
The Higgs Portal

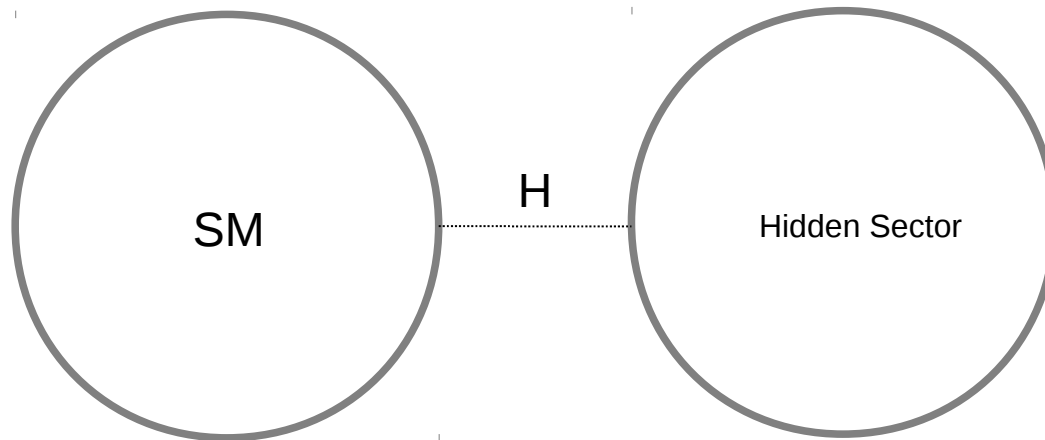
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-
- the Higgs and the hidden sector
 - the Higgs and dark matter
 - the Higgs and inflation
-

The Higgs and the hidden sector



Special role of the Higgs :

Silveira, Zee '85
Veltman, Yndurain '89
...

$|H|^2$ = the only gauge and Lorentz-inv. dim-2 operator

$$L = a |H|^2 S^2 + b |H|^2 S$$

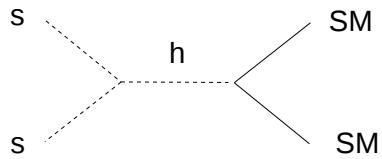
(S = "hidden" scalar)

b=0 (S has hidden charge):

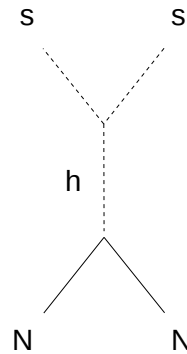
$$L = a |H|^2 S^2$$

"S" is stable and couples weakly to SM \Rightarrow **DARK MATTER** (?)

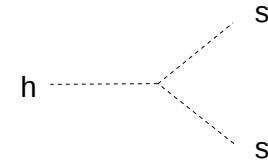
Dark matter:



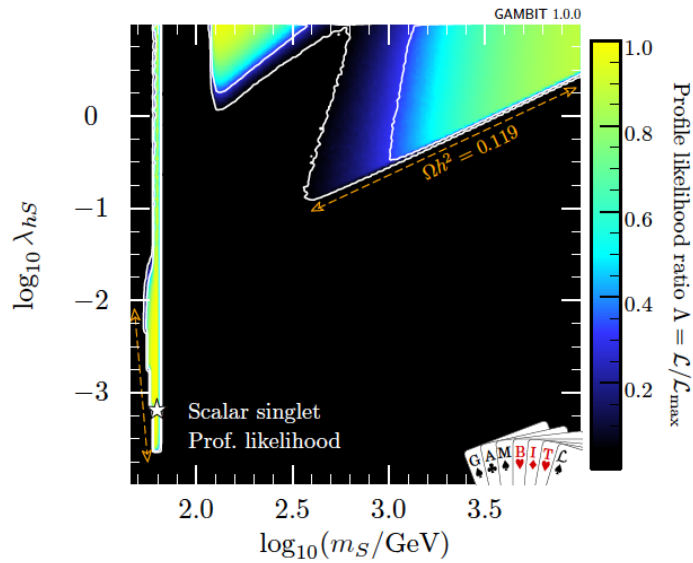
DM annihilation



DM direct detection



Higgs decay



GAMBIT collaboration
1705.07931

white contour = 2 σ bound

The Higgs portal and DD cancellation

Add a complex scalar S , require softly broken U(1) symmetry:

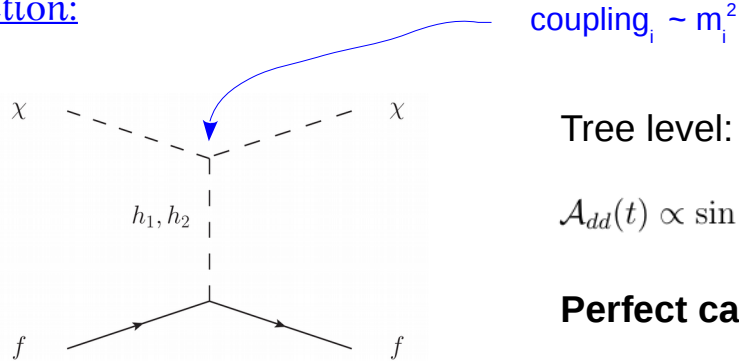
$$V = V_0 + V_{\text{soft}} ,$$

$$V_0 = -\frac{\mu_H^2}{2} |H|^2 - \frac{\mu_S^2}{2} |S|^2 + \frac{\lambda_H}{2} |H|^4 + \lambda_{HS} |H|^2 |S|^2 + \frac{\lambda_S}{2} |S|^4 ,$$

$$V_{\text{soft}} = -\frac{\mu_S'^2}{4} S^2 + \text{h.c.}$$

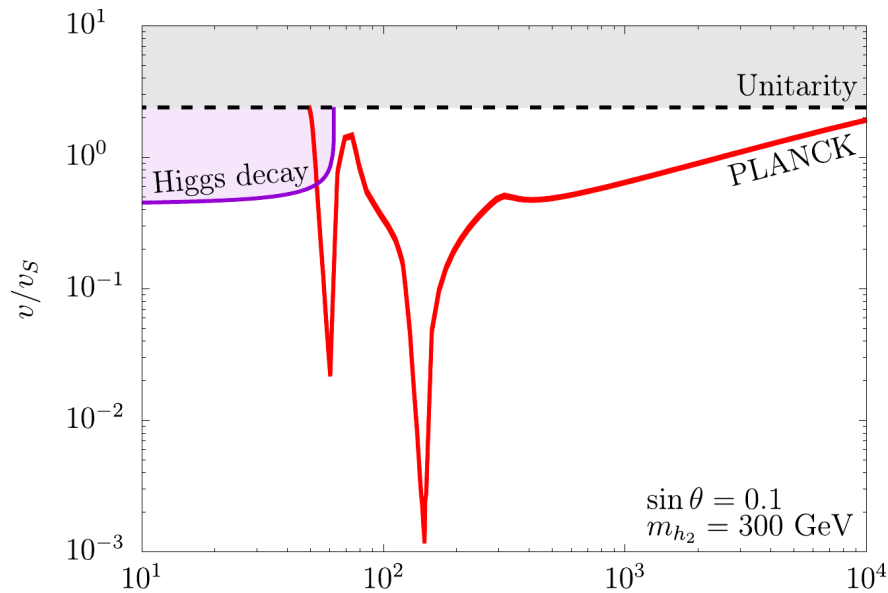
All parameters are real $\Rightarrow \langle S \rangle = \text{real}$, $S \rightarrow S^*$ symmetry

Im S = Dark Matter

Direct detection:

Tree level:

$$\mathcal{A}_{dd}(t) \propto \sin \theta \cos \theta \left(\frac{m_2^2}{t - m_2^2} - \frac{m_1^2}{t - m_1^2} \right) \rightarrow 0$$

Perfect cancellation for any parameter choice !

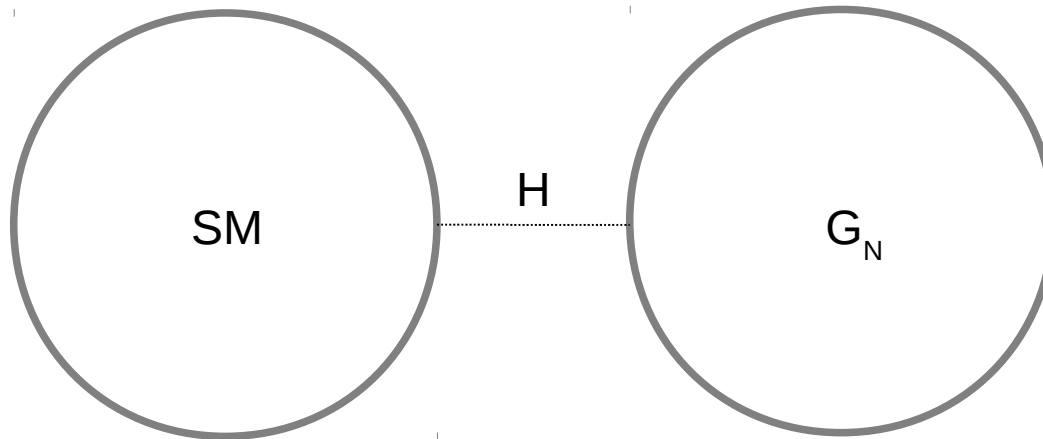
Direct detection = loop-suppressed

Annihilation = unsuppressed

**Excellent WIMP**

(from 60 GeV to 10 TeV)

The Higgs and vector dark matter



$$V \sim \bar{H} H \bar{S} S$$



H-S mixing



h couples to G_N

Lie groups possess discrete symmetries



gauge fields as dark matter

E.g. $U(1) : A_\mu \rightarrow -A_\mu$

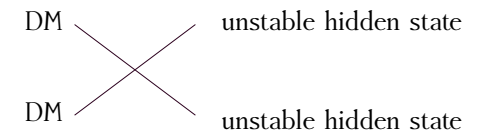


Minimal G_N breaking implies:

- Vector DM
- Multicomponent DM
- "Secluded" DM

$$A_\mu$$

$$A_\mu, \chi$$



(à la Pospelov et al. '07)

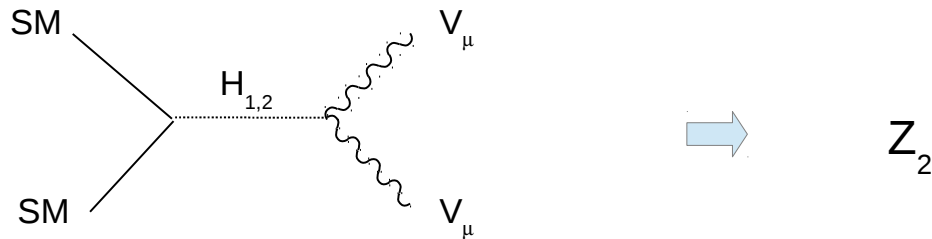
Higgs mechanism in the hidden sector :

$$L = -1/4 F_{\mu\nu} F^{\mu\nu} + D_{\mu} S^* D^{\mu} S - V(S) + \lambda/4 \bar{H} H S^* S$$

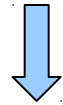
$$S \longrightarrow \text{VEV}$$



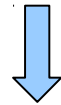
SM couplings:



gauge invariance (+ minimal field content)

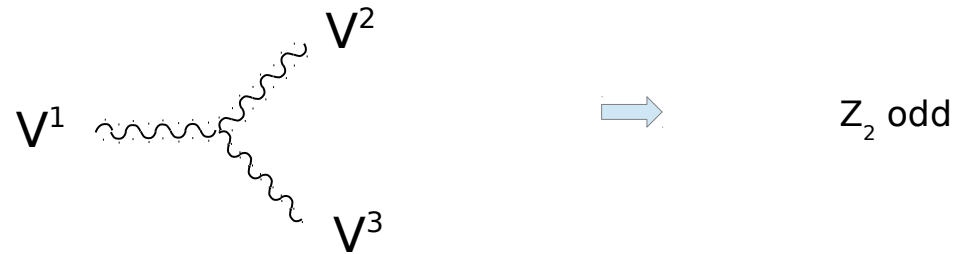


Z_2



gauge fields are natural DM candidates

Non-abelian case:



But there are 2 Z_2 's:

gauge transform

charge conjugation

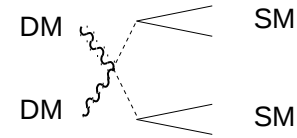
$$\begin{aligned} V^{1,2} &\rightarrow -V^{1,2} & , & & V^3 &\rightarrow V^3 \\ V^{1,3} &\rightarrow -V^{1,3} & , & & V^2 &\rightarrow V^2 \end{aligned}$$

$\Rightarrow V^a = \text{stable}$

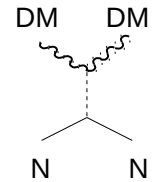
Advantage over the simplest Higgs portal DM:

there exists another hidden sector state which

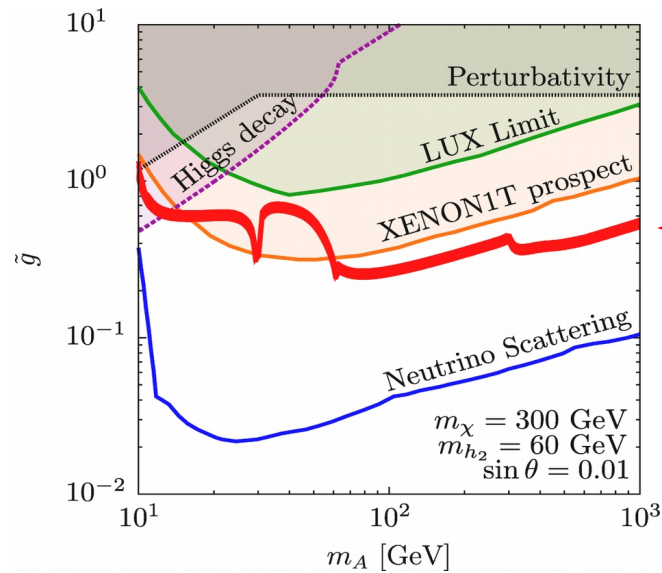
- **can be lighter than DM**
- **is unstable**



