

Cold Quark Matter, or “neutron” stars to 3 loops

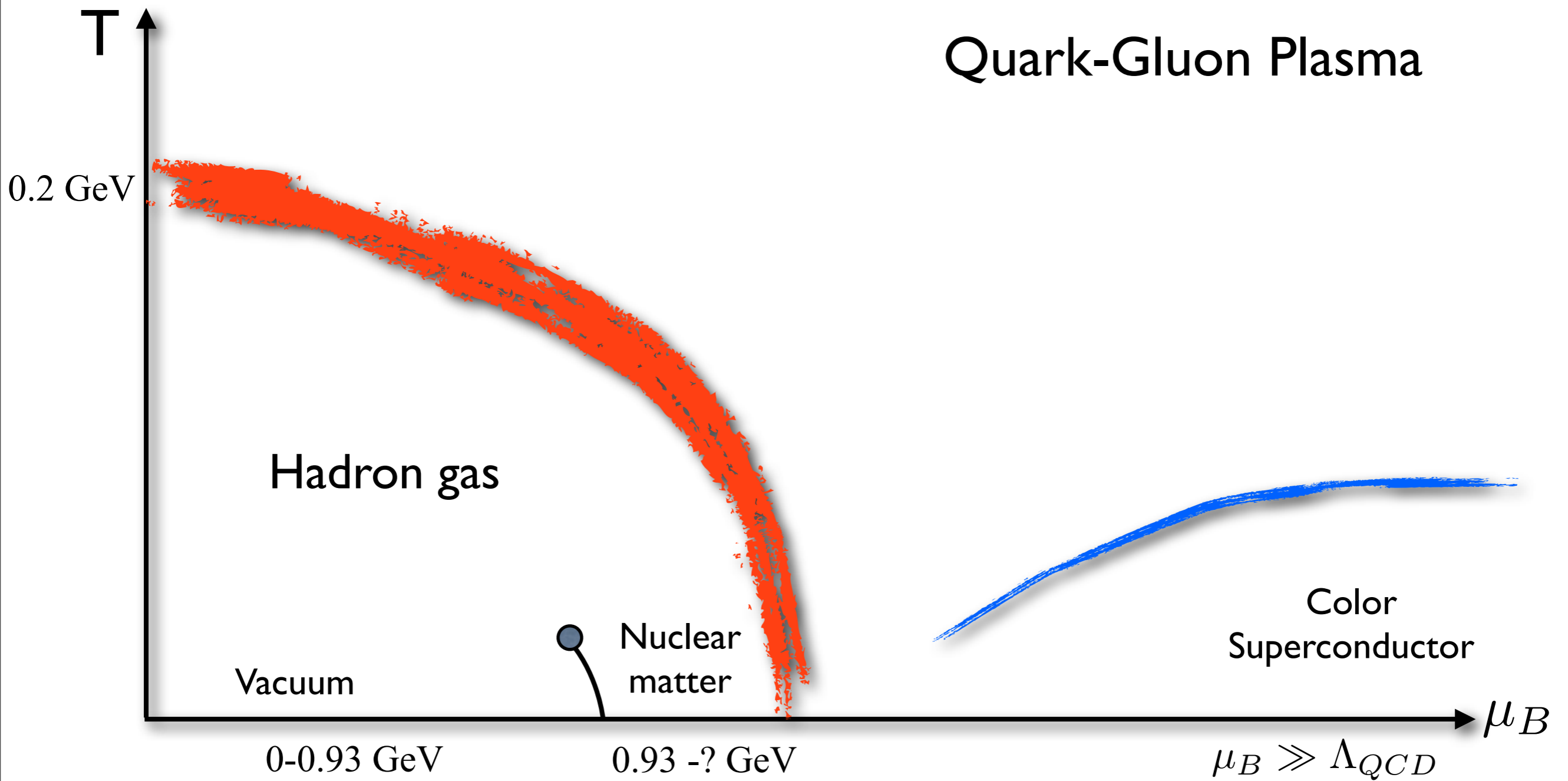
Alexi Kurkela
ETH Zürich,

with Paul Romatschke and Aleksi Vuorinen
arXiv: 0912.1856

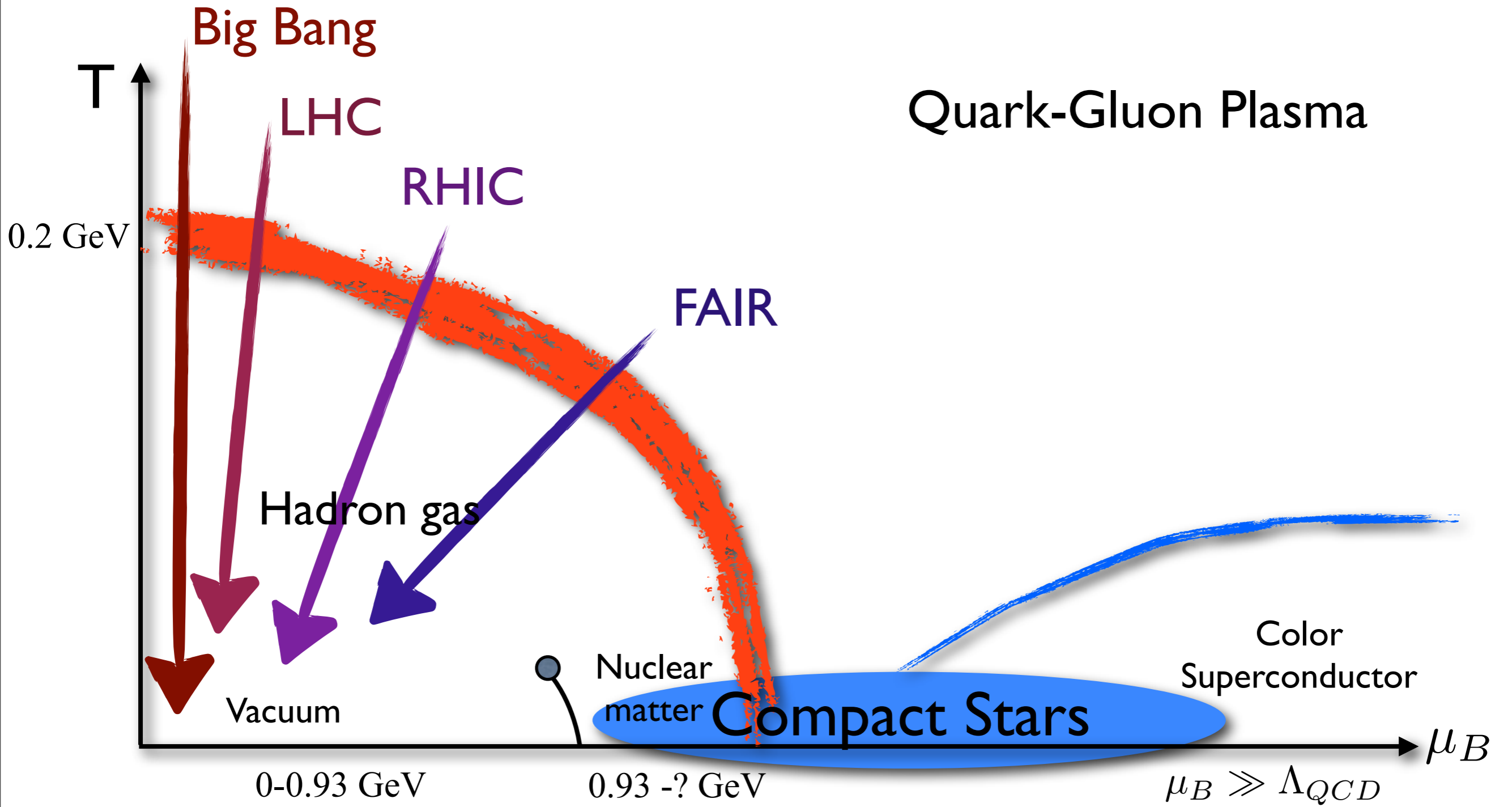
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Ultimate goal to understand strongly interacting matter throughout the phase diagram:



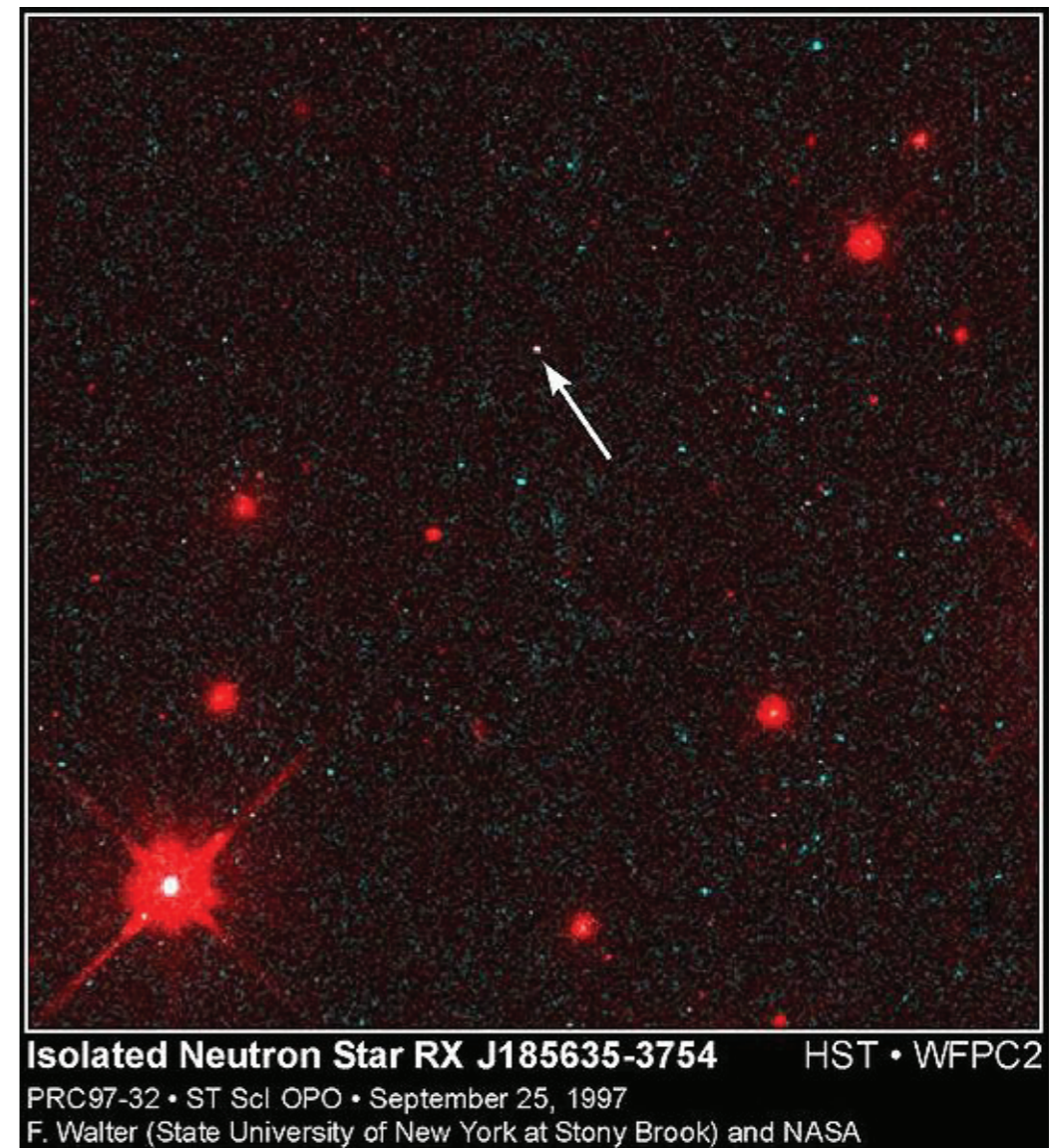
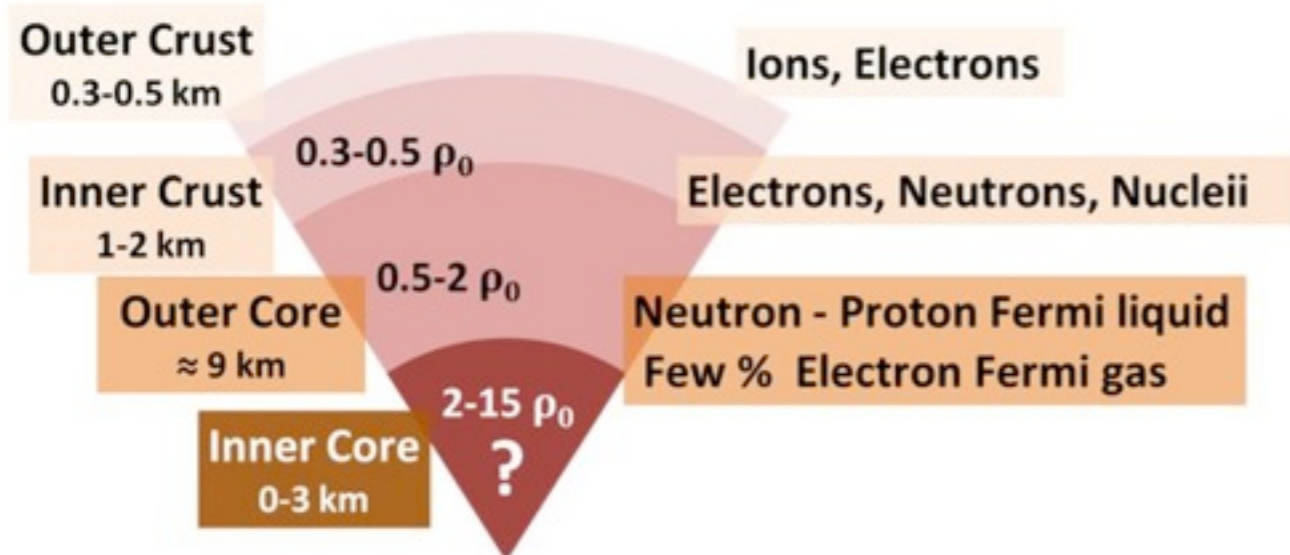
Ultimate goal to understand strongly interacting matter throughout the phase diagram:



Compact Stars:

Results from a core collapse of a star with $M \gtrsim 10M_{\odot}$

- Masses $\lesssim 2.0M_{\odot}$
- Radii $\sim 15\text{km}$
- $T < 10^6\text{K} < \text{KeV}$ ($T_0 \sim 30\text{MeV}$)
- $n \lesssim 15\rho_0$ ($\rho_0 = 0.16\text{fm}^{-3}$)



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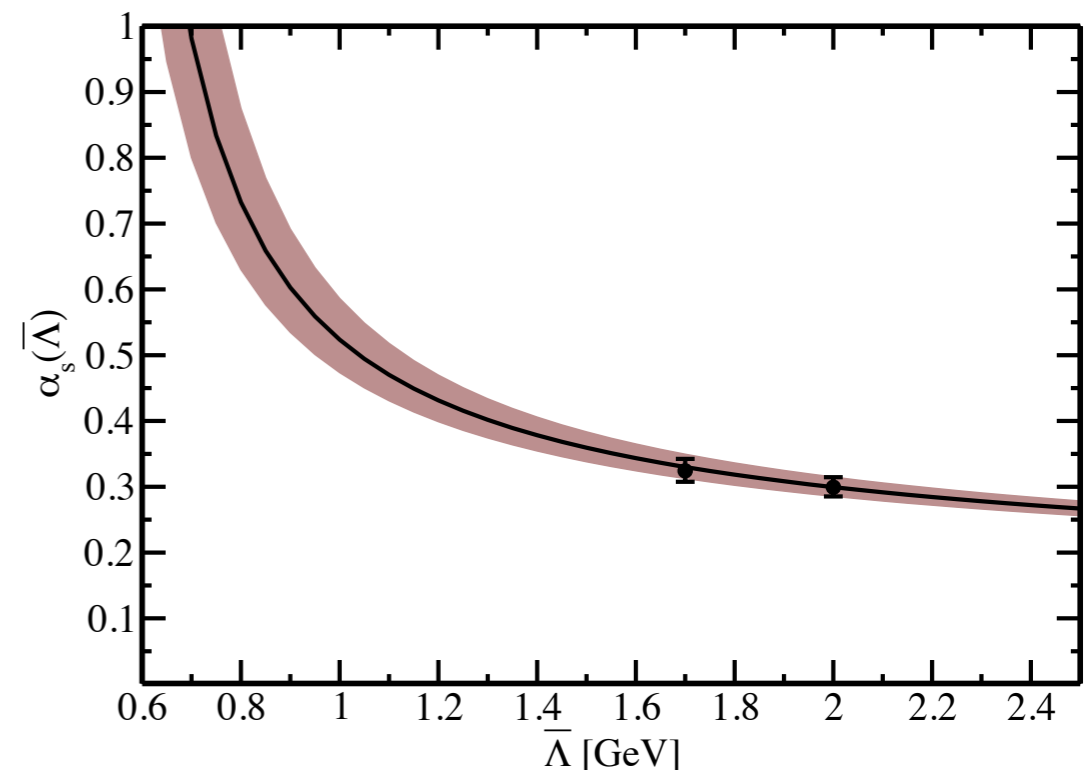
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- At $\mu > T$: Simulations become impossible due to **sign problem**.
- Resort to approximations ($T = 0$):
 - At low densities ($\rho_B < 0.16 \text{ fm}^{-3}$, $\mu_B \sim 1 \text{ GeV}$):
Quantum many-body theory. Dynamics of nucleons, hyperons, etc..
 - At (asymptotically) high densities:

$$\alpha_s(\mu) \sim 1/\log(\mu^2)$$

→ **Perturbation theory**

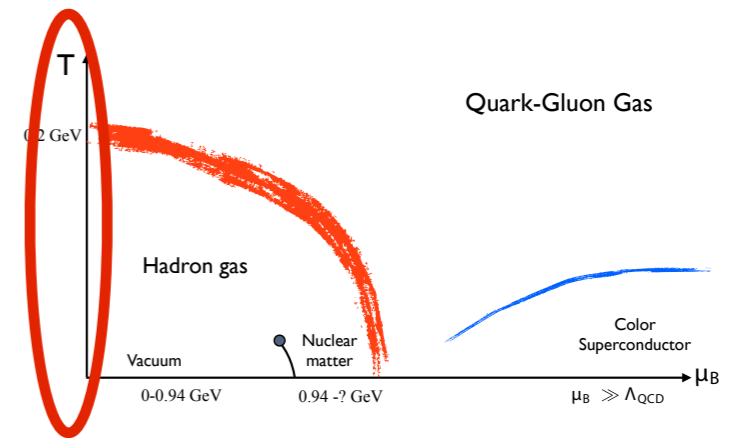


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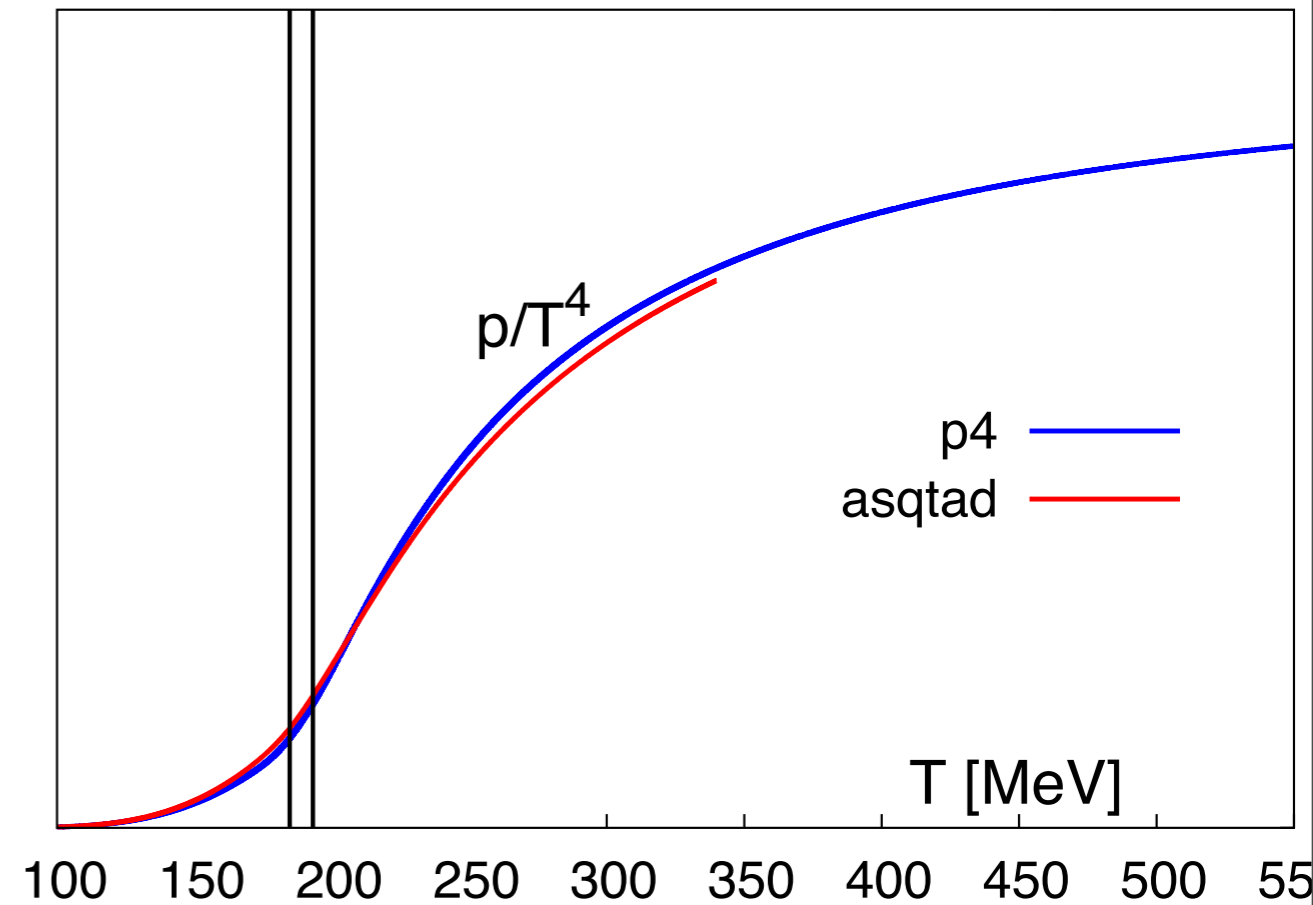
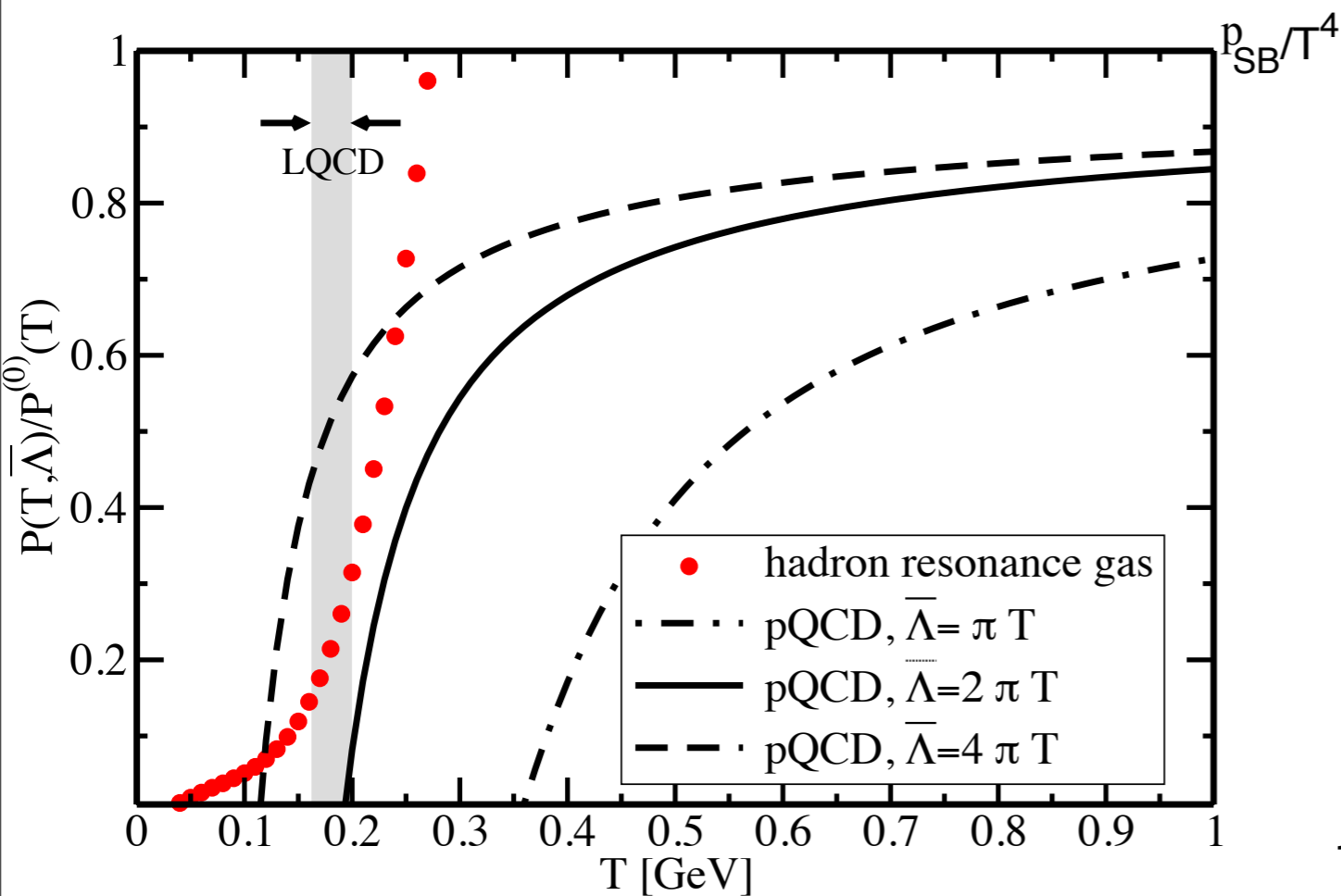
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→ **Perturbation theory**
- **Strategy**: Interpolate between high and low densities to get a unified description of EoS.
 - Relies on **not** having exotic phase between!

