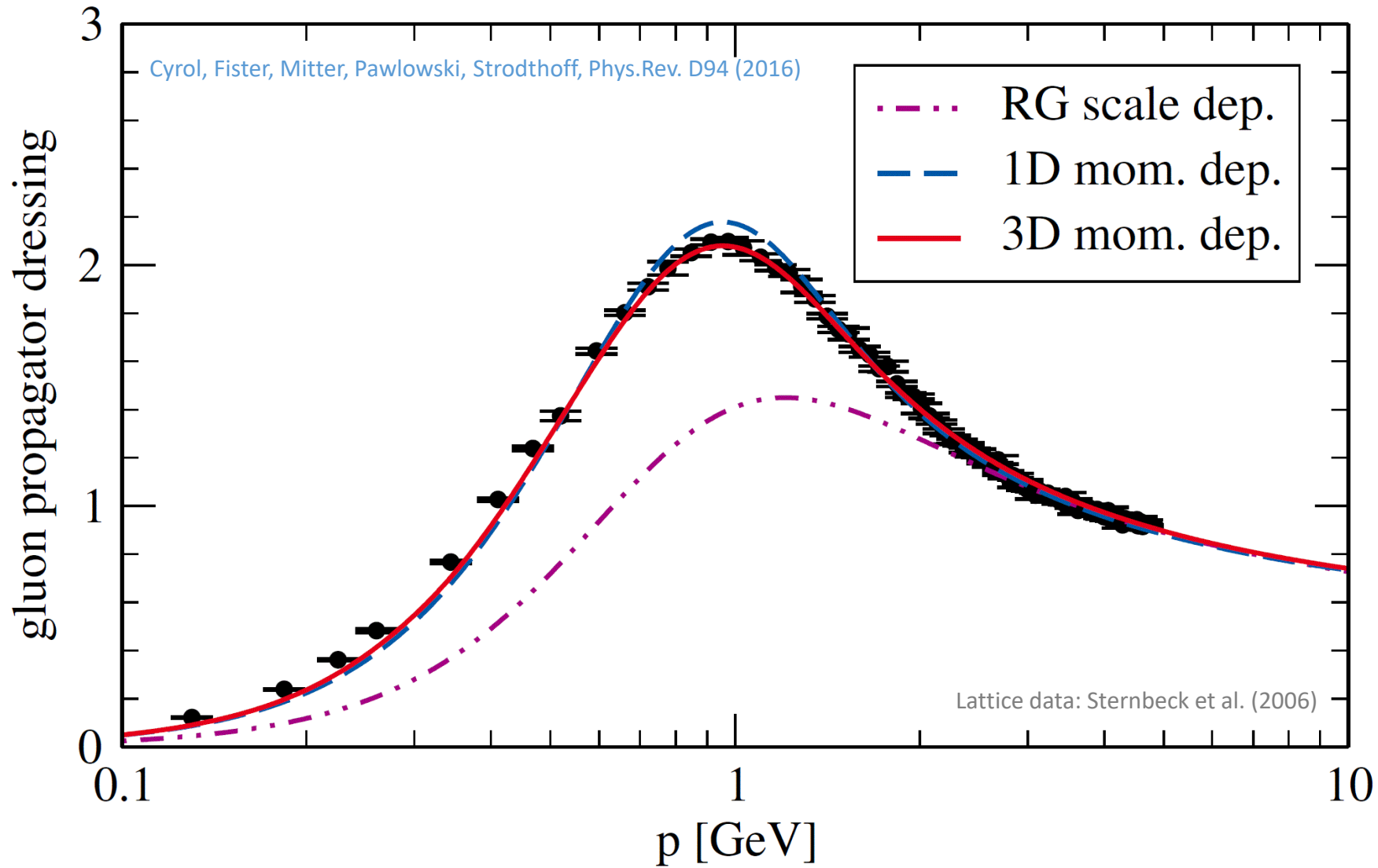
Finite temperature QCD, Extended truncations,...

[Cyrol, Mitter, Pawłowski, NW, work in progress](#)



# YM from the FRG

Aiming at apparent convergence



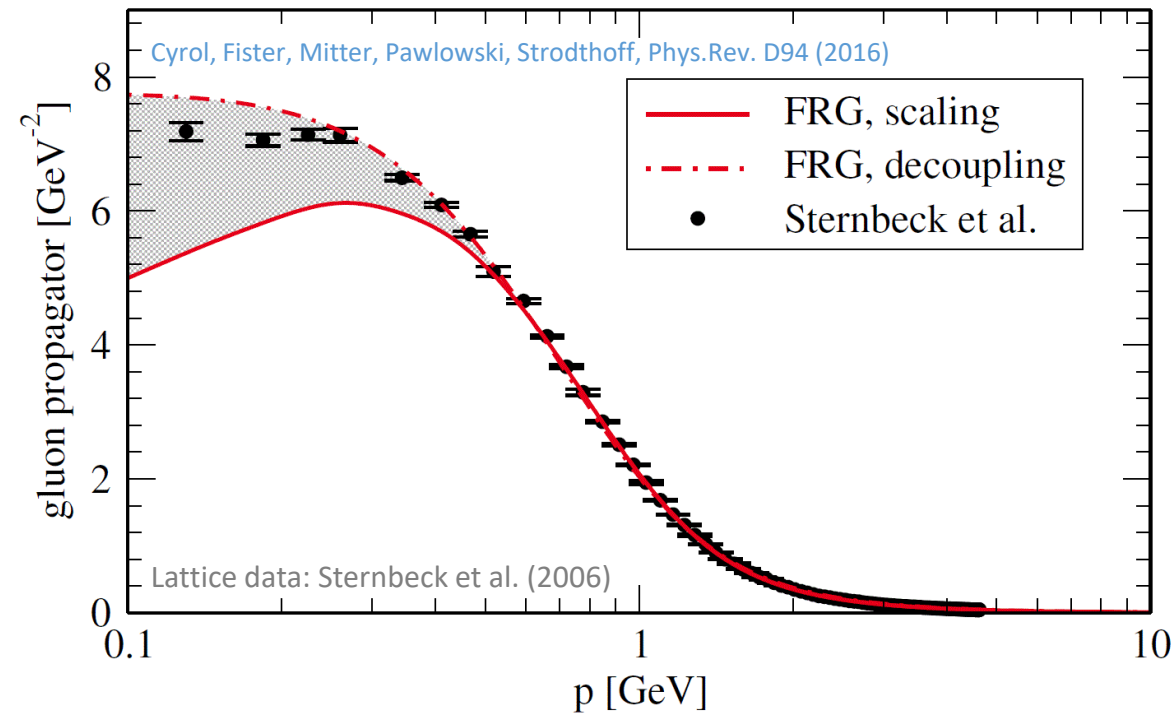
- Systematic improvement possible
- High numerical accuracy possible
- Direct computation of spectral functions possible in the near future

➔ Suitable for reconstruction methods

cf. poster by Anton Cyrol  
cf. talk by Mario Mitter

## Reconstructing spectral functions

Use as much prior knowledge as possible!



Cyrol, Pawlowski, Rothkopf, NW arxiv:1804.today

# Reconstructing spectral functions

1. Representation:

$$G_E(p) = \int_{\mu>0} \frac{d\mu}{2\pi} \frac{2\mu \rho(\mu)}{p^2 + \mu^2}$$

2. Normalization: Super-convergence property

Oehme, Zimmermann, Phys.Rev. D21 (1980)

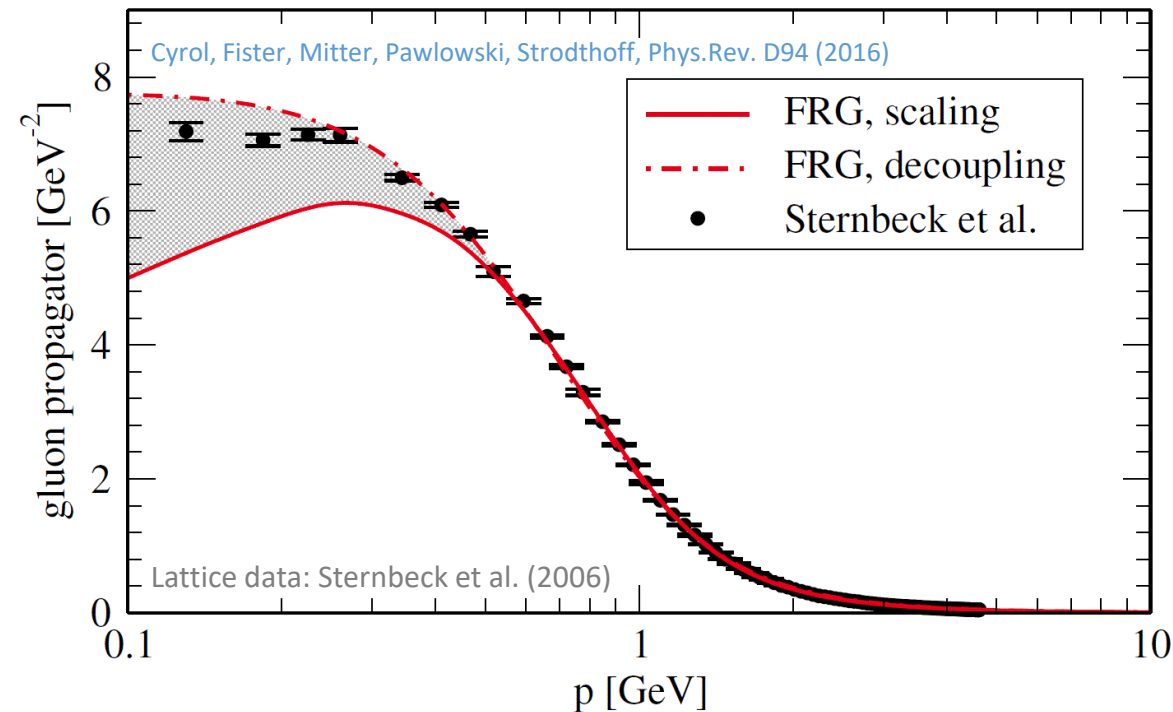
$$\int_{\mu>0} d\mu \mu \rho(\mu) = 0$$

3. UV behavior: Perturbation theory

$$\rho(\omega) \sim -\frac{Z_{UV}}{\omega^2 \ln(\omega^2)^{1+\gamma}}$$

4. IR behavior: **new**

Use as much prior knowledge as possible!



Cyrol, Pawłowski, Rothkopf, NW arxiv:1804.today

# Reconstructing spectral functions

Infrared behavior of spectral functions

Start from

$$G(p_0) = \int_0^\infty \frac{d\lambda}{\pi} \frac{\lambda \rho(\lambda)}{\lambda^2 + p_0^2}$$

# Reconstructing spectral functions

Infrared behavior of spectral functions

Start from

$$G(p_0) = \int_0^\infty \frac{d\lambda}{\pi} \frac{\lambda \rho(\lambda)}{\lambda^2 + p_0^2}$$

Derivative w.r.t.  $p_0$



$$\partial_{p_0} G(p_0) = - \int_{-\infty}^\infty \frac{d\lambda}{\pi} \lambda p_0 \frac{\rho(\lambda)}{(\lambda^2 + p_0^2)^2}$$

# Reconstructing spectral functions

Infrared behavior of spectral functions

Start from

$$G(p_0) = \int_0^\infty \frac{d\lambda}{\pi} \frac{\lambda \rho(\lambda)}{\lambda^2 + p_0^2}$$

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Limit  $p_0 \rightarrow 0$

$$\lim_{p_0 \rightarrow 0^+} \partial_{p_0} G(p_0) = -\frac{1}{2} \lim_{\omega \rightarrow 0^+} \partial_\omega \rho(\omega)$$

# Reconstructing spectral functions

Infrared behavior of spectral functions

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$$G(p_0) = \int_0^\infty \frac{d\lambda}{\pi} \frac{\lambda \rho(\lambda)}{\lambda^2 + p_0^2}$$

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Limit  $p_0 \rightarrow 0$

Gluon:

Scaling solution to Yang-Mills

$$\hat{G}_A^{(\text{sca})}(p_0) \sim Z_{\text{IR}} (\hat{p}_0^2)^{-1+2\kappa}$$

$$\lim_{p_0 \rightarrow 0^+} \partial_{p_0} G(p_0) = -\frac{1}{2} \lim_{\omega \rightarrow 0^+} \partial_\omega \rho(\omega)$$

# Reconstructing spectral functions

Infrared behavior of spectral functions

Start from

$$G(p_0) = \int_0^\infty \frac{d\lambda}{\pi} \frac{\lambda \rho(\lambda)}{\lambda^2 + p_0^2}$$

Derivative w.r.t.  $p_0$



$$\partial_{p_0} G(p_0) = - \int_{-\infty}^\infty \frac{d\lambda}{\pi} \lambda p_0 \frac{\rho(\lambda)}{(\lambda^2 + p_0^2)^2}$$



