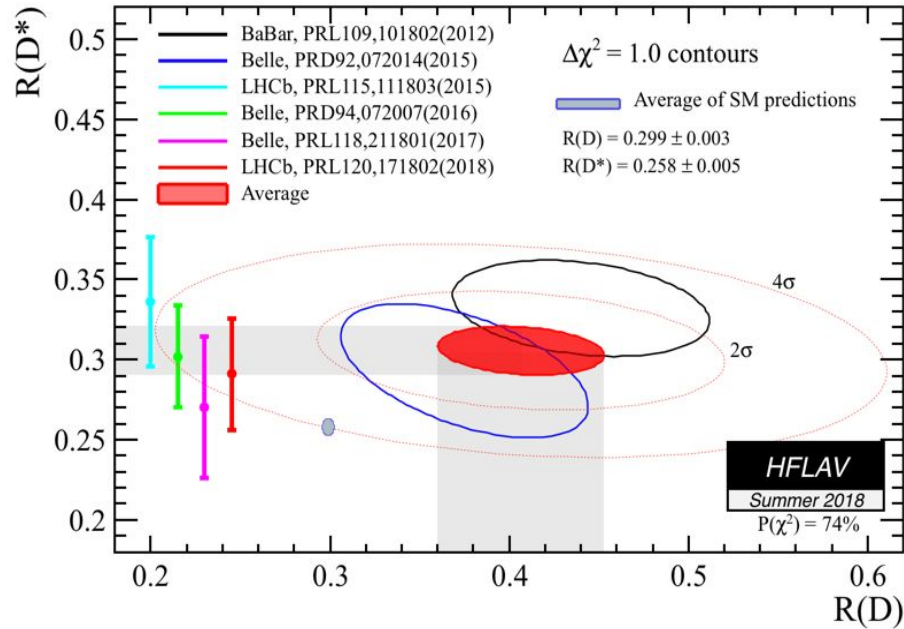


# Flavour Anomalies

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# Introduction to Flavour Physics

Flavour, Universality & Tests

# The Standard Model

- Gauge Group:  $G_{\text{SM}} = SU(3)_C \times SU(2)_L \times U(1)_Y$
- Lagrangian:  $\mathcal{L}_{\text{SM}} = \mathcal{L}_{\text{kin}} + \mathcal{L}_{\text{Higgs}} + \mathcal{L}_{\text{Yukawa}}$
- Fermions: (in 3 generations  $i = 1, 2, 3$ )

$$Q_{Li} (3, 2)_{+1/6}, \quad U_{Ri} (3, 1)_{+2/3}, \quad D_{Ri} (3, 1)_{-1/3}, \quad L_{Li} (1, 2)_{-1/2}, \quad E_{Ri} (1, 1)_{-1}$$

with doublets  $Q_{Li} = (U_{Li}, D_{Li})$  &  $L_{Li} = (\nu_{Li}, E_{Li})$

- Higgs inducing SSB:  $SU(2)_L \times U(1)_Y \rightarrow U(1)_{EM}$   
 $\phi (1, 2)_{+1/2}, \quad \phi = \frac{1}{\sqrt{2}}(0, v + H)$

# Flavour (Physics)

- **Flavour = species of fermion**

in SM: 6 quark and 6 lepton flavours:  $u, d, c, s, t, b, e, \mu, \tau, \nu_e, \nu_\mu, \nu_\tau$

- Kinetic terms induce couplings of flavours to gauge bosons through gauge covariant derivative
- After SSB: (focus on  $W^\pm$ ,  $Z$  and on coupling of left-handed fermions)

$$-\mathcal{L}_{\text{kin}}^q \supset \frac{g}{\sqrt{2}} \bar{U}_{Li} \gamma^\mu \delta_{ij} D_{Lj} W^+ + \text{h.c.} \quad \text{flavour mixing}$$
$$+ \frac{g}{\sqrt{2}} \bar{U}_{Li} \gamma^\mu \delta_{ij} U_{Lj} Z + \frac{g}{\sqrt{2}} \bar{D}_{Li} \gamma^\mu \delta_{ij} D_{Lj} Z \quad \text{no generation mixing!?!}$$

$\Rightarrow W^\pm$  can induce **flavour change** (no flavour changing neutral current (FCNC) via  $Z$  or gluons or photon), same holds true for lepton kinetic term

# (Flavour) Universality

- (Flavour) **Universality = flavour-independent coupling** to all gauge bosons

⇒ for  $E \gg m$ , we have

$$\Gamma\left(Z \text{ wavy} \begin{array}{l} \nearrow e \\ \searrow \bar{e} \end{array}\right) = \Gamma\left(Z \text{ wavy} \begin{array}{l} \nearrow \mu \\ \searrow \bar{\mu} \end{array}\right)$$

or

$$\Gamma\left(W^+ \text{ wavy} \begin{array}{l} \nearrow e \\ \searrow \bar{\nu}_e \end{array}\right) = \Gamma\left(W^+ \text{ wavy} \begin{array}{l} \nearrow \mu \\ \searrow \bar{\nu}_\mu \end{array}\right)$$

(for finite energies: mass dependence)

- We focus on  $SU(2)_L$ -sector

For leptons it has been measured:  $g_e = g_\mu = g_\tau$

- Compare previous form of gauge boson couplings: *looks* universal

But **universality is a basis independent property** ⇒ Go to mass eigenbasis

This is the basis we use when measuring particles

# Quark vs Lepton Flavour Universality: Quarks

Diagonalize Yukawa interaction

$$\mathcal{L}_{\text{Yukawa}}^q \supset Y_{ij}^U \bar{Q}_{Li} \tilde{\phi} U_{Rj} + Y_{ij}^D \bar{Q}_{Li} \phi D_{Rj} \rightarrow \text{diagonal mass terms}$$

Do this by *unitary* field transformation of left-handed doublet

$$U_{Li} \rightarrow V_{ij}^U U_{Lj}, \quad D_{Li} \rightarrow V_{ij}^D D_{Lj}, \quad (\text{some transf. of right-handed quarks})$$

The (so far diagonal) coupling term to  $W$ ,  $\bar{U}_{Li} \delta_{ij} D_{Lj} W^+$ , transforms to

$$\delta_{ij} \rightarrow (V^U)_{ik}^\dagger \delta_{kl} (V^D)_{lj} = (V_{\text{CKM}})_{ij}$$

$\Rightarrow$  **Non-universal** due to *independent* transf. of components of  $SU(2)_L$ -doublet

Note that e.g.  $\bar{U}_{Li} \delta_{ij} U_{Lj} Z$  stays diagonal/universal ( $\rightarrow$  still no FCNC)

# Quark vs Lepton Flavour Universality: Leptons

Repeat for lepton sector...

$\mathcal{L}_{\text{Yukawa}}^l \supset Y_{ij}^l \bar{L}_{Li} \phi E_{Rj} + \text{h.c.} \rightarrow$  diagonal mass terms

$L_{Li} \rightarrow V_{ij}^L L_{Lj}$ , (*some transf. of right-handed leptons*)

Important difference: Components of doublet transformed together  
(there is only *one* Yukawa matrix to be diagonalized)

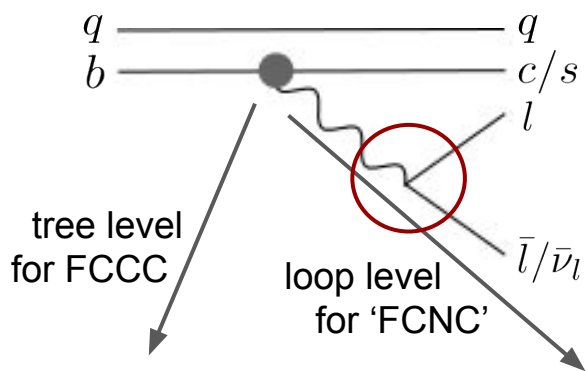
Therefore the  $W^\pm$ -coupling transforms like

$$\delta_{ij} \rightarrow (V^l)_{ik}^\dagger \delta_{kl} (V^l)_{lj} = \delta_{ij}$$

$\Rightarrow$  **Lepton Flavour Universality (LFU)** in the SM

# Testing LFU: B to D/K

Consider B-meson decays to D- or K-mesons:

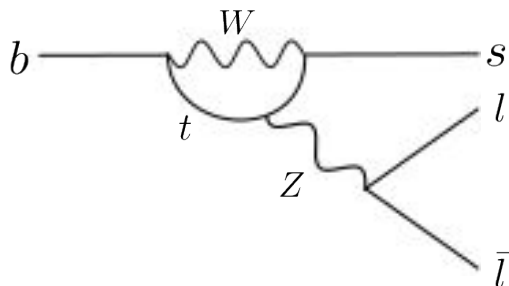
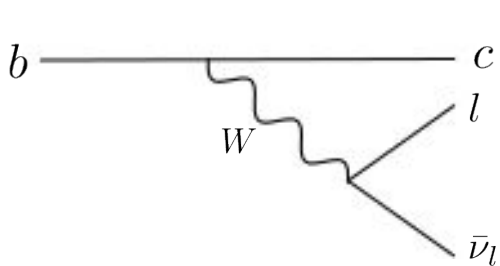


where e.g.