Inspiration from finite-T case:



Hot Quark Matter ($\mu=0$)

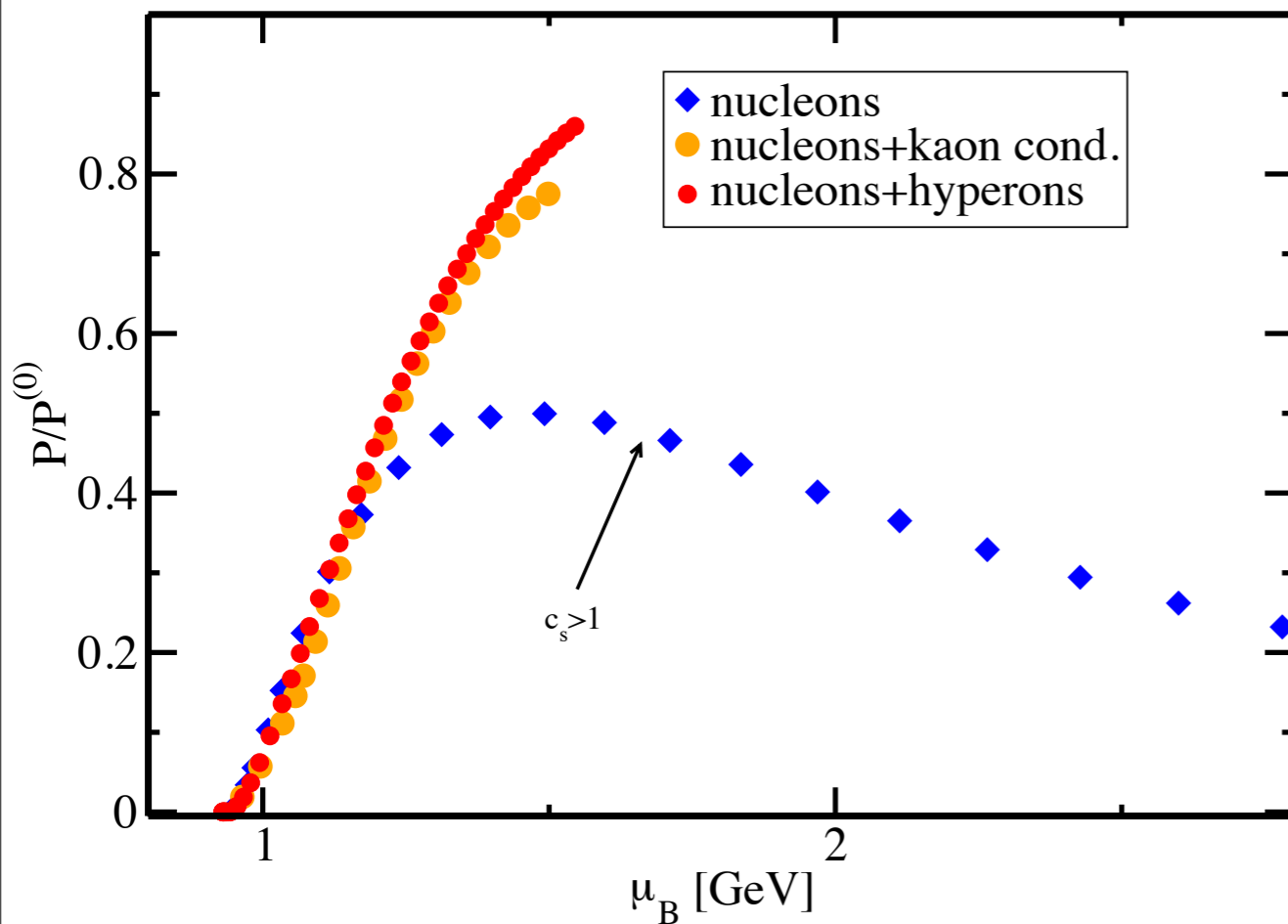
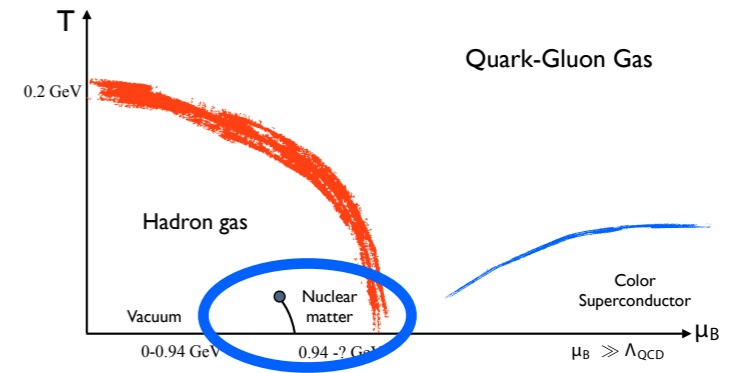


- Low-T : Hadron resonance gas
- High-T : Resummed perturbation theory ($\alpha_s^{5/2}$)

Lattice

Bazavov et al. 0903.4379

At low densities:



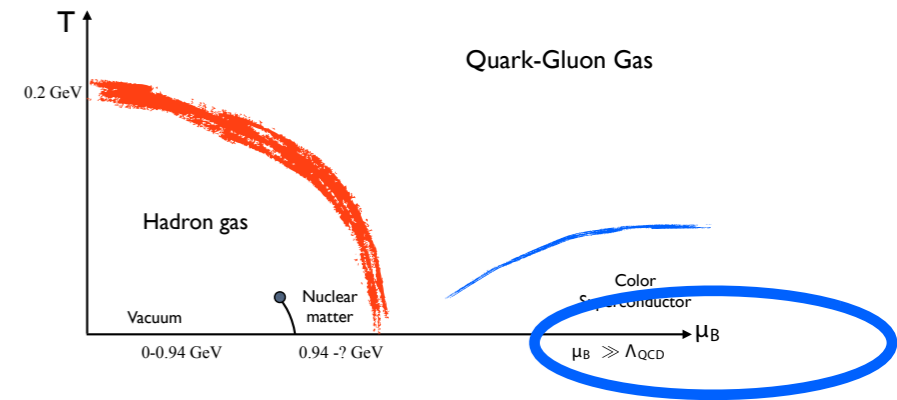
Nucleons (=APR nucl-th/9804027):

- Input:
 - 2-body: Argonne v_{18}
 - 3-body: Urbana IX
- Variational Chain Summation:
 - EoS for PNM and SNM
- Interpolate the correct neutron/proton fraction using Skyrme-model ansatz

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Perturbative evaluation of the EoS:



- Thermodynamics defined by the grand potential

$$\Omega(\mu_u, \mu_d, \mu_s, m_s) = -T \log \int \mathcal{D}\bar{\psi} \mathcal{D}\psi \mathcal{D}A_\mu e^{-\int d^4x \mathcal{L}_{\text{QCD}}}$$

$$\mathcal{L}_{\text{QCD}} = \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \bar{\psi}_i (\gamma_\mu D_\mu + m_i - \mu_i \gamma_0) \psi_i.$$

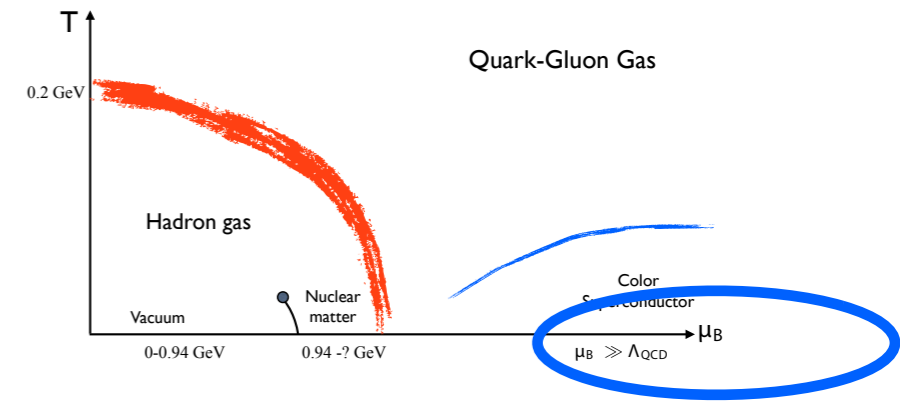
- All thermodynamical quantities derived from $\Omega(\mu_u, \mu_d, \mu_s, m_s)$

