

# The functional renormalisation group and applications to the phase structure of QCD

**Jan M. Pawłowski**

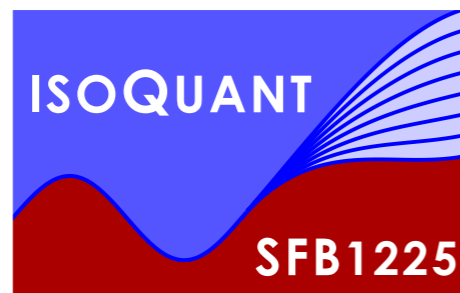
**Universität Heidelberg & ExtreMe Matter Institute**

**Trento, June 11<sup>th</sup> - 15<sup>th</sup> 2018**



GEFÖRDERT VOM

Bundesministerium  
für Bildung  
und Forschung



# Material

## Lecture notes

Non-perturbative methods in gauge theories

hand-written

Critical phenomena

hand-written

QCD

tex

## Topical reviews

Collection of reviews & lecture notes on the FRG & DSE

Structure of the FRG: Aspects of the FRG

JMP '05, Annals Phys.322:2831-2915,2007

## talks

The FRG approach to gauge theories & applications to QCD

JMP, ERG 2012 Aussois

Aspects of the QCD phase diagram and the EoS

B.-J. Schaefer, CompStar 2012 School Zadar

Schladming 2011: Physics at all scales: The Renormalization Group

Schladming 2013: The phase diagram of QCD: Thermodynamics, order parameters & dynamics

JMP, Schladming 2013

# Outline

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- **(I) Introduction to the phase diagram of QCD**
- **(II) Functional Renormalisation group for QCD**
- **(II) Phase structure of QCD & dynamics**

# (I) Introduction to the phase diagram of QCD

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- **Heavy ion collisions**

- Phases of a heavy ion collision

- **Phase structure of QCD**

- Perturbative QCD & asymptotic freedom
  - chiral symmetry breaking
  - confinement

# **(II) Functional Renormalisation group for QCD**

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- **Introduction to the functional renormalisation group**

- Derivation of the flow equation

- Expansion schemes

- Optimisation and error control\*

- **FRG for QCD**

- FRG for QCD and  $T=0$  Yang-Mills theories

- Dynamical hadronisation

- QCD correlation functions at  $T=0$

# **(III) Phase structure of QCD and dynamics**

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- **Yang-Mills theory at finite temperature**
  - Order parameter potential for confinement
  - Correlation functions at finite temperature
  - Polyakov loop from functional methods
  
- **Application to the phase structure of QCD and dynamics\***
  - QCD-assisted hydrodynamics\*
  - QCD-assisted transport\*
  - QCD at imaginary chemical potential\*

# **(I) Introduction to the phase diagram of QCD**

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- **Heavy ion collisions**

- Phases of a heavy ion collision

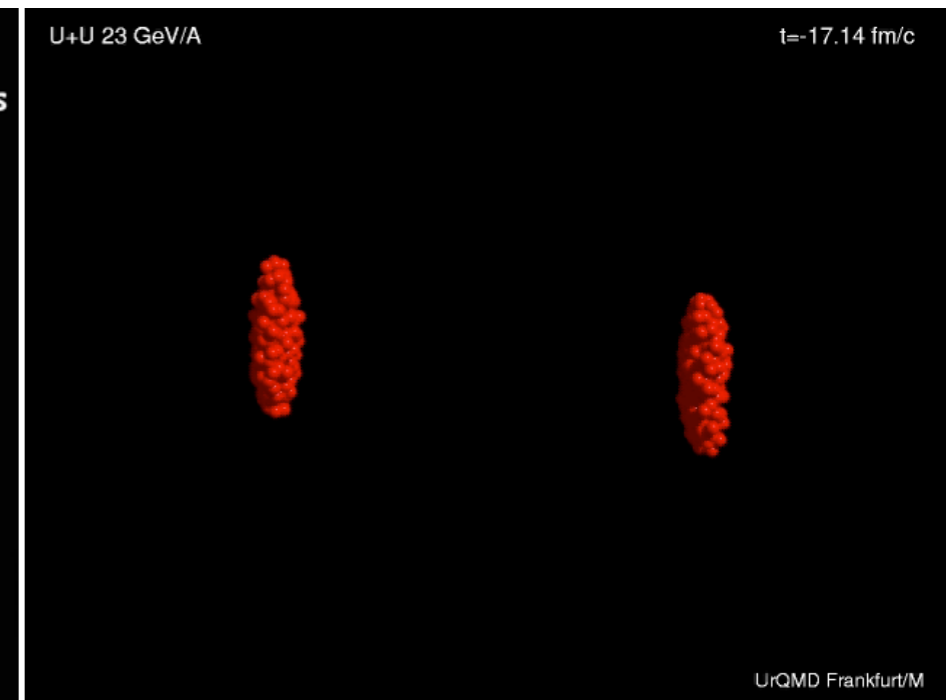
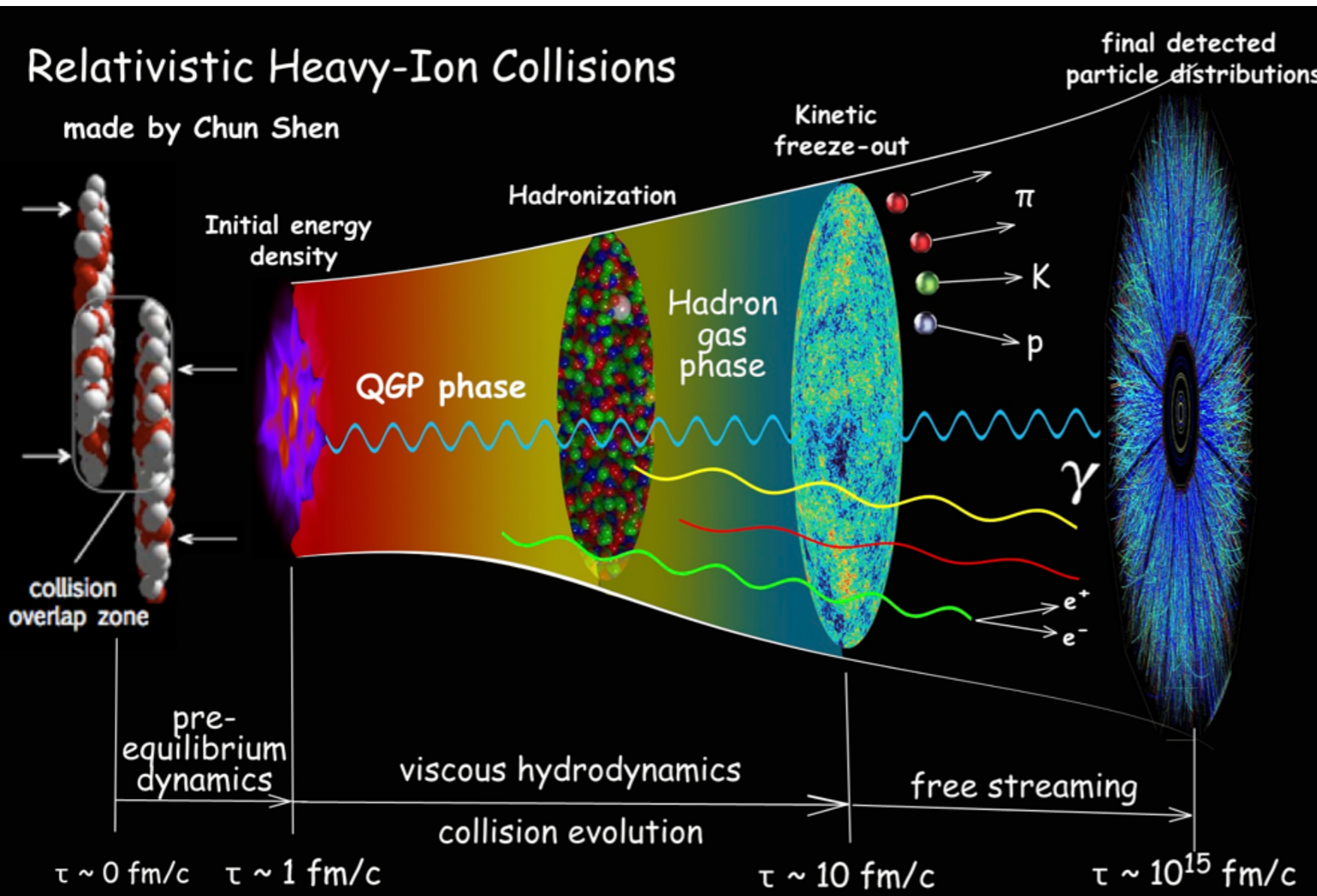
- **Phase structure of QCD**

- Perturbative QCD & asymptotic freedom
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# Heavy ion collisions

## 'Phases/EPOCHs' of a heavy ion collision & time scales

### Simulation of a heavy ion collision



UrQMD Frankfurt/M

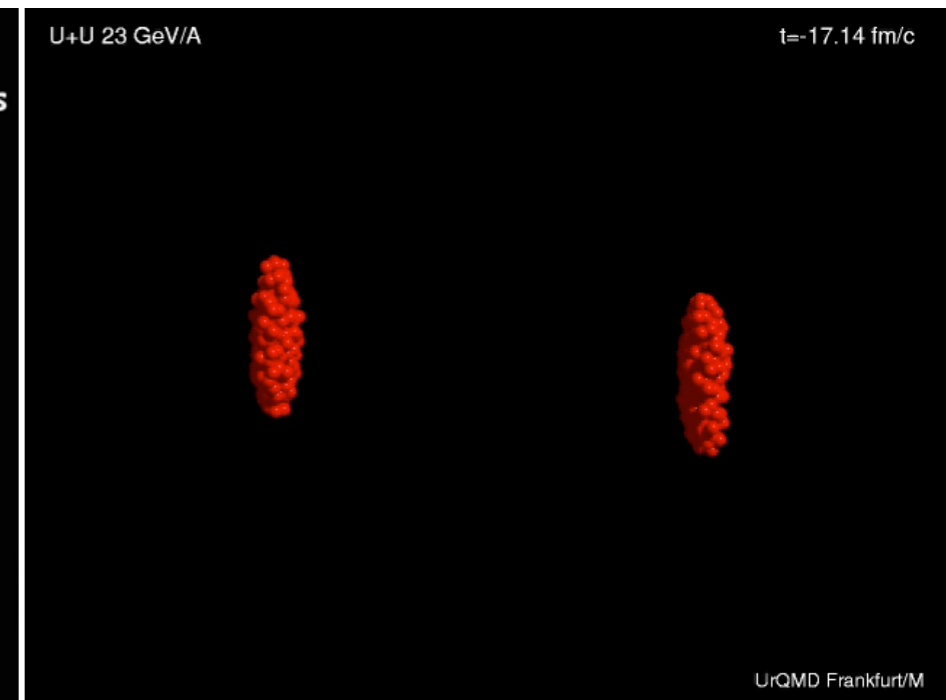
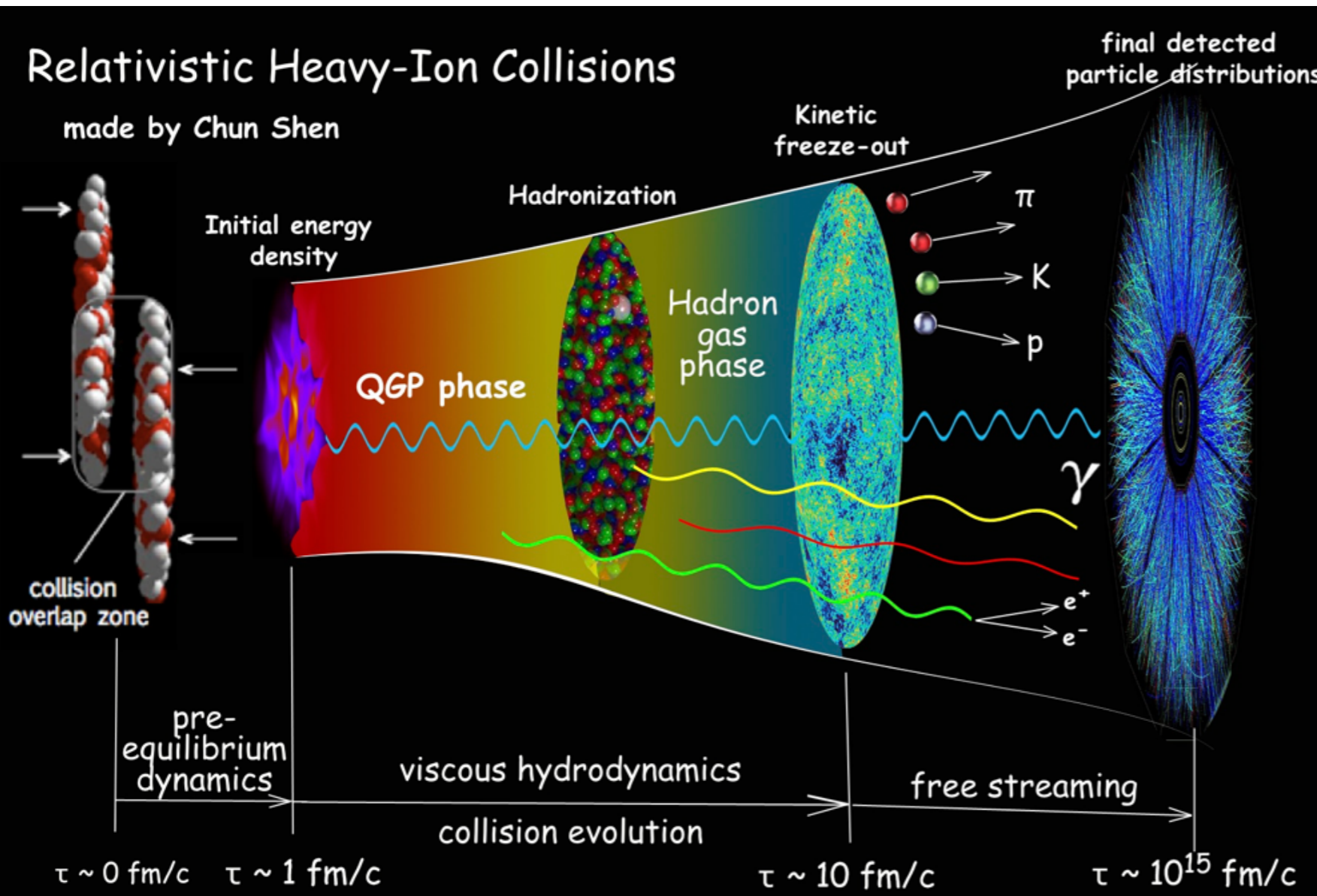
$$1 \text{ fm}/c \sim 3 \times 10^{-24} \text{ seconds}$$



# Heavy ion collisions

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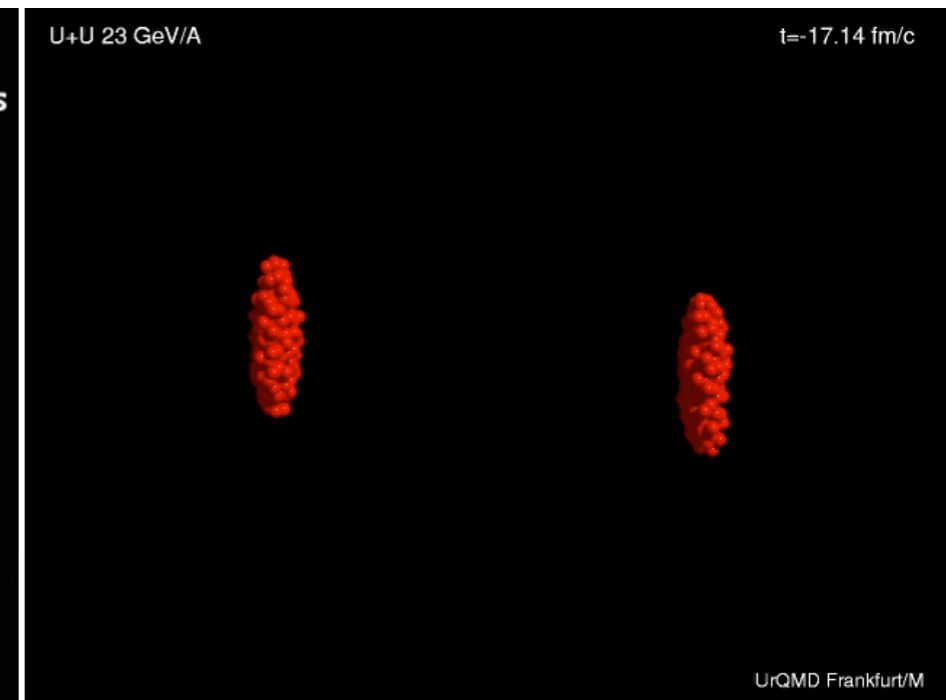
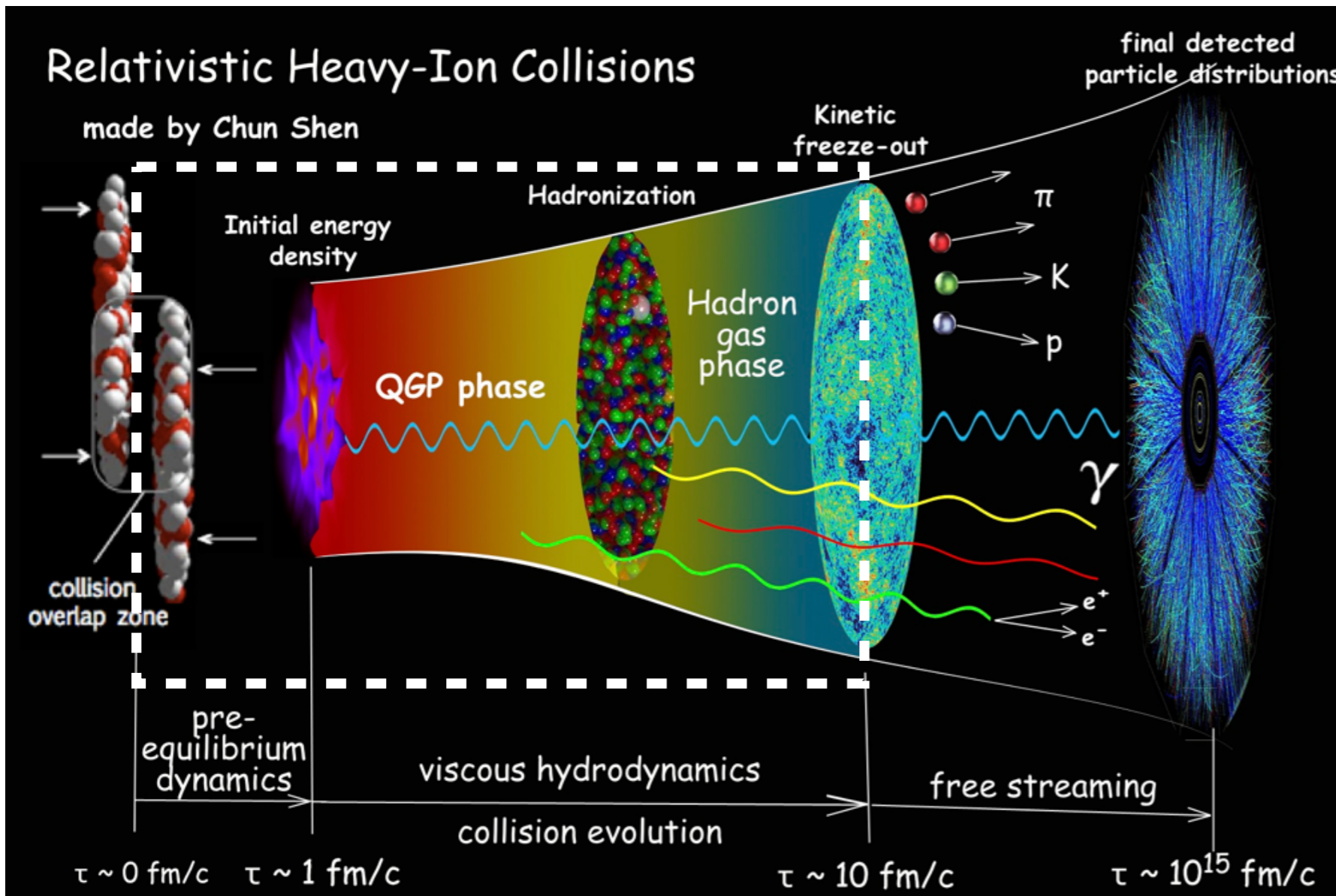
UrQMD Frankfurt/M

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# Heavy ion collisions

## 'Phases/EPOCHS' of a heavy ion collision & time scales

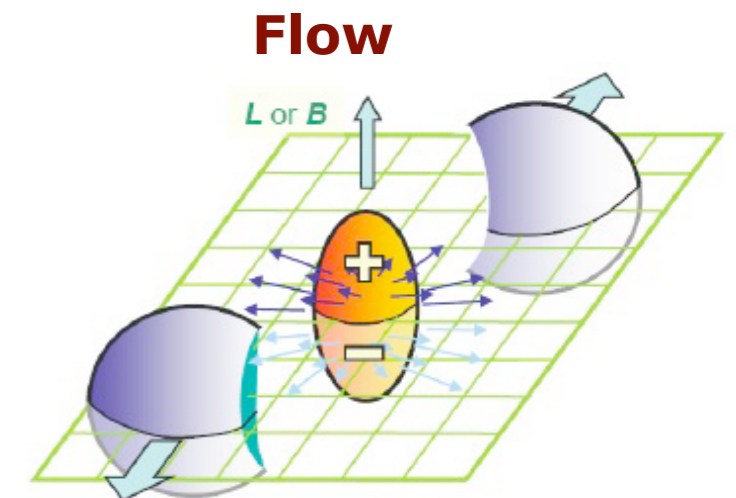
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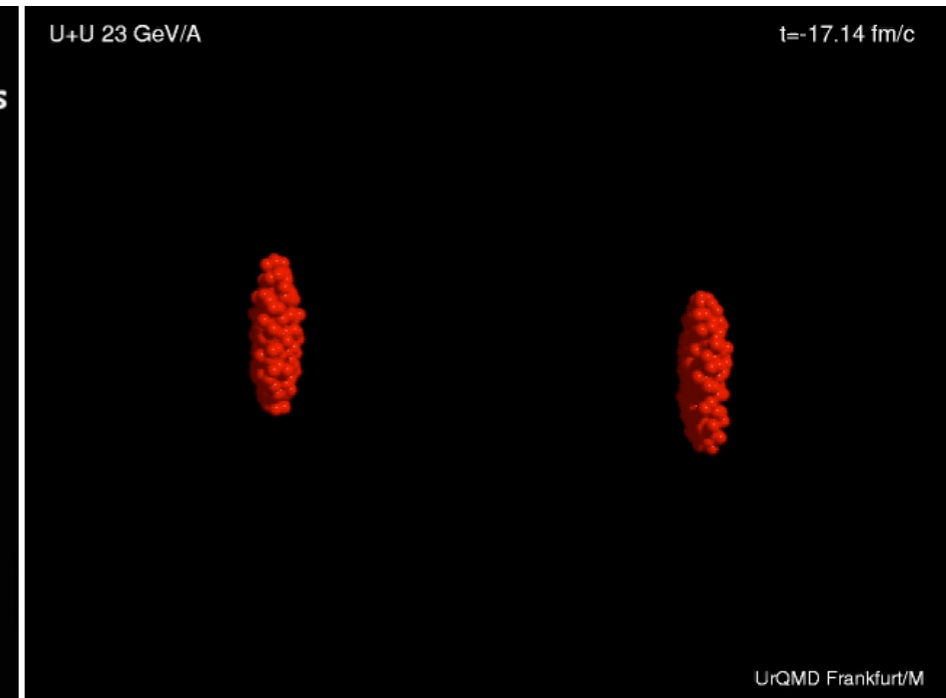
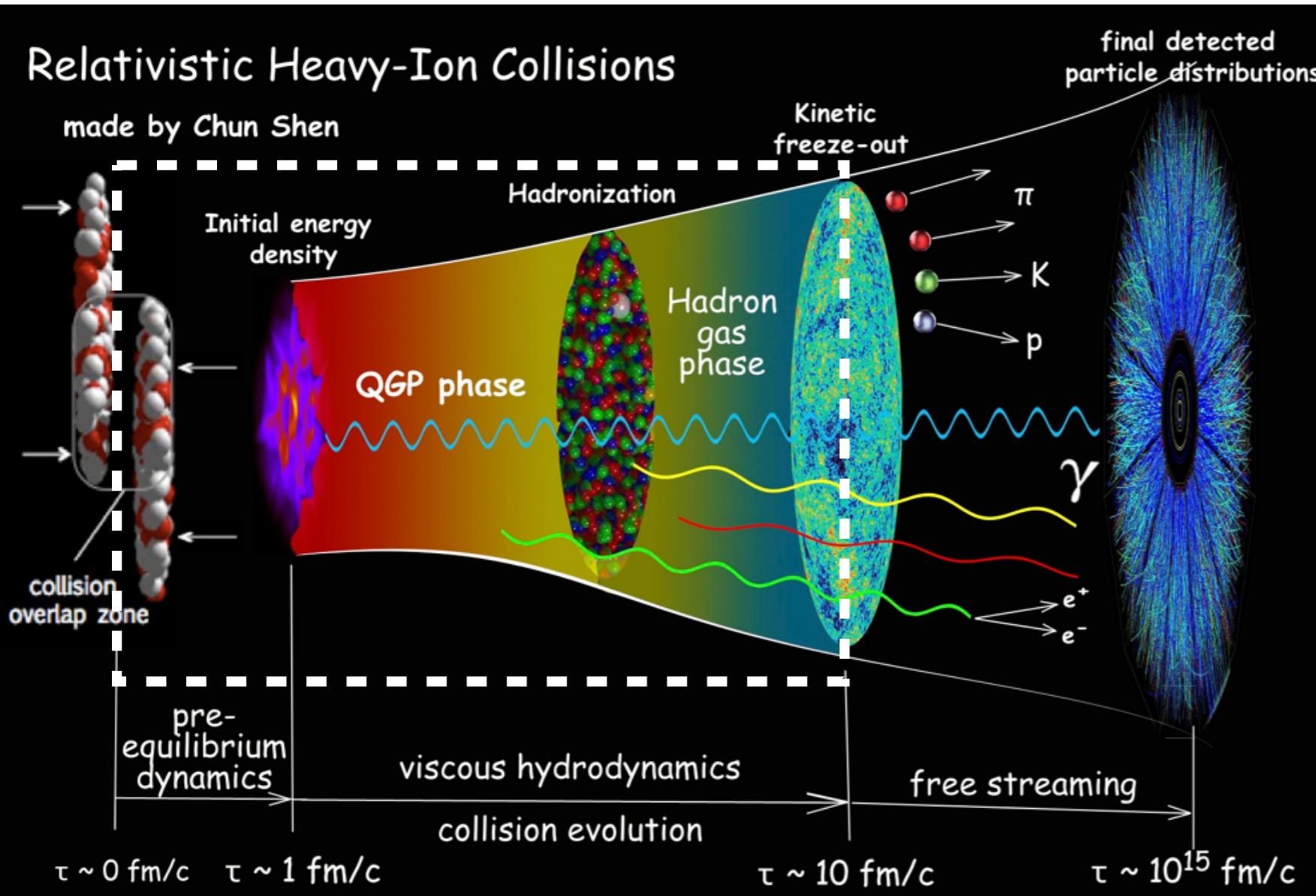


$$\frac{dN}{d(\varphi - \Psi_R)} = \frac{N_0}{2\pi} \left( 1 + 2 \sum_n v_n \cos[n(\varphi - \Psi_R)] \right)$$

# Heavy ion collisions

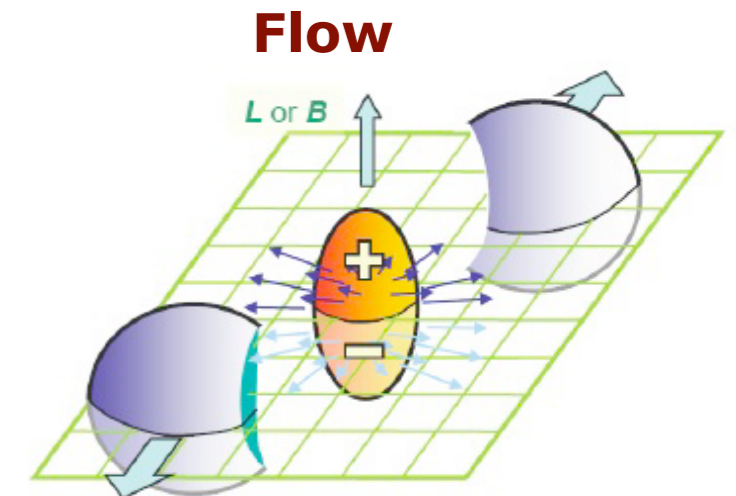
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### Simulation of a heavy ion collision



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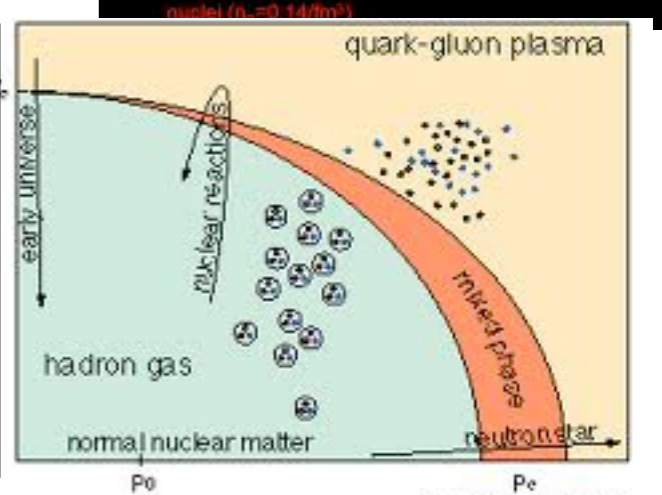
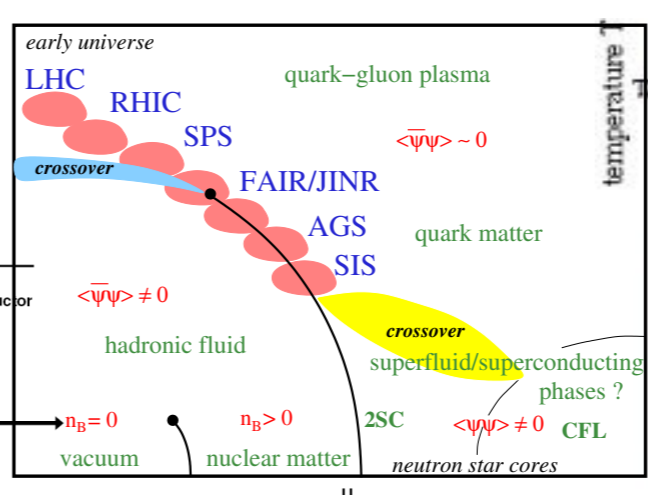
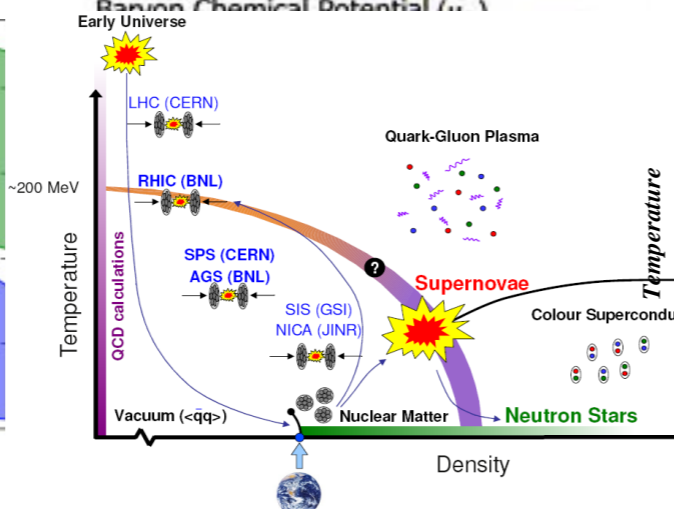
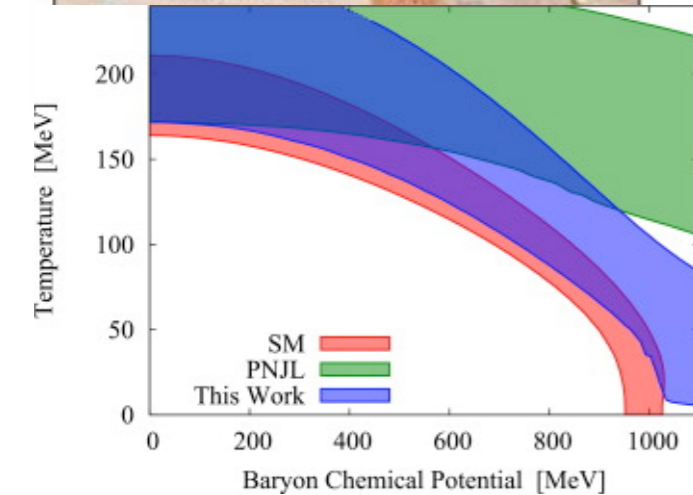
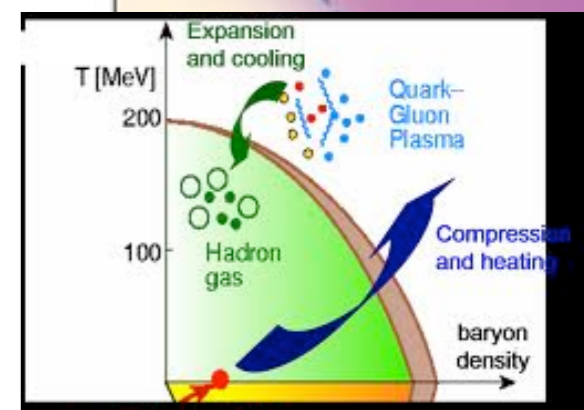
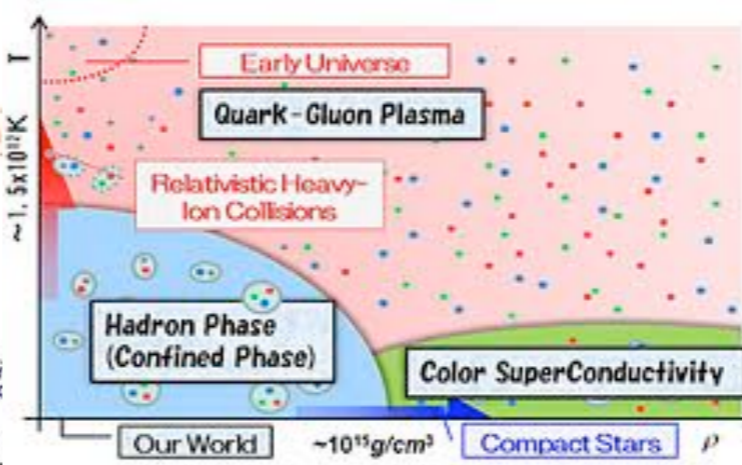
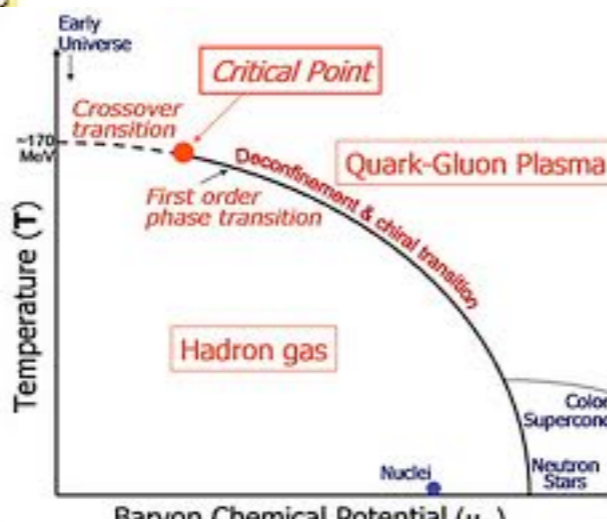
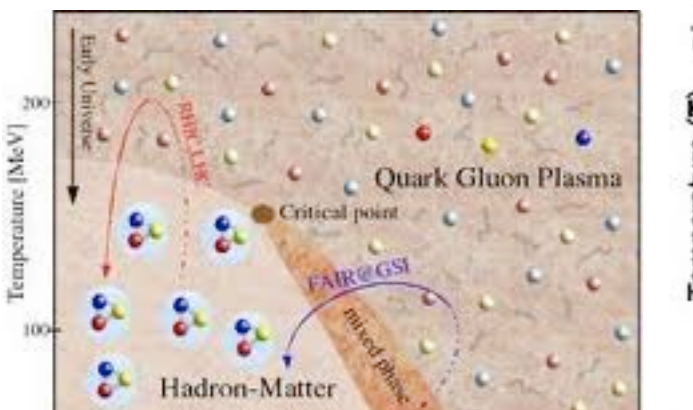
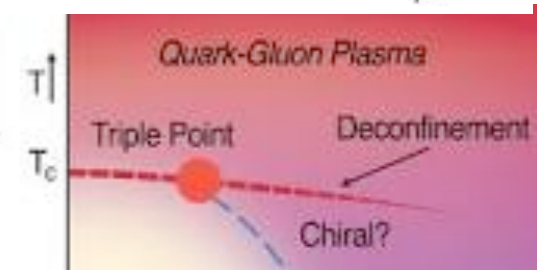
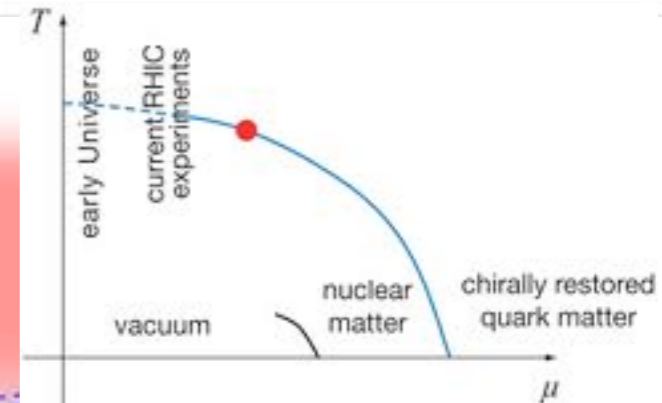
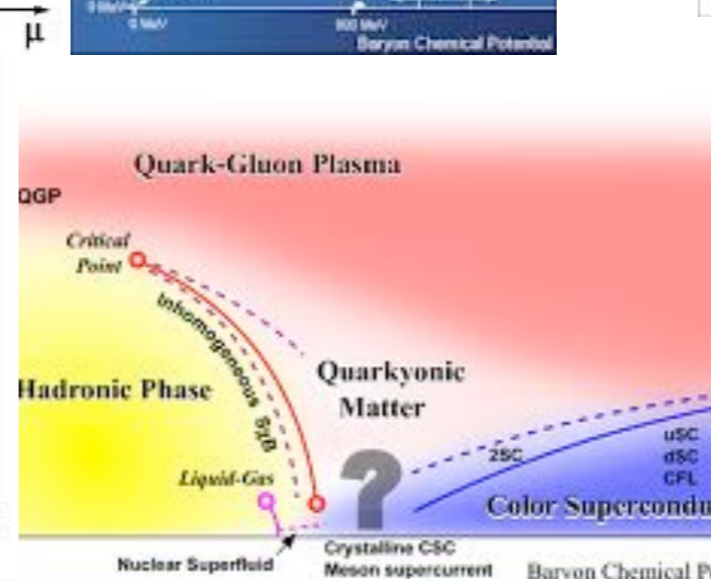
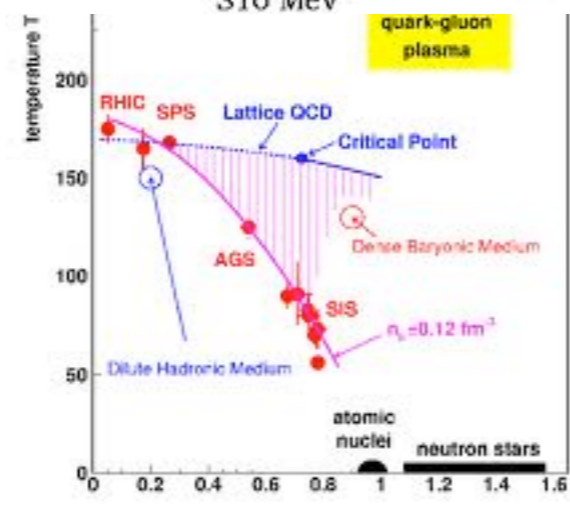
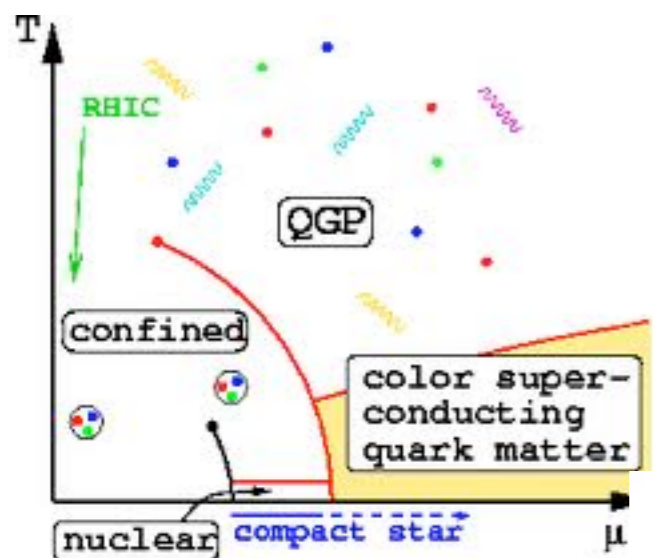
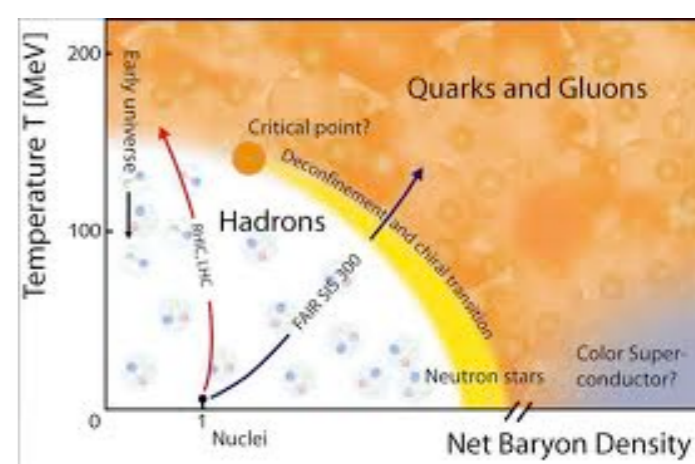
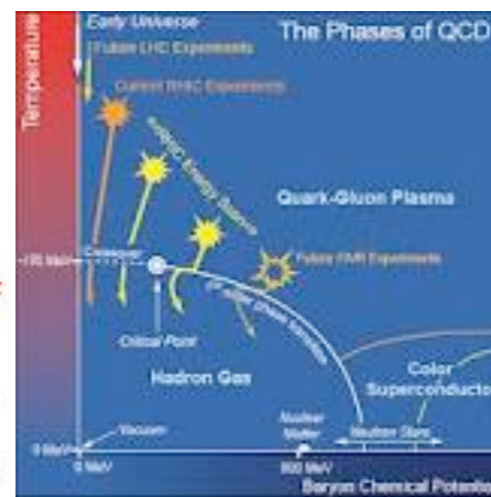
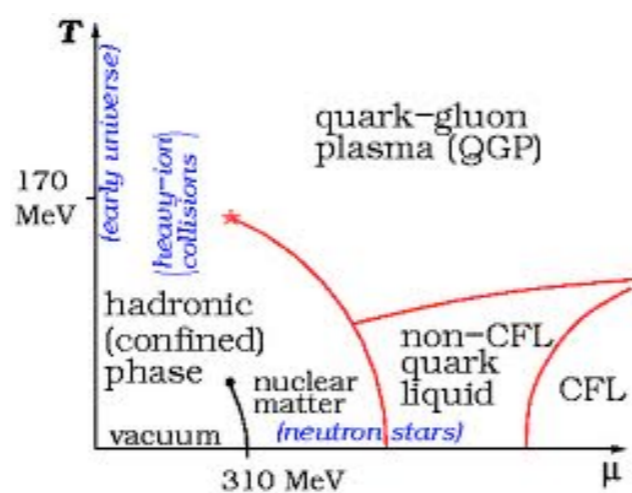
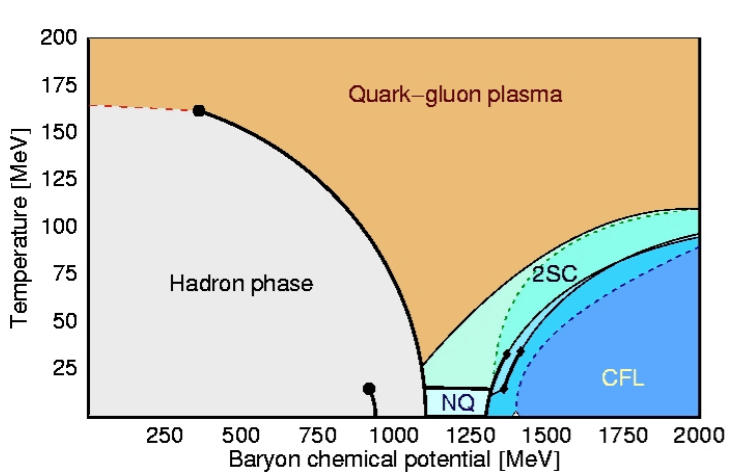
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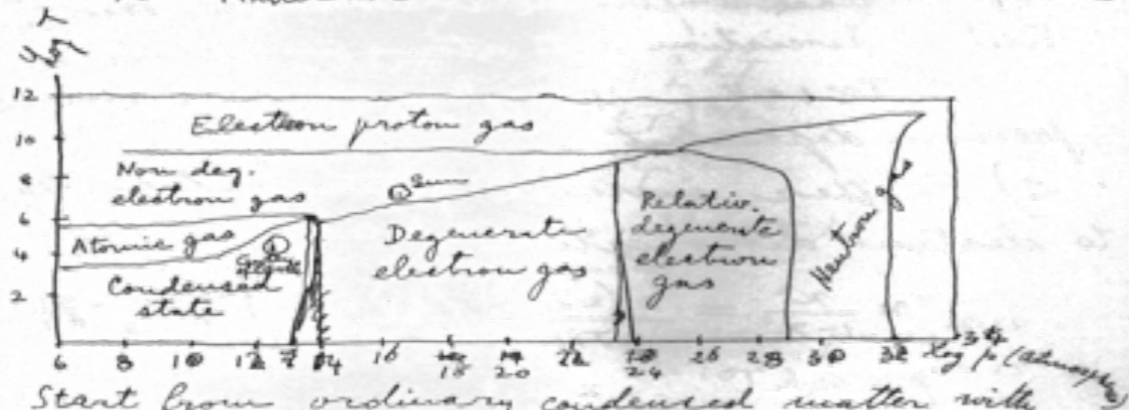


$$1 \text{ fm}/c \sim 3 \times 10^{-24} \text{ seconds}$$

**QCD:**  $v_n$

$$\frac{dN}{d(\varphi - \Psi_R)} = \frac{N_0}{2\pi} \left( 1 + 2 \sum_n v_n \cos[n(\varphi - \Psi_R)] \right)$$





Start from ordinary condensed matter with Fermi equation of state controlled by ordinary chemical forces.

a) Increase pressure at  $T < 1000$  until deg. electron energies exceeds 20 eV —

Condition  $\bar{w} = \frac{3}{40} \left(\frac{6}{\pi}\right)^{2/3} \frac{h^2 n^{2/3}}{2^{1/3} m}$       $p = \frac{2}{3} \bar{w} n$

$\bar{w} = 3.6 \times 10^{-27} n^{2/3} = 3.2 \times 10^{-11}$

$n \approx 10^{24}$       $p = \frac{2}{3} 3.2 \times 10^{-11} \times 10^{24} \approx 2 \times 10^{13} \approx 2 \times 10^7 \text{ atm}$