$$\Gamma(B \rightarrow D l \bar{\nu}_l) = \Gamma(b \rightarrow c l \bar{\nu}_l) \cdot F_{\text{QCD}}$$

→  $F_{\text{QCD}}$  is independent of  $l$  and  $\bar{\nu}_l$

⇒ consider ratios such that lepton-independent factors drop out





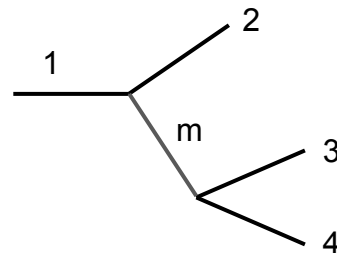
# Testing LFU: Factorization Question

Almost all the calculations of the branching ratios in flavour physics rely on the **narrow width approximation (NWA)**:

Intermediate particle *created on-shell* with subsequent decay

Works well when:

- Mass peak is narrow:  $\Gamma_m \ll m$ .
- Propagator is separable from matrix element.
- Sub-processes are kinematically allowed:  $\sqrt{s} \gg m + m_2$ ,  $m \gg m_3 + m_4$ .
- No interference.



$$\Gamma(1 \rightarrow 234) = \Gamma(1 \rightarrow 2m) \times \text{Br}(m \rightarrow 34)$$

## Testing LFU

$$\begin{aligned} \frac{\Gamma(B \rightarrow D l \bar{\nu}_l)}{\Gamma(B \rightarrow D l' \bar{\nu}_{l'})} &= \frac{Br(W \rightarrow l \bar{\nu}_l)}{Br(W \rightarrow l' \bar{\nu}_{l'})} \frac{\Gamma(b \rightarrow c W)}{\Gamma(b \rightarrow c W)} \frac{F_{\text{QCD}}}{F_{\text{QCD}}} \\ &= \frac{Br(W \rightarrow l \bar{\nu}_l)}{Br(W \rightarrow l' \bar{\nu}_{l'})} \stackrel{?}{=} 1 \end{aligned}$$

Test: Ratios of decay rates that only differ by final lepton content (e.g.  $B \rightarrow D l \bar{\nu}_l$ ) should be unity (up to lepton mass dependence).

# Experimental Signatures of Flavour Anomalies

# Signature Part - General Idea

- Measure B decays that only differ in final lepton content (Test LFU)

$$B \rightarrow X l \nu_l$$

$$B \rightarrow X l l$$

where X is meson under study

$$R_X \equiv \frac{BR(B \rightarrow X l l / l \nu)}{BR(B \rightarrow X l' l' / l' \nu')}$$

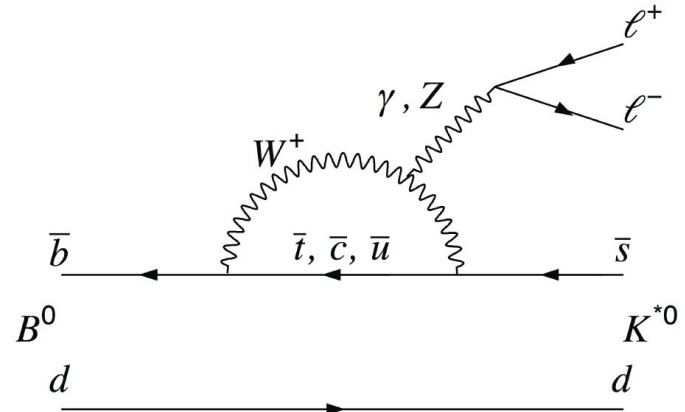
- rare loop induced b decays  $R_{K^*}(b \rightarrow s)$
- tree-level tauonic decays  $R_{D^*}, R_{J/\Psi}(b \rightarrow c)$

# FCNC (RK)

$$R_{K^{(*)}} \equiv \frac{BR(B \rightarrow K^{(*)} \mu^+ \mu^-)}{BR(B \rightarrow K^{(*)} e^+ e^-)}$$

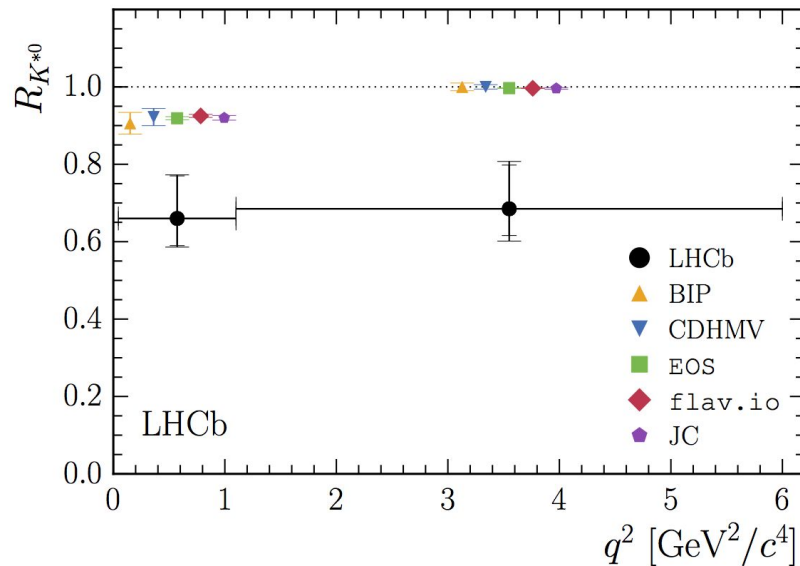
- Loop process, rare in SM, good chance for new physics
- Theoretical uncertainties factor out and cancel
- In measurement: double ratio to J/Psi, first order systematic cancellation

From SM:  $R_{K^{(*)}} = 1 + \text{phase space corr.}$



# FCNC discrepancies

- Two bins:
  - low- $q^2$   $0.0045 \text{ GeV}^2$ - $1.1 \text{ GeV}^2$
  - central- $q^2$   $1.1 \text{ GeV}^2$  -  $6 \text{ GeV}^2$
  - good theoretical description

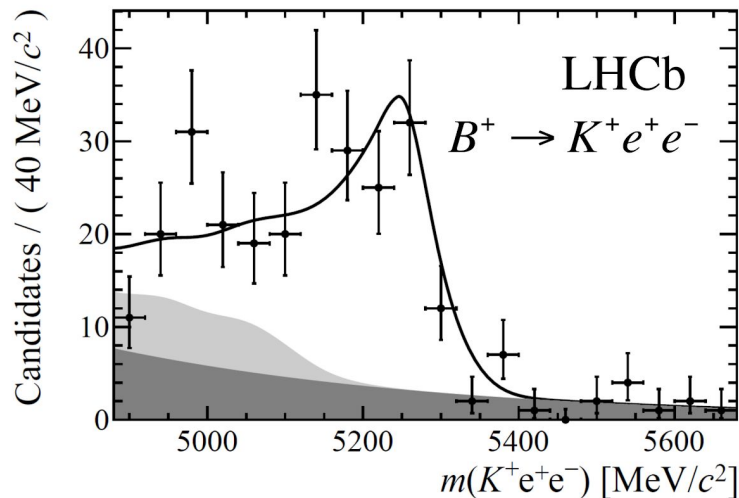
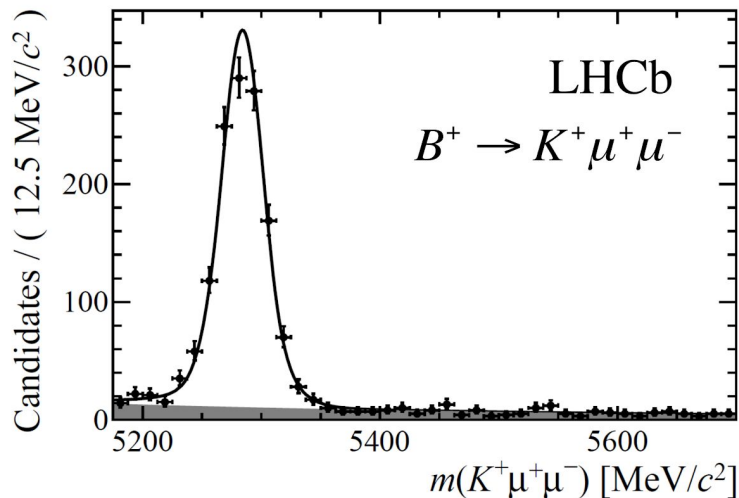


$$R_{K^*} = 0.66_{-0.07}^{+0.11} \text{ (stat)} \pm 0.03 \text{ (syst)} \quad \text{for } 0.045 < q^2 < 1.1 \text{ GeV}^2/c^4$$

$$R_{K^*} = 0.69_{-0.07}^{+0.11} \text{ (stat)} \pm 0.05 \text{ (syst)} \quad \text{for } 1.1 < q^2 < 6.0 \text{ GeV}^2/c^4$$