Limit  $p_0 \rightarrow 0$

Gluon:

Scaling solution to Yang-Mills

$$\hat{G}_A^{(\text{sca})}(p_0) \sim Z_{\text{IR}} (\hat{p}_0^2)^{-1+2\kappa}$$

$$\hat{\rho}_A^{(\text{sca})}(\omega) \sim -2 Z_{\text{IR}} \text{sgn}(\hat{\omega}) (\hat{\omega}^2)^{-1+2\kappa}$$

$$\lim_{p_0 \rightarrow 0^+} \partial_{p_0} G(p_0) = -\frac{1}{2} \lim_{\omega \rightarrow 0^+} \partial_\omega \rho(\omega)$$

Gluon spectral function is negative for small frequencies



# Reconstructing spectral functions

Infrared behavior of spectral functions

Start from

$$G(p_0) = \int_0^\infty \frac{d\lambda}{\pi} \frac{\lambda \rho(\lambda)}{\lambda^2 + p_0^2}$$

Derivative w.r.t.  $p_0$



$$\partial_{p_0} G(p_0) = - \int_{-\infty}^\infty \frac{d\lambda}{\pi} \lambda p_0 \frac{\rho(\lambda)}{(\lambda^2 + p_0^2)^2}$$



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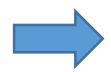
Gluon spectral function is negative for small frequencies

Similar result for decoupling scenarios,  
for more details see

Cyrol, Pawłowski, Rothkopf, NW arxiv:1804.today

## Reconstructing spectral functions

Constructing a basis for the reconstruction:



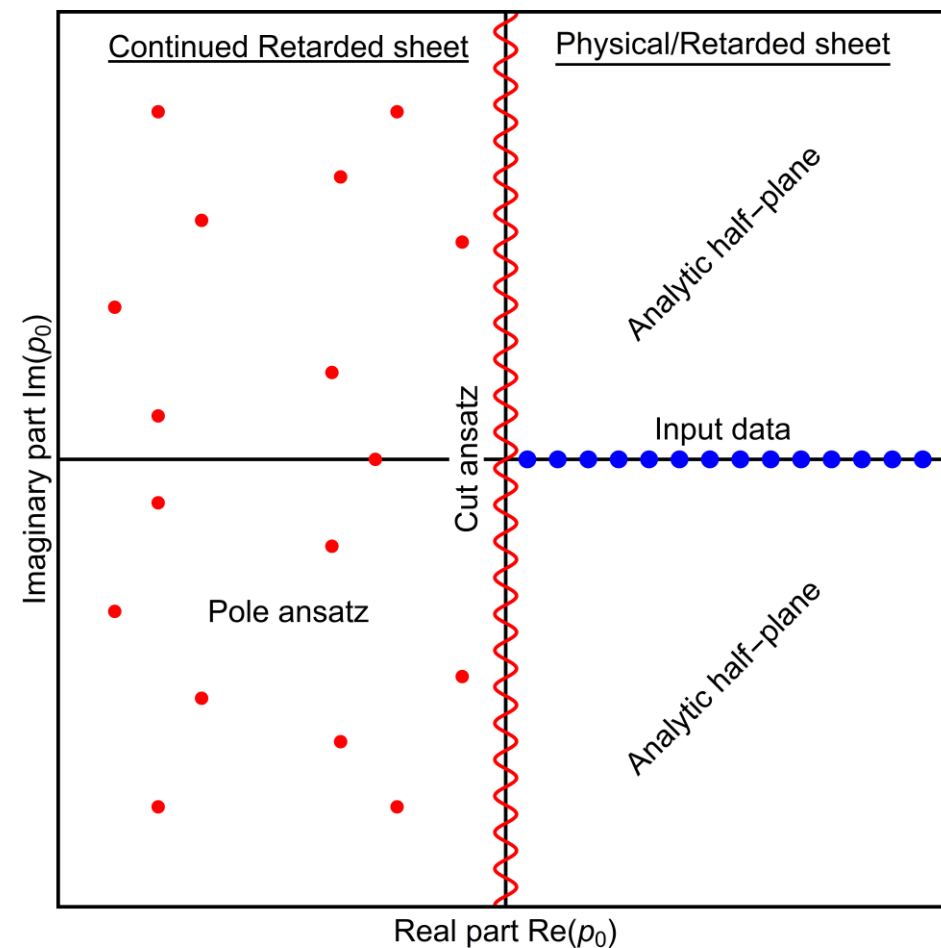
Analytically continue the retarded propagator to the entire complex plane

# Reconstructing spectral functions

Constructing a basis for the reconstruction:

➔ Analytically continue the retarded propagator to the entire complex plane

$$\rho(\omega, \vec{p}) = 2 \lim_{\varepsilon \rightarrow 0} \text{Im} G_E(-i(\omega + i\varepsilon), \vec{p})$$



Cyrol, Pawłowski, Rothkopf, NW arxiv:1804.today

# Reconstructing spectral functions

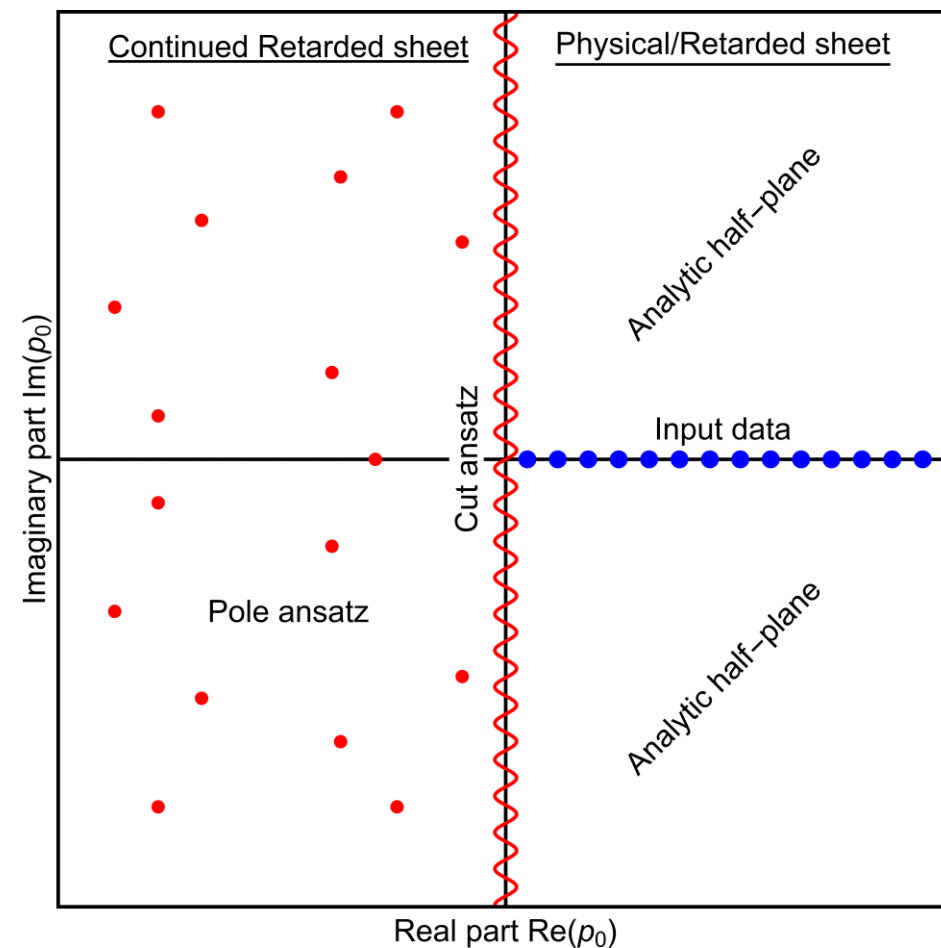
Constructing a basis for the reconstruction:

➔ Analytically continue the retarded propagator to the entire complex plane

Poles

$$\hat{G}_{\text{Ans}}^{\text{pole}}(p_0) = \sum_{k=1}^{N_{\text{ps}}} \prod_{j=1}^{N_{\text{pp}}^{(k)}} \left( \frac{\hat{N}_k}{(\hat{p}_0 + \hat{\Gamma}_{k,j})^2 + \hat{M}_{k,j}^2} \right)^{\delta_{k,j}}$$

$$\rho(\omega, \vec{p}) = 2 \lim_{\varepsilon \rightarrow 0} \text{Im} G_E(-i(\omega + i\varepsilon), \vec{p})$$



# Reconstructing spectral functions

Constructing a basis for the reconstruction:

➔ Analytically continue the retarded propagator to the entire complex plane

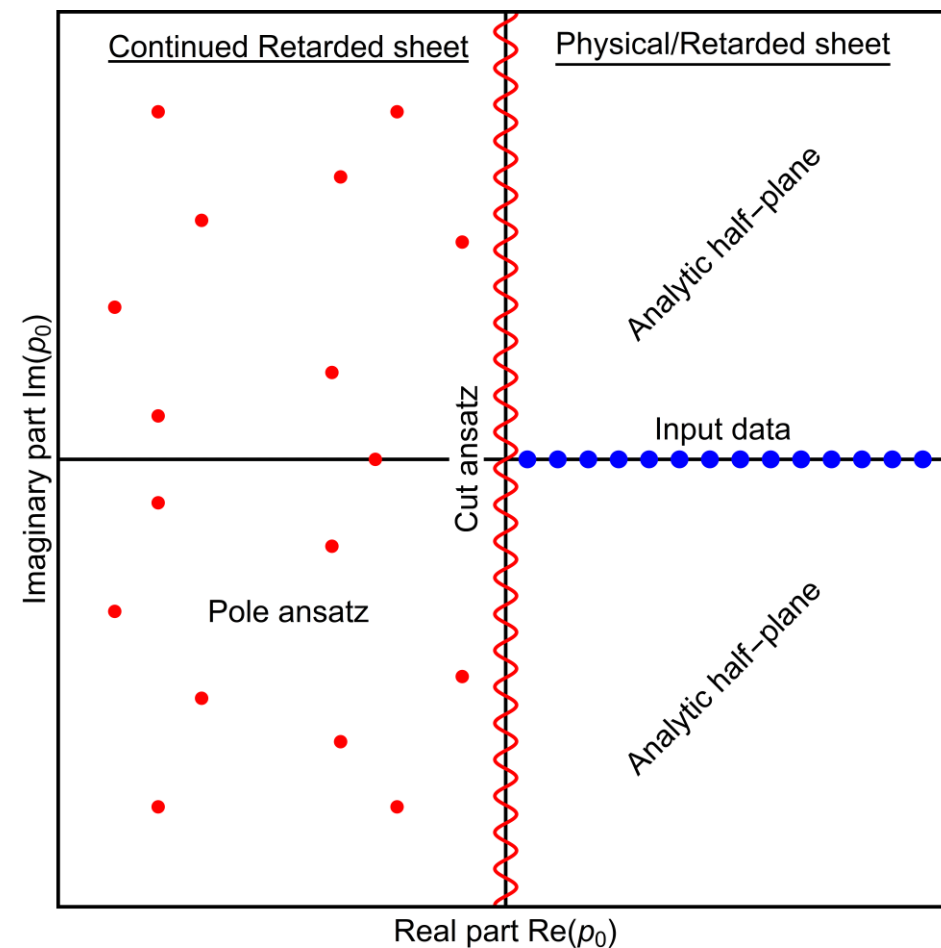
Poles

$$\hat{G}_{\text{Ans}}^{\text{pole}}(p_0) = \sum_{k=1}^{N_{\text{ps}}} \prod_{j=1}^{N_{\text{pp}}^{(k)}} \left( \frac{\hat{N}_k}{(\hat{p}_0 + \hat{\Gamma}_{k,j})^2 + \hat{M}_{k,j}^2} \right)^{\delta_{k,j}}$$

$$\hat{G}_{\text{Ans}}^{\text{poly}}(p_0) = \sum_{j=1}^{N_{\text{poly}}} \hat{a}_k (\hat{p}_0^2)^{\frac{j}{2}}$$

Polynomial

$$\rho(\omega, \vec{p}) = 2 \lim_{\varepsilon \rightarrow 0} \text{Im } G_E(-i(\omega + i\varepsilon), \vec{p})$$



# Reconstructing spectral functions

$$\rho(\omega, \vec{p}) = 2 \lim_{\varepsilon \rightarrow 0} \text{Im } G_E(-i(\omega + i\varepsilon), \vec{p})$$

Constructing a basis for the reconstruction:

➔ Analytically continue the retarded propagator to the entire complex plane

Poles

$$\hat{G}_{\text{Ans}}^{\text{pole}}(p_0) = \sum_{k=1}^{N_{\text{ps}}} \prod_{j=1}^{N_{\text{pp}}^{(k)}} \left( \frac{\hat{N}_k}{(\hat{p}_0 + \hat{\Gamma}_{k,j})^2 + \hat{M}_{k,j}^2} \right)^{\delta_{k,j}}$$