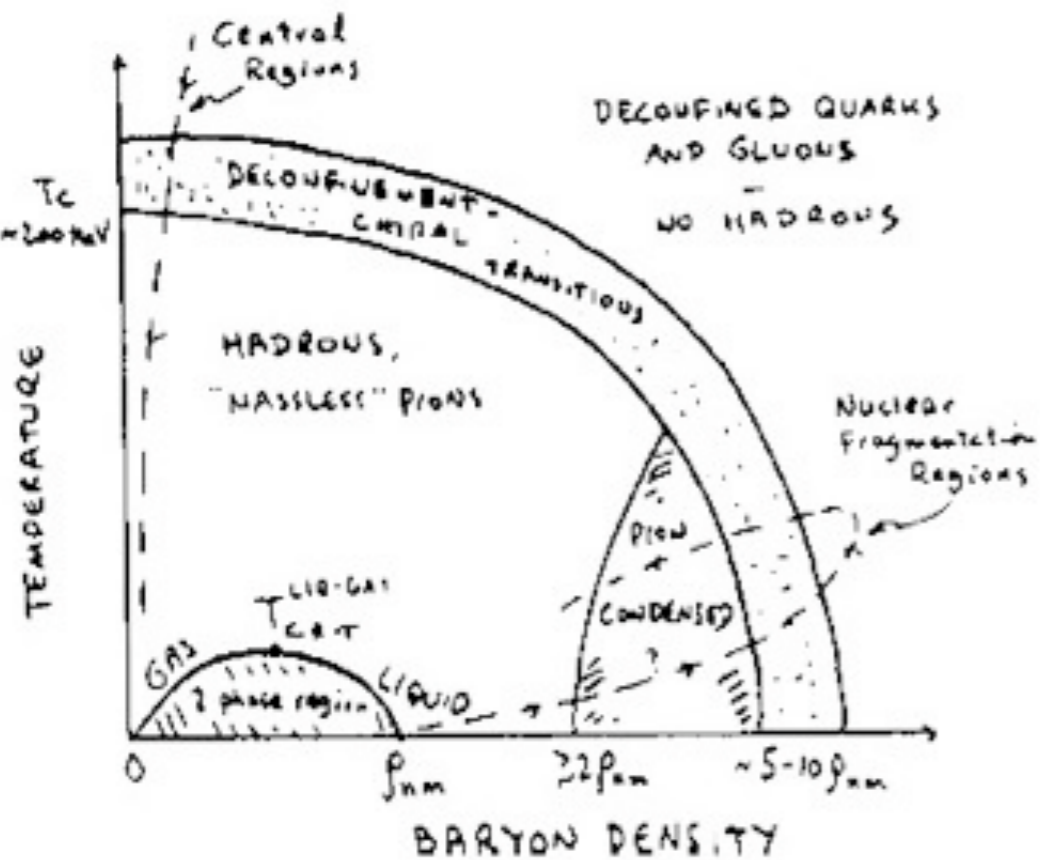
As pressure increases beyond this point

$p = 3.6 \times 10^{-27} n^{2/3} \Rightarrow n \times \frac{2}{3} = 2.4 \times 10^{-27} n^{5/3}$

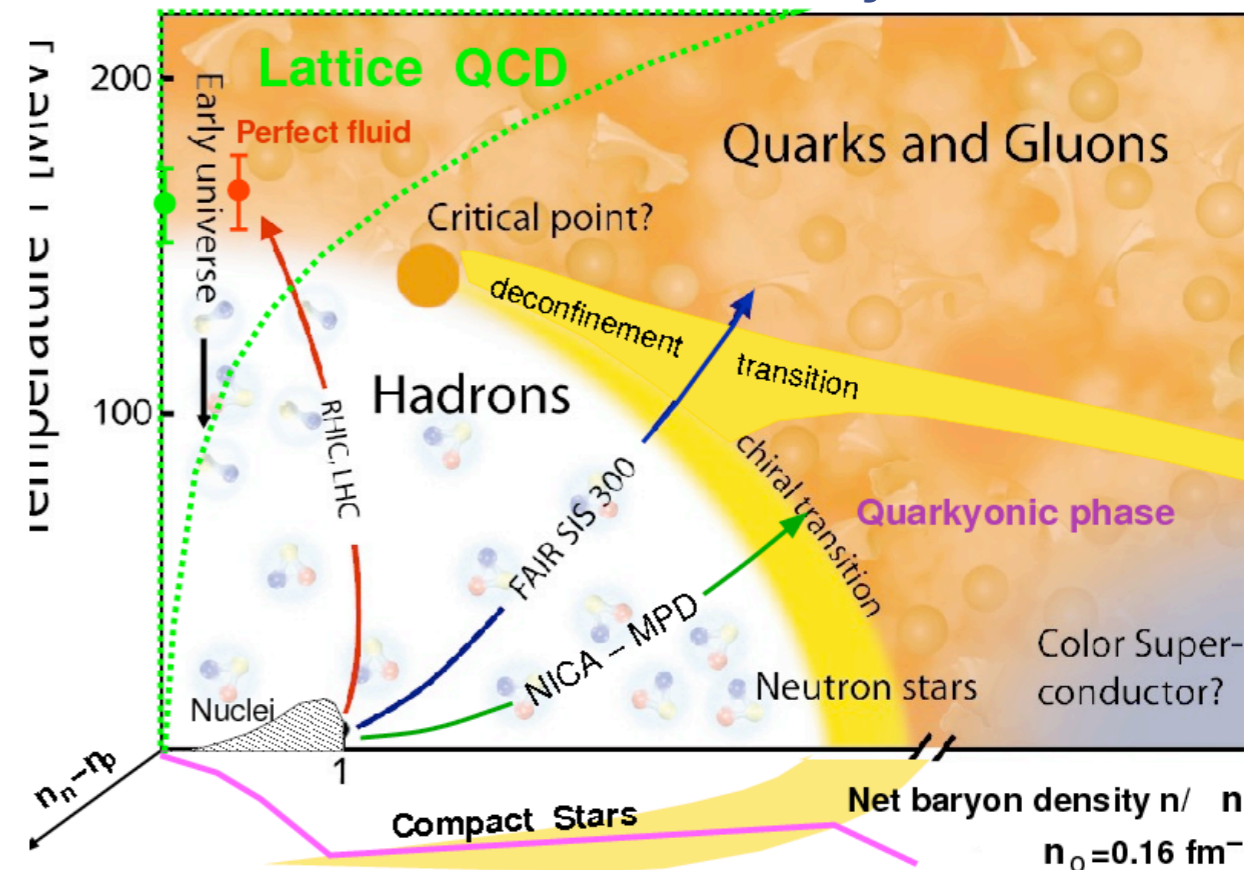
$n = 6 \times 10^{23} \frac{p}{\Lambda^3}$       $p = 10^{12.01} \left(\frac{p}{\Lambda}\right)^{5/3} \approx 2.2 \times 10^{12} p^{5/3}$

1953 Enrico Fermi

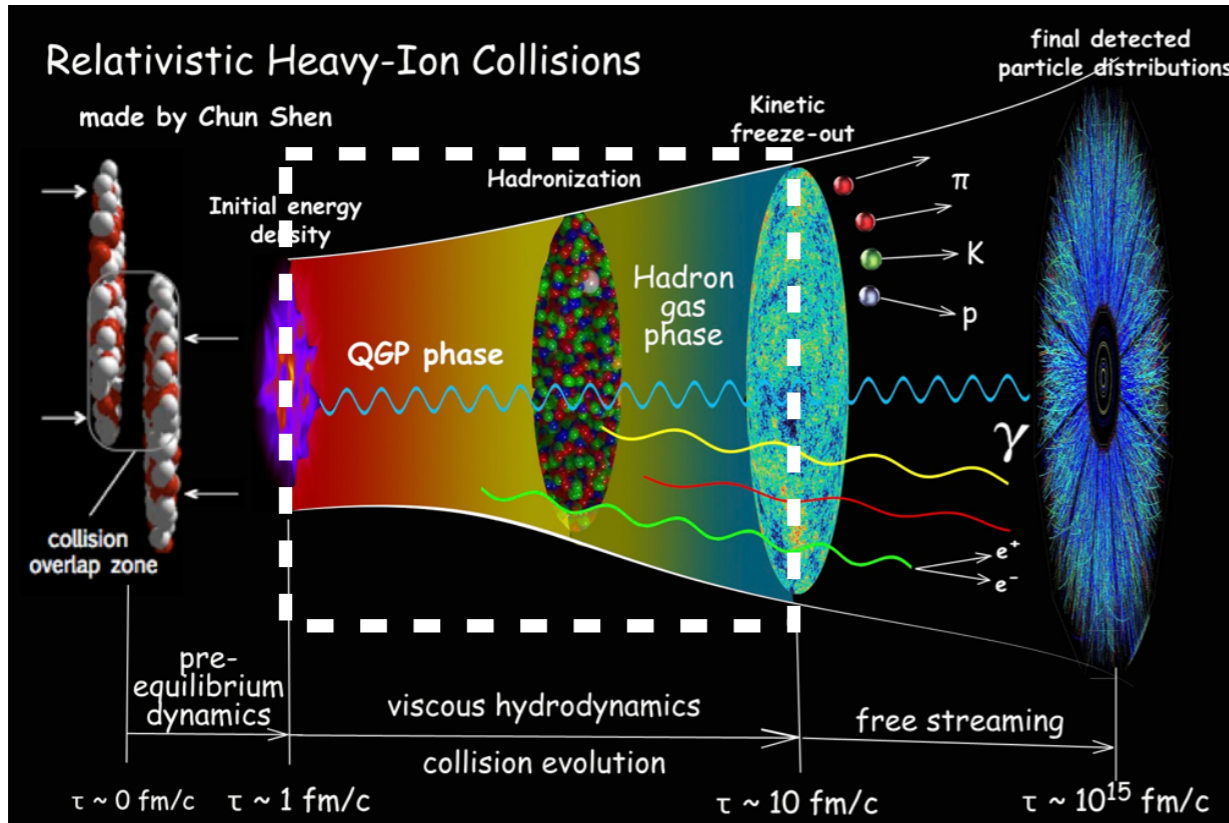
1983 US long range plan, Gordon Baym



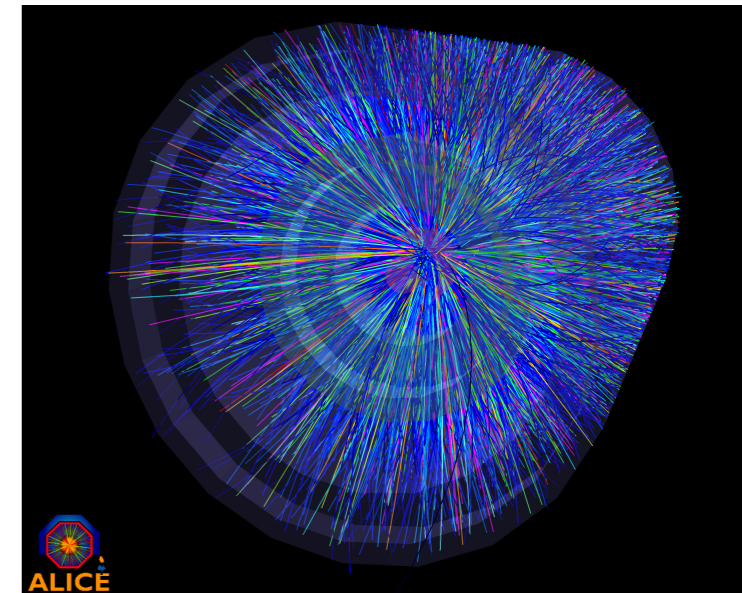
Larry McLerran '09



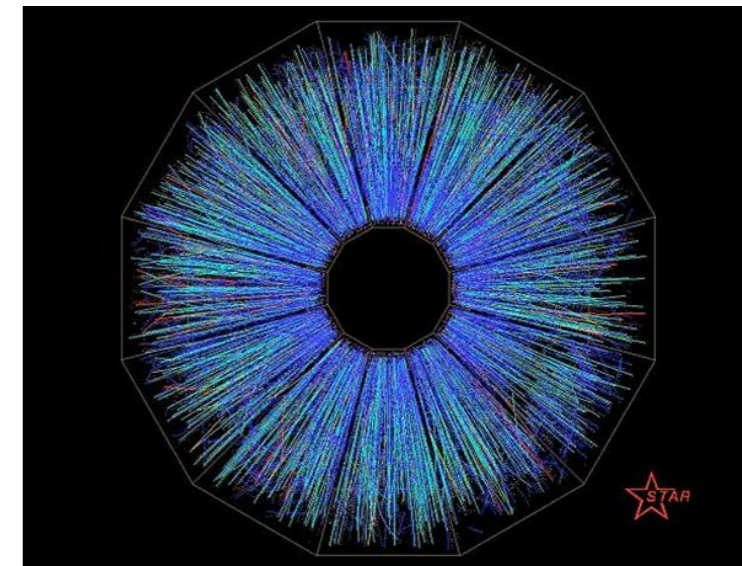
# Heavy ion collisions



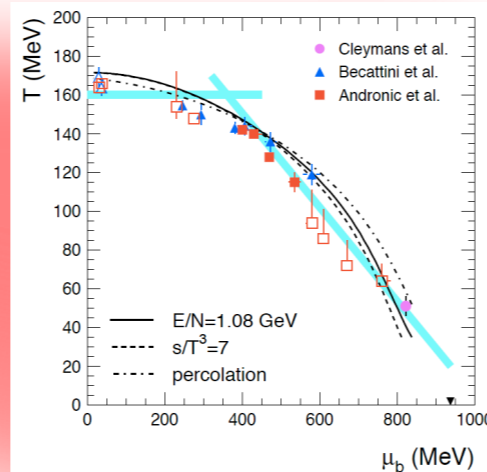
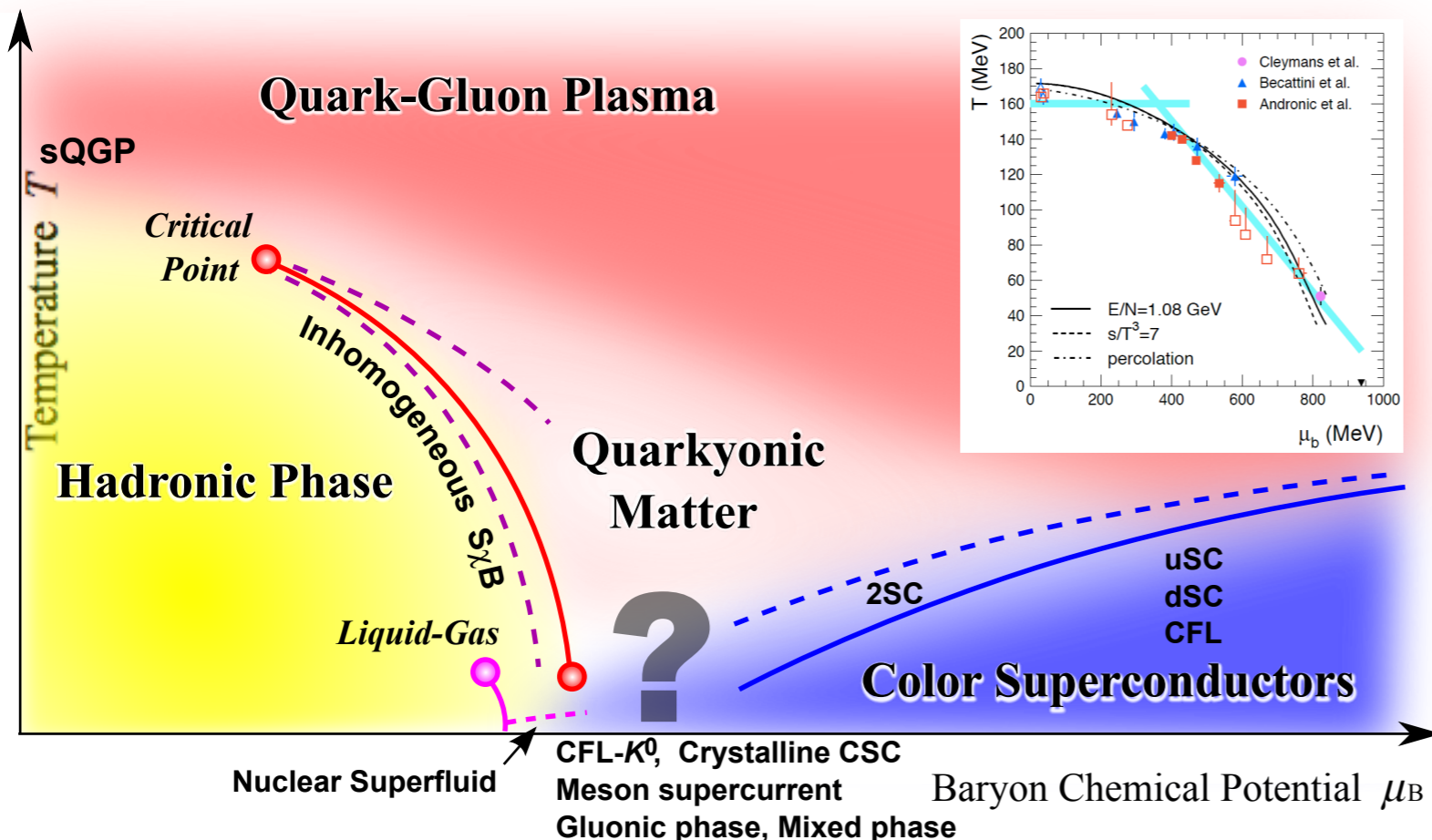
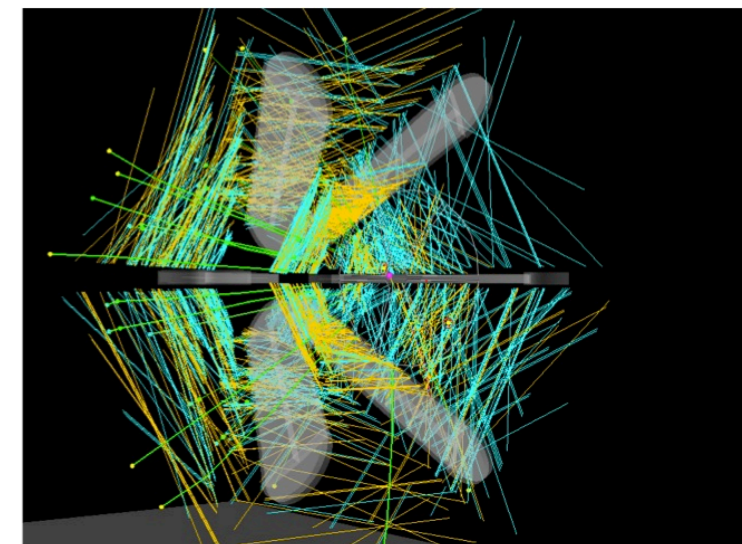
**LHC**



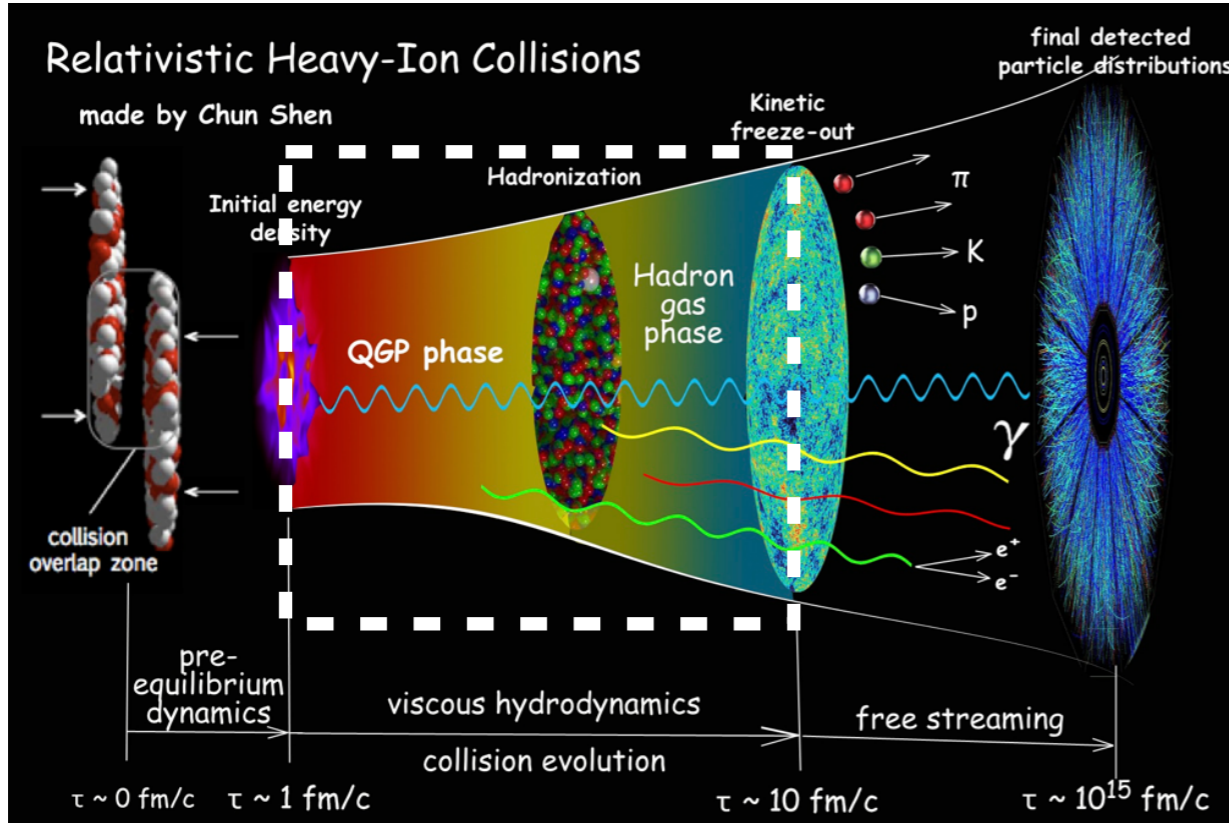
**RHIC**



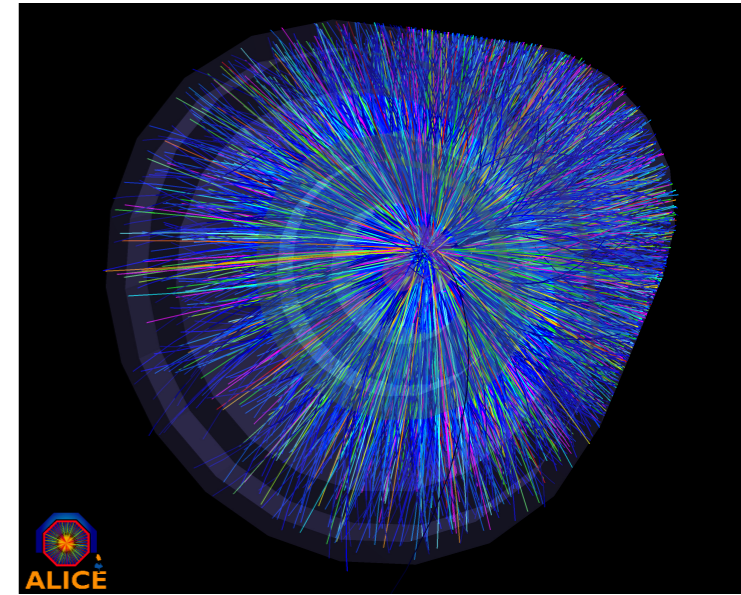
**HADES**



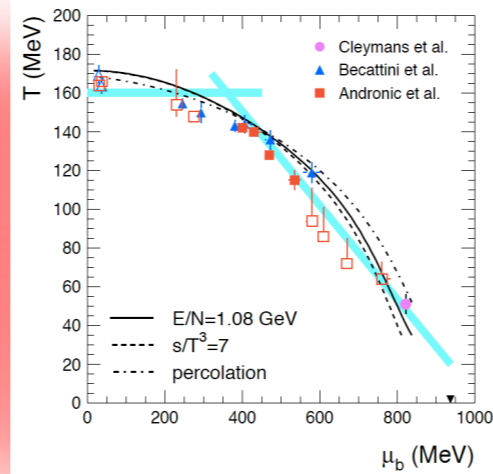
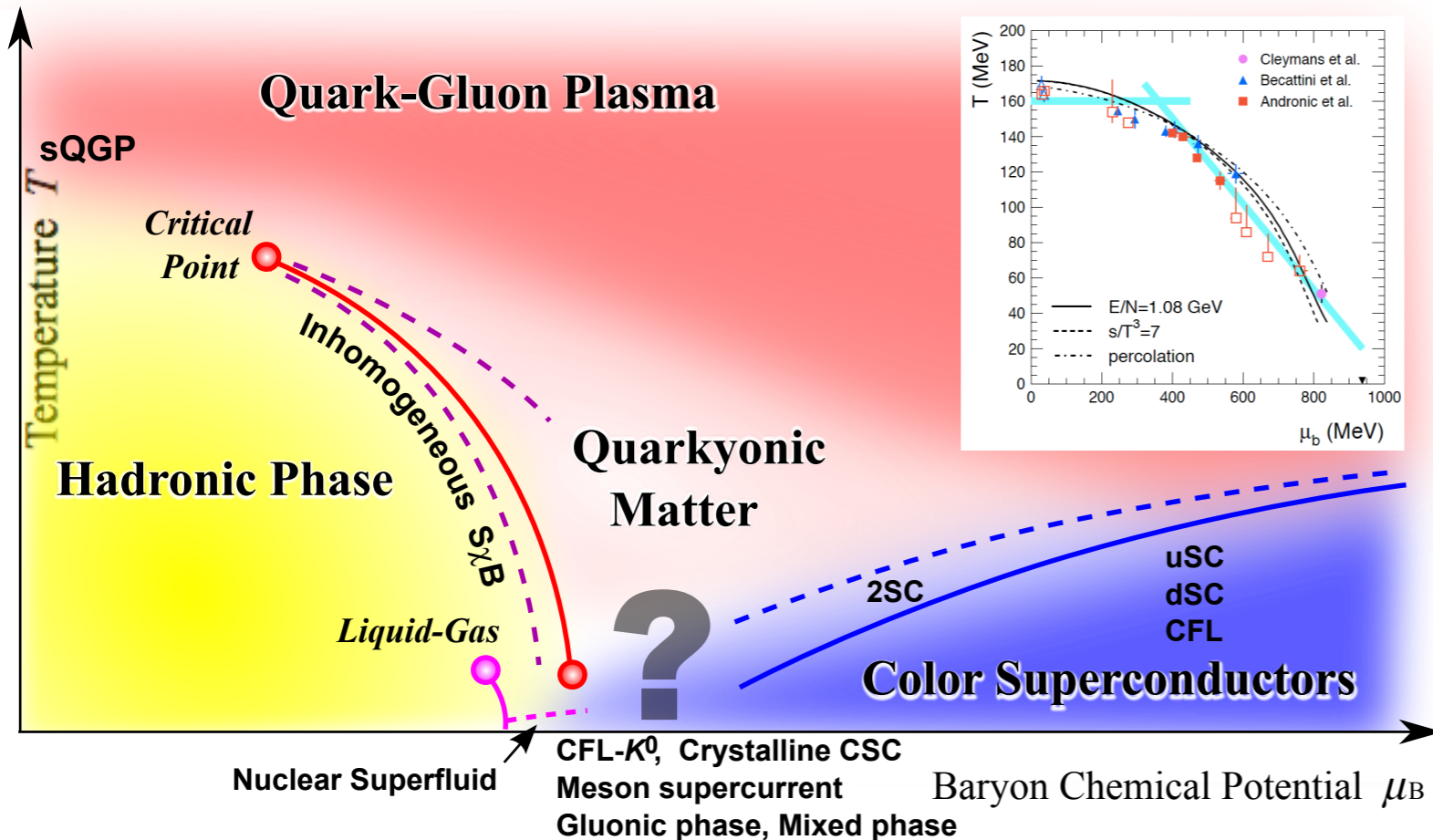
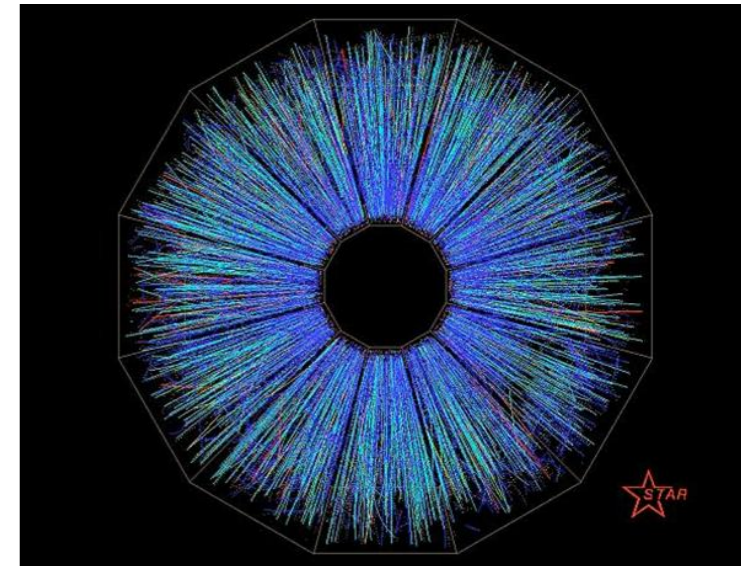
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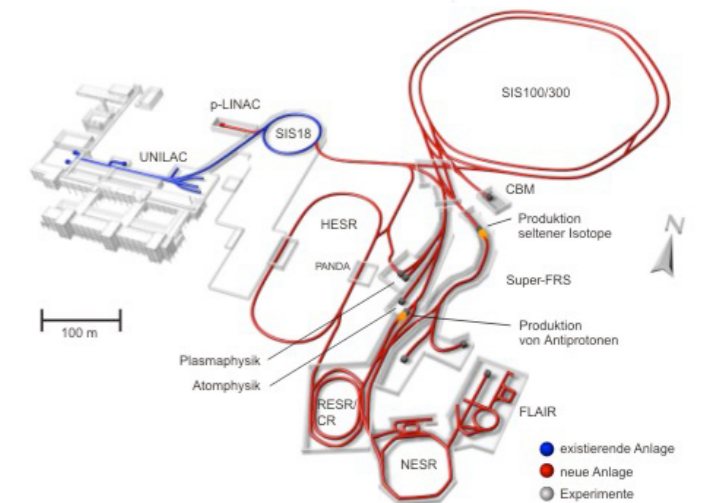
**LHC**



**RHIC**

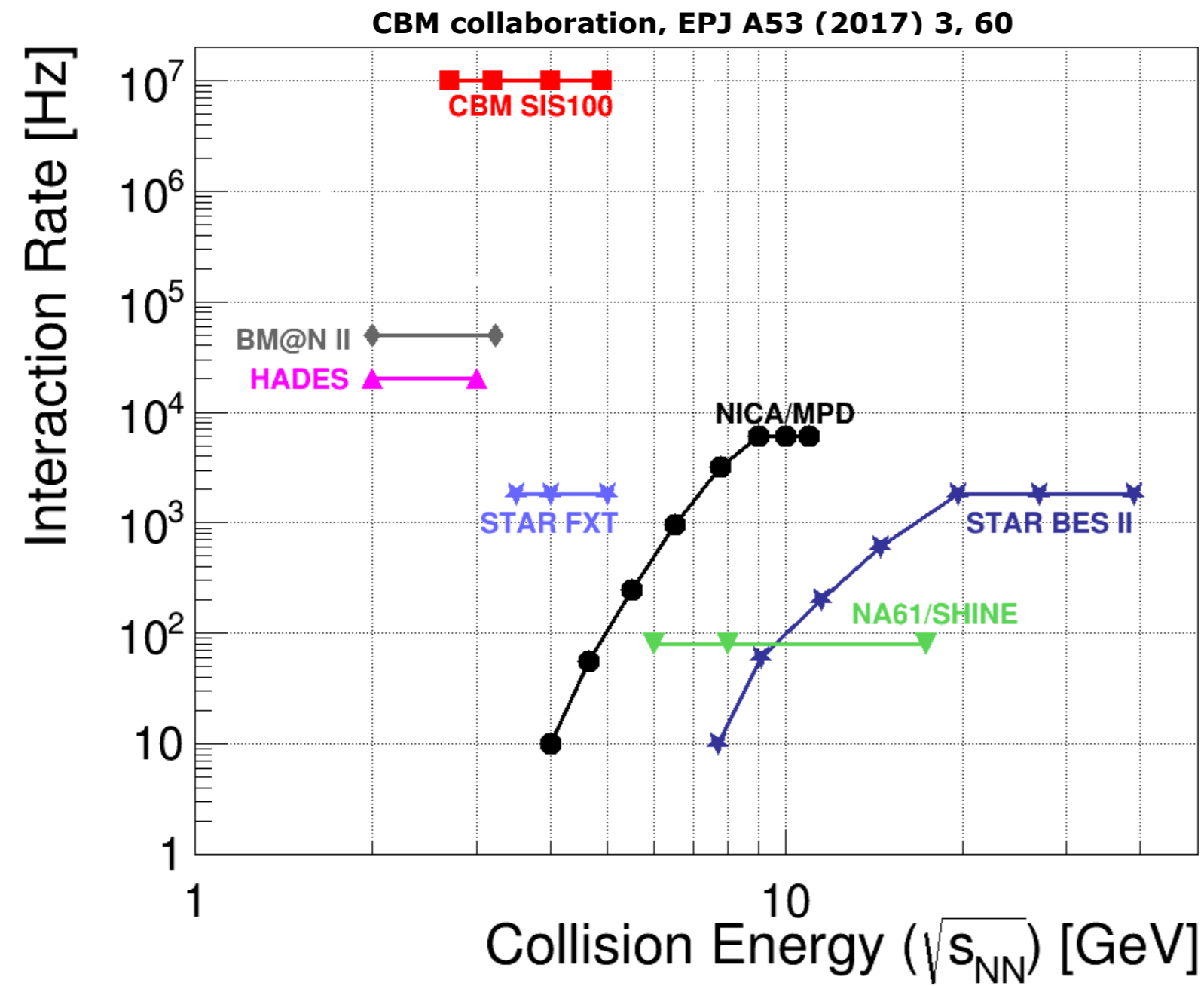


**FAIR**

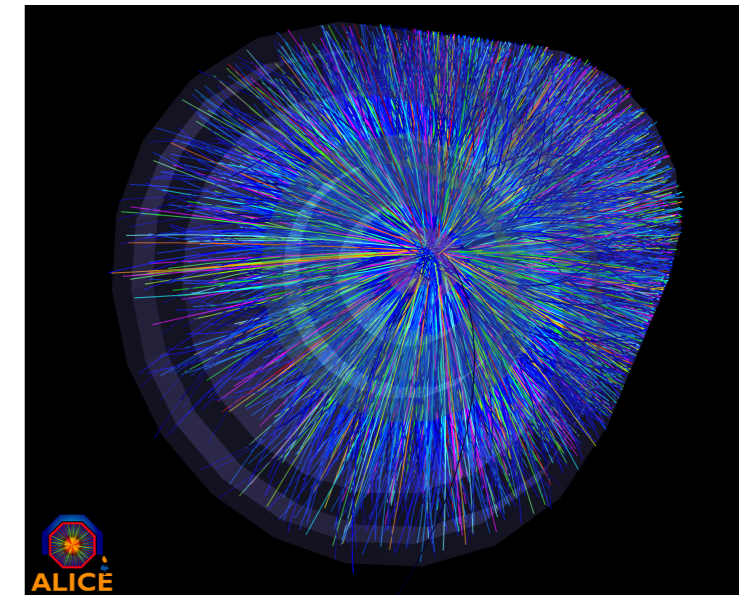


**NICA**

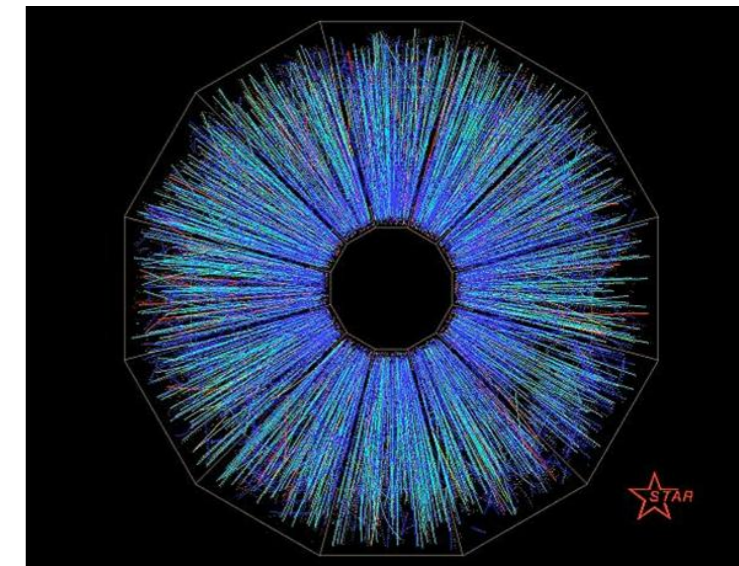
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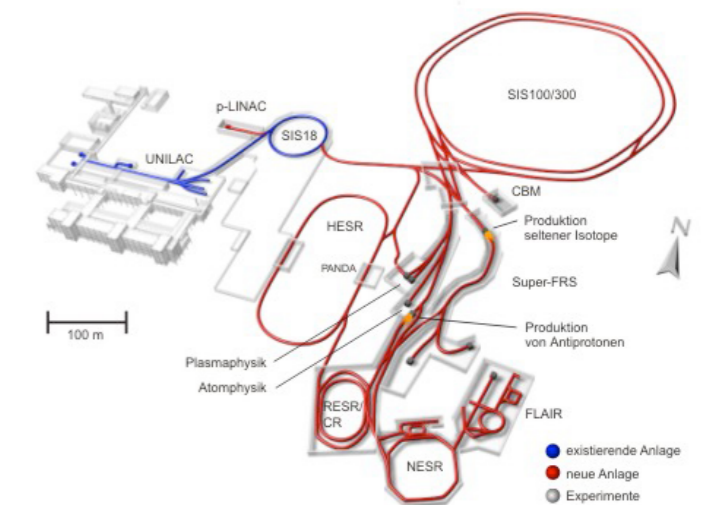
LHC



RHIC



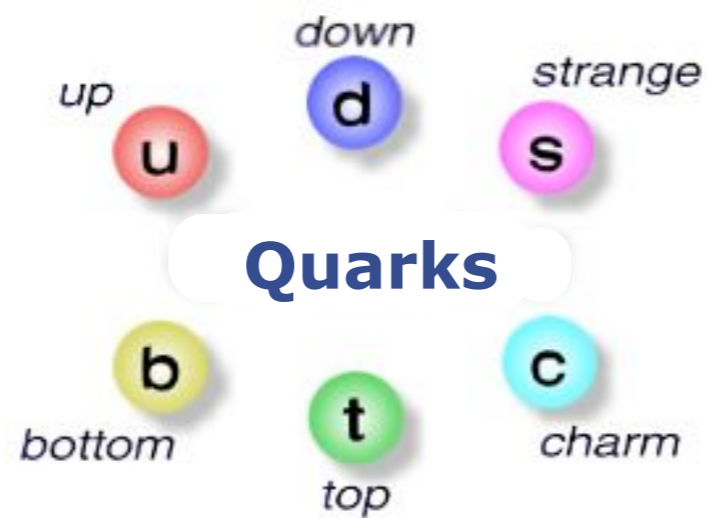
FAIR



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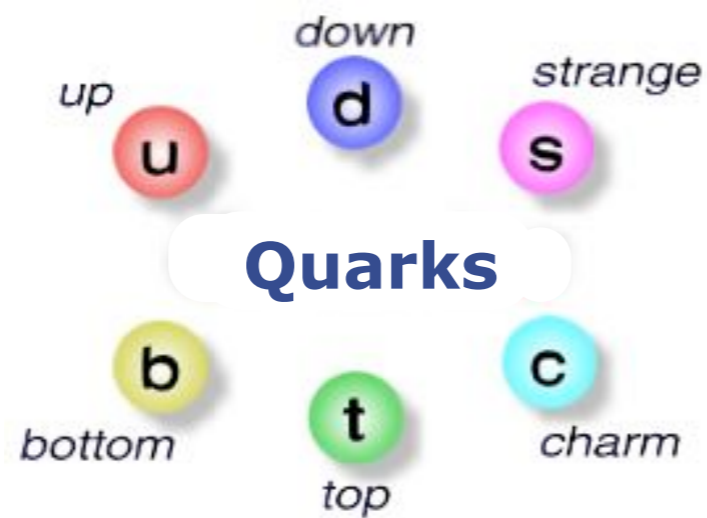


# Phase structure of QCD



**Gluons**

# Perturbative QCD & asymptotic freedom



**Gluons**

# QCD, asymptotic freedom and all that

## Action and interactions

### QCD action $S_{\text{QCD}}$

Yang-Mills

gauge fixing

$$\frac{1}{4} \int_x F_{\mu\nu}^a F_{\mu\nu}^a + \frac{1}{2\xi} \int_x (\partial_\mu A_\mu^a)^2 + \int_x \bar{c}^a \partial_\mu D_\mu^{ab} c^b + \int_x \bar{q} \cdot (i\not{D} + i m_\psi + i\mu\gamma_0) \cdot q$$

gluon

ghost

quarks

Pure gauge theory

matter sector

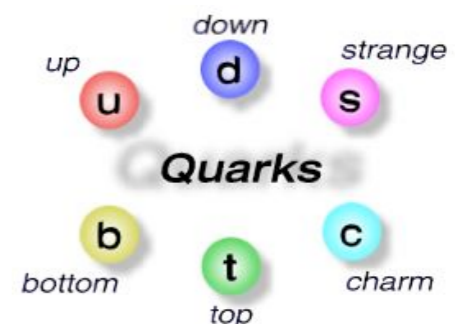
$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g f^{abc} A_\mu^b A_\nu^c$$

$$\not{D} = \gamma_\mu D_\mu$$

$$a, b, c = 1, \dots, N_c^2 - 1$$



$$N_f = 6$$



$$D_\mu(A) = \partial_\mu - i g A_\mu$$

# QCD, asymptotic freedom and all that

## Action and interactions

### QCD action $S_{\text{QCD}}$

#### Yang-Mills

#### gauge fixing

$$\frac{1}{4} \int_x F_{\mu\nu}^a F_{\mu\nu}^a + \frac{1}{2\xi} \int_x (\partial_\mu A_\mu^a)^2 + \int_x \bar{c}^a \partial_\mu D_\mu^{ab} c^b + \int_x \bar{q} \cdot (i\not{D} + i m_\psi + i\mu\gamma_0) \cdot q$$

**gluon**                      **ghost**                      **quarks**

#### Pure gauge theory

#### matter sector

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g f^{abc} A_\mu^b A_\nu^c$$

$$\not{D} = \gamma_\mu D_\mu$$

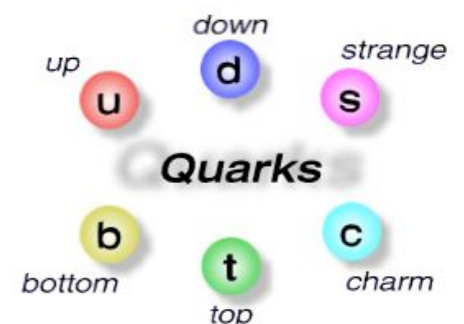
$$a, b, c = 1, \dots, N_c^2 - 1$$



**covariant derivative in adjoint representation**

$$D_\mu^{ab}(A) = \partial_\mu \delta^{ab} - g f^{abc} A_\mu^c$$

$$N_f = 6$$



# QCD, asymptotic freedom and all that

## Action and interactions

### QCD action $S_{\text{QCD}}$

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$$\frac{1}{4} \int_x F_{\mu\nu}^a F_{\mu\nu}^a + \frac{1}{2\xi} \int_x (\partial_\mu A_\mu^a)^2 + \int_x \bar{c}^a \partial_\mu D_\mu^{ab} c^b + \int_x \bar{q} \cdot (i\not{D} + i m_\psi + i\mu\gamma_0) \cdot q$$