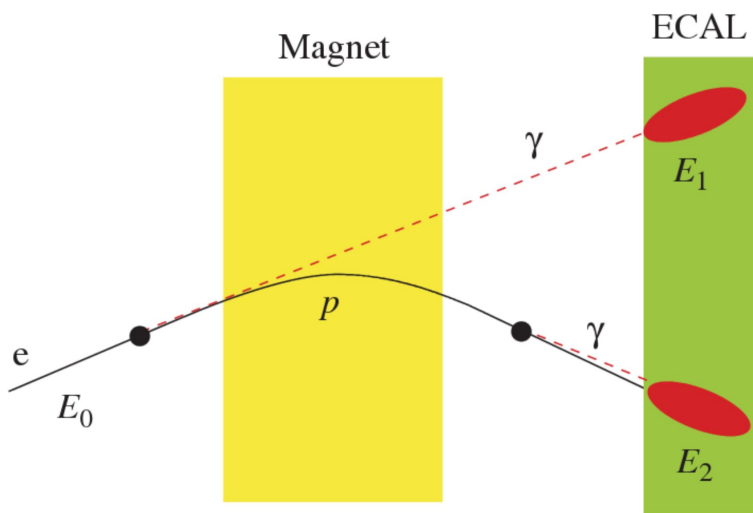
- SM compatibility at 2.2-2.5 $\sigma$  level

# Experimental Difficulties



- Muons very clean
- Electrons more problematic

# Difficulties in electron reconstruction



- Electron reconstruction difficult
- Bremsstrahlung affects resolution & efficiencies
- Can be partially corrected

Also: Higher Trigger Threshold for e-

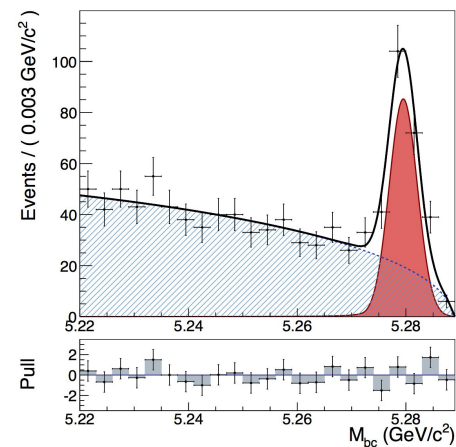
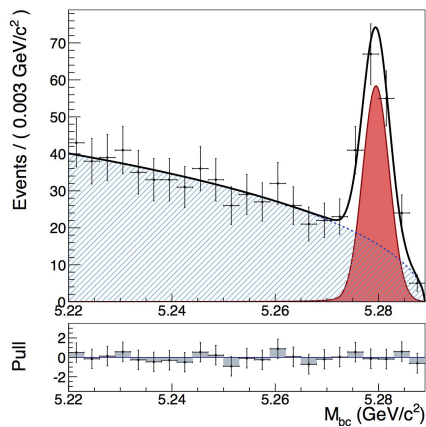
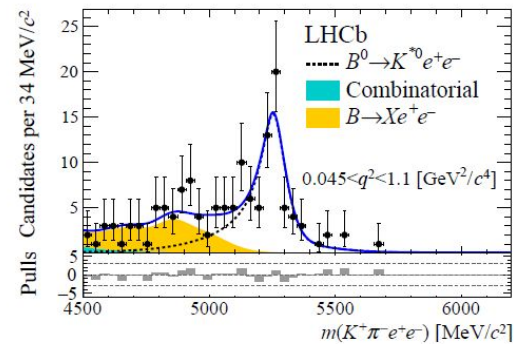
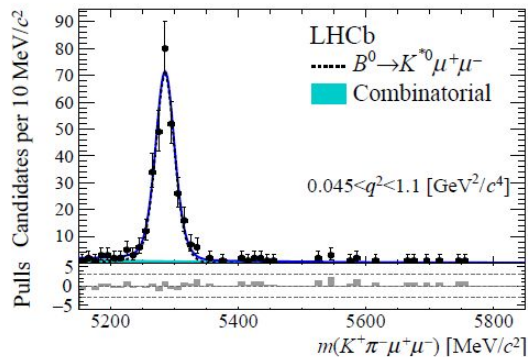


# Outlook for the RK(\*) anomaly

→ Higher statistics from LHCb

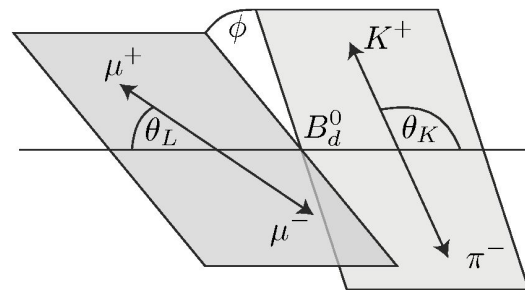
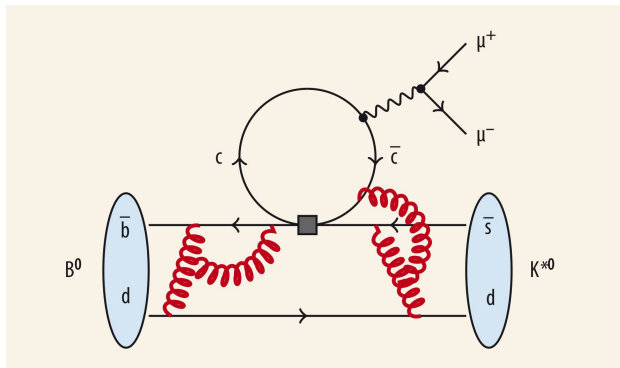
→ New experiment: Belle II

→ Improved resolution in electron channel



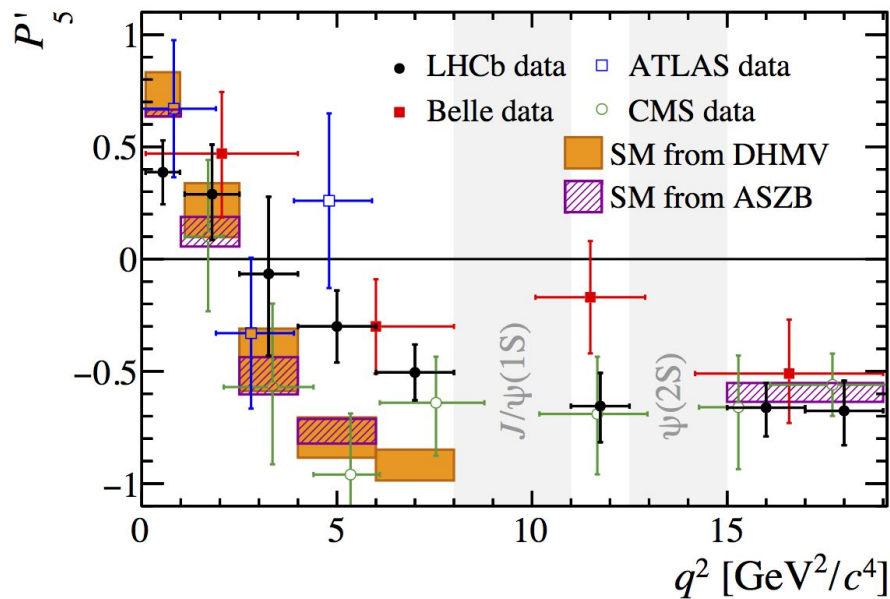
# Angular Observable for FCNC

- angular observable  $P_5'$   
 -> form factor uncertainties cancel  
 at leading order
- significant tension of 3.4 sigma
- $J/\Psi$ : theo. prediction difficult



$$d^4\Gamma$$

$$d \cos \theta_\ell d \cos \theta_K d\phi dq^2$$



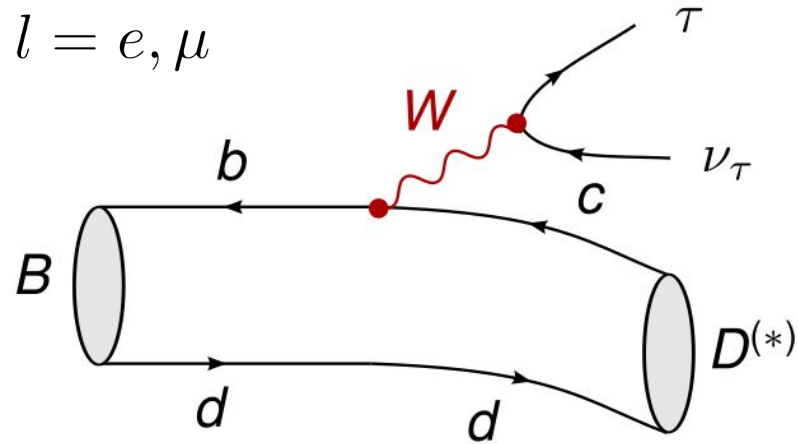
# Tree Level (RD, RD\*)

$$R_{D^{(*)}} \equiv \frac{BR(B \rightarrow D^{(*)} \tau \bar{\nu}_\tau)}{BR(B \rightarrow D^{(*)} l \bar{\nu}_l)}, \quad \text{with } l = e, \mu$$

