

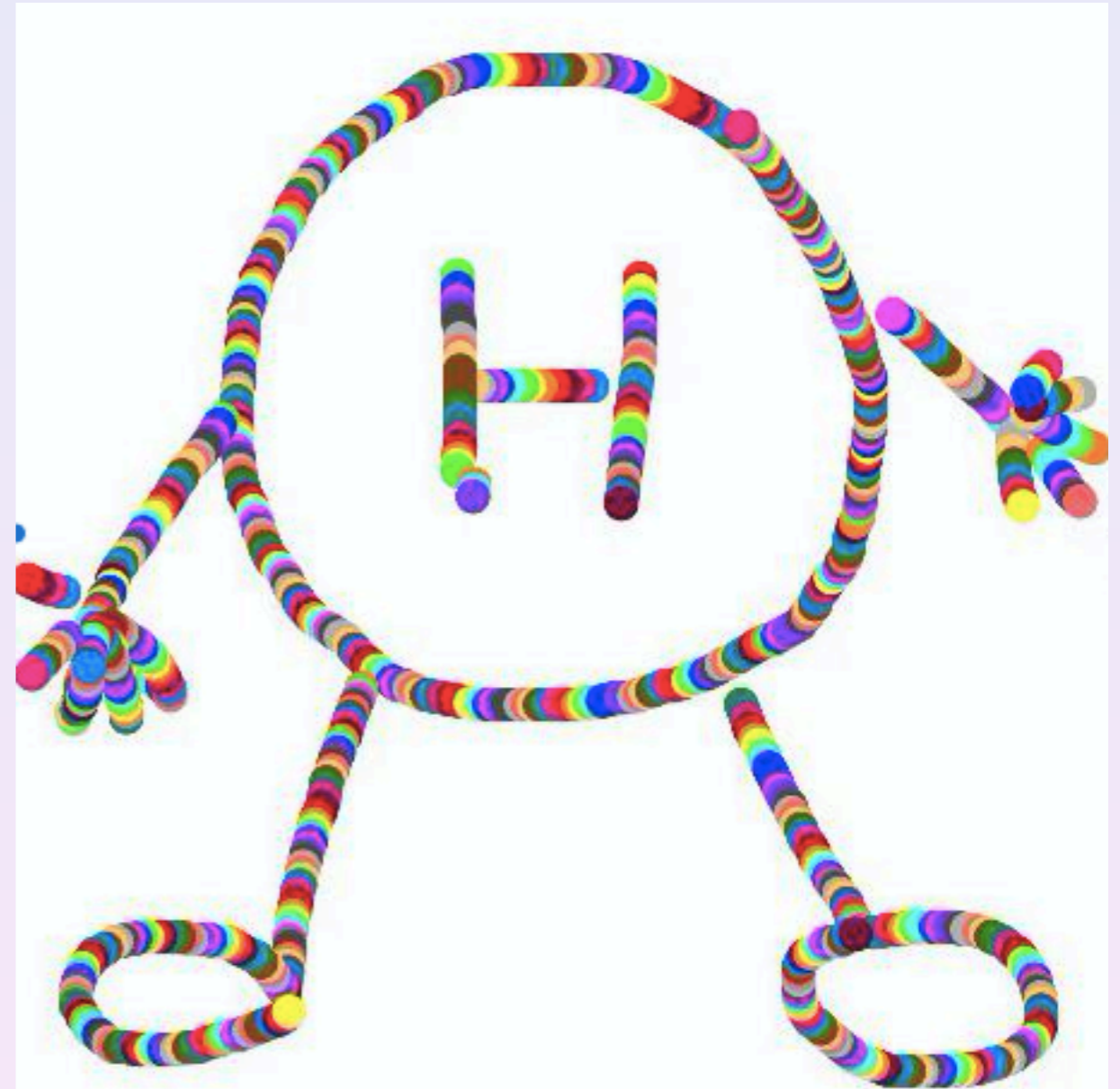
Composite Higgs Models

Ramona Gröber

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Durham University

Higgs Couplings

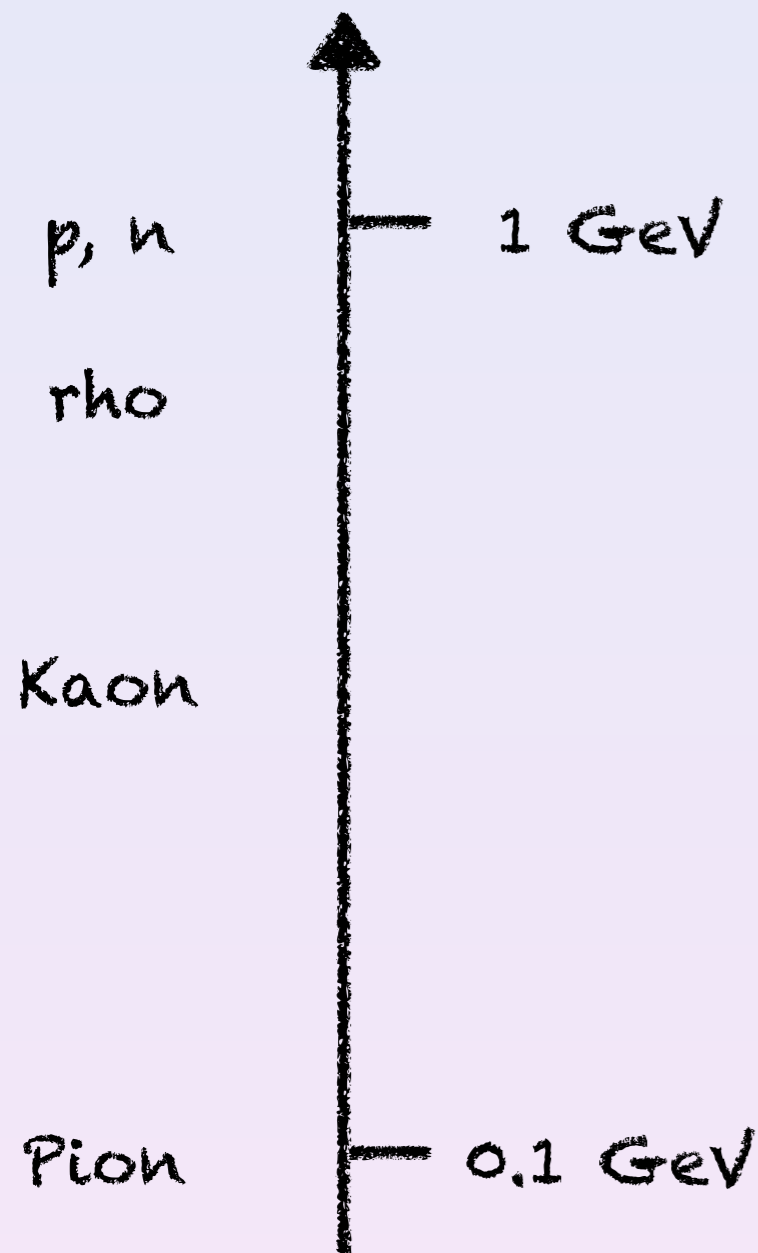
10/11/2017



Outline

- Introduction to Composite Higgs Models
 - Composite scalars in QCD
 - Composite Higgs Models
 - Lagrangian
 - Partial compositeness
- How can we test Composite Higgs Models at the LHC?
 - Higgs couplings
 - Resonances
 - Non-linearities
- Non-minimal Composite Higgs Models
 - Dark matter

Composite Scalars in QCD



Spectrum of QCD:

Scalars without hierarchy problem

The QCD scale cuts their quantum corrections naturally off

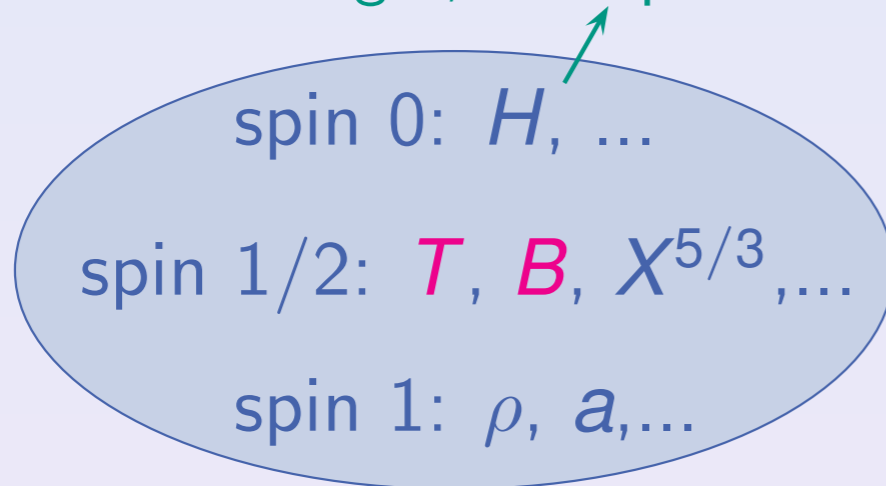
Pions as pseudo-Nambu Goldstone bosons, naturally much lighter than QCD scale

QCD scale is natural, can be largely separated from any high scale due to logarithmical running

