

Physics Teams W17/18:

Gravitational Waves

Part I: Peter Reimitz, Patrick Foldenauer and Christian Reichelt

Universität Heidelberg
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Examples of black hole evidence

Circumstantial evidence

1969: Donald Lynden-Bell:

Explains that quasars could be powered by supermassive black holes

1971: Louise Webster, Paul

Murdin, and Thomas Bolton:

Discover 15 solar mass invisible companion (Cygnus X-1) to a star



NASA



ESA/Hubble

Gravitational waves → Direct evidence for black holes

Soon: Event Horizon Telescope
Shadow of Sagittarius A*



EHT/Hotaka Shiokawa

What are gravitational waves?

Everything starts with the Einstein field equations

$$G_{\mu\nu} \equiv R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = 8\pi GT_{\mu\nu}$$

Weak field approximation (linearise) - expanding around flat space

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, \quad |h_{\mu\nu}| \ll 1$$

Three simple steps:

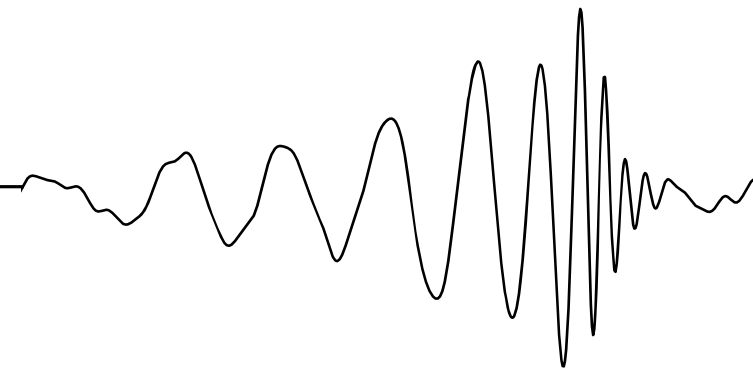
1. Linearise equation

2. Choose gauge fixing (Lorentz gauge $\square x^\mu = 0$)

3. Define the trace-reversed metric perturbation: $\bar{h}_{\mu\nu} = h_{\mu\nu} - \frac{1}{2}\eta_{\mu\nu}h$

Gives a **relativistic wave equation** for a symmetric two-tensor:

$$\square \bar{h}_{\mu\nu} = -16\pi GT_{\mu\nu}$$



Solutions to the wave equation

Consider vacuum solutions $T_{\mu\nu} = 0$, and write a solution

$$\bar{h}_{\mu\nu} = C_{\mu\nu} e^{ik_\sigma x^\sigma}$$

where k_σ is the wave vector.

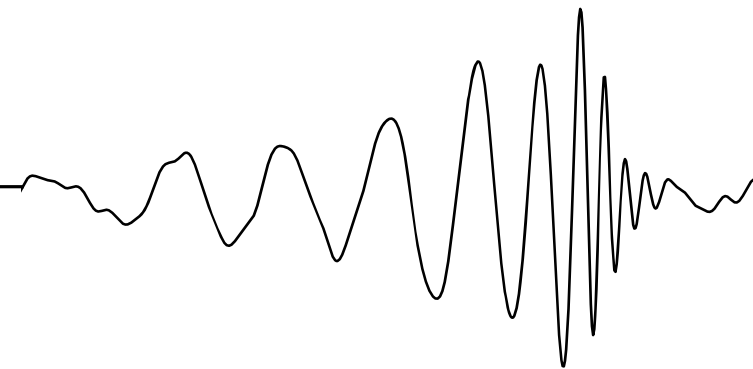
What do we learn from the E.O.M. and gauge fixing?

- $k_\mu k^\mu = 0$: Gravitational waves propagate at the **speed of light**
- Parameters of $C_{\mu\nu}$ reduced from 10 to 2: **Two polarisations**

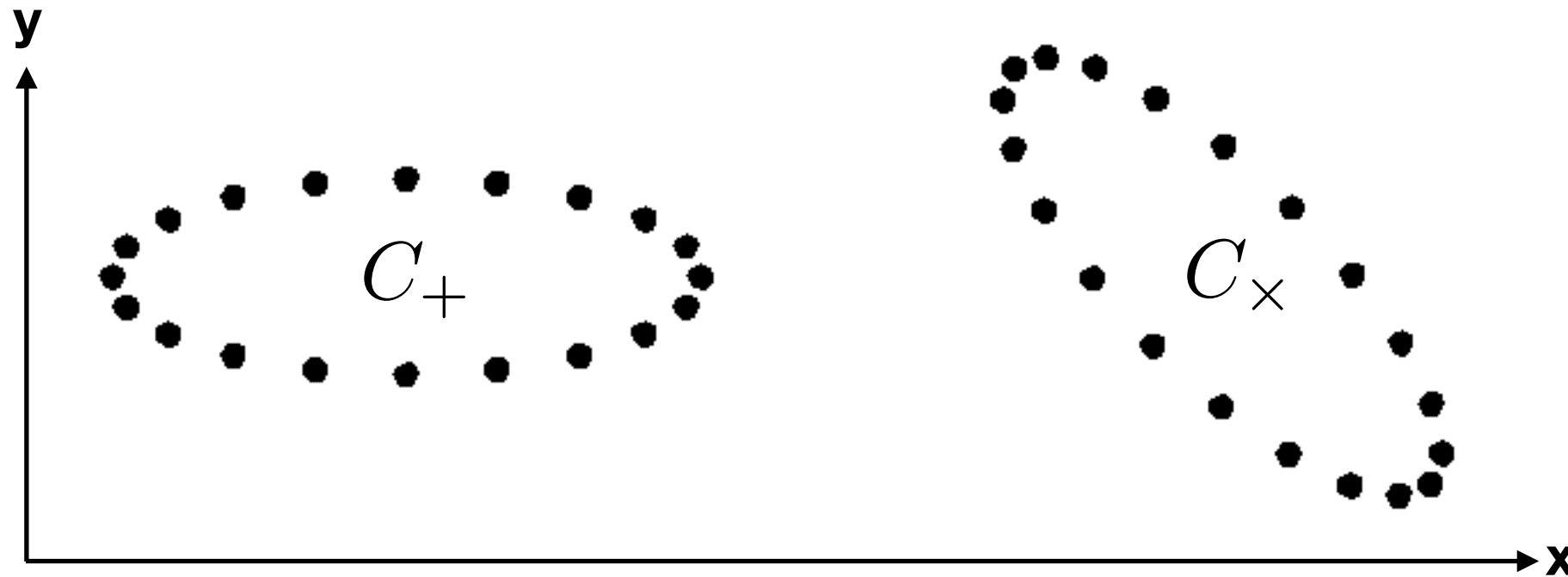
For a wave travelling in x^3 direction ($k^\mu = (\omega, 0, 0, \omega)$) we get:

$$C_{\mu\nu} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & C_+ & C_\times & 0 \\ 0 & C_\times & -C_+ & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

Solutions to the wave equation



Effect on separation of particles in the two modes:



Note:
 Modes are invariant
 under 180° rotation
 ↓
Spin 2 particle

Generation of gravitational waves? Couple to matter $T_{\mu\nu} \neq 0$

Solution:

$$\bar{h}_{ij}(t, \mathbf{x}) = \frac{2G}{rc^4} \frac{d^2}{dt^2} q_{ij}(t - r/c)$$

Lesson: **Gravitational waves generated by the quadrupole moment**

$$q_{ij} = \int \rho \left(x^i x^j - \frac{1}{3} \delta^{ij} r^2 \right) d^3x$$

The quadrupole

A moment of hand-waving:

Why is the leading contribution the quadrupole?



GW should fall off as $h_{ij} \sim 1/r$ (energy conservation)

Multipole expansion of gravitational potential

$$\Phi = -\frac{Gm}{r} + \frac{Gq_i}{r^2} + \frac{Gq_{ij}}{r^3} + \dots$$

So by dimensional analysis

$$h_{ij} \sim \frac{Gm}{c^2 r} + \frac{G}{c^3 r} \frac{\partial q_i}{\partial t} + \frac{G}{c^4 r} \frac{\partial^2 q_{ij}}{\partial t^2} + \dots$$

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No fluctuations:
Mass conservation

$$\frac{\partial q_i}{\partial t} \sim p$$

No fluctuations:
Momentum conservation

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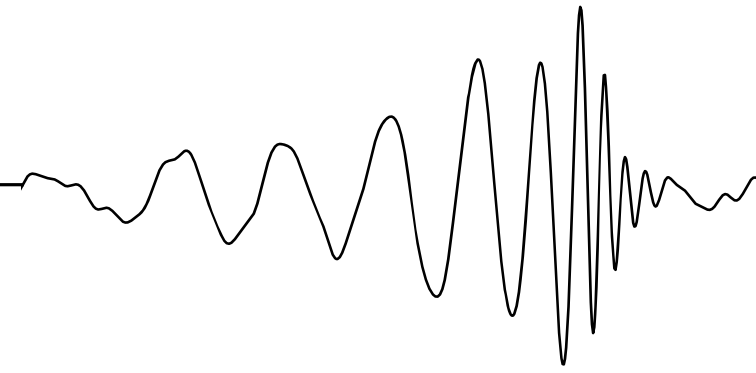
$$\frac{\partial q_i}{\partial t} \sim p$$

No fluctuations:
Momentum conservation

 What about the equivalent “magnetic dipole”?
Forbidden by conservation of angular momentum

In general quanta with spin s has multipoles $l \geq s$

Estimating the amplitude of GW



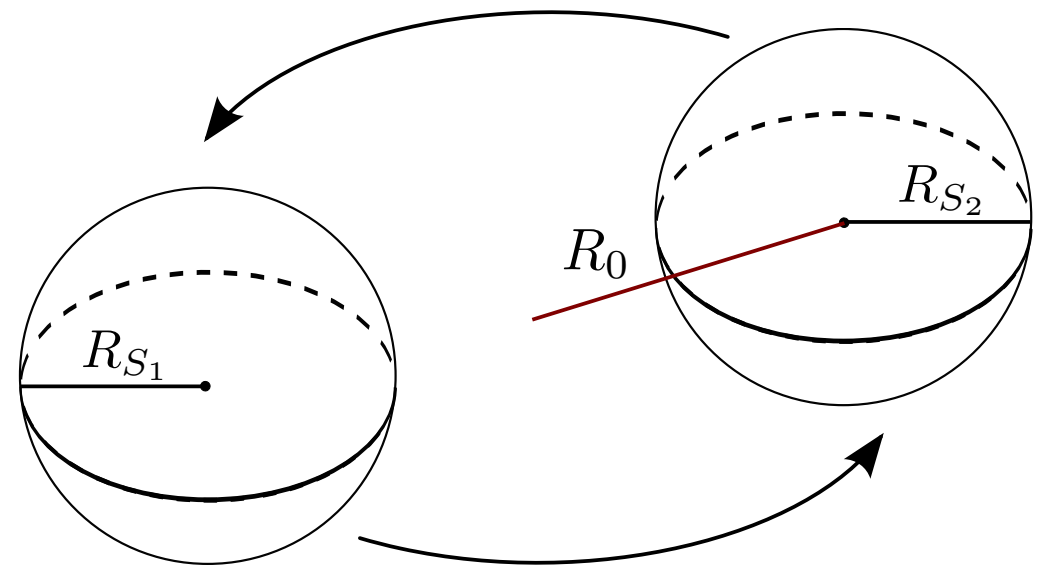
Going back to the solution

$$\bar{h}_{ij}(t, \mathbf{x}) = \frac{2G}{rc^4} \frac{d^2}{dt^2} q_{ij}(t - r/c)$$

These perturbations are tiny: $2G/c^4 \sim 10^{-44}$

Crude estimate for a black hole binary:

$$|h| \sim \frac{R_{S_1} R_{S_2}}{R_0 r}$$



Assuming close by BHs, masses M_\odot
and located in the Virgo cluster ($r \approx 15$ Mpc)

$$|h| \sim 10^{-21}$$

How can these tiny fluctuations be measured?