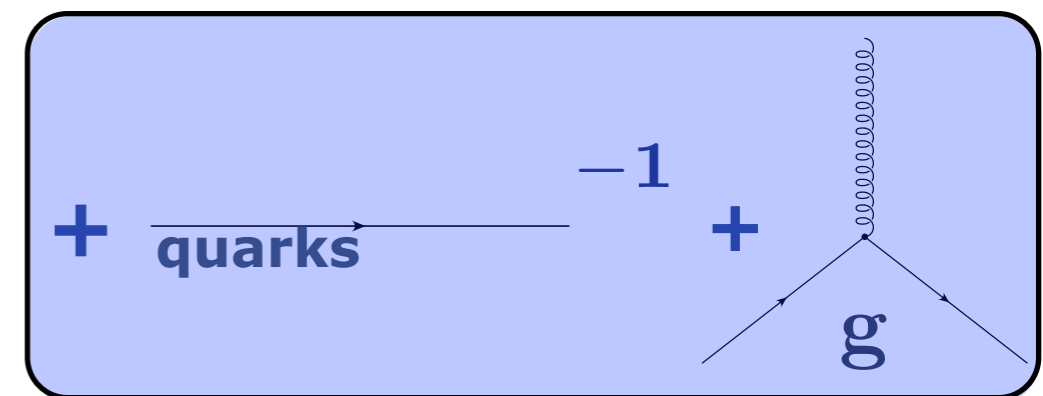
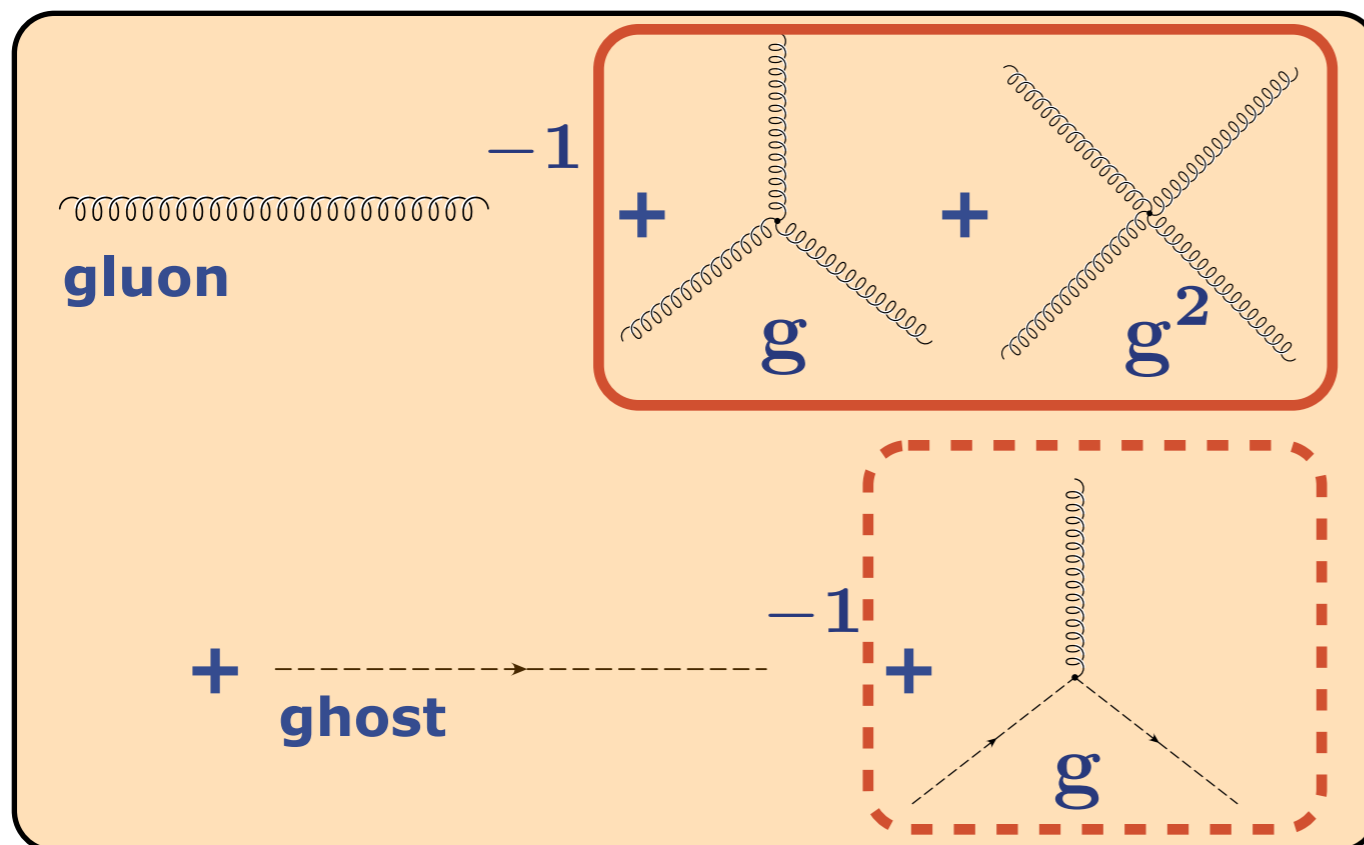
gluon

ghost

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Pure gauge theory

matter sector



parameters

- 1 coupling  $g$
  - mass matrix  $m_\psi$
- $N_f \times N_f$

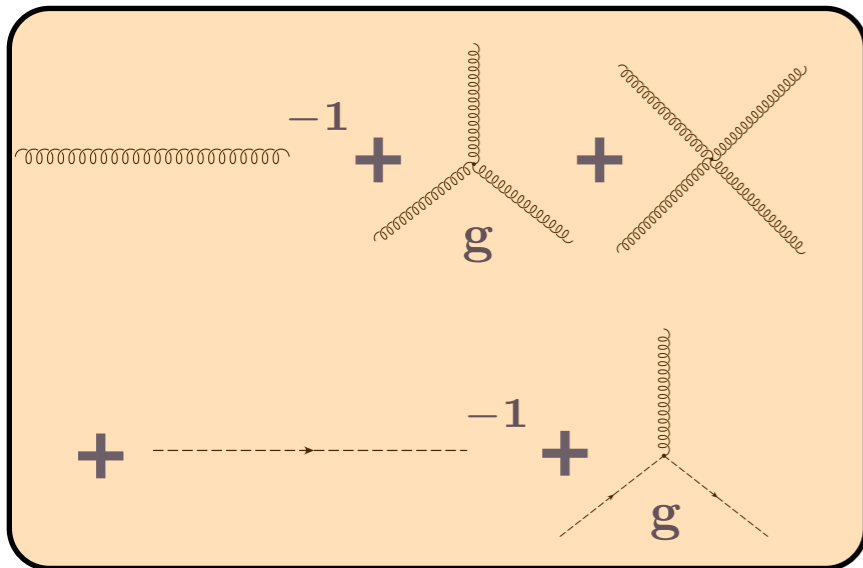
# QCD, asymptotic freedom and all that

## Running coupling at low and high energies

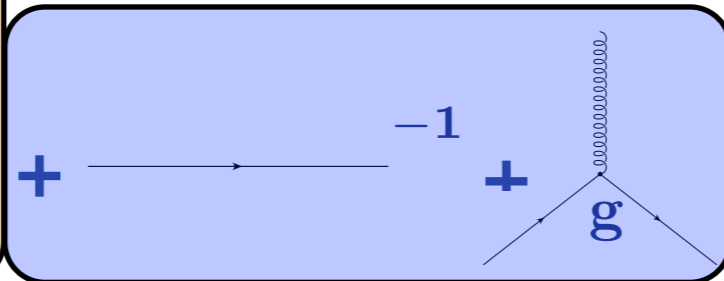
$$\alpha_s(Q) = \frac{g^2(Q)}{4\pi}$$

Millenium Prize 1 Mio \$

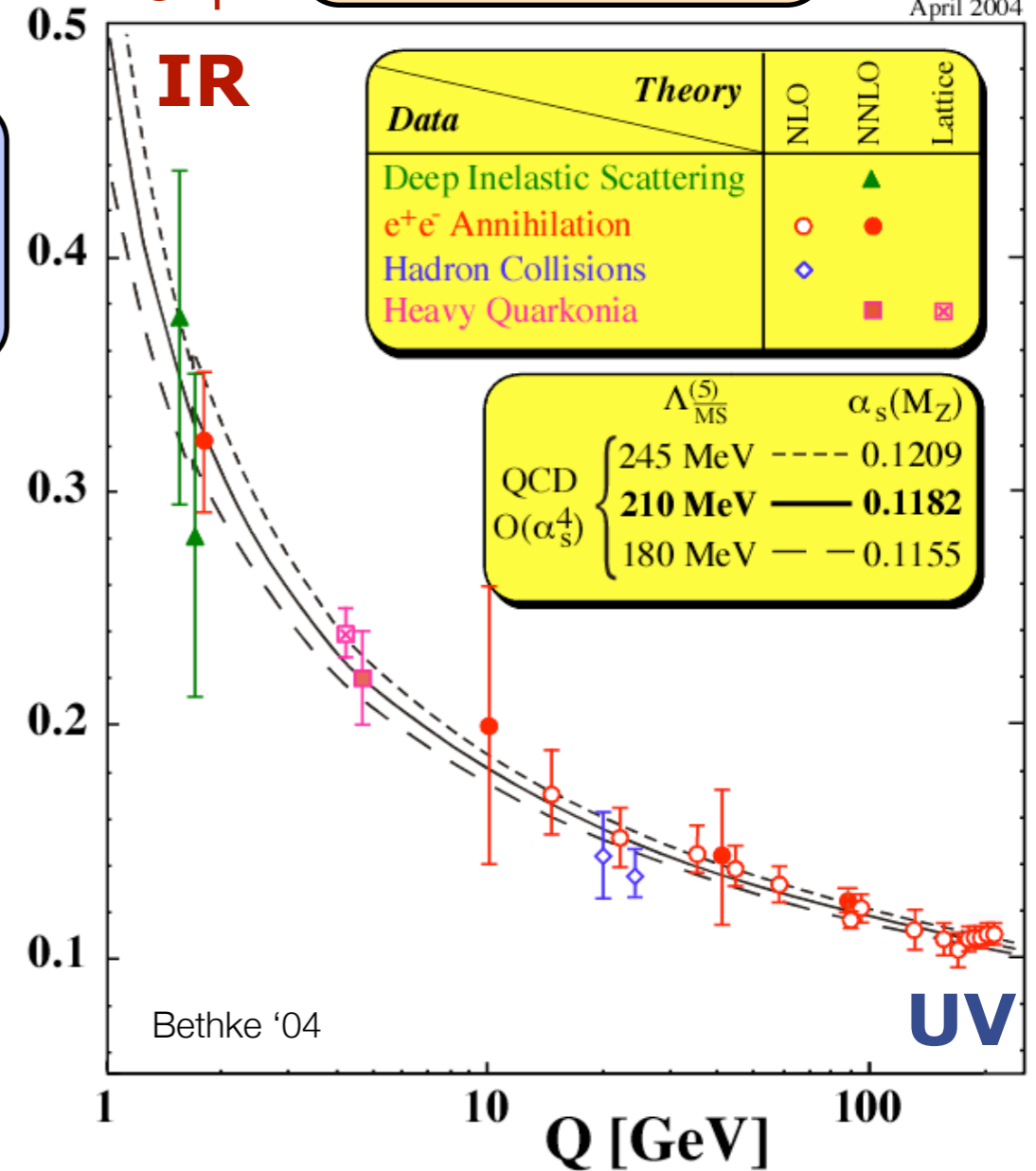
April 2004



Pure gauge theory



matter sector



Bethke '04

Nobel Prize '04

Gross, Politzer, Wilczek

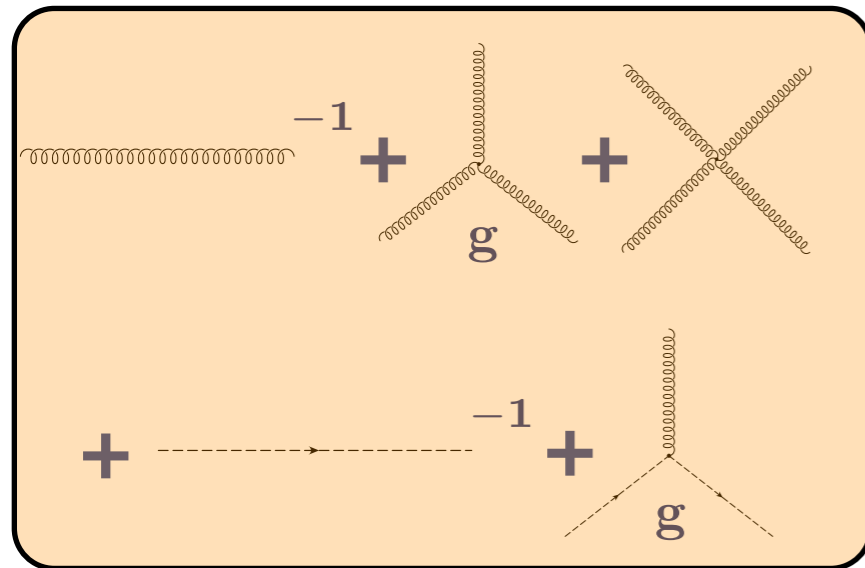
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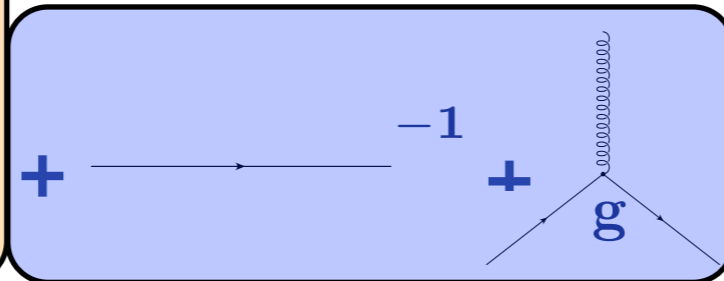
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April 2004



Pure gauge theory



matter sector

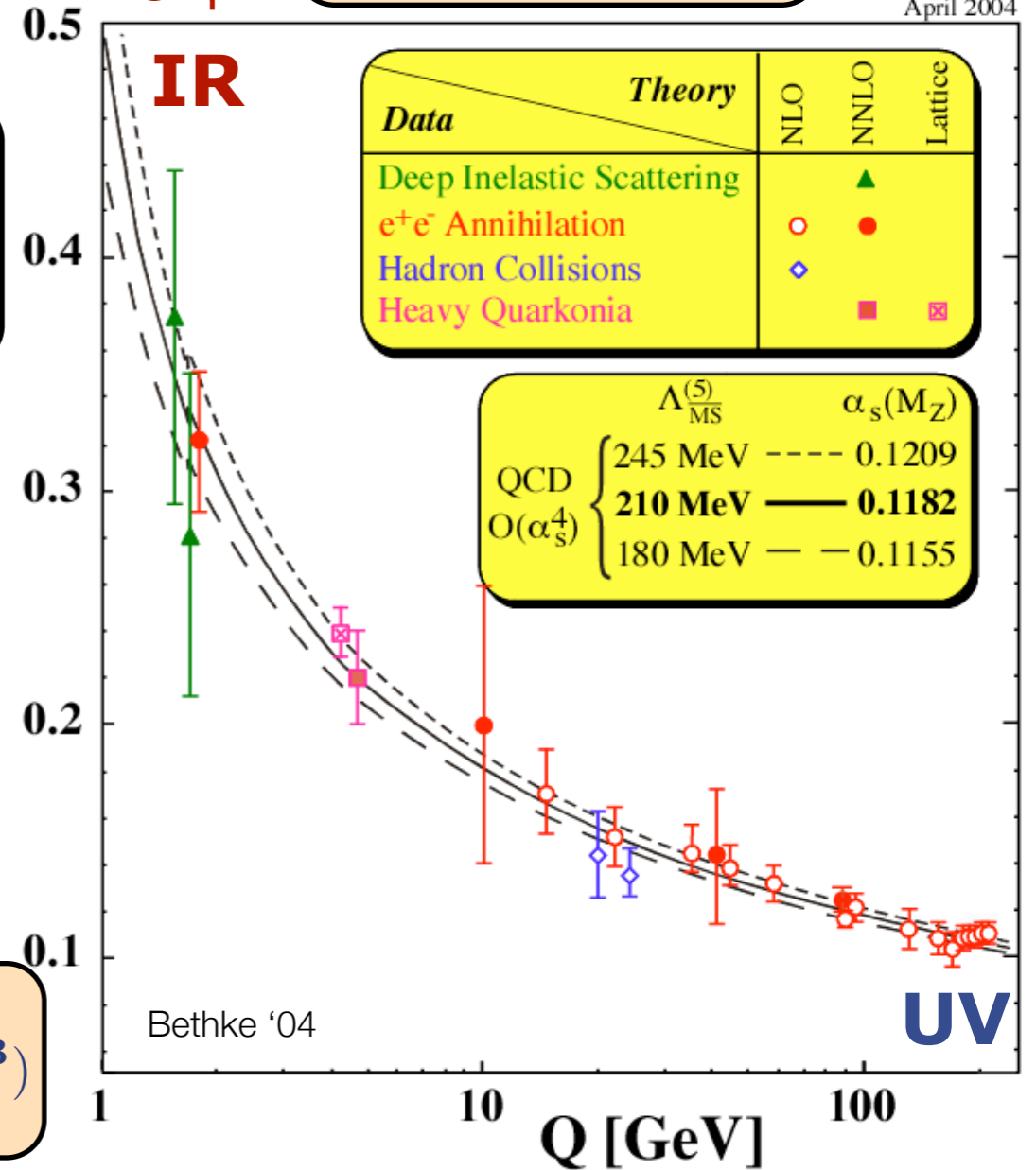
- running coupling (1-loop)

$$\alpha_s(Q) = \frac{\alpha_s(\mu)}{1 - \alpha_s(\mu)\beta_0 \log(Q^2/\mu^2)}$$

- beta function

$$\beta = Q^2 \frac{\partial \alpha_s(Q)}{\partial Q^2} = \beta_0 \alpha_s(\mu)^2 + \mathcal{O}(\alpha_s(\mu)^3)$$

$$\beta = -\frac{1}{12\pi} (33 - 2N_f) \alpha_s^2 + \mathcal{O}(\alpha_s^3)$$



Nobel Prize '04  
Gross, Politzer, Wilczek

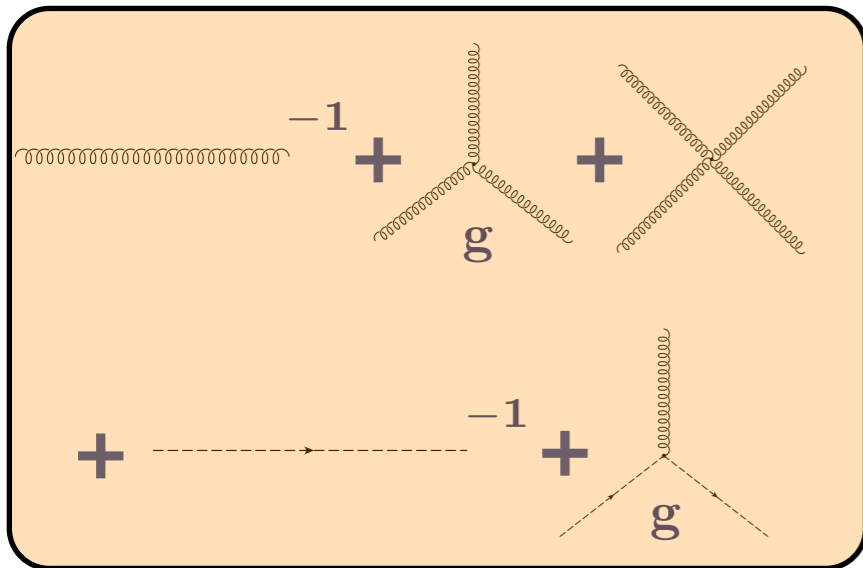
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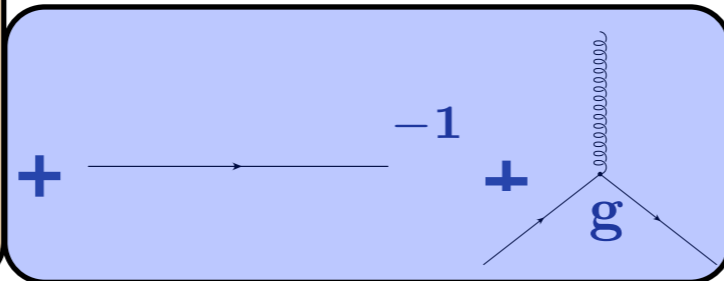
Millenium Prize 1 Mio \$

$$\alpha_s(Q) = \frac{g^2(Q)}{4\pi}$$

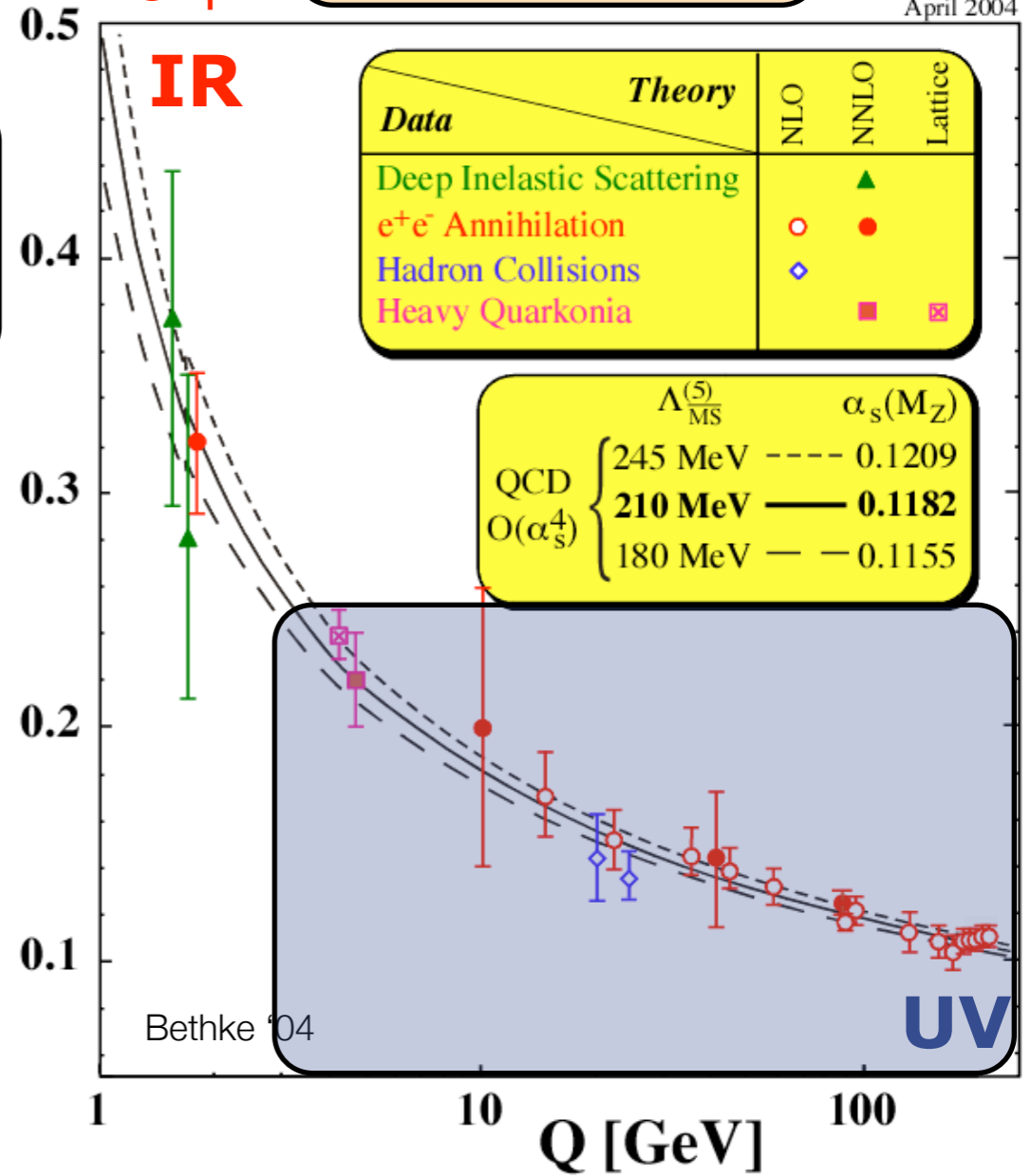
April 2004



Pure gauge theory



matter sector



- running coupling (1-loop)

$$\alpha_s(Q) = \frac{\alpha_s(\mu)}{1 - \alpha_s(\mu)\beta_0 \log(Q^2/\mu^2)}$$

- UV: asymptotic freedom

$$\alpha_s(Q \rightarrow \infty) = 0$$

Nobel Prize '04  
Gross, Politzer, Wilczek



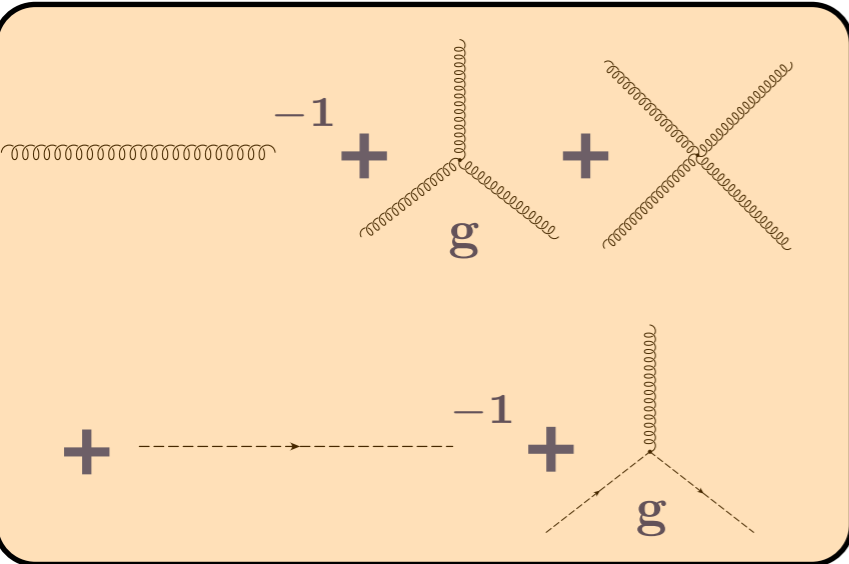
# QCD, asymptotic freedom and all that

## Running coupling at low and high energies

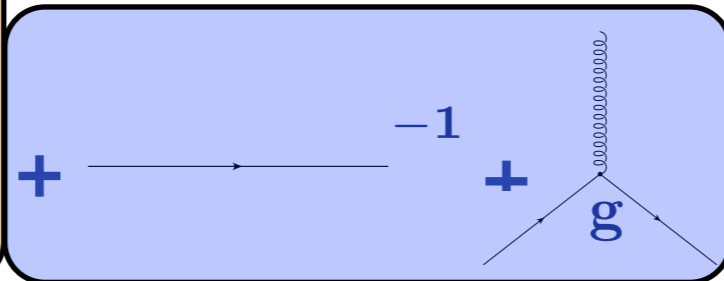
Millenium Prize 1 Mio \$

$$\alpha_s(Q) = \frac{g^2(Q)}{4\pi}$$

April 2004



Pure gauge theory



matter sector

- running coupling (1-loop)

$$\alpha_s(Q) = \frac{\alpha_s(\mu)}{1 - \alpha_s(\mu)\beta_0 \log(Q^2/\mu^2)}$$

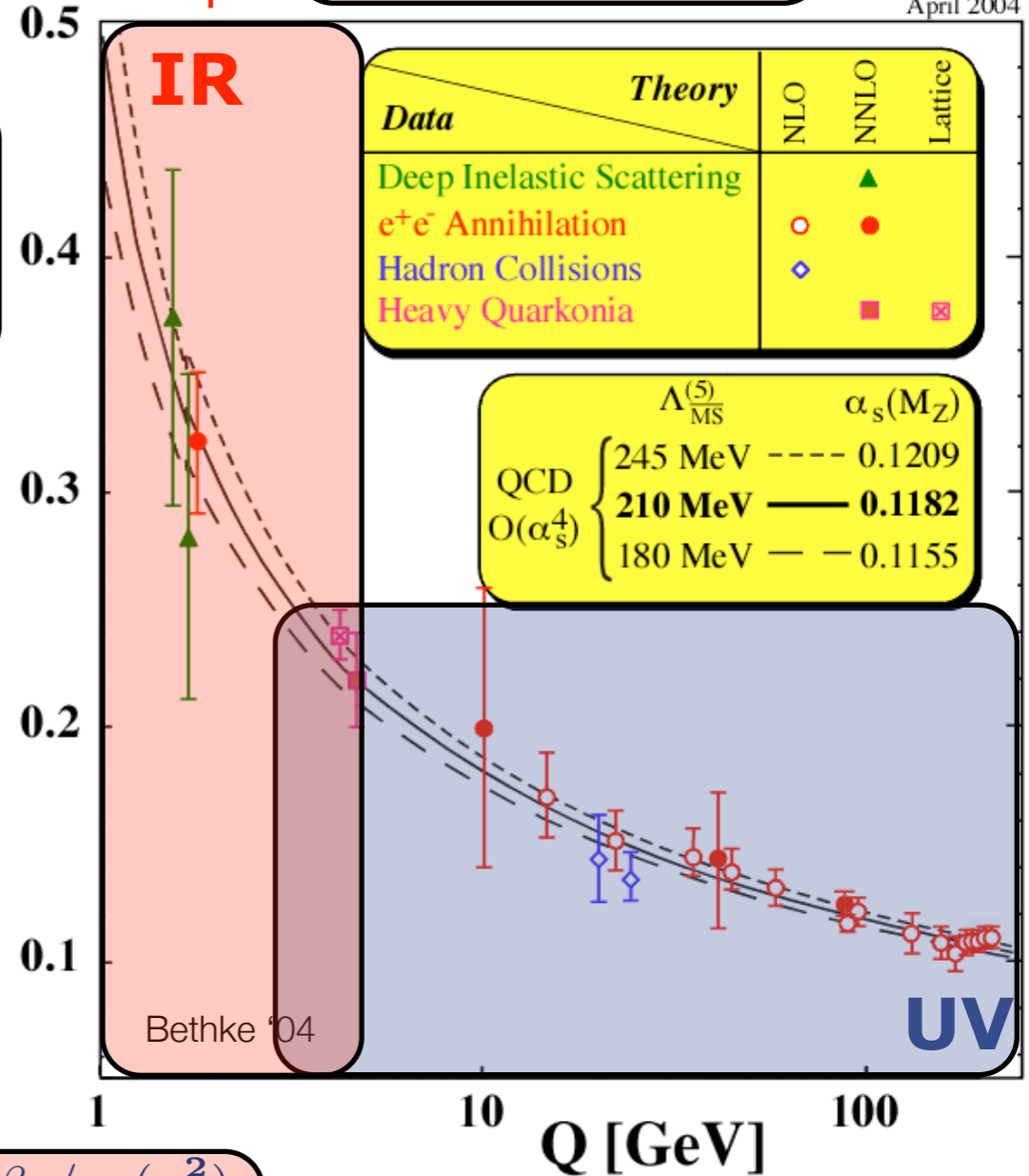
- IR: failure of perturbation theory

$$\alpha_s(\Lambda_{\text{QCD}}^2) = \infty$$

at

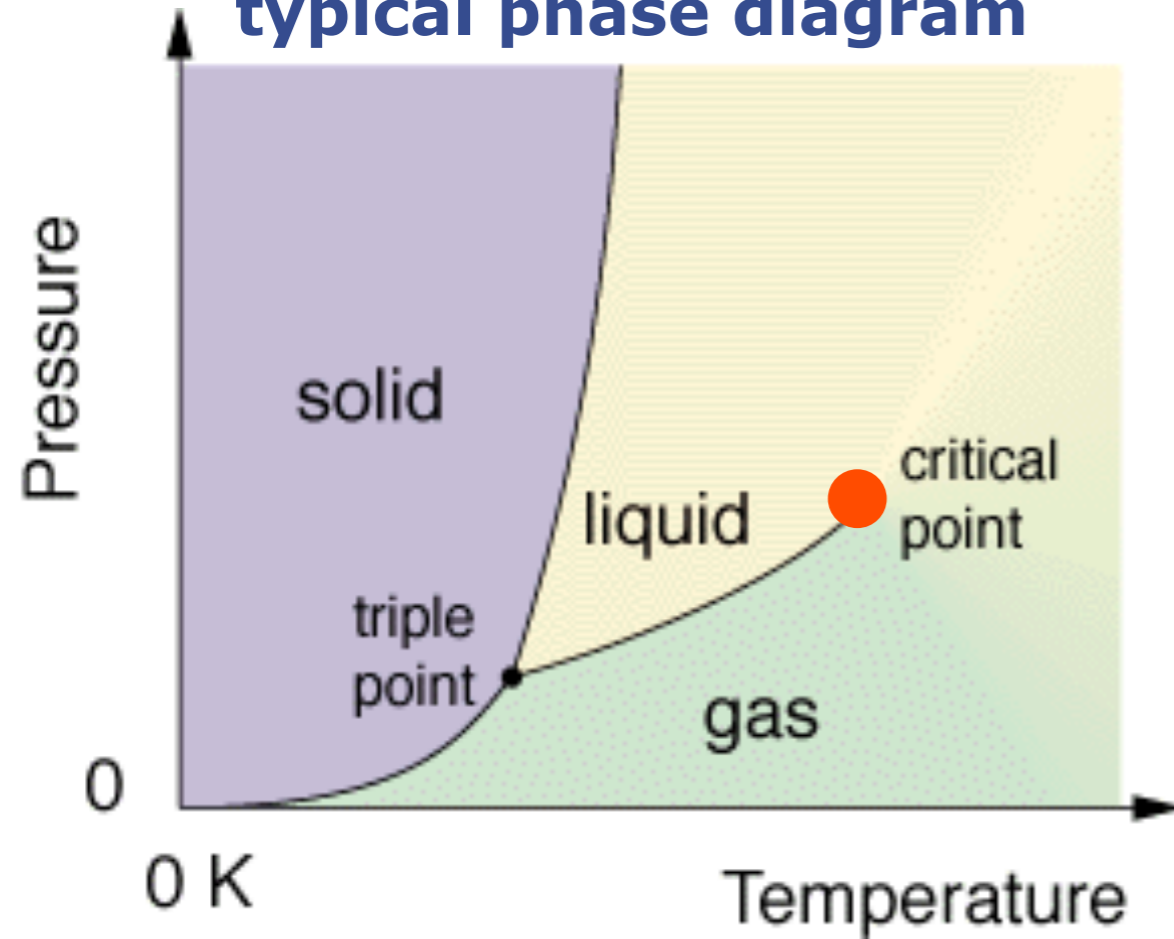
$$\Lambda_{\text{QCD}}^2 = \mu^2 e^{-\beta_0/\alpha_s(\mu^2)}$$

$$\Lambda_{\text{QCD}} = 217_{-23}^{+25} \text{ MeV}$$



Nobel Prize '04  
Gross, Politzer, Wilczek

typical phase diagram



<http://tl.tkk.fi/research/theory/TypicalPD.gif>

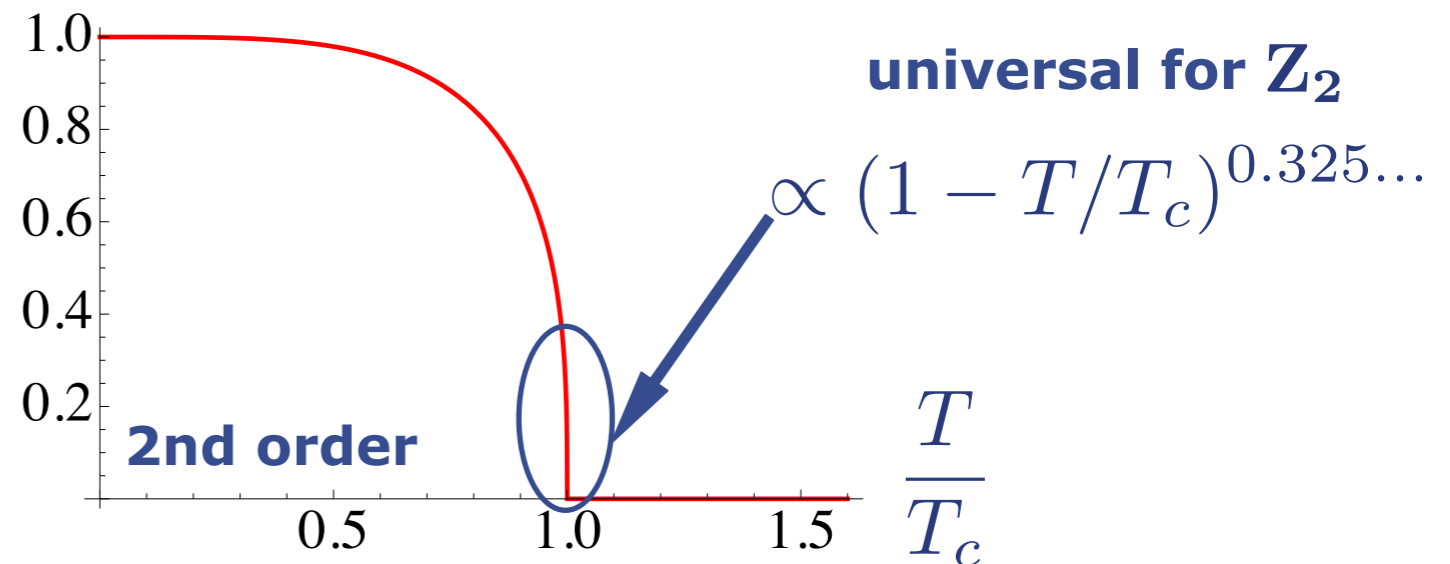
Order parameter: density  $\rho$

density jumps	1st order phase transition
derivative of density jumps	2nd order phase transition
density smooth	cross-over

Ising model in 3d: (  $\downarrow$   $\uparrow$  )-spin system

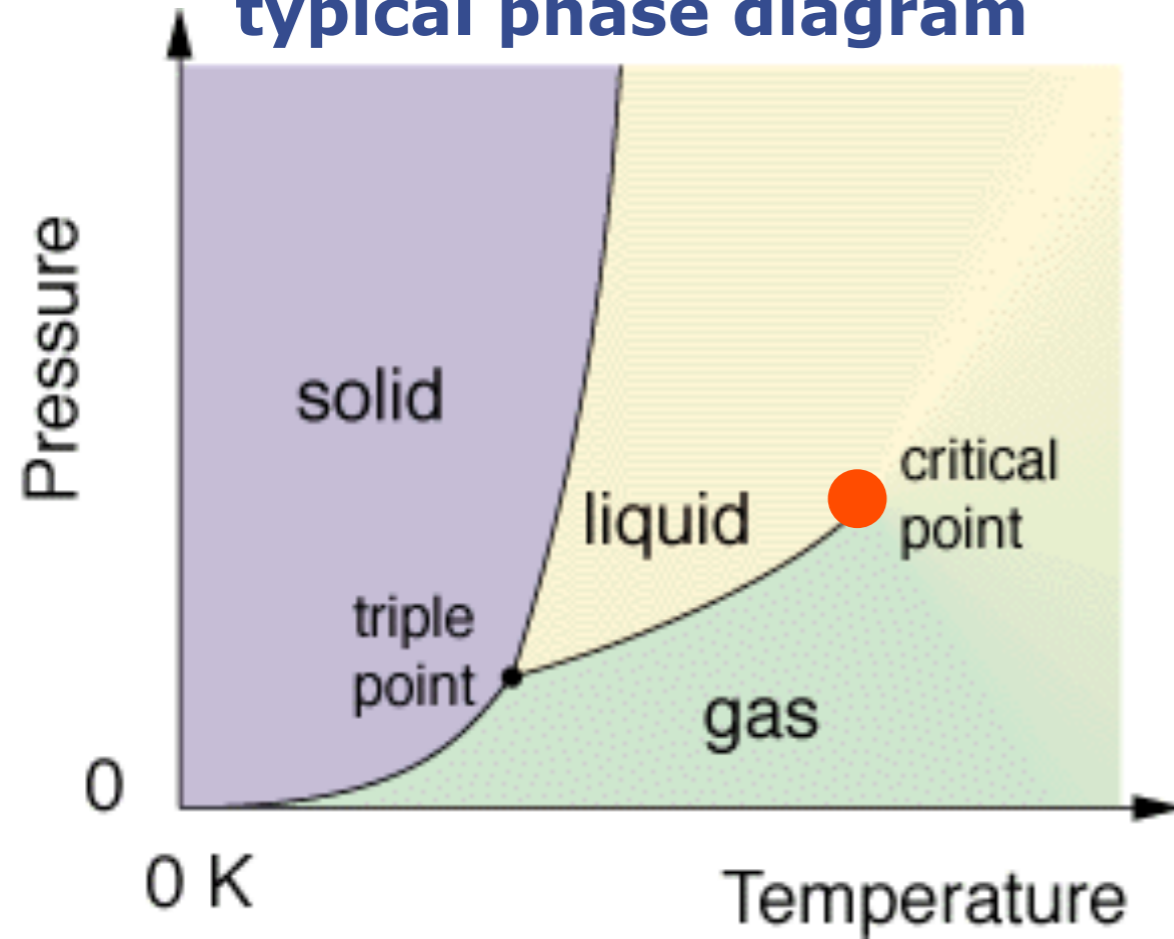
Order parameter:  $\langle \uparrow \rangle$