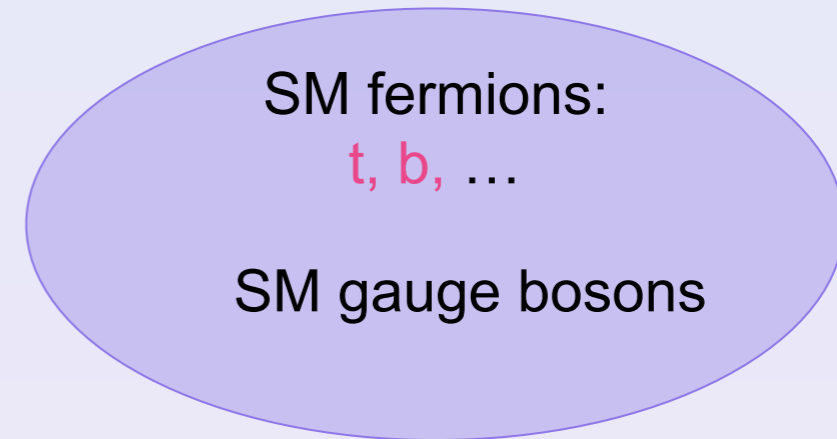
Composite Higgs Models

[Georgi, Kaplan '84]

light, since pseudo-Goldstone boson



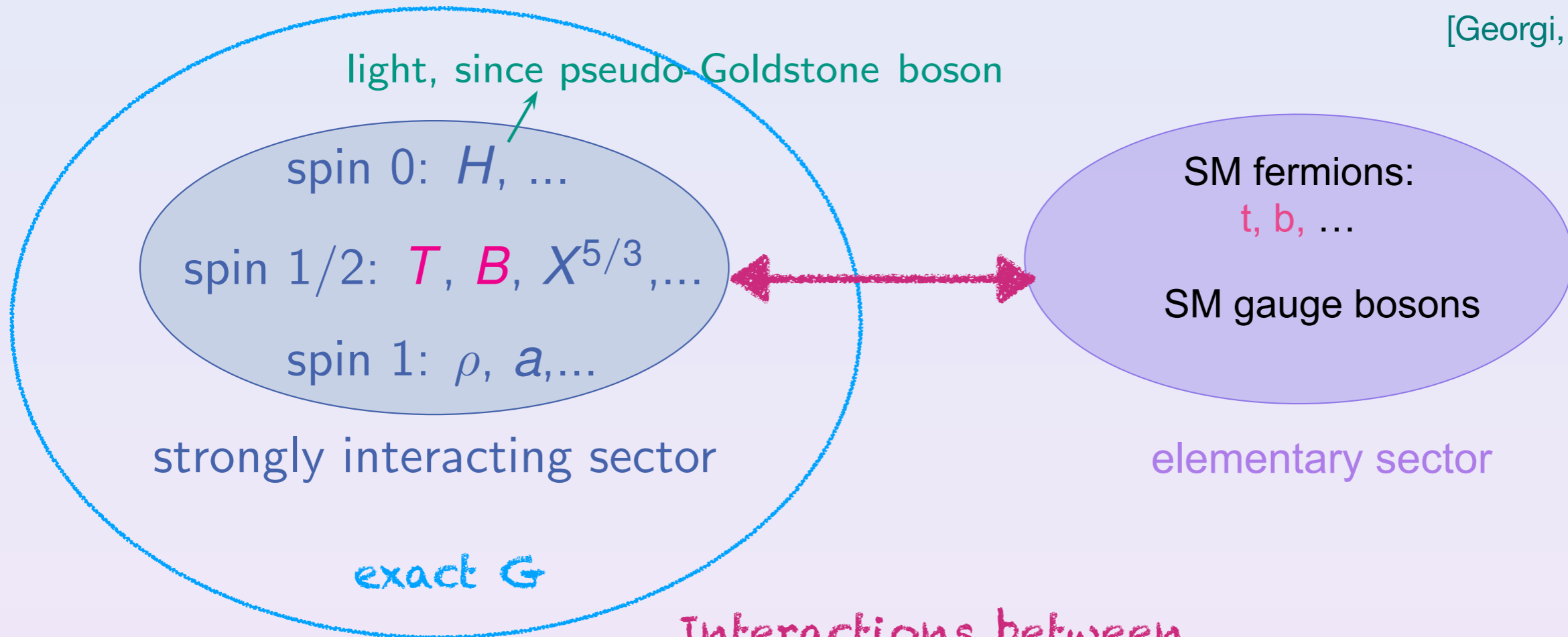
strongly interacting sector



elementary sector

Composite Higgs Models

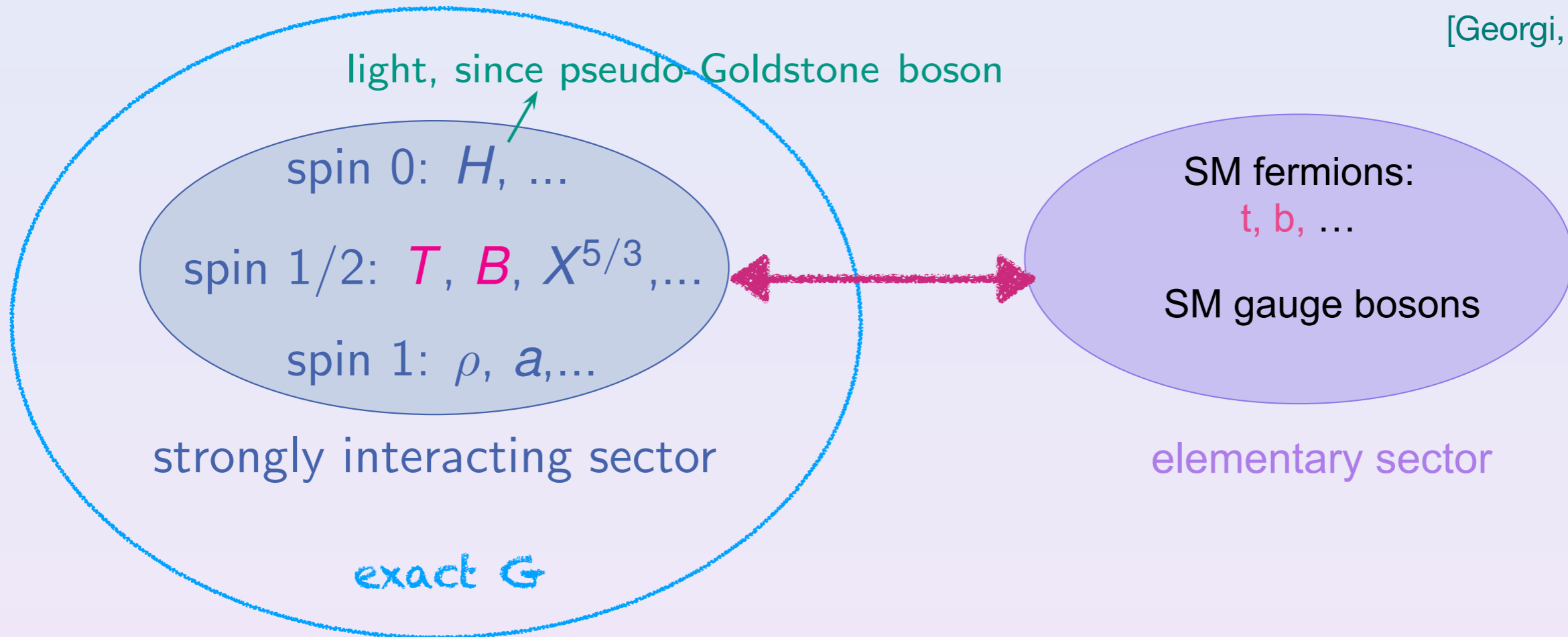
[Georgi, Kaplan '84]



Interactions between elementary and strongly interacting sector break G at scale f ; coset G/H should contain the SM gauge group

Composite Higgs Models

[Georgi, Kaplan '84]



Higgs mass generated by quantum corrections

Partial compositeness: top quark mass generated by linear interactions with strongly interacting sector

Lagrangian for a Composite Higgs

Requirements:

- the unbroken group H contains the SM gauge group G

$$G \rightarrow H \subset SU(2) \times U(1)$$

- The coset space G/H contains at least 4 Goldstone bosons (corresponding to the Higgs doublet)
- Custodial symmetry: to be conform with EWPTs

$$H \subset SU(2)_R \times SU(2)_L$$

Minimal model: [Agashe, Contino, Pomarol '04]

$$SO(5) \times U(1) / SO(4) \times U(1)$$

contains 4 Goldstone bosons (3 the usual Goldstones + Higgs boson)

Lagrangian for a Composite Higgs

Lagrangian from CCWZ construction: [Callan, Coleman, Wess, Zumino '69]

Goldstone matrix (X denote the broken generators)

$$U = e^{-i\frac{\sqrt{2}}{f}\pi^{\hat{a}}X^a}$$

Defining

$$iU^{-1}\partial_{\mu}U = d_{\mu,a}X^a + e_{\mu,a}T^a$$

the Lagrangian is

$$\mathcal{L} = \frac{f^2}{4}\text{Tr}(d_{\mu}d^{\mu})$$

Minimal model:

more details in [Panico, Wulzer '15]

We take a $\Sigma_0 = (0, 0, 0, 0, 1)^T$ to project on coset space.

Then

$$\mathcal{L} = \frac{f^2}{2}(D_{\mu}\Sigma)^{\dagger}D^{\mu}\Sigma, \quad \text{with } \Sigma = U\Sigma_0,$$

$$\mathcal{L} = \frac{1}{2}\partial_{\mu}h\partial^{\mu}h + \frac{g^2}{4}f^2\sin^2\left(\frac{h}{f}\right)W_{\mu}^{+}W^{\mu-} + \frac{g^2}{8c_W^2}f^2\sin^2\left(\frac{h}{f}\right)Z_{\mu}Z^{\mu}.$$

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Leads to non-linearities

$$\xi = \frac{v^2}{f^2} = \sin^2\frac{\langle h \rangle}{f}$$

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Leads to non-linearities

$$\xi = \frac{v^2}{f^2} = \sin^2\frac{\langle h \rangle}{f}$$

$$\frac{g_{hVV}}{g_{hVV}^{SM}} = \sqrt{1 - \xi}$$

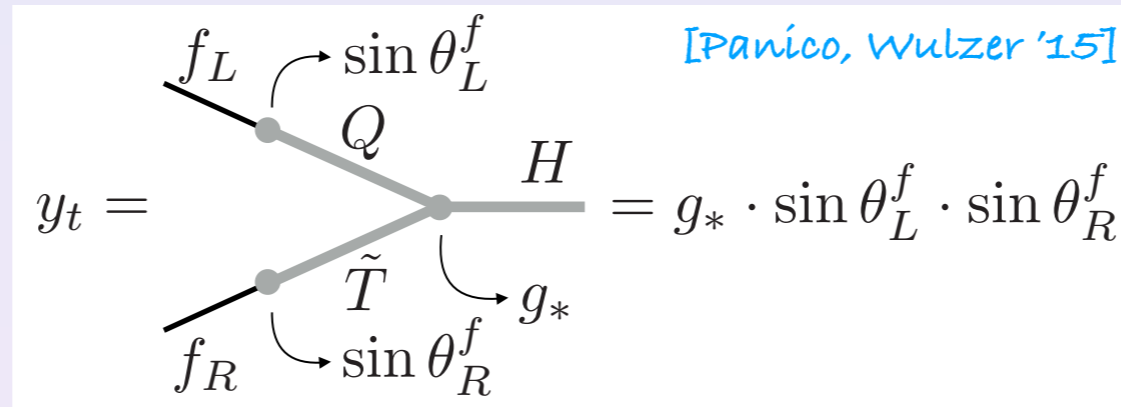
Partial compositeness

[Kaplan '91]