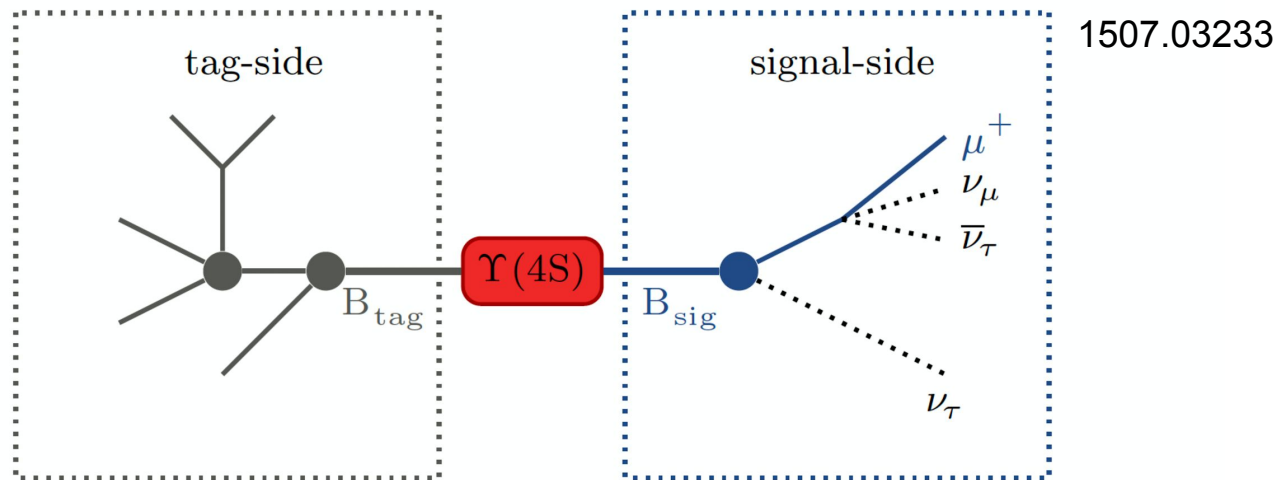
$$\rightarrow R_D \cong 0.3, R_{D^*} \cong 0.25$$

Problem:  $\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau$  or  $\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau$

→ Similar final states in numerator & denominator



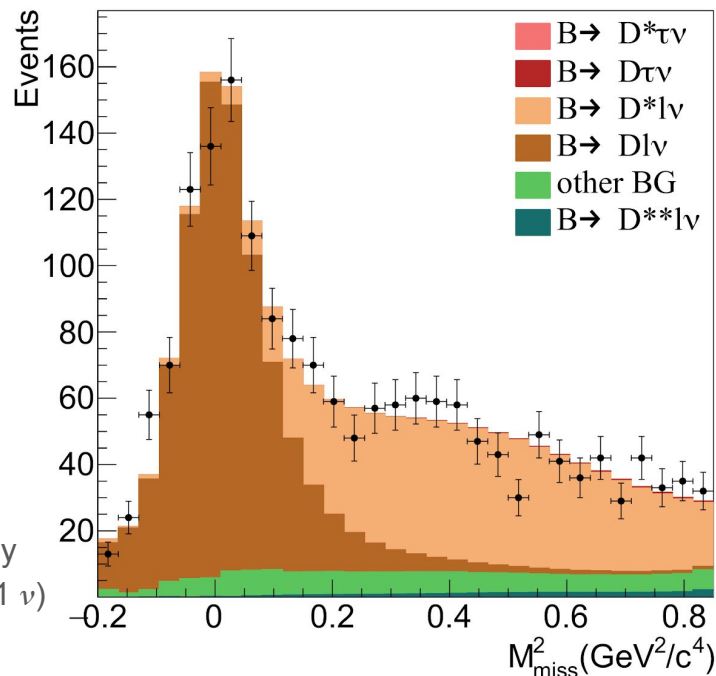
# Interlude: Advantages of Belle



E.g. hadronic:  $p_{\text{invisible}} = p_{\text{beam}} - p_{\text{tag}} - p_{\text{visible}}$

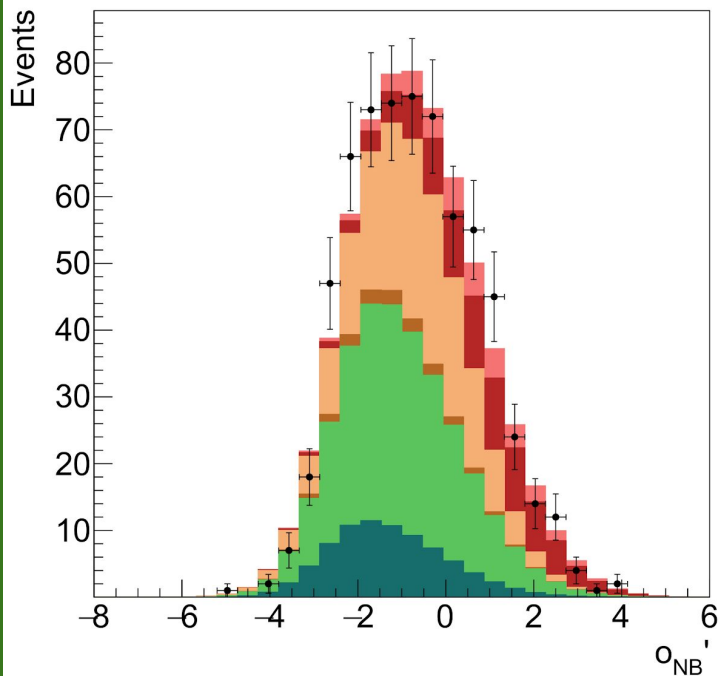
→ Different  $p_{\text{invisible}}$  for numerator (3  $\nu$ ) and denominator (1  $\nu$ )

# $M_{\text{miss}}$ distribution sg/bkg



dominated by  
denominator ( $1 \nu$ )

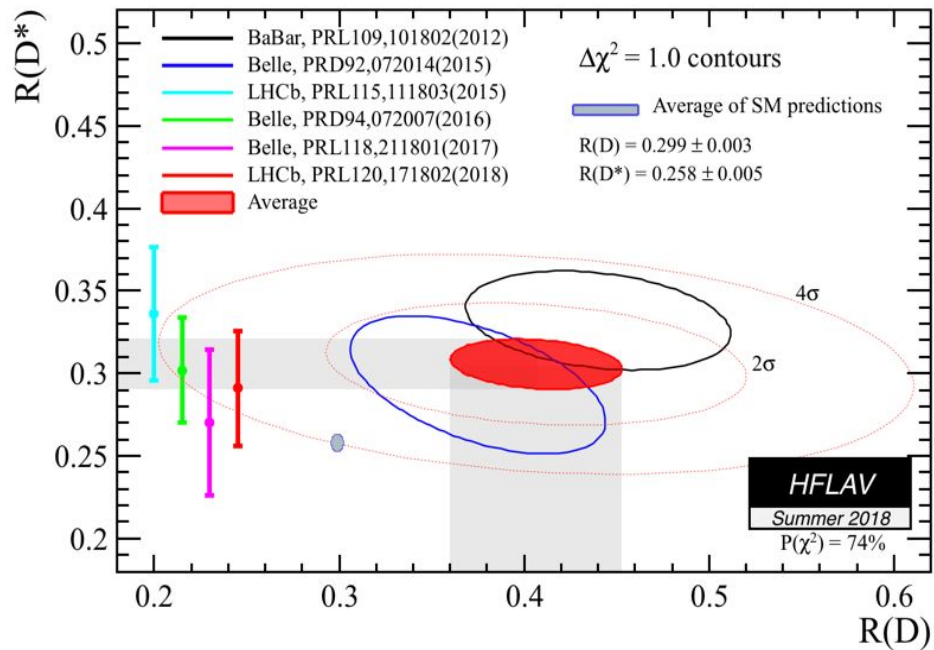
$M_{\text{miss}}^2 < 0.85$



dominated by  
numerator ( $3 \nu$ )

$M_{\text{miss}}^2 > 0.85$  (+ neural net)

# Results (RD, RD\*) anomaly



# Theory & Model-building

## **$b \rightarrow s$ anomalies**

Found by **LHCb** (and perhaps  
hinted by **Belle**)

Many observables: global pattern

Neutral current

1-loop (and CKM-suppressed)  
in the SM

The New Physics can be heavy

## **$b \rightarrow c$ anomalies**

Found by several experiments  
(**LHCb**, **BaBar** and **Belle**)