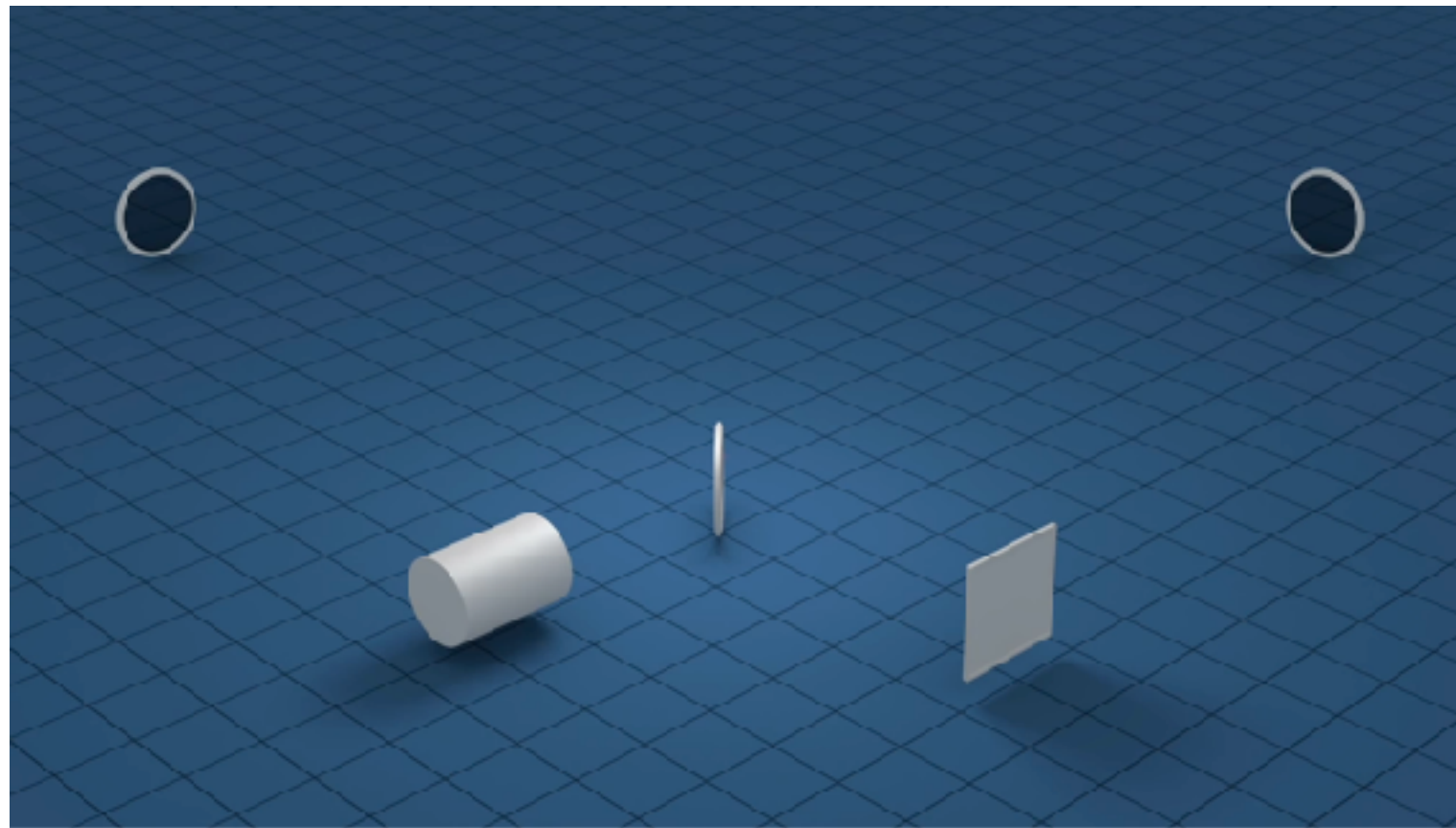
Measurement principle

Proposition: Time difference of light travelling in one arm?

In principle possible, but less attractive

➔ **Laser interferometry:**

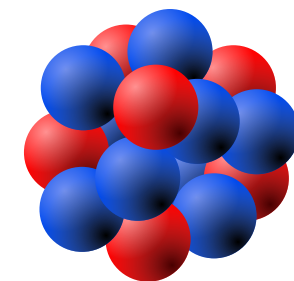
- Quadrupole nature of waves makes Michelson interferometer ideal



LIGO/T. Pyle

Need long distances and high precision since

$$\Delta L = L \cdot h \approx L \cdot 10^{-21}$$

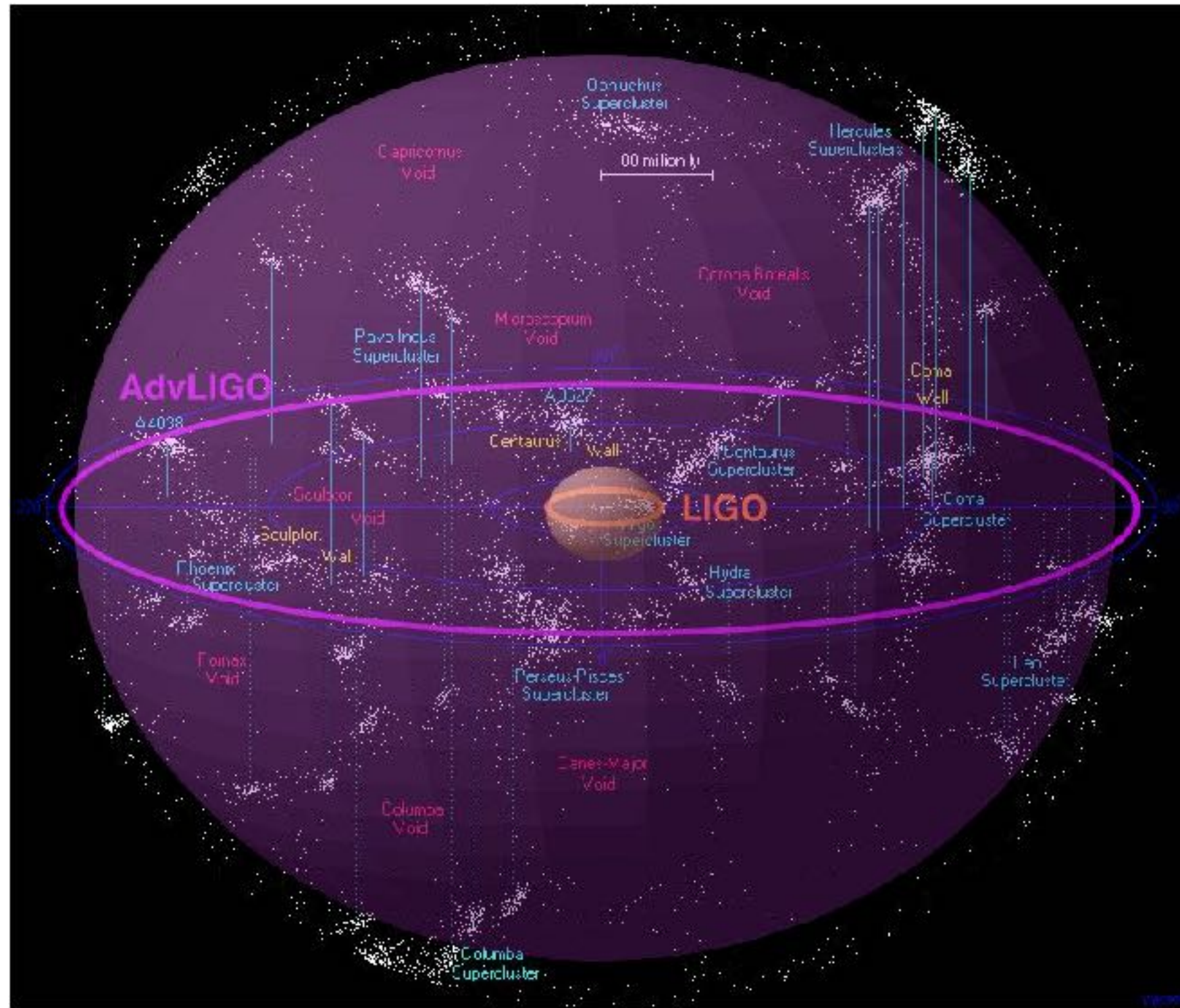
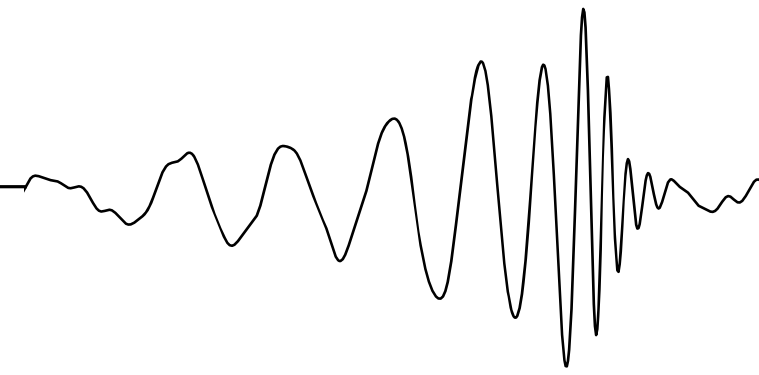


$$\sim 10^{-15} \text{ m}$$

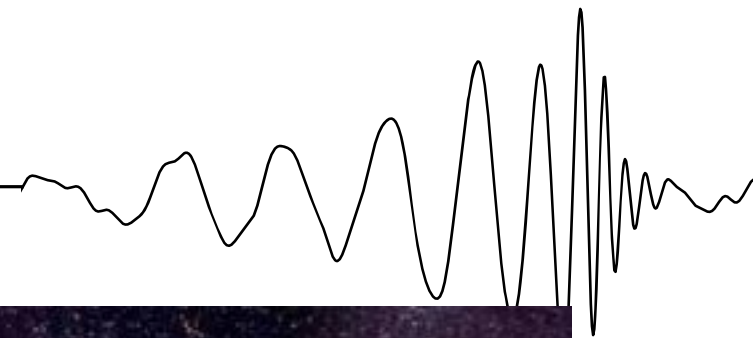
How are gravitational waves produced

Need the following ingredients:

- Very compact
- Relativistic
- Non spherically symmetric mass distribution
- Strongly time-varying quadrupole moment
- Close enough for observation



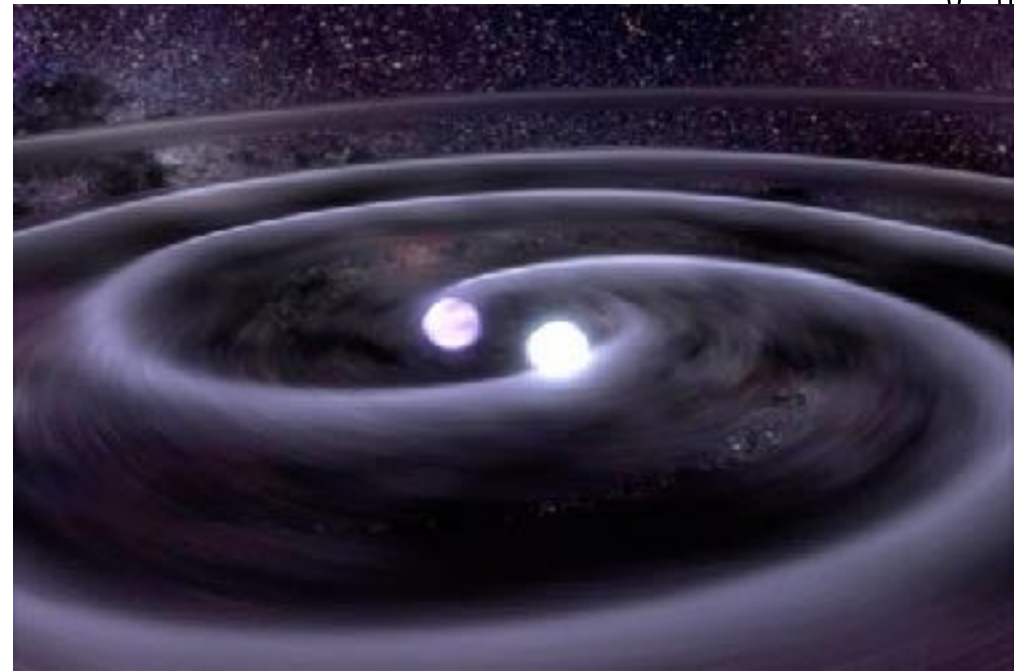
Sources of Gravitational Waves



Well-defined

Compact binary coalescence:

Short-lived signal from inspiraling of two compact objects like e.g. neutron stars or black holes (first GW observation)



Continuous wave:

Long-lived signal from e.g. spinning (non-axisymmetric) neutron stars



Sources of Gravitational Waves

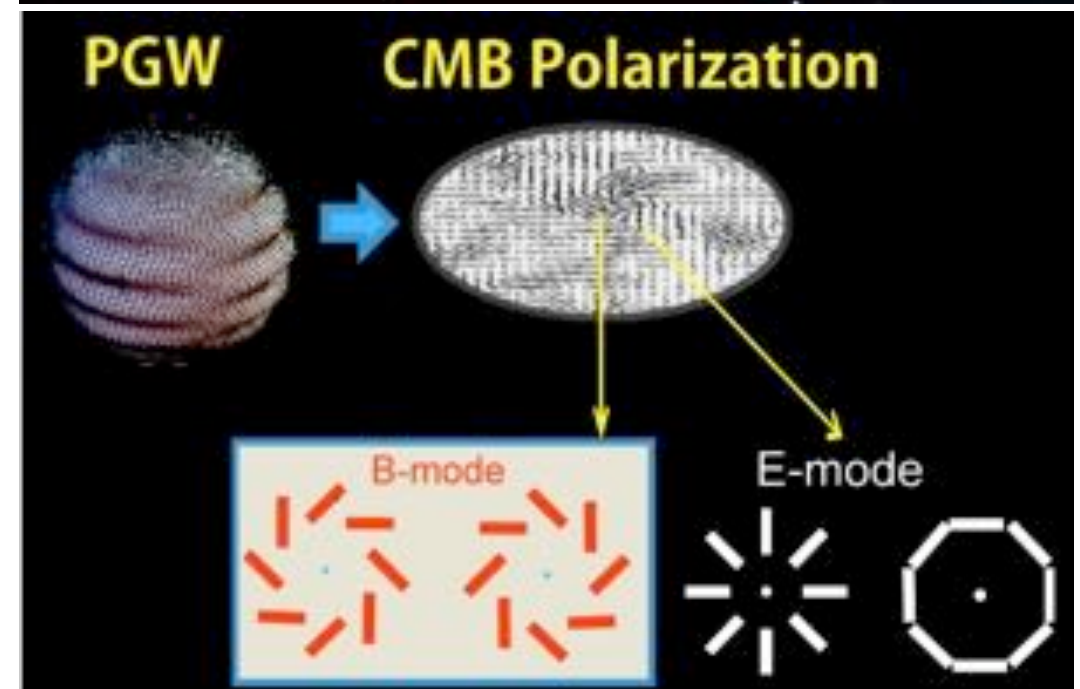
Not so well-defined ...

Bursts:

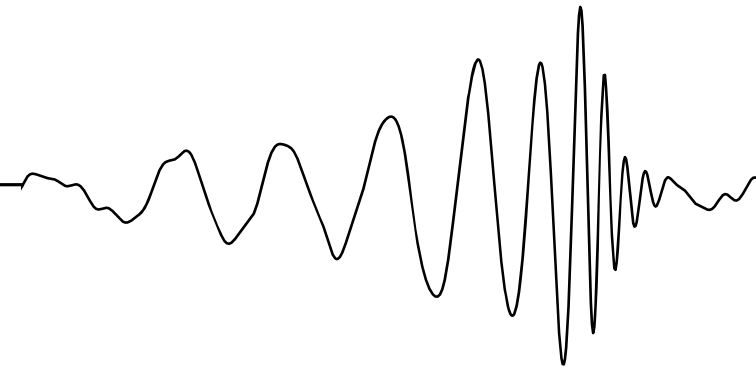
Very short-lived, poorly known transients as e.g. produced from supernovae

Stochastic background:

Long-lived superposition of incoherent sources, e.g. primordial gravitational waves from CMB (see second part)



First discovery: Black hole merger



To see how a coalescence signal is produced we have to look at the total energy luminosity in the radiation zone

$$\mathcal{L} = \frac{G}{5c^5} \langle \ddot{q}_{ij} \ddot{q}^{ij} \rangle \quad \left(\mathcal{L} = \frac{dE}{dt} \right)$$

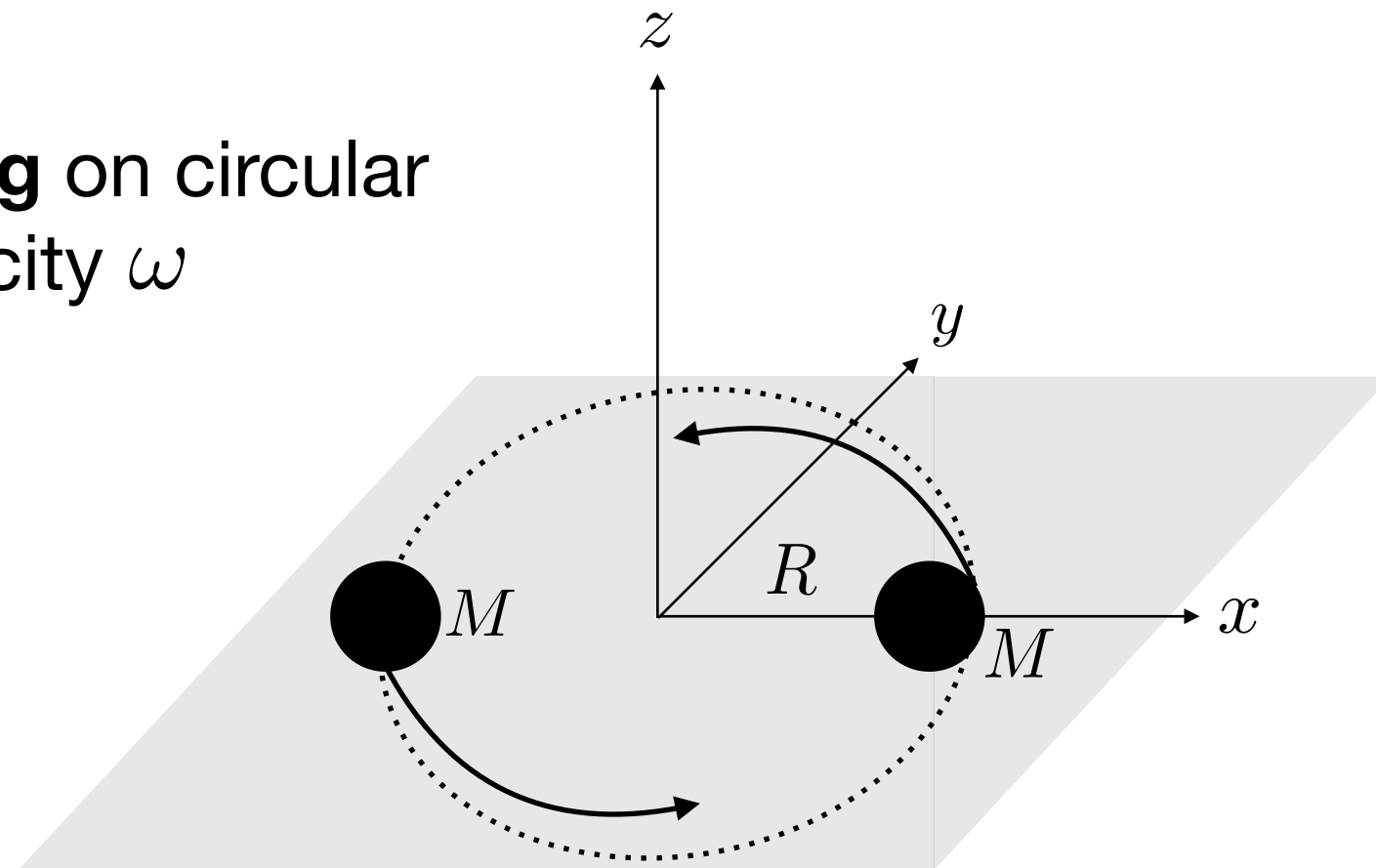
Information encoded in third time derivative of quadrupole moment q^{ij}

(Over)simple Example:

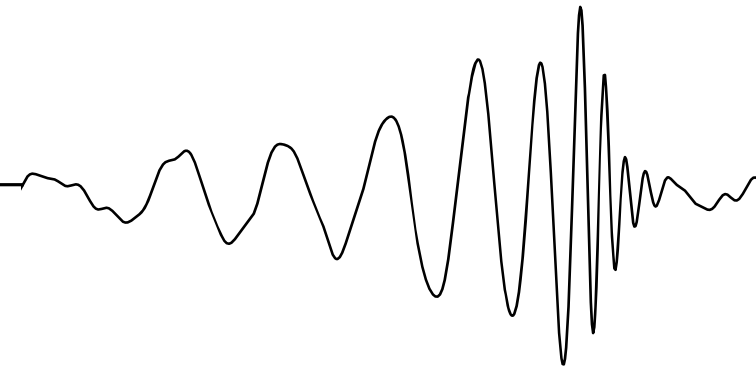
Two objects of mass M **inspiring** on circular orbit of radius R and angular velocity ω

Total energy $E_{\text{tot}} = -\frac{GM^2}{4R(r)}$

Energy loss $\frac{dE}{dt} = \frac{GM^2}{4R^2} \frac{dR}{dt}$



First discovery: Black hole merger



One finds for the quadrupole moment:

$$q^{ij} = M \sum_{n=1}^2 x_n^i x_n^j = MR^2 \begin{pmatrix} 1 + \cos(2\omega t) & \sin(2\omega t) & 0 \\ \sin(2\omega t) & 1 - \cos(2\omega t) & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

Using Kepler's third law $\omega^2 = \frac{GM}{4R^3}$ one can find a differential

equation for the radius $R(t)$ using luminosity and energy loss:

$$\mathcal{L} = \frac{dE}{dt} \Rightarrow R^3 \frac{dR}{dt} = -\frac{8}{5} G^3 M^3$$

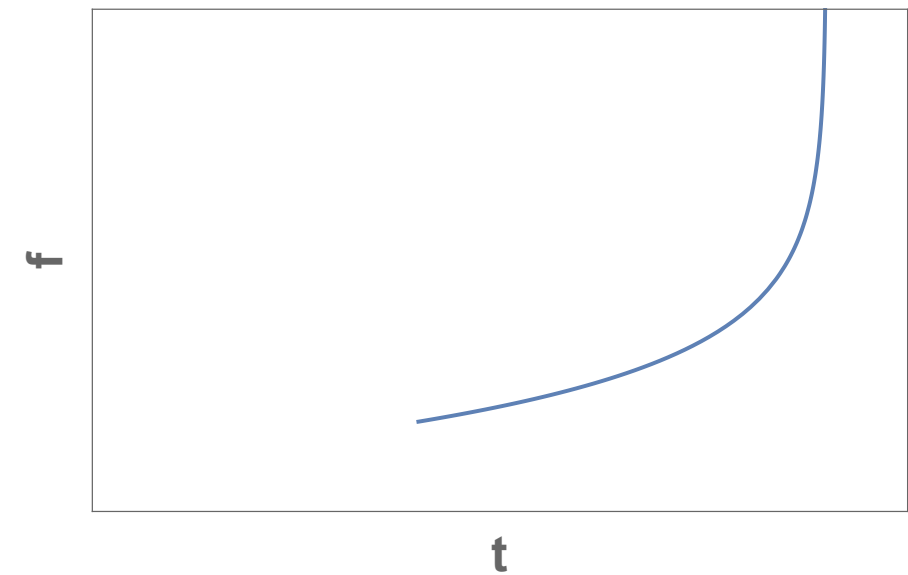
$$R(t) = \left[\frac{32G^3 M^3}{5} (t_{\text{coal}} - t) \right]^{\frac{1}{4}}$$

First discovery: Black hole merger



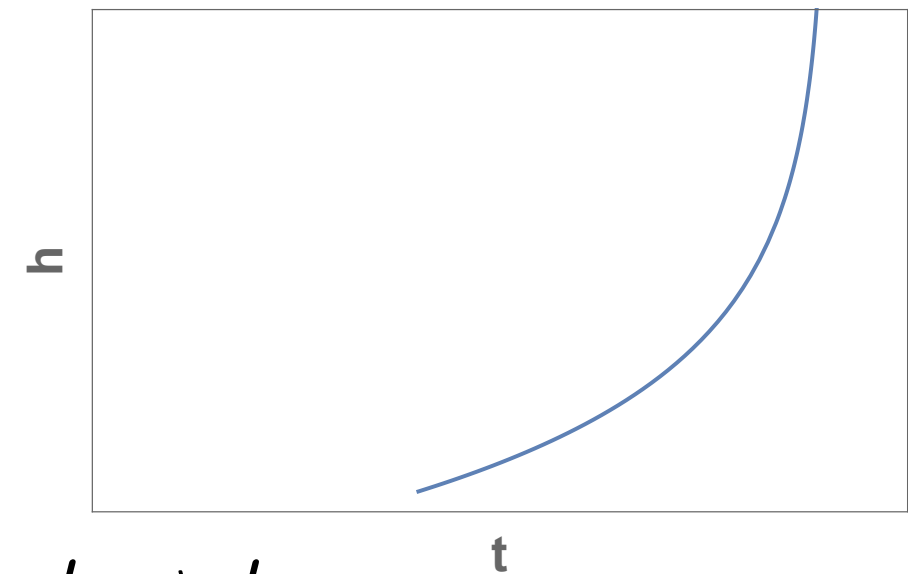
With the time dependent radius we can obtain the frequency

$$f_{\text{GW}} = \frac{2\omega}{2\pi} = \frac{[2 \cdot 5^3]^{\frac{1}{8}}}{8\pi(GM)^{\frac{5}{8}}(t_{\text{coal}} - t)^{\frac{3}{8}}}$$



and the amplitude

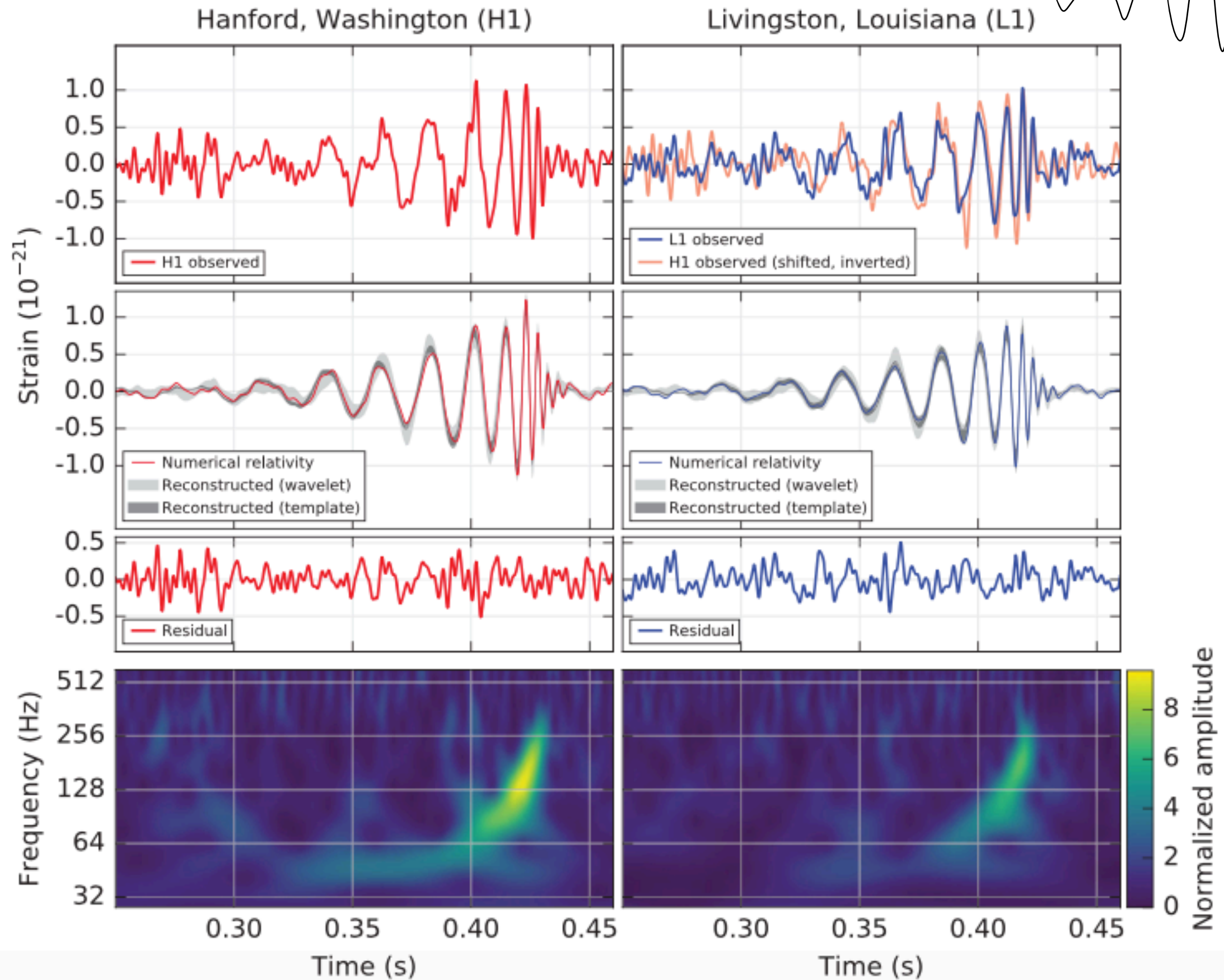
$$h_0 = \frac{1}{r} \left[\frac{5G^5 M^5}{2} \right]^{\frac{1}{4}} \frac{1}{(t_{\text{coal}} - t)^{\frac{1}{4}}}$$