Elementary fermions mix with strong-interacting sector by linear couplings

$$\mathcal{L} = \lambda_L \bar{q}_L \mathcal{O}_R + \lambda_R \bar{t}_R \mathcal{O}_L$$

Mixing with top partners generate the top Yukawas



If we are only interested in the non-linearities:

Example: fermions transforming in the fundamental of $SO(5)$

$$Q_L^{2/3} = \frac{1}{\sqrt{2}} \left(d_L, -id_L, u_L, iu_L, 0 \right)^T, \quad U_R = \frac{1}{\sqrt{2}} \left(0, 0, 0, 0, \sqrt{2}u_R \right)^T$$

$$\mathcal{L}_Y = f \left[-y_u (\bar{U}_R \Sigma) (\Sigma^T Q_L^{2/3}) - y_d (\bar{D}_R \Sigma) (\Sigma^T Q_L^{-1/3}) \right] + \text{h.c.}$$

$$\longrightarrow \text{MCHM5: } \frac{g_{h\bar{f}f}}{g_{h\bar{f}f}^{SM}} = \frac{1-2\xi}{\sqrt{1-\xi}} \quad \text{MCHM10: } \frac{g_{h\bar{f}f}}{g_{h\bar{f}f}^{SM}} = \frac{1-2\xi}{\sqrt{1-\xi}} \quad \text{MCHM4: } \frac{g_{h\bar{f}f}}{g_{h\bar{f}f}^{SM}} = \sqrt{1-\xi}$$

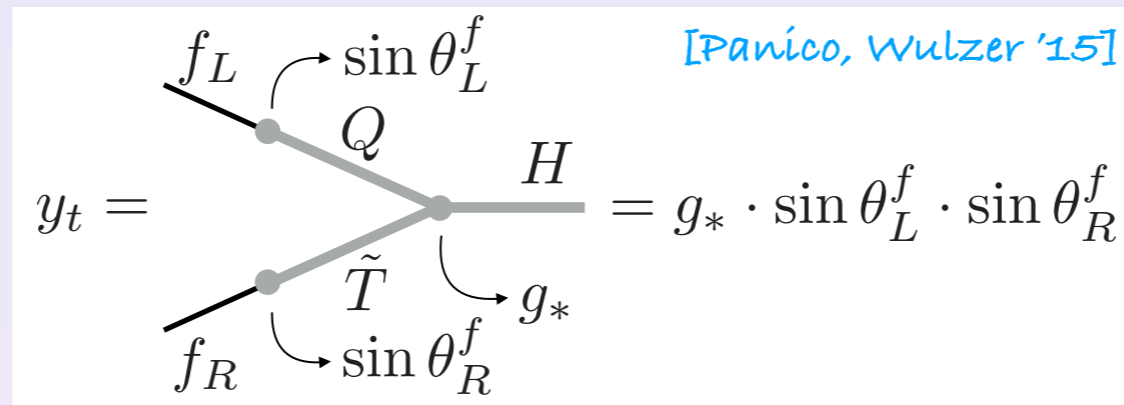
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The couplings break the global symmetry explicitly

→ Connection between top partner masses and Higgs mass

$$\frac{m_h^2}{m_t^2} \simeq \frac{N_c}{\pi^2} \frac{m_\psi^2}{f^2}$$

Low tuning = light top partner masses

[Matsedonskyi, Panico, Wulzer '12,
Marzocca, Serone, Shu '12, Pomarol, Riva '12,
Panico, Redi, Tesi, Wulzer '12,
Pappadopulo, Thamm, Torre '13]

How to test whether the Higgs is Composite?

- Higgs coupling modifications
- more resonances from the strongly-interacting sector
- non-linearities in Higgs interactions

How to test whether the Higgs is Composite?

- Higgs coupling modifications

In $SO(5)/SO(4)$ models:

$$\frac{g_{hVV}}{g_{hVV}^{SM}} = \sqrt{1 - \xi}$$

Holds true also in other models,
e.g. $SO(6)/SO(5)$
with dark matter candidate

Higgs self-couplings and couplings to fermion depend on fermion embedding

- more resonances from the strongly-interacting sector

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Probe couplings of n Higgs bosons

$$c_{tt} \neq \frac{3}{2}(c_t - 1)$$

$$c_g \neq c_{gg}$$

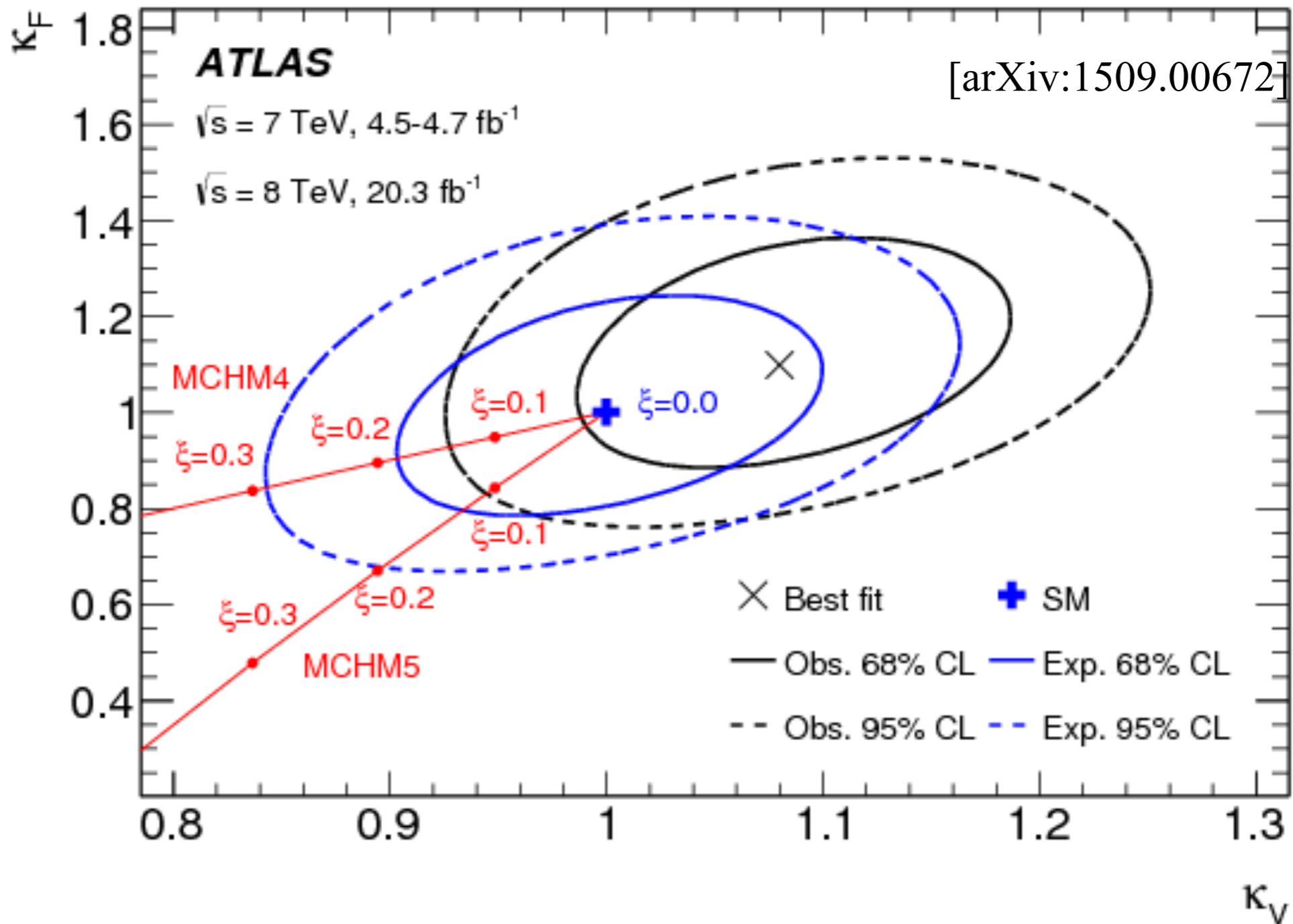
Higgs couplings

MCHM4:

$$\kappa_F = \sqrt{1 - \xi}$$

MCHM5:

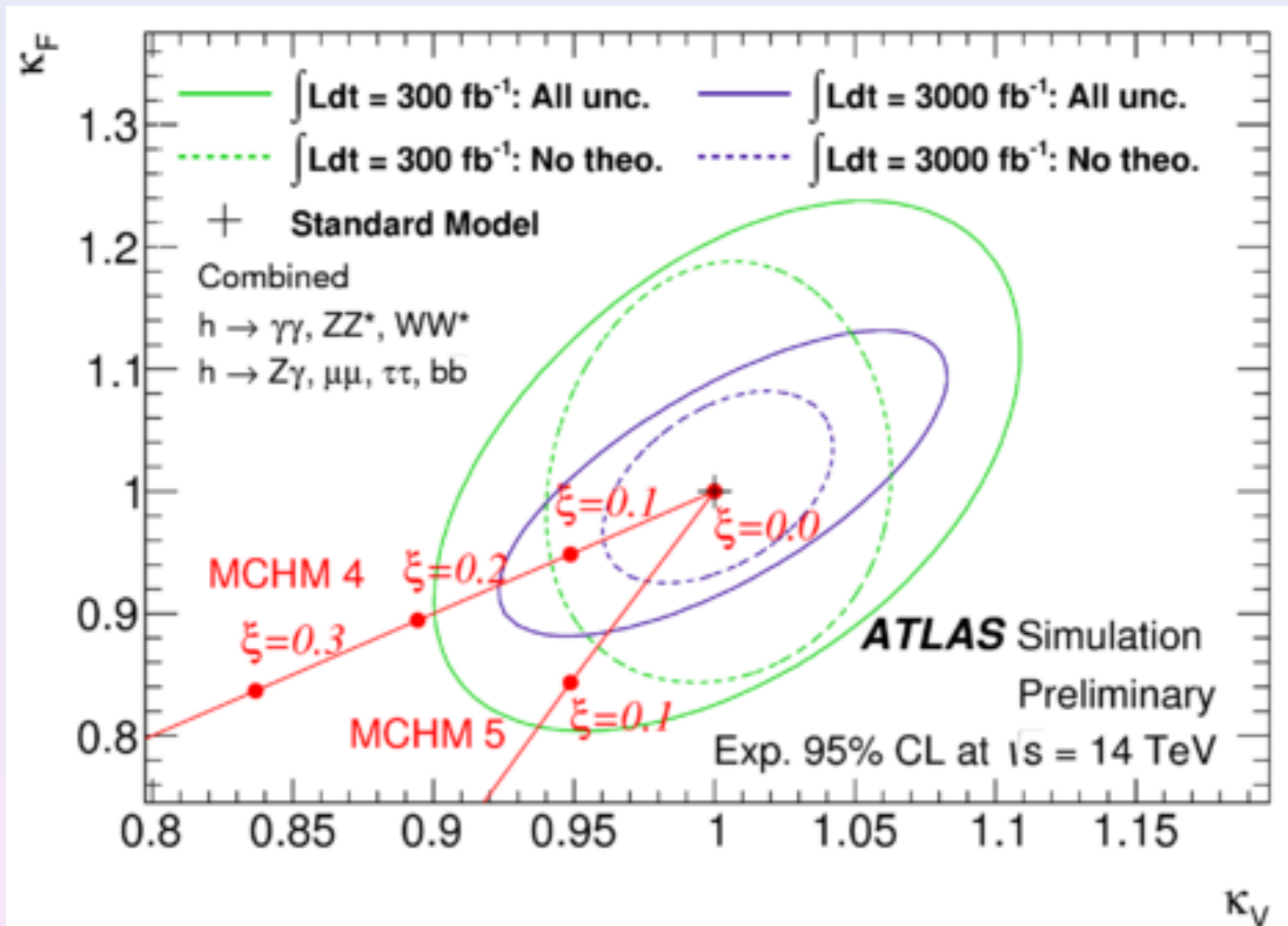
$$\kappa_F = \frac{1 - 2\xi}{\sqrt{1 - \xi}}$$