Two observables:  $R(D)$  and  $R(D^*)$

Charged current

Tree-level in the SM

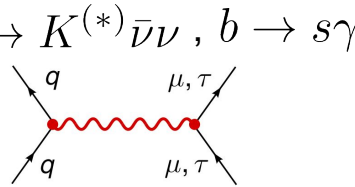
The New Physics must be light

# General consideration and remarks

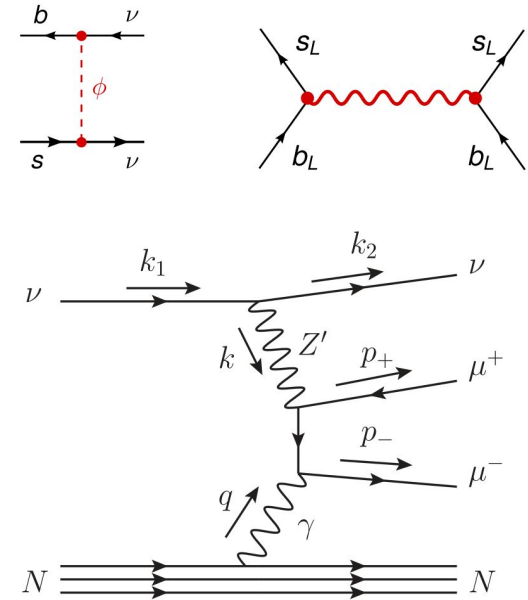
- angular and BR anomalies can be faked by hadronic uncertainties -> QCD effect?
- LFU ratios are “clean” (cannot be mimiced by hadronic physics) -> deviation still below  $3\sigma$

Long list of experimental constraints:

- other flavor observables: Bs-mixing,  $B \rightarrow K^{(*)} \bar{\nu} \nu$ ,  $b \rightarrow s \gamma$
- direct LHC search:  $pp \rightarrow \mu\mu, \tau\tau$
- lepton universality test:  $Z \rightarrow ll$
- neutrino trident production
- precision EW data



Anomalies can go away





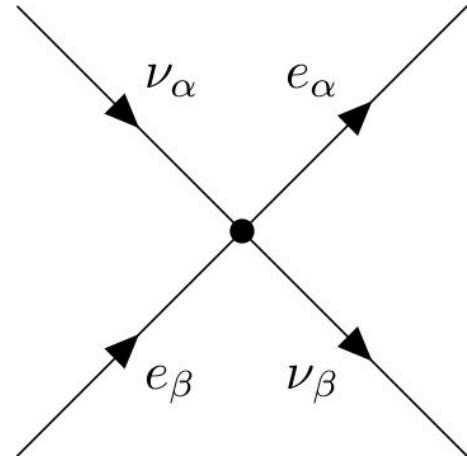
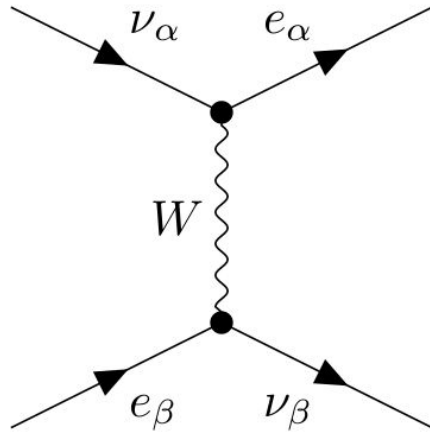
# EFT as model-independent approach

Assume:

1. Anomalies caused by New physics
2. new states are “heavy”:  $\Lambda \gg m_b$



perfect playground for EFT!



$$\frac{g^2}{8m_W^2} \longrightarrow \frac{G_F}{\sqrt{2}}$$

# weak EFT for (b-s) anomaly:

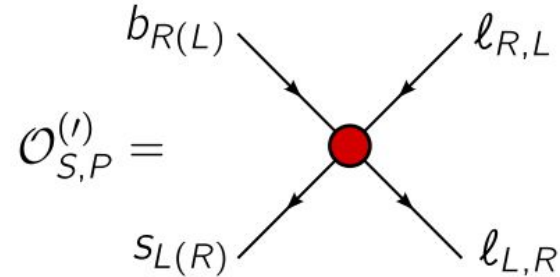
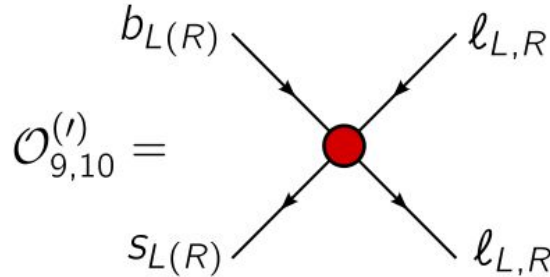
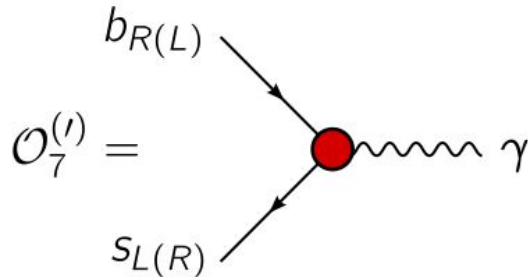
- non-renormalisable operators  $O_i$  + Wilson coefficients  $C_i$
- $C_i$  receive contributions from SM and NP
- SM reaction calculable and known with high precision
- **important** for anomaly:  $C_9, C_{10}$

$$\mathcal{H}_{\text{eff}} = -\frac{4 G_F}{\sqrt{2}} \frac{e^2}{16\pi^2} V_{tb} V_{ts}^* \sum_i C_i O_i + \text{h.c.}$$

$$O_7^{(\prime)} = \frac{m_b}{e} (\bar{s} \sigma_{\mu\nu} P_{R(L)} b) F^{\mu\nu}$$

$$O_9^{(\prime)} = (\bar{s} \gamma_\mu P_{L(R)} b) (\bar{\ell} \gamma^\mu \ell) \quad O_{10}^{(\prime)} = (\bar{s} \gamma_\mu P_{L(R)} b) (\bar{\ell} \gamma^\mu \gamma_5 \ell)$$

$$O_S^{(\prime)} = (\bar{s} P_{R(L)} b) (\bar{\ell} \ell) \quad O_P^{(\prime)} = (\bar{s} P_{R(L)} b) (\bar{\ell} \gamma_5 \ell)$$



# Gauge-invariant EFT approach: SMEFT

- non-gauge invariant EFTs **miss** relations among operators
- formulate EFT in terms of gauge-invariant operators
  - up to dim-6
  - 2499 real parameters
  - full 1-loop RGEs computed

SMEFT operator	Definition	Matching	Order
$[Q_{\ell q}^{(1)}]_{aa23}$	$(\bar{\ell}_a \gamma_\mu \ell_a) (\bar{q}_2 \gamma^\mu q_3)$	$\mathcal{O}_{9,10}$	Tree
$[Q_{\ell q}^{(3)}]_{aa23}$	$(\bar{\ell}_a \gamma_\mu \tau^I \ell_a) (\bar{q}_2 \gamma^\mu \tau^I q_3)$	$\mathcal{O}_{9,10}$	Tree
$[Q_{qe}]_{23aa}$	$(\bar{q}_2 \gamma_\mu q_3) (\bar{e}_a \gamma^\mu e_a)$	$\mathcal{O}_{9,10}$	Tree
$[Q_{\ell d}]_{aa23}$	$(\bar{\ell}_a \gamma_\mu \ell_a) (\bar{d}_2 \gamma^\mu d_3)$	$\mathcal{O}'_{9,10}$	Tree
$[Q_{ed}]_{aa23}$	$(\bar{e}_a \gamma_\mu e_a) (\bar{d}_2 \gamma^\mu d_3)$	$\mathcal{O}'_{9,10}$	Tree
$[Q_{\varphi \ell}^{(1)}]_{aa}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{\ell}_a \gamma^\mu \ell_a)$	$\mathcal{O}_{9,10}$	1-loop
$[Q_{\varphi \ell}^{(3)}]_{aa}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi) (\bar{\ell}_a \gamma^\mu \tau^I \ell_a)$	$\mathcal{O}_{9,10}$	1-loop
$[Q_{\ell u}]_{aa33}$	$(\bar{\ell}_a \gamma_\mu \ell_a) (\bar{u}_3 \gamma^\mu u_3)$	$\mathcal{O}_{9,10}$	1-loop
$[Q_{\varphi e}]_{aa}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{e}_a \gamma^\mu e_a)$	$\mathcal{O}_{9,10}$	1-loop
$[Q_{eu}]_{aa33}$	$(\bar{e}_a \gamma_\mu e_a) (\bar{u}_3 \gamma^\mu u_3)$	$\mathcal{O}_{9,10}$	1-loop

$$(\bar{\ell}_a \gamma_\mu \ell_a) (\bar{q}_2 \gamma^\mu q_3) \propto \mathcal{O}_9 - \mathcal{O}_{10}$$

One  
gauge-invariant  
operator

Two  
non-gauge-invariant  
operators

NP  
scale  
 $\Lambda$

?

SMEFT

$$[Q_{\ell q}^{(1)}]_{2223} = (\bar{\ell}_2 \gamma_\mu \ell_2) (\bar{q}_2 \gamma^\mu q_3)$$

EW  
scale  
 $\mu_{EW}$

Matching

WET

$$\mathcal{O}_{9\mu} = (\bar{s} \gamma_\mu P_L b) (\bar{\mu} \gamma^\mu \mu)$$

⋮

# Global fits

- same Wilson coefficients enter several observables
- use pattern of deviations to extract “best” value

↳ NP preferred over SM by more than **4-5 $\sigma$**  !