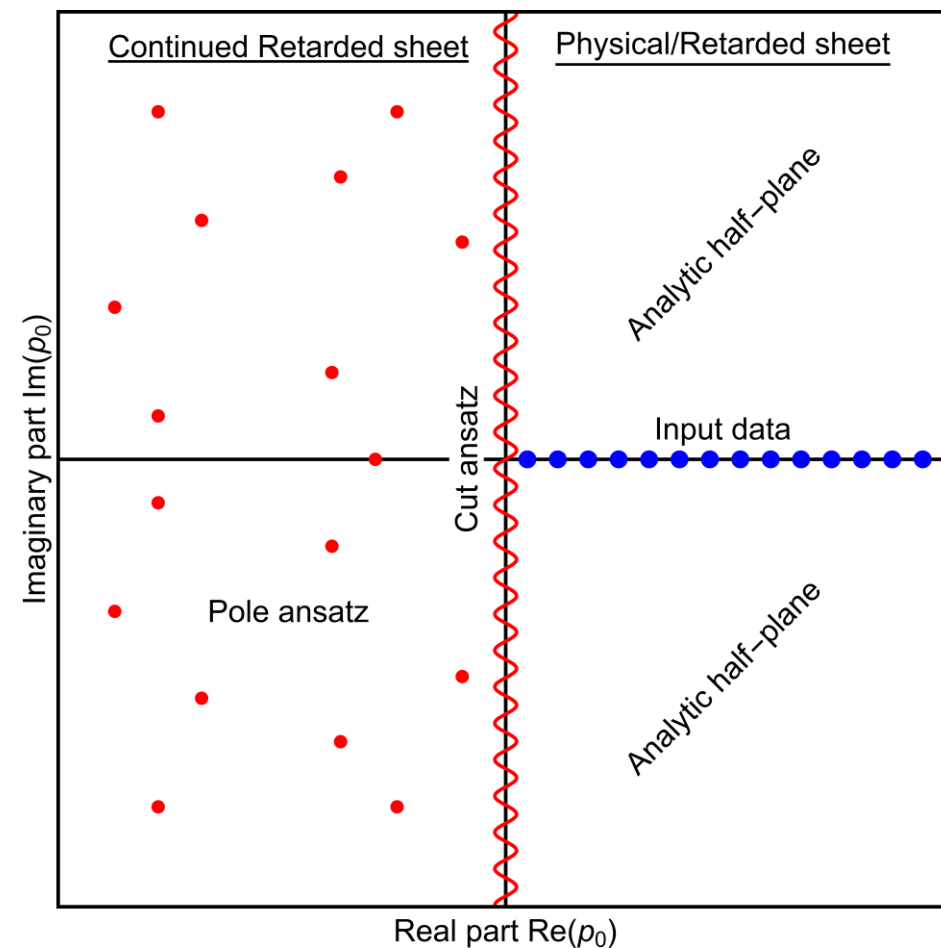
$$\hat{G}_{\text{Ans}}^{\text{poly}}(p_0) = \sum_{j=1}^{N_{\text{poly}}} \hat{a}_k (\hat{p}_0^2)^{\frac{j}{2}}$$

Polynomial

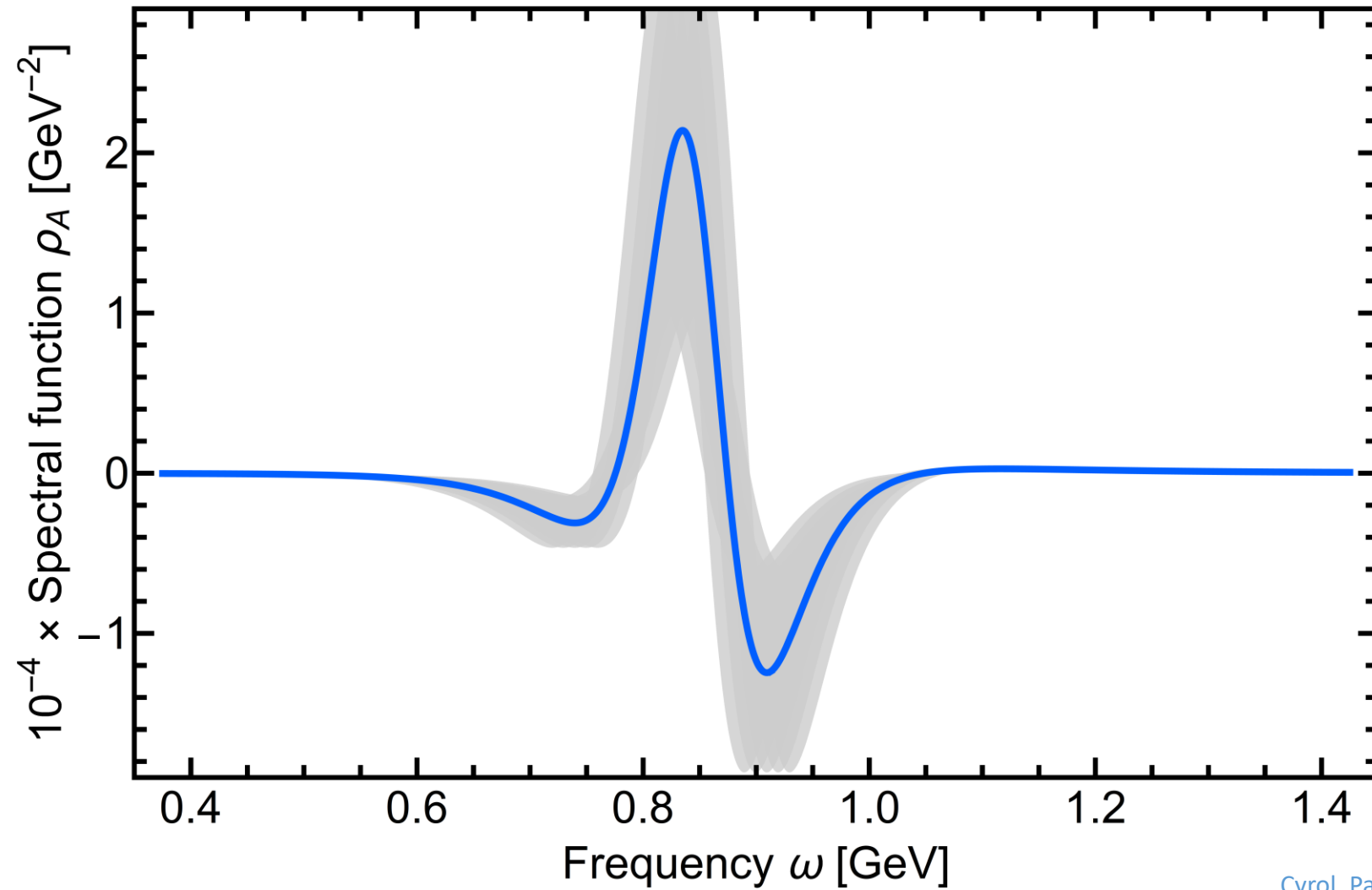
$$\hat{G}_{\text{Ans}}^{\text{asy}}(p_0) = (\hat{p}_0^2)^{-1-2\alpha} \left[ \log \left( 1 + \frac{\hat{p}_0^2}{\hat{\lambda}^2} \right) \right]^{-1-\beta}$$

Cuts

Full ansatz  $G_{\text{Ans}}(p_0) = \mathcal{K} \hat{G}_{\text{Ans}}^{\text{pole}}(p_0) \hat{G}_{\text{Ans}}^{\text{poly}}(p_0) \hat{G}_{\text{Ans}}^{\text{asy}}(p_0)$