MCHM4/5: $\kappa_V = \sqrt{1 - \xi}$

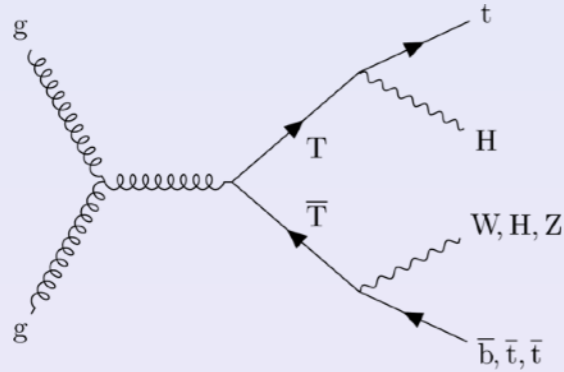
Higgs couplings

talk of M. Testa, HL-LHC workshop

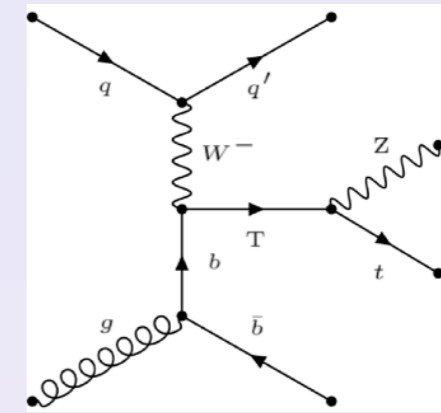


New resonances: Vector-like fermions

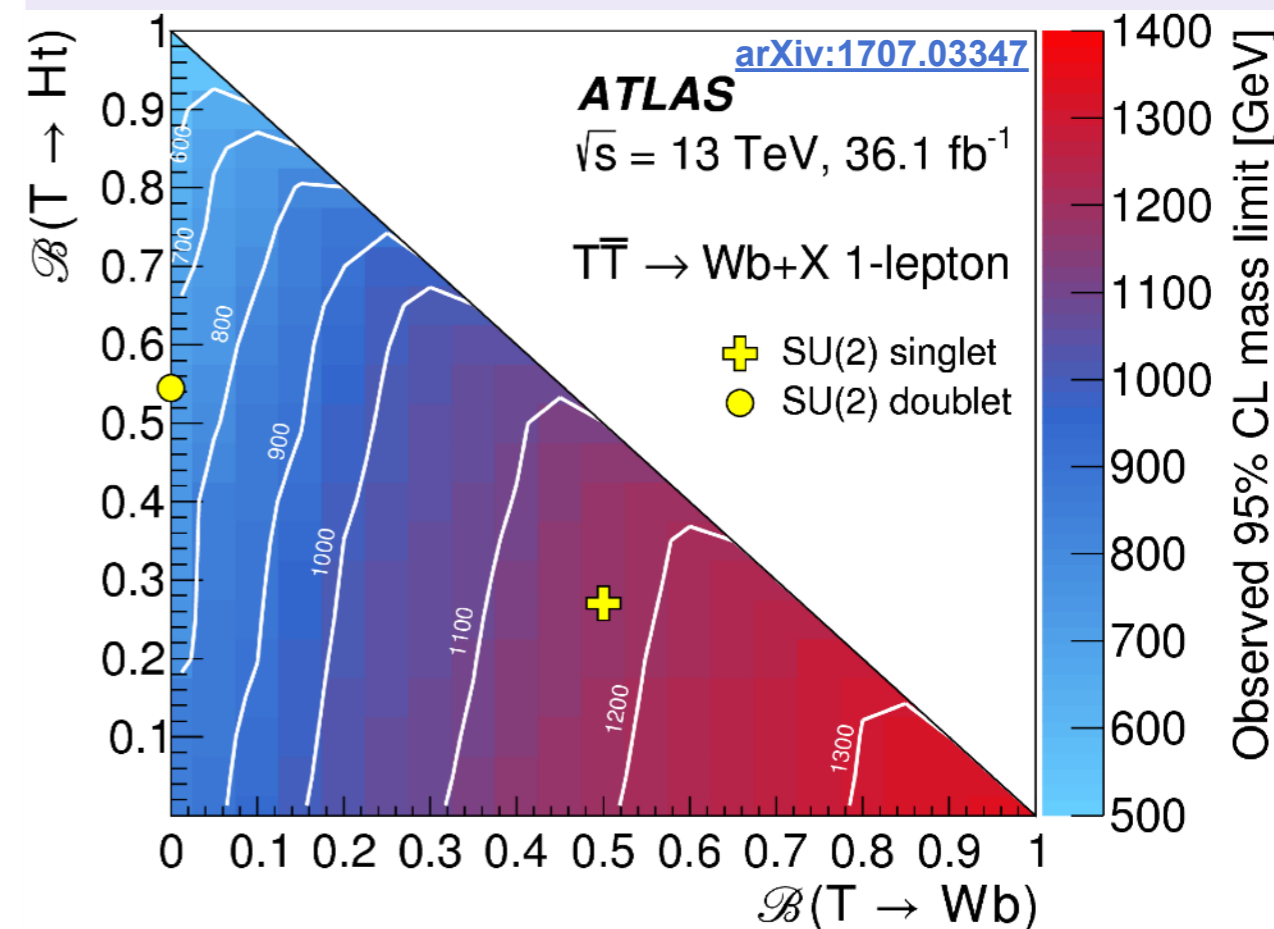
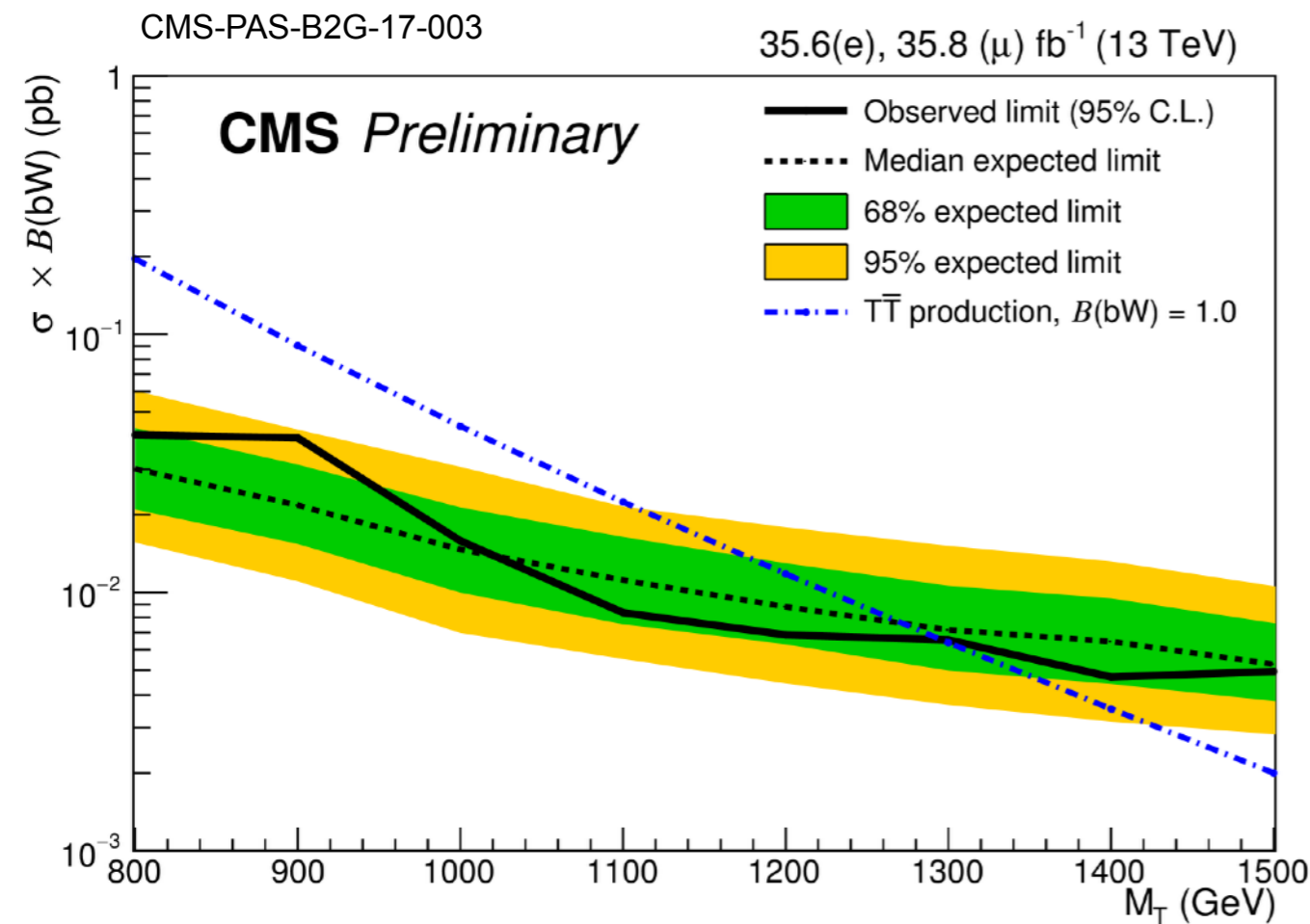
Pair production of VLQs:
pure QCD process, model independent



