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# NLO QCD corrections to Higgs boson pair production via gluon fusion

Seraina Glaus

Theory Group LTP, Paul Scherrer Institute & University Zurich In collaboration with M. Spira, M. Mühlleitner, J. Baglio, F. Campanario, J. Streicher

# Outline

- Motivation
- Objective
- Previous work
- NLO Cross section
  - Virtual Corrections
  - Real Corrections
- Numerical Analysis
- Conclusions/Outlook



#### Detection of a Higgs boson with a mass ~ 125 GeV





- Further investigations of the properties of the detected particle for a unique association to a model
- Higgs mass, coupling strengths, spin and CP already determined
- self-coupling strength still unknown



$$V = \frac{\lambda}{2} (\phi^{\dagger}\phi - \frac{v^2}{2})^2$$

$$\lambda = \frac{M_H^2}{2v^2}$$

$$\lambda_{H^3} = 3 \frac{M_H^2}{v}$$

 $\lambda_{H^4} = 3$ 



#### Higgs boson pair production

#### **Production channel**

#### Cross section





Baglio, Djouadi, Gröber, Mühlleitner, Quevillon, Spira



#### **Uncertainties:**



Baglio, Djouadi, Gröber, Mühlleitner, Quevillon, Spira



#### Gluon fusion gg $\rightarrow$ HH: loop induced

Complete calculation of the NLO QCD corrections (2-loop) considering the top- und bottom mass dependences in the context of the Standard Model  $\rightarrow$  2-loop integrals with 3 kinematical parameter ratios





# **Previous work**

- Virtual & real (N)NLO QCD corrections in large top mass limit: ~100%
   Dawson,Dittmaier,Spira de Florian,Mazzitelli
- Large top mass expansion: ~ ±10%

$$\sigma = \sigma_0 + \frac{\sigma_1}{m_t^2} + \dots + \frac{\sigma_4}{m_t^8}$$

Grigo, Hoff, Melnikov, Steinhauser

Grigo, Melnikov, Steinhauser

- NLO mass effects of the real NLO correction alone
  - ~ -10 %

Frederix, Frixione, Hirschi, Maltoni, Mattelaer, Torrielli, Vryonidou, Zaro

 NLO QCD corrections including the full top mass dependence
 Borowka, Greiner, Heinrich, Jones, Kerner, S

Borowka, Greiner, Heinrich, Jones, Kerner, Schlenk, Schubert, Zirke



# **NLO Corrections**





#### **NLO Corrections**

$$\sigma_{\text{NLO}}(pp \to HH + X) = \sigma_{\text{LO}} + \Delta \sigma_{\text{virt}} + \Delta \sigma_{gg} + \Delta \sigma_{gq} + \Delta \sigma_{q\bar{q}},$$