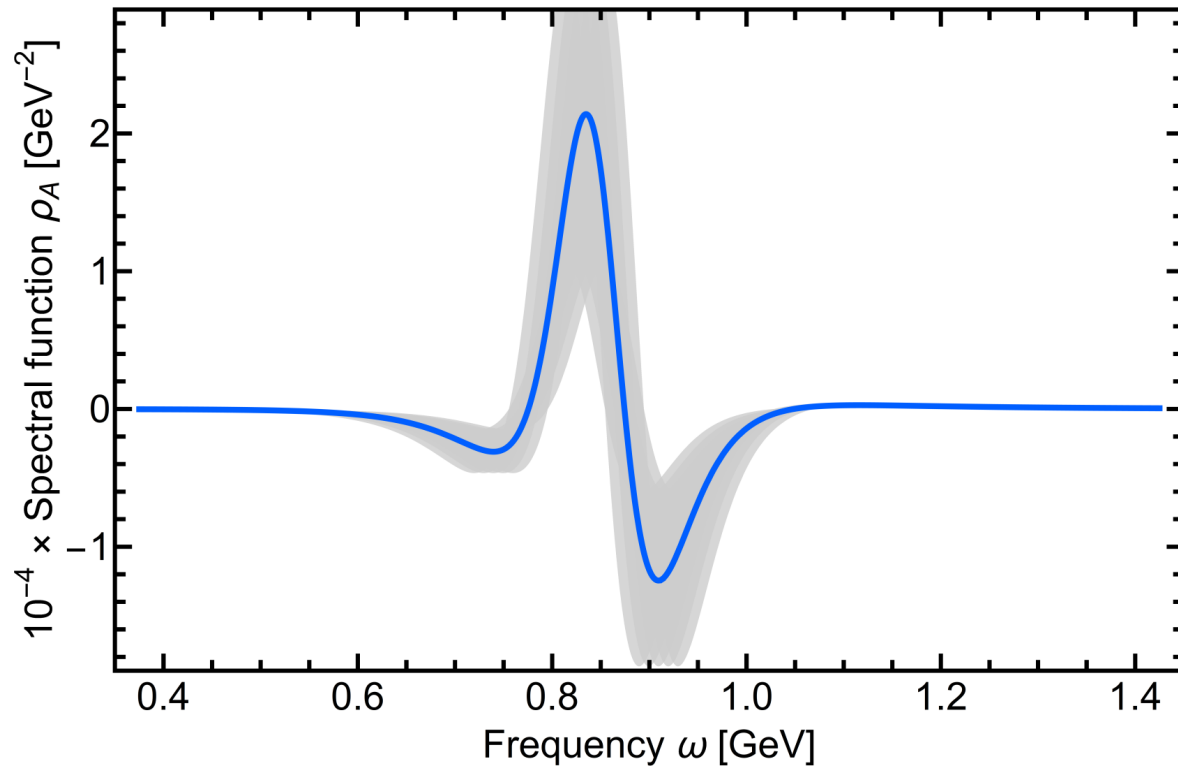


## Gluon spectral function

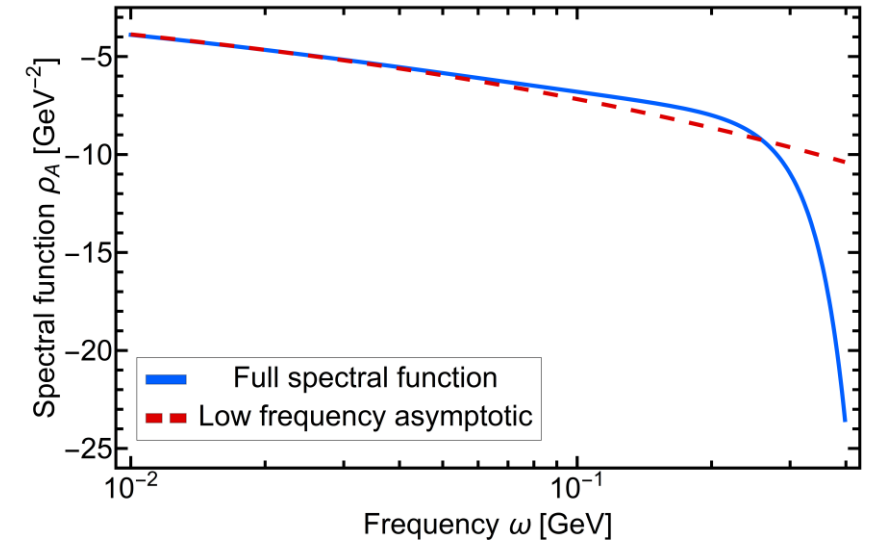


Cyrol, Pawlowski, Rothkopf, NW arxiv:1804.today

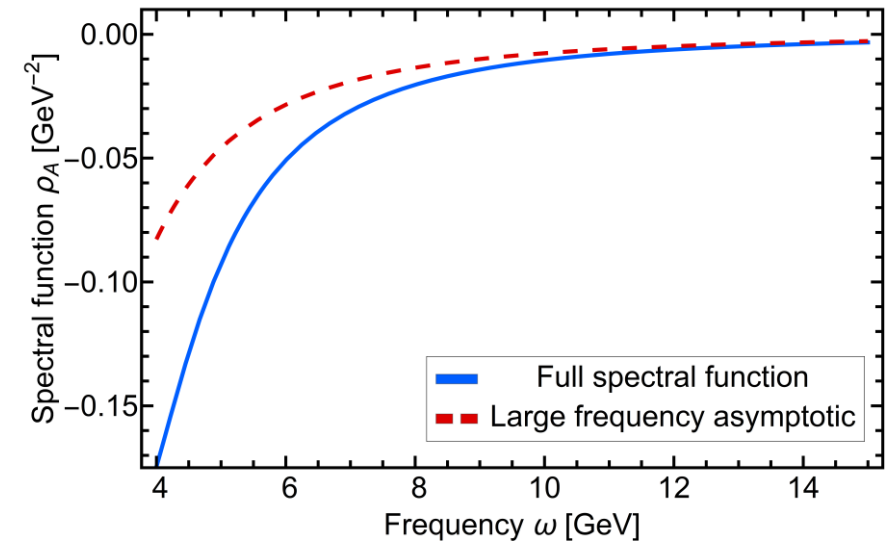
## Gluon spectral function



IR



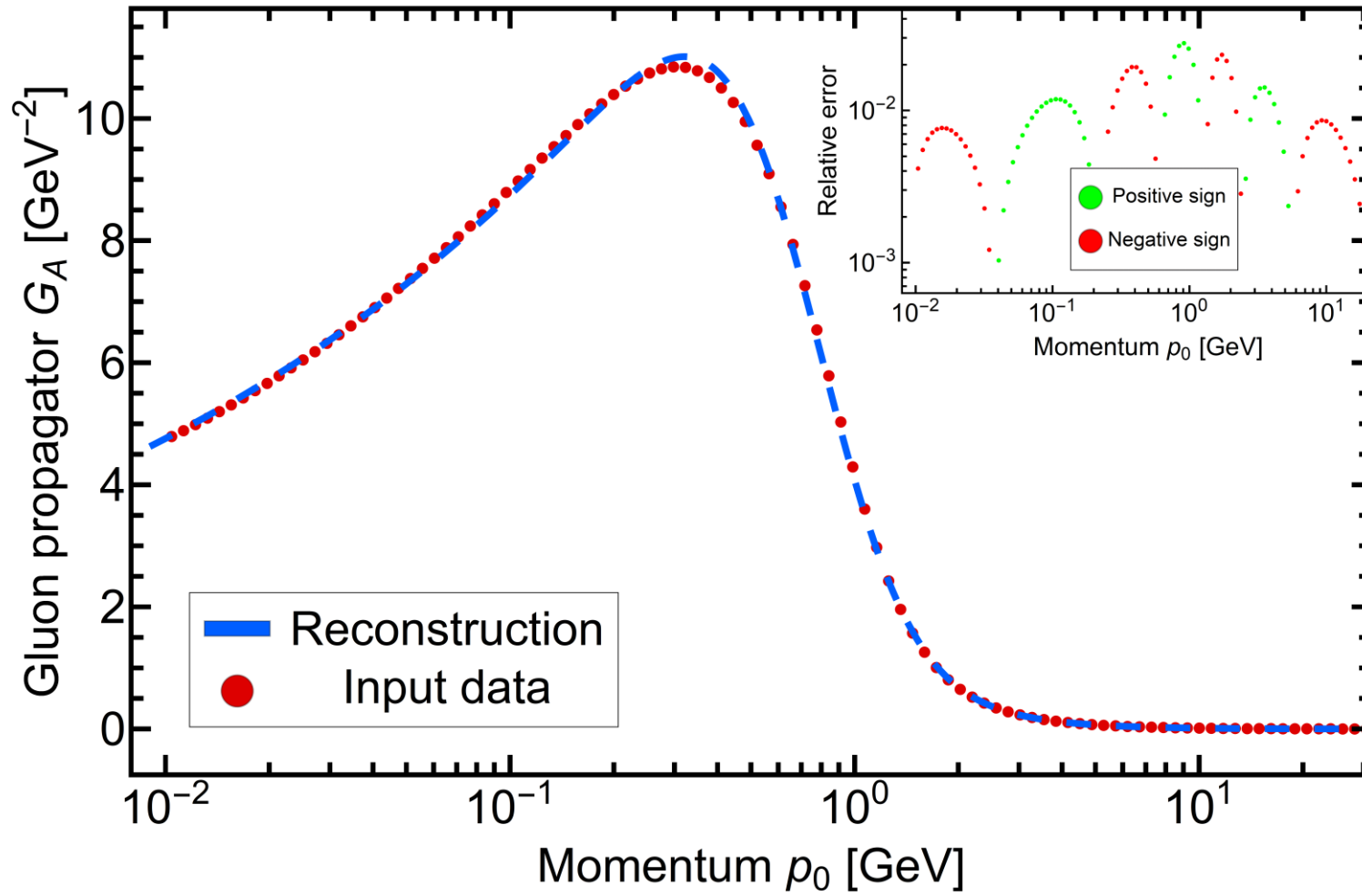
UV



Cyrol, Pawlowski, Rothkopf, NW arxiv:1804.today



## Gluon spectral function



## Summary & Outlook

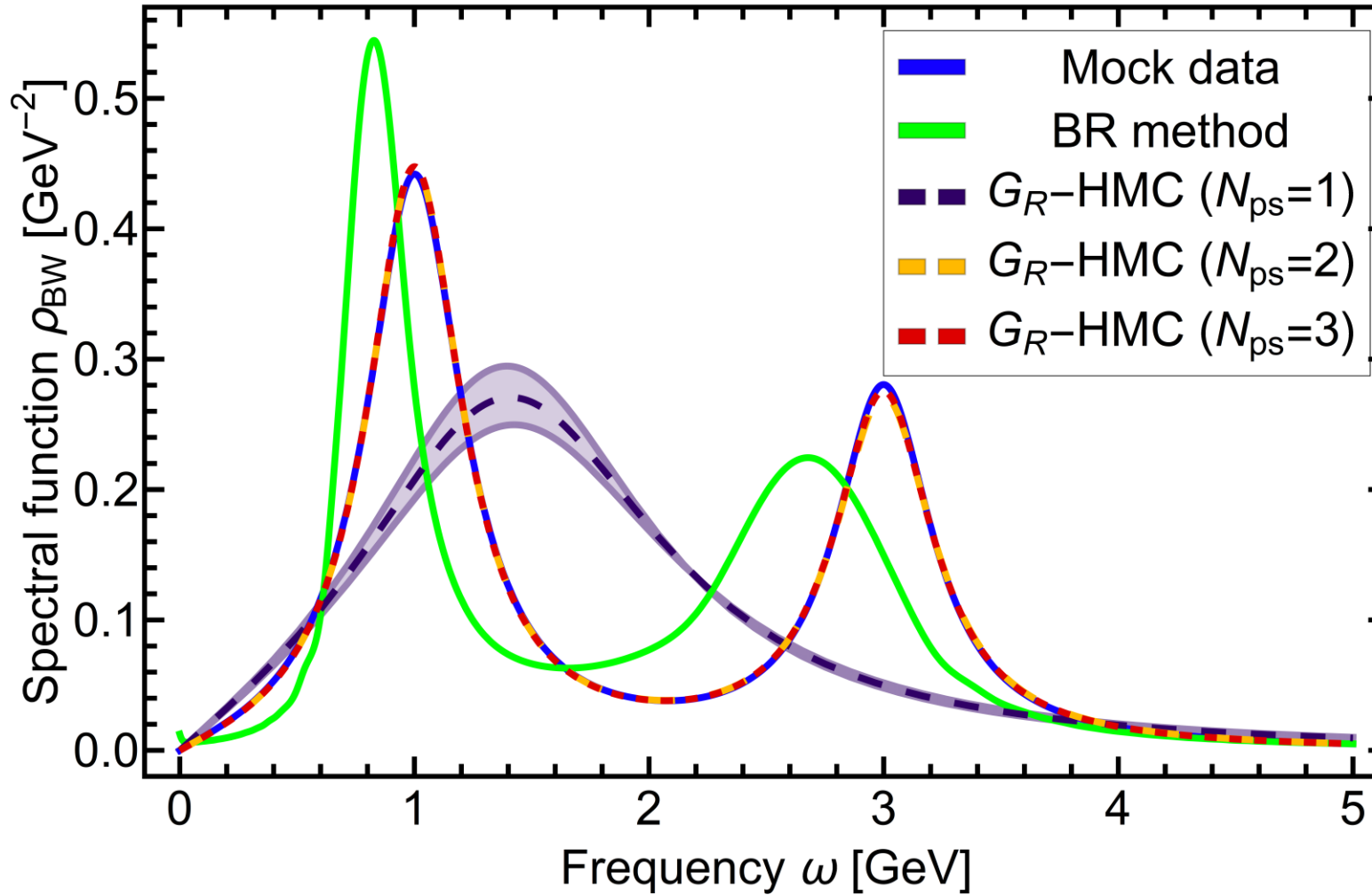
- Direct calculation of spectral functions in functional methods
- Reconstruction of the gluon spectral function with all priors

Thank you for your attention!