Single production of VLQs:
potentially higher reach, model-dependence



Limits from pair production:



Higgs non-linearities

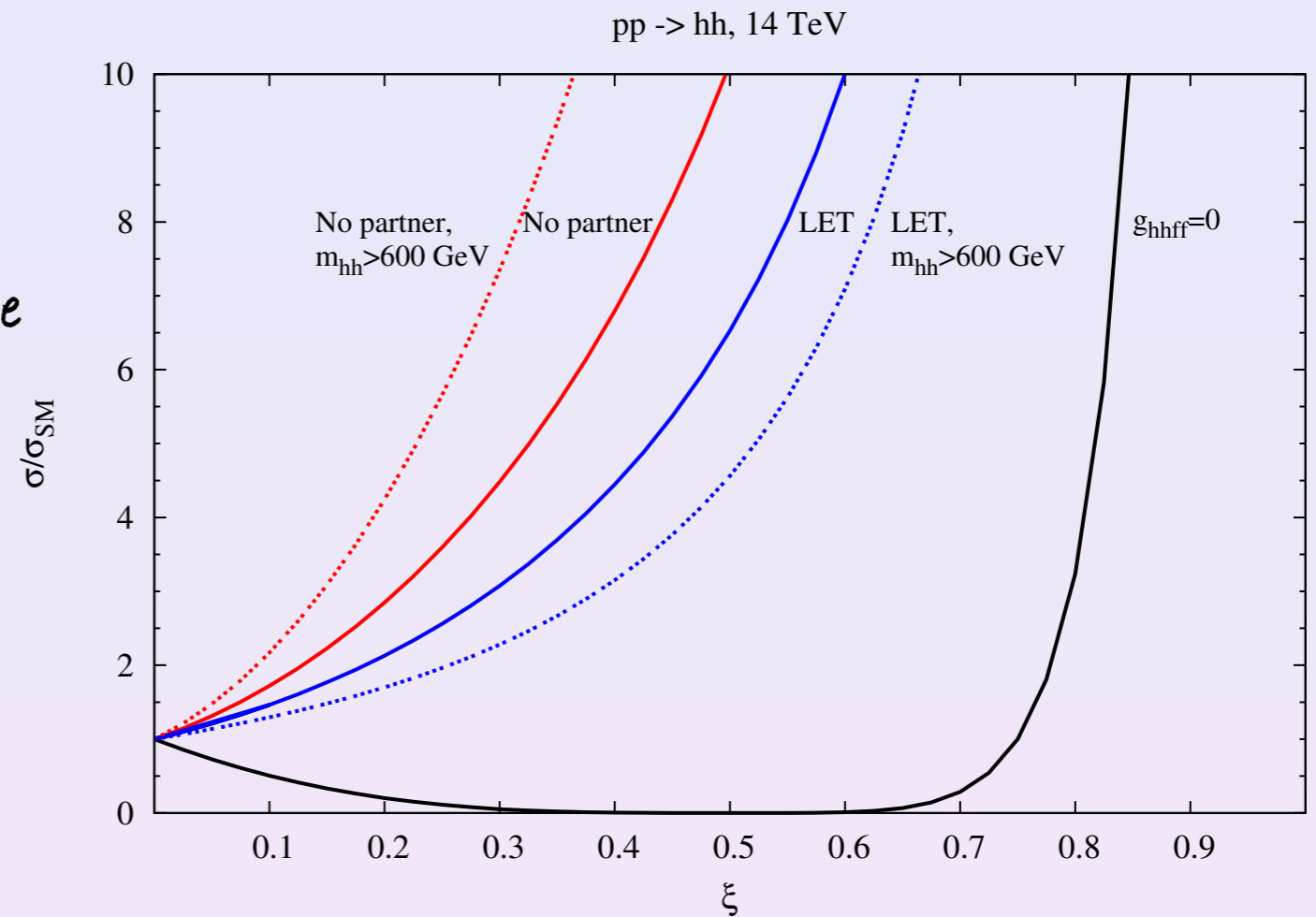
Needs to be probed in multi-Higgs interactions

[Gillioz, RG, Grojean, Mühlleitner, Salvioni '12]

Di-Higgs production via Gluon fusion

New $hhff$ coupling leads to large increase of cross section

[RG, Mühlleitner '10, Contino, Ghezzi, Moretti, Panico, Piccinini, Wulzer '12]



Higgs non-linearities

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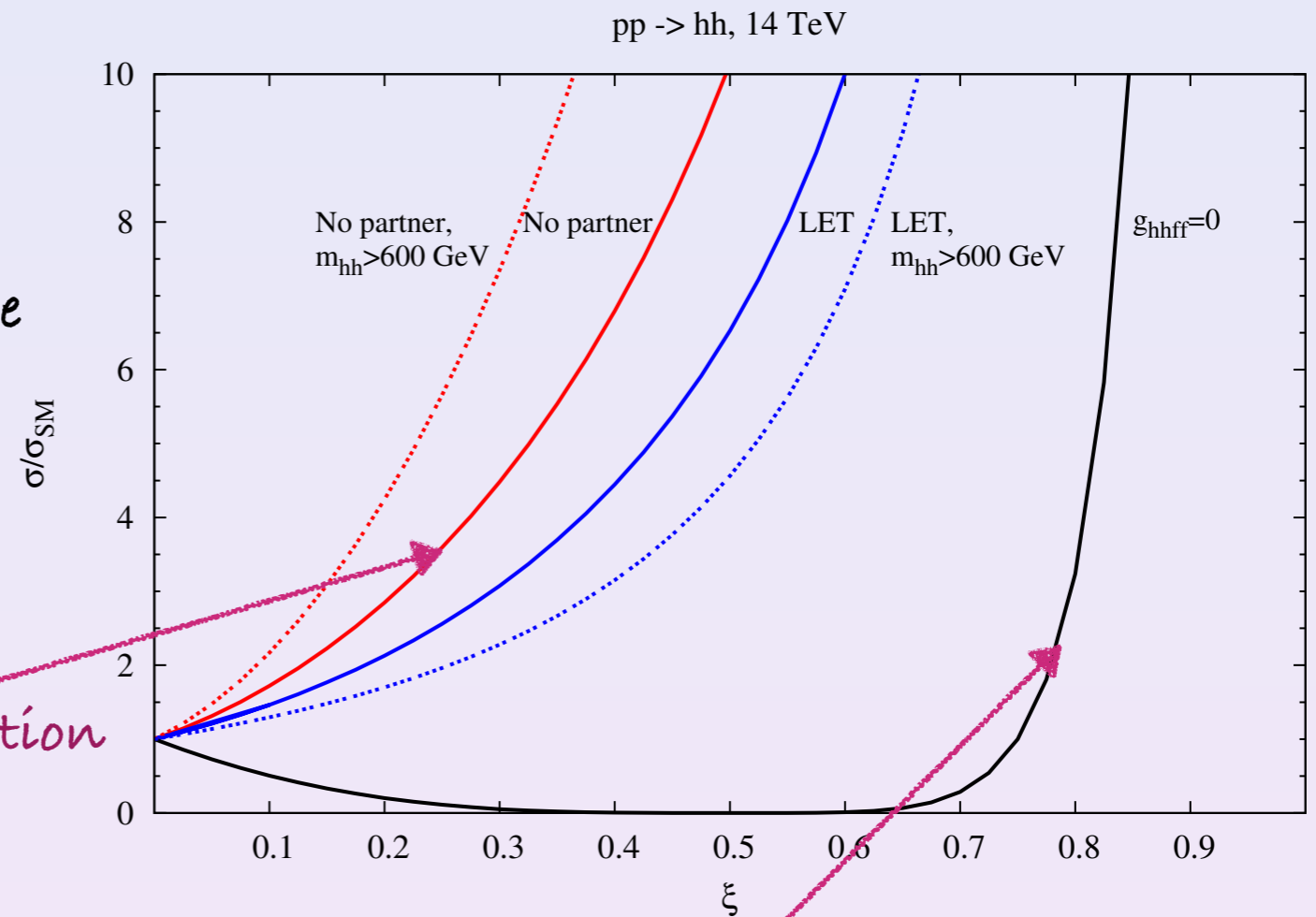
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Full LO MCHM5 $pp \rightarrow hh$ cross section