$$\frac{\langle \uparrow \rangle}{\langle \uparrow \rangle_0}$$



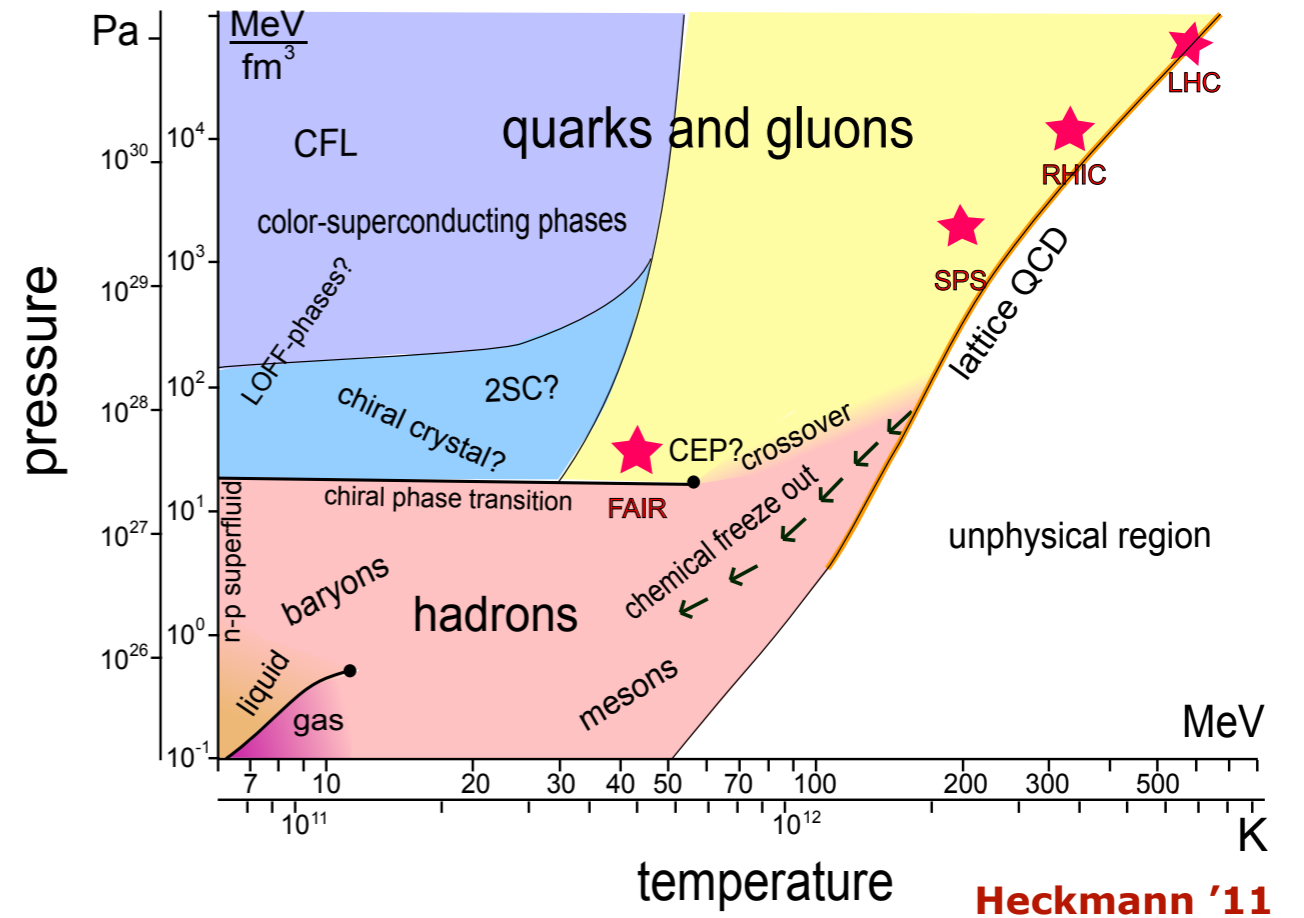
# Phase diagrams & order parameters

## typical phase diagram



<http://tli.tkk.fi/research/theory/TypicalPD.gif>

## phase diagram of QCD



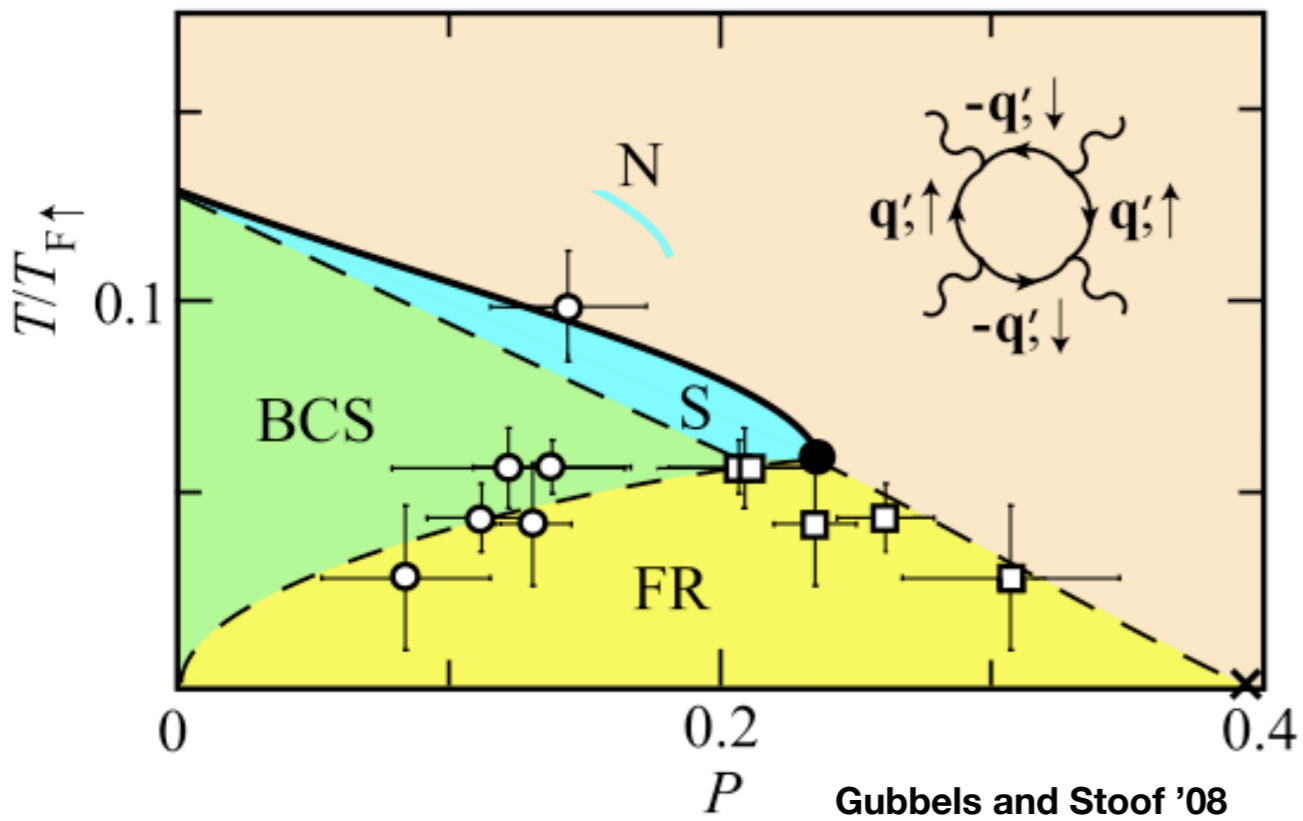
Heckmann '11

## Phases in QCD

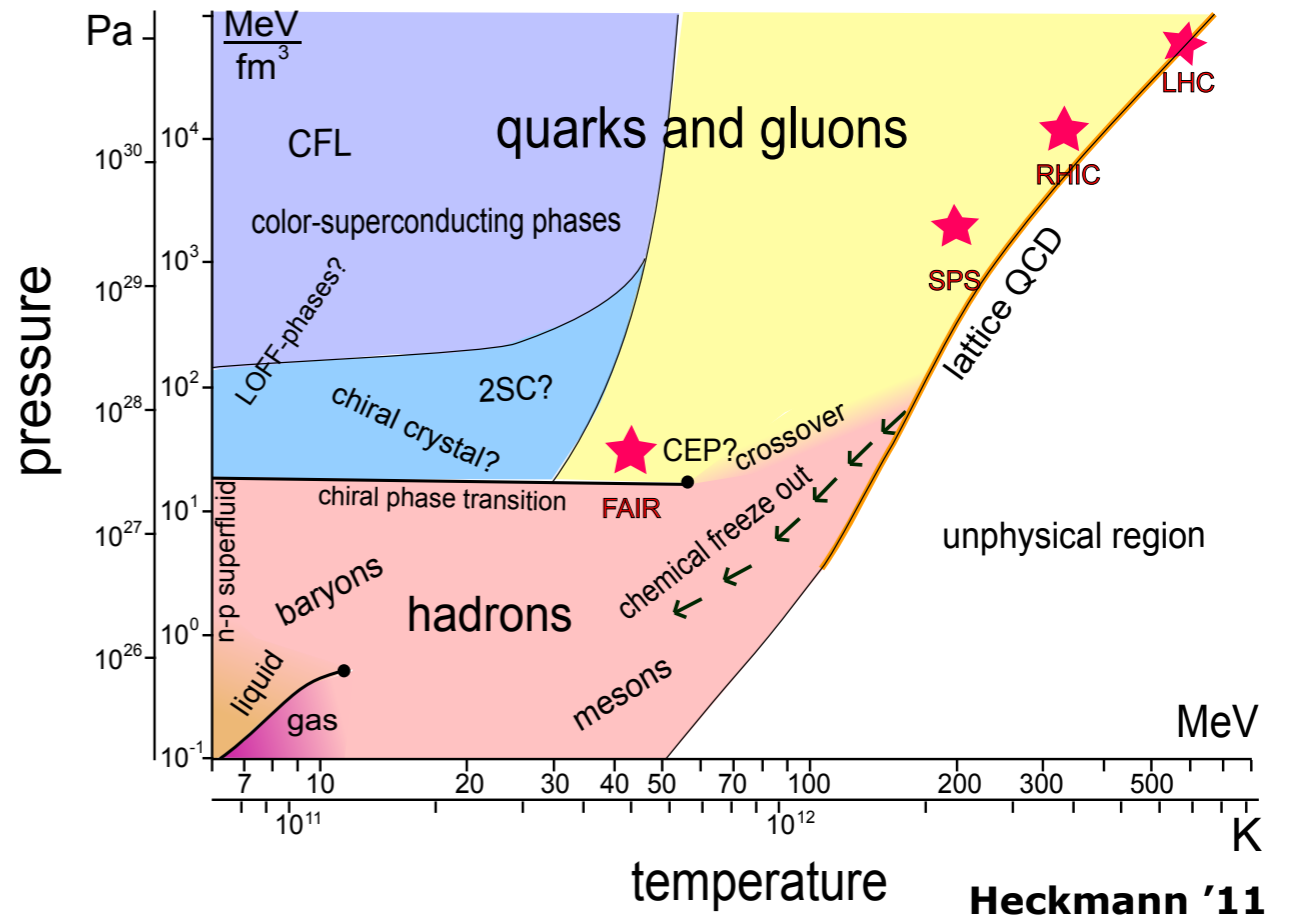
quarks massless - massive

quarks confined - deconfined

### Phase diagram of cold atoms



### phase diagram of QCD

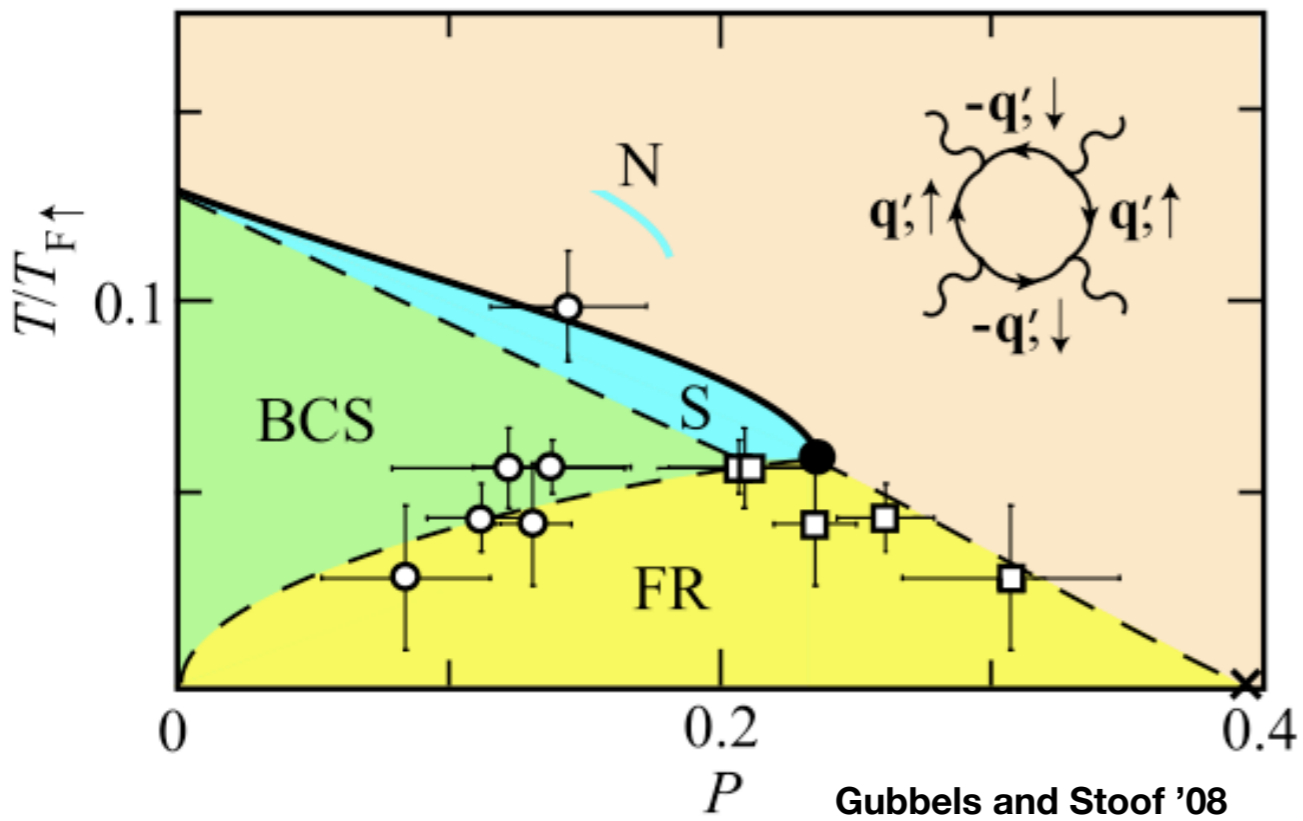


### Phases in QCD

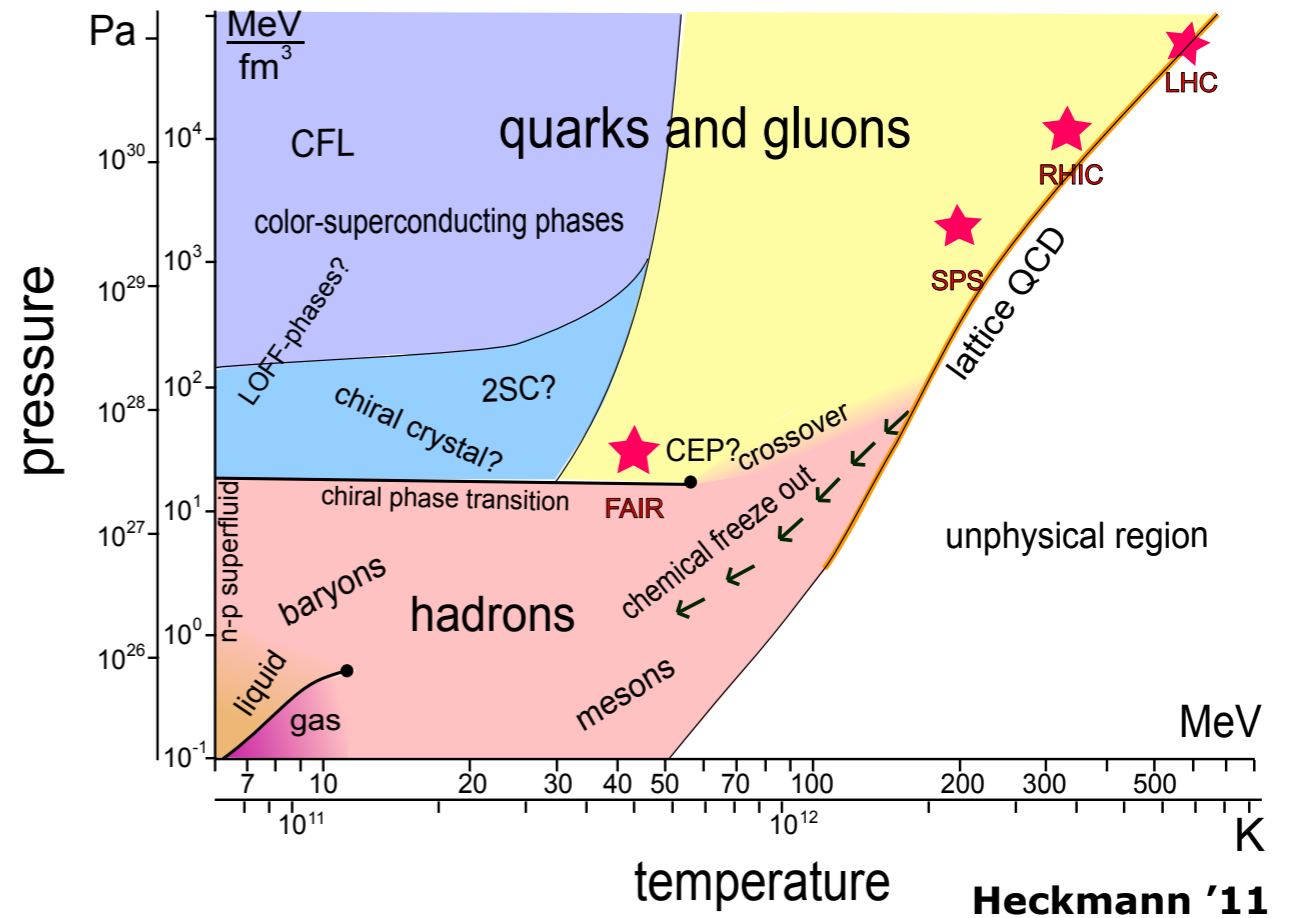
quarks massless - massive

quarks confined - deconfined

## Phase diagram of cold atoms



## phase diagram of QCD



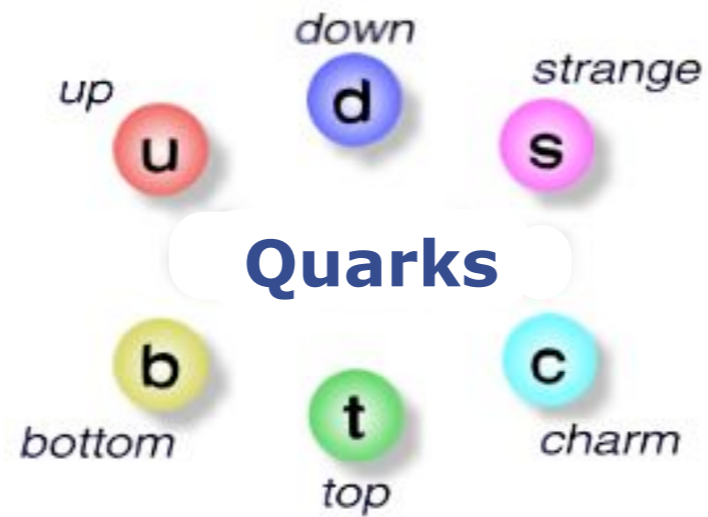
## Phases in QCD

quarks massless - massive

quarks confined - deconfined

Strong chiral symmetry breaking makes up for 99% of the mass of the visible part of matter in the universe

# Chiral symmetry breaking



# Chiral symmetry breaking

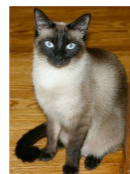
strong chiral symmetry breaking  $\Delta m_{\chi SB} \approx 400 \text{ MeV}$



Generation	first	second	third	Charge
Mass [MeV]	1.5-4	1150-1350	$170 \times 10^3$	
Quark	u	c	t	$\frac{2}{3}$
Quark	d	s	b	$-\frac{1}{3}$
Mass [MeV]	4-8	80-130	$(4.1-4.4) \times 10^3$	

two light flavours and one heavy flavour: 2+1

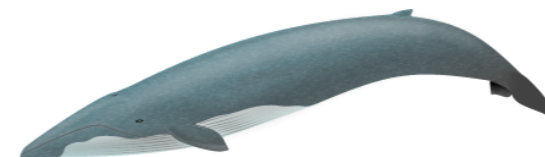
up



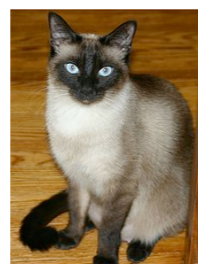
charm



top



down



strange



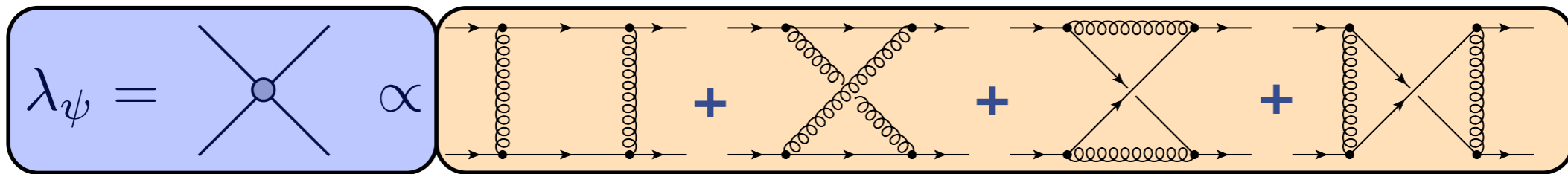
bottom



# Chiral symmetry breaking

- Perturbative four-fermi coupling

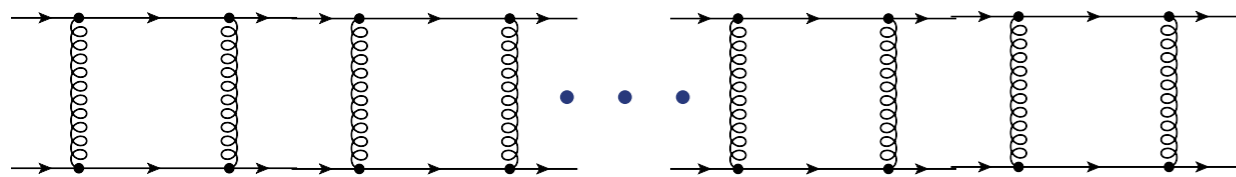
$$\frac{\lambda_\psi}{2} \int [(\bar{q}q)^2 + (i\bar{q}\gamma_5 \vec{\tau}q)^2]$$



$$\lambda_\psi \propto \alpha_s^2$$

$$N_f = 2 : \vec{\tau} = (\sigma_1, \sigma_2, \sigma_3)$$

- Fermionic mass term for  $\langle \bar{q}q \rangle \neq 0$



$$\frac{\lambda_\psi}{2} \int (\bar{q}q)^2 \longrightarrow \frac{\lambda_\psi}{2} \int \langle \bar{q}q \rangle \bar{q}q$$

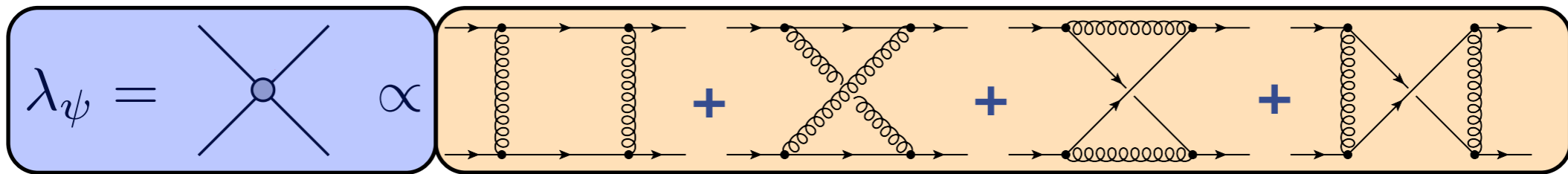
mean field



# Chiral symmetry breaking

- **Perturbative four-fermi coupling**

$$\frac{\lambda_\psi}{2} \int [(\bar{q}q)^2 + (i\bar{q}\gamma_5\vec{\tau}q)^2]$$

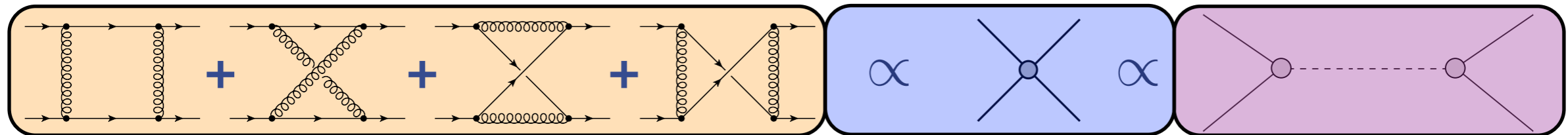


$$\lambda_\psi \propto \alpha_s^2$$

$$N_f = 2 : \vec{\tau} = (\sigma_1, \sigma_2, \sigma_3)$$

- **Bosonisation (Hubbard-Stratonovich)**

$$\langle \sigma \rangle \neq 0$$

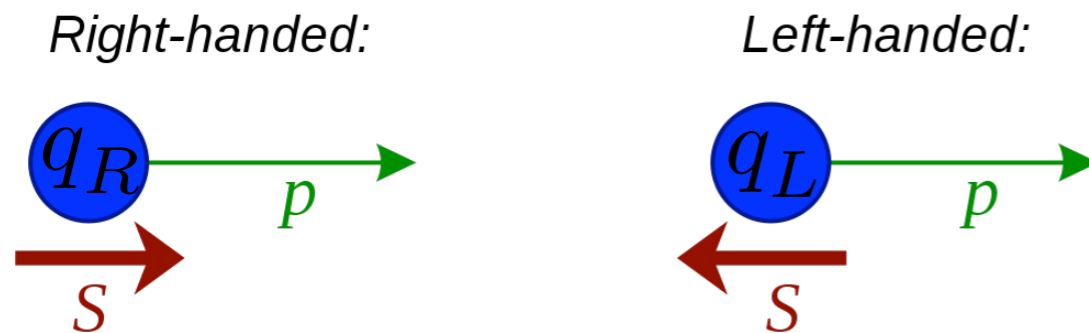


$$\frac{\lambda_\psi}{2} \int [(\bar{\psi}\psi)^2 + (i\bar{\psi}\gamma_5\vec{\tau}\psi)^2] = \frac{m_\sigma^2}{2} \int_x (\sigma^2 + \vec{\pi}^2) + i h \int_x \bar{\psi}(\sigma + i\gamma_5\vec{\tau}\vec{\pi})\psi$$

EOM( $\sigma$ )

# Chiral symmetry breaking

- Chirality for massless particles



- Order parameter

$$\sigma \simeq \langle \bar{q}q \rangle$$

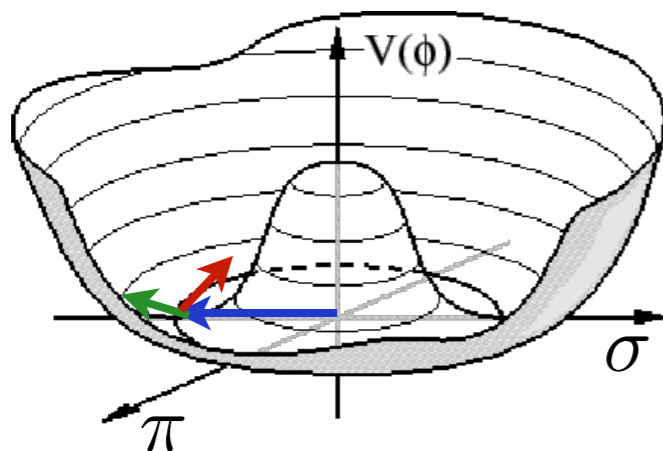
chiral condensate

$$\bar{q}q = q_R^\dagger q_L + q_L^\dagger q_R$$

- Chiral symmetry  $\sigma = 0$

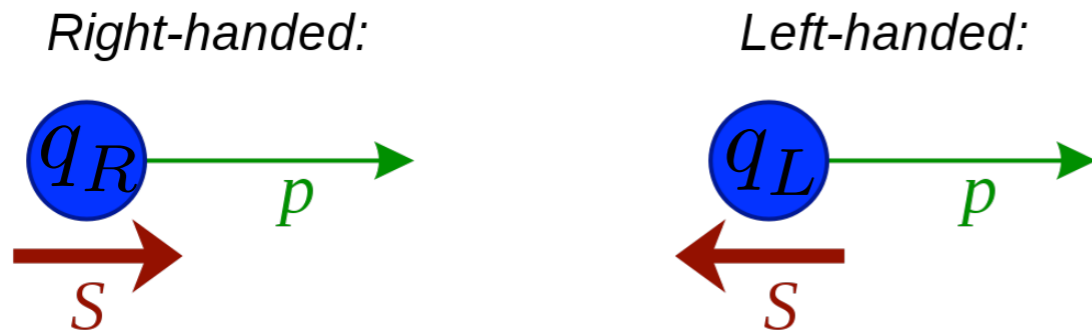
- Symmetry broken  $\sigma \neq 0$

- Meson potential



# Chiral symmetry breaking

- Chirality for massless particles



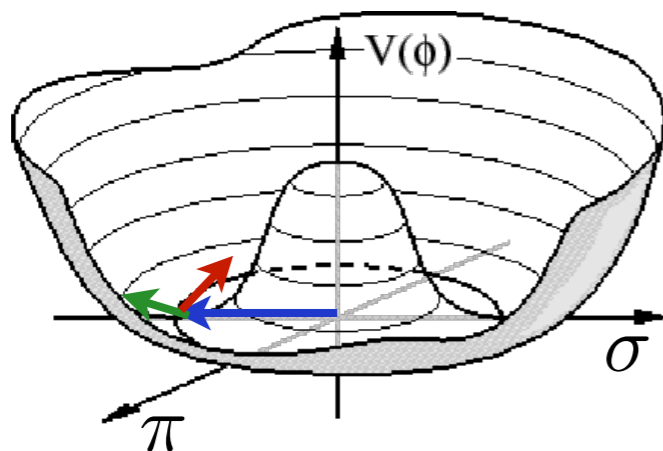
- Order parameter

$$\sigma \simeq \langle \bar{q}q \rangle \quad \text{chiral condensate}$$

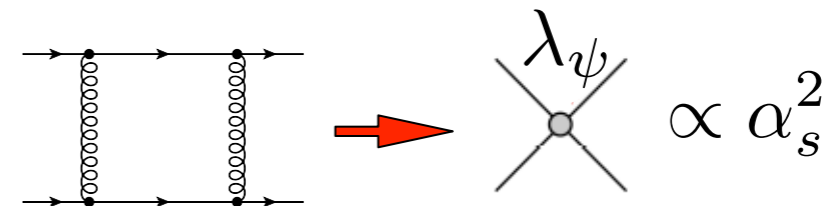
- Chiral symmetry  $\sigma = 0$

- Symmetry broken  $\sigma \neq 0$

- Meson potential

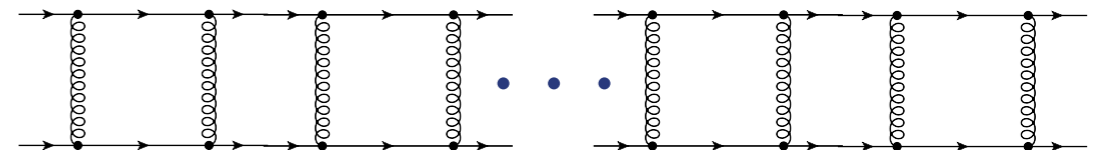


## chiral symmetry



$$\int d^4x \lambda_\psi [(\bar{q}q)^2 - (\bar{q}\gamma_5 q)^2]$$

## strong correlations

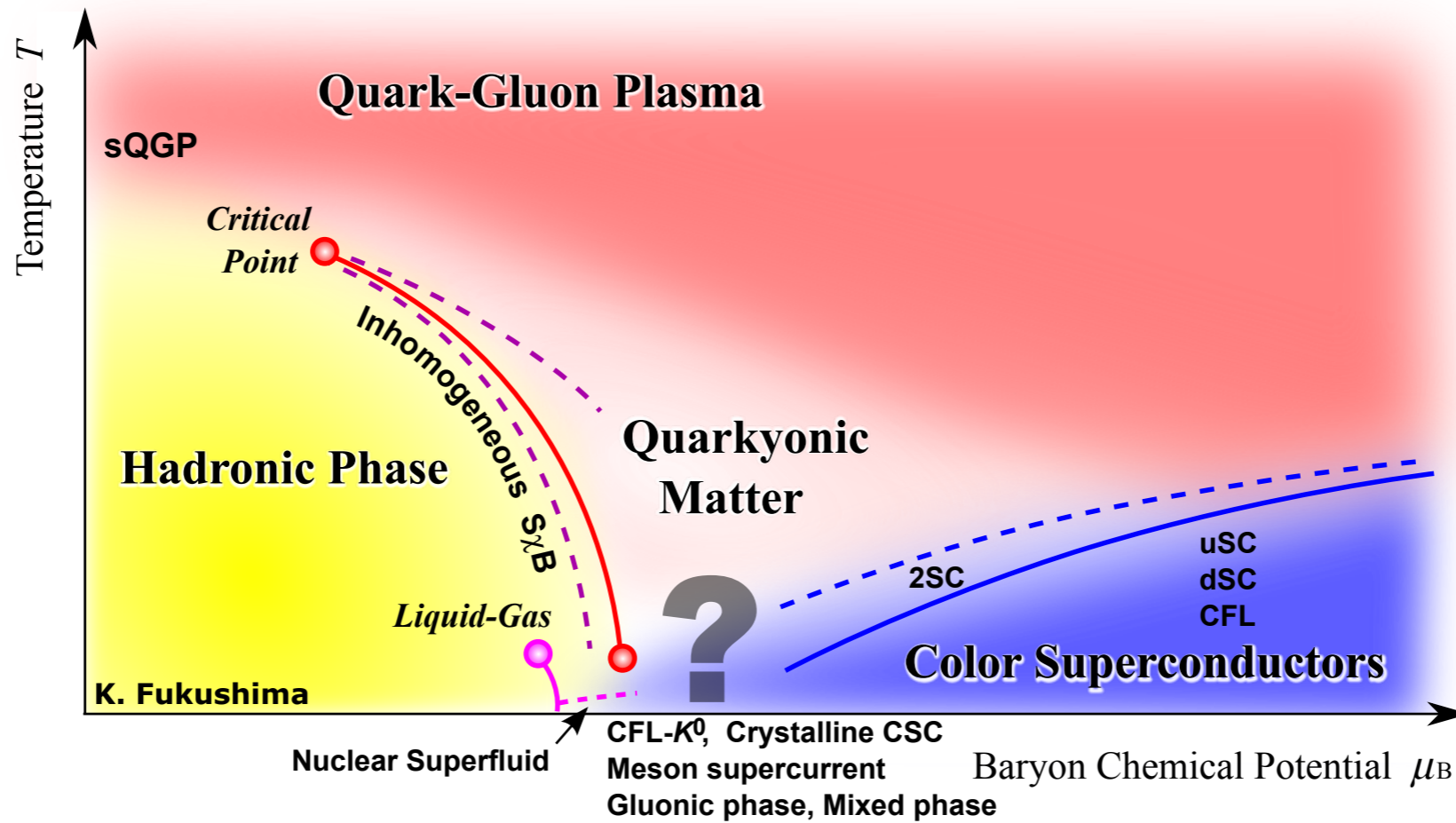


$$\langle \bar{q}q \rangle \neq 0$$

mass term:  $\langle \bar{q}q \rangle \bar{q}q$

chiral symmetry broken

# Phase diagram & order parameters



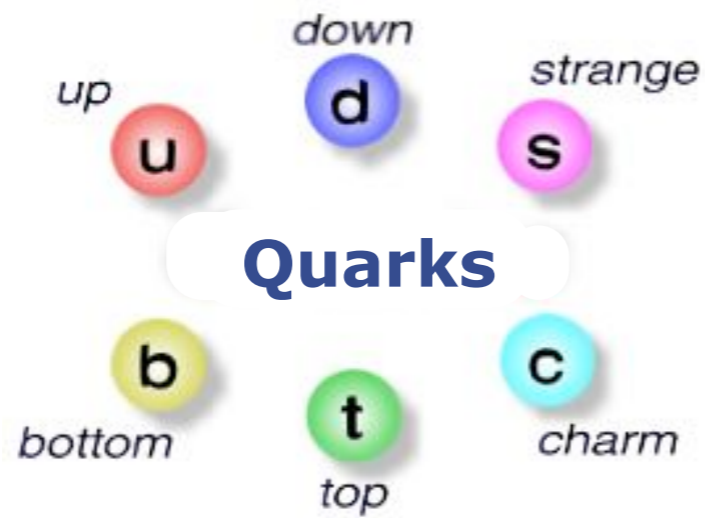
## Phases in QCD

quarks massless - massive

quarks confined - deconfined

chiral condensate  $\int_{\vec{x}} \langle \bar{q}(\mathbf{x})q(\mathbf{x}) \rangle$

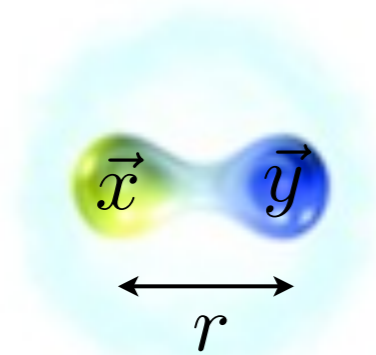
# Confinement



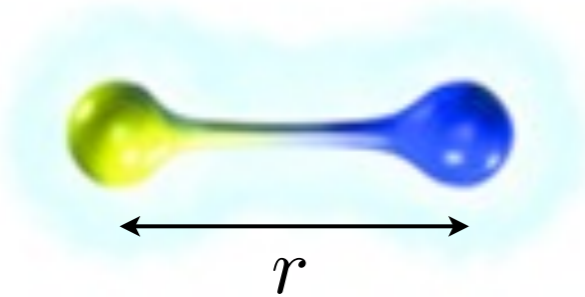
**Gluons**

# Confinement

Free energy  $F_{q\bar{q}}$  of a quark - antiquark pair

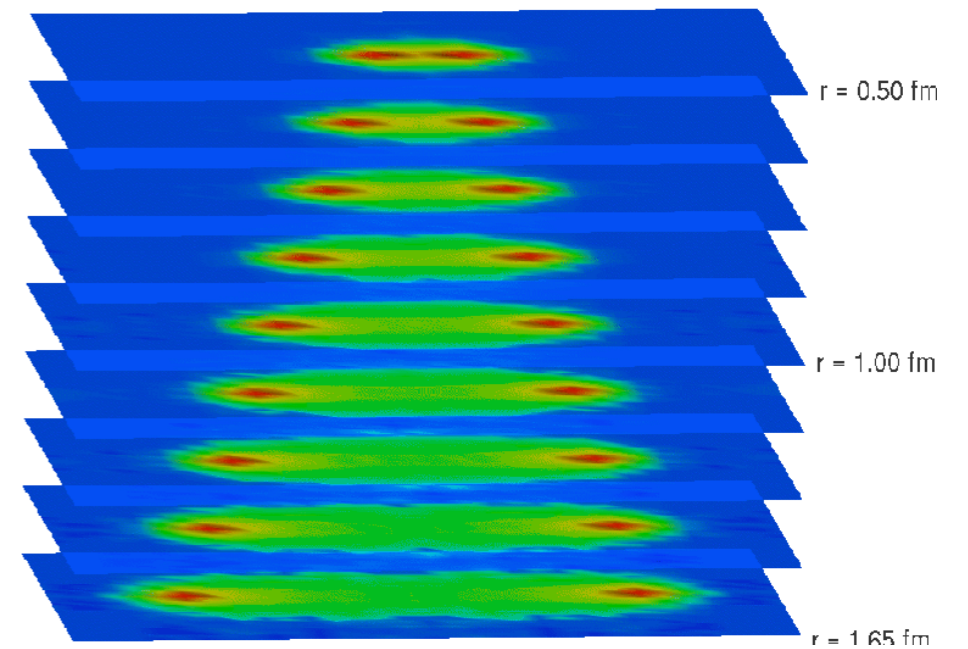


$$F_{q\bar{q}} \simeq -\frac{1}{r}$$



$$F_{q\bar{q}} \simeq \sigma r$$

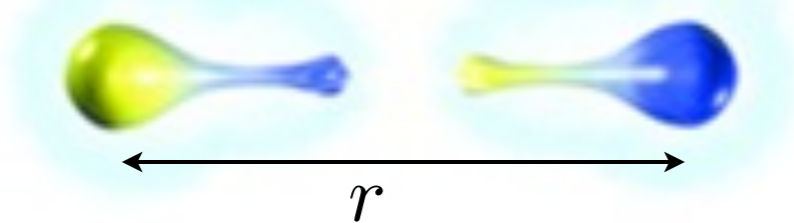
pure gauge theory



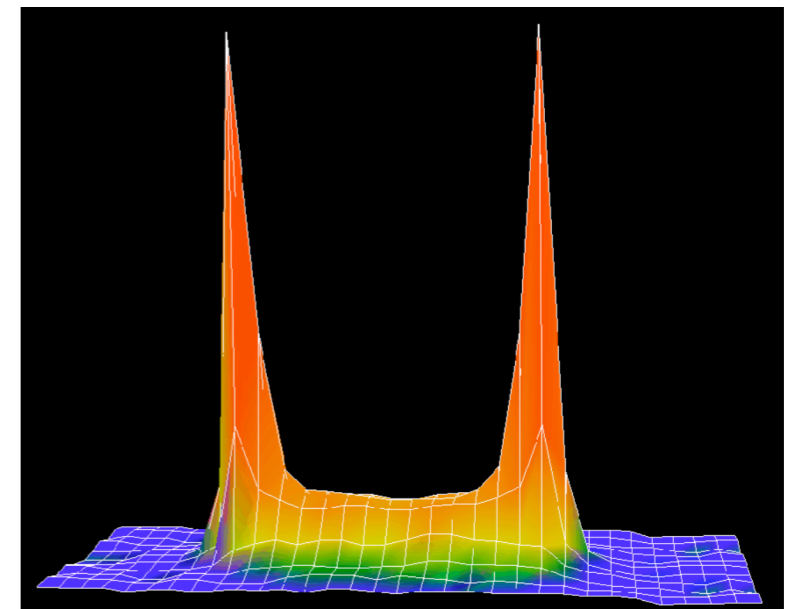
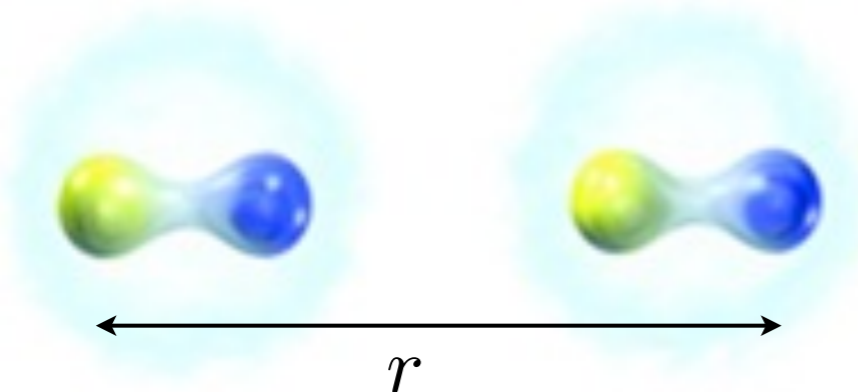
Energy density

Bali et al. '94

string breaking at  $r \approx 1\text{fm}$

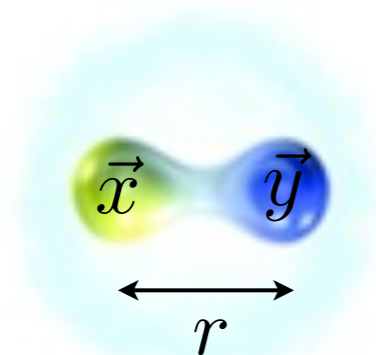


$$F_{q\bar{q}} \simeq \text{const.}$$

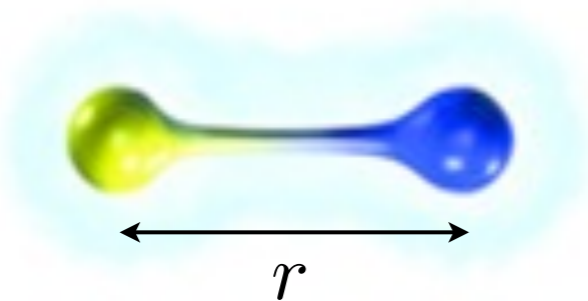


# Confinement

Free energy  $F_{q\bar{q}}$  of a quark - antiquark pair

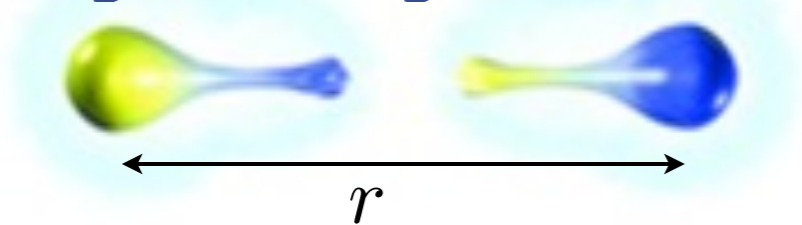


$$F_{q\bar{q}} \simeq -\frac{1}{r}$$

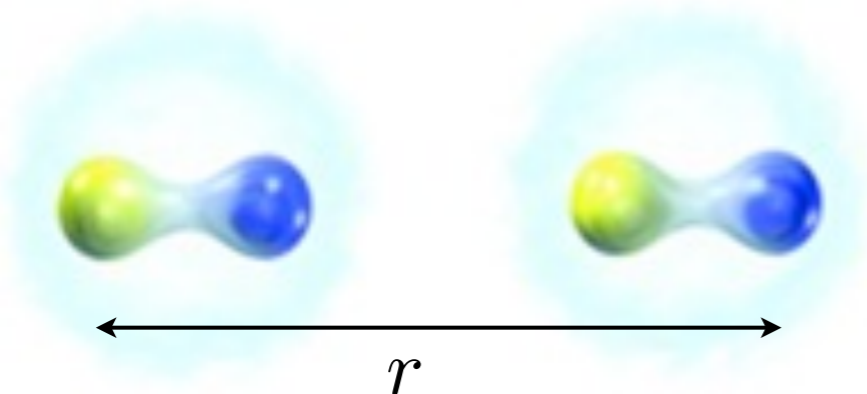


$$F_{q\bar{q}} \simeq \sigma r$$

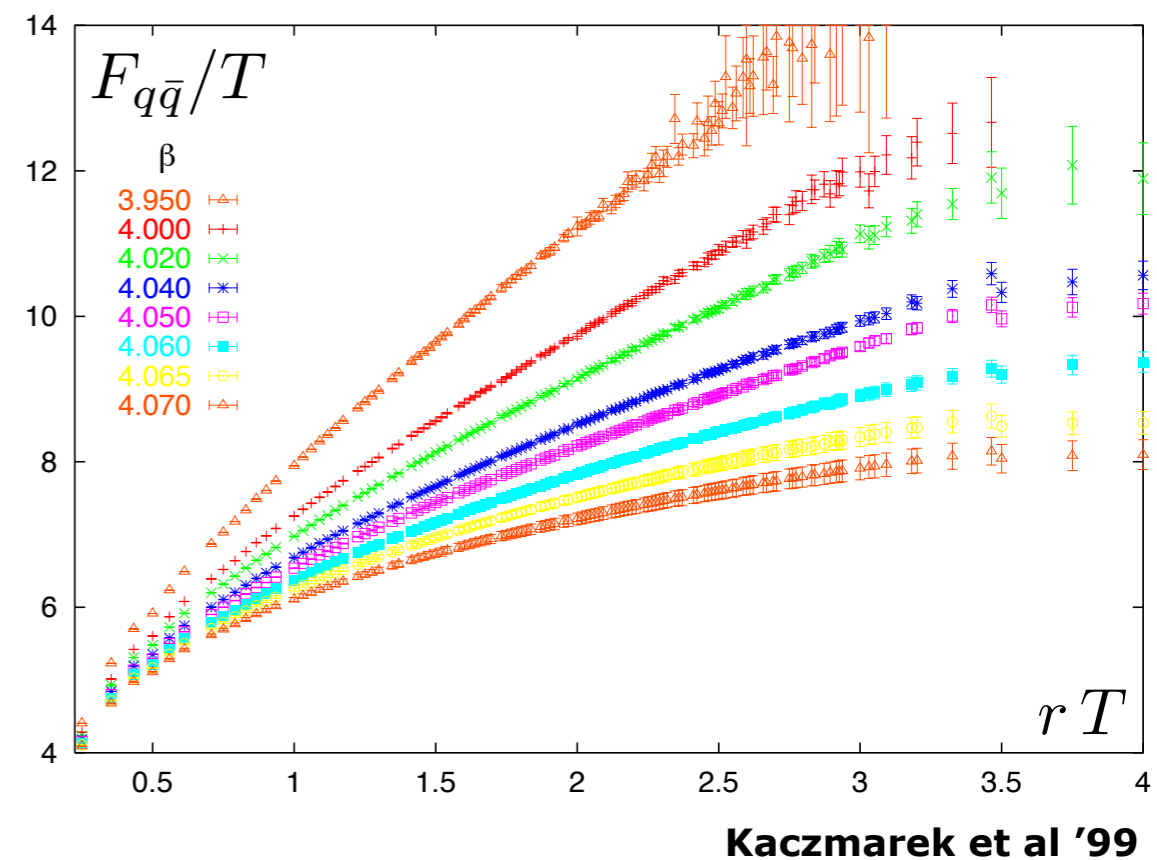
string breaking at  $r \approx 1\text{fm}$



$$F_{q\bar{q}} \simeq \text{const.}$$

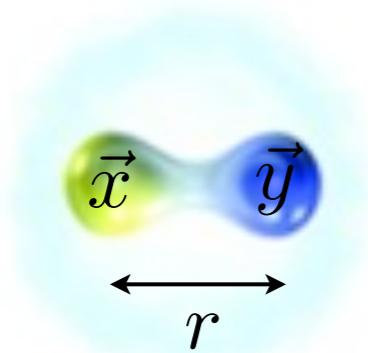


pure gauge theory

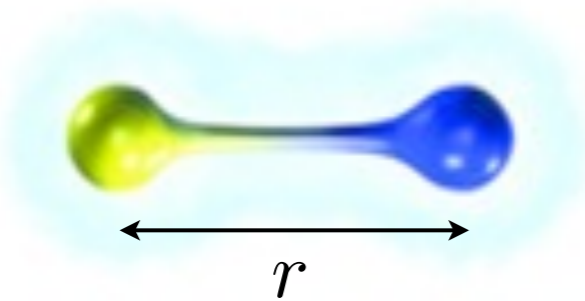


# Confinement

Free energy  $F_{q\bar{q}}$  of a quark - antiquark pair

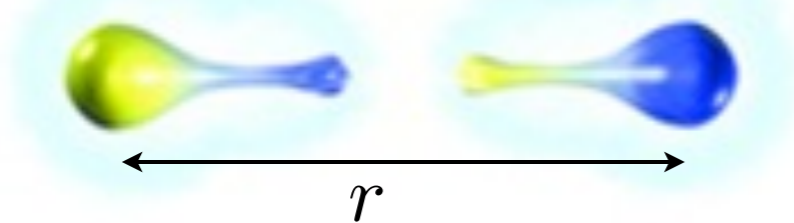


$$F_{q\bar{q}} \simeq -\frac{1}{r}$$

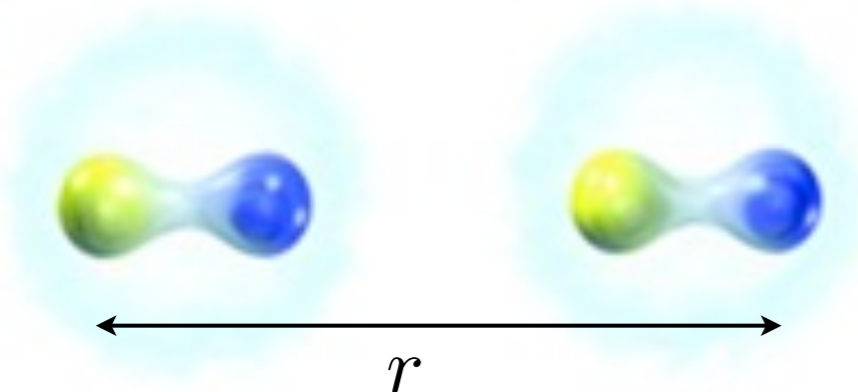


$$F_{q\bar{q}} \simeq \sigma r$$

string breaking at  $r \approx 1\text{fm}$



$$F_{q\bar{q}} \simeq \text{const.}$$



Order parameter  $\sim \langle q \rangle$

$$\Phi = e^{-\frac{1}{2T} F_{q\bar{q}}(\infty)}$$

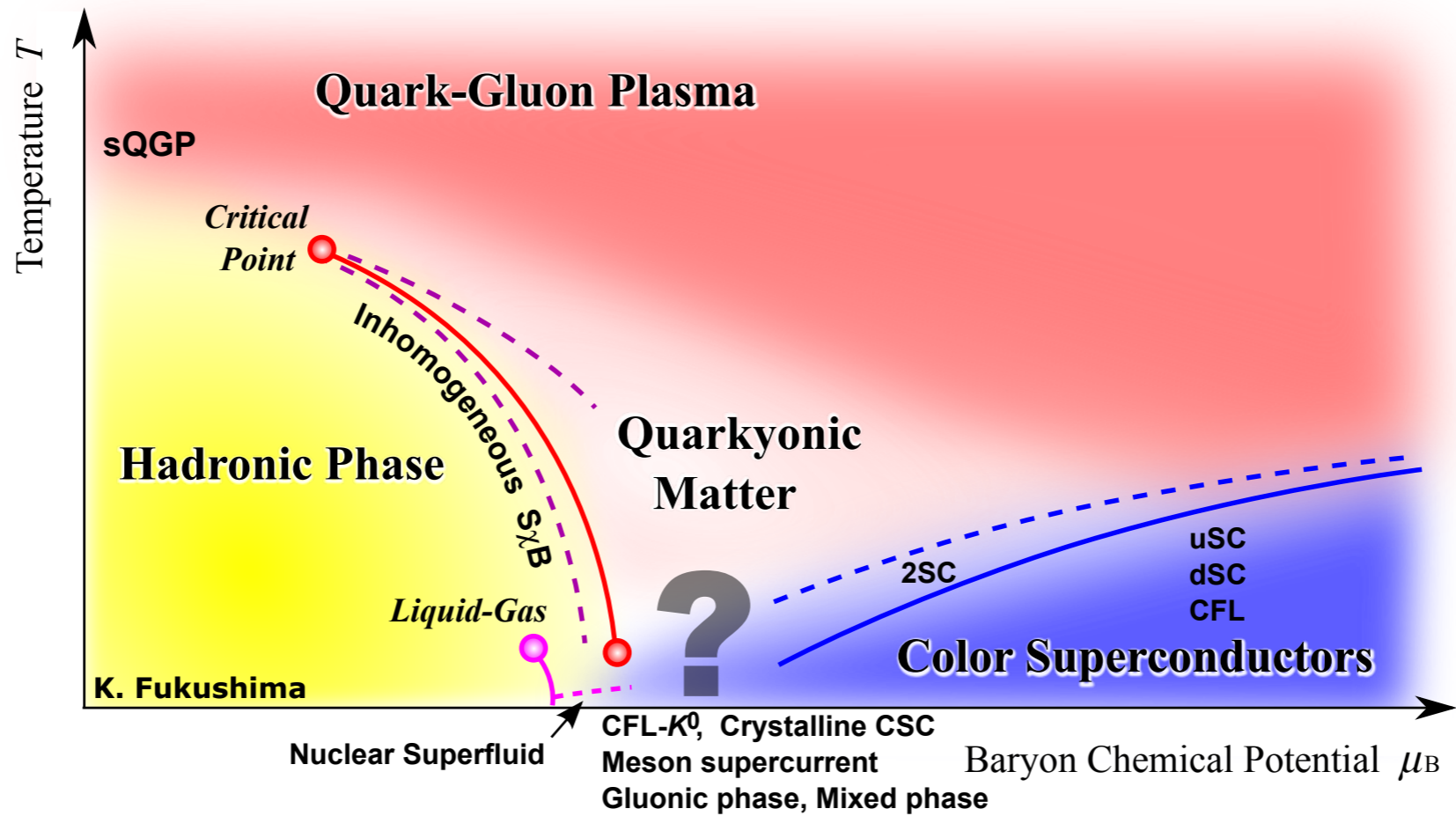
- **Confinement**  $\Phi = 0$
- **Deconfinement**  $\Phi \neq 0$

**Polyakov loop**

$$\Phi = \frac{1}{3} \langle \text{Tr } \mathcal{P} \exp\{ig \int_0^{1/T} dx_0 A_0\} \rangle$$



# Phase diagram & order parameters



## Phases in QCD

quarks massless - massive

chiral condensate  $\int_{\vec{x}} \langle \bar{q}(\mathbf{x})q(\mathbf{x}) \rangle$

quarks confined - deconfined

Polyakov loop  $\Phi = \frac{1}{N_c} \langle \text{tr } \mathcal{P} e^{i g \int_0^\beta A_0(\mathbf{x})} \rangle$

# **(II) Functional Renormalisation group for QCD**

---

- **Introduction to the functional renormalisation group**

- Derivation of the flow equation

- Expansion schemes

- Optimisation and error control\*

- **FRG for QCD**

- FRG for QCD and  $T=0$  Yang-Mills theories

- Dynamical hadronisation

- QCD correlation functions at  $T=0$

# **Functional Renormalisation Group**

# Functional Renormalisation Group

---

## Generating functional $Z$

$$Z[J] = \frac{1}{\mathcal{N}} \int d\varphi e^{-S[\varphi] + \int_x J\varphi}$$

**partition function**

$$\langle \varphi \rangle_J = \phi$$

$$S[\varphi] = \frac{1}{2} \int_x \left[ \partial_\mu \varphi \partial_\mu \varphi + m^2 \varphi^2 + \frac{\lambda}{4} \varphi^4 \right]$$