- Finite temperature gluon spectral functions
- Transport coefficients

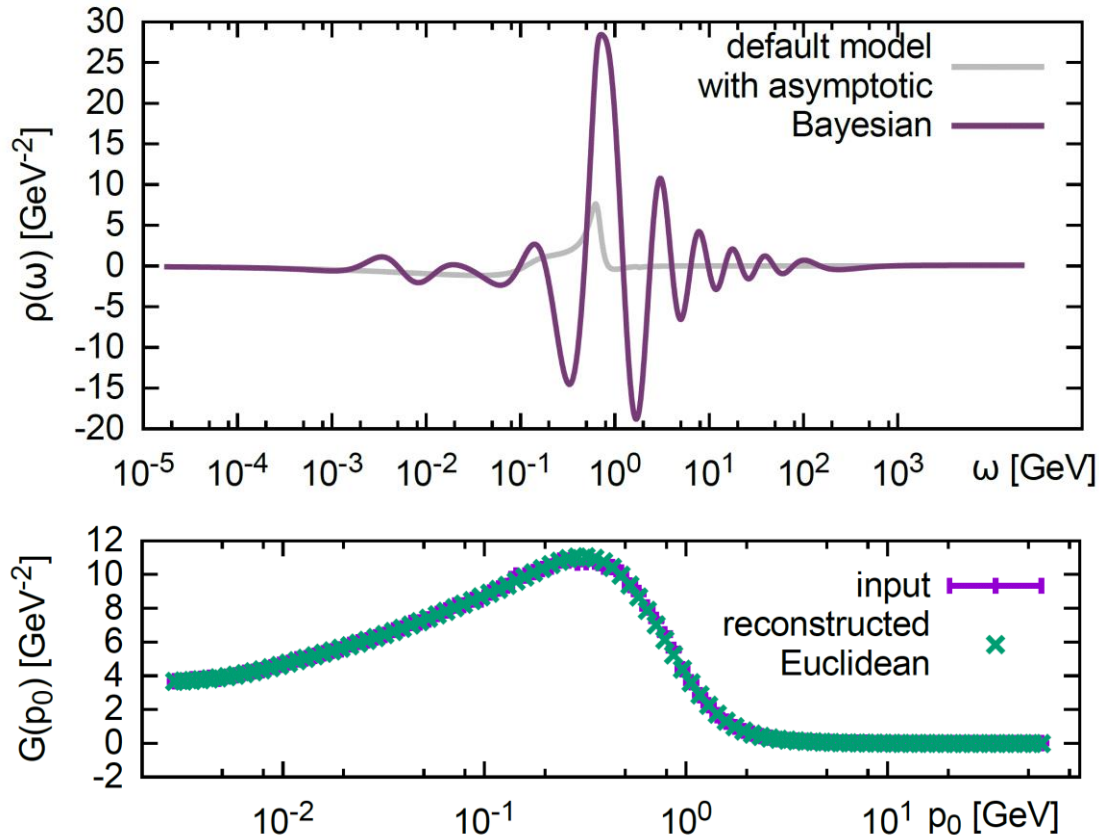
Backup slides

# Breit-Wigner benchmark

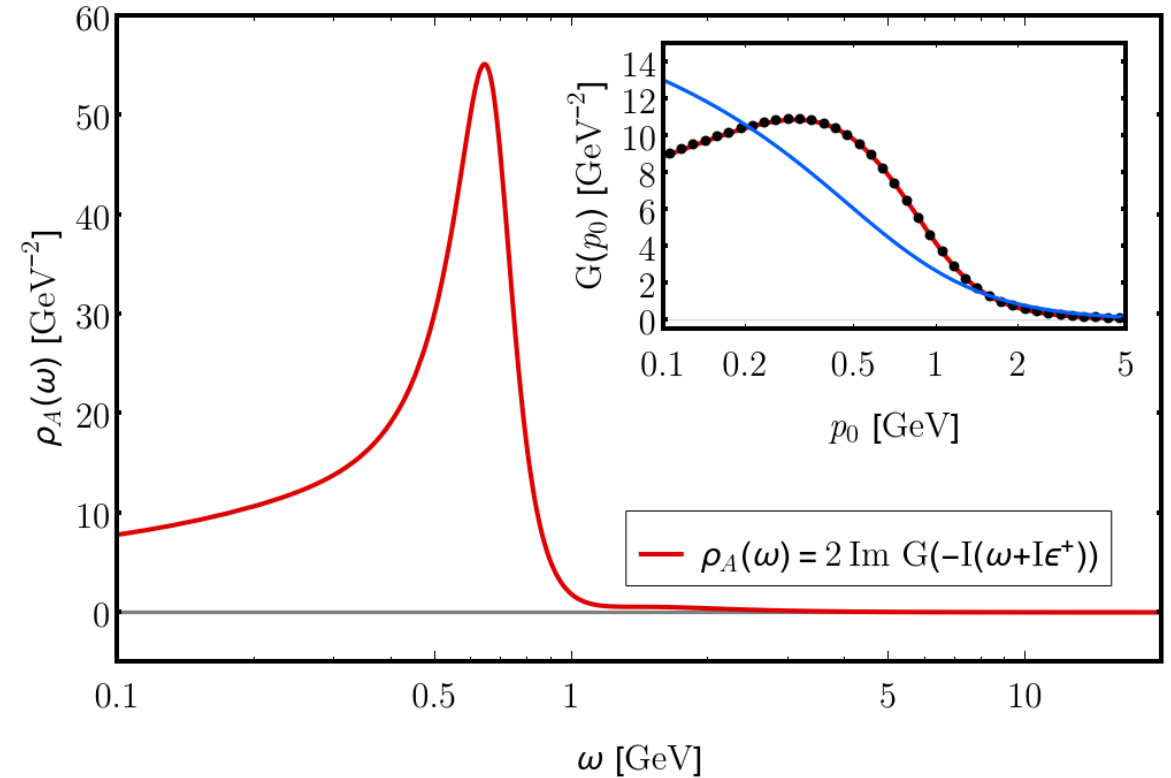


# Comparison with other works

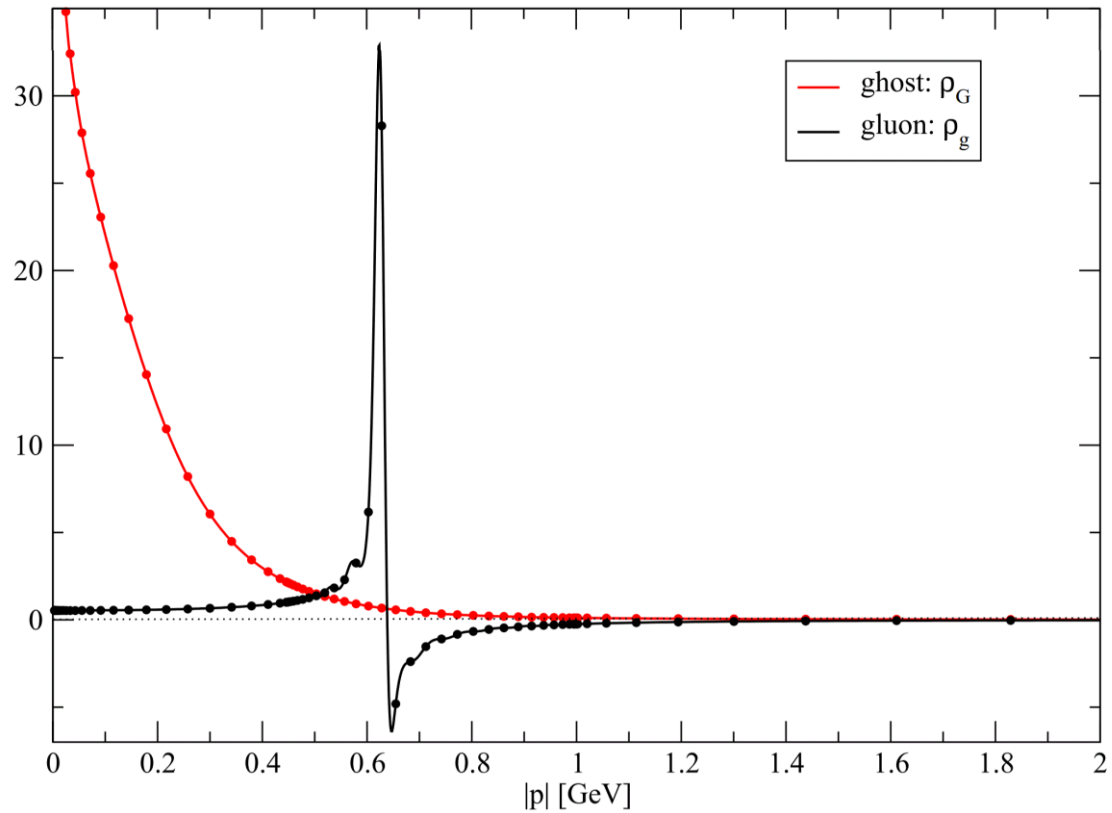
## Bayesian reconstruction



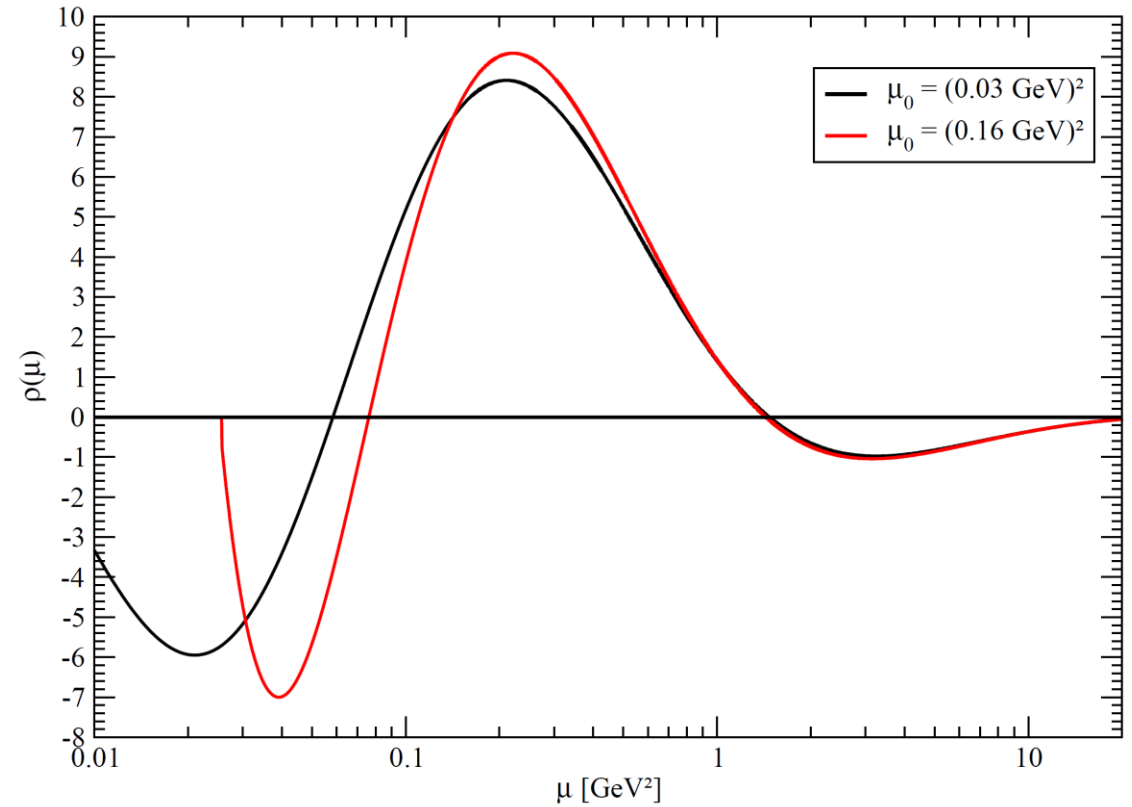
## Pade/Schlessinger point



# Comparison with other works



Strauss, Fischer, Kellermann, Phys.Rev.Lett. 109 (2012)

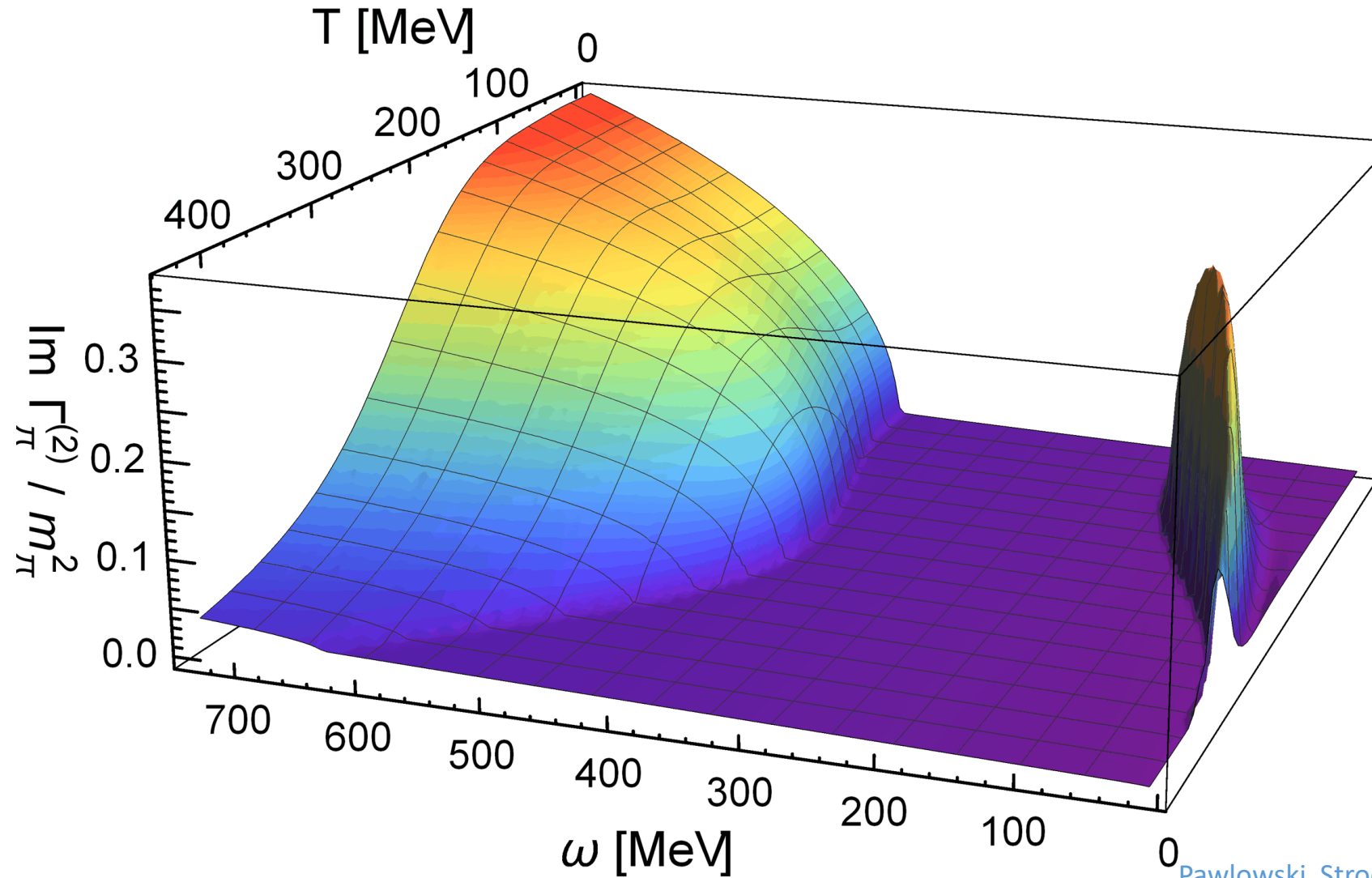


Dudal, Oliveira, Silva, Phys.Rev. D89 (2014)