MCHM5 $pp \rightarrow hh$ cross section setting $hhff$ coupling to zero

Higgs non-linearities

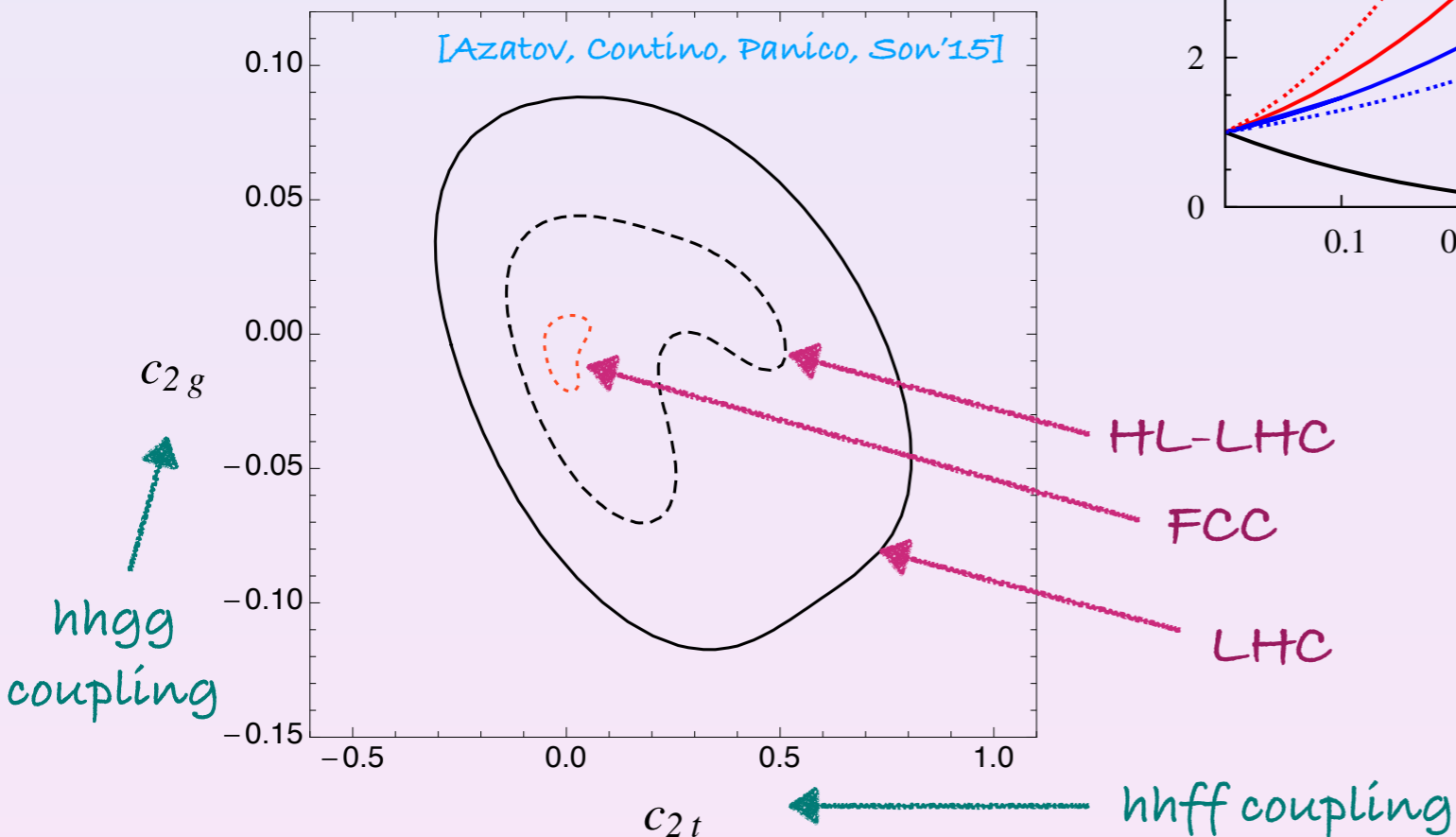
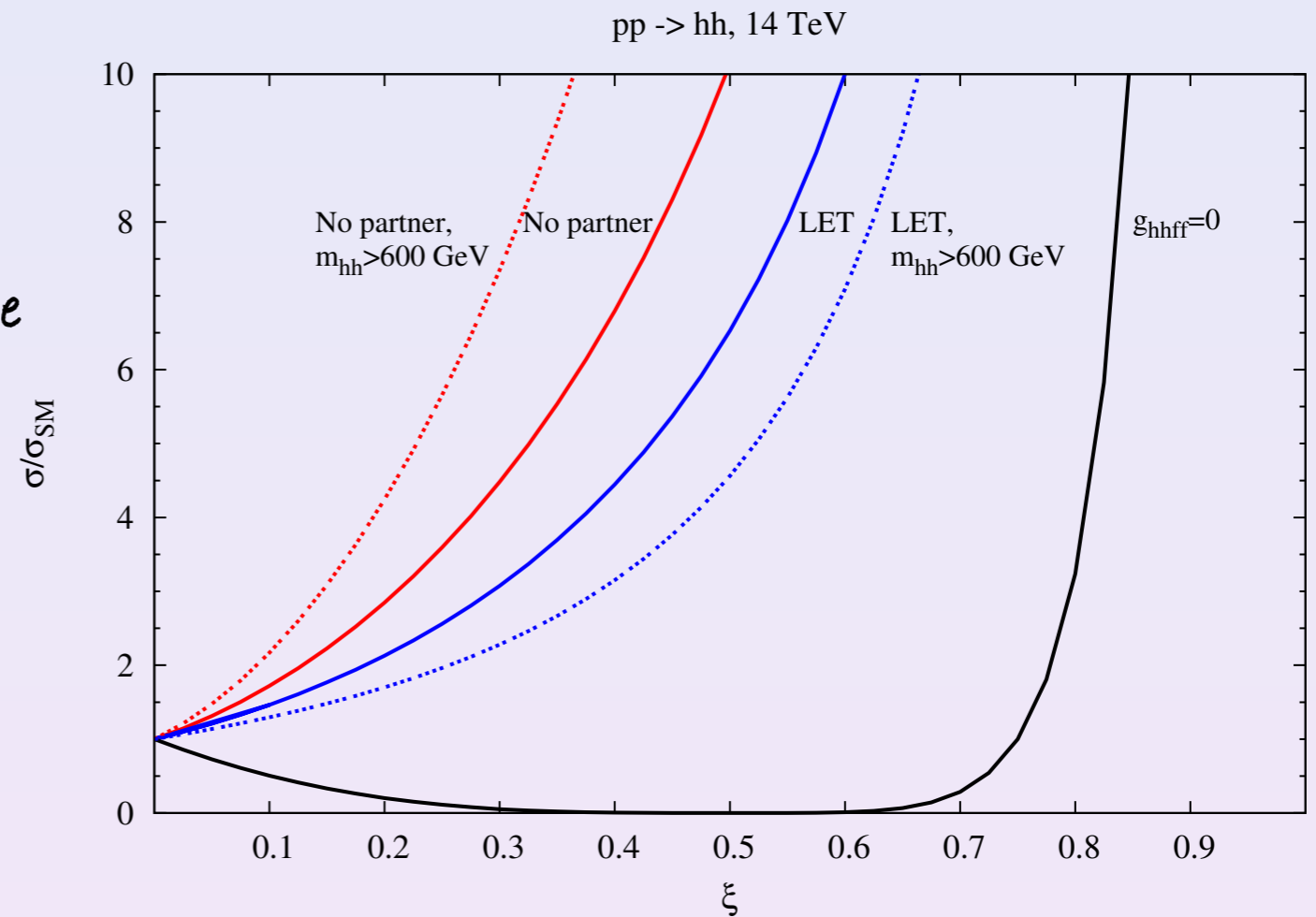
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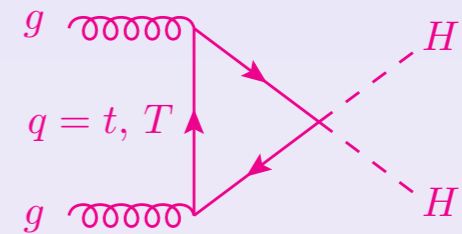
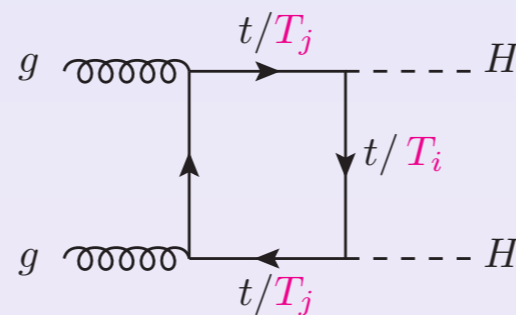
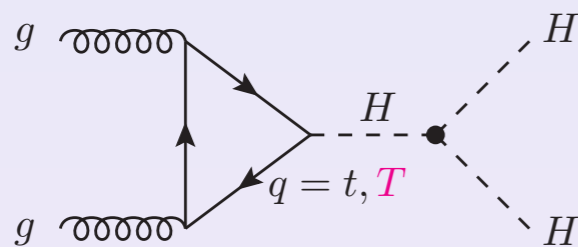


Composite Di-Higgs production

Can the increase in the cross section even be so large that New Physics will be first observed in di-Higgs production?

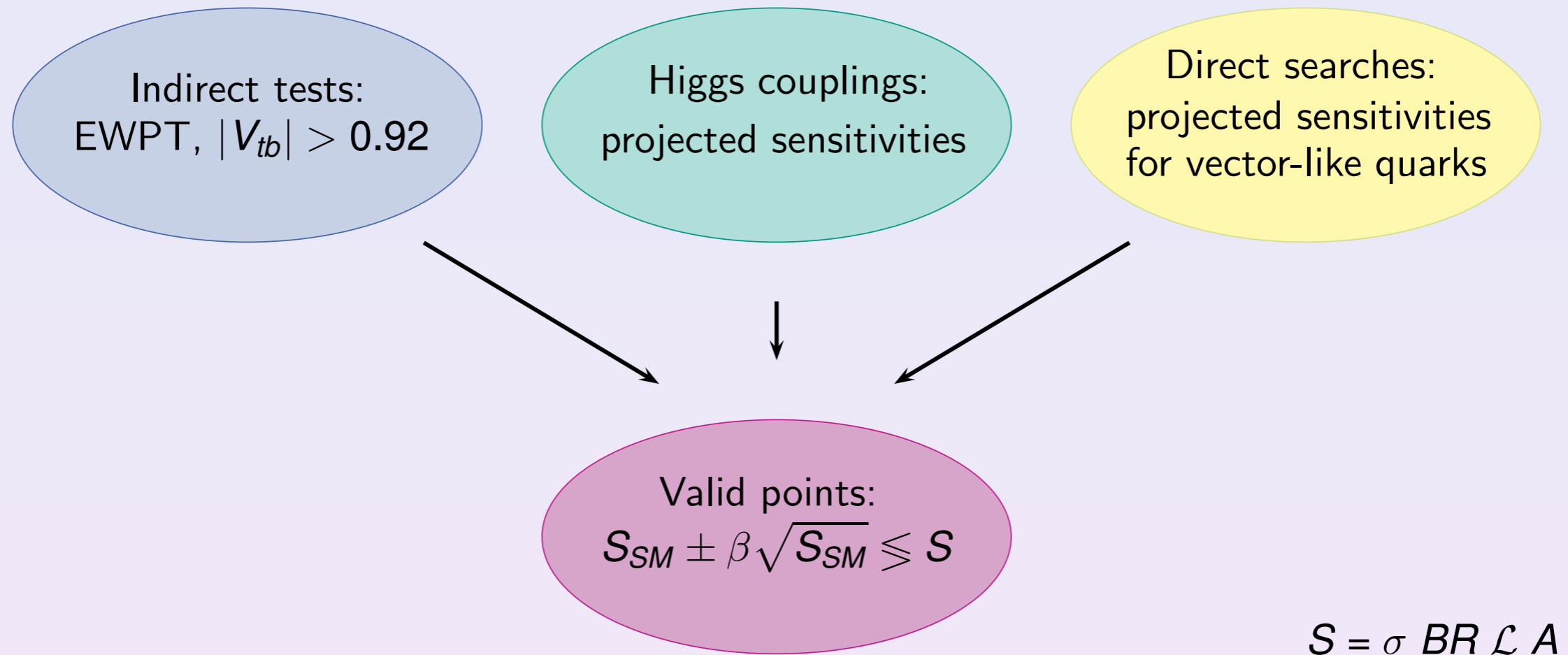
More freedom necessary than just new parameter ξ