**classical action**

**zero-dimensional example: 'Functional' flows for integrals**

# Functional Renormalisation Group

---

## Generating functional $Z$

$$Z[J] = \frac{1}{\mathcal{N}} \int d\varphi e^{-S[\varphi] + \int_x J\varphi}$$

partition function

$$\langle \varphi \rangle_J = \phi$$

## Effective action $\Gamma$

$$\Gamma[\phi] = -\log \int d\hat{\varphi} e^{-S[\hat{\varphi} + \phi] + \int_x \hat{\varphi} \frac{\delta\Gamma[\phi]}{\delta\phi}}$$

free energy

$$\varphi = \hat{\varphi} + \phi$$

$$\langle \hat{\varphi} \rangle_{\frac{\delta\Gamma}{\delta\phi}} = 0$$

$$J = \frac{\delta\Gamma}{\delta\phi}$$

# Functional Renormalisation Group

## Generating functional $Z$

$$Z[J] = \frac{1}{\mathcal{N}} \int d\varphi e^{-S[\varphi] + \int_x J\varphi}$$

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free energy

$$\varphi = \hat{\varphi} + \phi$$

$$\langle \hat{\varphi} \rangle_{\frac{\delta\Gamma}{\delta\phi}} = 0$$

$$J = \frac{\delta\Gamma}{\delta\phi}$$

$$\Gamma[\phi] = \sup_J \left( \int_x J \cdot \phi - \log Z[J] \right)$$

Legendre transform

# Functional Renormalisation Group

## Generating functional $Z$

$$Z[J] = \frac{1}{\mathcal{N}} \int d\varphi e^{-S[\varphi] + \int_x J\varphi}$$

partition function

$$\langle \varphi \rangle_J = \phi$$

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free energy

$$\varphi = \hat{\varphi} + \phi$$
$$\langle \hat{\varphi} \rangle_{\frac{\delta\Gamma}{\delta\phi}} = 0$$

$$J = \frac{\delta\Gamma}{\delta\phi}$$

## Dyson-Schwinger equation

$$\frac{\delta\Gamma[\phi]}{\delta\phi(x)} = \left\langle \frac{\delta S[\hat{\varphi} + \phi]}{\delta\phi(x)} \right\rangle$$

quantum equation of motion

# Functional Renormalisation Group

---

## Dyson-Schwinger equation

$$\frac{\delta\Gamma[\phi]}{\delta\phi(x)} = \left\langle \frac{\delta S[\hat{\phi} + \phi]}{\delta\phi(x)} \right\rangle$$

## Diagrammatics

$$S[\phi] = \frac{1}{2} \int_x \left[ \partial_\mu \phi \partial_\mu \phi + m^2 \phi^2 + \frac{\lambda}{4} \phi^4 \right]$$



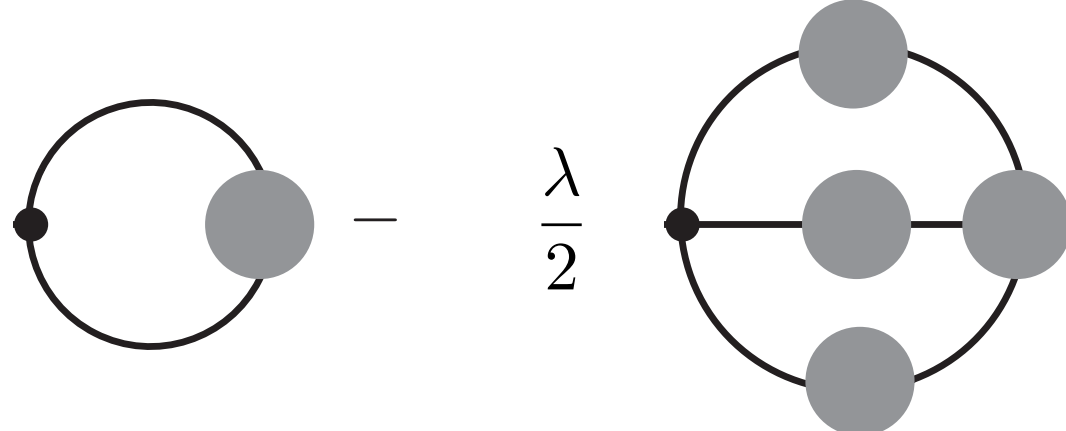
# Functional Renormalisation Group

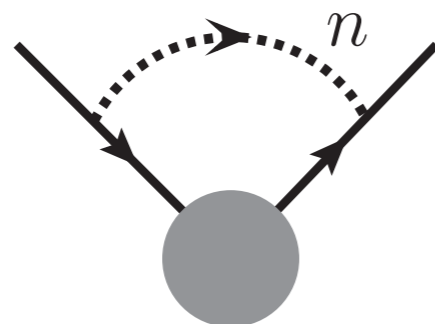
## Dyson-Schwinger equation

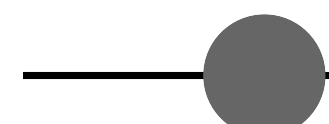
$$\frac{\delta\Gamma[\phi]}{\delta\phi(x)} = \left\langle \frac{\delta S[\hat{\varphi} + \phi]}{\delta\phi(x)} \right\rangle$$

## Diagrammatics

$$S[\phi] = \frac{1}{2} \int_x \left[ \partial_\mu \phi \partial_\mu \phi + m^2 \phi^2 + \frac{\lambda}{4} \phi^4 \right]$$

$$\frac{\lambda}{2} \langle [\hat{\varphi}(x) + \phi(x)]^3 \rangle = \frac{\lambda}{2} \phi^3(x) + \frac{3\lambda}{2} \phi(x) \text{ (loop with 1 vertex)} - \frac{\lambda}{2} \text{ (loop with 3 vertices)}$$


$$\Gamma^{(n)} = \text{Diagram with } n \text{ external lines and a loop}$$


$$\mathbf{G} = \text{Diagram with 2 external lines and a vertex} = \langle \hat{\varphi}(x) \hat{\varphi}(y) \rangle$$


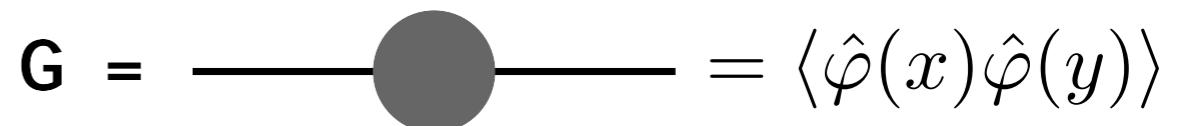
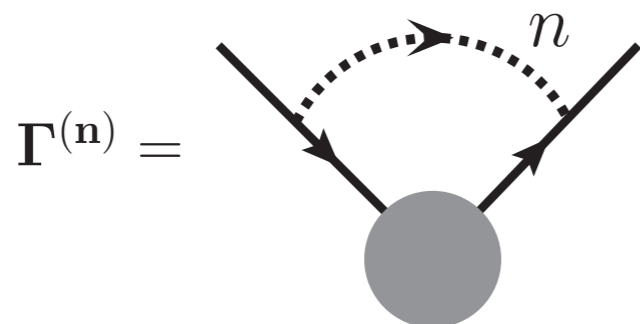
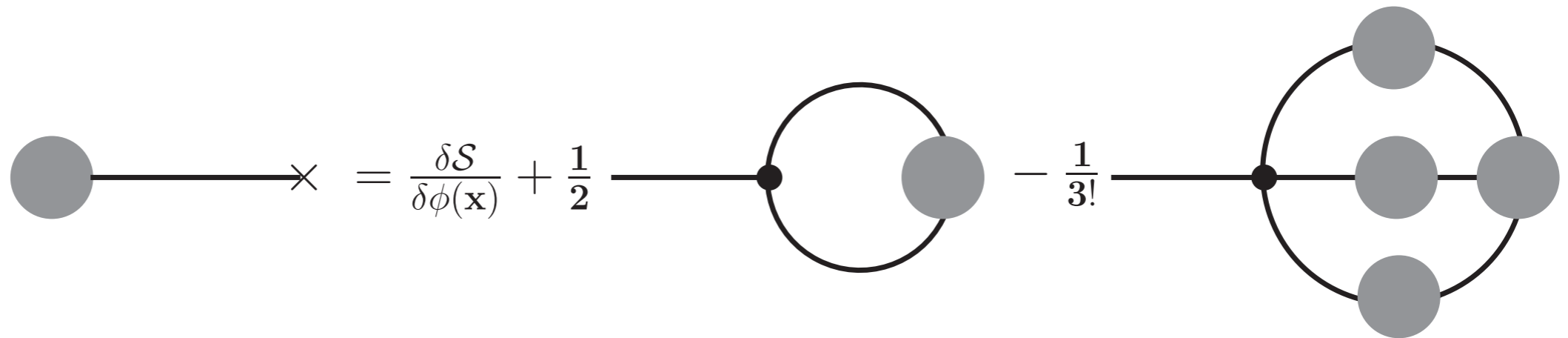
# Functional Renormalisation Group

## Dyson-Schwinger equation

$$\frac{\delta\Gamma[\phi]}{\delta\phi(x)} = \left\langle \frac{\delta S[\hat{\phi} + \phi]}{\delta\phi(x)} \right\rangle$$

## Diagrammatics

$$S[\phi] = \frac{1}{2} \int_x \left[ \partial_\mu \phi \partial_\mu \phi + m^2 \phi^2 + \frac{\lambda}{4} \phi^4 \right]$$



# Functional Renormalisation Group

---

## Effective action $\Gamma$

$$\Gamma[\phi] = -\log \int d\hat{\varphi} e^{-S[\hat{\varphi}+\phi] + \int_x \hat{\varphi} \frac{\delta\Gamma[\phi]}{\delta\phi}}$$

## No quantum fluctuations

$$\Gamma[\phi] = -\log e^{-S[\phi]} = S[\phi]$$

# Functional Renormalisation Group

## Effective action $\Gamma$

$$\Gamma[\phi] = -\log \int d\hat{\varphi} e^{-S[\hat{\varphi}+\phi] + \int_x \hat{\varphi} \frac{\delta\Gamma[\phi]}{\delta\phi}}$$

UV quantum fluctuations up to  $p^2 = k^2$



# Functional Renormalisation Group

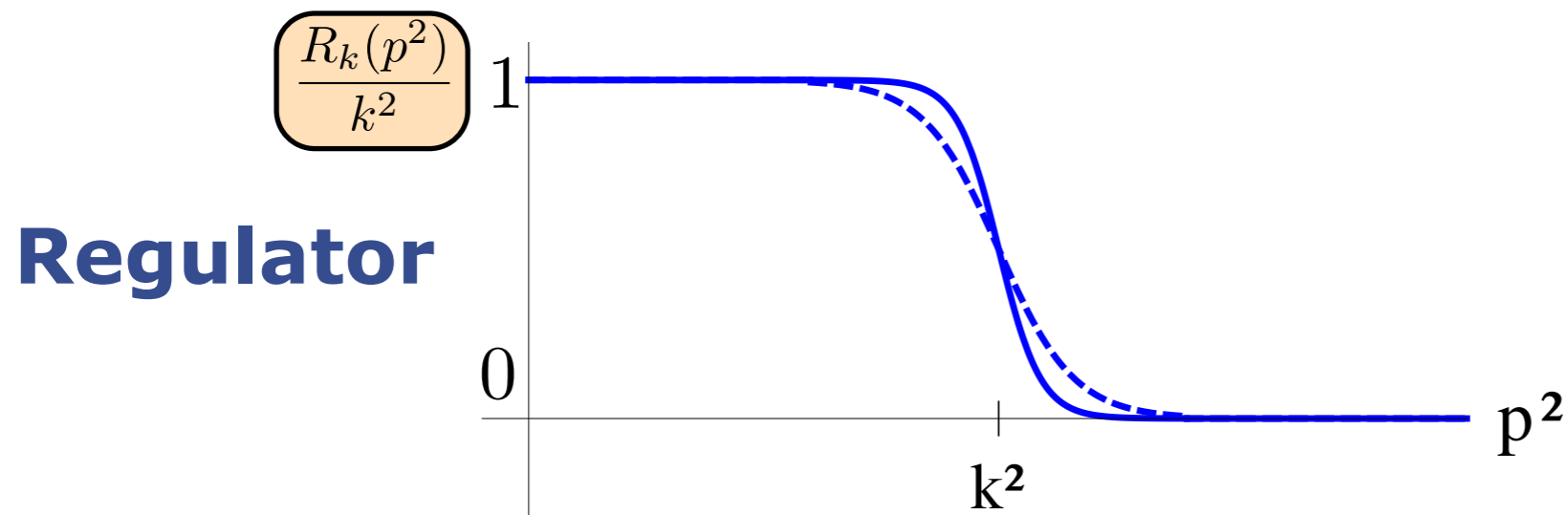
## Effective action $\Gamma_k$

$$\Gamma_k[\phi] = -\log \int d\hat{\varphi} e^{-S[\hat{\varphi}+\phi] + \frac{1}{2} \int_p \hat{\varphi}(p) R_k(p^2) \hat{\varphi}(-p) + \int_x \hat{\varphi} \frac{\delta \Gamma_k[\phi]}{\delta \phi}}$$

$$\frac{\delta \Gamma_k[\phi]}{\delta \phi(x)} = \left\langle \frac{\delta S[\hat{\varphi} + \phi]}{\delta \phi(x)} \right\rangle$$

**DSE**

UV quantum fluctuations up to  $p^2 = k^2$

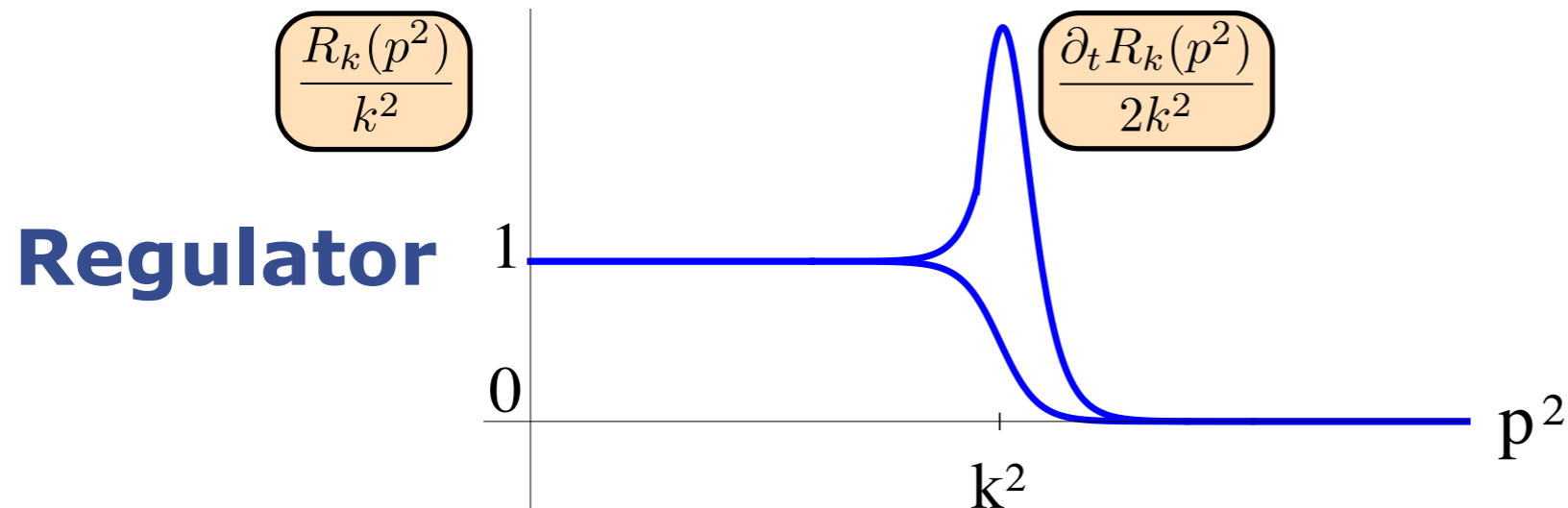
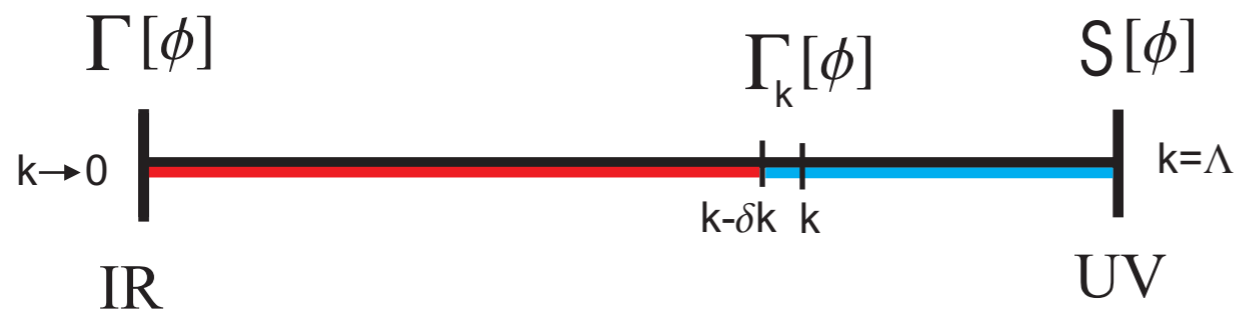


# Functional Renormalisation Group

## Effective action $\Gamma_k$

$$\Gamma_k[\phi] = -\log \int d\hat{\varphi} e^{-S[\hat{\varphi}+\phi] + \frac{1}{2} \int_p \hat{\varphi}(p) R_k(p^2) \hat{\varphi}(-p) + \int_x \hat{\varphi} \frac{\delta \Gamma_k[\phi]}{\delta \phi}}$$

## UV quantum fluctuations up to $p^2 = k^2$



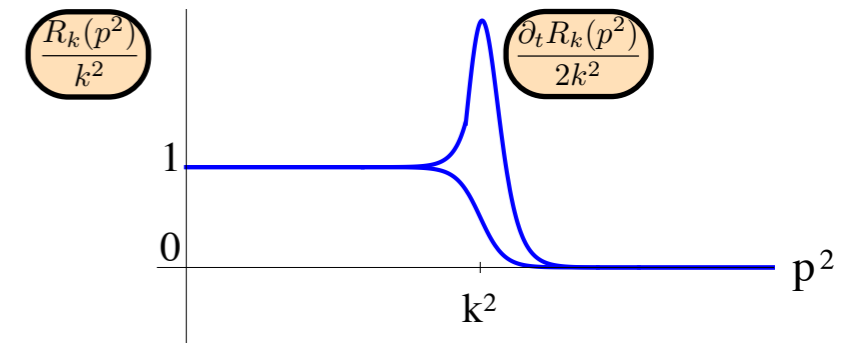
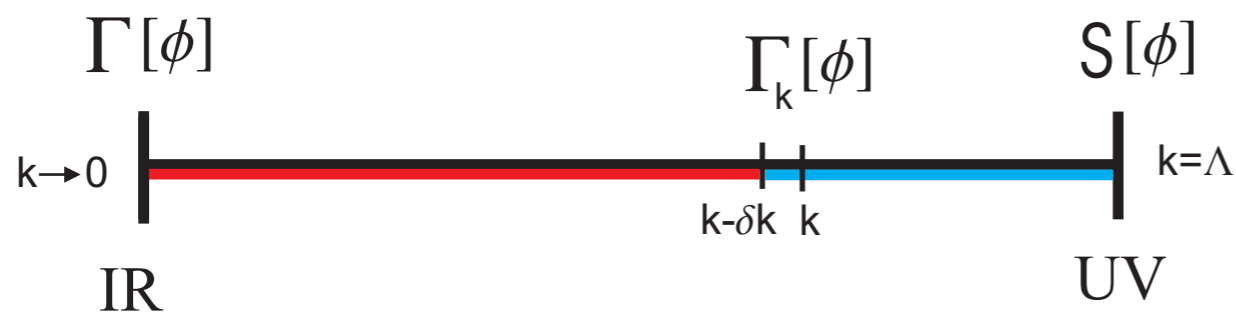
$$t = \log \frac{k}{\Lambda}$$

# Functional Renormalisation Group

## Effective action $\Gamma_k$

$$\Gamma_k[\phi] = -\log \int d\hat{\varphi} e^{-S[\hat{\varphi}+\phi] + \frac{1}{2} \int_p \hat{\varphi}(p) R_k(p^2) \hat{\varphi}(-p) + \int_x \hat{\varphi} \frac{\delta \Gamma_k[\phi]}{\delta \phi}}$$

## UV quantum fluctuations up to $p^2 = k^2$



## Flow

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \int \frac{d^4 p}{(2\pi)^4} \langle \hat{\varphi}(p) \hat{\varphi}(-p) \rangle \partial_t R_k(p^2)$$

$$t = \log \frac{k}{\Lambda}$$

# Functional Renormalisation Group

**Flow**

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \int \frac{d^4 p}{(2\pi)^4} \langle \hat{\varphi}(p) \hat{\varphi}(-p) \rangle \partial_t R_k(p^2)$$

$$t = \log \frac{k}{\Lambda}$$

**Propagator**

$$\mathbf{G} = \text{---} \bullet \text{---} = \langle \hat{\varphi}(x) \hat{\varphi}(y) \rangle$$



# Functional Renormalisation Group

## Flow

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \int \frac{d^4 p}{(2\pi)^4} \langle \hat{\varphi}(p) \hat{\varphi}(-p) \rangle \partial_t R_k(p^2)$$

$$t = \log \frac{k}{\Lambda}$$

## Propagator

$$G = \text{---} \bullet \text{---} = \langle \hat{\varphi}(x) \hat{\varphi}(y) \rangle$$

$$\frac{\delta \Gamma_k[\phi]}{\delta \phi(x)} = \left\langle \frac{\delta S[\hat{\varphi} + \phi]}{\delta \phi(x)} \right\rangle$$

## DSE

$$\Gamma_k^{(2)}[\phi] - S^{(2)}[\phi] = \frac{1}{2} \text{---} \bullet \text{---} + \frac{1}{2} \text{---} \bullet \text{---} \text{---} \bullet \text{---} + \frac{1}{3!} \text{---} \bullet \text{---} \text{---} \bullet \text{---} \text{---} \bullet \text{---}$$

$$\Gamma_k[\phi] = -\log \int d\hat{\varphi} e^{-S[\hat{\varphi} + \phi] + \frac{1}{2} \int_p \hat{\varphi}(p) R_k(p^2) \hat{\varphi}(-p) + \int_x \hat{\varphi} \frac{\delta \Gamma_k[\phi]}{\delta \phi}}$$

# Functional Renormalisation Group

## Flow

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \int \frac{d^4 p}{(2\pi)^4} \langle \hat{\varphi}(p) \hat{\varphi}(-p) \rangle \partial_t R_k(p^2)$$

$$t = \log \frac{k}{\Lambda}$$

## Propagator

$$G = \text{---} \bullet \text{---} = \langle \hat{\varphi}(x) \hat{\varphi}(y) \rangle$$

$$G^{-1}[\phi] = \Gamma_k^{(2)}[\phi] + R_k$$

$$\frac{\delta \Gamma_k[\phi]}{\delta \phi(x)} = \left\langle \frac{\delta S[\hat{\varphi} + \phi]}{\delta \phi(x)} \right\rangle$$

$$\Gamma_k[\phi] = -\log \int d\hat{\varphi} e^{-S[\hat{\varphi} + \phi] + \frac{1}{2} \int_p \hat{\varphi}(p) R_k(p^2) \hat{\varphi}(-p) + \int_x \hat{\varphi} \frac{\delta \Gamma_k[\phi]}{\delta \phi}}$$

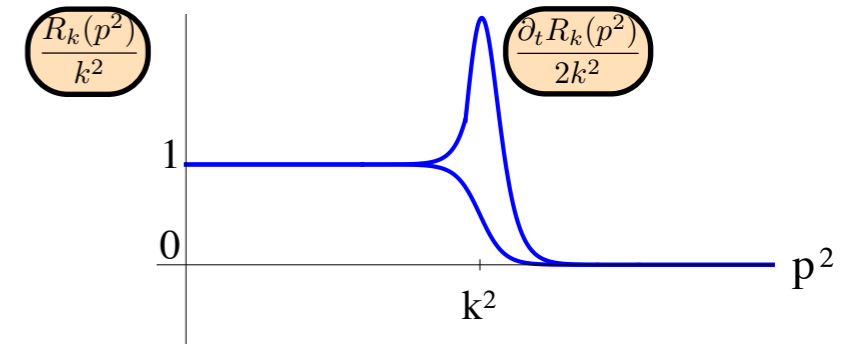
## DSE

$$\Gamma_k^{(2)}[\phi] - S^{(2)}[\phi] = \frac{1}{2} \text{---} \bullet \text{---} + \frac{1}{2} \text{---} \bullet \text{---} \bullet \text{---} + \frac{1}{3!} \text{---} \bullet \text{---} \bullet \text{---} \bullet \text{---}$$

# Functional Renormalisation Group

## Flow

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$



## Diagrammatics

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \left[ \text{Diagram} \right]$$

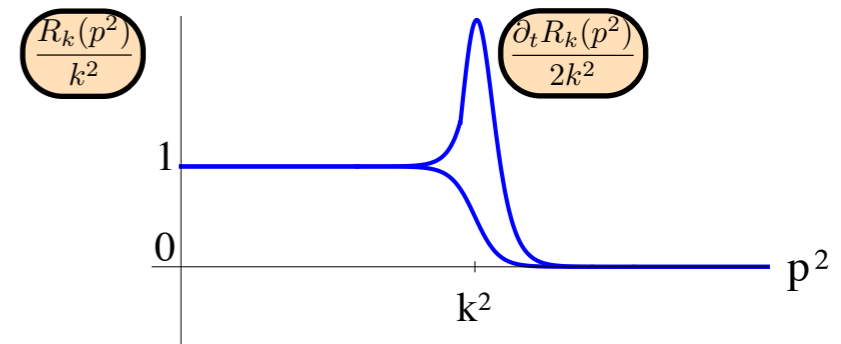
## Propagator

$$\partial_t \Gamma_k^{(2)}[\phi] = -\frac{1}{2} \frac{\delta}{\delta \phi} \left[ \text{Diagram 1} \right] = -\frac{1}{2} \left[ \text{Diagram 2} \right] + \left[ \text{Diagram 3} \right]$$

# Functional Renormalisation Group

## Flow

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$



## Propagator

$$\partial_t \Gamma_k^{(2)}[\phi] = -\frac{1}{2} \frac{\delta}{\delta \phi} \left[ \text{Diagram 1} \right] = -\frac{1}{2} \left[ \text{Diagram 2} \right] + \left[ \text{Diagram 3} \right]$$

The diagrams represent Feynman diagrams for the two-point function. Diagram 1 is a circle with a cross on top and a vertical line at the bottom. Diagram 2 is a circle with a cross on top and two lines meeting at the bottom. Diagram 3 is a circle with a cross on top and two lines extending from the right side.

FRG

DSE

$$\Gamma^{(2)}[\phi] - S^{(2)}[\phi] = \frac{1}{2} \left[ \text{Diagram 4} \right] + \frac{1}{2} \left[ \text{Diagram 5} \right] + \frac{1}{3!} \left[ \text{Diagram 6} \right]$$

The diagrams represent Feynman diagrams for the two-point function. Diagram 4 is a tadpole diagram with a loop and a vertical line. Diagram 5 is a circle with two lines extending from the right side. Diagram 6 is a circle with two lines extending from the right side and a vertical line at the bottom.

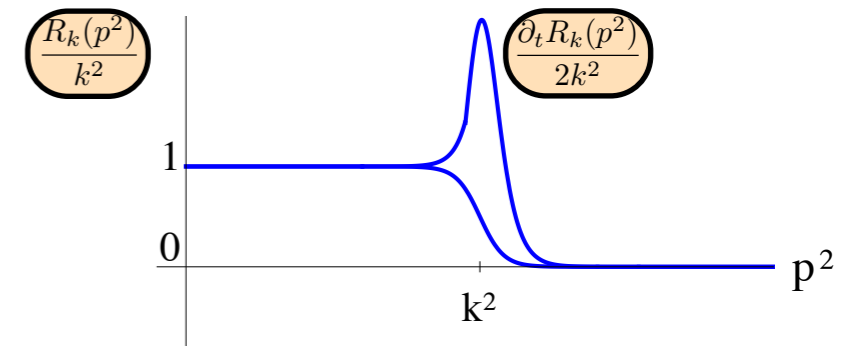
$$\partial_t \Gamma^{(n)} = \text{Flow}_n[\Gamma^{(m)}; m = 2, \dots, n + 2]$$

$$\Gamma^{(n)} = \text{DSE}_n[S^{(m)}, \Gamma^{(m)}; m = 2, \dots, n + 2]$$

# Functional Renormalisation Group

## Flow

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$



## Properties

	FRG	DSE	2PI	3PI	4PI
• 1-loop exact	✓	—			
• closed	✓	✓			
• RG-scaling	✓	—	—	—	✓
• Energy/particle-number conserv.	—	—	✓	✓	✓

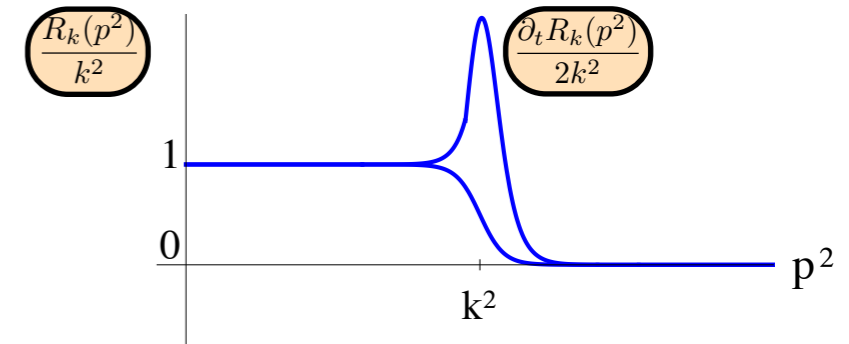
✓ automatic

— only in specific approximation schemes

# Functional Renormalisation Group

## Flow

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$



## Properties

### FunMethods

- 1-loop exact ✓
- closed ✓
- RG-scaling ✓
- Energy/particle-number conserv. ✓

✓ automatic

— only in specific approximation schemes

# Approximation schemes

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$

$$\partial_t \Gamma^{(n)} = \text{Flow}_n[\Gamma^{(m)}; m = 2, \dots, n + 2]$$

## Derivative expansion

- Expansion in powers of momenta
- controlled in the presence of a mass gap  $m_{\text{gap}}$

- Expansion parameter

$$\frac{p^2}{\max(k^2, m_{\text{gap}}^2)}$$

## Vertex expansion

- Expansion in number  $n$  of external fields
- controlled in perturbation theory/presence of symmetries
- Expansion parameter  $n$

Mixtures, exact resummation schemes, ....

# Approximation schemes

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$

## Derivative expansion

- Expansion in powers of momenta
- controlled in the presence of a mass gap  $m_{\text{gap}}$

- Expansion parameter

$$\frac{p^2}{\max(k^2, m_{\text{gap}}^2)}$$

## Lowest order: 0th order

$$\Gamma_k[\phi] = \frac{1}{2} \int_p \phi p^2 \phi + \int_x V_k(\phi) + O(p^2)$$

$$\Gamma_k^{(2)}[\phi](p, q) = (p^2 + V_k''(\phi)) (2\pi)^d \delta(p - q)$$



# Approximation schemes

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$

## Derivative expansion

Lowest order: 0th order

$$\Gamma_k[\phi] = \frac{1}{2} \int_p \phi p^2 \phi + \int_x V_k(\phi) + O(p^2)$$

$$\Gamma_k^{(2)}[\phi](p, q) = (p^2 + V_k''(\phi)) (2\pi)^d \delta(p - q)$$

$$R_{k,\text{opt}}(p^2) = (k^2 - p^2) \theta(k^2 - p^2)$$

$$\partial_t R_{k,\text{opt}}(p^2) = 2k^2 \theta(k^2 - p^2)$$

Flow

$$\partial_t V_k[\phi] = \frac{1}{2d} \frac{\Omega_d}{(2\pi)^d} k^d \frac{k^2}{k^2 + V''(\phi)}$$

$$\Omega_d = \frac{2\pi^{d/2}}{\Gamma(d/2)}$$

# Approximation schemes

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$

## Derivative expansion

Lowest order: 0th order

$$\Gamma_k[\phi] = \frac{1}{2} \int_p \phi p^2 \phi + \int_x V_k(\phi) + O(p^2)$$

$$\Gamma_k^{(2)}[\phi](p, q) = (p^2 + V_k''(\phi)) (2\pi)^d \delta(p - q)$$

$$\Gamma^{(2)}[\phi](p) + R_{k,\text{opt}}(p^2) = [k^2 + V''(\phi)] \theta(k^2 - p^2) + (p^2 + V''(\phi)) \theta(p^2 - k^2)$$

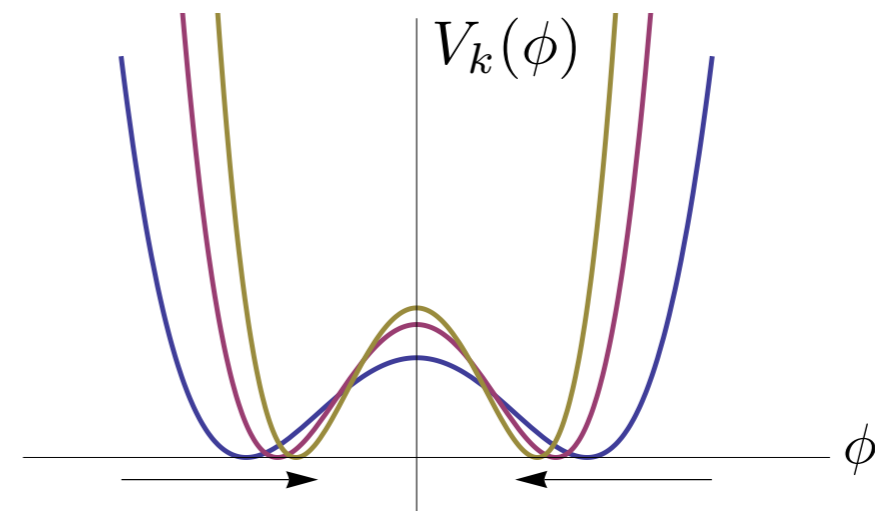
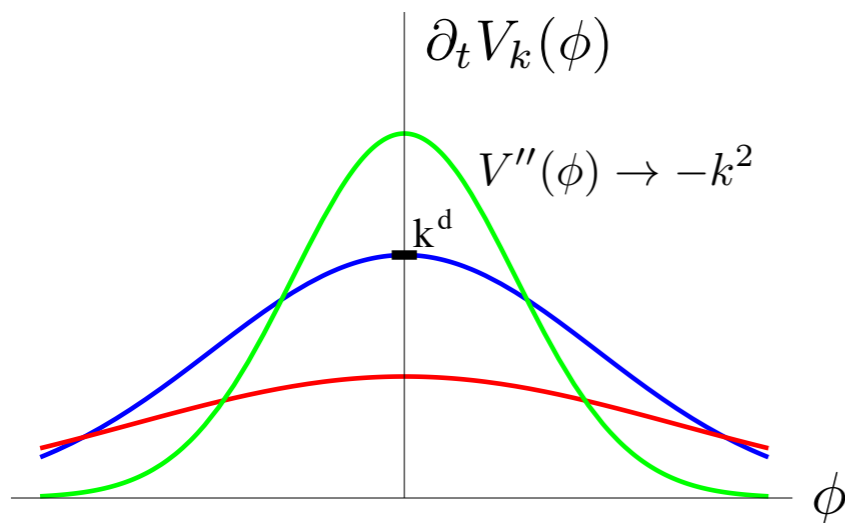
Flow

$$\partial_t V_k[\phi] = \frac{1}{2d} \frac{\Omega_d}{(2\pi)^d} k^d \frac{k^2}{k^2 + V''(\phi)}$$

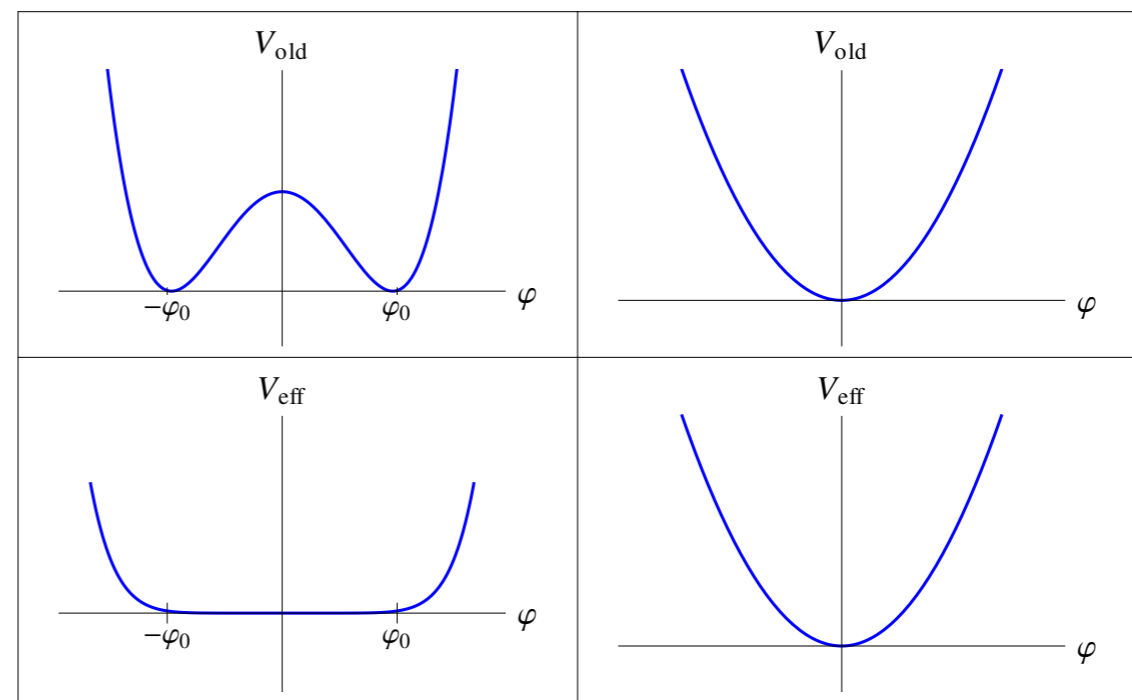
$$\Omega_d = \frac{2\pi^{d/2}}{\Gamma(d/2)}$$

# Approximation schemes & phase structure

$$\partial_t V_k[\phi] = \frac{1}{2d} \frac{\Omega_d}{(2\pi)^d} k^d \frac{k^2}{k^2 + V''(\phi)}$$

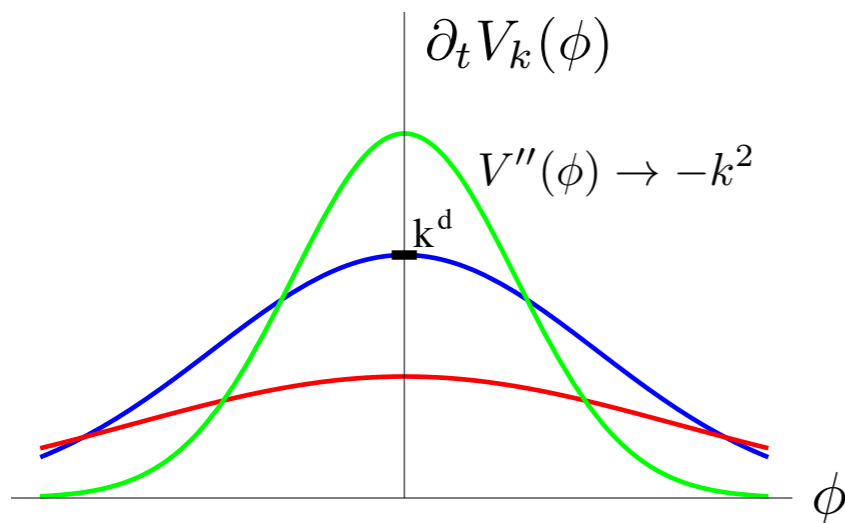


- **bosonic flow is symmetry-restoring**
- **flow guarantees convexity**

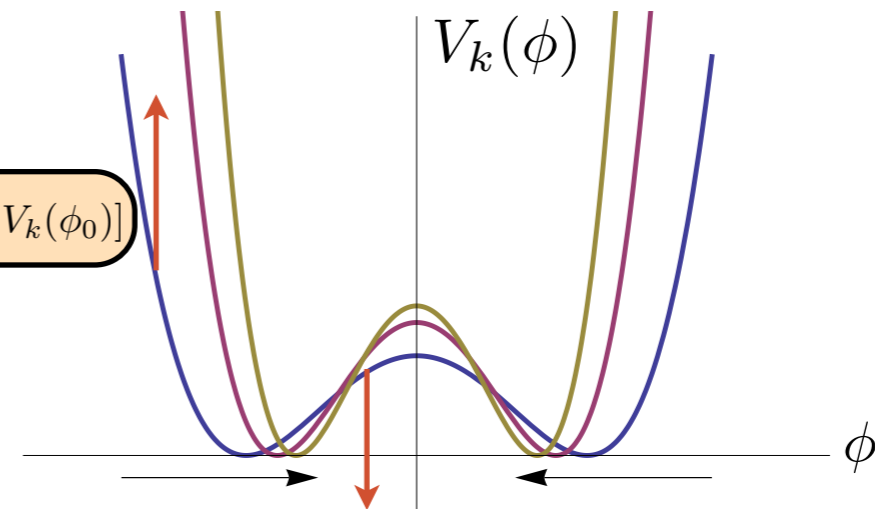


# Approximation schemes & phase structure

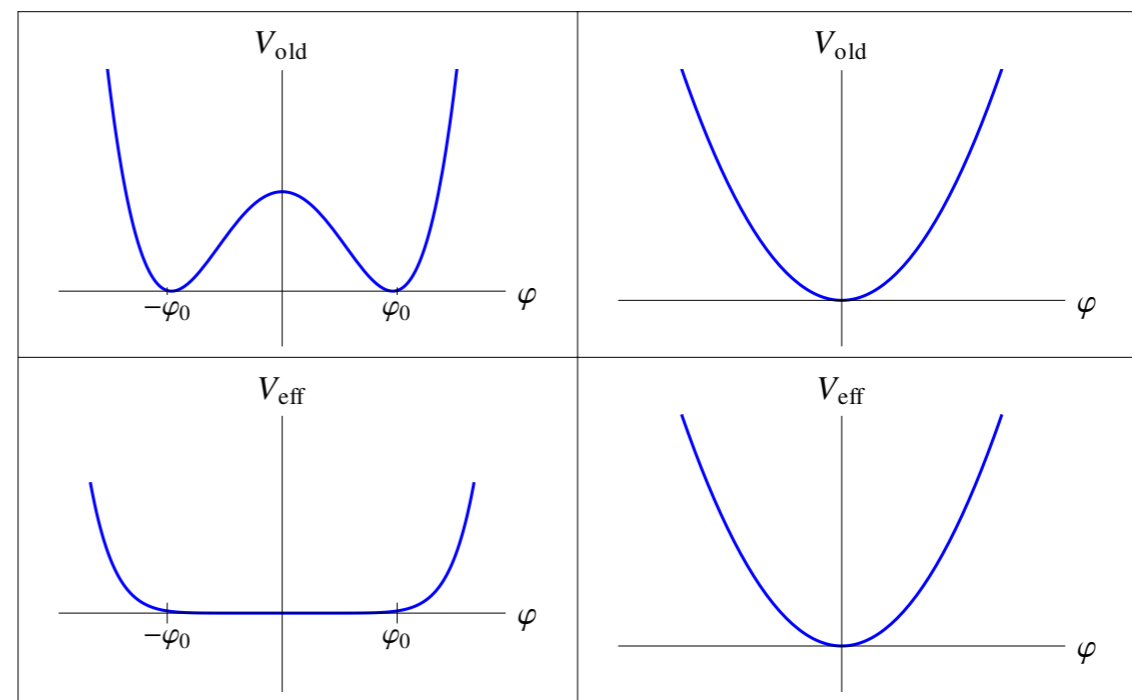
$$\partial_t V_k[\phi] = \frac{1}{2d} \frac{\Omega_d}{(2\pi)^d} k^d \frac{k^2}{k^2 + V''(\phi)}$$



$$-\frac{\Delta k}{k} [\partial_t V_k(\phi) - \partial_t V_k(\phi_0)]$$

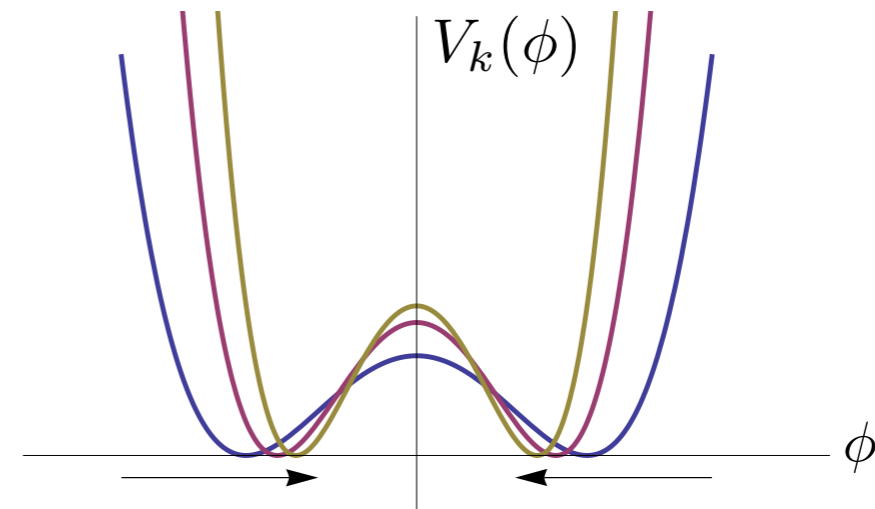
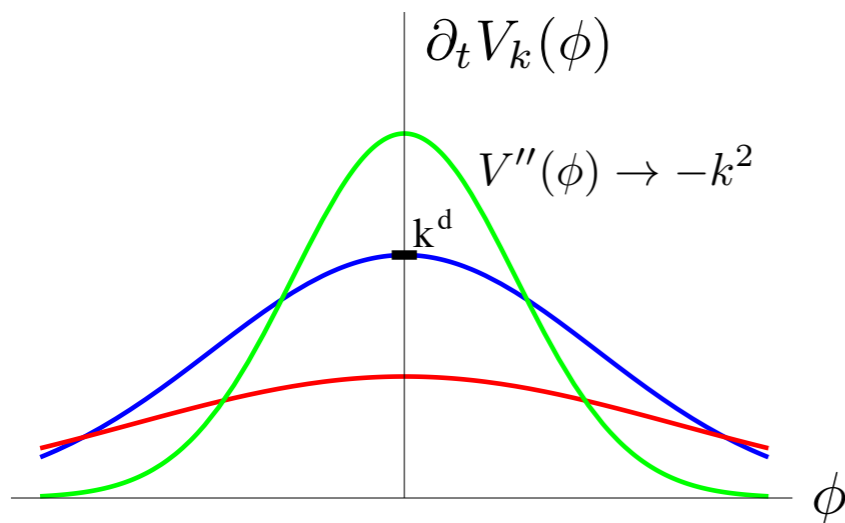


- bosonic flow is symmetry-restoring
- flow guarantees convexity



# Approximation schemes & phase structure

$$\partial_t V_k[\phi] = \frac{1}{2d} \frac{\Omega_d}{(2\pi)^d} k^d \frac{k^2}{k^2 + V''(\phi)}$$



- bosonic flow is symmetry-restoring
- flow guarantees convexity of effective action

Litim, JMP, Vergara '06

## Example: 3d critical exponents with FRG

$$\Gamma_k[\phi] = \frac{1}{2} \int_p Z_k \phi p^2 \phi + \int_x V_k(\phi)$$

$$V_k(\phi) = \sum_{n=1}^{N_{\max}} \frac{\lambda_n}{n!} (\phi^2 - \phi_{0,k}^2)^n$$

$$N = 1 : \nu_{\text{Ising}} = 0.630\dots$$

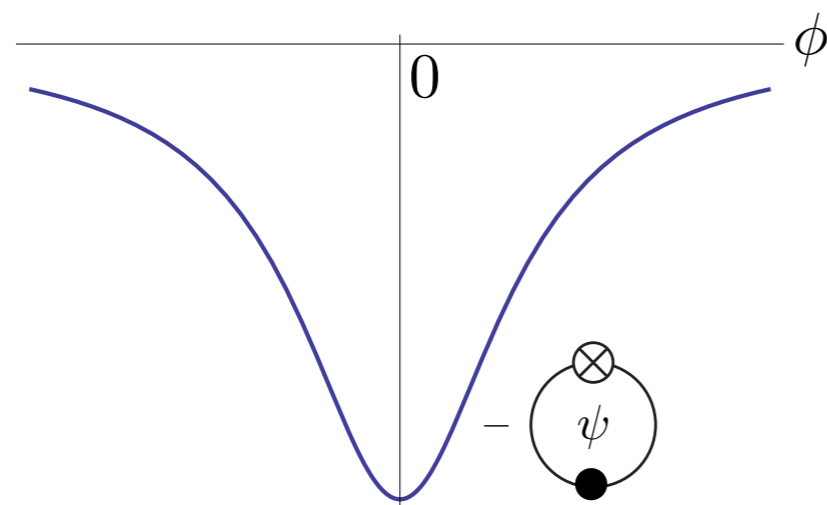
$$N = 1 : \nu_{\text{Ising}} = 0.637\dots$$