Modulation of both frequency and amplitude as $t \rightarrow t_{\text{coal}}$

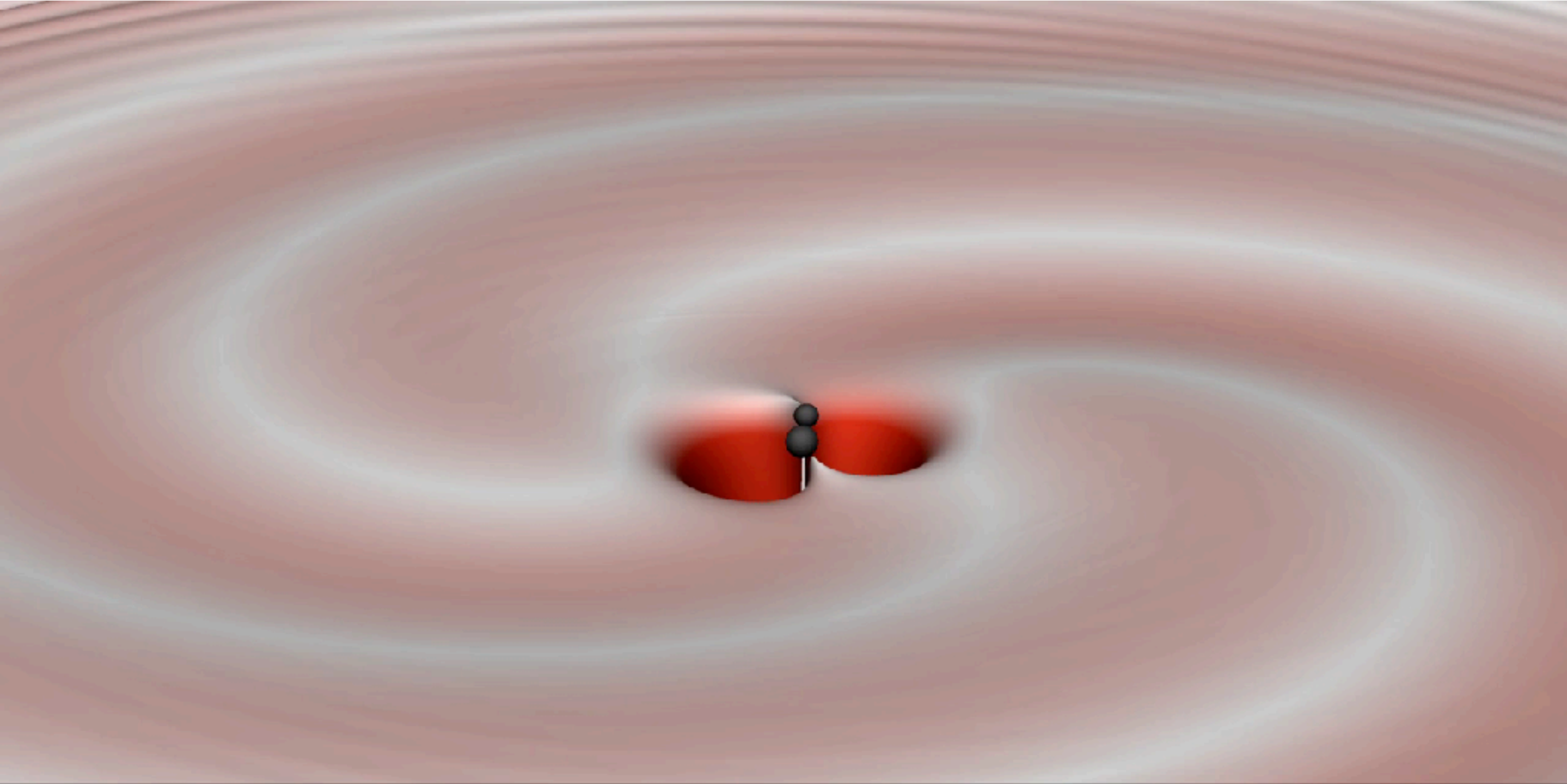
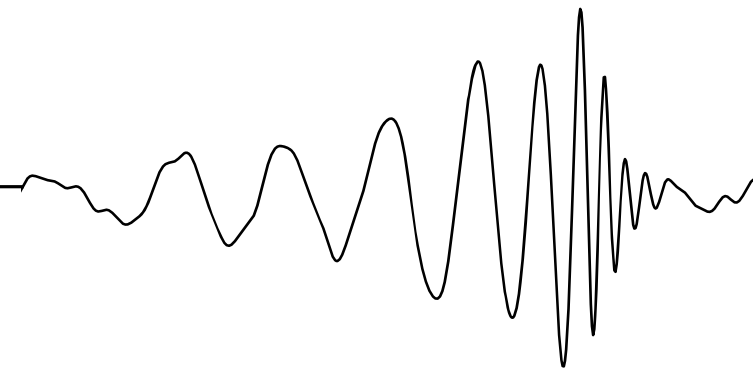
Can extract M from f_{GW} and the distance r to us from h_0

First discovery: The Signal

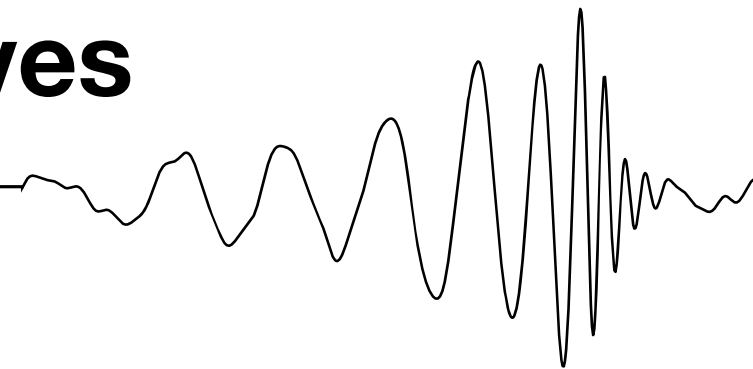


Simulation of GW150914

First black hole merger event detected



Detectable effects of gravitational waves

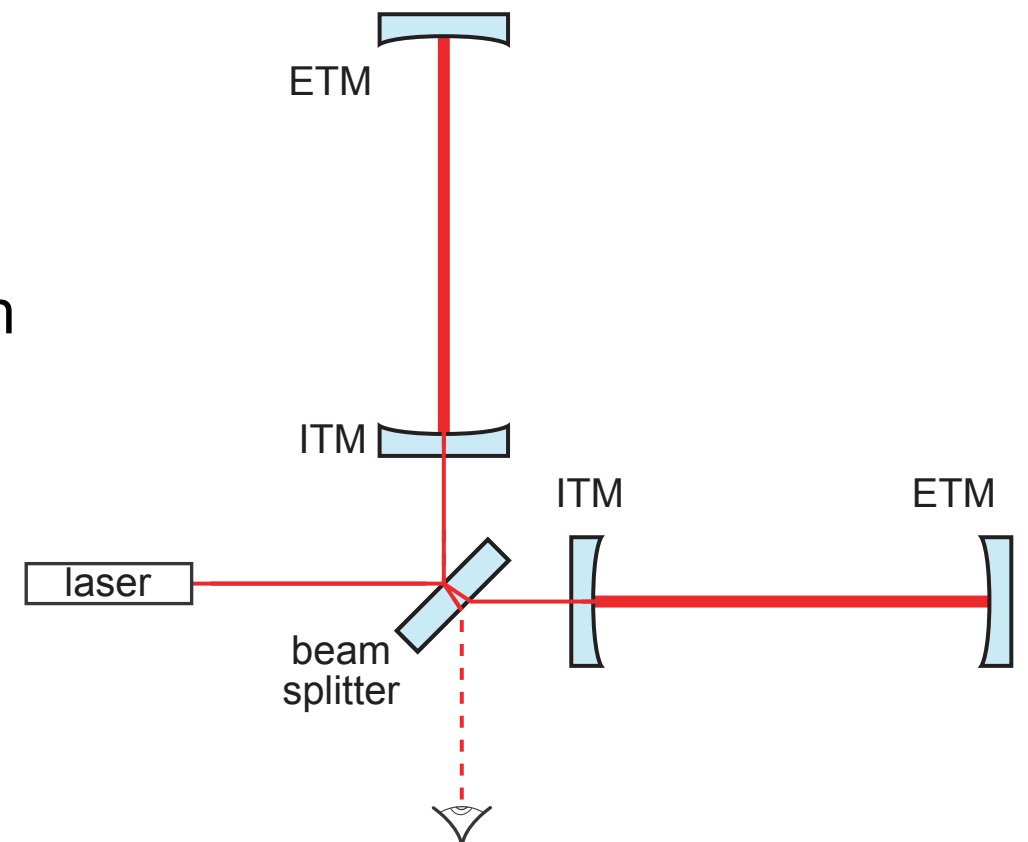


Effect of gravitational waves:

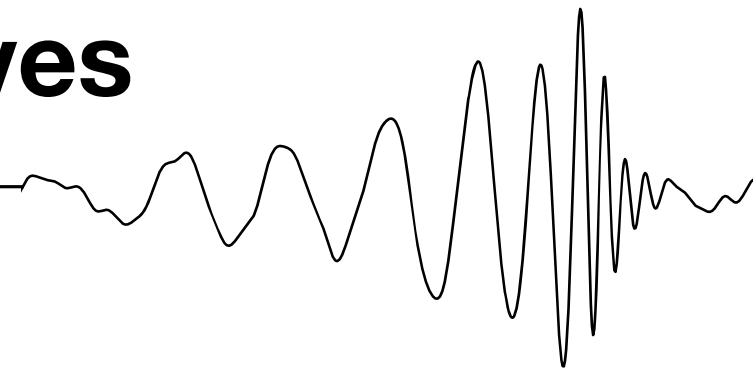
Change of separation of two masses in space, i.e. the arm length of our detector

$$\Delta L \sim \frac{hL}{2} \quad \Rightarrow \quad h \sim \frac{2\Delta L}{L} \lesssim 10^{-21}$$

- ⇒ longer detectors imply larger measurable separations ΔL
- ⇒ very small GW amplitudes h can only be resolved if separation of masses is long enough
- ⇒ effective arm length is increased by bounces between mirrors



Detectable effects of gravitational waves



Effect of gravitational waves:

Change of separation of two masses in space, i.e. the arm length of our detector

$$\Delta L \sim \frac{hL}{2} \quad \Rightarrow \quad h \sim \frac{2\Delta L}{L} \lesssim 10^{-21}$$

Signal-to-noise ratio:

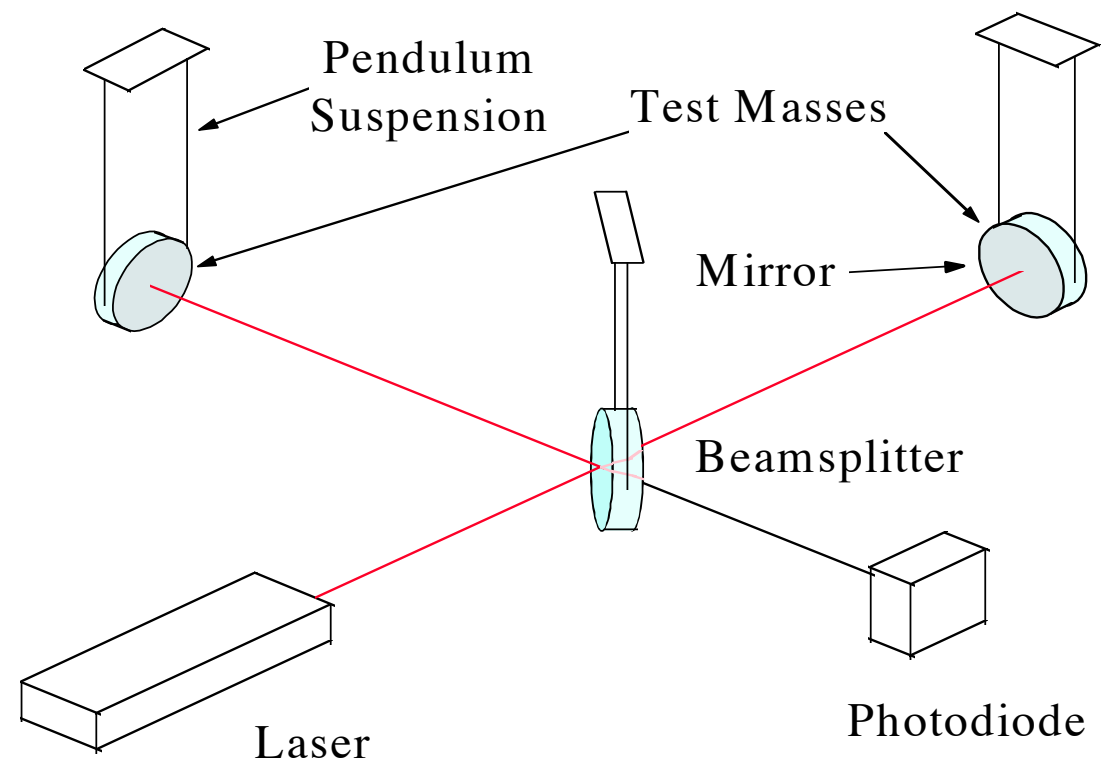
$$\left(\frac{S}{N}\right)^2 \sim \int \frac{|\tilde{h}(f)|^2}{S_n(f)} df$$