$\sim \cos \theta$



$\sim \sin \theta$



correct relic density

- DM annihilation efficient
- Direct detection suppressed

à la Pospelov et al.'07



Higgs portal DM = viable WIMP

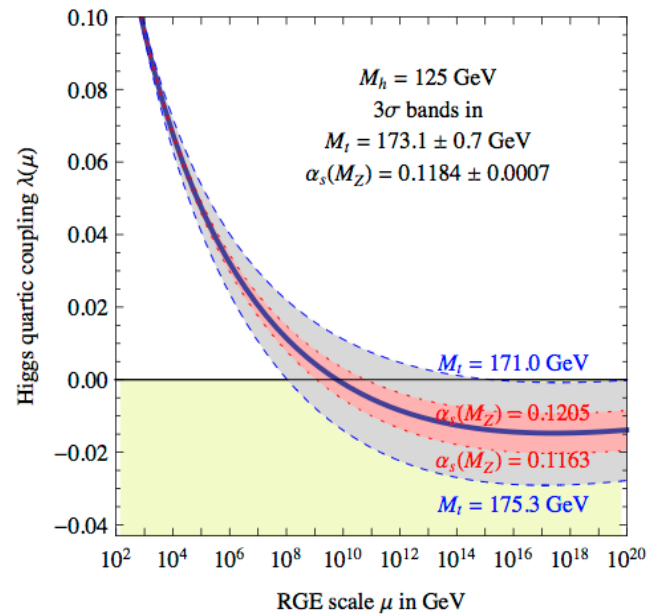
(especially if there's more than one state in the hidden sector)

The Higgs and inflation

Buttazzo et al.'13

SM stability bound:

$$m_h > (129.6 \pm 1.5) \text{ GeV}$$

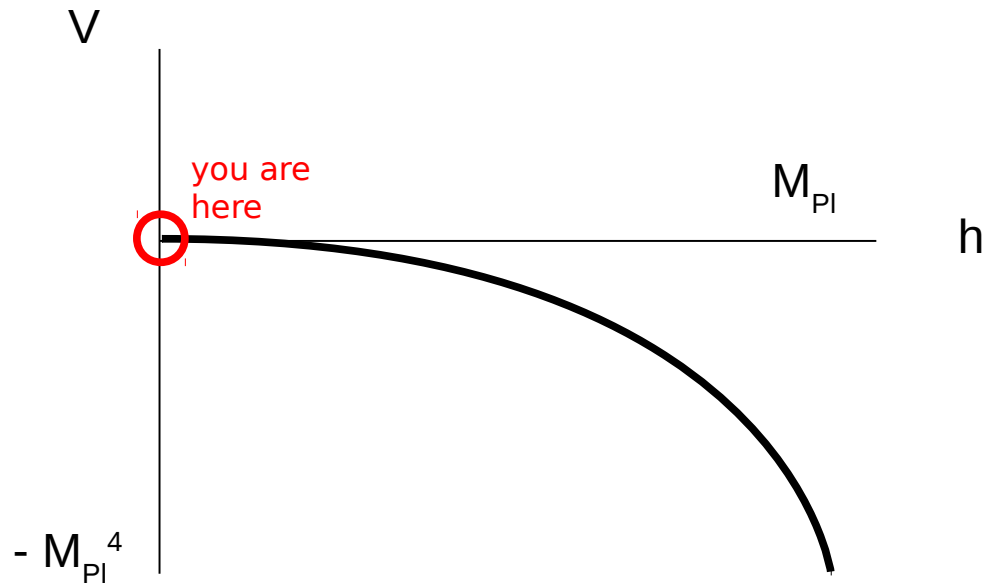


(not settled : Alekhin et al. '12
Bezrukov et al. '12)

$$h \gg \Lambda \sim 10^{10} \text{ GeV}$$



$$V \sim \frac{1}{4} \lambda(h) h^4, \quad \lambda(h) < 0$$



$$\Lambda = 10^{-8} M_{\text{Pl}}$$

,

$$\text{barrier} = 10^{-32} M_{\text{Pl}}^4$$

Problems :

- how did the Universe end up at $h \sim 0$?
- why did it stay there during inflation ?

(large fluctuations: [Espinosa, Guidice, Riotto'07](#))

Solutions :

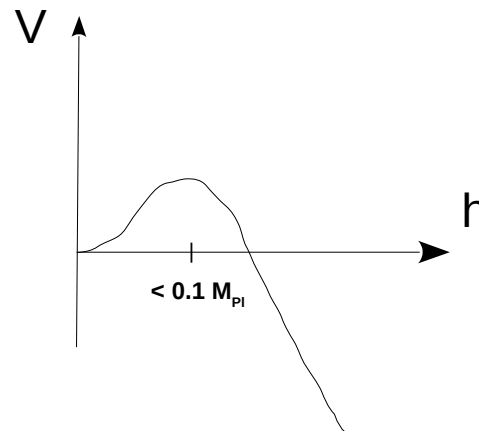
- modify the Higgs potential during inflation
 - just modify the Higgs potential
-

Minimal solution:

Higgs-inflaton coupling:

$$\Delta V = \frac{1}{2} \lambda_{h\phi} h^2 \phi^2$$

("Higgs portal" coupling)