↳  $C_{9\mu}$  seems to be crucial !

## Inclusive

$$B \rightarrow X_s \gamma \text{ (BR)} \text{ ..... } C_7^{(\prime)}$$

$$B \rightarrow X_s \ell^+ \ell^- \text{ (dBR/dq}^2\text{)} \text{ ..... } C_7^{(\prime)}, C_9^{(\prime)}, C_{10}^{(\prime)}$$

## Exclusive leptonic

$$B_s \rightarrow \ell^+ \ell^- \text{ (BR)} \text{ ..... } C_{10}^{(\prime)}$$

## Exclusive radiative/semileptonic

$$B \rightarrow K^* \gamma \text{ (BR, S, A}_1\text{)} \text{ ..... } C_7^{(\prime)}$$

$$B \rightarrow K \ell^+ \ell^- \text{ (dBR/dq}^2\text{)} \text{ ..... } C_7^{(\prime)}, C_9^{(\prime)}, C_{10}^{(\prime)}$$

$$B \rightarrow K^* \ell^+ \ell^- \text{ (dBR/dq}^2, \text{ angular obs.)} \text{ ... } C_7^{(\prime)}, C_9^{(\prime)}, C_{10}^{(\prime)}$$

$$B_s \rightarrow \phi \ell^+ \ell^- \text{ (dBR/dq}^2, \text{ angular obs.)} \text{ ... } C_7^{(\prime)}, C_9^{(\prime)}, C_{10}^{(\prime)}$$

[A. Vicente, Post-FPCP School 2018]

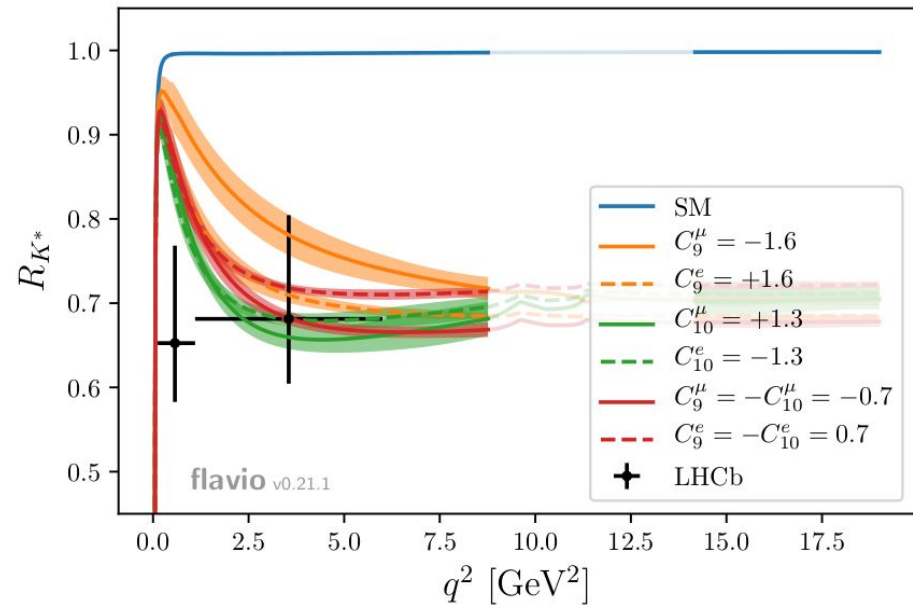
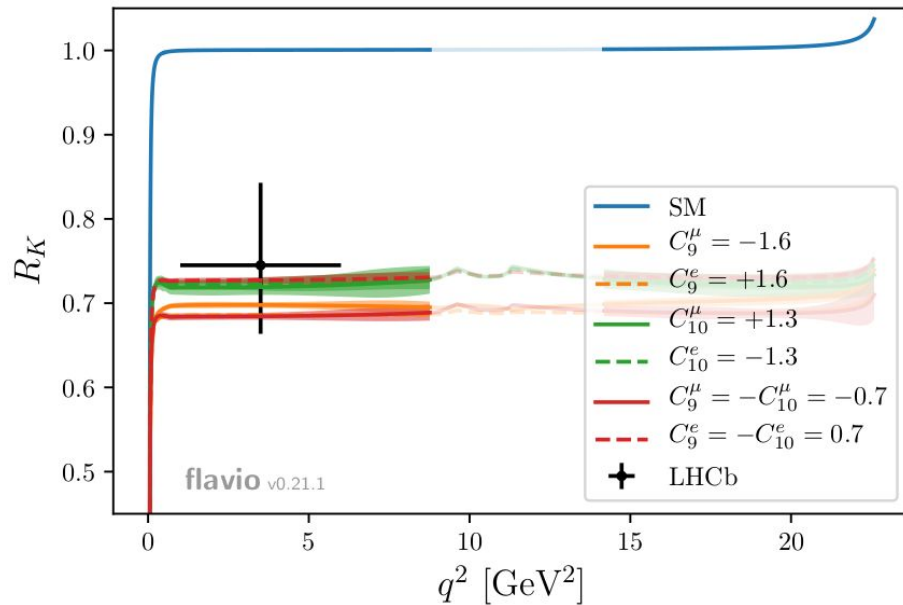
1D Hyp.	All					LFUV				
	Best fit	1 $\sigma$	2 $\sigma$	Pull <sub>SM</sub>	p-value	Best fit	1 $\sigma$	2 $\sigma$	Pull <sub>SM</sub>	p-value
$C_{9\mu}^{\text{NP}}$	-1.10	[-1.27, -0.92]	[-1.43, -0.74]	5.7	72	-1.76	[-2.36, -1.23]	[-3.04, -0.76]	3.9	69
$C_{9\mu}^{\text{NP}} = -C_{10\mu}^{\text{NP}}$	-0.61	[-0.73, -0.48]	[-0.87, -0.36]	5.2	61	-0.66	[-0.84, -0.48]	[-1.04, -0.32]	4.1	78
$C_{9\mu}^{\text{NP}} = -C'_{9\mu}$	-1.01	[-1.18, -0.84]	[-1.33, -0.65]	5.4	66	-1.64	[-2.12, -1.05]	[-2.52, -0.49]	3.2	31
$C_{9\mu}^{\text{NP}} = -3C_{9e}^{\text{NP}}$	-1.06	[-1.23, -0.89]	[-1.39, -0.71]	5.8	74	-1.35	[-1.82, -0.95]	[-2.38, -0.59]	4.0	71

All observables  
“clean” + “dirty”

Only LFUV observables  
“clean”

# Model-independent fits to $C_{9,10}^{(')}$

e.g. in context of  $R_{K(*)}$



More observables needed for discrimination among different best-fit scenarios !

# UV models: difficulties & common features

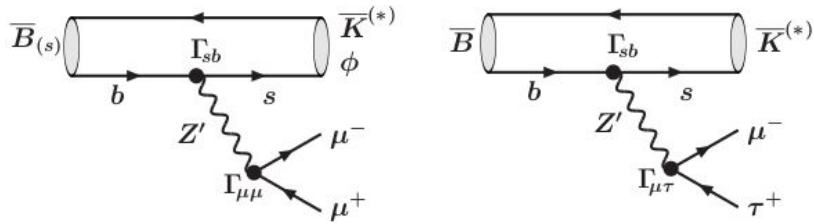
- Loop suppression of neutral currents with respect to the charged ones.
- NP:  $J_{\text{quark}} \times J_{\text{lepton}}$  with no traces in  $J_{\text{quark}} \times J_{\text{quark}}$  (constraints from  $B_s$  mixing) and  $J_{\text{lepton}} \times J_{\text{lepton}}$  (constraints from pure LFV/LUV decays).
- Most models involve:
  - New charged (coloured) states.
  - Mass  $\sim \text{TeV}$  (to explain relatively large effects).
  - Significant coupling to the 3<sup>rd</sup>-generation SM fermions (constraints from resonances decaying to  $\tau \tau$  pairs).

**Typical UV complete theory contains new states that are:**

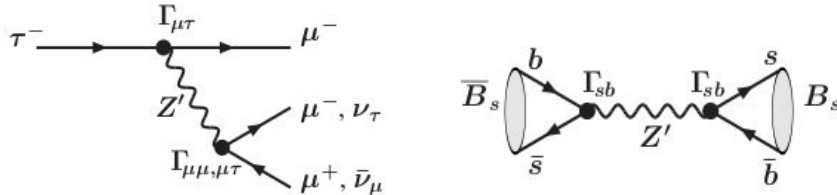
- Lorentz scalars/vectors
- $SU(3)_c$ : singlet/triplet;  $SU(2)_L$ : singlet/doublet/triplet.