$$pV = -\Omega(\mu_u, \mu_d, \mu_s, m_s)$$

$$n_i = -\partial_{\mu_i} \Omega(\mu_u, \mu_d, \mu_s, m_s)$$

$$\varepsilon = -p + n_i \mu_i$$

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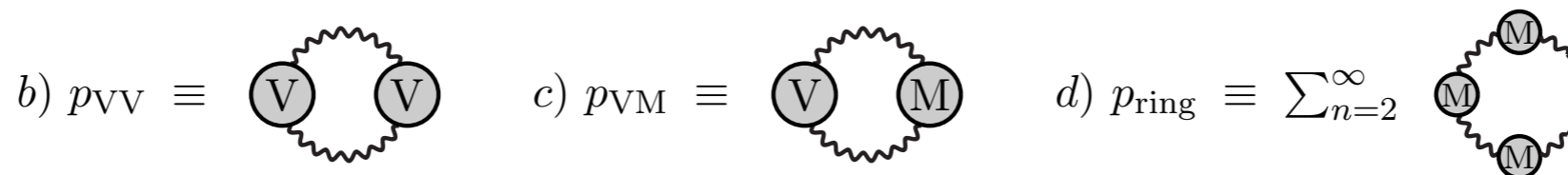
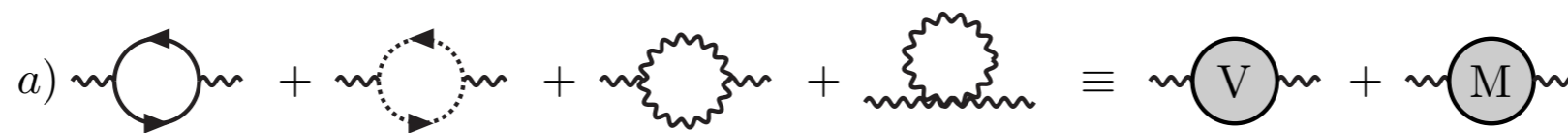
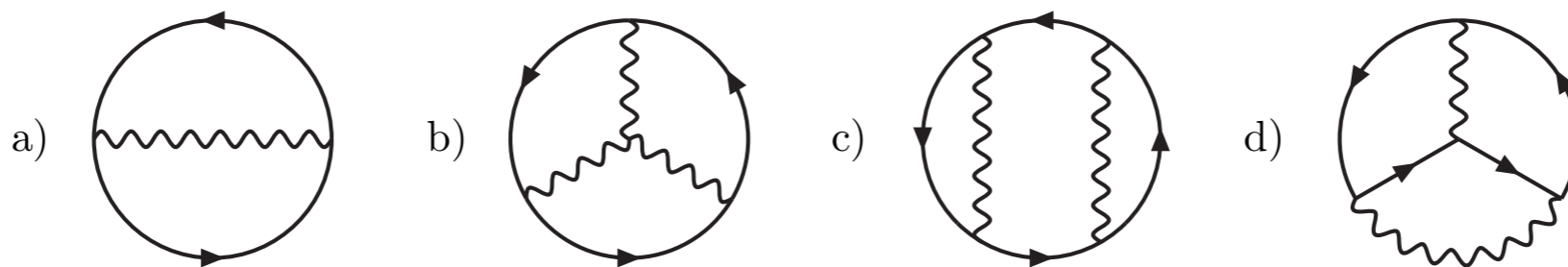


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- $\Omega(\mu_u, \mu_d, \mu_s, m_s)$ available through 1PI diagrams:



Physical EoS at high density:

- Electric neutrality and beta-equilibrium fix all but one chemical potentials:

$$p(\mu_u(\mu), \mu_d(\mu), \mu_s(\mu)) = p(\mu)$$

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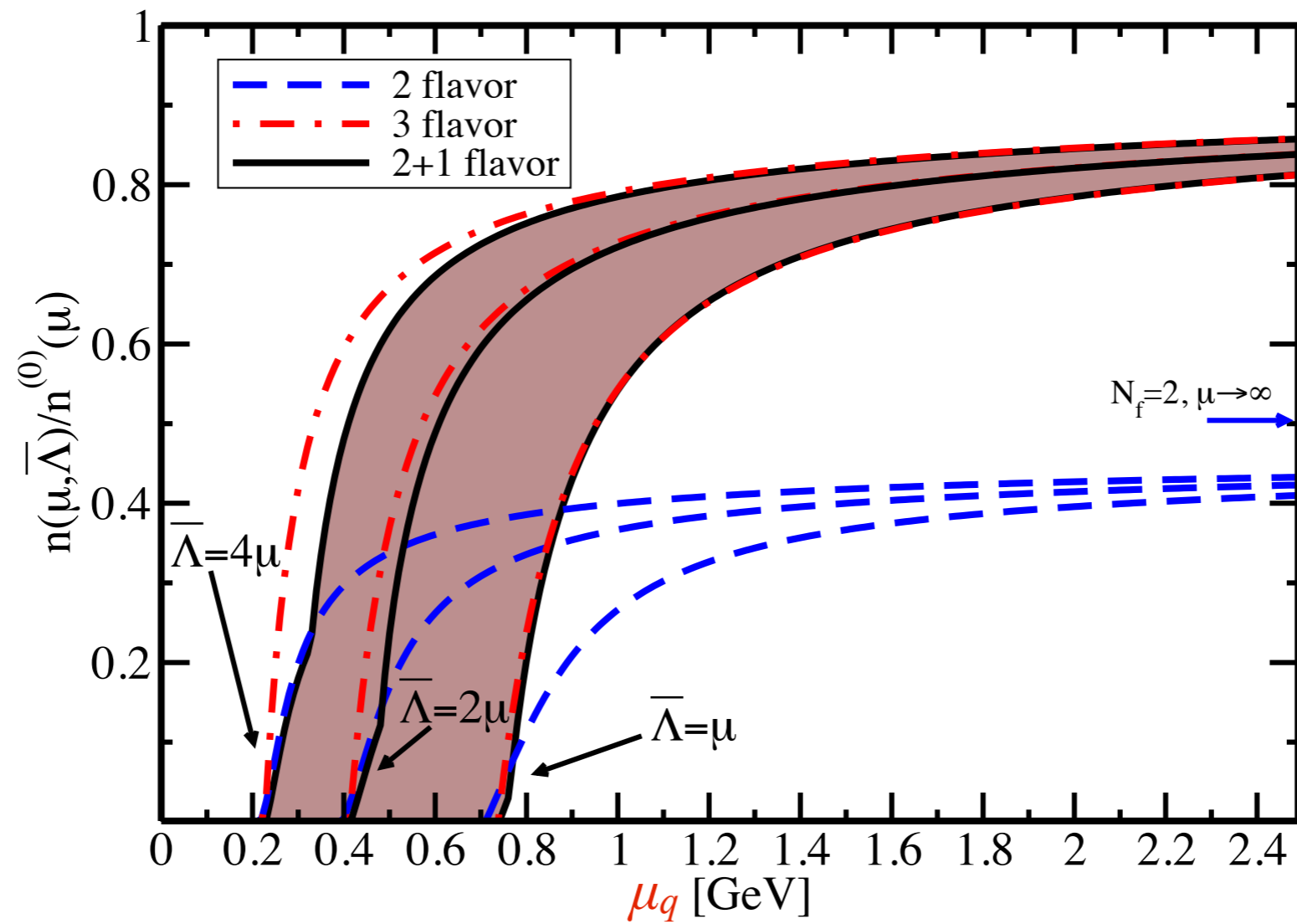
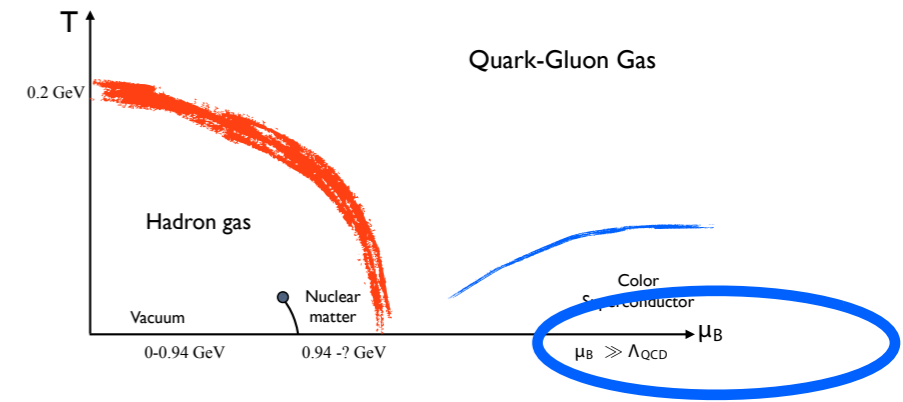
- Define $p(\mu = 0) = 0$, outside perturbative reach obviously...
- The theory has a pairing instability \rightarrow Non-perturbative term in $p(\mu)$

$$p(\mu) = p^{(\text{pert})}(\mu, \bar{\Lambda}) - B + \frac{\Delta^2 \mu_B^2}{3\pi^2}$$

With $\Delta=0,\dots,100$ MeV ($\Delta \ll k_F$)

And vary B for all allowed values

At high densities:

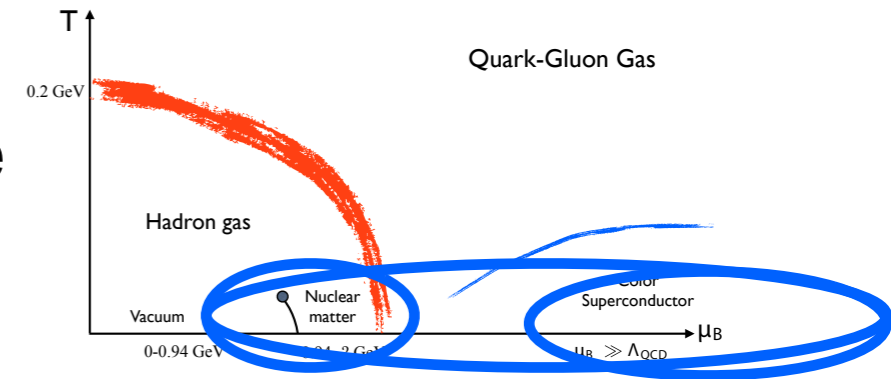


$$n(\mu_B, \bar{\Lambda}) = \partial_{\mu_B} p(\mu_B, \bar{\Lambda})$$

Outline:

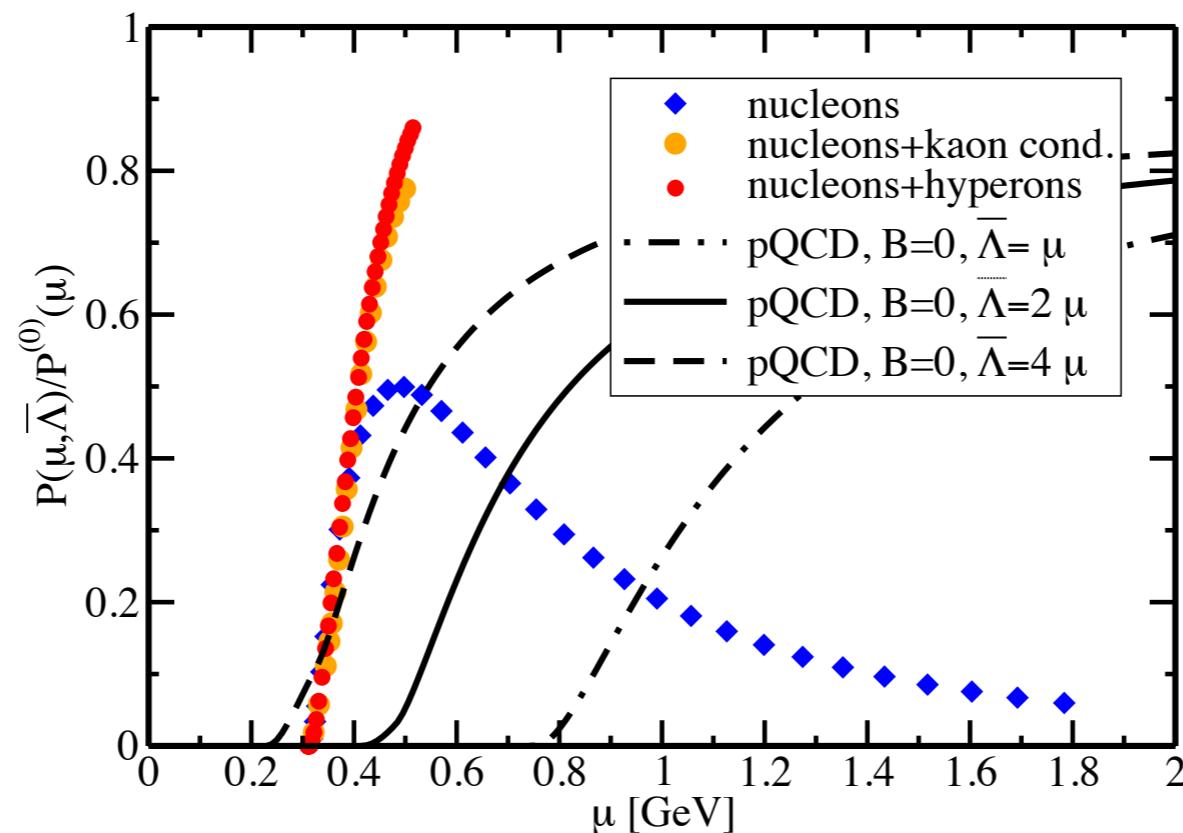
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If there are no **exotic phases**, there will be a phase transition between **hadronic** and **quark matter** phases at some μ_{pt}