 $\Delta V + V_{SM} :$ 

For all initial values of h up to $0.1 M_{Pl}$, the h -potential is convex
(higher h -values \rightarrow Planckian density)

Constraints:

$$\left\{ \begin{array}{ll} - \text{should not affect } V_{\text{infl}} & \rightarrow \lambda_{h\phi} < 10^{-6} \\ - \Delta V + V_{\text{SM}} > 0 & \rightarrow \phi_0 > 20 M_{\text{Pl}} \end{array} \right.$$

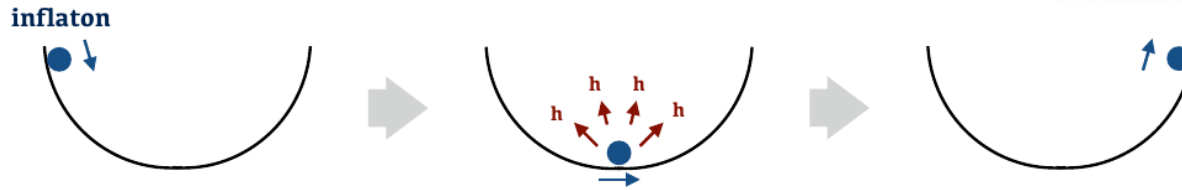
Large effective mass term $\sim \lambda_{h\phi} \phi^2 \Rightarrow$ $h(t) \sim h(0) \exp(-3/2 Ht)$

Higgs field is driven to zero during inflation !

Higgs vacuum destabilization through preheating

$$V(\phi) = \frac{1}{2}m^2\phi^2$$

$$\phi \simeq \Phi \cos mt \quad \text{with} \quad \Phi \sim \Phi_0 a^{-3/2}$$



Kofman, Linde, Starobinsky '98

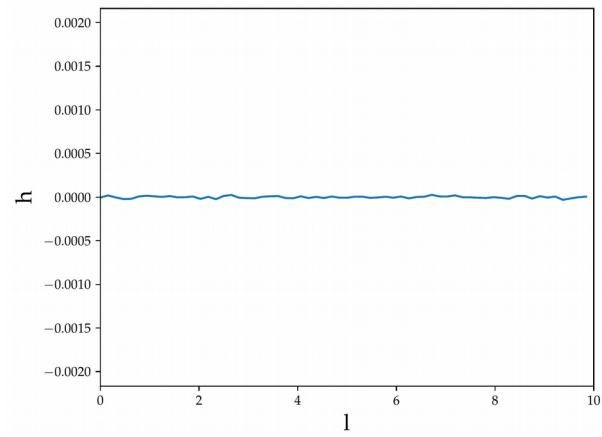
$$\Delta V = \frac{1}{2} \lambda_{h\phi} h^2 \phi^2 \quad \Rightarrow \quad \text{parametric resonance}$$

$$\langle h^2 \rangle \propto \text{Number of Higgs quanta}$$

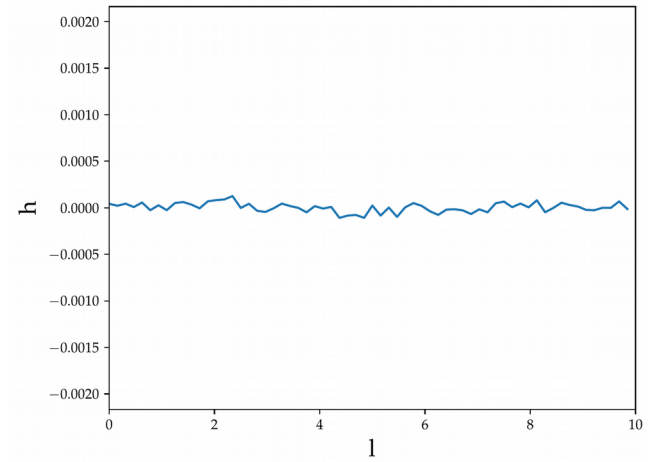


Lattice results:

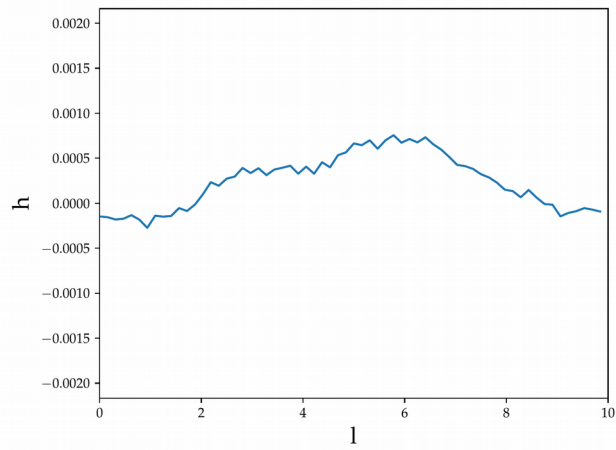
t=0



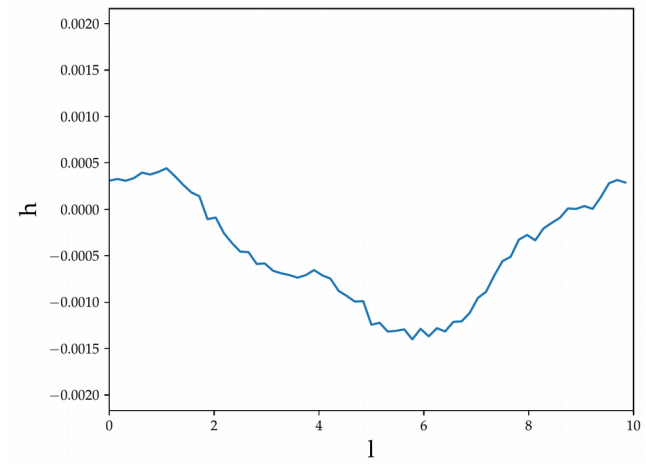
t=15



t=27



t=30



General case: including all relevant operators

$$\begin{aligned}
 -\mathcal{L}_{hR} &= \xi H^\dagger H \hat{R} , \\
 -\mathcal{L}_{h\phi} &= \lambda_{h\phi} H^\dagger H \phi^2 + \sigma H^\dagger H \phi
 \end{aligned}$$

Einstein frame:

$$g^{\mu\nu} = \Omega^{-1} \hat{g}^{\mu\nu} \qquad \Omega(h) \equiv 1 - \xi h^2$$

Potential:

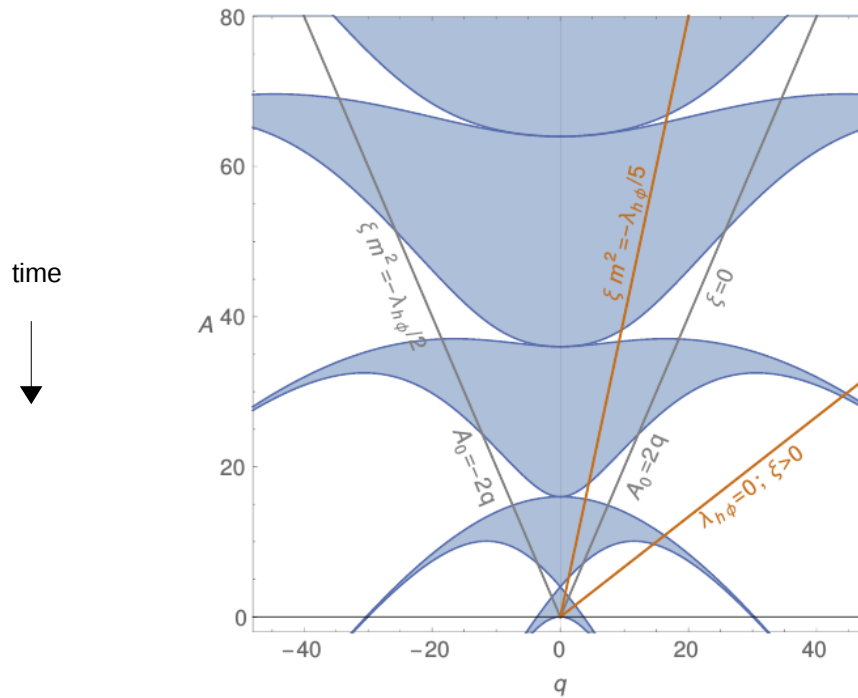
$$V(\phi, h_c) = \frac{1}{2} m^2 \phi^2 + \frac{1}{2} (\lambda_{h\phi} + 2\xi m^2) \phi^2 h_c^2 + \frac{1}{2} \sigma \phi h_c^2 + \frac{1}{4} \lambda h_c^4 + \dots$$

$$h_c \simeq h \left[1 + \left(\xi + \frac{1}{6} \right) \xi h^2 \right]$$

Mode EOM for the Higgs (simplified):