# Example #1: $Z'$

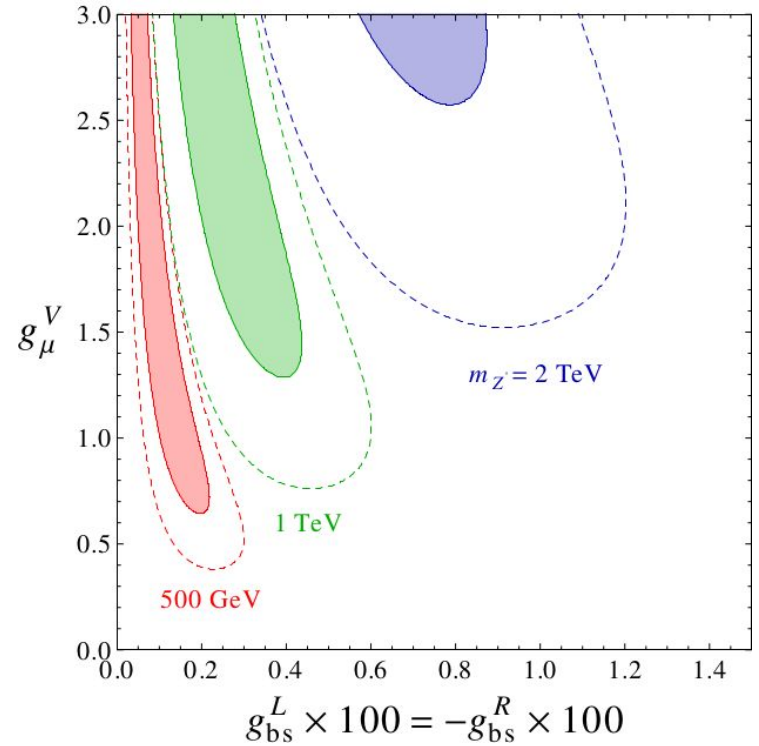
Additional  $U(1)_X$  generates  $O_9, O_{10}$ :



But also:



Explains  $B \rightarrow K$  anomalies for  $m_{Z'} \sim \text{TeV}$



# Example #2: leptoquark

New scalar field:  $SU(3)_c$ -triplet,  $SU(2)_L$ -singlet.



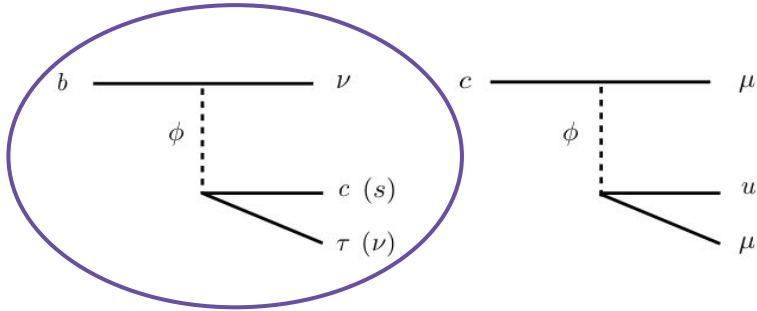
Can it explain both  $B \rightarrow K$  &  $B \rightarrow D$  anomalies?

Requiring also electric charge equal to  $-1/3$ ...

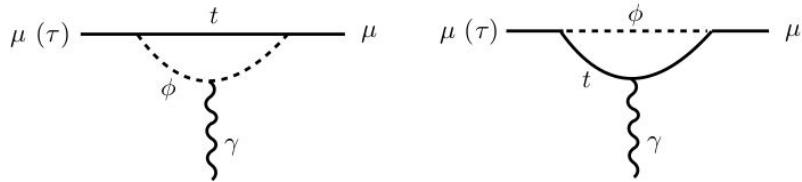


# Example #2: leptoquark

$B \rightarrow D^* \tau \nu$ : tree-level

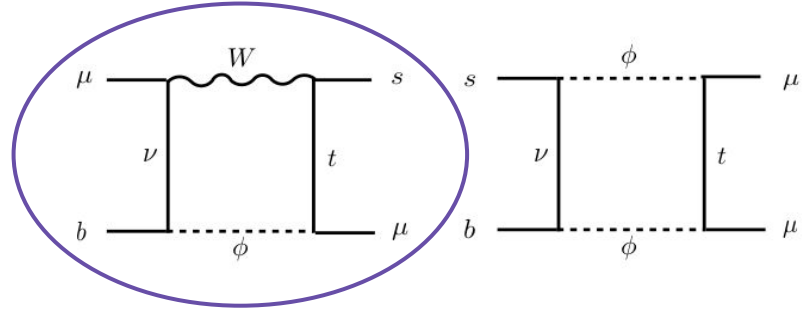


But also:



(Bounds from  $B_S - \bar{B}_S$  mixing,  $D \rightarrow \mu^+ \mu^-$ ,  $\tau \rightarrow \mu \gamma$ )

$B \rightarrow K^* \ell \ell$ : only at 1-loop level



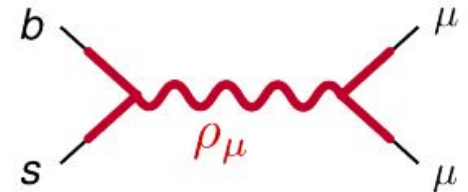
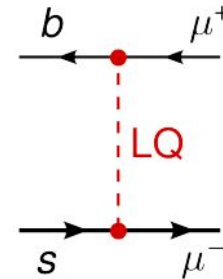
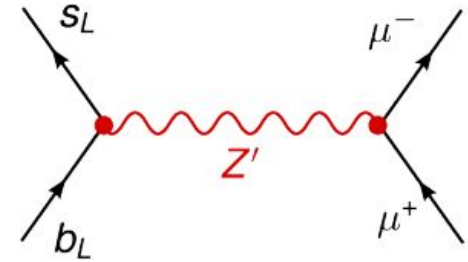
[1511.01900]

Explains both  $B \rightarrow K$  &  $B \rightarrow D$   
anomalies for  $m_\phi \sim \text{TeV}^*$

\*According to 1608.07583, accurate calculation of the loop-induced effects makes  $R_D^{\tau/\ell}$  inconsistent with data.

# Mainstream models:

1.  $Z'$ 
  - flavor-changing coupling to LH quarks
  - VL couplings to leptons
  - flavor violation or non-universality in lepton sector
2. Leptoquarks
  - scalar or vector
  - not simult. lepton non-universal and L conserving
3. Compositeness
  - neutral resonance, coupling to muons (part. composite)
  - lepton flavor violating couplings
  - constrained by LEP (Z-width) and  $B_s$ - $B_s$ -mixing



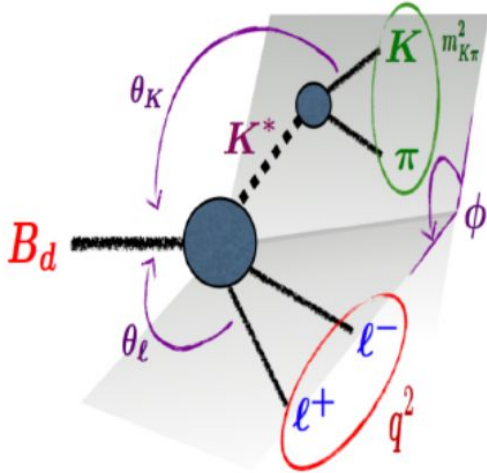
# Summary: Flavor could be around the corner!

- SM prediction: LFU!
- Several anomalies in B physics
  - $b \rightarrow s \mu \mu$  BR &  $P_5'$  - hadr. uncertainties, but significant
  - $R_K^{(*)}$  - theo. clean but not too significant
  - $R_D^{(*)}$  - theo. clean and significant
- NP highly constrained, but combined NP solution for all anomalies possible!
- More data and new experiments crucial
  - LHC Run 2
  - Belle-II experiment

Backup

# Angular observables

$$\frac{d^4\Gamma}{dq^2 d\cos\theta_K d\cos\theta_l d\phi} = \frac{9}{32\pi} \left[ J_{1s} \sin^2 \theta_K + J_{1c} \cos^2 \theta_K + (J_{2s} \sin^2 \theta_K + J_{2c} \cos^2 \theta_K) \cos 2\theta_l \right. \\ \left. + J_3 \sin^2 \theta_K \sin^2 \theta_l \cos 2\phi + J_4 \sin 2\theta_K \sin 2\theta_l \cos \phi + J_5 \sin 2\theta_K \sin \theta_l \cos \phi \right. \\ \left. + (J_{6s} \sin^2 \theta_K + J_{6c} \cos^2 \theta_K) \cos \theta_l + J_7 \sin 2\theta_K \sin \theta_l \sin \phi \right. \\ \left. + J_8 \sin 2\theta_K \sin 2\theta_l \sin \phi + J_9 \sin^2 \theta_K \sin^2 \theta_l \sin 2\phi \right]$$



[Figure borrowed from Javier Virto]

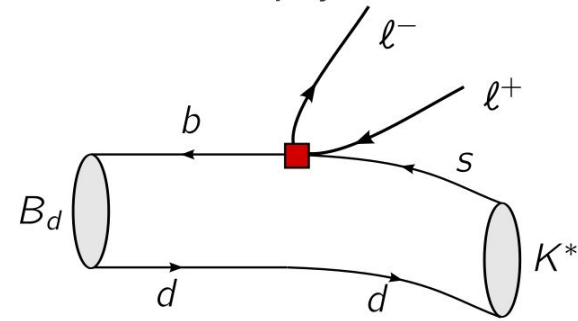
$J_i$  : functions of  $q^2$ ,  $C_i$ , FF