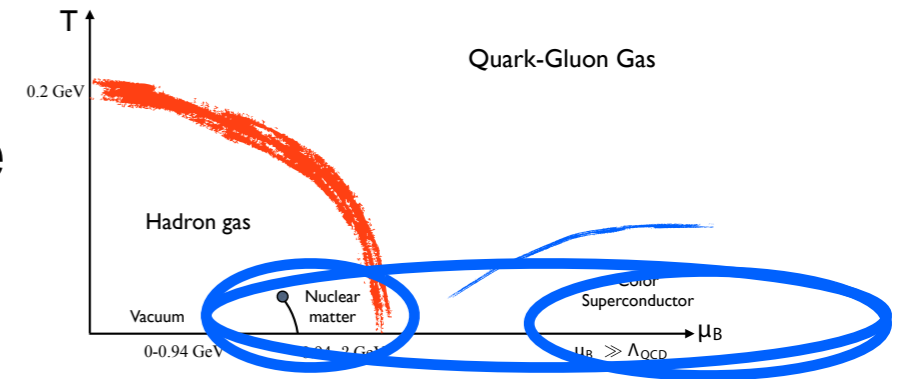


- Extrapolate the low- and high-density EoS to the intermediate region
- Catalog all possible self consistent EoS's ($B(\mu_{pt}), \bar{\Lambda}$)
 - Equal pressure at phase transition
 - Monotonically increasing energy density

$$\bar{\Lambda}_{\overline{\text{MS}}}=378 \text{ MeV}$$

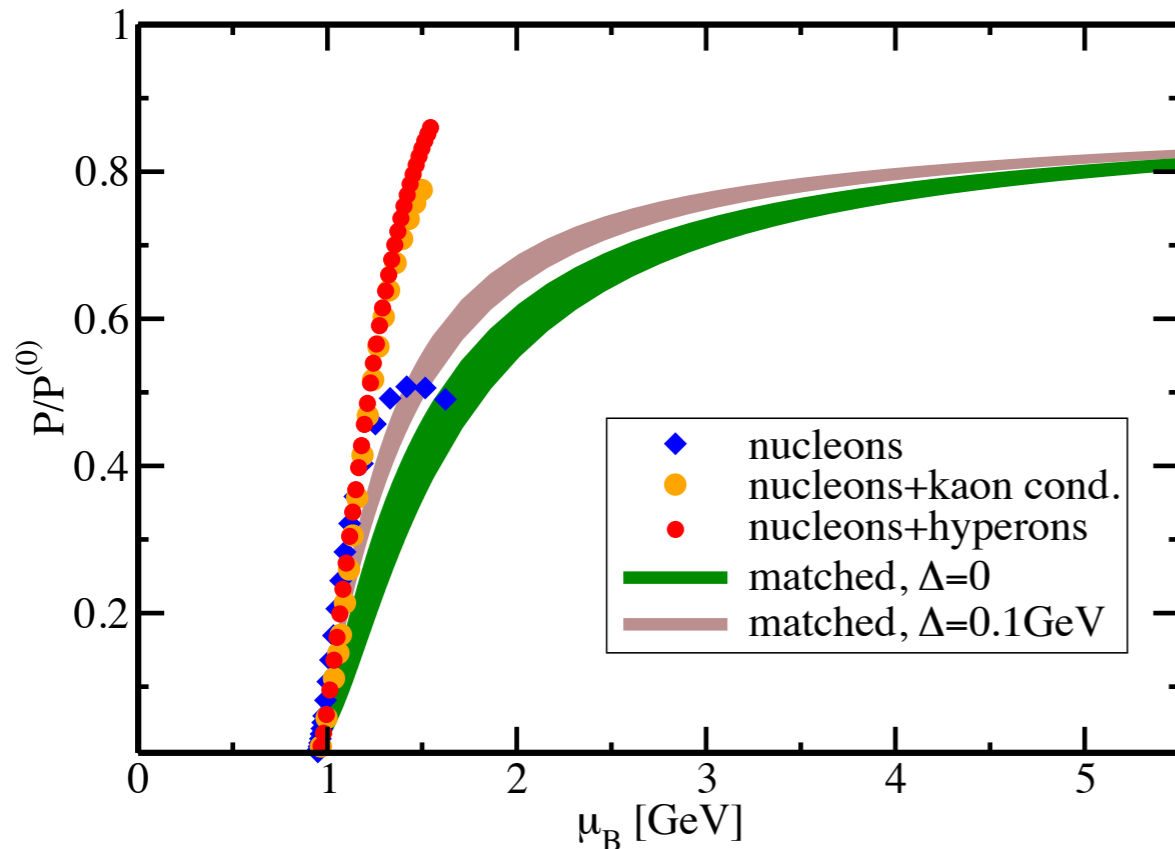


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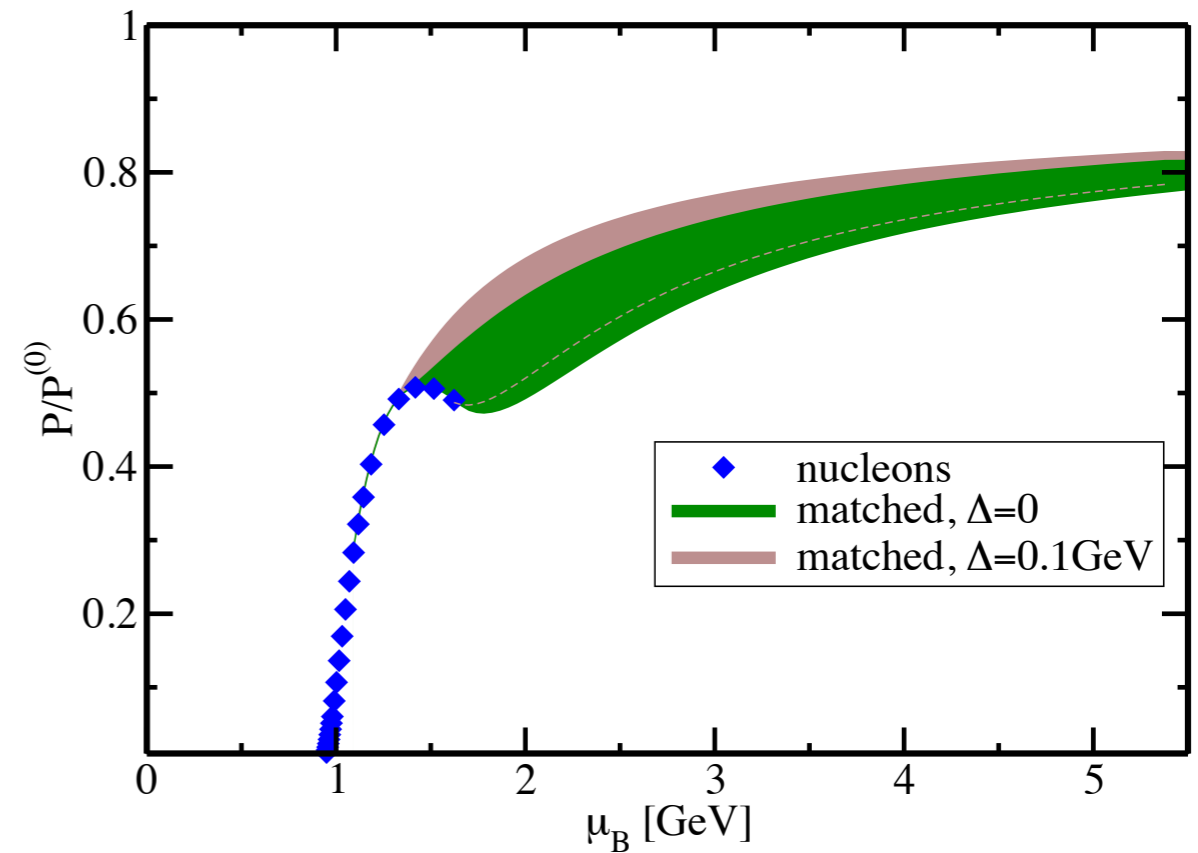
- Matching possible in two disjoint regions:

Case I ('low μ ') Matching



$$0.16\text{fm}^{-3} < \rho_B^c \lesssim 0.32\text{fm}^{-3}$$

Case II ('high μ ') Matching



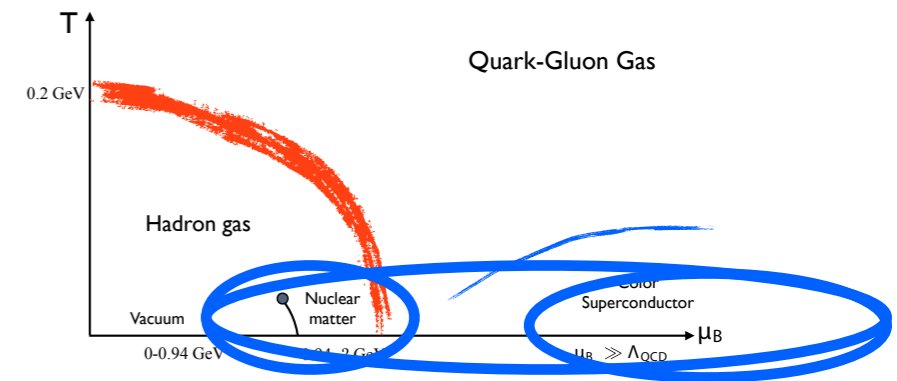
$$\rho_B^c > 0.64\text{fm}^{-3}$$

Represents the best educated guess available for the true EoS on full μ -range
 The location of μ_{pt} not available via perturbation theory, only via matching.

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Strange quark matter hypothesis:



If the energy per baryon in quark matter is less than

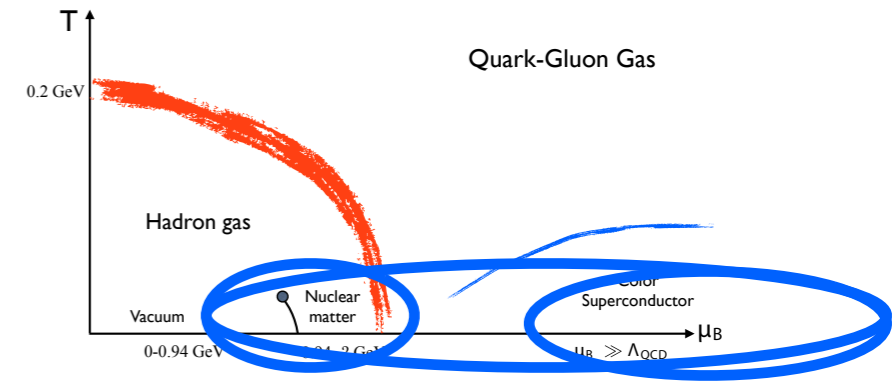
$$E/A = 3\mu_c = 0.93\text{GeV} \quad {}^{56}\text{Fe}$$

then quark matter is the true ground state \rightarrow Nuclear matter metastable.