Mathieu Equation: $X_k'' + (A_k + 2q \cos 4z) X_k = 0 \quad z = mt/2$

inflaton-induced Higgs mass²

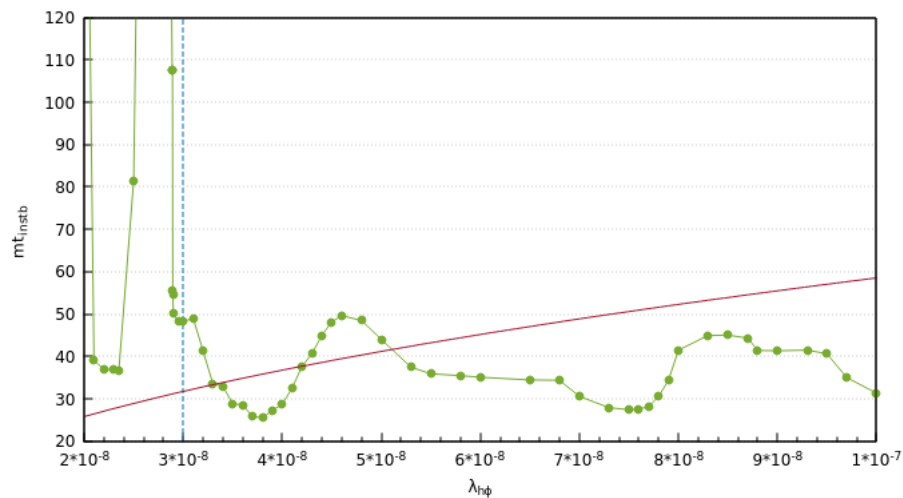


$A, q \sim t^{-2}$

white = *unstable*
shaded = *stable*

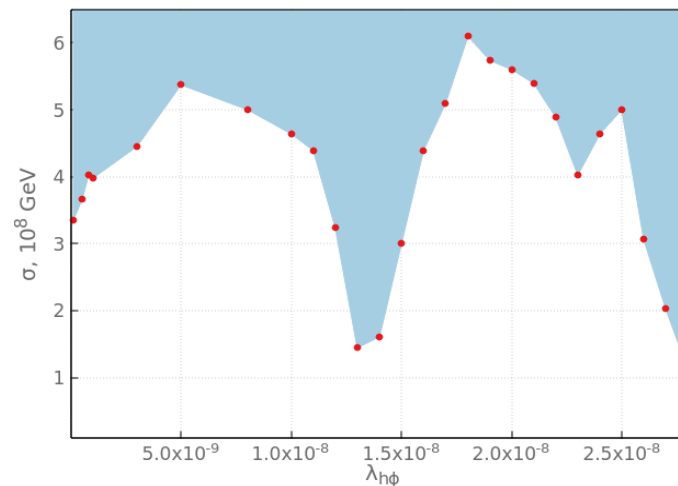
Bounds on the Higgs-inflaton couplings :

quartic :



$$\lambda_{h\phi} < 3 \times 10^{-8}$$

trilinear :



$$\sigma < 10^8 \text{ GeV}$$

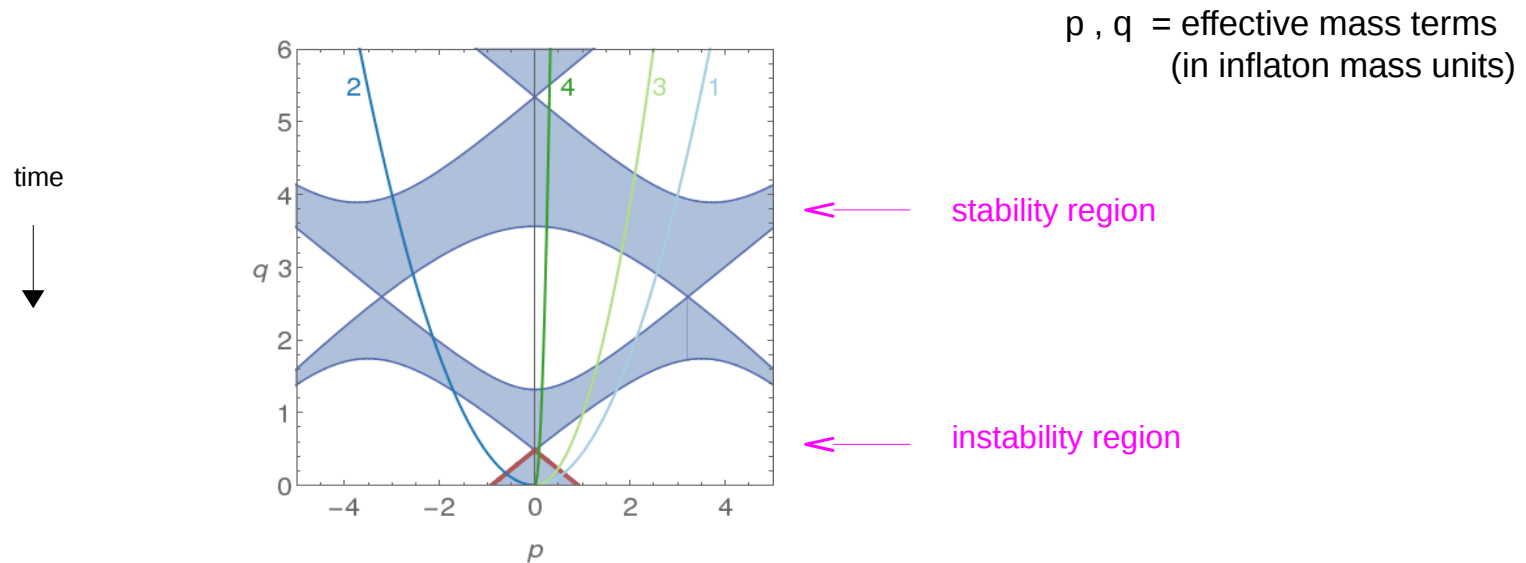
Conclusion

- Higgs sector is special
 - key to the hidden sector / DM / inflation
 - Higgs portal WIMP DM is viable
 - Higgs-inflaton interaction is crucial
-