$$\begin{split} \sigma_{\text{LO}} &= \int_{\tau_0}^1 d\tau \; \frac{d\mathcal{L}^{gg}}{d\tau} \; \hat{\sigma}_{\text{LO}}(Q^2 = \tau s) \\ \Delta \sigma_{\text{virt}} &= \frac{\alpha_s(\mu)}{\pi} \int_{\tau_0}^1 d\tau \; \frac{d\mathcal{L}^{gg}}{d\tau} \; \hat{\sigma}_{\text{LO}}(Q^2 = \tau s) \; C \\ \Delta \sigma_{gg} &= \frac{\alpha_s(\mu)}{\pi} \int_{\tau_0}^1 d\tau \; \frac{d\mathcal{L}^{gg}}{d\tau} \int_{\tau_0/\tau}^1 \frac{dz}{z} \; \hat{\sigma}_{\text{LO}}(Q^2 = z\tau s) \left\{ -z P_{gg}(z) \log \frac{M^2}{\tau s} \right. \\ &+ d_{gg}(z) + 6[1 + z^4 + (1 - z)^4] \left( \frac{\log(1 - z)}{1 - z} \right)_+ \right\} \\ \Delta \sigma_{gq} &= \frac{\alpha_s(\mu)}{\pi} \int_{\tau_0}^1 d\tau \sum_{q,\bar{q}} \frac{d\mathcal{L}^{gq}}{d\tau} \int_{\tau_0/\tau}^1 \frac{dz}{z} \; \hat{\sigma}_{\text{LO}}(Q^2 = z\tau s) \left\{ -\frac{z}{2} P_{gq}(z) \log \frac{M^2}{\tau s(1 - z)^2} \right. \\ &+ d_{gg}(z) \right\} \\ \Delta \sigma_{q\bar{q}} &= \frac{\alpha_s(\mu)}{\pi} \int_{\tau_0}^1 d\tau \sum_q \frac{d\mathcal{L}^{q\bar{q}}}{d\tau} \int_{\tau_0/\tau}^1 \frac{dz}{z} \; \hat{\sigma}_{\text{LO}}(Q^2 = z\tau s) \; d_{q\bar{q}}(z) \\ \left. C \to \pi^2 + \frac{11}{2} + C_{\Delta\Delta}, \qquad d_{gg} \to -\frac{11}{2}(1 - z)^3, \qquad d_{gq} \to \frac{2}{3}z^2 - (1 - z)^2, \qquad d_{q\bar{q}} \to \frac{32}{27}(1 - z)^3 \end{split}$$



#### Set-up

- Associate the 47 two-loop box diagrams to similar topologies
- Generate matrix elements for all possible diagrams (by hand → *Reduce, Mathematica*)
- Use dimensional regularisation:  $n = 4 2\epsilon$
- Perform Feynman parametrisation  $\rightarrow$  6-dimensional integrals

$$\frac{1}{A_1^{\alpha_1}\cdots A_n^{\alpha_n}} = \frac{\Gamma(\alpha_1+\ldots+\alpha_n)}{\Gamma(\alpha_1)\cdots\Gamma(\alpha_n)} \int_0^1 du_1 \int_0^{1-u_1} du_2\cdots \int_0^{1-u_1-\ldots-u_{n-2}} du_{n-1} \frac{u_1^{\alpha_1-1}\cdots u_{n-1}^{\alpha_n-1}(1-u_1-\ldots-u_{n-1})^{\alpha_n-1}}{[u_1A_1+\ldots+u_{n-1}A_{n-1}+(1-u_1-\ldots-u_{n-1})A_n]^{\alpha_1+\ldots+\alpha_n}},$$



#### Divergences

 Extraction of the ultraviolet divergences of the matrix elements using endpoint subtractions of the 6-dimensional Feynman integrals

$$\int_0^1 dx \, \frac{f(x)}{(1-x)^{1-\epsilon}} = \int_0^1 dx \, \frac{f(1)}{(1-x)^{1-\epsilon}} + \int_0^1 dx \, \frac{f(x) - f(1)}{(1-x)^{1-\epsilon}}$$
$$= \frac{f(1)}{\epsilon} + \int_0^1 dx \, \frac{f(x) - f(1)}{1-x} + \mathcal{O}(\epsilon)$$

 Extraction of the infrared and collinear divergences using a proper subtraction of the integrand

denominator:  $N = ar^2 + br + c$   $N_0 = br + c$   $a = \mathcal{O}(\rho)$   $b = 1 + \mathcal{O}(\rho)$  $c = -\rho_s x(1-x)(1-s)t$ 

$$\int_{0}^{1} \mathrm{d}\vec{x} \mathrm{d}r \frac{rH(\vec{x},r)}{N^{3+2\epsilon}} = \int_{0}^{1} \mathrm{d}\vec{x} \mathrm{d}r \Big\{ \Big( \frac{rH(\vec{x},r)}{N^{3+2\epsilon}} - \frac{rH(\vec{x},0)}{N_{0}^{3+2\epsilon}} \Big) + \frac{rH(\vec{x},0)}{N_{0}^{3+2\epsilon}} \Big\}$$