Life time:

- **Nucleons \rightarrow 2 Flavor quark matter:**
 - Equilibration through Strong interactions $\rightarrow t_{eq} \sim 1/\Lambda_{\text{QCD}}$
 - Short lived. Ruled out by “experiment”
- **Nucleons \rightarrow strange quark matter:**
 - Equilibration through Weak interactions $t_{eq} \sim 10^{60}$ years for $A > 6$
 - Adding d.o.f's increases pressure \rightarrow more likely to be stable
 - Experimentally plausible, lets find out what the theory says!

Strange quark matter hypothesis:

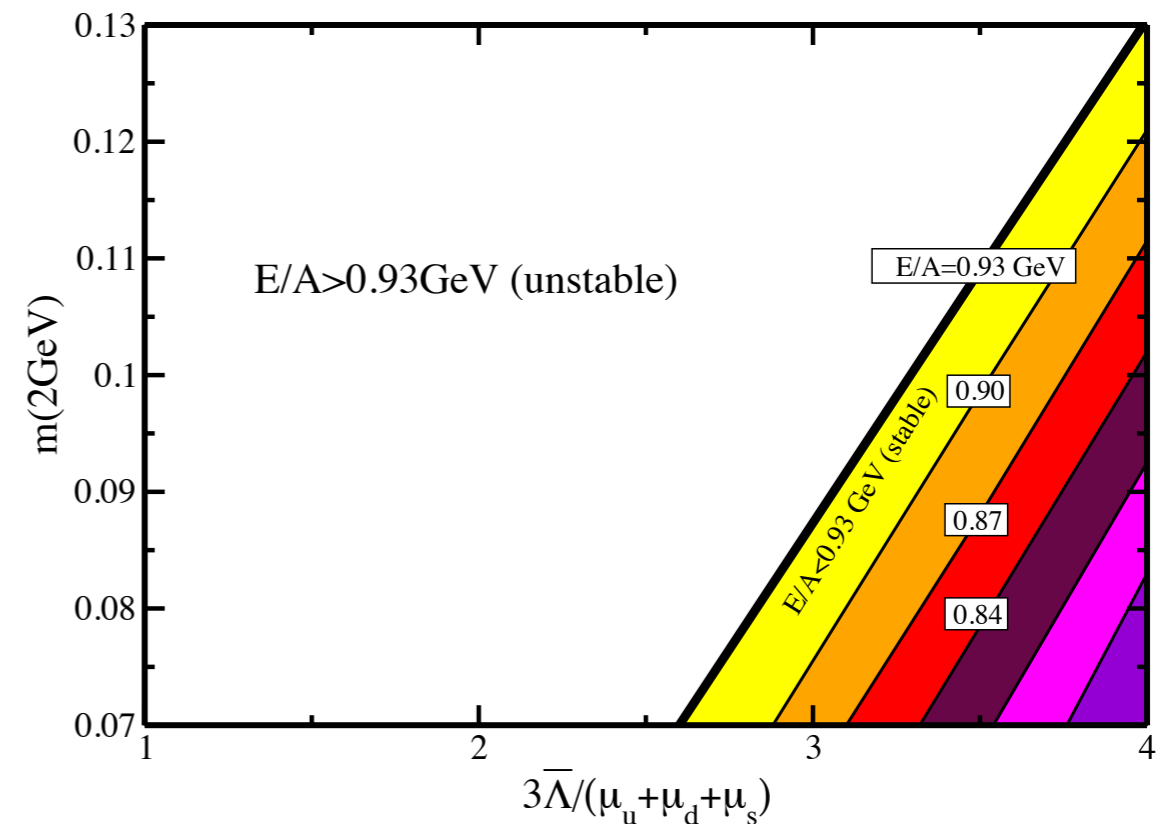
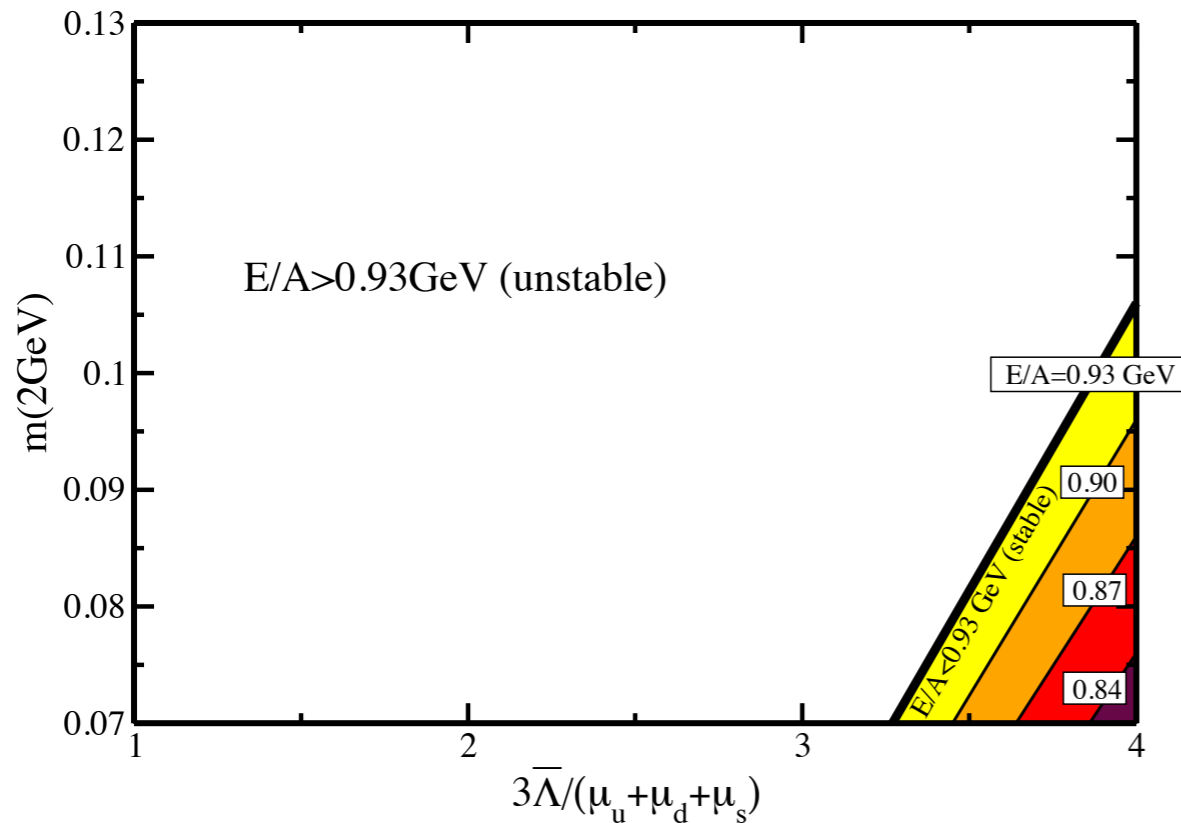


Strategy: Find out if SQM stable in the parameter space $(B, \bar{\Lambda})$ with

- $n_s > 0$, quark mass essential!
- $E/A > 0.93 \text{ GeV}$

Normal Quark matter, $\Delta=0$, $\Lambda_{\overline{\text{MS}}}=0.378\text{GeV}$

CSC, $\Delta=100 \text{ MeV}$, $\Lambda_{\overline{\text{MS}}}=0.378\text{GeV}$



- Parameter space very **hostile**, adding CSC makes bit more plausible
- Absolutely stable SQM disfavored but not ruled out

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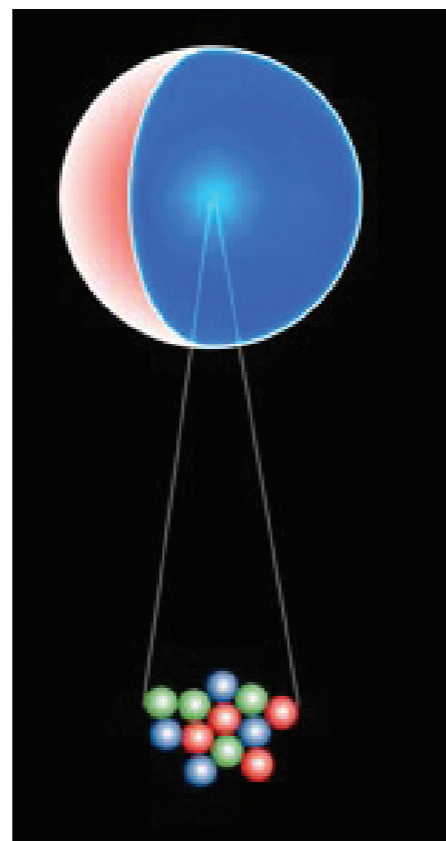
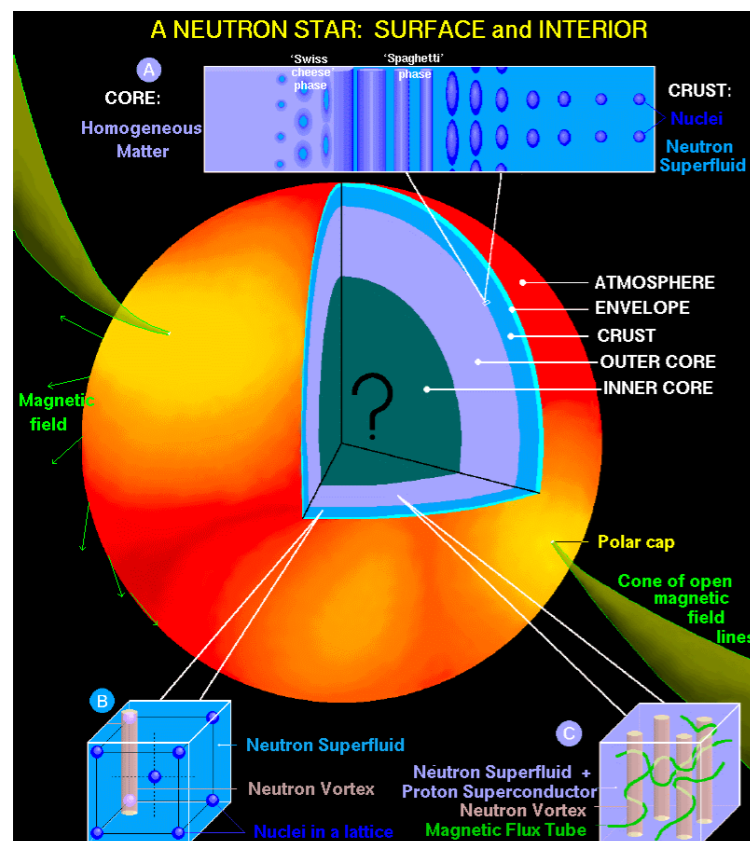
Compact stars:

- Compact stars formed of degenerate matter form a sequence in M-R plane (unlike white dwarfs)
- The M-R relation is very sensitive to the EoS
- M-R relation solved from the TOV-equations:

$$dM(r) = 4\pi r^2 \varepsilon(r) dr,$$

$$dP(r) = -\frac{G(P(r) + \varepsilon(r))(M(r) + 4\pi r^3 P(r))}{r(r - 2GM(r))} dr,$$