Take a model with top and bottom partners



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for $b\bar{b}\gamma\gamma$ and $b\bar{b}\tau\bar{\tau}$ final state [Baglio, Djouadi, RG, Mühlleitner, Quevillon, Spira '12]

EWPTs from [Gillioz, RG, Kapuvari, Mühlleitner '14]

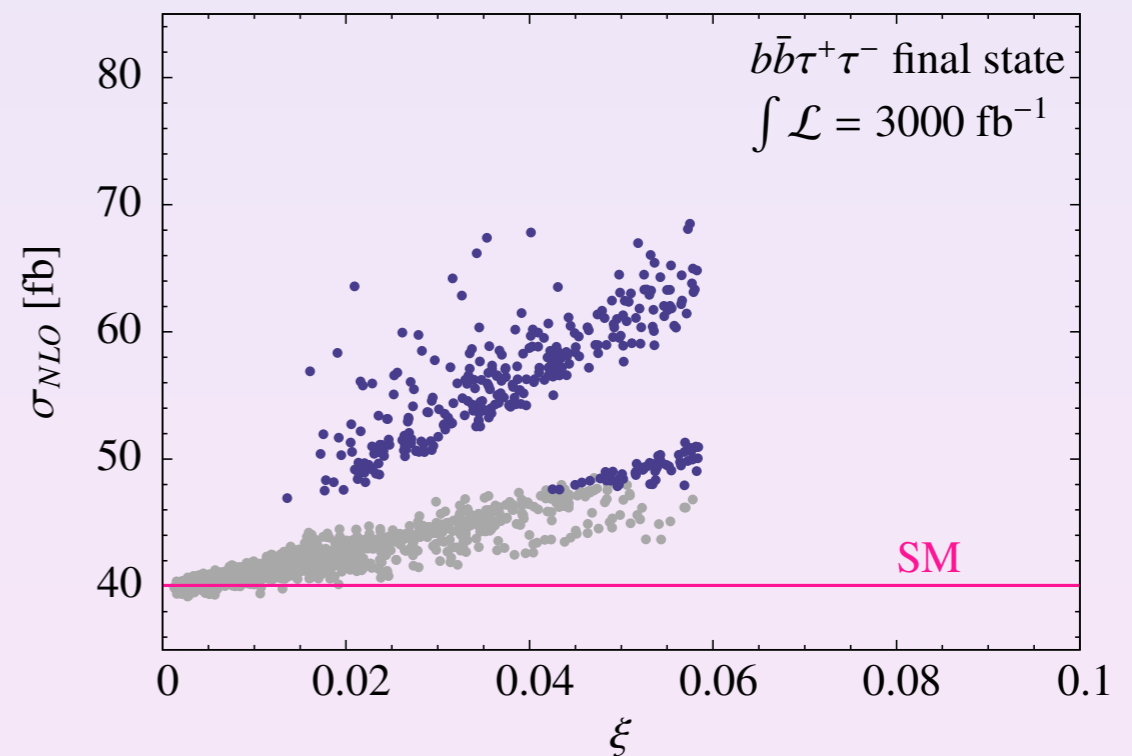
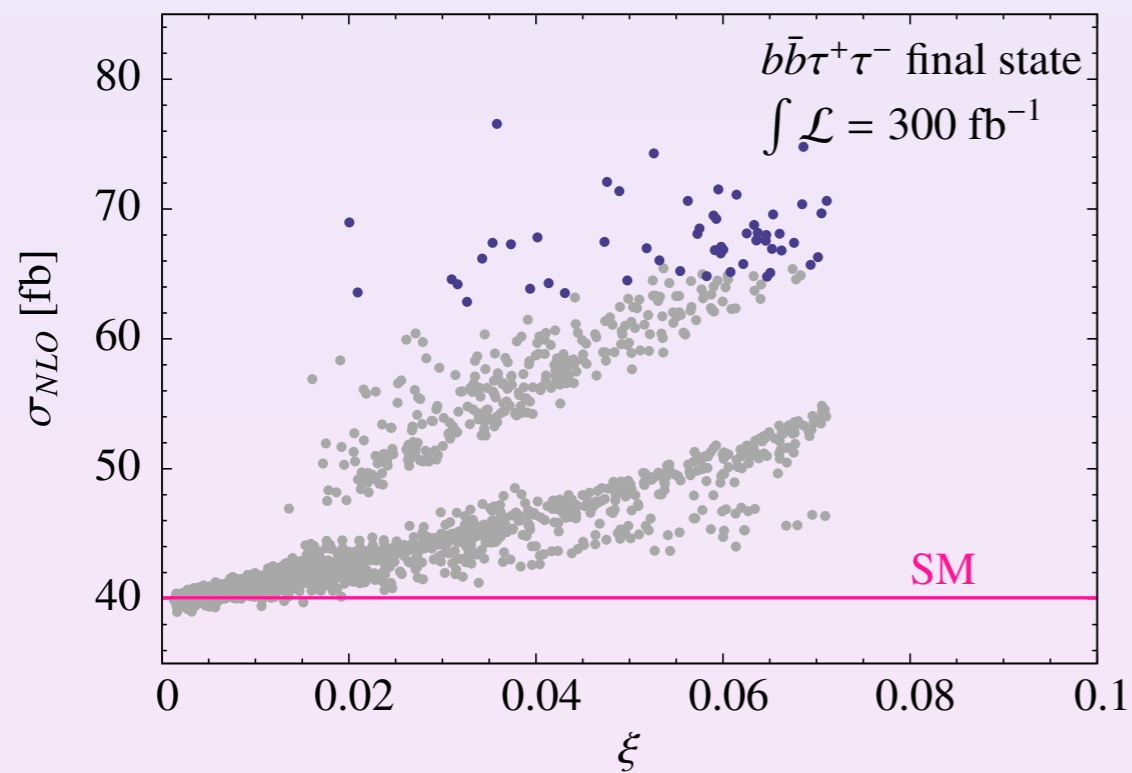
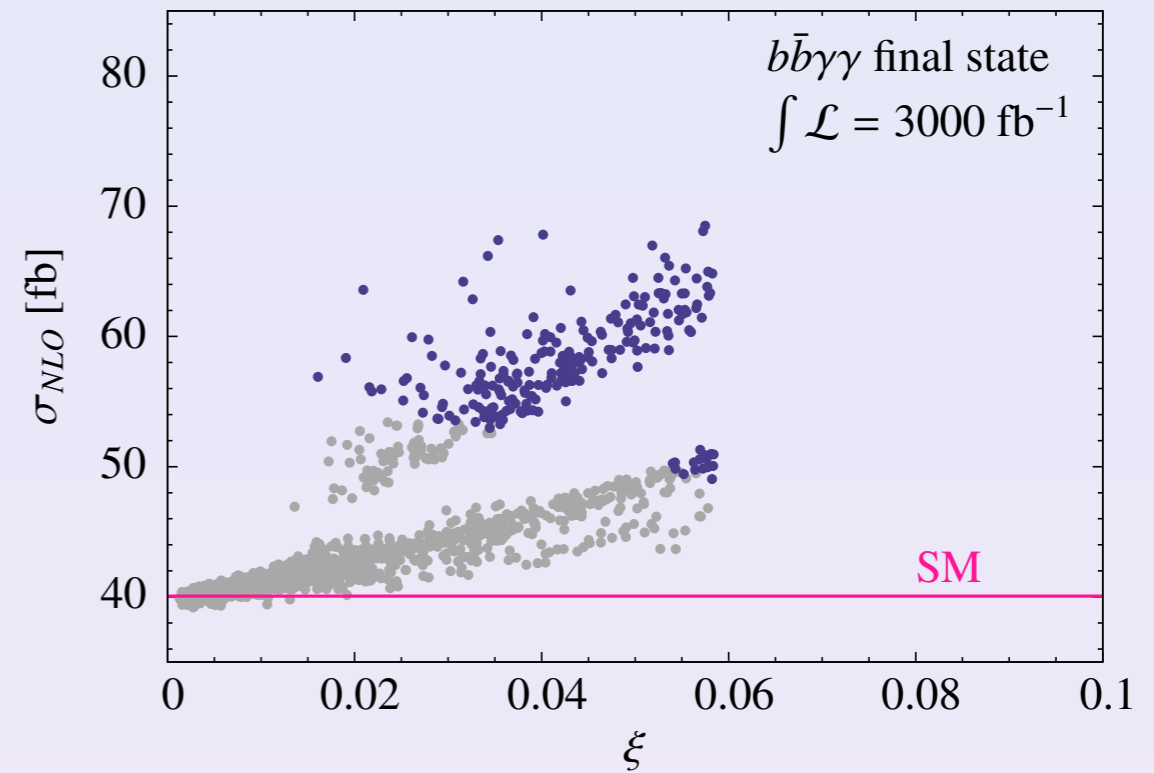
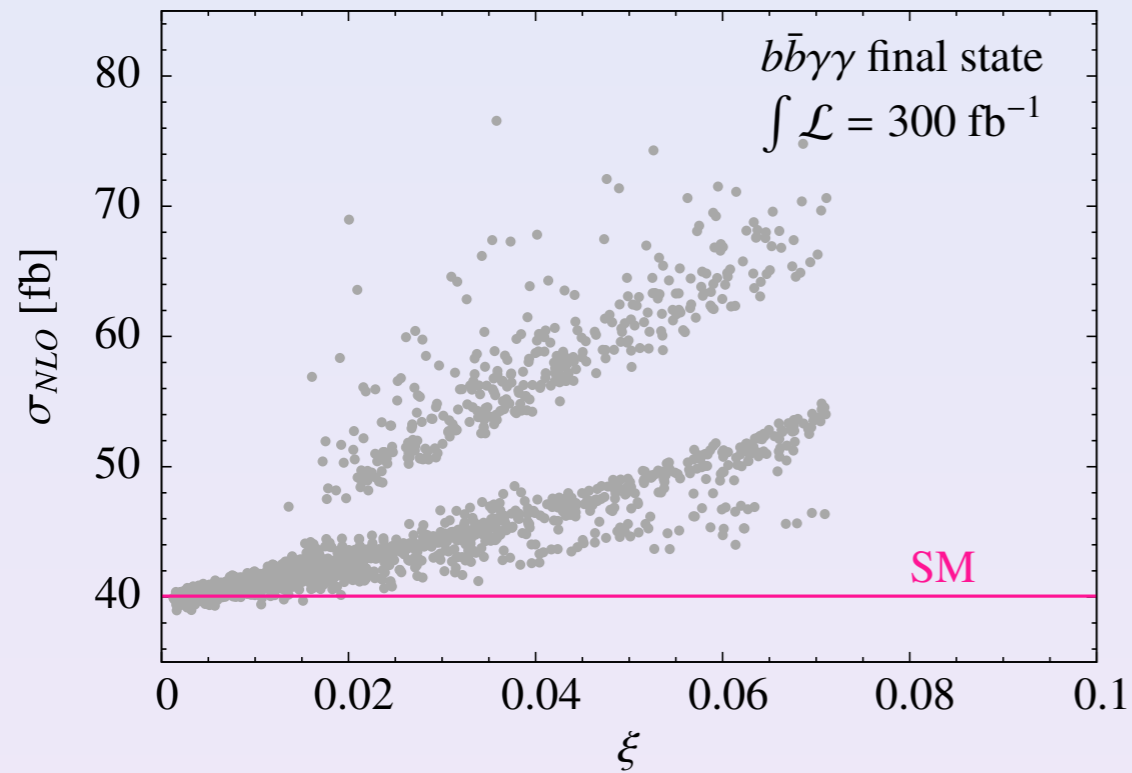
Projected sensitivities for Higgs couplings from