A simple program to compute critical exponents in O(N)-models with the Wetterich equation

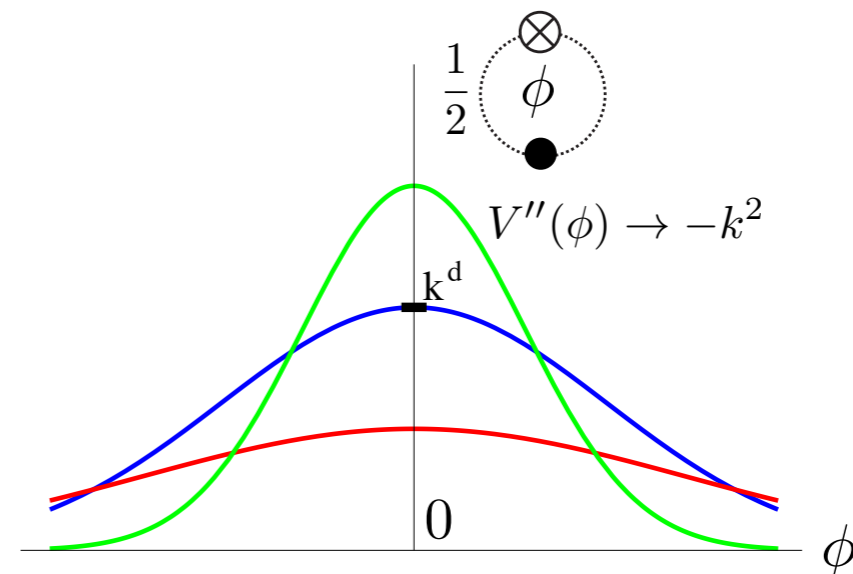
Michael Scherer

# Approximation schemes & phase structure

$$\partial_t V_k(\phi) = - \text{[Feynman diagram with } \psi \text{ loop]} + \frac{1}{2} \text{[Feynman diagram with } \phi \text{ loop]}$$



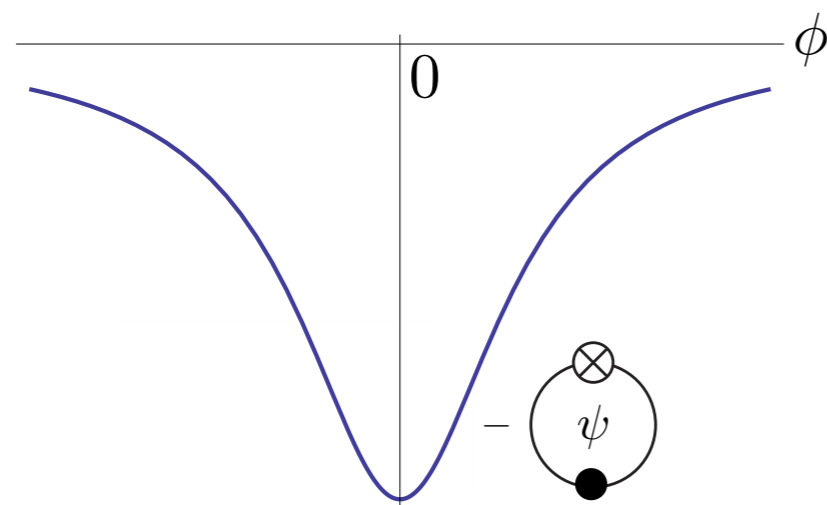
+



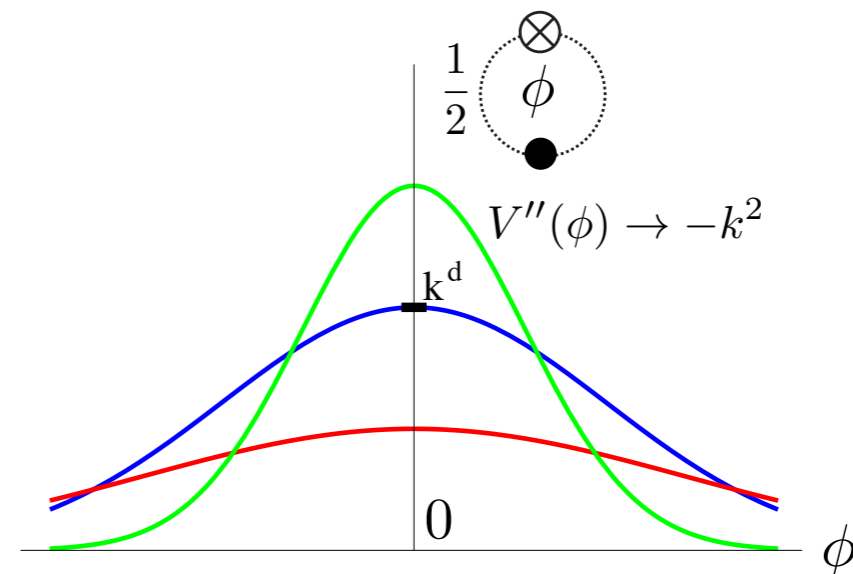
- bosonic flow is symmetry-restoring
- fermionic flow is symmetry-breaking
- competing dynamics decides about fate of symmetries
- flow guarantees convexity

# Approximation schemes & phase structure

$$\partial_t V_k(\phi) = - \text{[diagram: circle with } \psi \text{ and } \otimes \text{]} + \frac{1}{2} \text{[diagram: circle with } \phi \text{ and } \otimes \text{]}$$



+



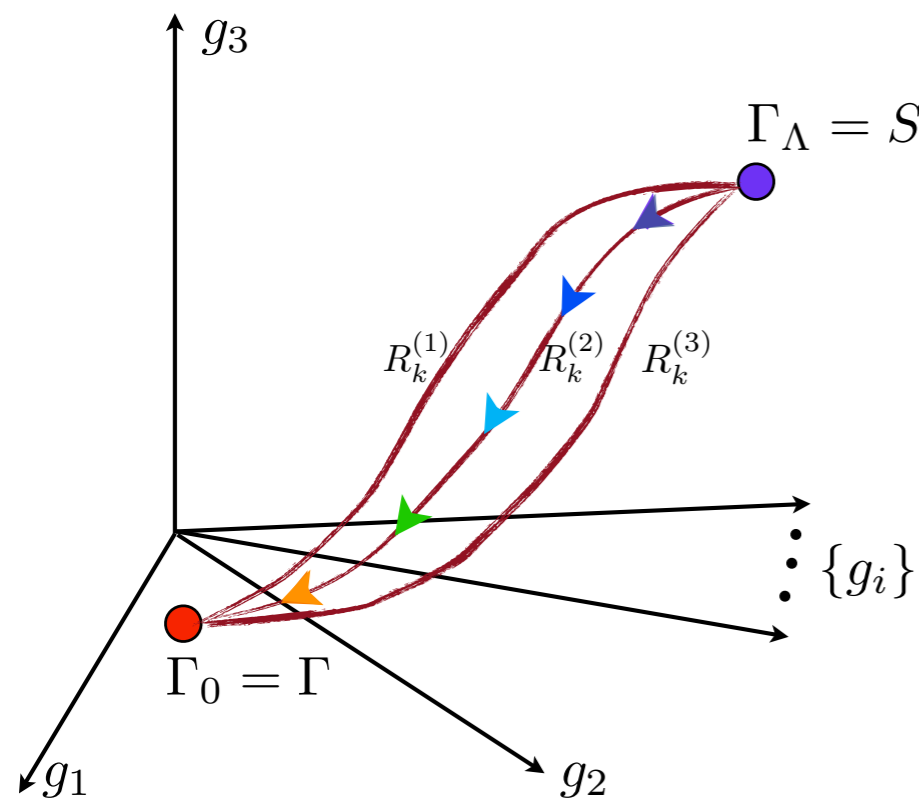
- bosonic flow is symmetry-restoring
- fermionic flow is symmetry-breaking
- competing dynamics decides about fate of symmetries
- flow guarantees convexity

'governs general phase structures'

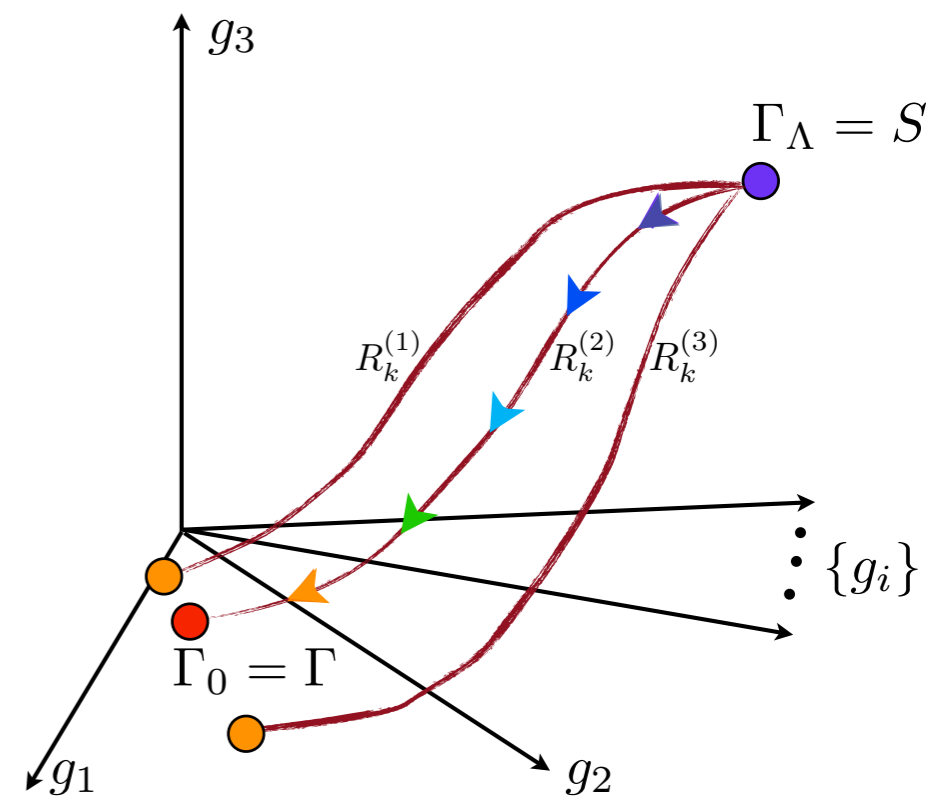
# Approximation schemes & error control

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$

Theory space



full flow



approximated flow

Optimisation: find  $R_k^{(2)}$  !

Litim '01: most rapid convergence

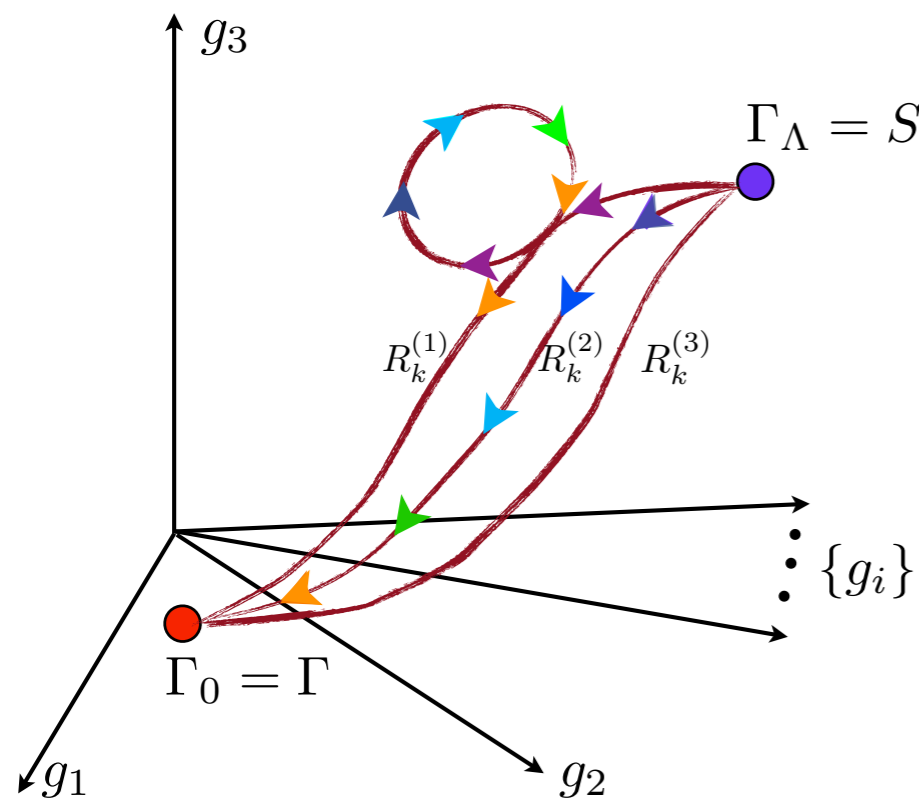
JMP '05: integrability



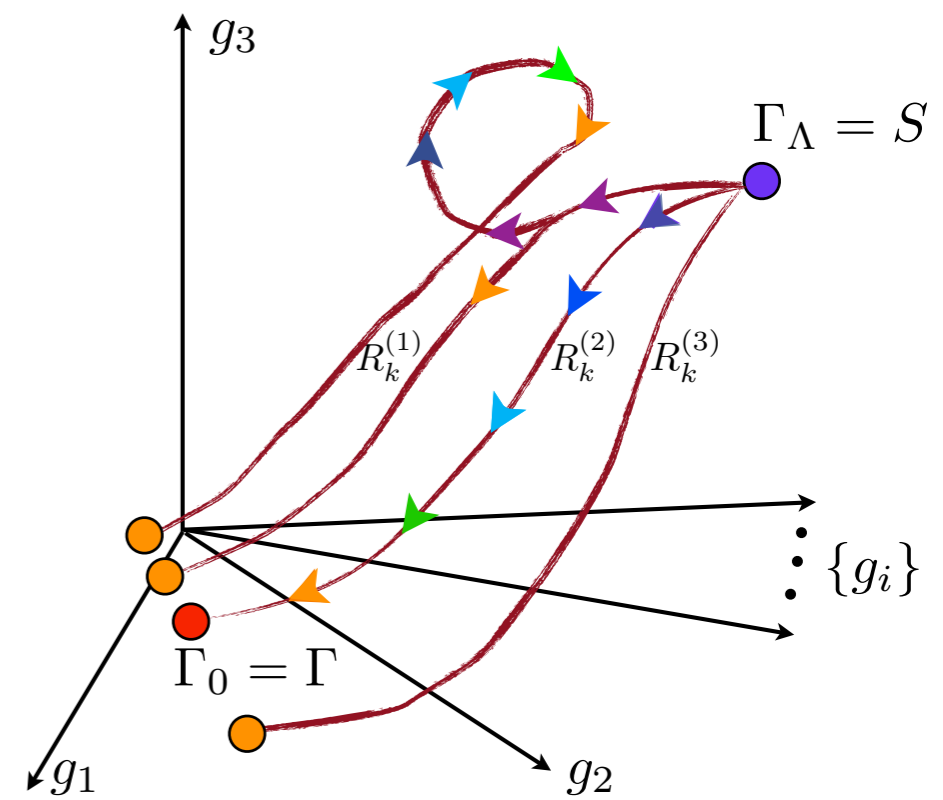
# Approximation schemes & error control

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$

Theory space



full flow



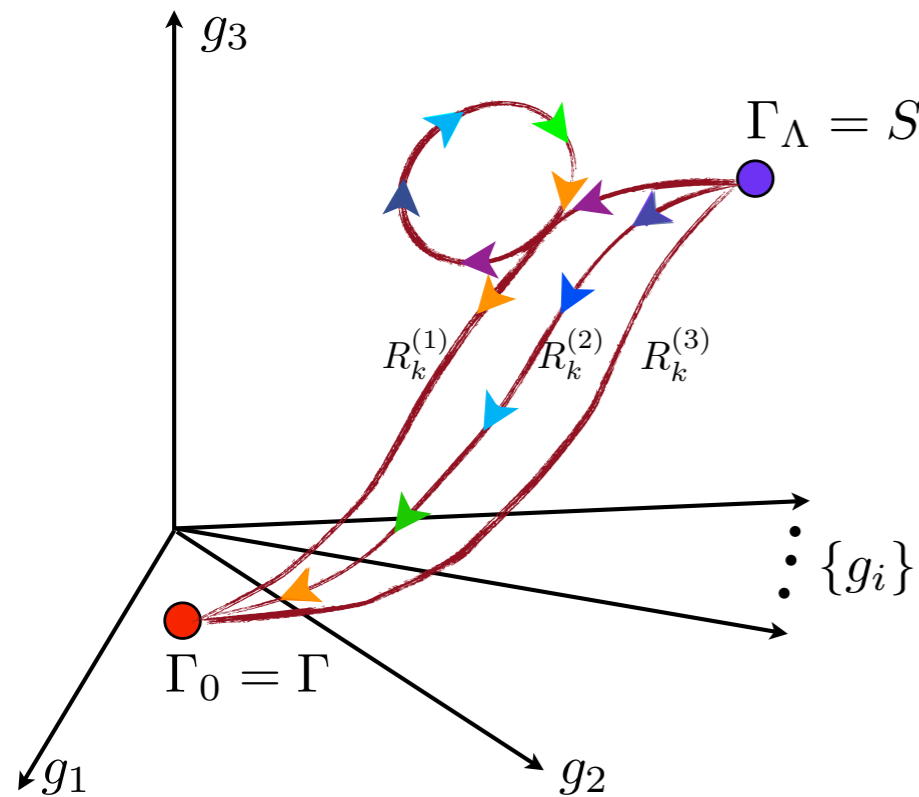
approximated flow

Optimisation: find  $R_k^{(2)}$  !

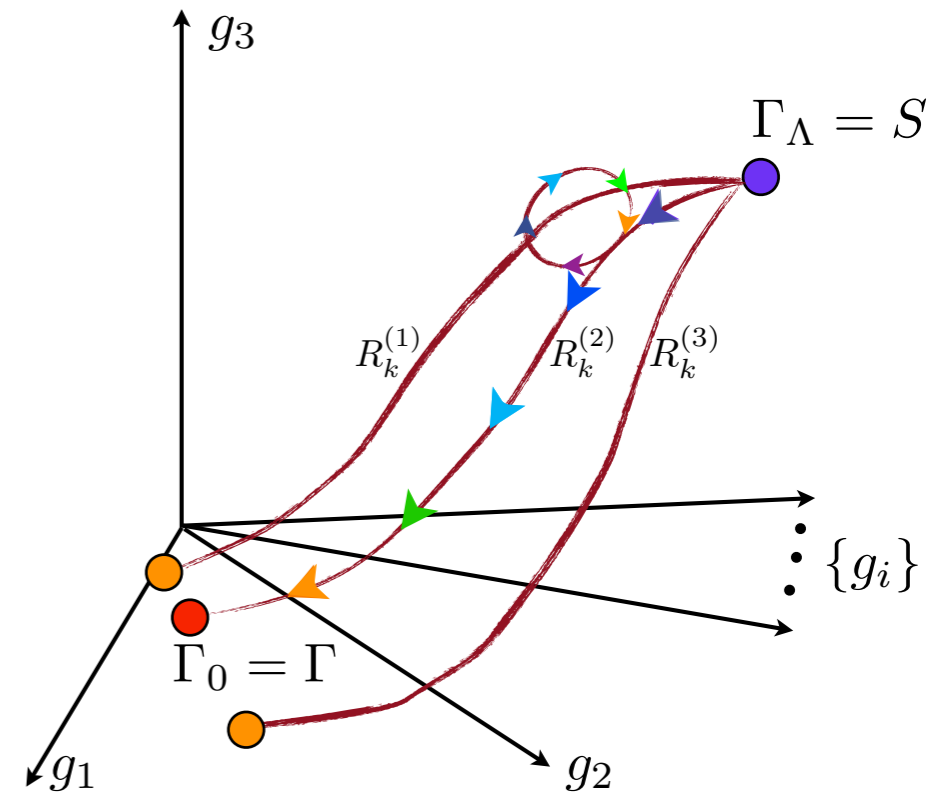
# Approximation schemes & error control

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k$$

Theory space



full flow



optimised flow

Optimisation: find  $R_k^{(2)}$  !

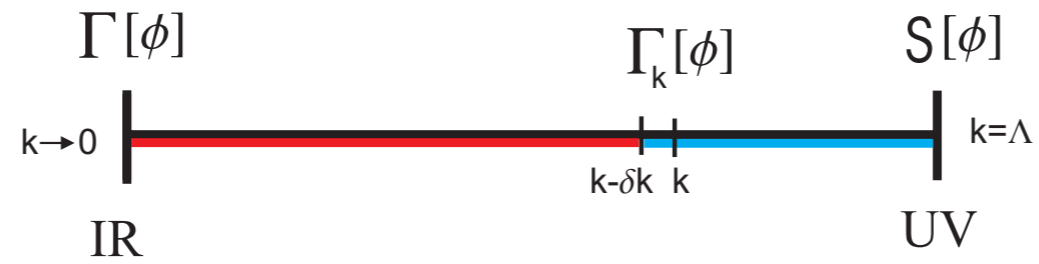
$$\lim_{L \rightarrow 0} \frac{1}{L} \text{cycle} \rightarrow 0$$

# FRG for QCD

# Functional RG for QCD

eg. JMP, AIP Conf.Proc. 1343 (2011)  
NPA 931 (2014) 113

free energy at momentum scale  $k$



**ab initio**

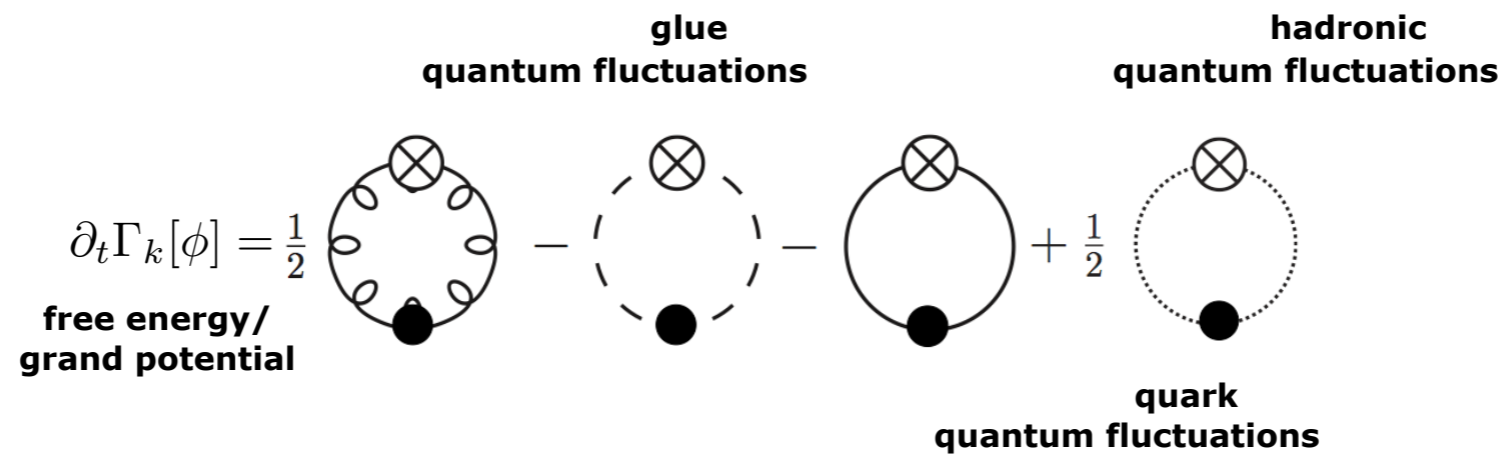
# Functional RG for QCD

eg. JMP, AIP Conf.Proc. 1343 (2011)  
NPA 931 (2014) 113

free energy at momentum scale  $k$



**ab initio**



RG-scale  $k$ :  $t = \ln k$

**closed form**

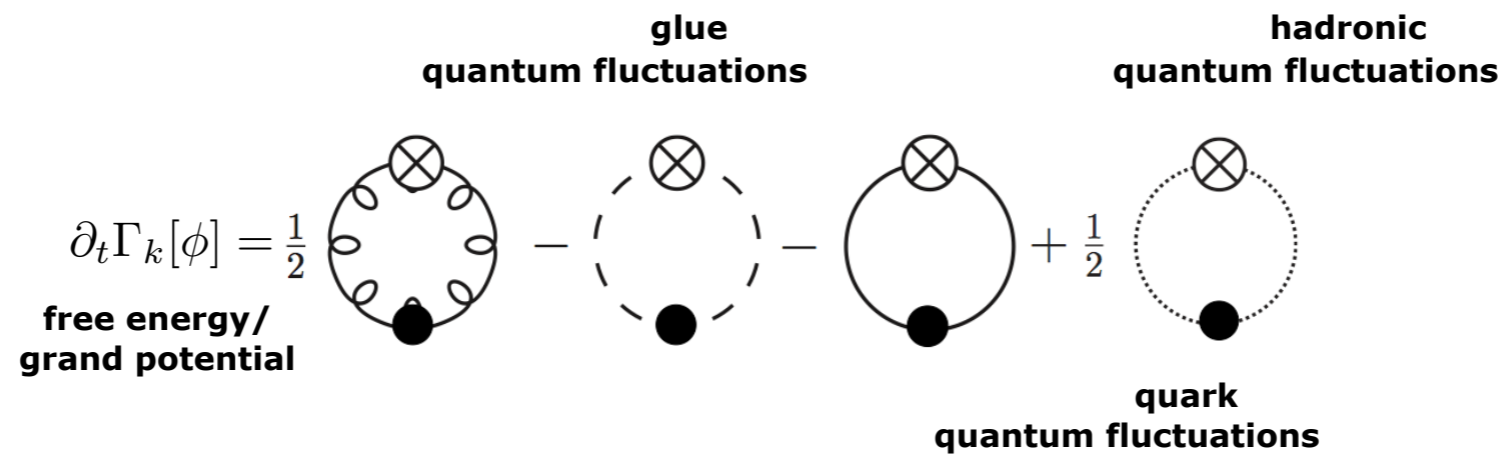
# Functional RG for QCD

eg. JMP, AIP Conf.Proc. 1343 (2011)  
NPA 931 (2014) 113

free energy at momentum scale  $k$




**ab initio**



RG-scale  $k$ :  $t = \ln k$

**properties**

**closed form**

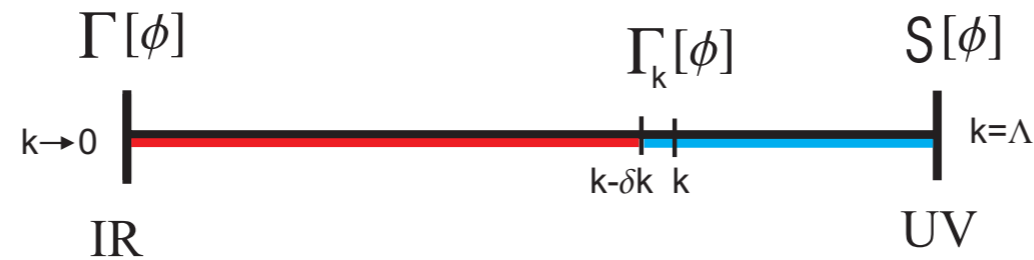
- **access to physics** 
- **numerically tractable, also at real time**  
**no sign problem**  
**systematic error control via closed form**
- **low energy models naturally incorporated**



# Functional RG for QCD

eg. JMP, AIP Conf.Proc. 1343 (2011)  
NPA 931 (2014) 113

free energy at momentum scale  $k$



**ab initio**

**glue quantum fluctuations**
**hadronic quantum fluctuations**

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \left( \text{glue loop} - \text{ghost loop} - \text{quark loop} \right) + \frac{1}{2} \left( \text{hadronic loop} \right)$$

free energy/  
grand potential
**quark quantum fluctuations**

RG-scale  $k$ :  $t = \ln k$

**closed form**

functional DSE :

$$\frac{\delta(\Gamma - S)}{\delta A_0} = \frac{1}{2} \left( \text{glue loop} - \text{ghost loop} - \text{quark loop} \right) - \frac{1}{6} \left( \text{hadronic loop} \right) + \dots$$

$A_0$  : background field

# Functional RG for QCD

**fQCD collaboration:** J. Braun, L. Corell, A. Cyrol, W.-j. Fu, M. Leonhardt, M. Mitter, JMP, M. Pospiech, F. Rennecke, N. Wink

Heidelberg, Dalian, Darmstadt

## Agenda

### QCD at finite T & $\mu$

Phase structure

Fluctuations

Phenomenology

### Real time correlation functions

Hadron spectrum & decays

Transport coefficients

Dynamics



# Functional RG for QCD

**fQCD collaboration:** J. Braun, L. Corell, A. Cyrol, W.-j. Fu, M. Leonhardt, M. Mitter, JMP, M. Pospiech, F. Rennecke, N. Wink

Heidelberg, Dalian, Darmstadt

## Agenda

### QCD at finite $T$ & $\mu$

Phase structure

Fluctuations

Phenomenology

### Real time correlation functions

Hadron spectrum & decays

Transport coefficients

Dynamics

### Selection of papers

**quenched QCD:**

Cyrol, Fister, Mitter, JMP, Strodthoff, PRD 94 (2016) 054005

**unquenched QCD:**

Braun, Fister, Haas, JMP, Rennecke, PRD 94 (2016) 034016

Cyrol, Mitter, JMP, Strodthoff, PRD 97 (2018) 054006

*vector mesons:* Rennecke, PRD 92 (2015) 076012

**pure glue:**

Mitter, JMP, Strodthoff, PRD 91 (2015) 054035

*finite  $T$ :* Cyrol, Mitter, JMP, Strodthoff, PRD 97 (2018) 054015

**finite density:** *fluctuations:* Fu, JMP, Schaefer, Rennecke, PRD 94 (2016) 11, 116020

*phase structure:* Braun, Leonhardt, Pospiech, PRD 96 (2017) 7, 076003

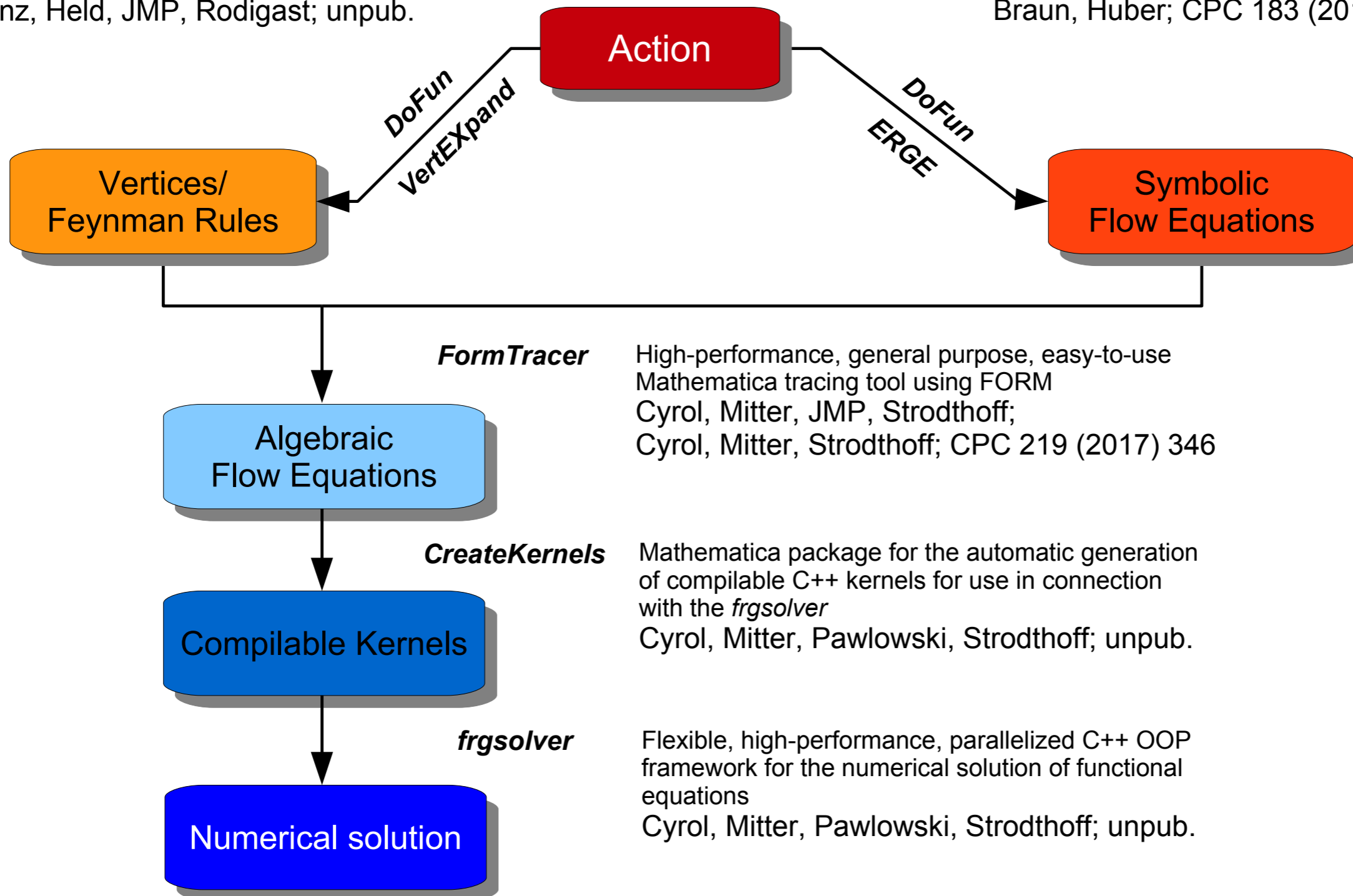
# fOCD: workflow

## **VertEXpand**

Mathematica package for the derivation of vertices from a given action using FORM  
Denz, Held, JMP, Rodigast; unpub.

## **DoFun**

Mathematica package for the derivation of functional equations  
Braun, Huber; CPC 183 (2012) 1290



### **FormTracer**

High-performance, general purpose, easy-to-use  
Mathematica tracing tool using FORM  
Cyrol, Mitter, JMP, Strodthoff;  
Cyrol, Mitter, Strodthoff; CPC 219 (2017) 346

### **CreateKernels**

Mathematica package for the automatic generation  
of compilable C++ kernels for use in connection  
with the *frgsolver*  
Cyrol, Mitter, Pawlowski, Strodthoff; unpub.

### **frgsolver**

Flexible, high-performance, parallelized C++ OOP  
framework for the numerical solution of functional  
equations  
Cyrol, Mitter, Pawlowski, Strodthoff; unpub.

GEFÖRDERT VOM



**FWF**

Der Wissenschaftsfonds.

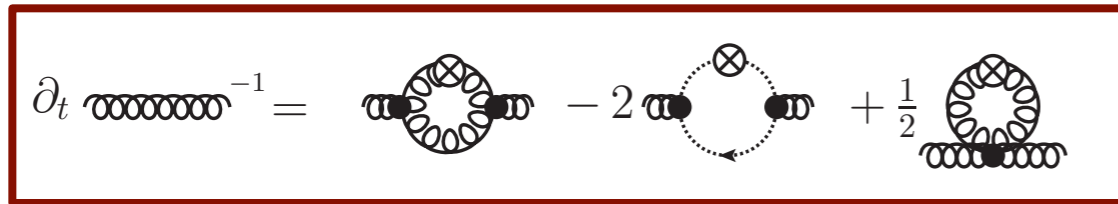


European Research Council

Established by the European Commission

# YM-theory: gluonic correlation functions

$$\langle A A \rangle(p^2)$$

$$\partial_t \text{gluon}^{-1} = \text{gluon loop} - 2 \text{ghost loop} + \frac{1}{2} \text{gluon tadpole}$$


# YM-theory: gluonic correlation functions

$$\partial_t \text{---}^{-1} = \text{---} \circlearrowleft + \text{---} \circlearrowright$$

$$\partial_t \text{---}^{-1} = \text{---} \circlearrowleft - 2 \text{---} \circlearrowright + \frac{1}{2} \text{---} \circlearrowleft$$

$$\partial_t \text{---} = - \text{---} - \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} + 2 \text{---} - \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} - \text{---} + 2 \text{---} - \text{---} + \text{perm.}$$

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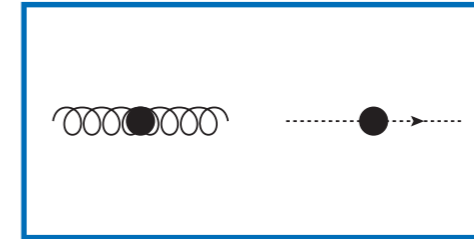
$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---}$$

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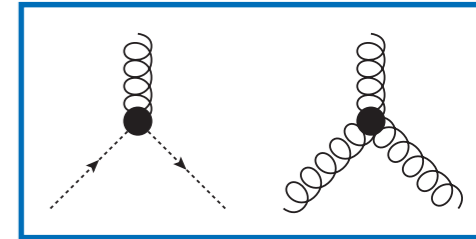
$$\partial_t \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

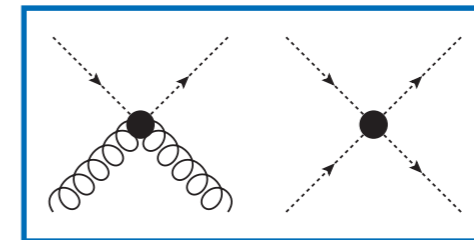
$$\partial_t \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$



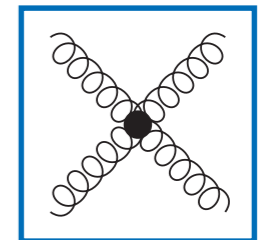
full. mom. dep.



full. mom. dep.  
classical tensor structures



mom. dep. needed by tadpoles  
full tensor basis



sym. point mom. dep. and  
mom. dep. needed by tadpole  
classical tensor structure

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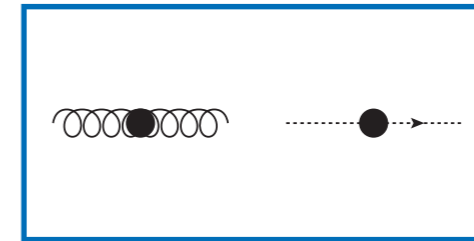
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$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} - 2 \text{---} \text{---} \text{---} + \frac{1}{2} \text{---} \text{---} \text{---}$$

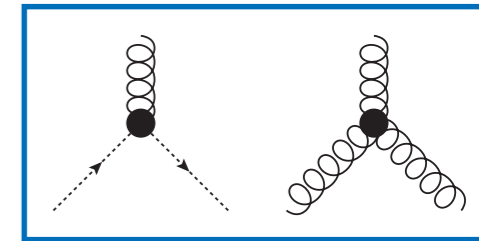
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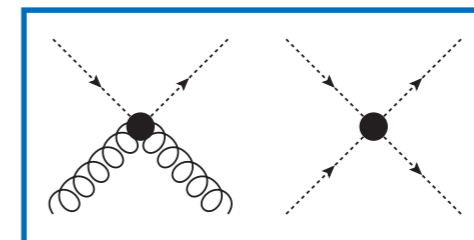
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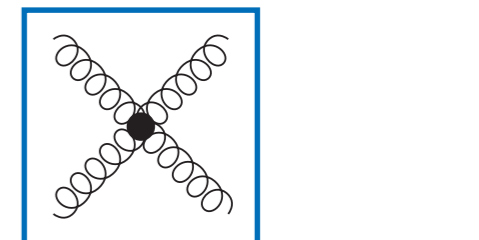
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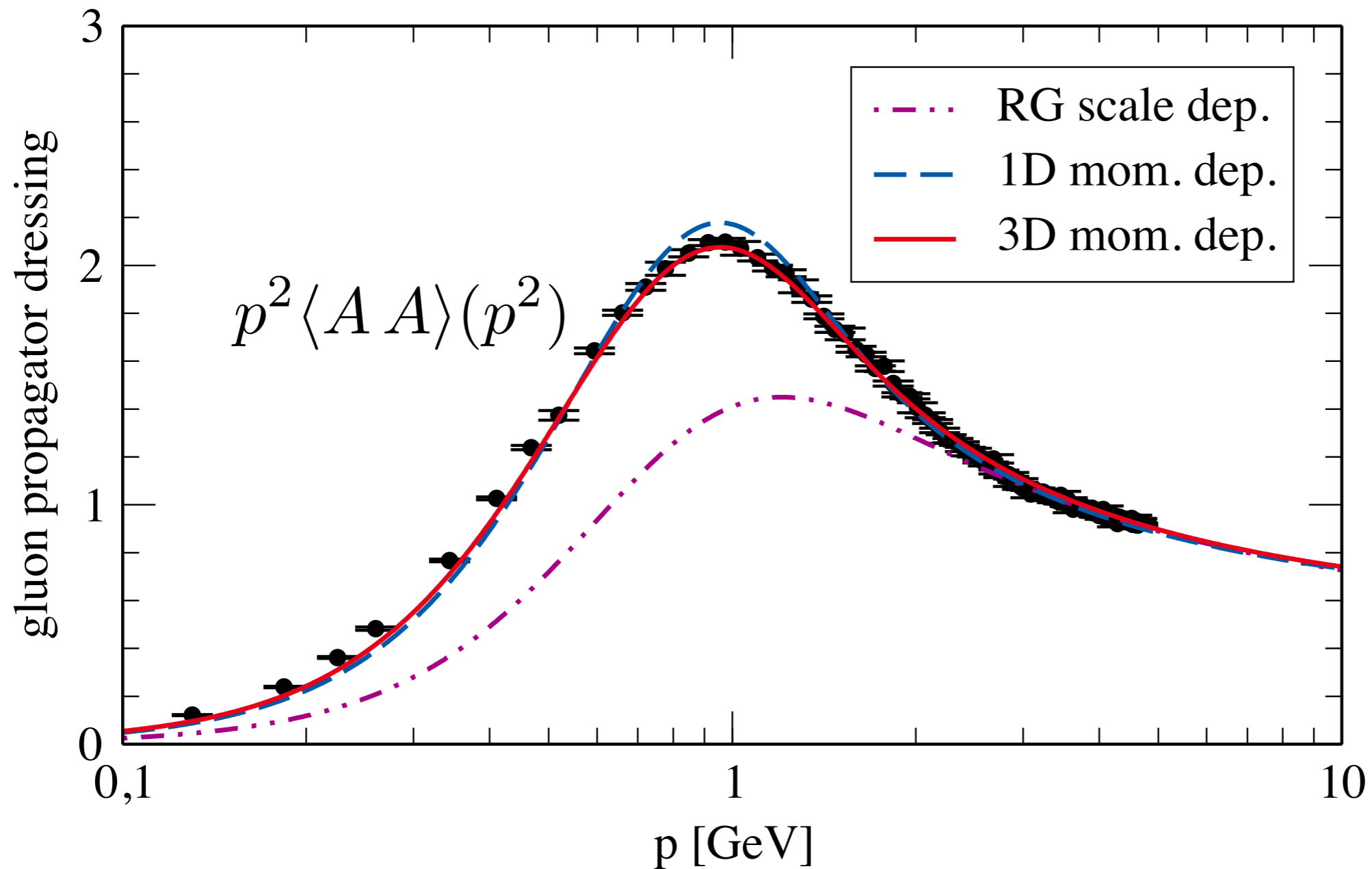


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**Aiming at apparent convergence**

# YM-theory: Euclidean gluon propagator

Functional Renormalisation Group

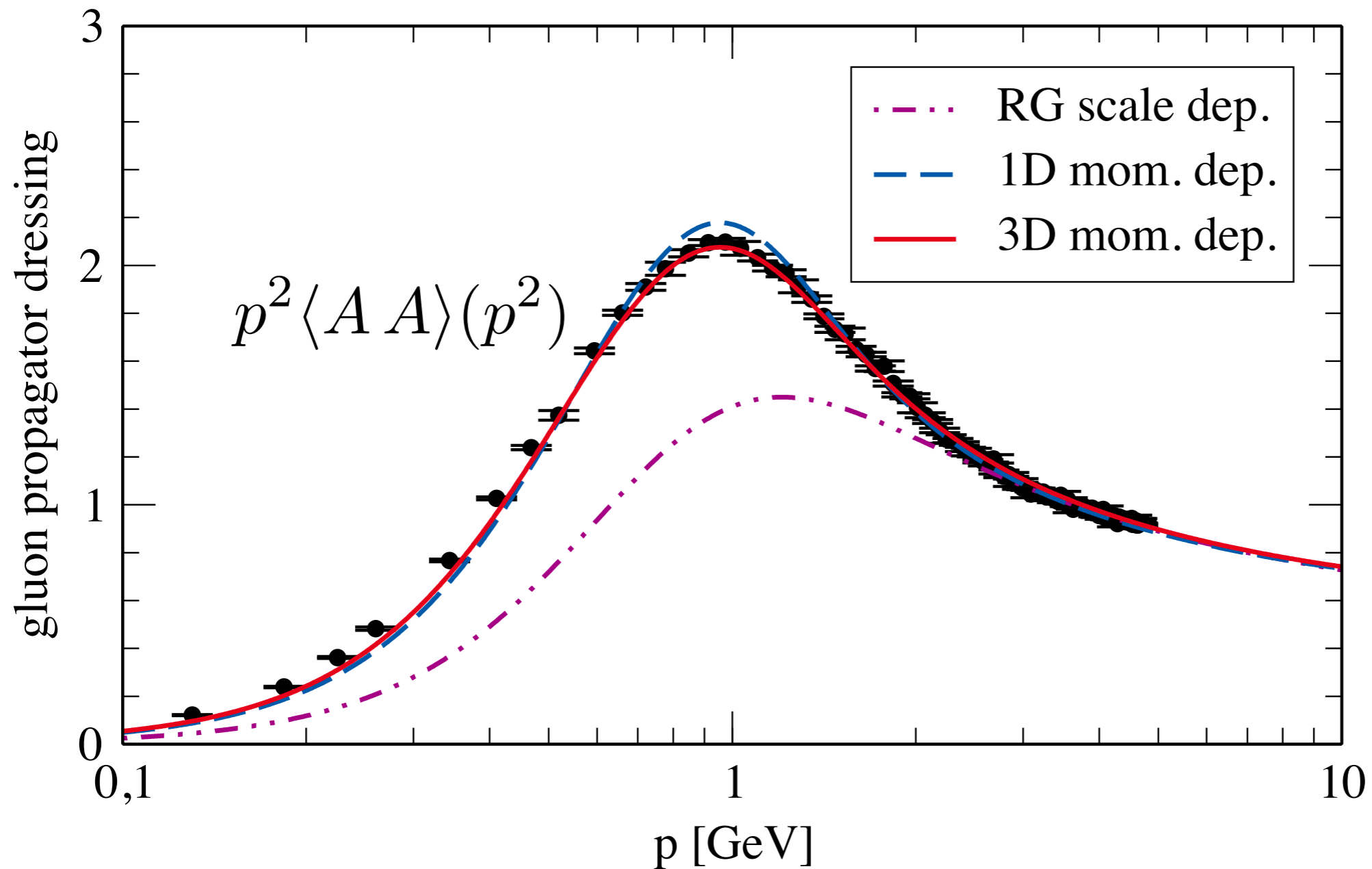


Lattice: Sternbeck, Ilgenfritz, Müller-Preussker, Schiller, Bogolubsky, PoS LAT2006, 076

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up to date pinch technique:  
Aguilar, Binosi, Papavassiliou, PRD 89 (2014) 085032