$$\uparrow$$
Taylor expansion in  $\epsilon$ 
analytical r-integration



#### Divergences

- Integration by parts due to numerical instabilities at the thresholds

 $M^2_{HH} > 4m^2_Q \Rightarrow m^2_Q \to m^2_Q (1 - i\bar{\epsilon}) \text{ with } \bar{\epsilon} \ll 1$ 

$$\int_0^1 dx \frac{f(x)}{(a+bx)^3} = \frac{f(0)}{2a^2} - \frac{f(1)}{2(a+b)^2} + \int_0^1 dx \frac{f'(x)}{2(a+bx)^2}$$

- gluon rescattering: threshold for  $M_{HH}^2 > 0$ 





#### Renormalisation

 $\alpha_{\text{S}}$  and  $m_{\text{Q}}$  need to be renormalised

 $\rightarrow$   $\alpha_{s}$  in  $\overline{MS}$  with N<sub>F</sub> = 5

→ mq on shell

$$\delta\sigma = \delta\alpha_s \frac{\delta\sigma_{LO}}{\delta\alpha_s} + \delta m_t \frac{\delta\sigma_{LO}}{\delta m_t}$$

Subtraction of the heavy-top limit  $\rightarrow$  virtual mass effects only (infrared finite)  $\Delta C_{mass} = C^0 - C_{HTL}^0$ 

Adding back the results of HPAIR (heavy-top limit)

$$C = C_{HTL} + \Delta C_{mass}$$

$$\uparrow$$
HPAIR



#### **Remaining Diagrams**

- Triangular diagrams
  - single Higgs case



#### One-particle reducible diagrams

- → analytical results
  - $(H \to Z\gamma)$

see e.g. Degrassi, Giardino, Gröber





# **Real Corrections**

#### Processes: $gg,q\bar{q} \rightarrow HHg; gq \rightarrow HHq$



- Full matrix elements generated with LoopTools
- Matrix elements in the heavy-top limit rescaled locally by massive LO matrix elements (with adjusted kinematics) 
   → subtracted 
   → free of divergences
- Adding back the results of HPAIR (heavy-top limit)



# **Numerical Analysis**

Use Vegas for numerical integration P. Lepage

 $\frac{d\sigma}{dQ}$ Calculation of differential cross section

$$\frac{\sigma}{Q^2}$$
,  $(Q^2 = m_{HH}^2)$ 

$$Q^{2} \frac{d\Delta \sigma_{virt}}{dQ^{2}} = \tau \frac{d\mathcal{L}^{gg}}{d\tau} \hat{\sigma}_{virt}(Q^{2}) \big|_{\tau = \frac{Q^{2}}{s}}$$

$$\uparrow$$
partonic cross section

(seven-dimensional integrals)

- Thresholds:  $m_Q^2 \rightarrow m_Q^2(1-i\overline{\epsilon})$ different  $\overline{\epsilon} \longrightarrow$  Richardson extrapolation  $\longrightarrow \overline{\epsilon} = 0$  (NWA)
  - Integration of real corrections straight forward



# **Conclusions and Outlook**

- Full NLO calculation close to finalisation
   numerical results soon
- NLO mass effects expected in the 10-20 % range
- Independent cross check of existing results

Outlook:

 Extension of the calculation to BSM-Higgs scenarios (dim 6, 2HDM)





# **Back-up**

#### **Richardson Extrapolation**

 $M_{2}[f(h), f(2h)] = 2f(h) - f(2h)$   $M_{4}[f(h), f(2h), f(4h)] = (8f(h) - 6f(2h) + f(4h))/3$  $M_{8}[f(h), f(2h), f(4h), f(8h)] = (64f(h) - 56f(2h) + 14f(4h) - f(8h))/21$ 