For $h \sim 10^{-21}$ the noise level must have an amplitude spectral density

$$\sqrt{S_n(f)} \simeq 10^{-23} \text{ Hz}^{-1/2}$$

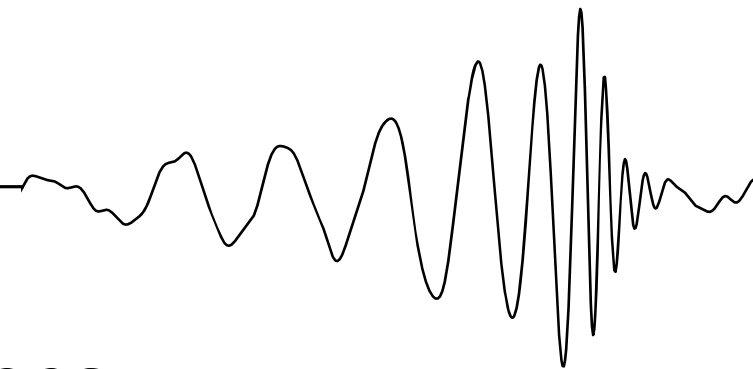
in frequency range of interest

⇒ sensitivity of gravitational wave detectors limited by noise



- System of **freely** suspended masses
- Resonant frequencies of pendulums should be smaller than frequencies of waves

Main noise sources



- **Residual gas noise** ➔ eliminated by vacuum pipes
- **Photoelectron shot noise (PSN)** ≈ 200 Hz

$$h_{\min} \sim \frac{1}{bL} \left(\frac{\hbar c \tilde{\lambda}}{\tau I_0} \right)^{1/2}$$

Minimize w.r.t. I_0 !

- **Radiation pressure noise (RPN)** ~ 10 - 50 Hz

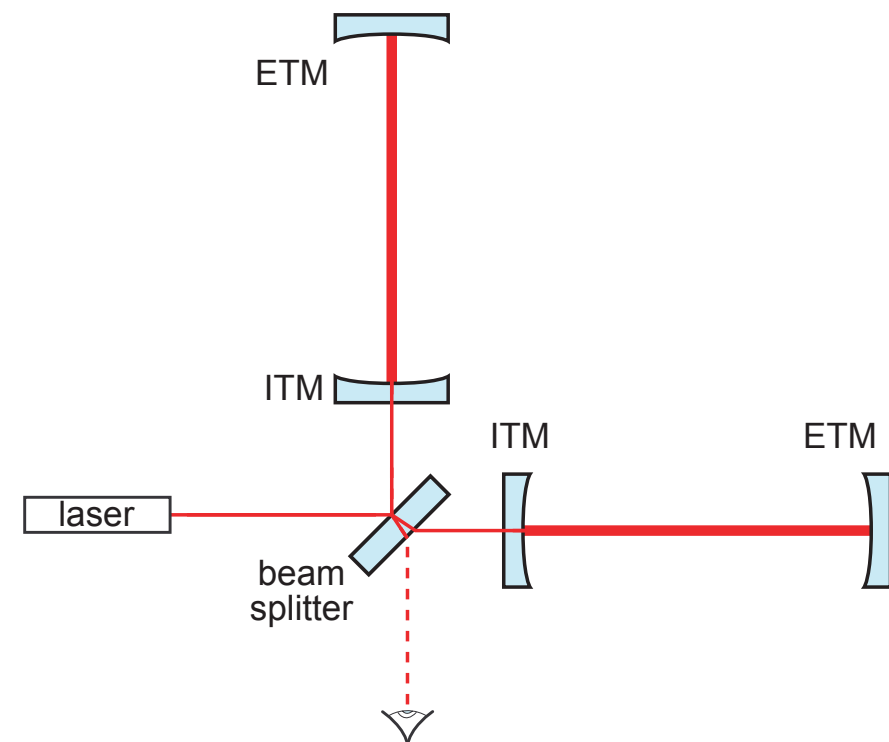
$$h_{\min} \sim \frac{\tau}{m} \frac{b}{L} \left(\frac{\tau \hbar I_0}{c \tilde{\lambda}} \right)^{1/2}$$

High-power Laser needed:

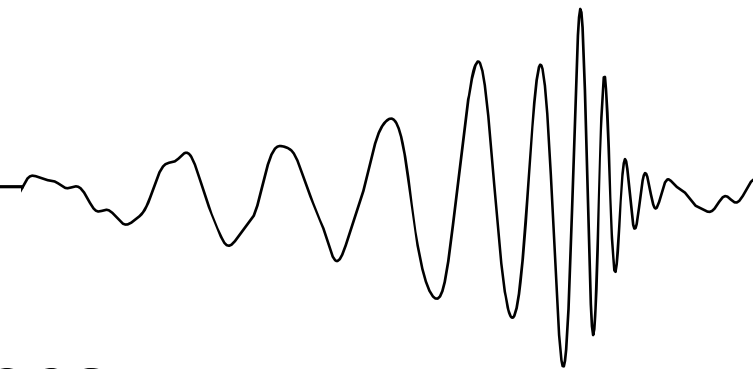
1. Use of Fabry-Pérot resonant cavities

⇒ required laser power reduced to kW

⇒ still very large!



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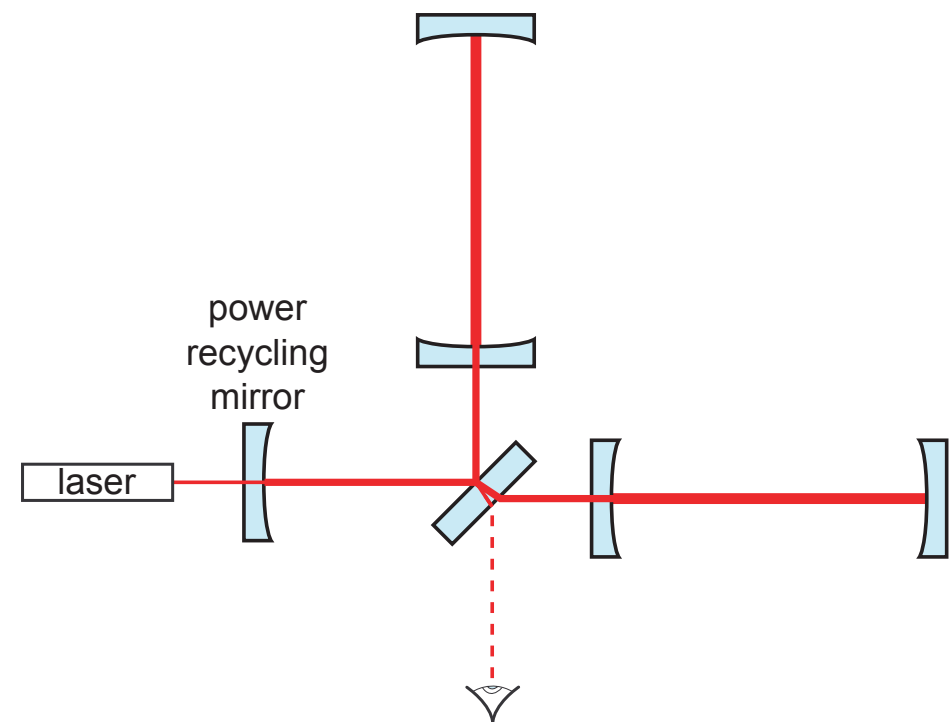
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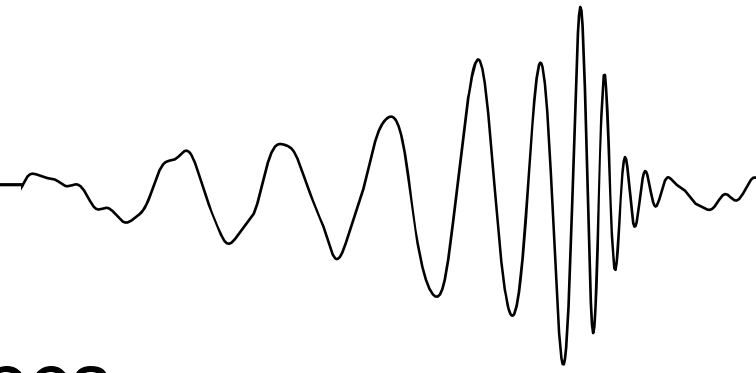
1. Use of Fabry-Pérot resonant cavities
2. Power recycling:

Power recycling mirror in front of laser
⇒ power built-up inside interferometer

⇒ required power reduced to 10 W



Main noise sources



- **Residual gas noise** ➔ eliminated by vacuum pipes
- **Photoelectron shot noise (PSN)** ≥ 200 Hz

$$h_{\min} \sim \frac{1}{bL} \left(\frac{\hbar c \tilde{\lambda}}{\tau I_0} \right)^{1/2}$$

Minimize w.r.t. I_0

- **Radiation pressure noise (RPN)** ~ 10 - 50 Hz

$$h_{\min} \sim \frac{\tau}{m} \frac{b}{L} \left(\frac{\tau \hbar I_0}{c \tilde{\lambda}} \right)^{1/2}$$

$$h_{\min} \sim \frac{1}{L} \left(\frac{\tau \hbar}{m} \right)^{1/2}$$

Standard Quantum Limit!

Note: Operation with optimal intensity for minimal PSN and RPN only depends on test mass!