## Application to the O(N)-Model

## Pion

Imaginary part of the retarded two-point function

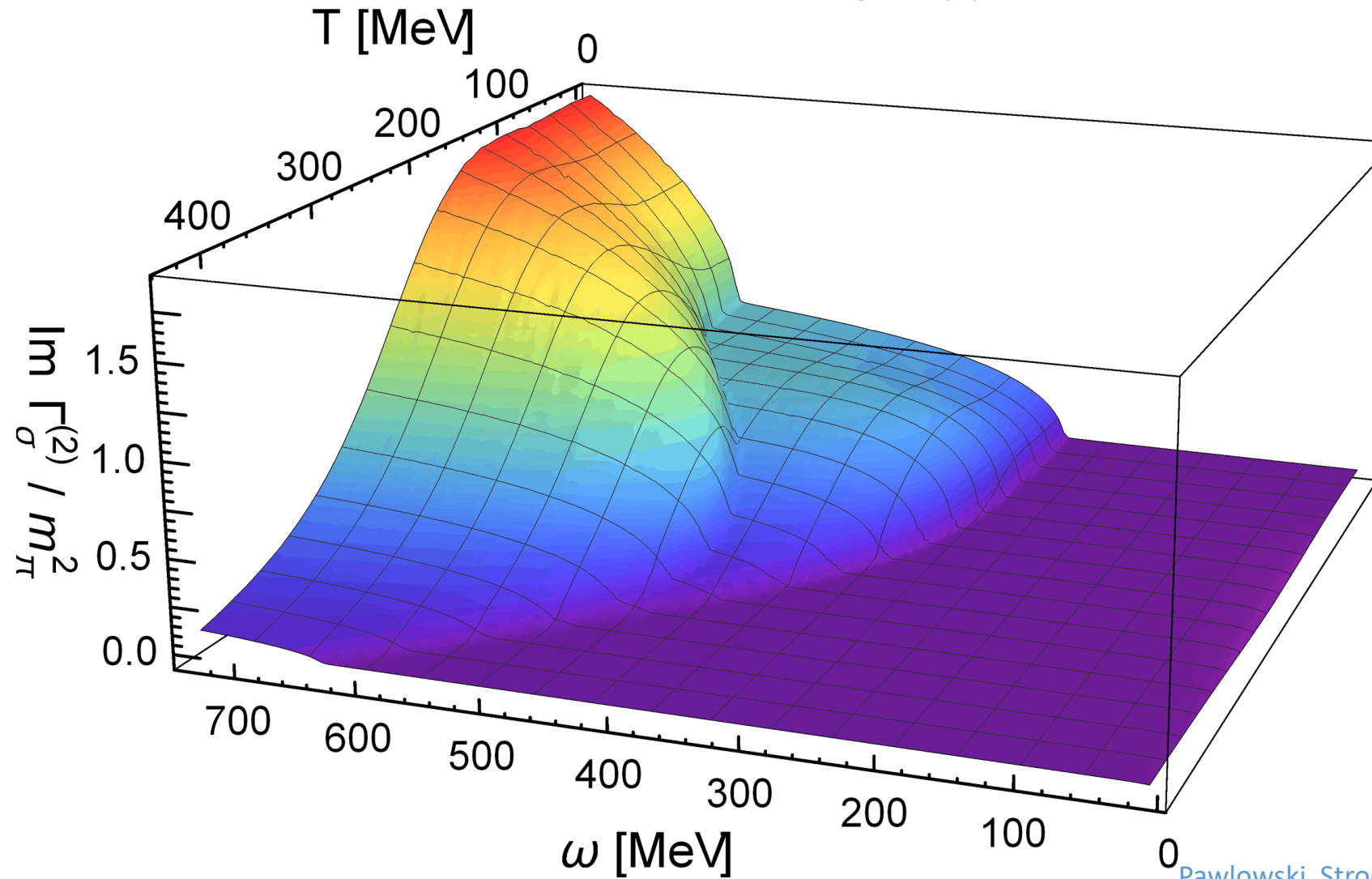


Pawlowski, Strodthoff, NW, arxiv:1711.07444

## Application to the O(N)-Model

## Sigma meson

Imaginary part of the retarded two-point function



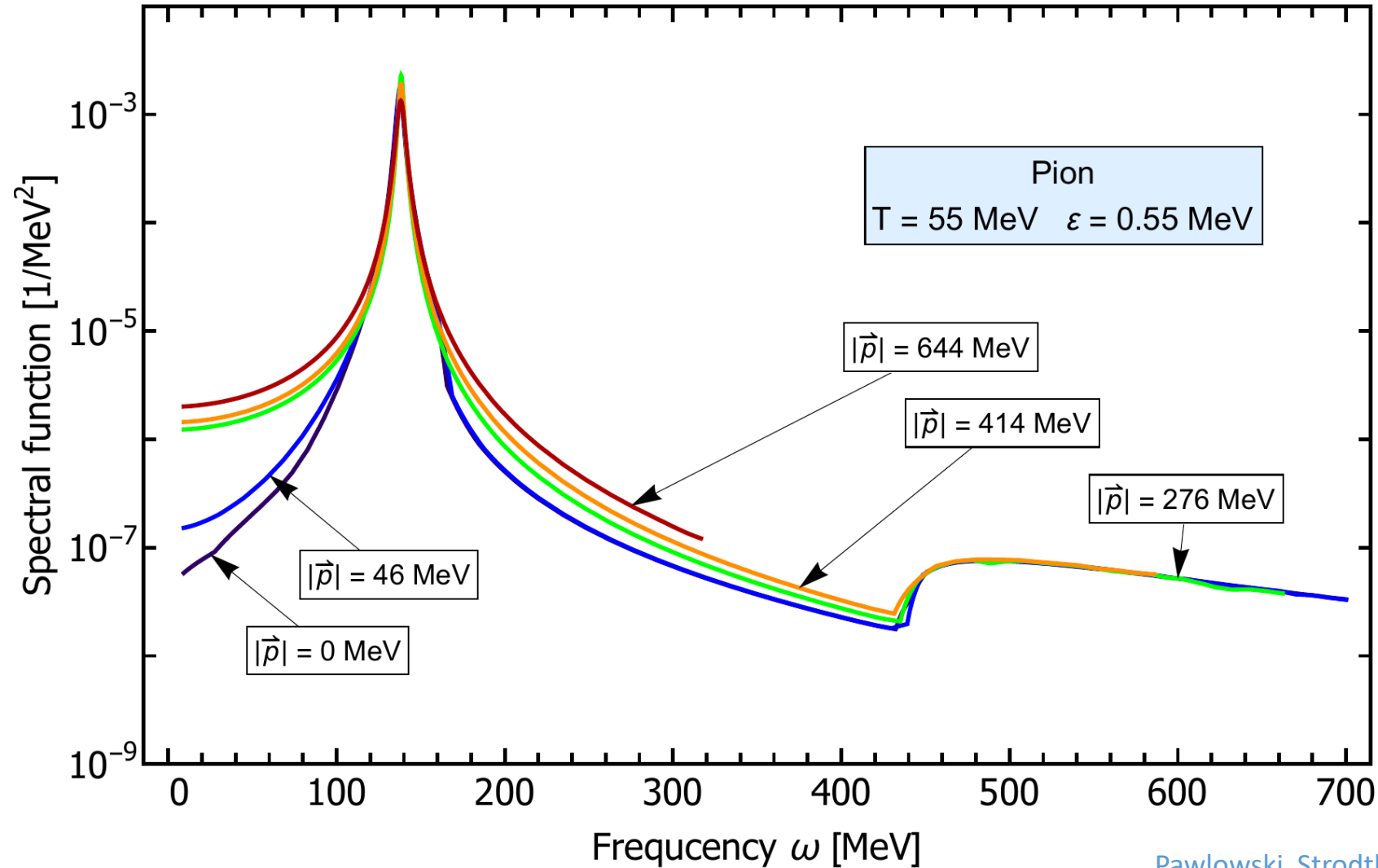
Pawlowski, Strodthoff, NW, arxiv:1711.07444