Full "mixed" resonance:

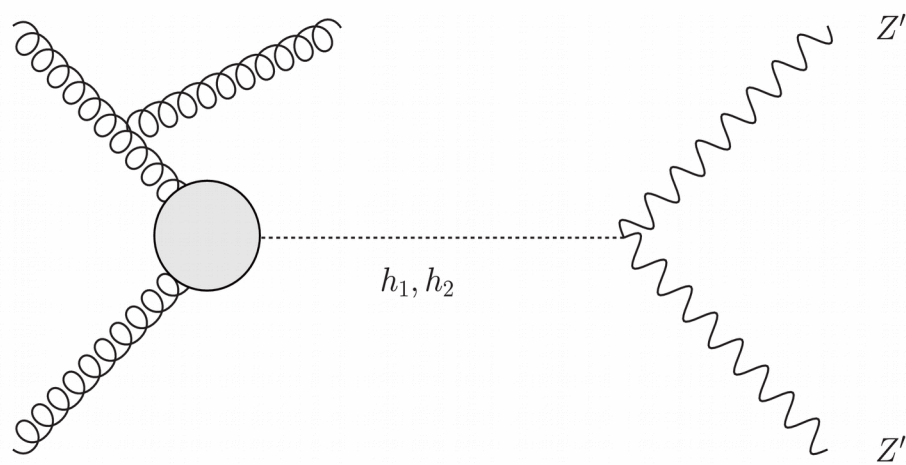
Whittaker–Hill equation (*no expansion limit*)

$$X'' + 2p \cos(2t) X + 2q \cos(4t) X = -A X$$



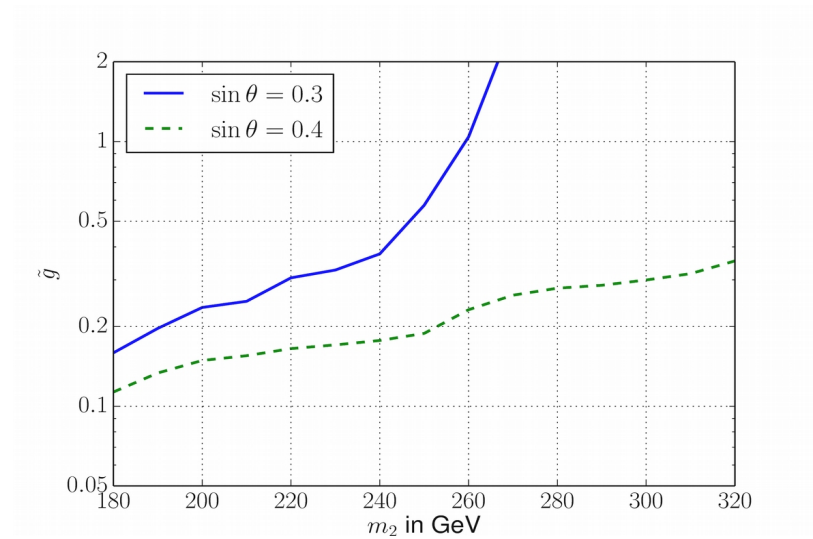
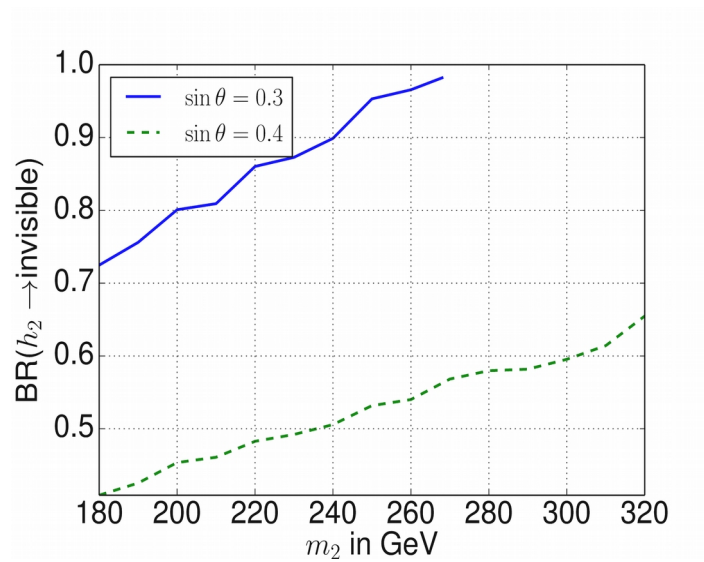
Higgsophilic gauge bosons at the LHC:

monojets (or VBF) + missing E_T



efficient for $m_2 > 2 m_{Z'}$.

Monojet results with 600 fb^{-1} :



+ similar results for VBF with already 300 fb^{-1}