Optimized observables  
[Descotes-Genon et al, 2012, 2013]

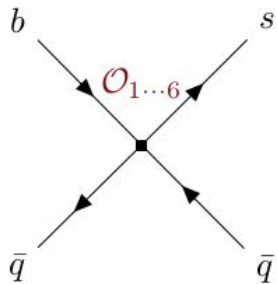
$$P'_5 = \frac{J_5}{2\sqrt{-J_{2s}J_{2c}}}$$

# weak EFT operators

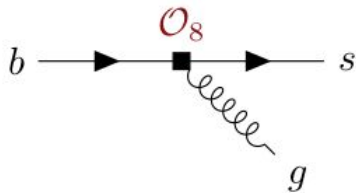
Operator set for  $b \rightarrow s$  transitions:



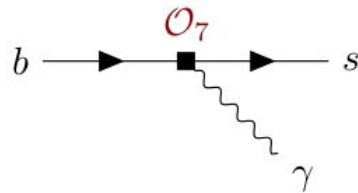
4-quark operators



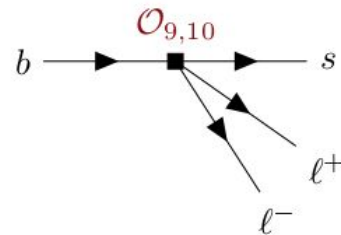
chromomagnetic dipole operator



electromagnetic dipole operator



semileptonic operators



$$\mathcal{O}_{1,2} \propto (\bar{s}\Gamma_\mu c)(\bar{c}\Gamma^\mu b)$$

$$\mathcal{O}_8 \propto (\bar{s}\sigma^{\mu\nu} T^a P_R) G_{\mu\nu}^a$$

$$\mathcal{O}_7 \propto (\bar{s}\sigma^{\mu\nu} P_R) F_{\mu\nu}^a$$

$$\mathcal{O}_9^l \propto (\bar{s}\gamma^\mu b_L)(\bar{l}\gamma_\mu l)$$

$$\mathcal{O}_{3,4} \propto (\bar{s}\Gamma_\mu b)\sum_q(\bar{q}\Gamma^\mu q)$$

$$\mathcal{O}_{10}^l \propto (\bar{s}\gamma^\mu b_L)(\bar{l}\gamma_\mu \gamma_5 l)$$

+ the chirality flipped counter-parts of the above operators,  $\mathcal{O}'_i$

# Typical EFT scales

All scales  $\Lambda_i$  probed so far appear to be rather large:

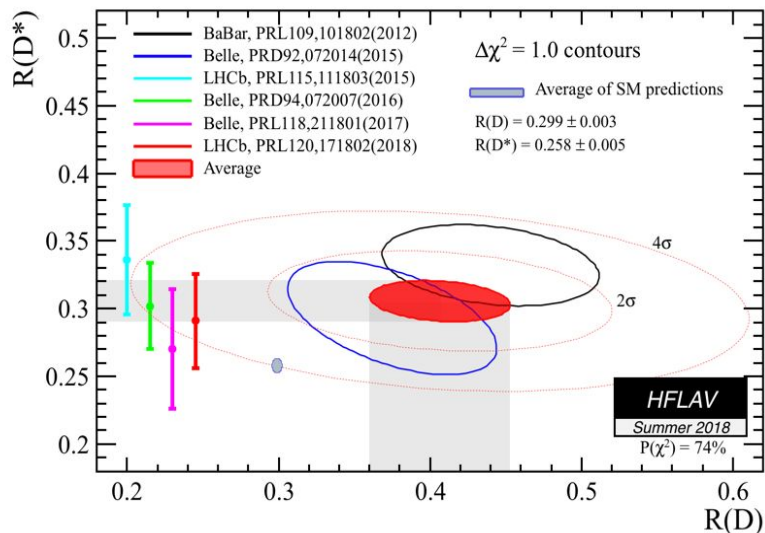
Order	Observable	New-physics scale for $g=O(1)$
D=5	Neutrino oscillations	$\Lambda \sim 10^9$ TeV
D=6	Proton decay	$\Lambda > 10^{12}$ TeV
D=6	Flavor physics	$\Lambda > 1-10^5$ TeV
D=6	EWPT	$\Lambda > 1$ TeV
D=6	Higgs couplings	$\Lambda > 0.5-1$ TeV

# combination of measurements

- “orthogonal” systematic uncertainties
- test different regions of parameter space
- plot
- combined significance...
- future improvements and prospects
- projected uncertainty and limitations

-> final comment: LHCb+Belle -> final data samples will be sufficient to confirm discovery of anomalies or rule it out

-> hot topic that could guide us to new physics -> Theory consideration part

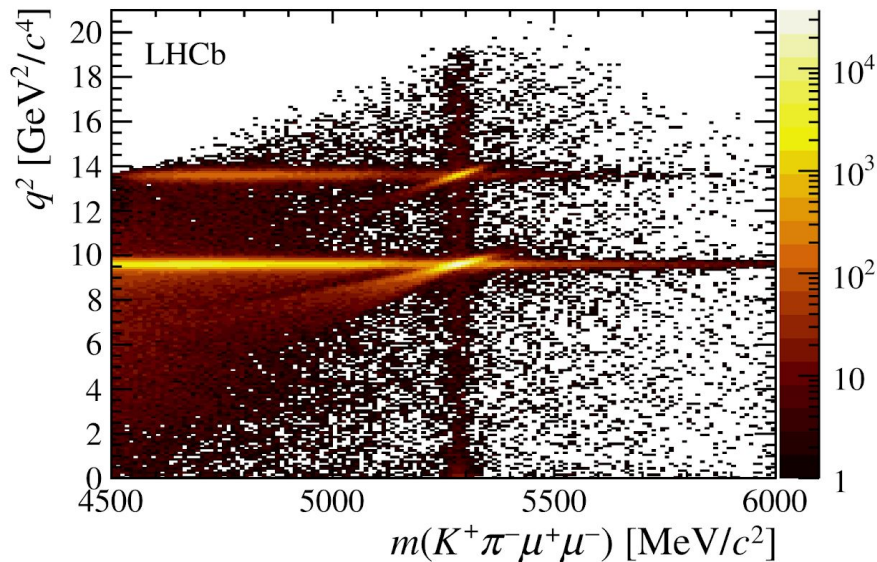




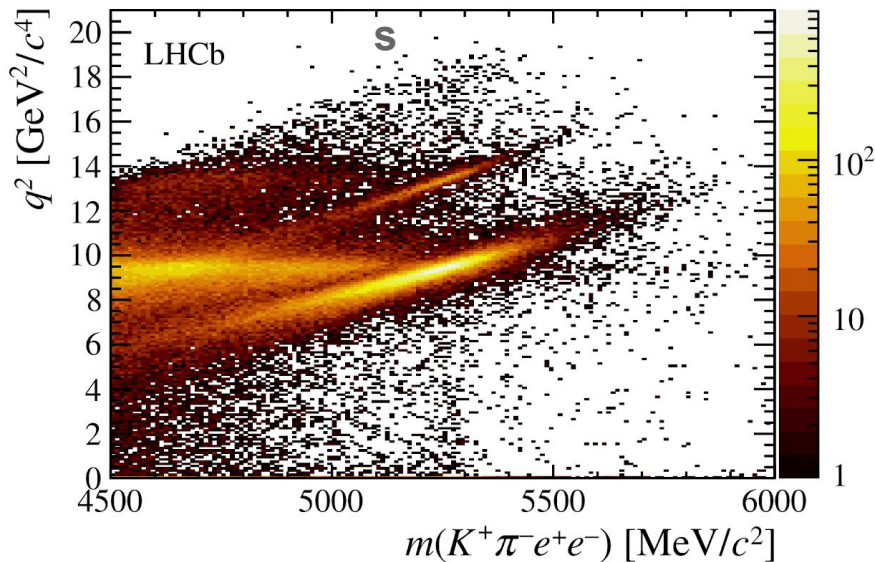
# PID - $q^2$ versus mass

JHEP08(2017)055

Muons



Electron



- J/Psi & Y(2s) visible as horizontal lines
- Vertical line: B → K\* l<sup>+</sup> l<sup>-</sup>

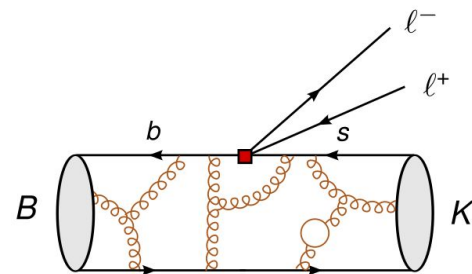
# Challenges on both sides ...

Experimental measurements:

- tbd
- tbd

Theoretical calculations:

- form factors: require non-perturbative calculations



- “non-factorisable” hadronic effects: problematic since easily generated at tree-level