[Englert, Freitas, Mühlleitner, Plehn, Rauch, Spira, Walz '14]

Projected sensitivities for direct searches of VLQs $m \lesssim 1.5$ TeV

Composite Di-Higgs production

[RG, Mühlleitner, Spira '16]



Non-minimal Composite Higgs Models

[Mrazek et al '11]

G	H	N_G	NGBs rep. $[H] = \text{rep.}[SU(2) \times SU(2)]$
SO(5)	SO(4)	4	$\mathbf{4} = (\mathbf{2}, \mathbf{2})$
SO(6)	SO(5)	5	$\mathbf{5} = (\mathbf{1}, \mathbf{1}) + (\mathbf{2}, \mathbf{2})$ [Gripaios et al '09; Frigerio et al '12; Marzocca, Urbano '14]
SO(6)	SO(4) \times SO(2)	8	$\mathbf{4}_{+2} + \bar{\mathbf{4}}_{-2} = 2 \times (\mathbf{2}, \mathbf{2})$ [Mrazek et al '11, De Curtis et al '16]
SO(7)	SO(6)	6	$\mathbf{6} = 2 \times (\mathbf{1}, \mathbf{1}) + (\mathbf{2}, \mathbf{2})$ [Balkin et al '17]
SO(7)	G_2	7	$\mathbf{7} = (\mathbf{1}, \mathbf{3}) + (\mathbf{2}, \mathbf{2})$ [Chala '12; Ballesteros et al '16]
SO(7)	SO(5) \times SO(2)	10	$\mathbf{10}_0 = (\mathbf{3}, \mathbf{1}) + (\mathbf{1}, \mathbf{3}) + (\mathbf{2}, \mathbf{2})$
SO(7)	$[SO(3)]^3$	12	$(\mathbf{2}, \mathbf{2}, \mathbf{3}) = 3 \times (\mathbf{2}, \mathbf{2})$
Sp(6)	Sp(4) \times SU(2)	8	$(\mathbf{4}, \mathbf{2}) = 2 \times (\mathbf{2}, \mathbf{2}), (\mathbf{2}, \mathbf{2}) + 2 \times (\mathbf{2}, \mathbf{1})$
SU(5)	SU(4) \times U(1)	8	$\mathbf{4}_{-5} + \bar{\mathbf{4}}_{+5} = 2 \times (\mathbf{2}, \mathbf{2})$
SU(5)	SO(5)	14	$\mathbf{14} = (\mathbf{3}, \mathbf{3}) + (\mathbf{2}, \mathbf{2}) + (\mathbf{1}, \mathbf{1})$

Larger coset space



extended Higgs sector

Non-minimal Composite Higgs Models

if new scalar stable



possible dark matter candidate

i.e. if there is a Z_2 or $U(1)$
symmetry

General parameterisation:

$$L = |D_\mu H|^2 \left[1 - a_1 \frac{S^2}{f^2} \right] + \frac{a_2}{f^2} \partial_\mu |H|^2 (S \partial_\mu S) + \frac{1}{2} (\partial_\mu S)^2 \left[1 - 2a_3 \frac{|H|^2}{f^2} \right] \\ - m_\rho^2 f^2 \frac{N_c y_t^2}{(4\pi)^2} \left[-\alpha \frac{|H|^2}{f^2} + \beta \frac{|H|^4}{f^4} + \gamma \frac{S^2}{f^2} + \delta \frac{S^2 |H|^2}{f^4} \right] + \left[i\epsilon \frac{y_t}{f^2} S^2 \overline{q_L} H t_R + \text{h.c.} \right] + \dots$$

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Annihilation cross section dominated by $H^2 S^2$ interactions

$$a = a_1 = a_3 = a_3$$

Relic density (if $\gamma > \delta/10$):

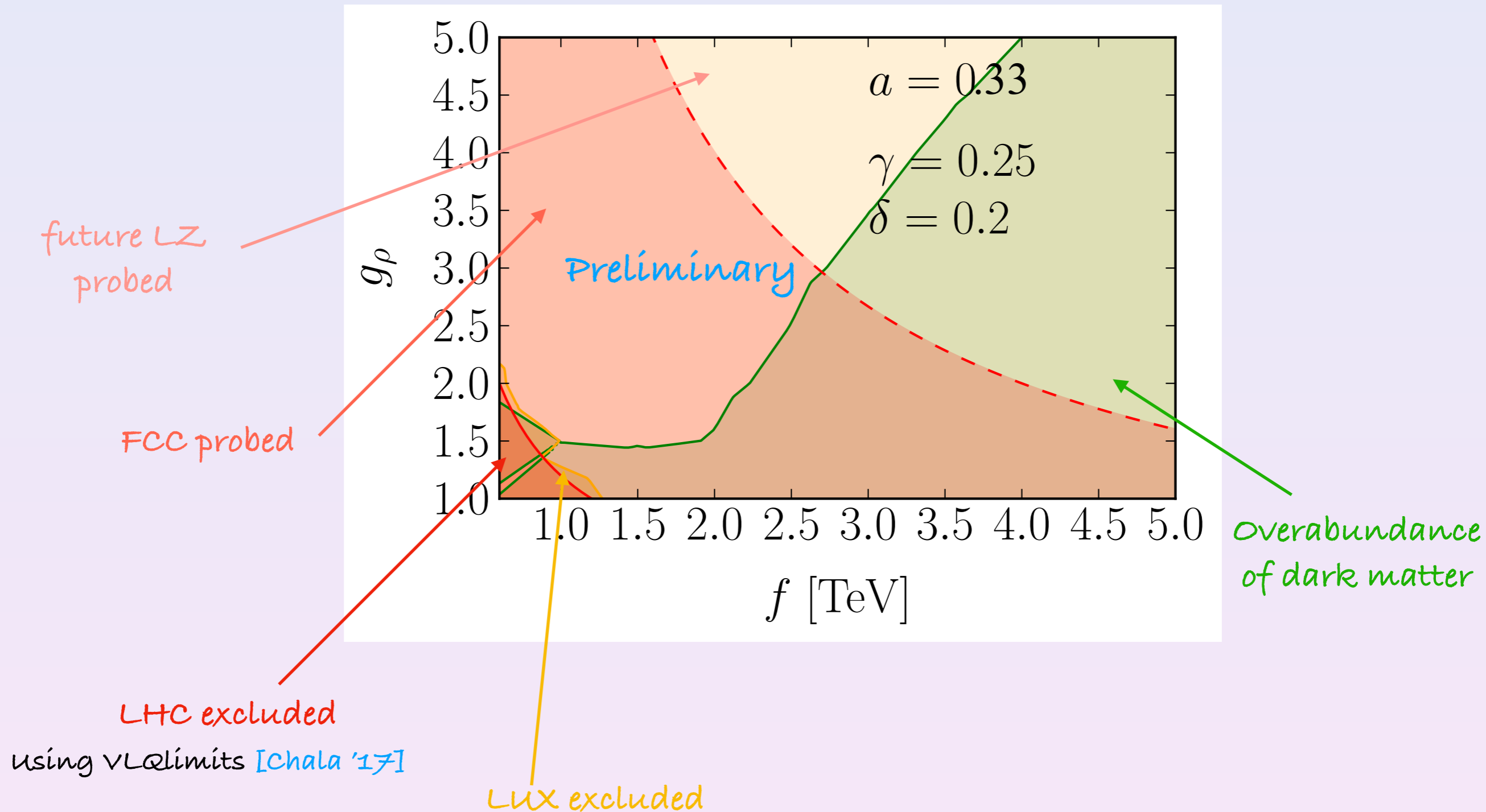
$$\Omega h^2 \propto \frac{f^2}{g_\rho^2} \quad \text{with} \quad m_\rho = g_\rho f$$

Main difference to non-composite case: derivative interactions

Non-minimal Composite Higgs Models

$SO(6)/SO(5)$ model

[Chala, RG, Spannowsky preliminary]



Conclusion

- Composite Higgs Models viable alternative to Standard Model
- Testable at LHC by
 - Higgs coupling measurements
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 - Non-linearities
- Composite Dark matter scenarios very predictive
 - Hints to values of f above tuning expectation
 - future colliders and dark matter direct detection experiments will shed light on such scenarios

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Thanks for your attention!