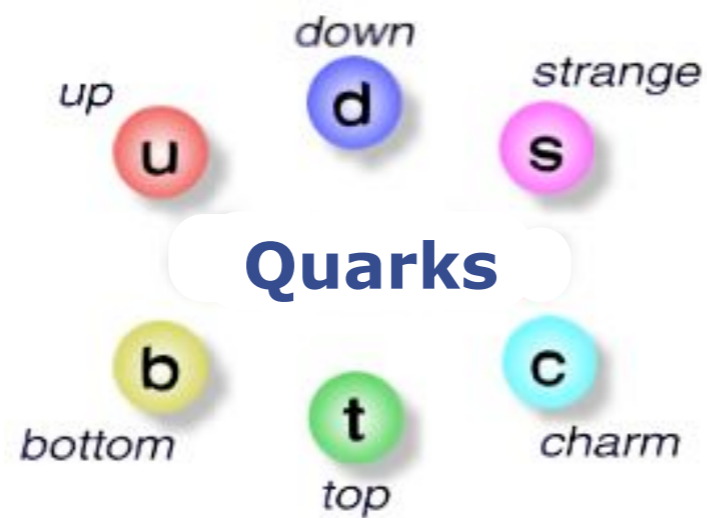
up to date DSE:  
Cyrol, Huber, Smekal, EPJ C75 (2015) 102

Cyrol, Fister, Mitter, JMP, Strodthoff, PRD 94 (2016) 054005

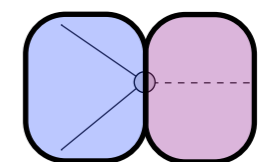


# Dynamical hadronisation

Gies, Wetterich '01  
JMP '05  
Flörchinger, Wetterich '09



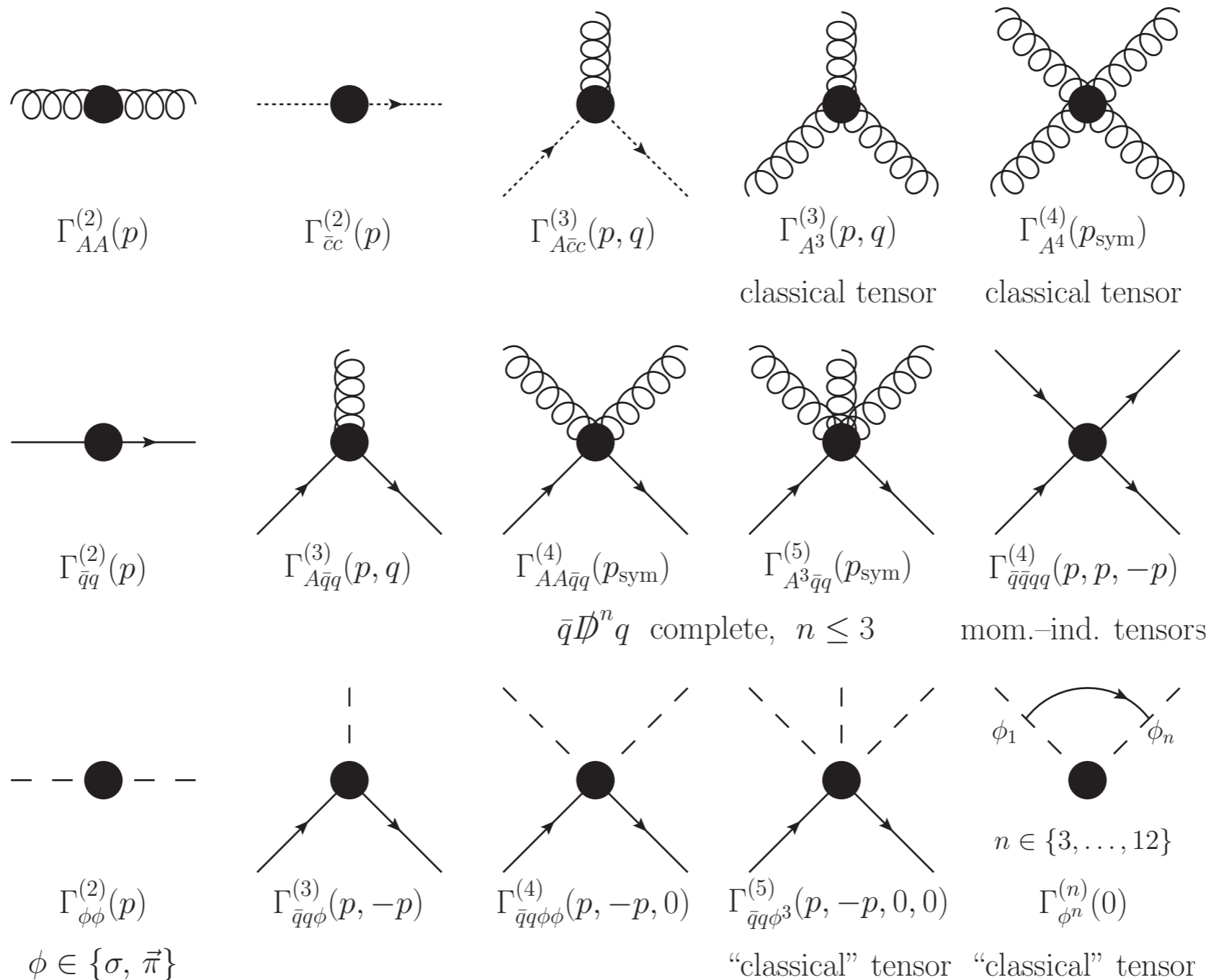
**Gluons**



dynamical

**Hadrons**

# QCD: current set of correlation functions



**Aiming at apparent convergence**

Cyrol, Mitter, JMP, Strodthoff, PRD 97 (2018) 054006,  
 PRD 97 (2018) 054015

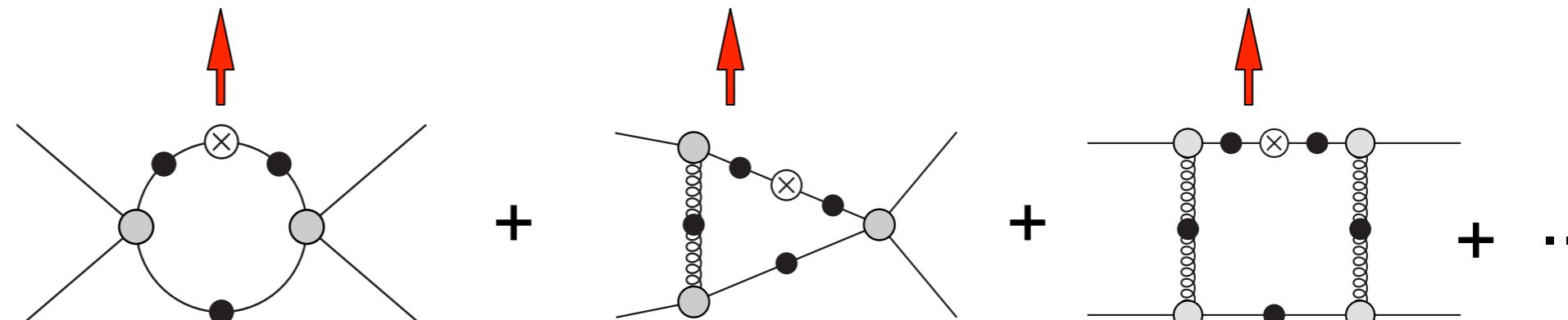
Cyrol, Fister, Mitter, JMP, Strodthoff, PRD 94 (2016) 054005

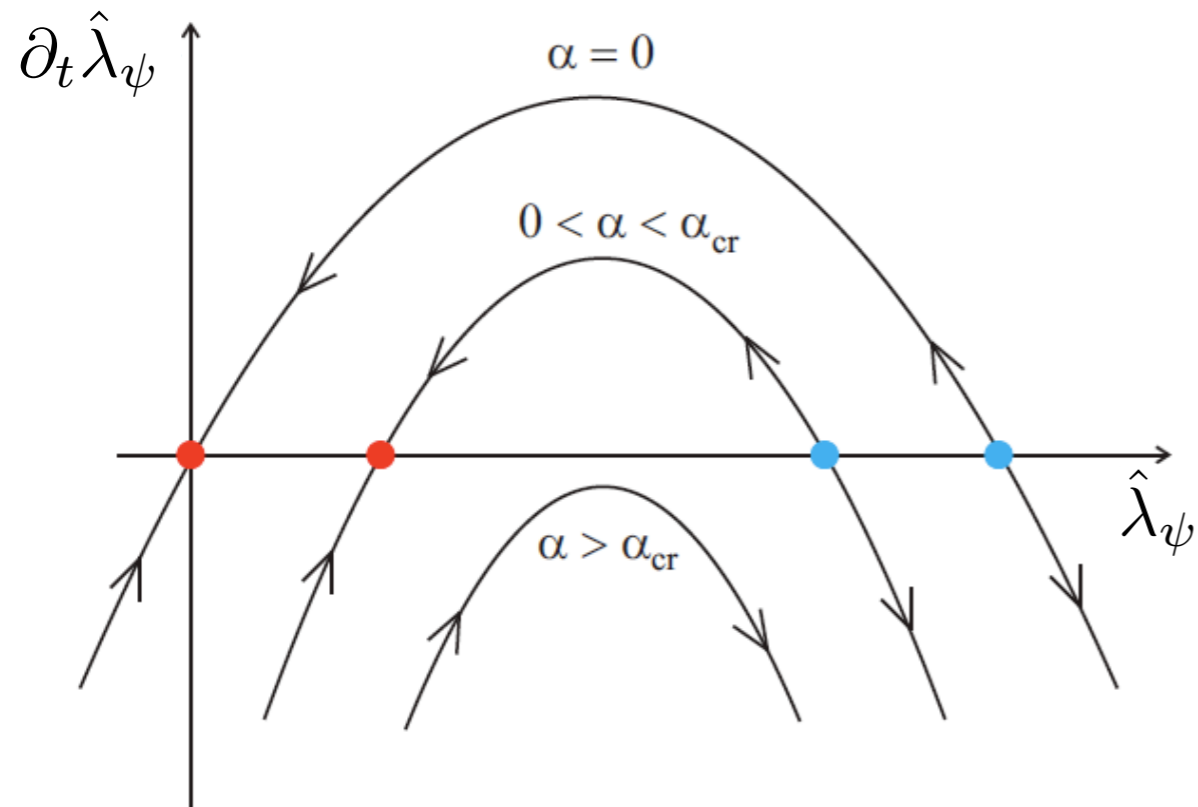
Mitter, JMP, Strodthoff, PRD 91 (2015) 054035

# Chiral symmetry breaking

## A glimpse at chiral symmetry breaking in QCD within the FRG

Flow for four-fermion coupling  $\hat{\lambda}_\psi = \lambda_\psi k^2$  with infrared scale  $k$

$$k \partial_k \hat{\lambda}_\psi = 2 \hat{\lambda}_\psi + A \left( \frac{T}{k} \right) \hat{\lambda}_\psi^2 + B \left( \frac{T}{k} \right) \hat{\lambda}_\psi \alpha_s + C \left( \frac{T}{k} \right) \alpha_s^2 + \dots$$


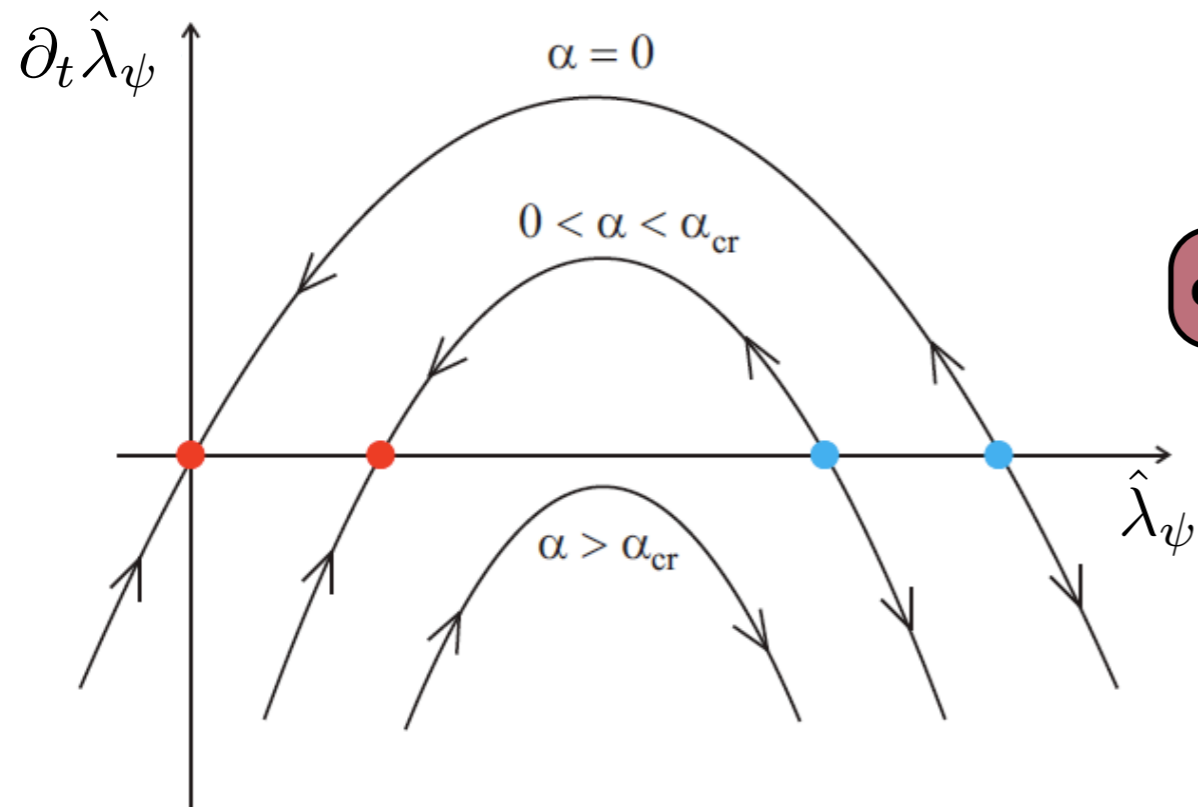


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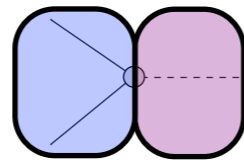
Gies, Wetterich '01  
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Hubbard-Stratonovich



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$$\frac{\partial}{\partial t} \Big|_\phi \Gamma_k[\phi] = \frac{1}{2} G_{k,\phi} \dot{R}_{k,\phi} + R_k G_{k,\phi} \frac{\delta \dot{\phi}}{\delta \phi} - \frac{\delta \Gamma}{\delta \phi} \dot{\phi}$$

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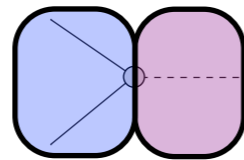
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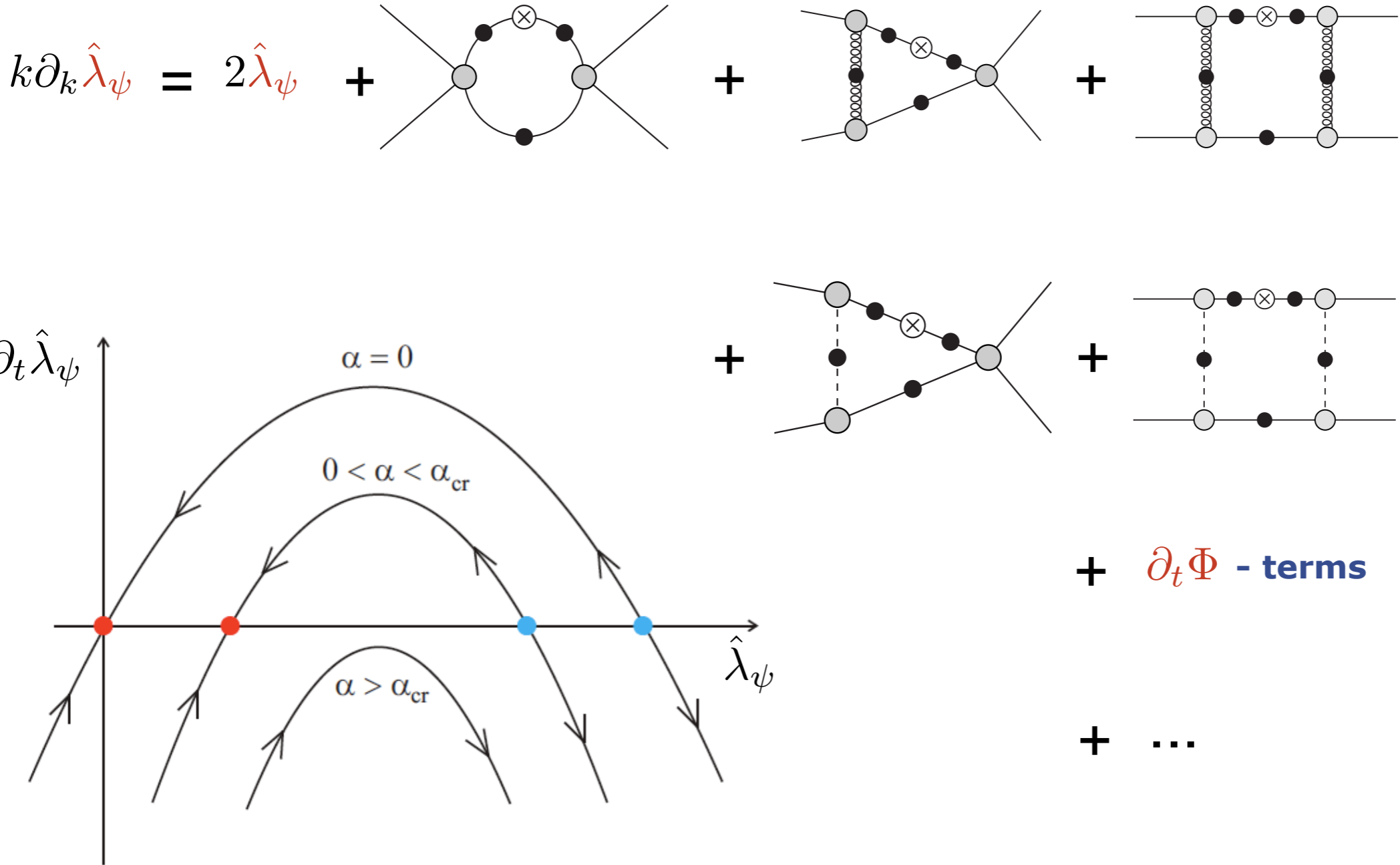
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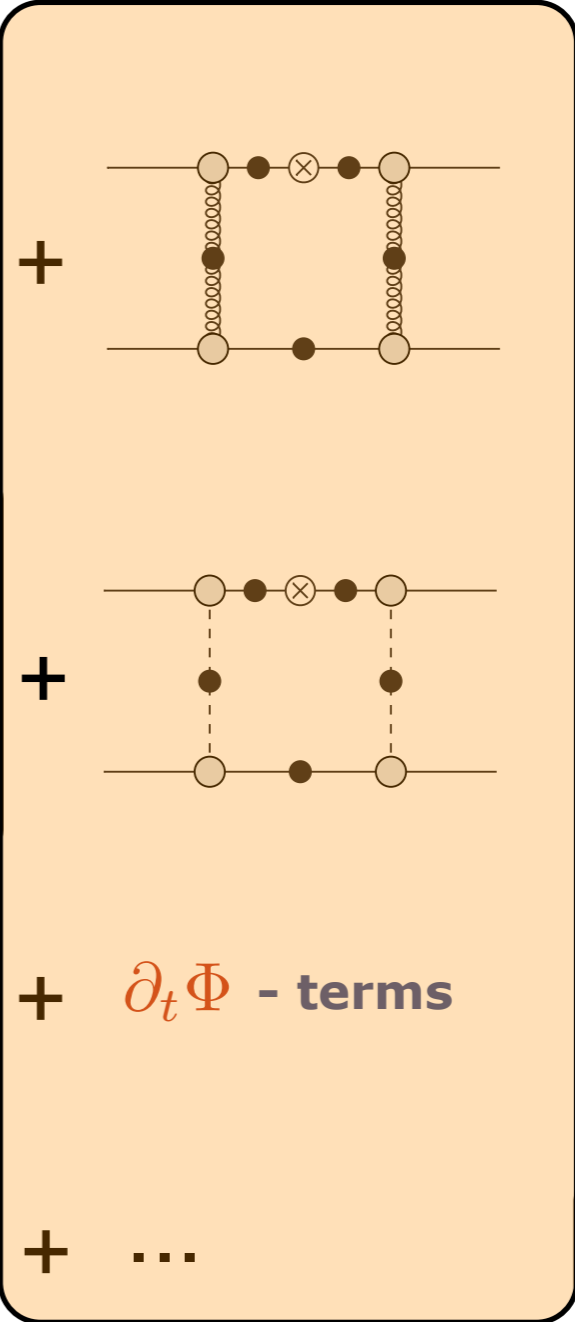
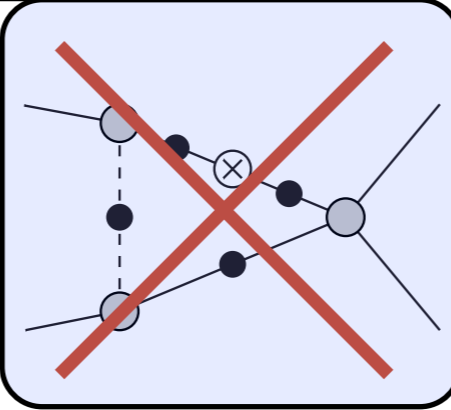
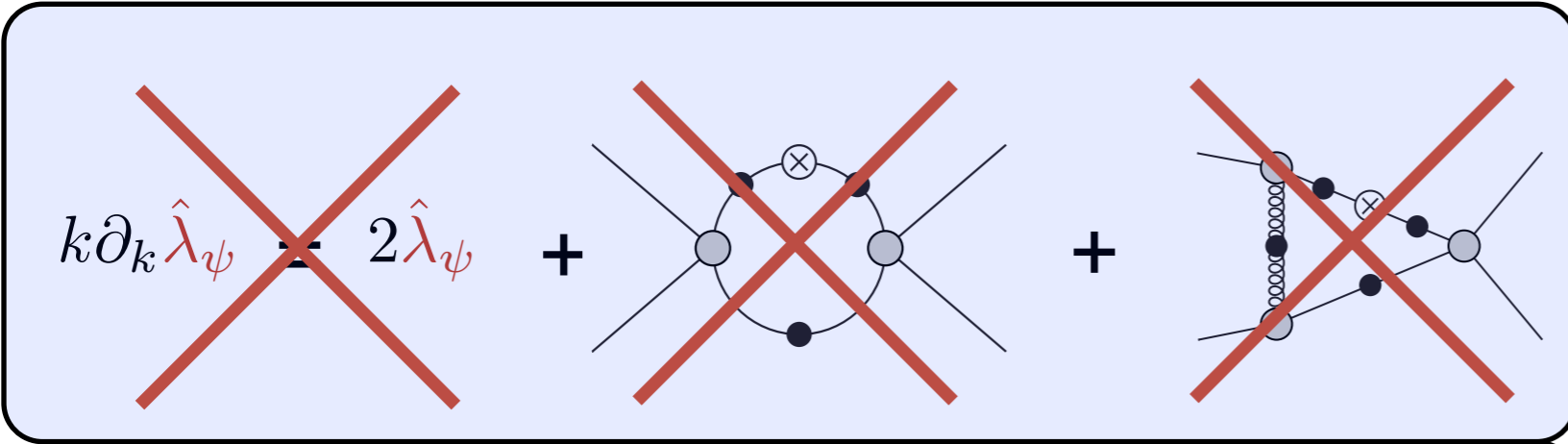
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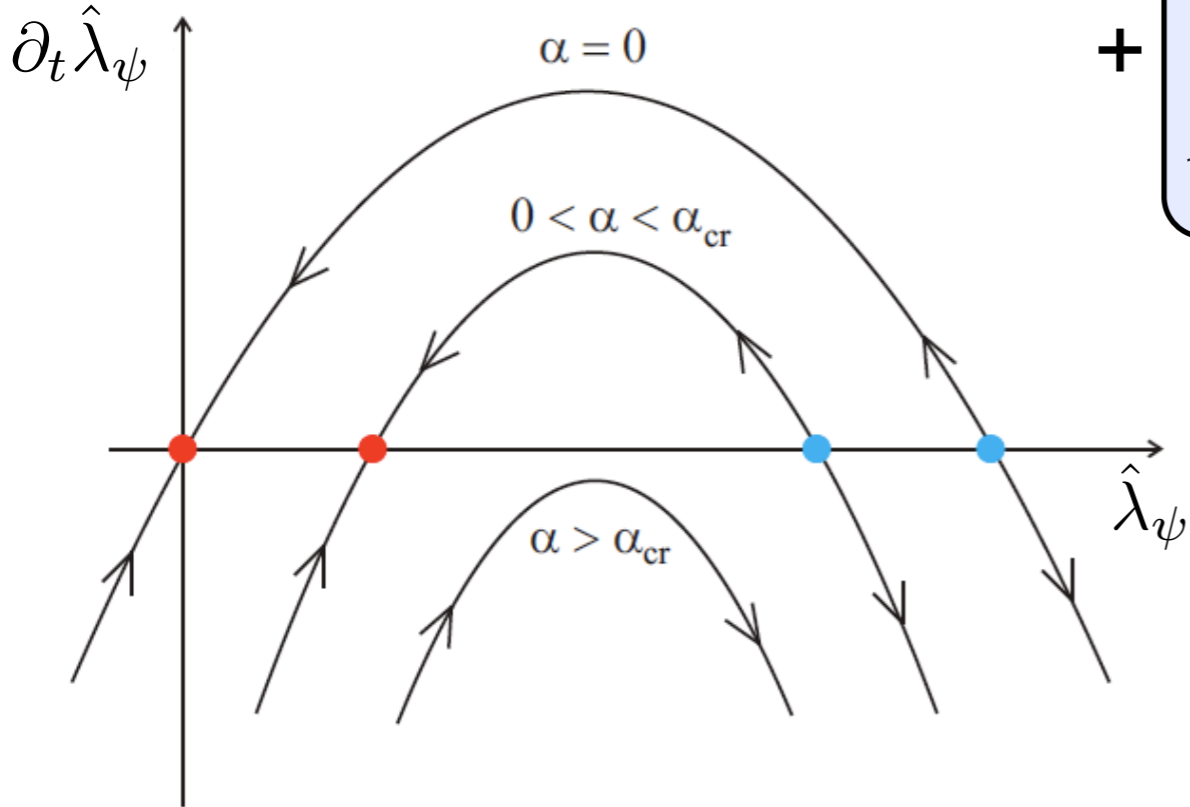


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Full bosonisation  $\hat{\lambda}_\psi = 0$



$= 0$

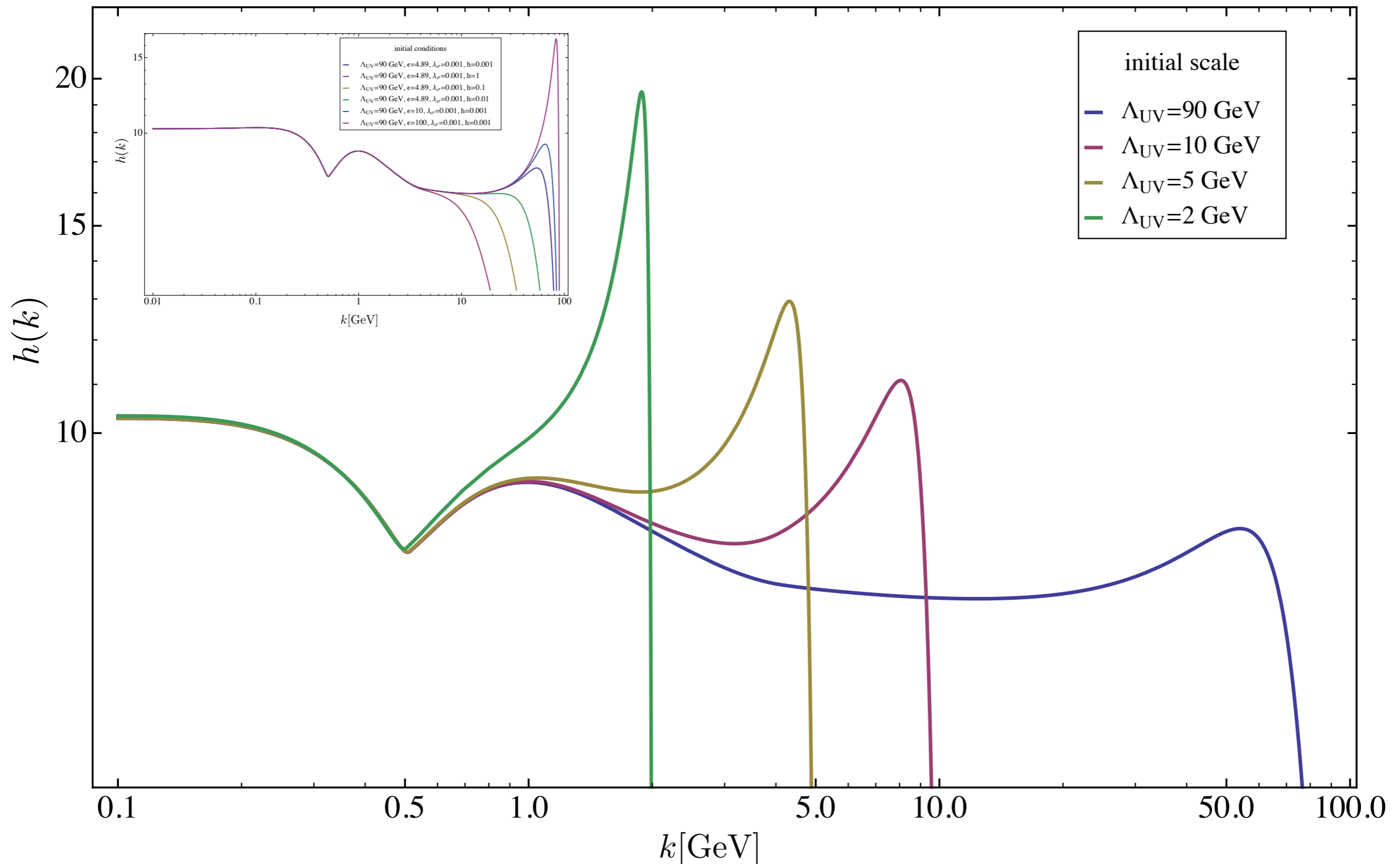




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Braun, Fister, Haas, JMP, Rennecke '14  
Mitter, JMP, Strodthoff '14  
Cyrol, Mitter, JMP, Strodthoff '17

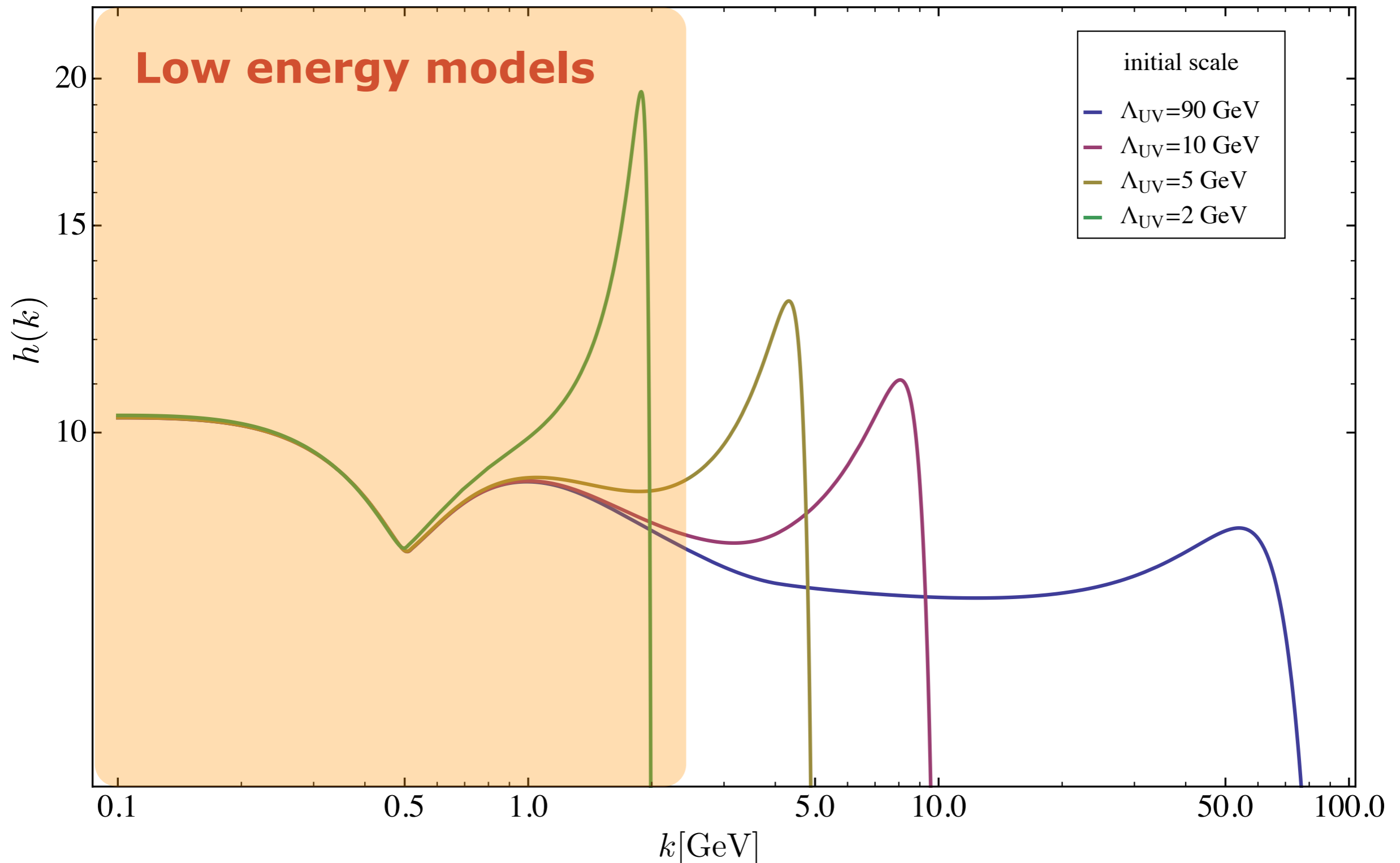
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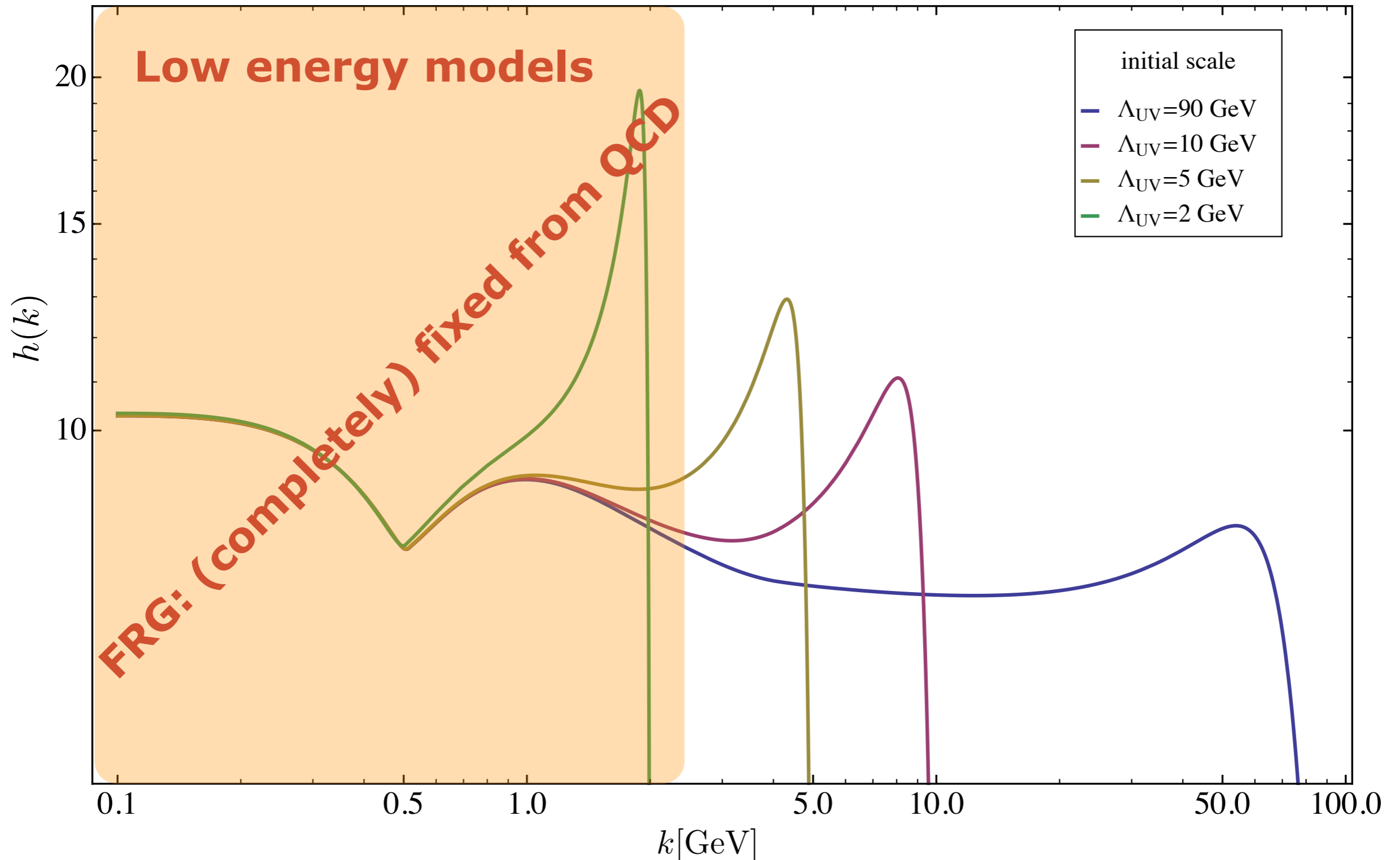


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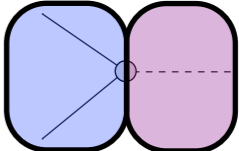
$$\hat{\lambda}_\psi = 0$$



# QCD results at T=0



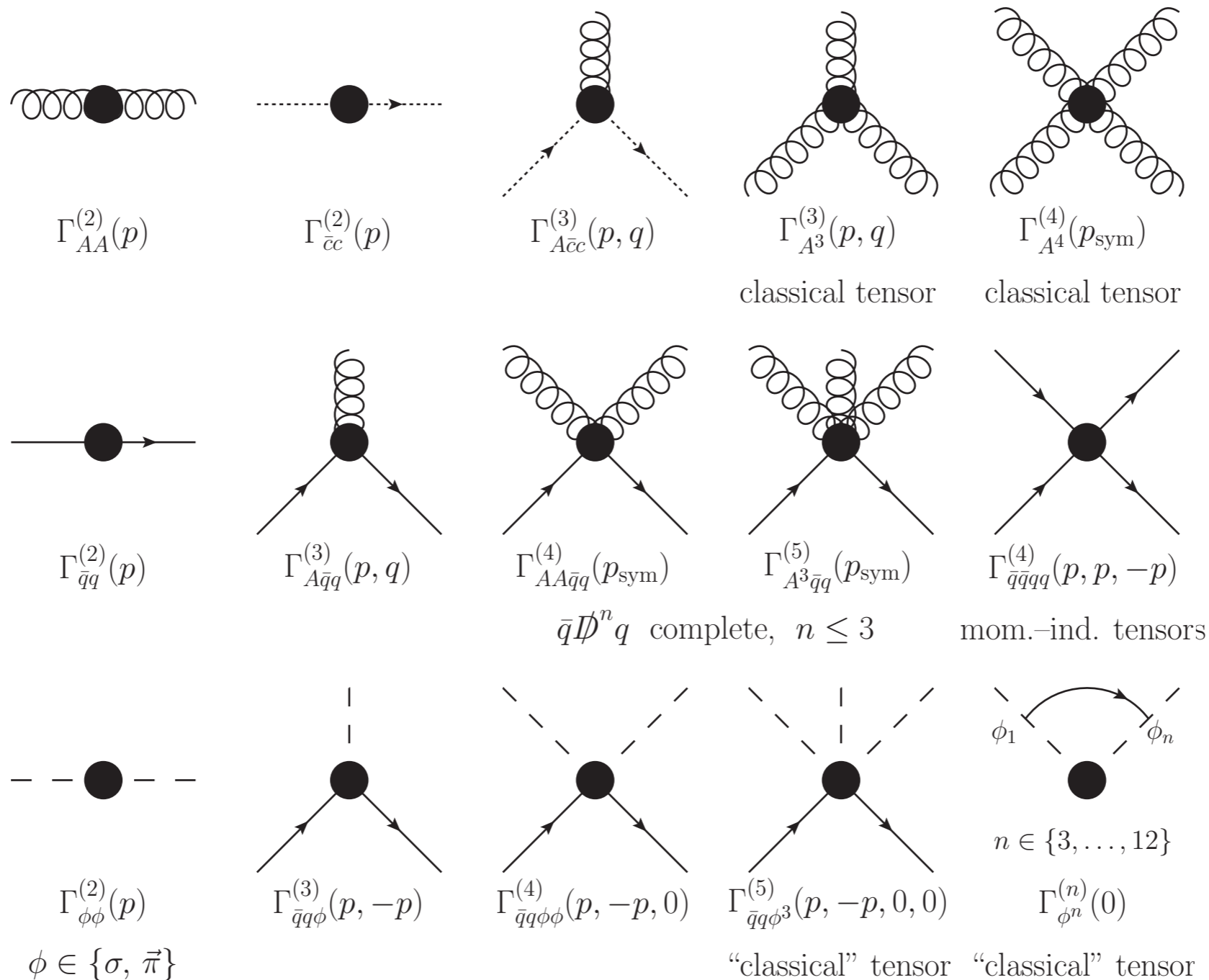
**Gluons**



**dynamical**

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**Aiming at apparent convergence**

Cyrol, Mitter, JMP, Strodthoff, PRD 97 (2018) 054006,  
 PRD 97 (2018) 054015

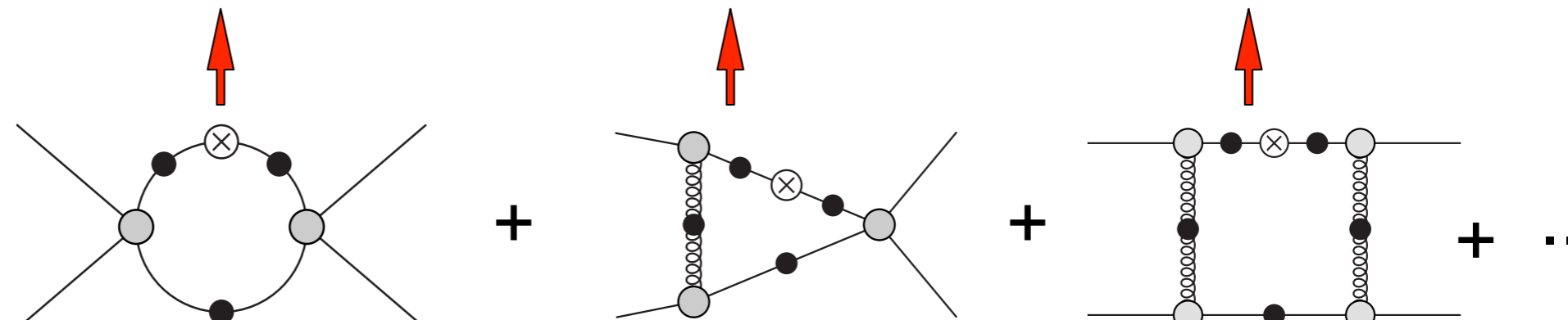
Cyrol, Fister, Mitter, JMP, Strodthoff, PRD 94 (2016) 054005

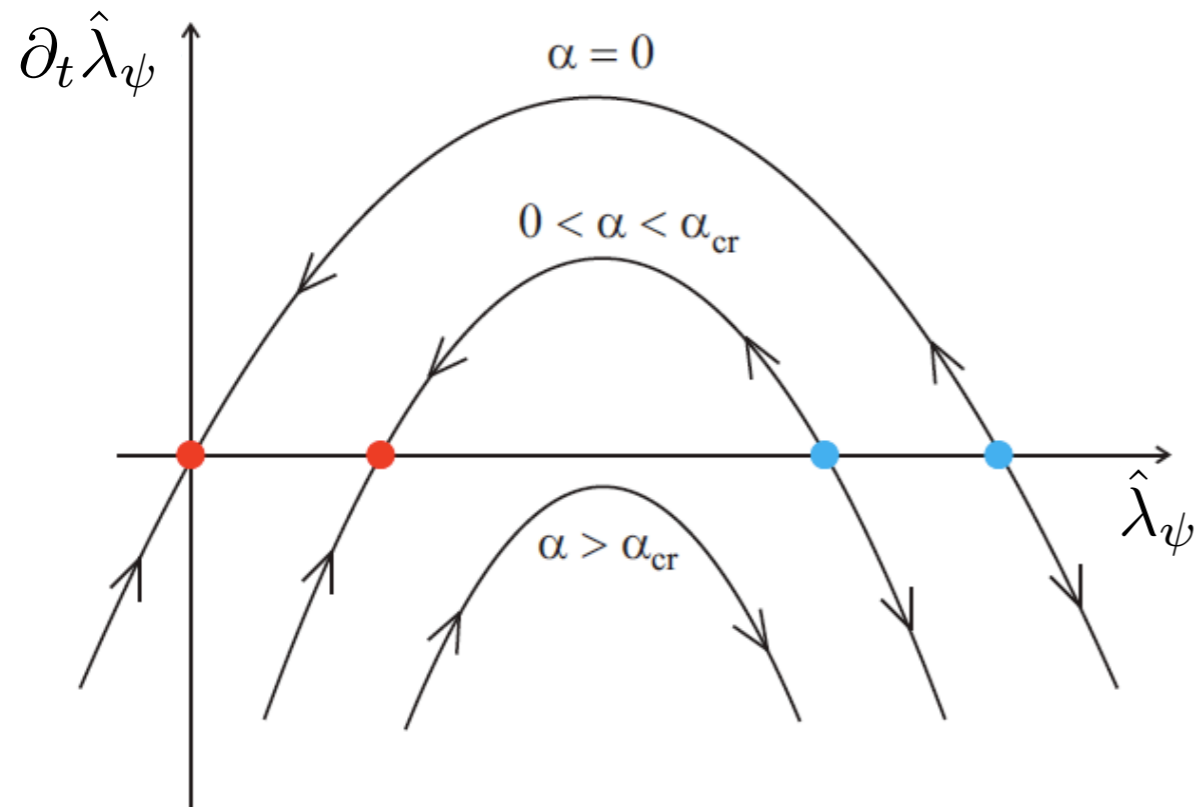
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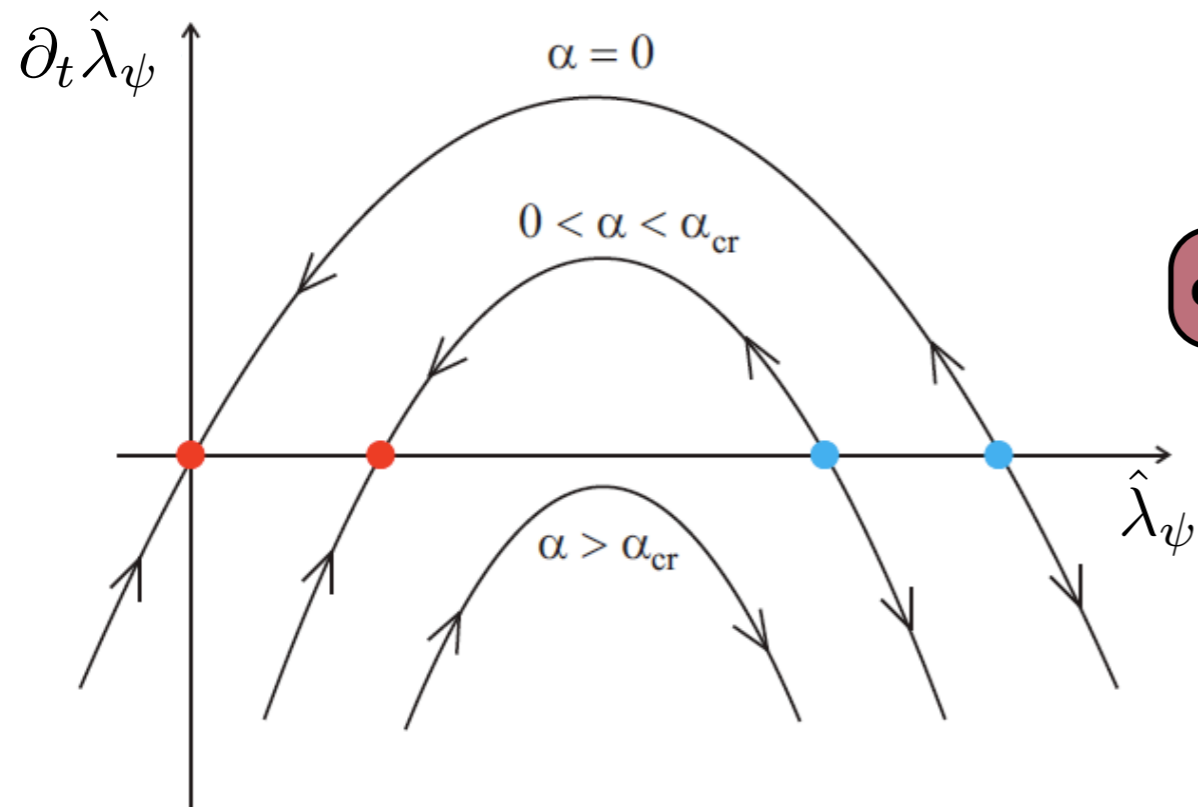