Heisenberg Uncertainty: $\Delta L \Delta p = \Delta L m \frac{\Delta L}{\tau} = \hbar$

$$\Rightarrow \Delta L^2 = \frac{\tau \hbar}{m}$$

Main noise sources

$$h_{\min} \sim \frac{1}{L} \left(\frac{\tau \hbar}{m} \right)^{1/2}$$

E.g. for Ligo:

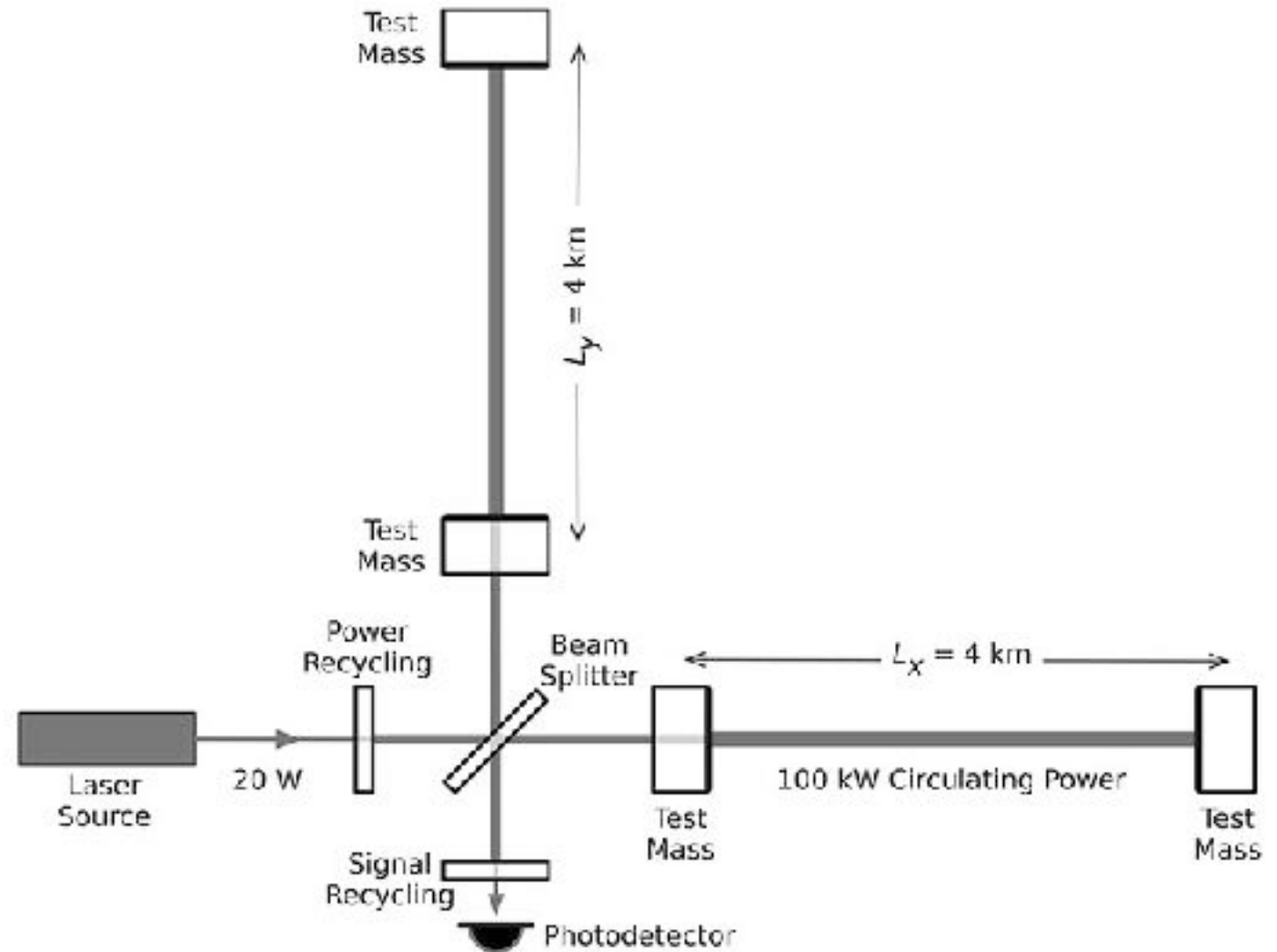
$$m = 100 \text{ kg}, L = 4 \text{ km}, \tau = 1 \text{ ms}$$

$$\Rightarrow h_{\min} \sim 10^{-23}$$

Controlled very well!

Signal recycling:

- Further enhances sensitivity of detector
- “resonating” signal in optical cavities
- Narrowing detection bandwidth by choosing suitable reflexivity
- center of frequency band set by length of formed cavity



Phys. Rev. Lett. **116**, 061102

Main noise sources

- **Thermal noise:**

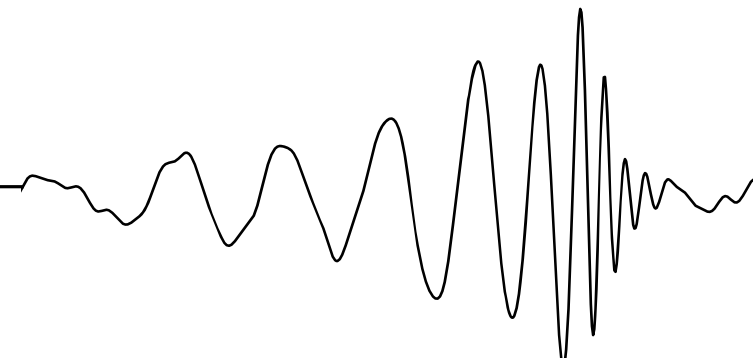
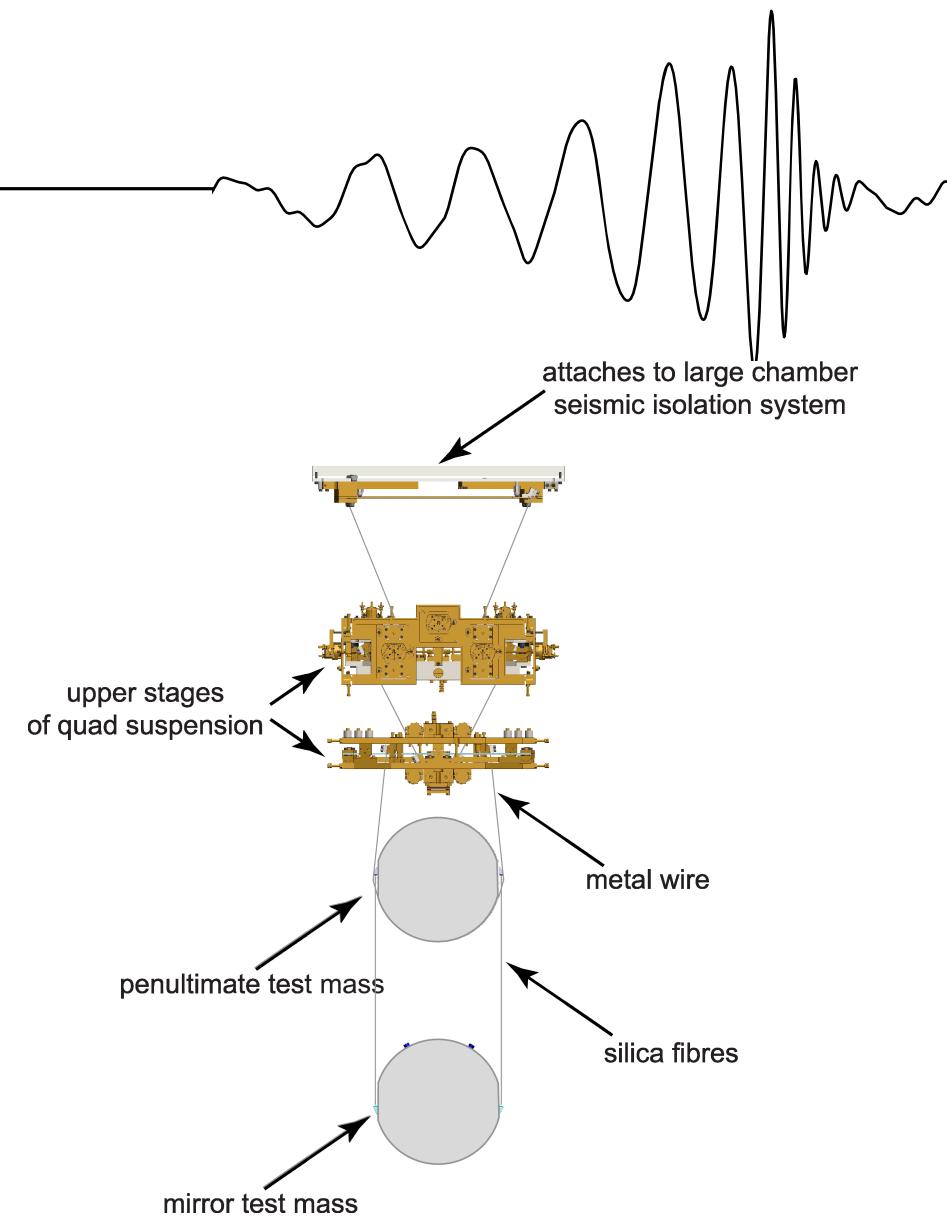
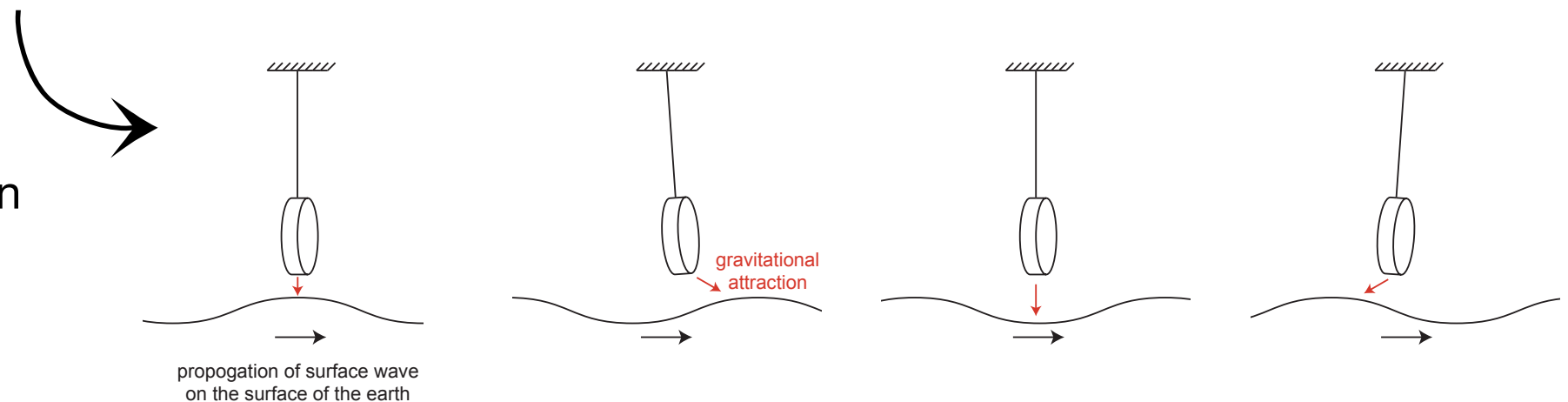
1. resonant modes of pendulum oscillator
2. Internal material oscillator modes

- **Seismic noise ≈ 60 Hz**

1. Vibrations of the ground
2. Gravity gradient noise (caused by seismic surface waves)

Options:

1. Monitor and subtraction method
2. Far from sources
 - ⇒ quiet location
 - ⇒ underground
 - ⇒ space missions



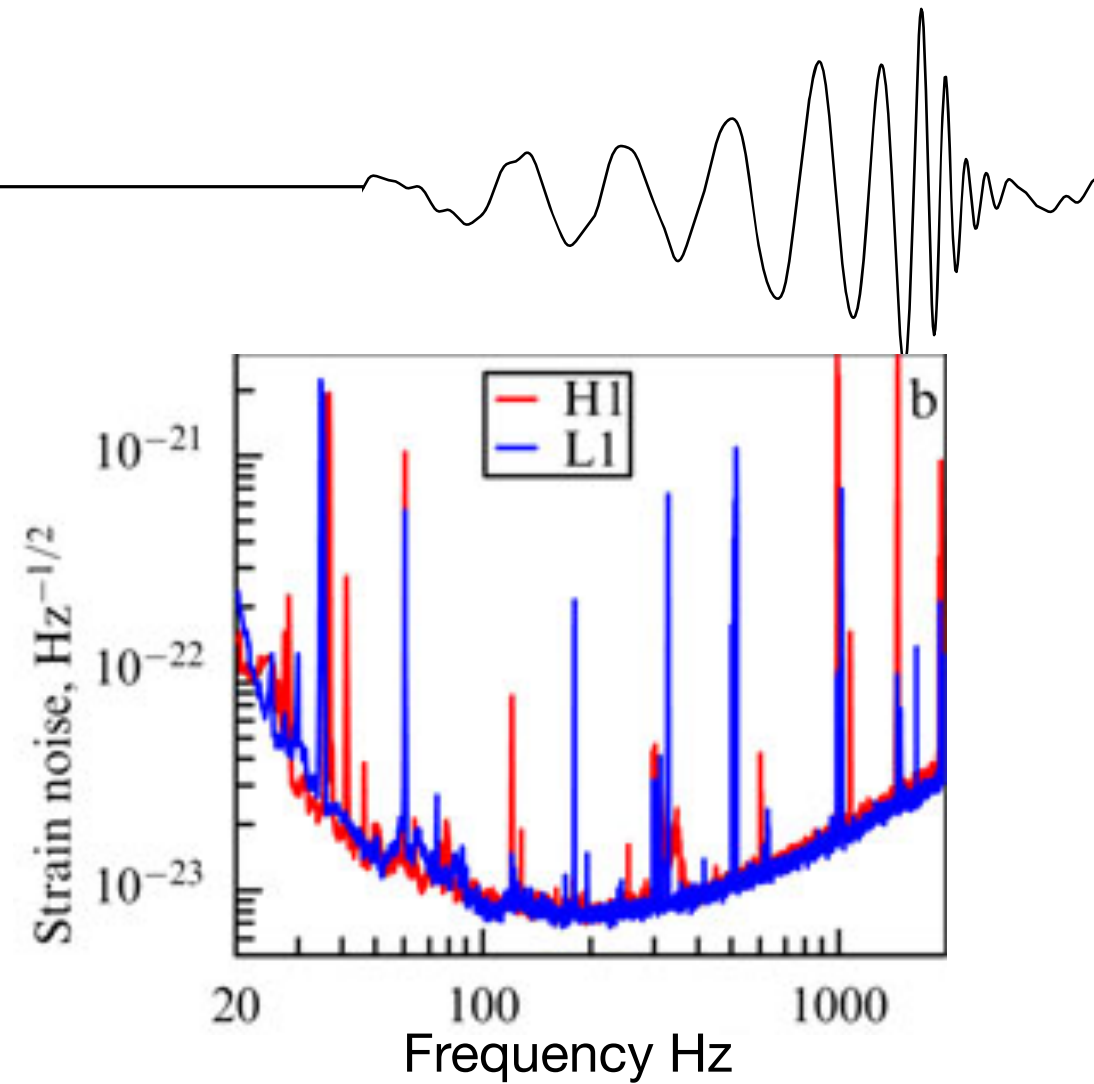
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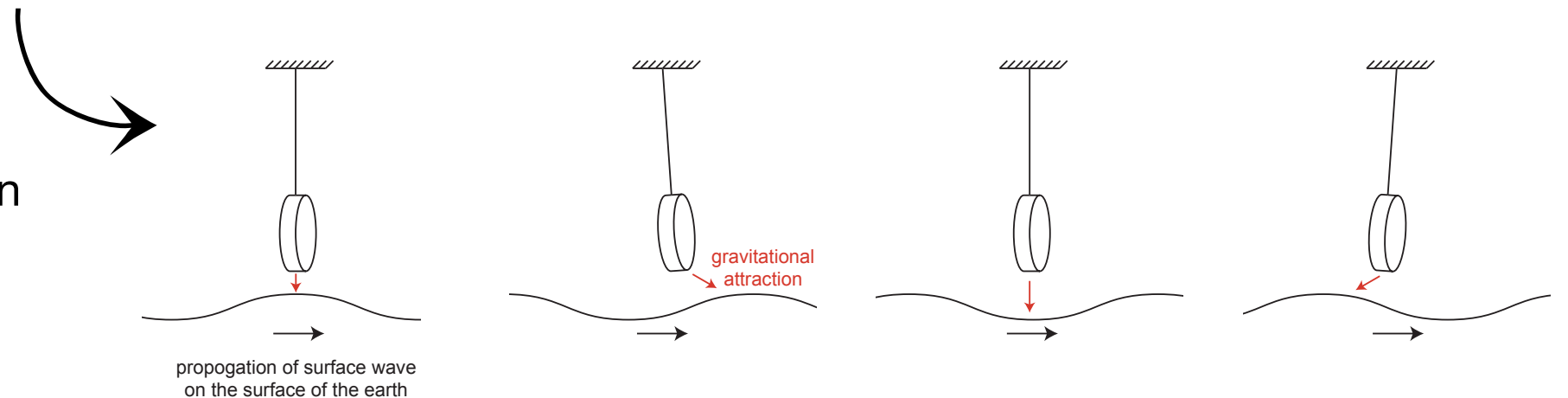
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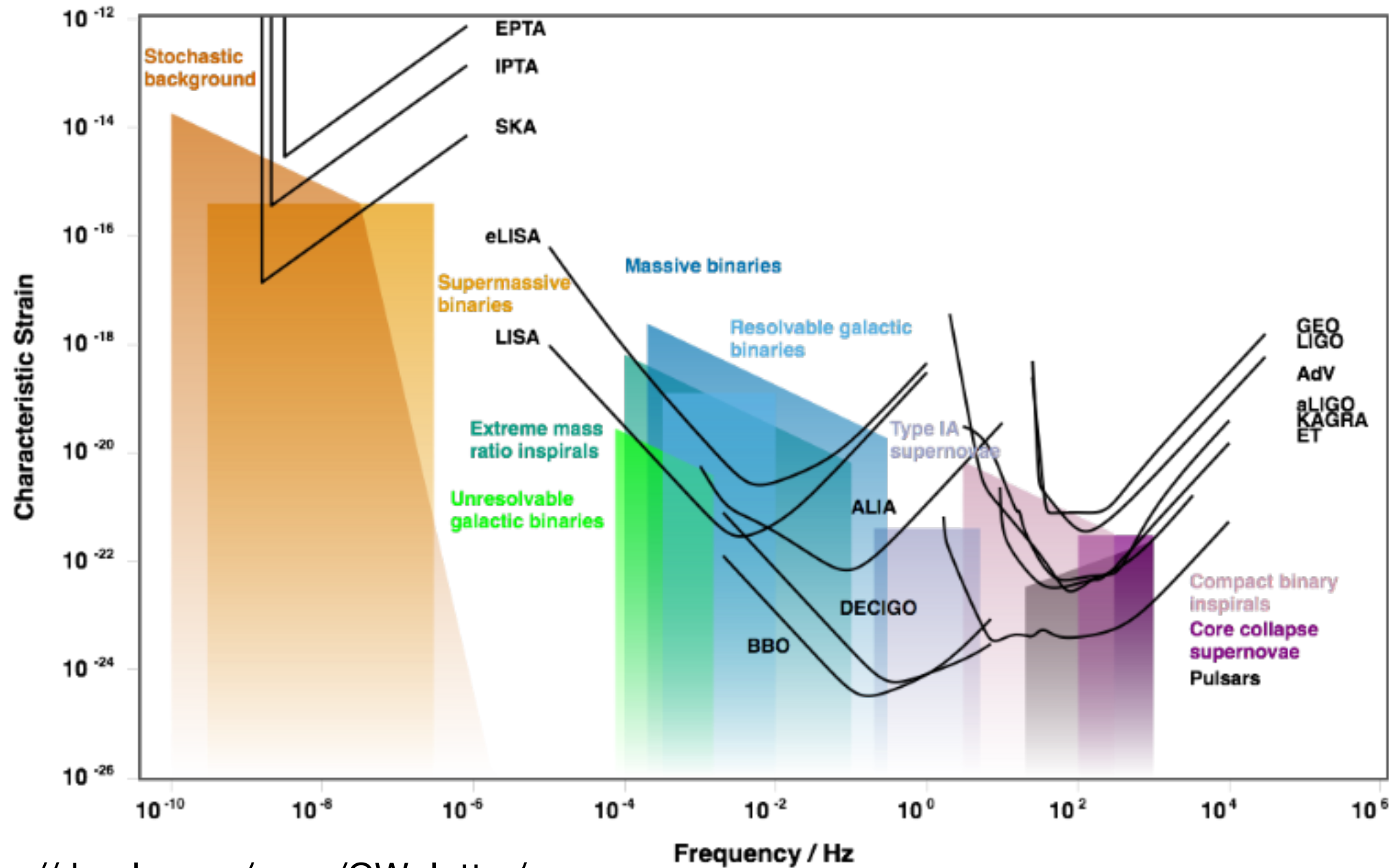
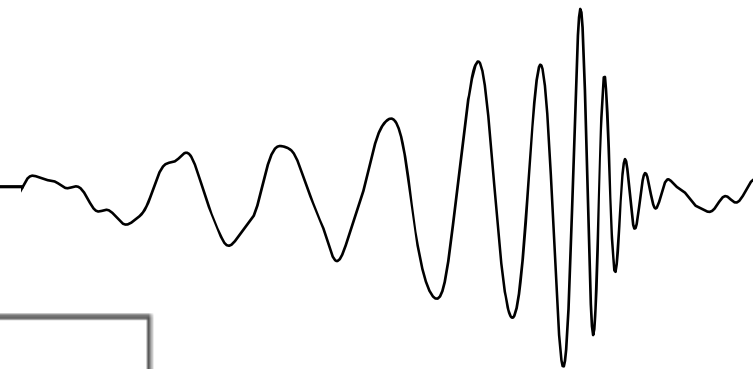


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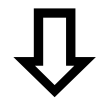
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Reach of gravitational wave detectors



Technology
(and going underground)



Objects further away

More massive objects



From earth to space

<http://rhcole.com/apps/GWplotter/>

Conclusions

- Entered a new area of gravitational wave astronomy
- First direct black hole evidence with aLIGO (5 BH events by now)
- Space-based missions will increase sensitivity to new sources like rotating neutron stars, etc...
- A bright future lies ahead of us