## Pion meson

## Application to the O(N)-Model

Finite temperature spectral function for various external momenta

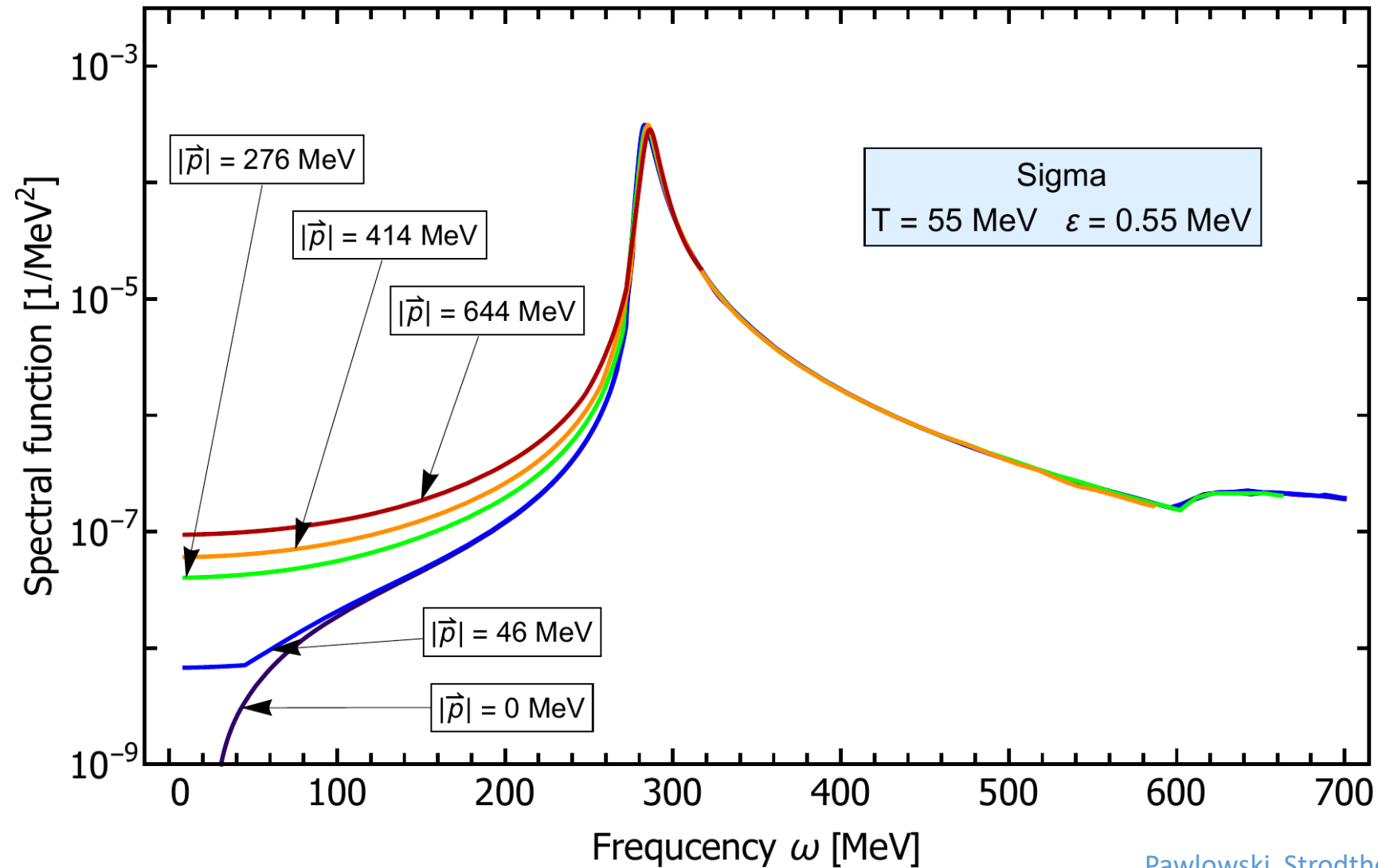


Pawlowski, Strodthoff, NW, arxiv:1711.07444

## Sigma meson

## Application to the O(N)-Model

Finite temperature spectral function for various external momenta

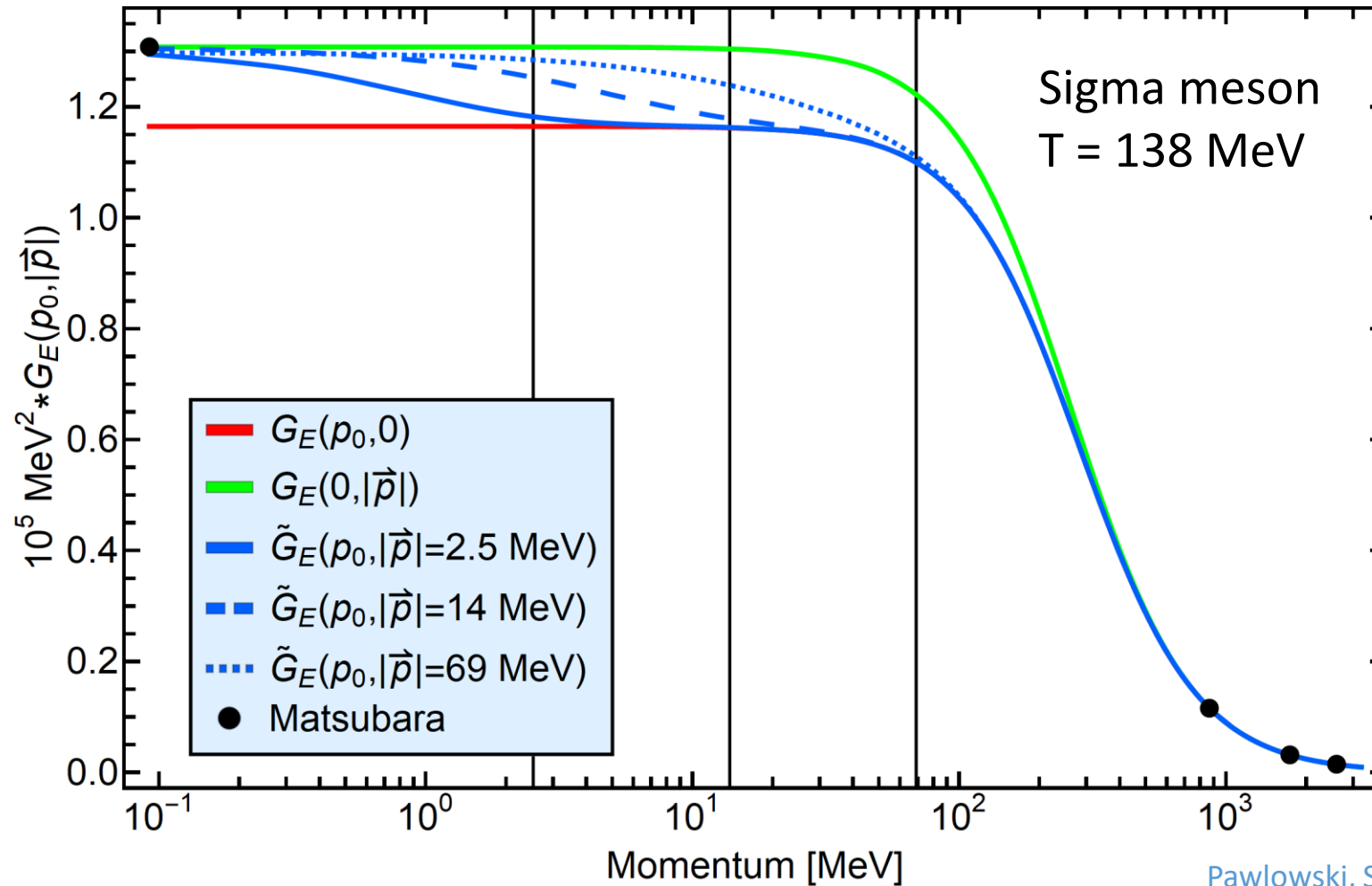


Pawlowski, Strodthoff, NW, arxiv:1711.07444

## Application to the O(N)-Model

In medium non-commuting limits

$$\lim_{\vec{p} \rightarrow 0} \lim_{p_0 \rightarrow 0} \Gamma^{(2)}(p_0, \vec{p}) \neq \lim_{p_0 \rightarrow 0} \lim_{\vec{p} \rightarrow 0} \Gamma^{(2)}(p_0, \vec{p})$$



Pawlowski, Strodthoff, NW, arxiv:1711.07444

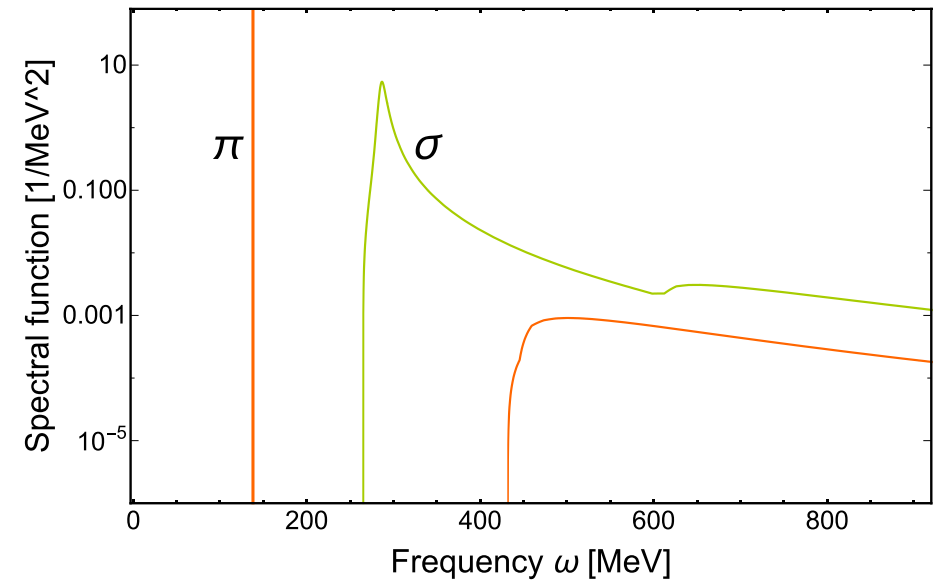
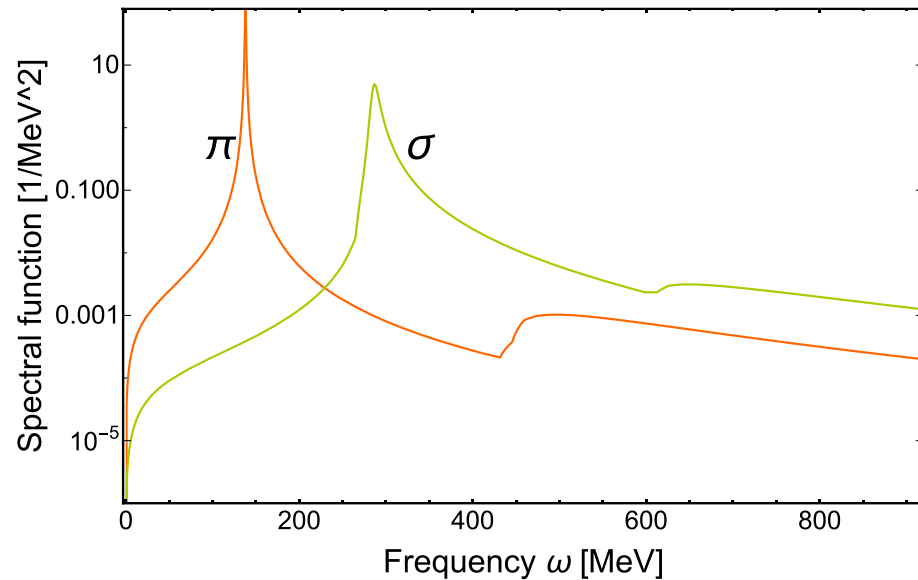
# Retarded Greens function

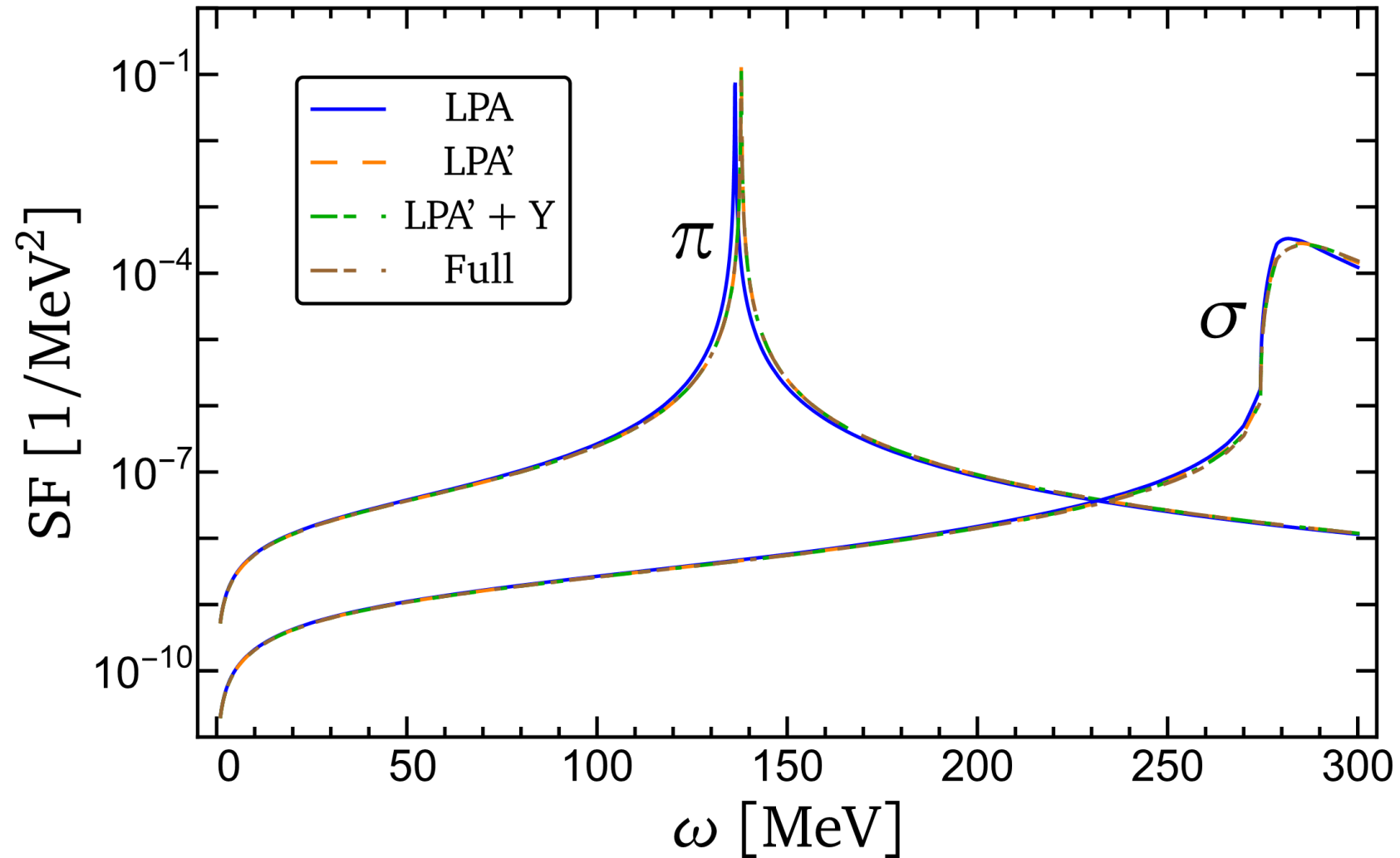
Retarded Greens function  $\lim_{\varepsilon \rightarrow 0} G(-i(\omega + i\varepsilon))$

Take limit analytically

Numerical extrapolation

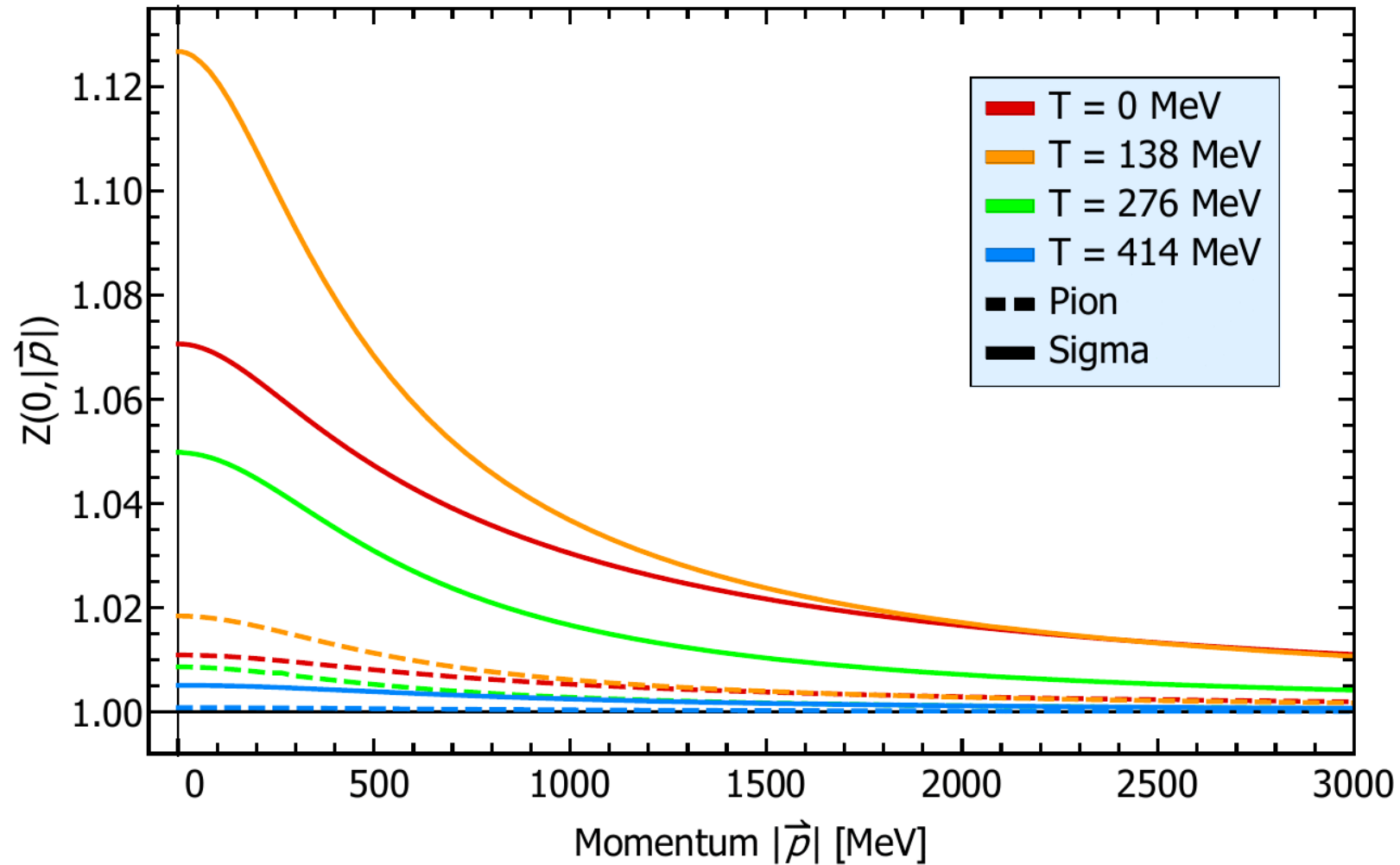
$$\rho(\omega, \vec{p}) = -2 \operatorname{Im} G_{\text{R}}(\omega, \vec{p})$$