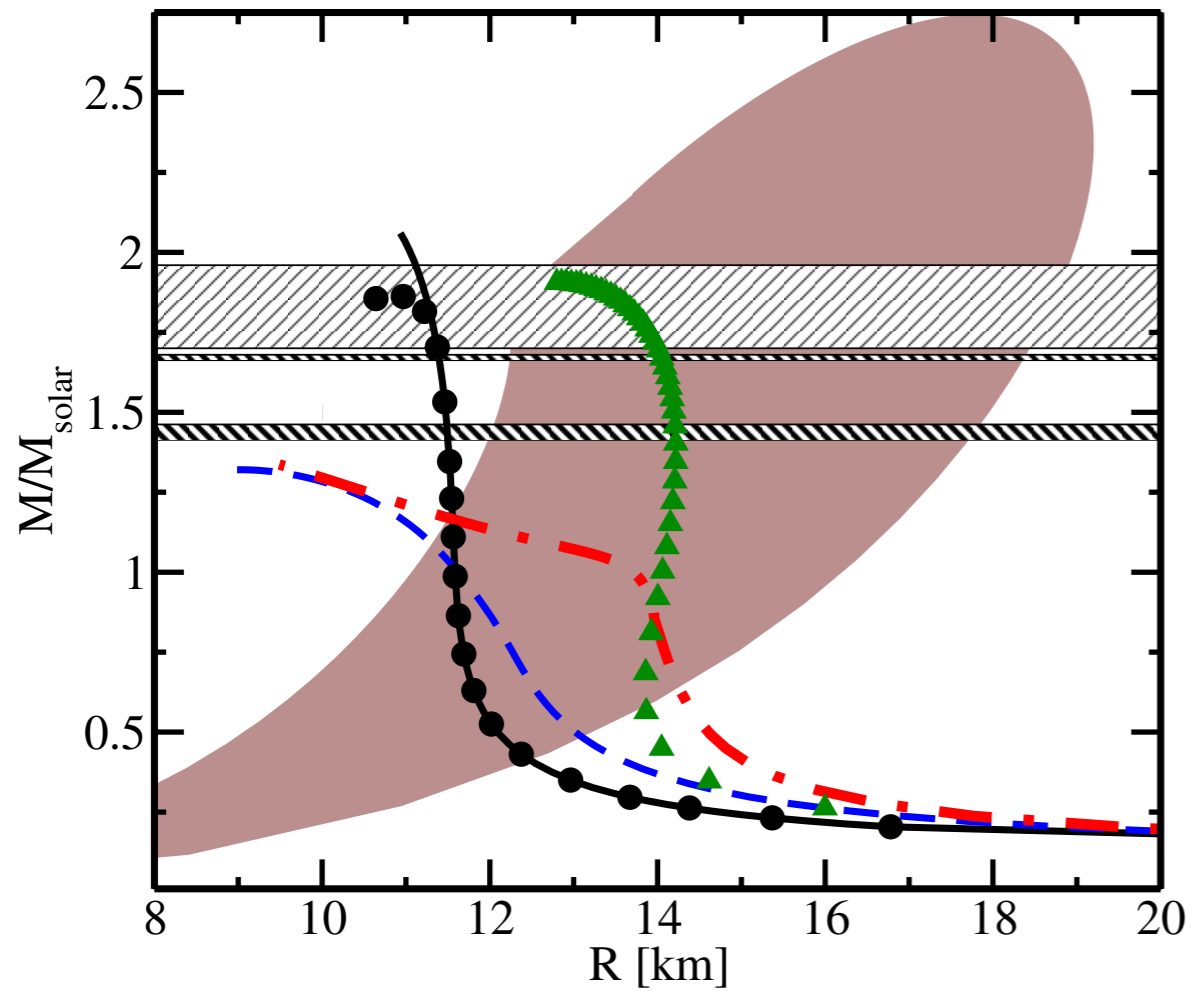
- Takes $\varepsilon(p)$ as input



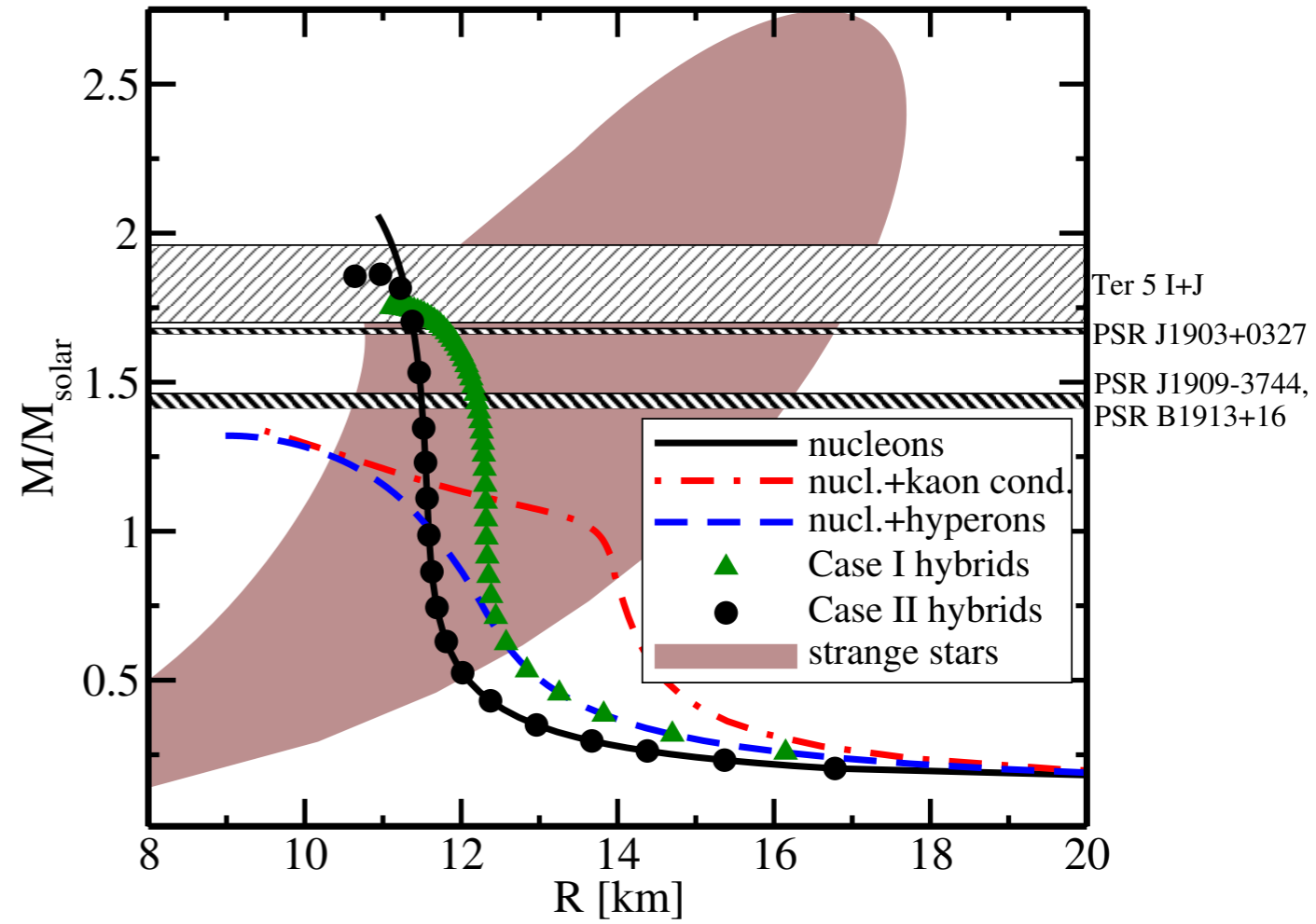
Let's consider compact stars made of:

- Pure Nuclear matter
- Pure Quark matter
- Hybrid stars with
 - large quark core with thin nucleonic crust (case I)
 - small quark core with thick nucleonic crust (case II)

Normal Quark Matter ($\Delta=0$)

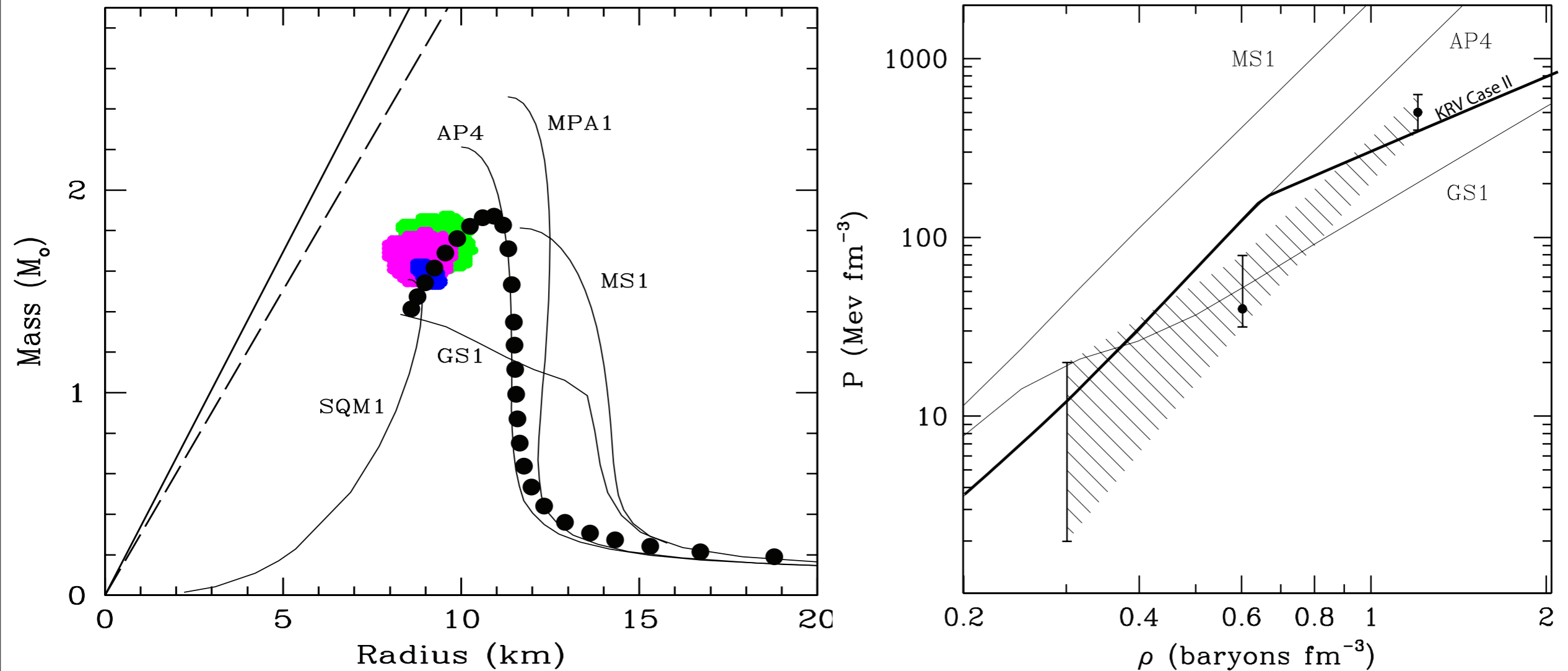


CSC, $\Delta=100$ MeV



- Effect of CSC very small.
- Hyperonic/Kaonic EoS ruled out
- For neutron/hybrid stars $M_{\max} \sim 2M_{\text{solar}}$
- Cannot exclude very massive **strange stars**
 - Dense quark stars ruled out

New observations from February:



- Case II agrees with the data better than any of the standard EoS
- Overestimates the radius
 - Accounting for the (possible) 2-component admixture in the transition:
 - Reduces radius, doesn't affect the maximum mass.
 - Smoothens EoS around transition
- Superconductivity reduces radius → Improve the treatment of CFL

Conclusions:

- The grand potential of QCD at finite density with finite m_s computed to α_s^2 .
 - Needed to create new perturbation theory machinery to overcome technical challenges
- Modeled the EoS in full range of μ_B (three logical possibilities):
 - **Hadron / quark matter transition:**
 - Realistic description for full range of μ_B
 - **Absolutely stable strange quark matter**
 - is disfavored but not ruled out
 - ...but an observation of a $M > 2M_{\text{solar}}$ would be a strong evidence in the opposite direction.
 - **Exotic (non-CSC) phases between hadrons and quark matter**
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Our Case II matching seems to perform better than any of the standard EoS in describing the recent experimental data.