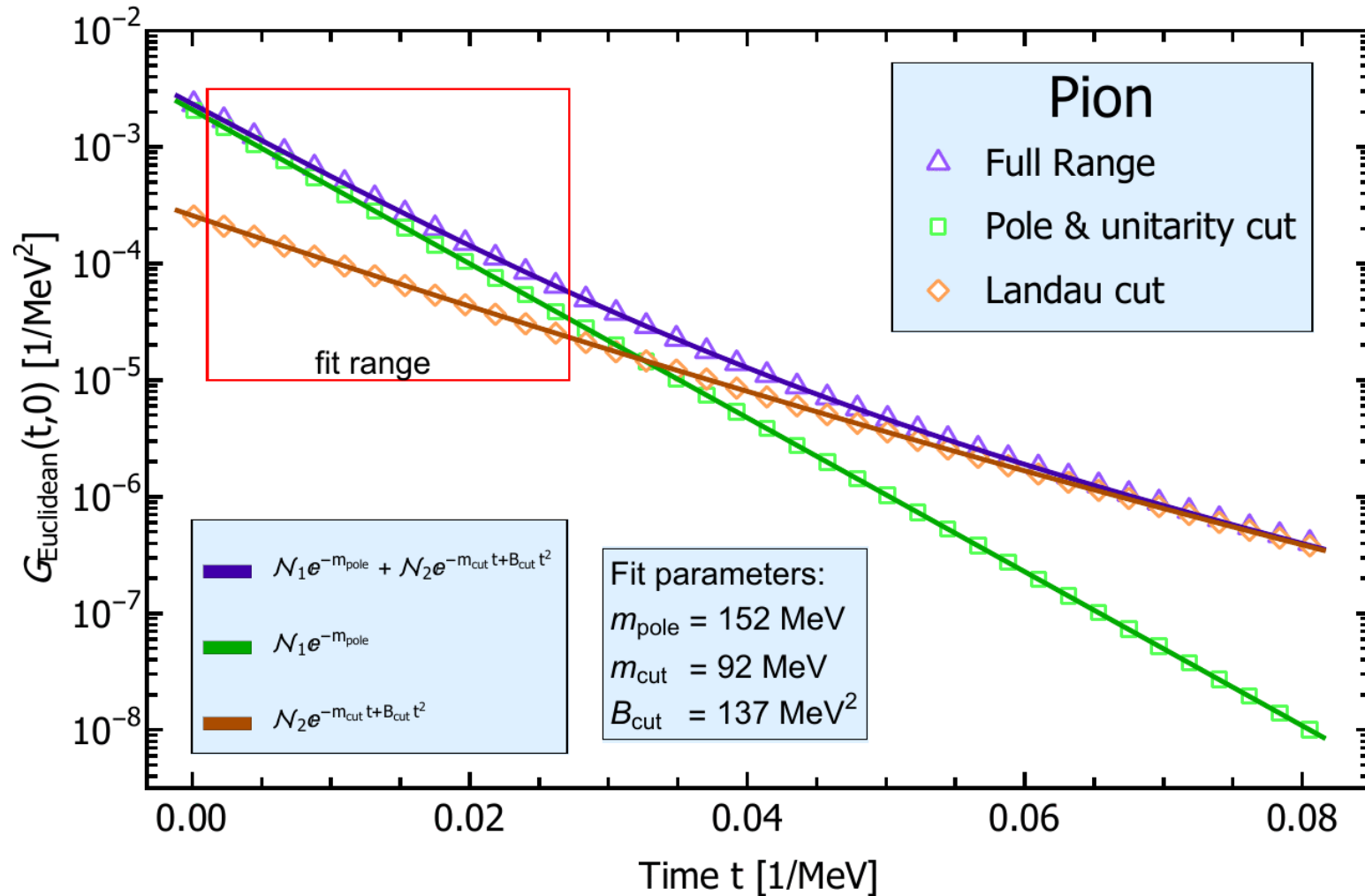
## Application to the O(N)-Model

## Euclidean momentum dependent dressings



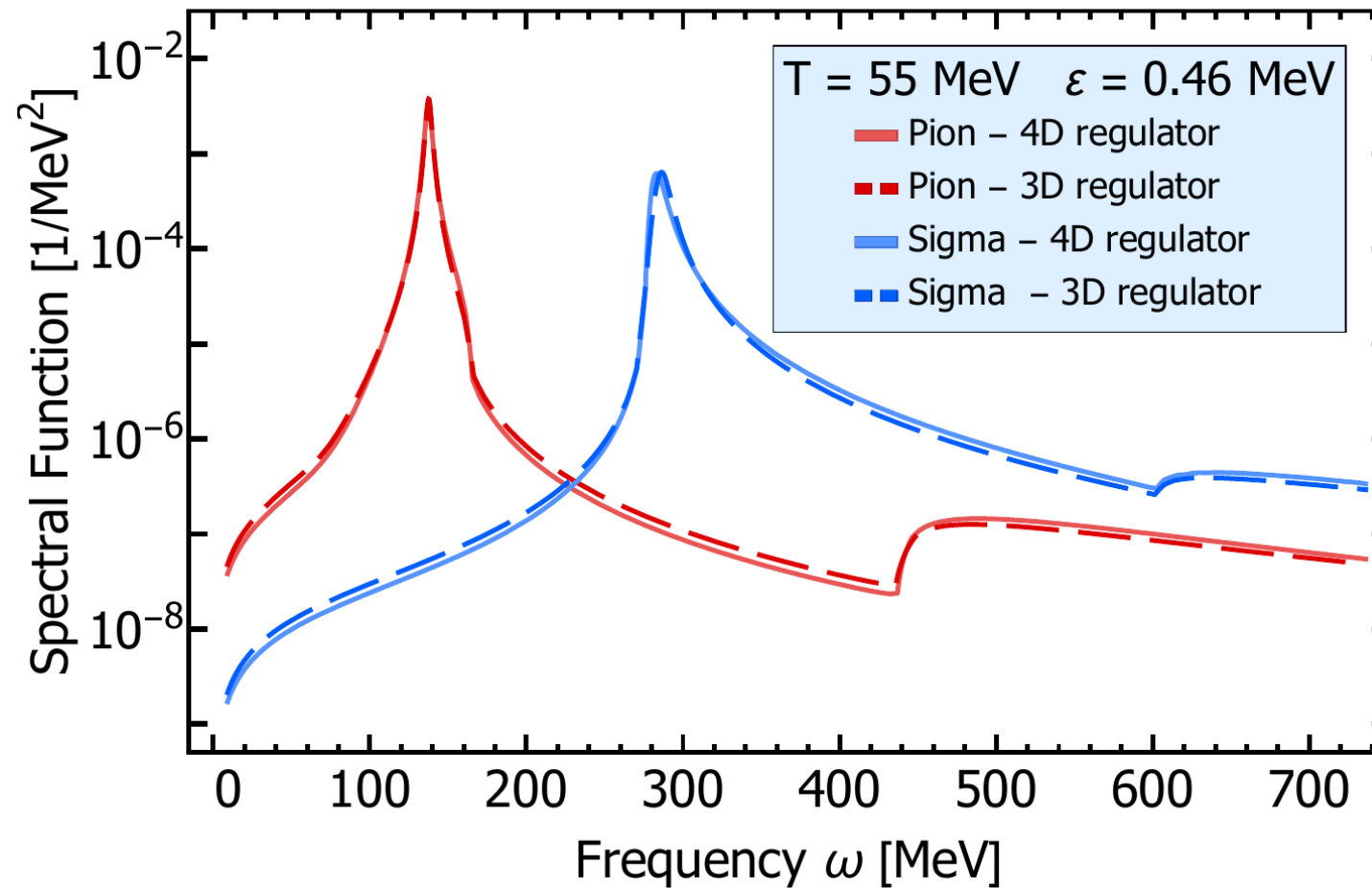
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## Application to the O(N)-Model



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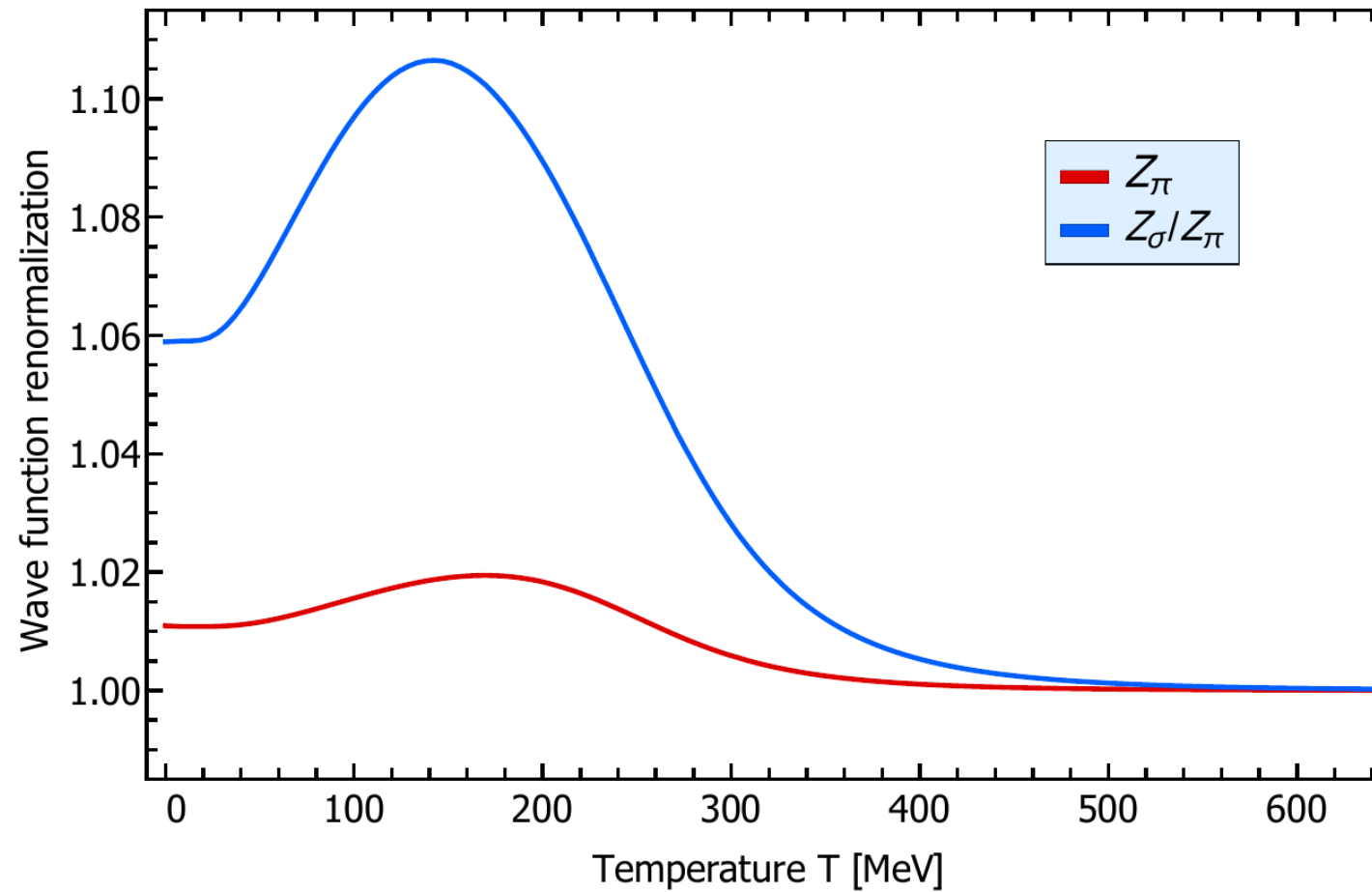
## Application to the O(N)-Model



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## Application to the O(N)-Model



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$$\partial_t \Gamma_k[\Phi] = \frac{1}{2} \left( \text{orange coiled circle with } \otimes \text{ at top} - \text{dashed circle with } \otimes \text{ at top} \right)$$