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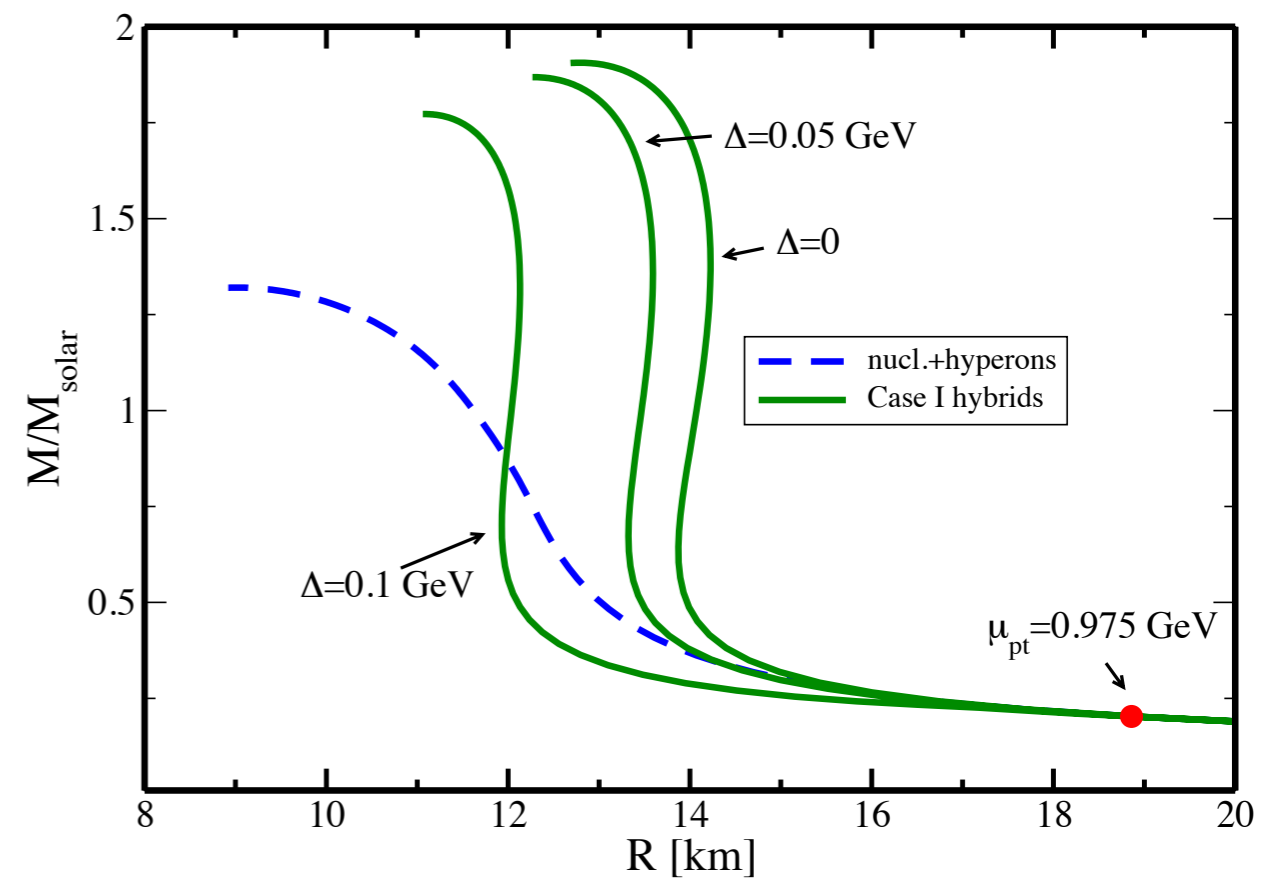
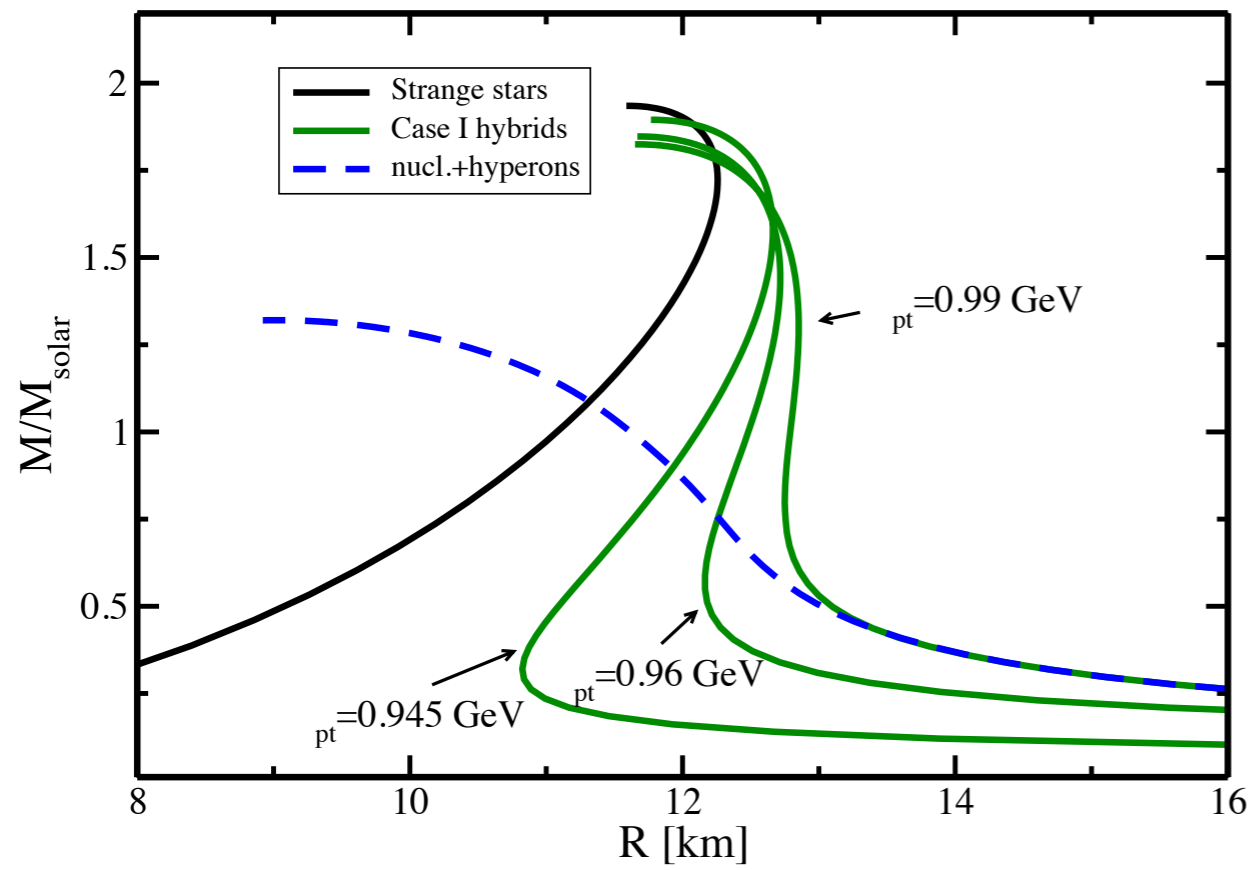
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 - By computing the mismatch of the fermi spheres
 - Assessing the different possibilities for CSC: CFL, 2CS...
- **Improve the perturbative calculation:**
 - $\alpha_s^2 \log(\alpha_s)$: Only ring diagrams involved
 - α_s^3 : Major undertaking
- **Improve the astrophysical modeling:**
 - Two-component mixtures of hadronic and quark matter
 - Moment of inertia, glitches
 - Neutron star oscillations
 - Rotating stars, r-modes
 - Cooling rates and transport effects

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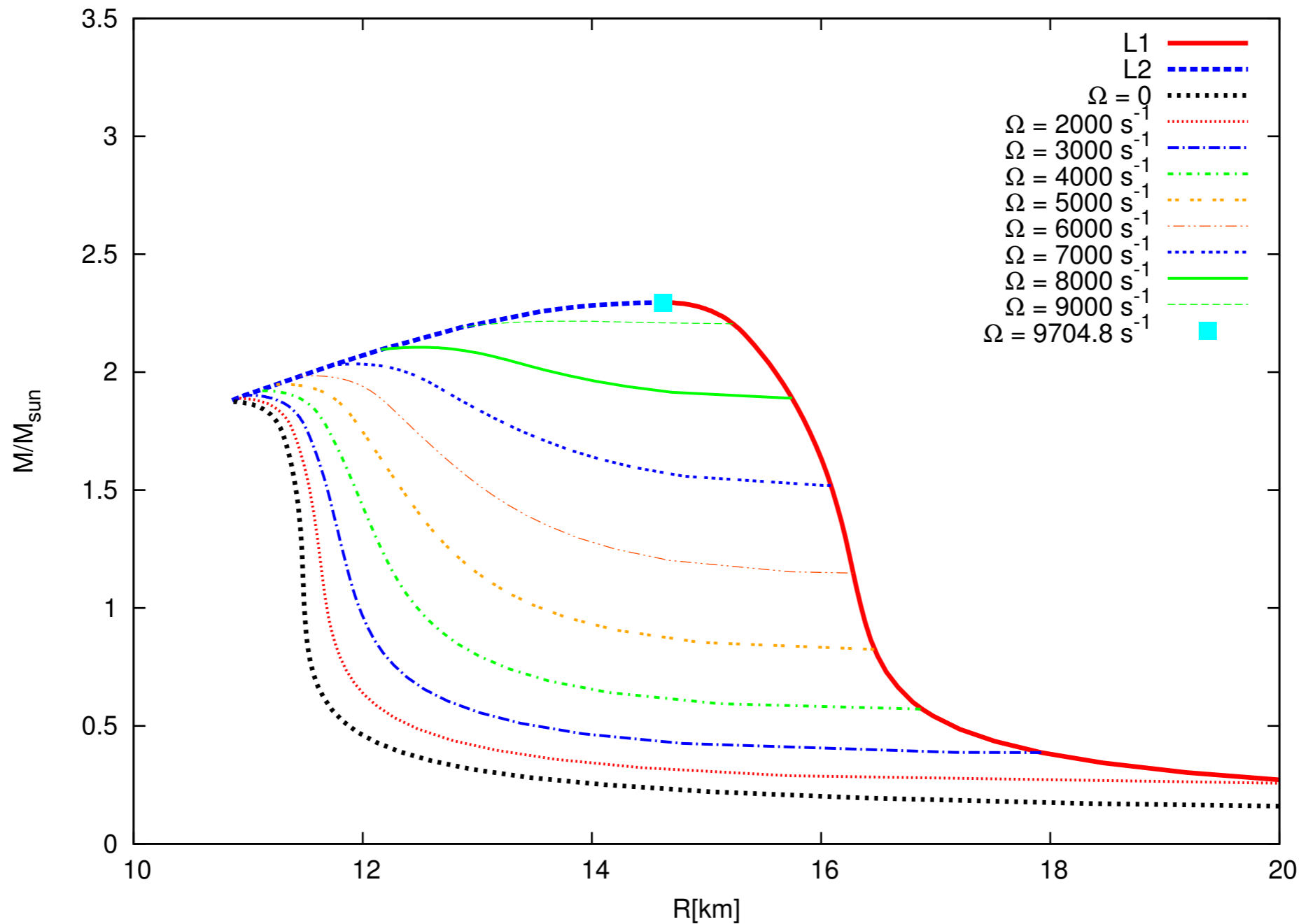
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and of course: the observations are advancing **very** fast,
new data expected to come anytime!!

Effect of the matching density and CSC:



Effect of rotation to the MR-curves (Case II matching)



PRELIMINARY!!! With Bin Wu