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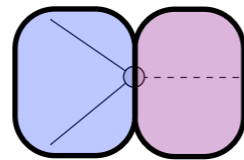
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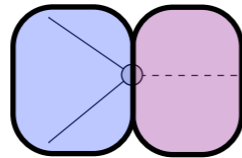
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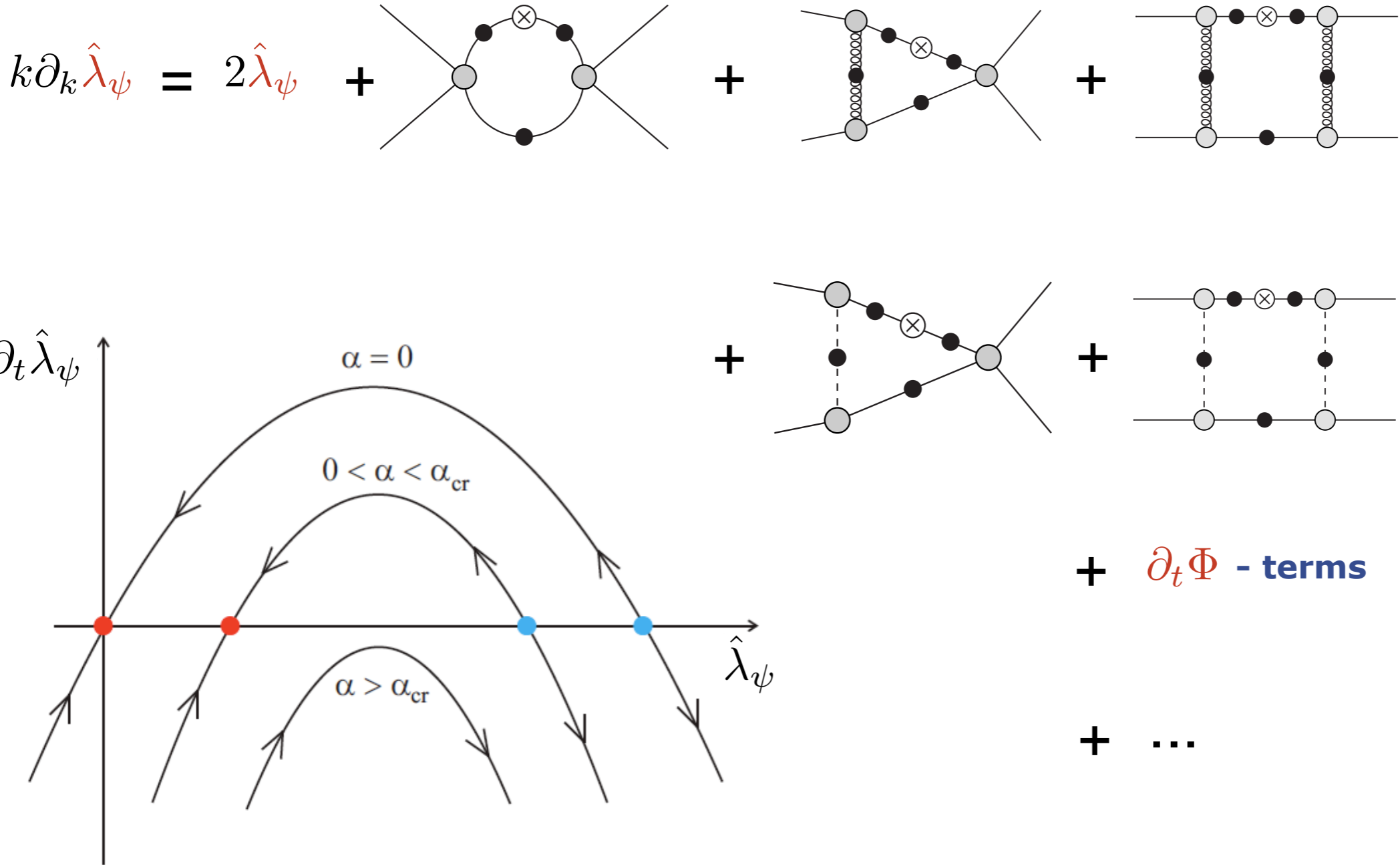
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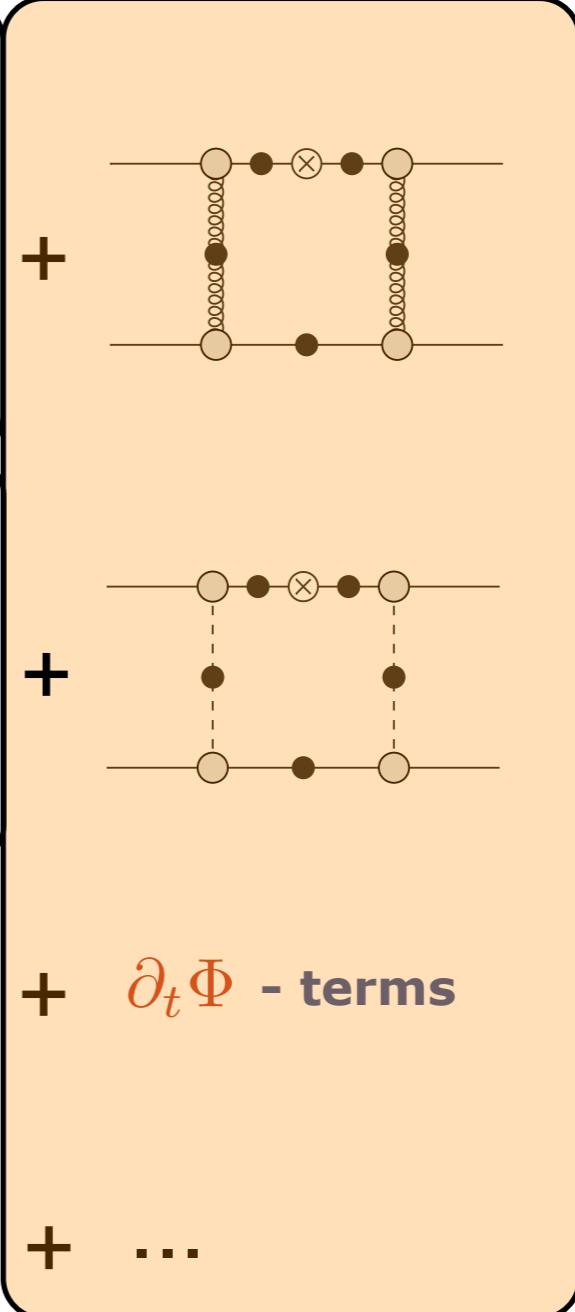
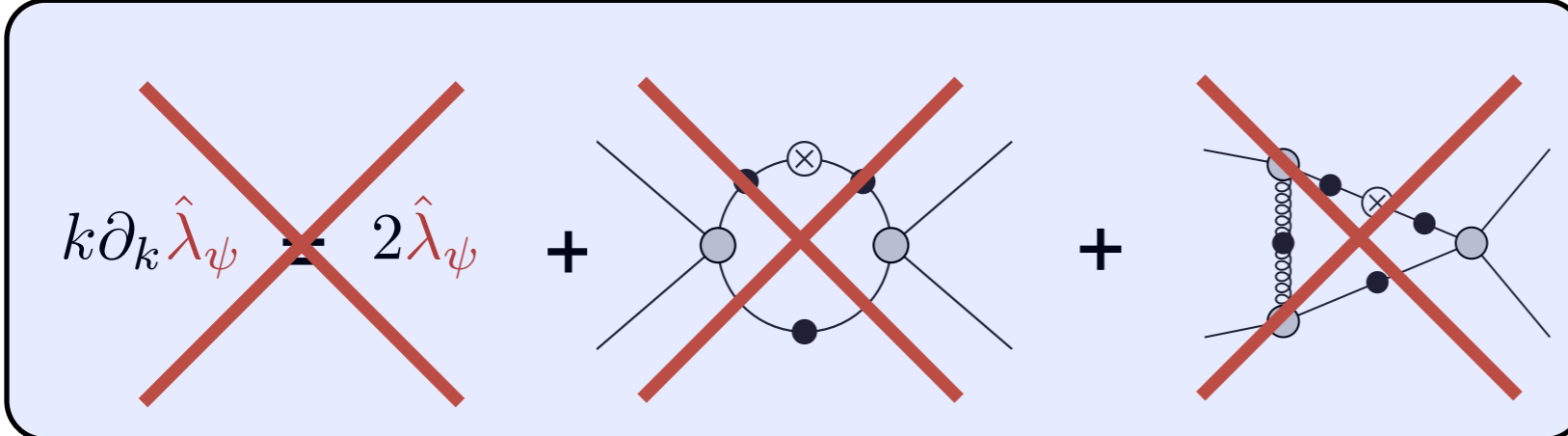
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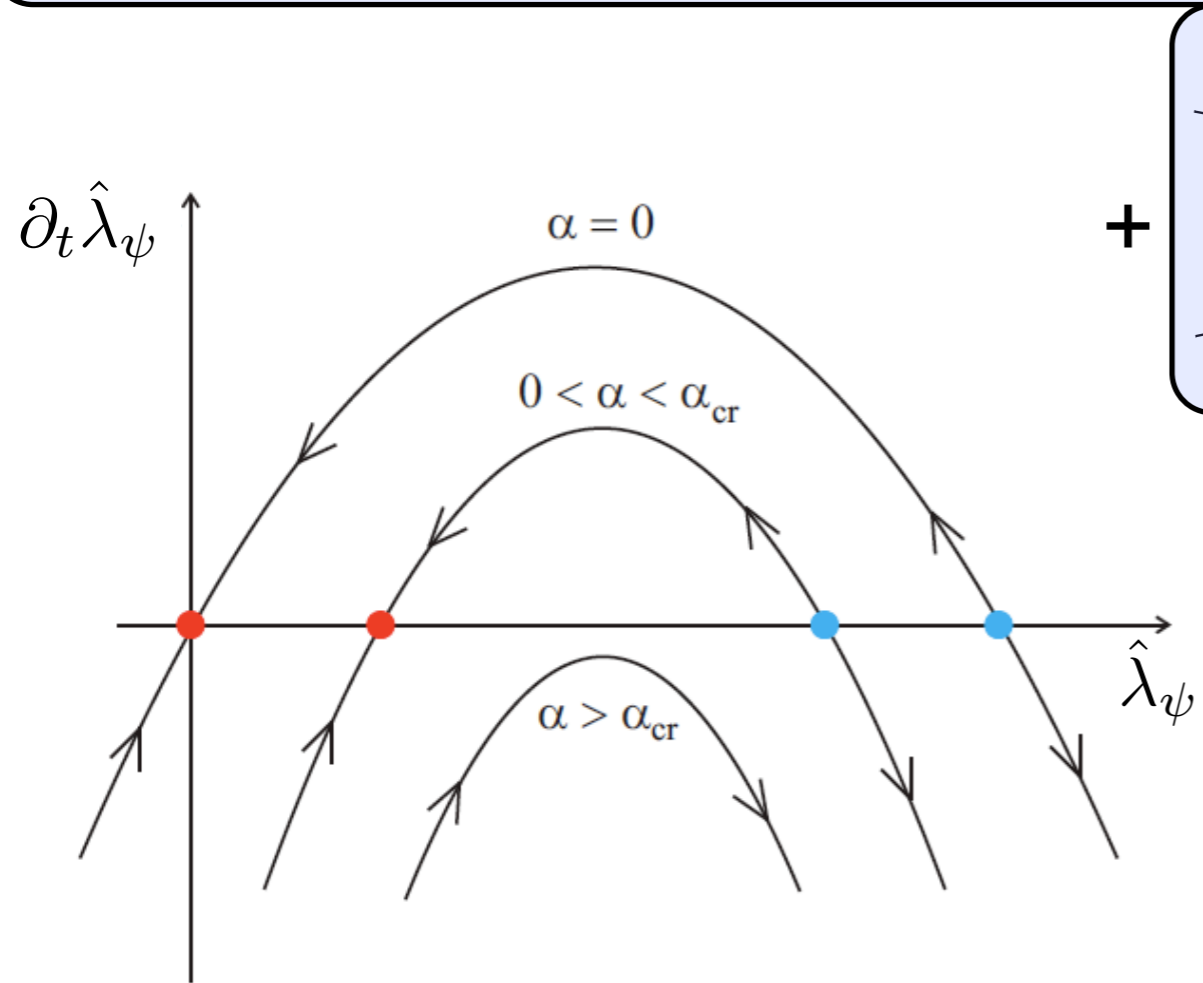


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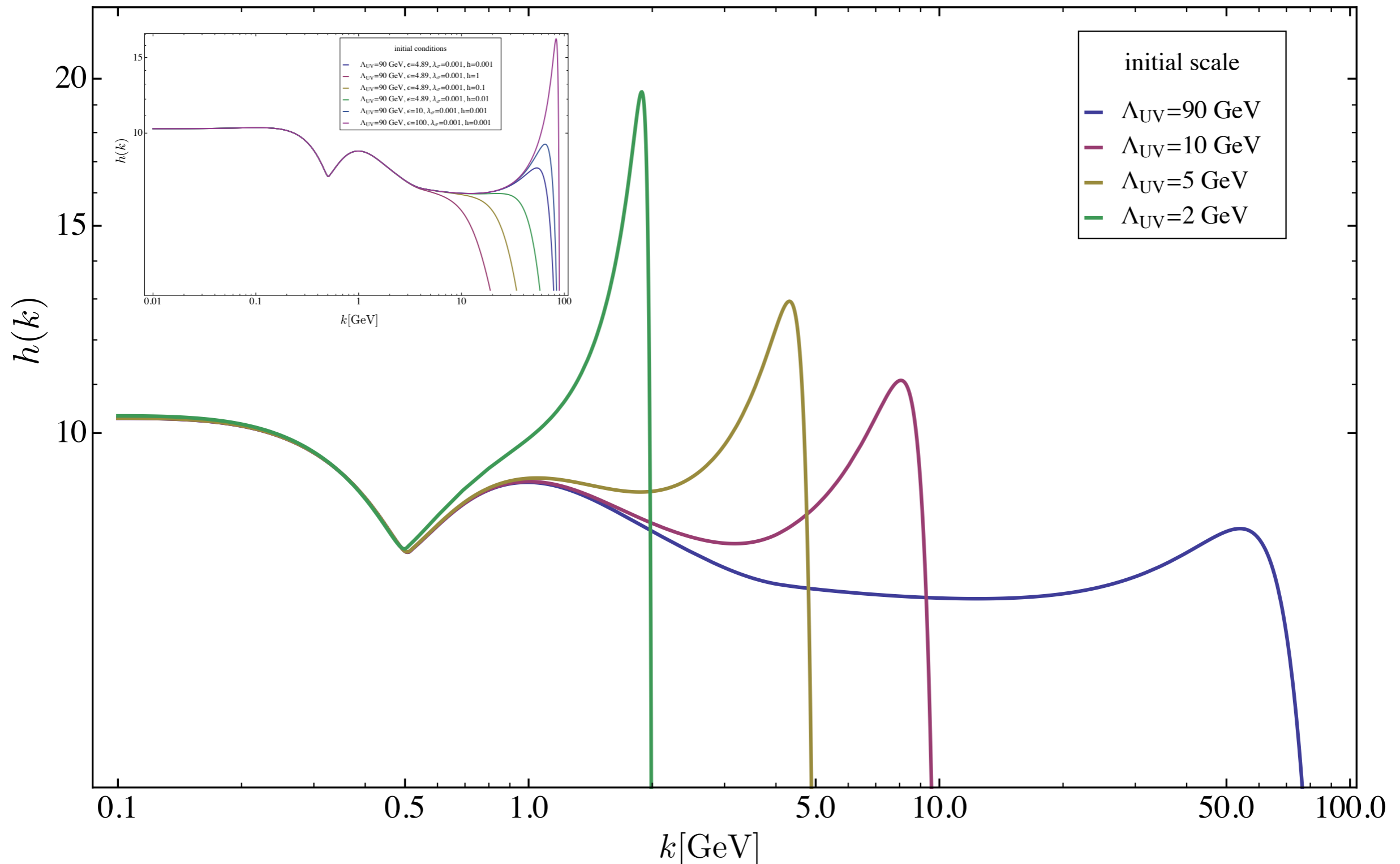
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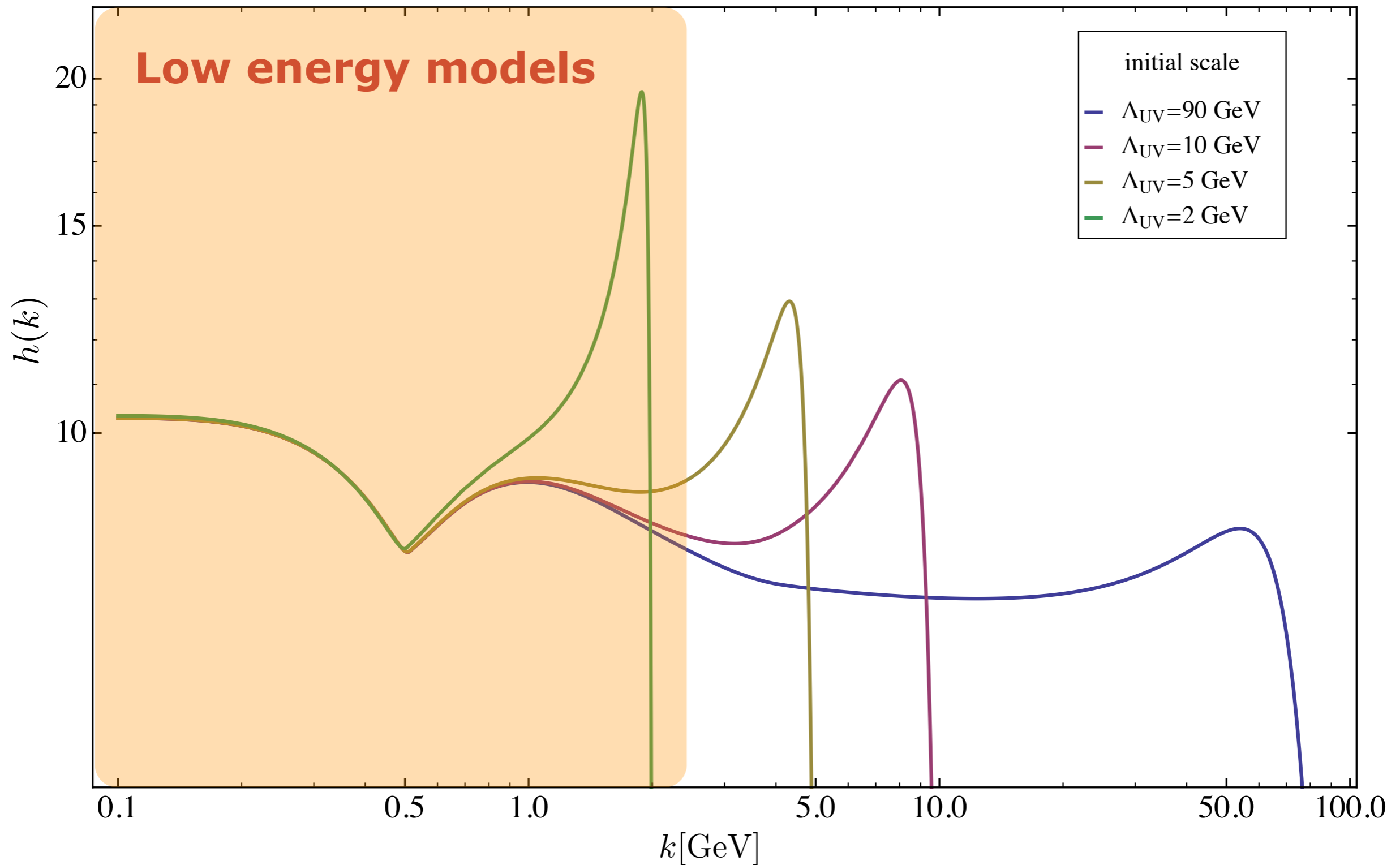
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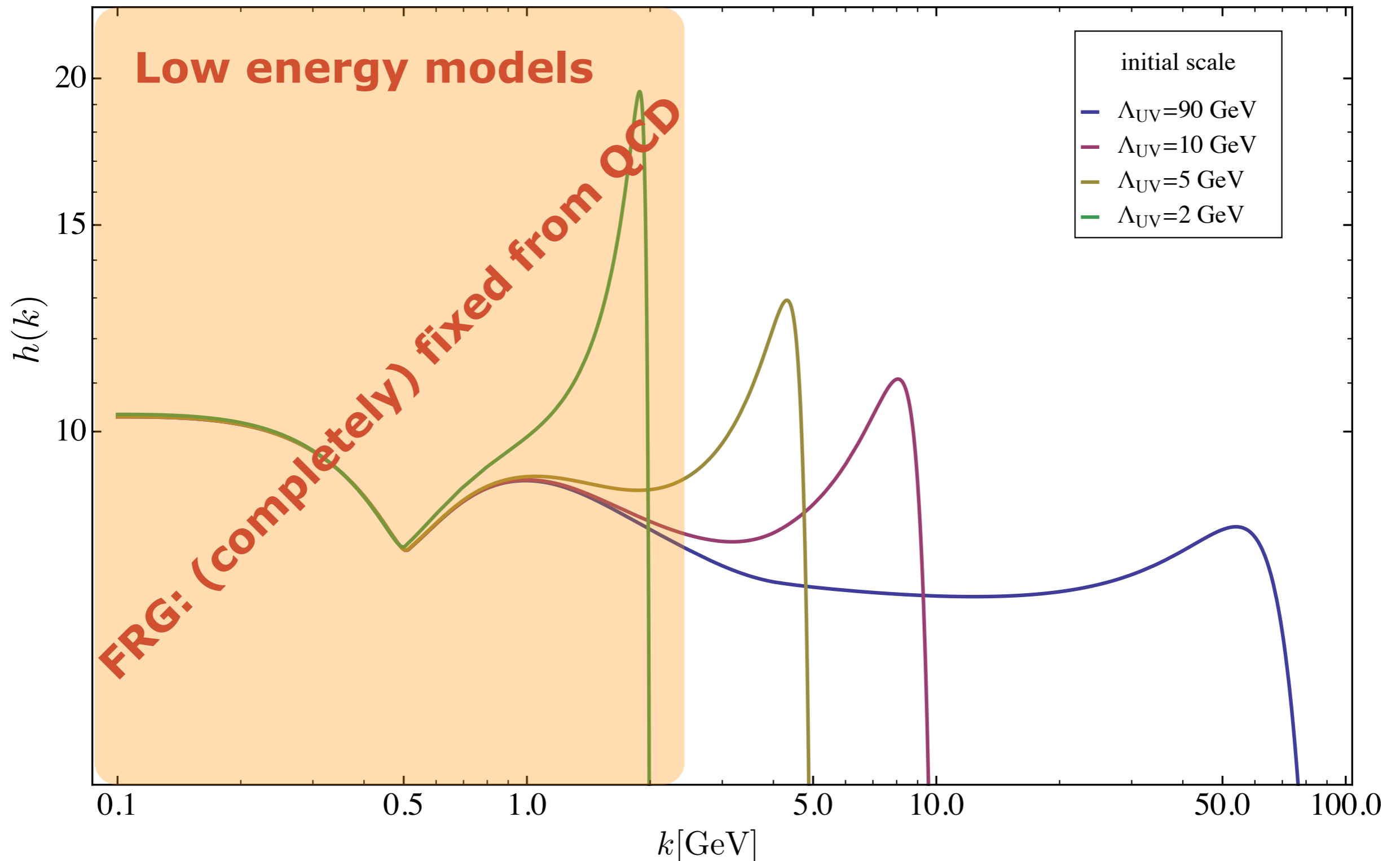


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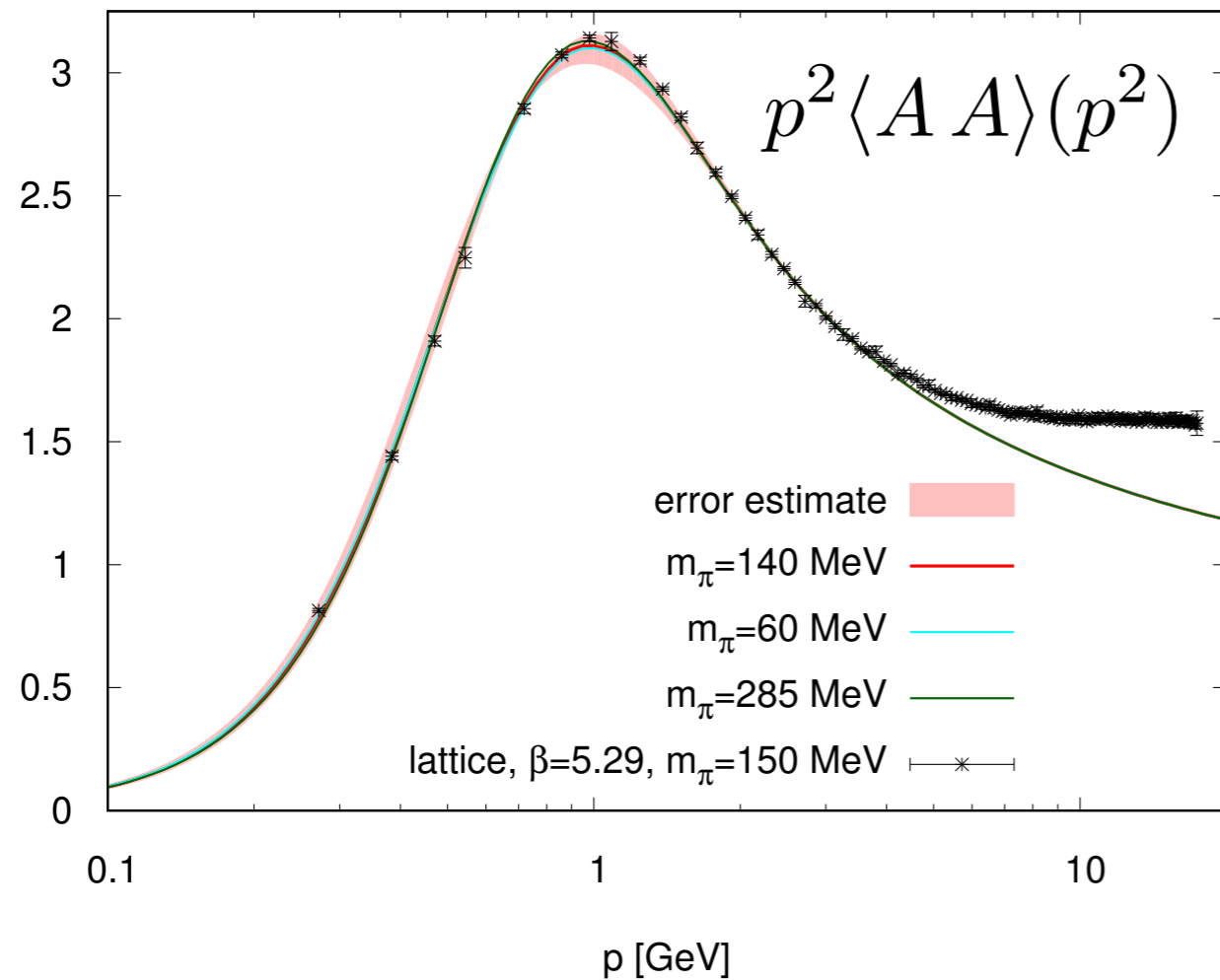
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Mitter, JMP, Strodthoff '14  
Cyrol, Mitter, JMP, Strodthoff '17

Full bosonisation

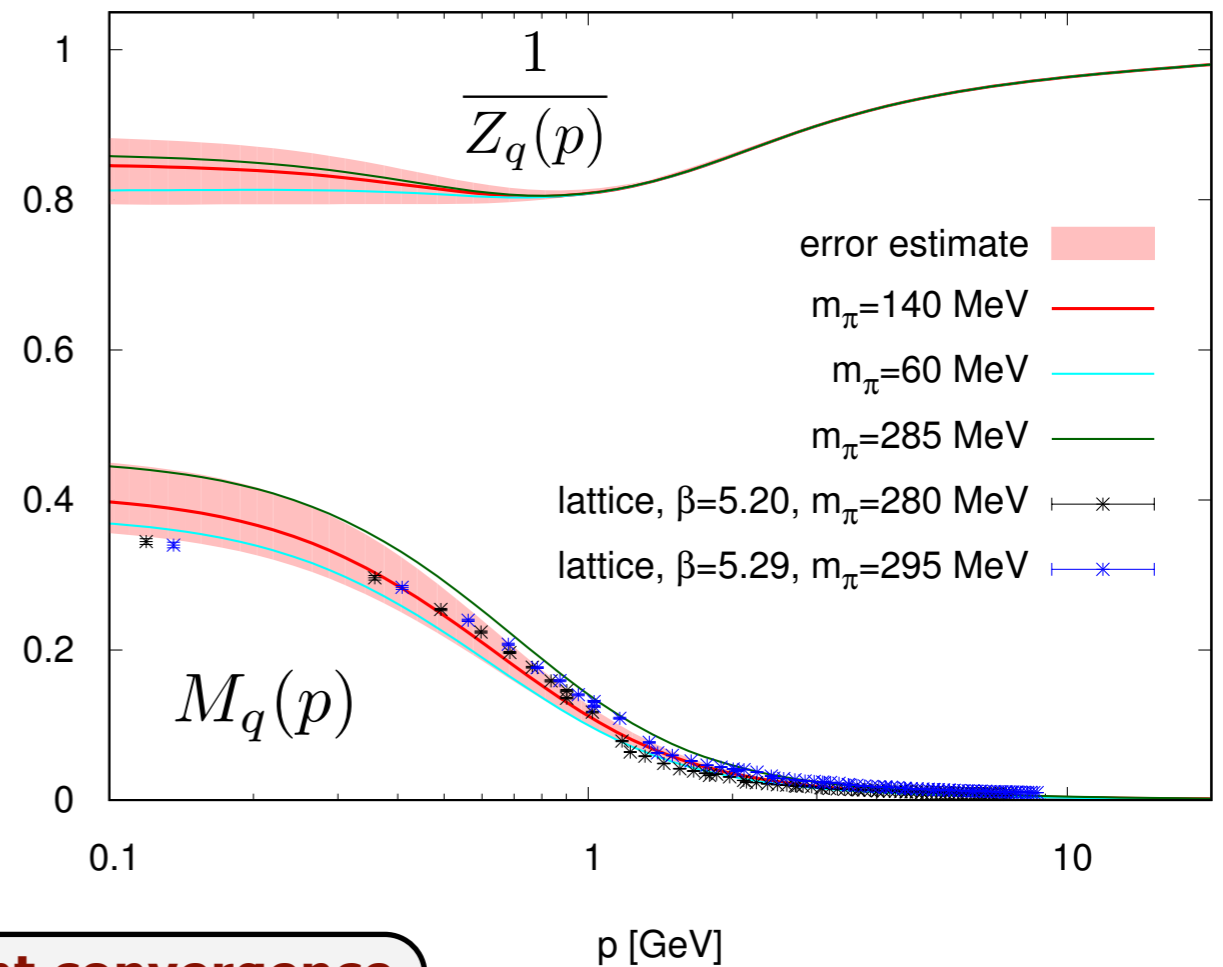
$$\hat{\lambda}_\psi = 0$$



# QCD: Euclidean propagators



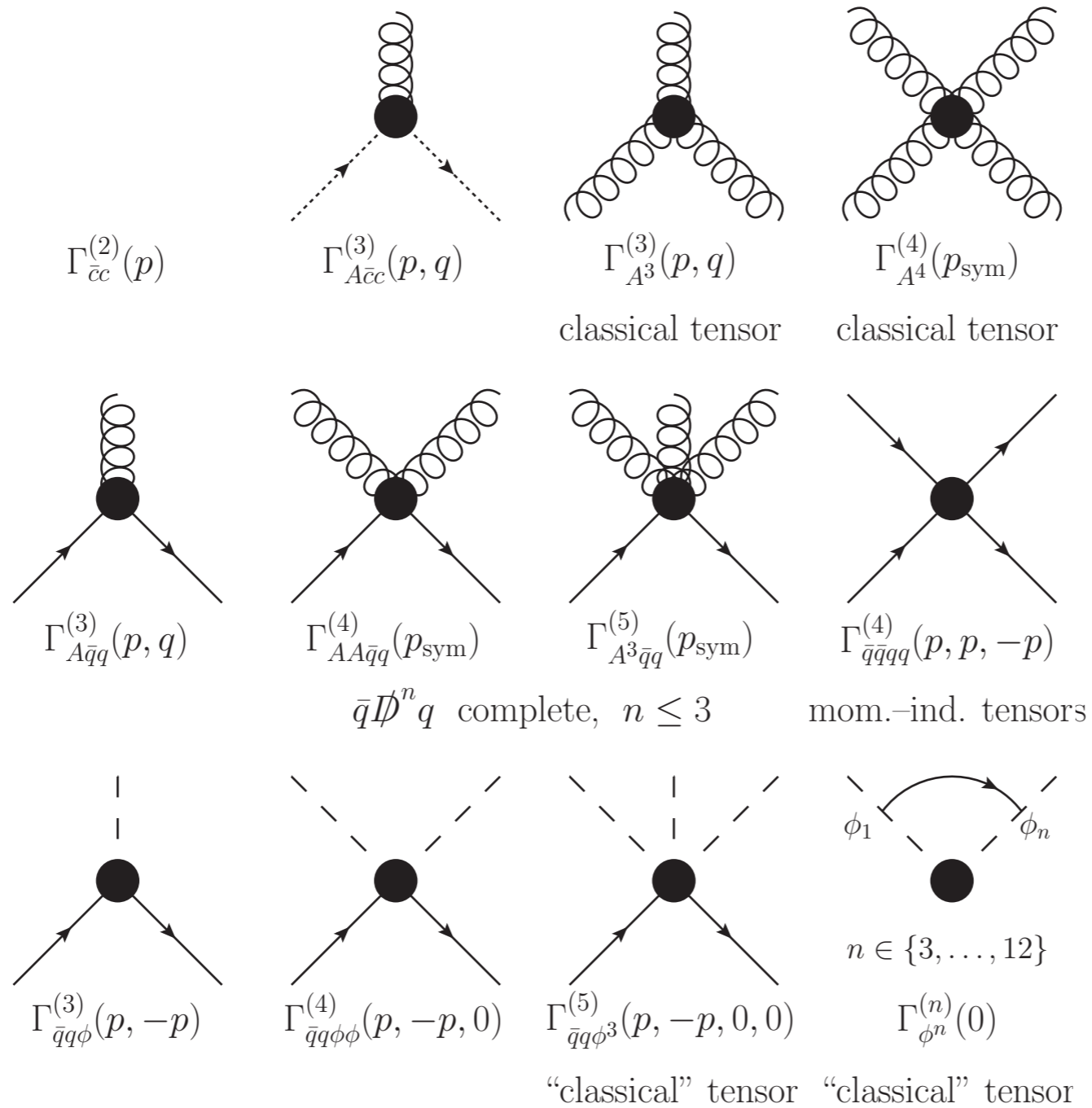
$$\frac{1}{Z_q(p)} \frac{1}{i \not{p} + M_q(p)}$$



**Aiming at apparent convergence**

lattice, e.g.: Oliviera et al, Acta Phys.Polon.Supp. 9 (2016) 363  
 Sternbeck et al, PoS LATTICE2016 (2017)  
 A. Athenodorou et al, PLB 761 (2016) 444

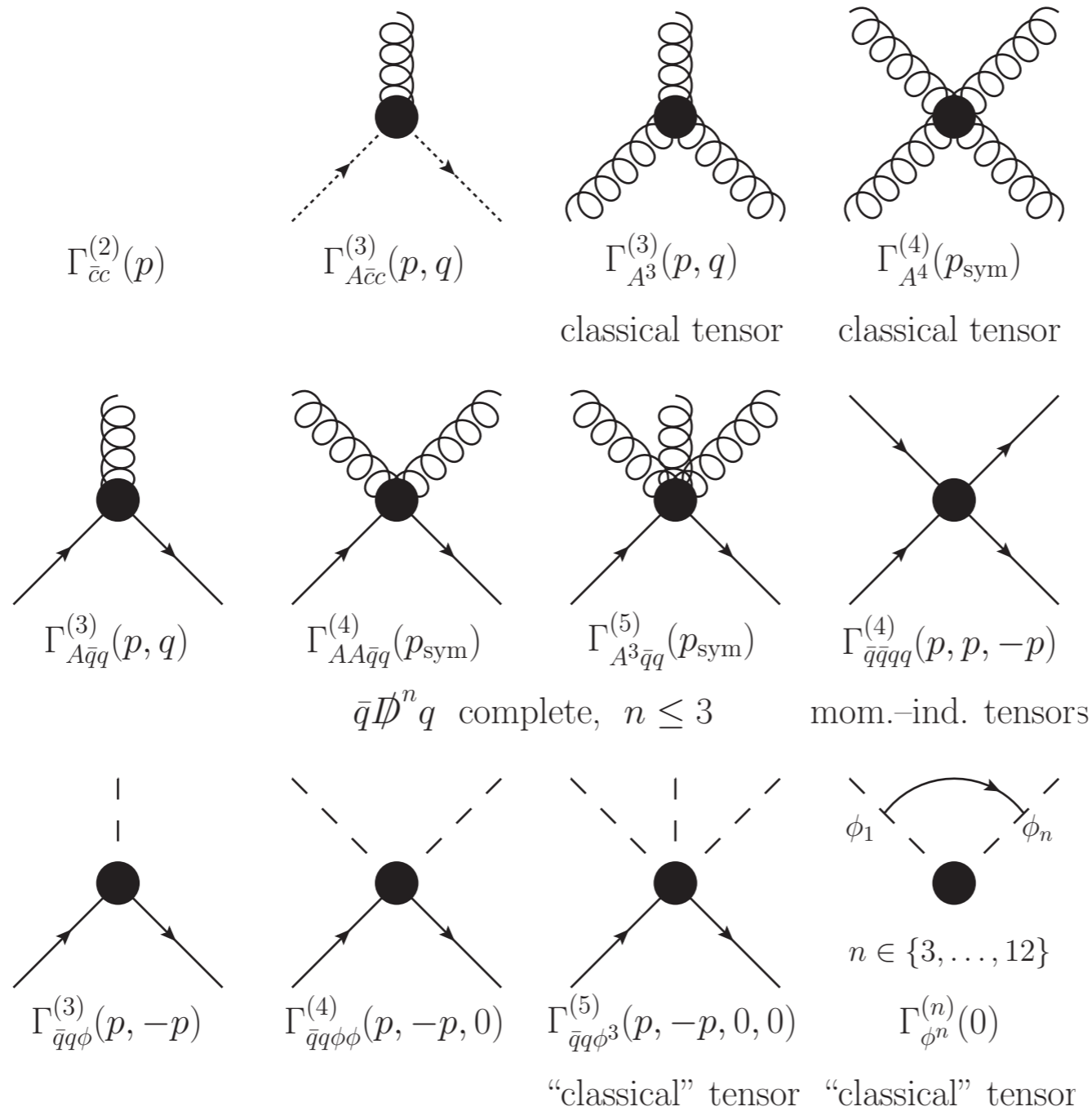
# QCD: Vertices



**Aiming at apparent convergence**



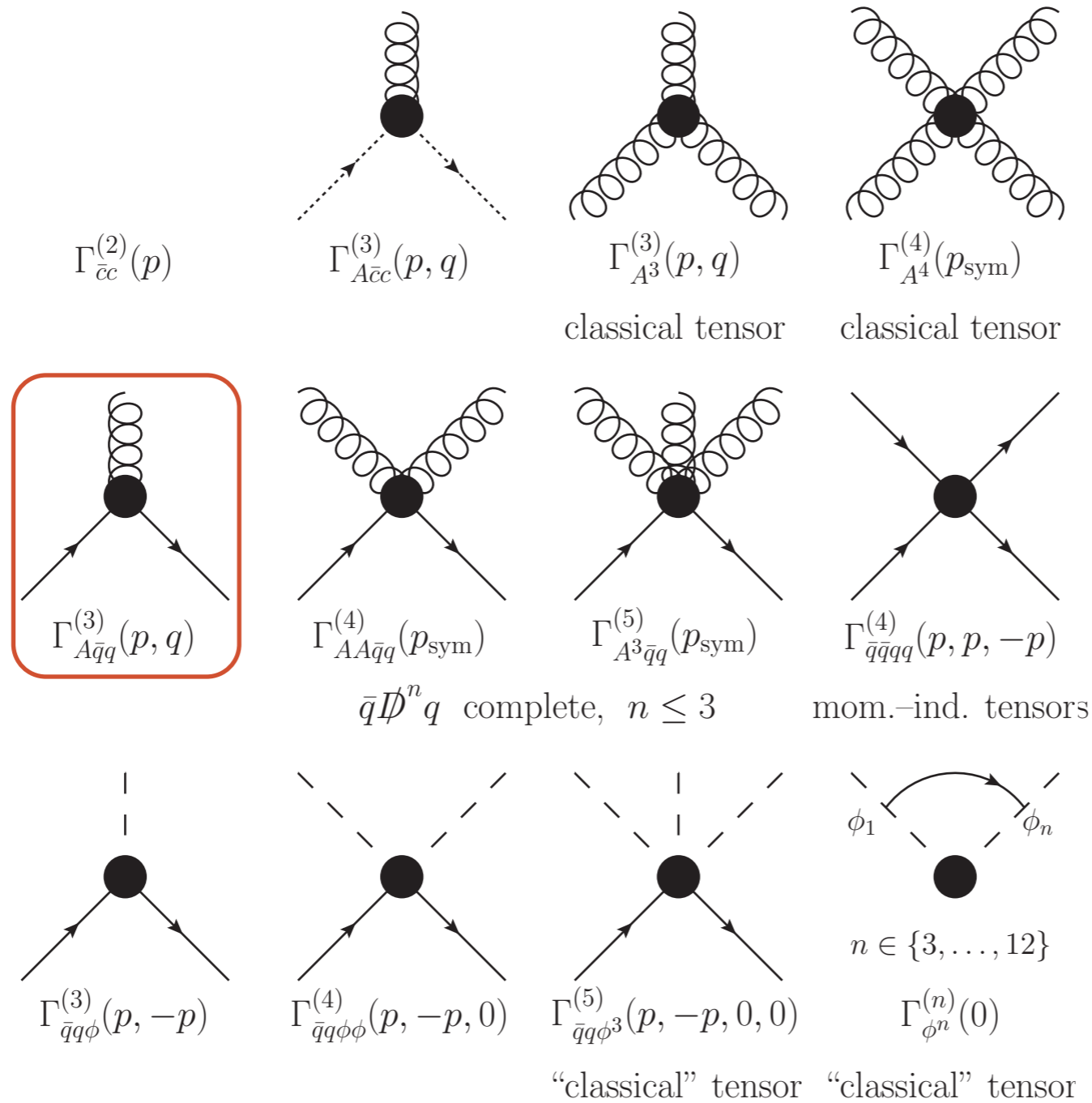
# QCD: Vertices



*Welches Schweinderl hätten's denn gerne?*

**Aiming at apparent convergence**

# QCD: Vertices



*Welches Schweinderl hätten's denn gerne?*

**Aiming at apparent convergence**

# Quark-gluon vertex

$$\left[ \Gamma_{\bar{q}qA}^{(3)} \right]_{\mu}^a(p, q) = 1_{2 \times 2}^{\text{flav}} T^a \sum_{i=1}^8 \lambda_i(p, q) \left[ \mathcal{T}_{\bar{q}qA}^{(i)} \right]_{\mu}(p, q)$$

## covariant expansion scheme

$$\bar{q}\not{D}q : \left[ \mathcal{T}_{\bar{q}qA}^{(1)} \right]_{\mu}(p, q) = -i \gamma_{\mu}$$

$$\bar{q}\not{D}^2q : \left[ \mathcal{T}_{\bar{q}qA}^{(2)} \right]_{\mu}(p, q) = (p - q)_{\mu} 1_{4 \times 4}$$

$$\bar{q}\not{D}^3q : \left[ \mathcal{T}_{\bar{q}qA}^{(5)} \right]_{\mu}(p, q) = i (\not{p} + \not{q})(p - q)_{\mu}$$

$$\left[ \mathcal{T}_{\bar{q}qA}^{(3)} \right]_{\mu}(p, q) = (\not{p} - \not{q})\gamma_{\mu}$$

$$\left[ \mathcal{T}_{\bar{q}qA}^{(6)} \right]_{\mu}(p, q) = i (\not{p} - \not{q})(p - q)_{\mu}$$

$$\left[ \mathcal{T}_{\bar{q}qA}^{(4)} \right]_{\mu}(p, q) = (\not{p} + \not{q})\gamma_{\mu}$$

$$\left[ \mathcal{T}_{\bar{q}qA}^{(7)} \right]_{\mu}(p, q) = \frac{i}{2} [\not{p}, \not{q}] \gamma_{\mu}$$

**Aiming at apparent convergence**

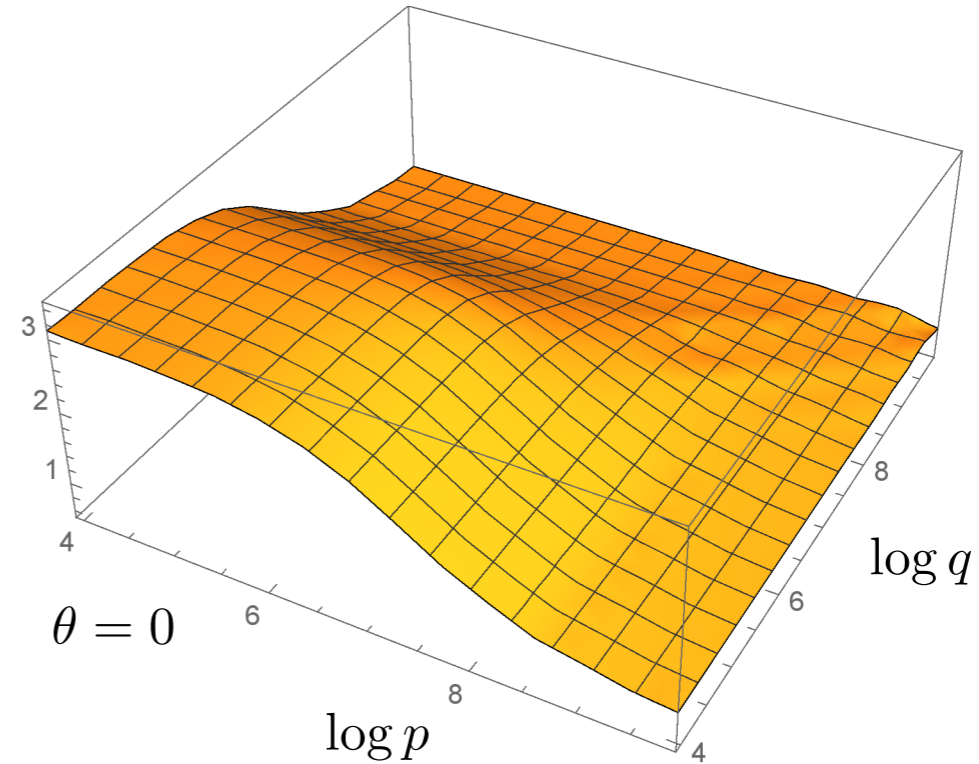
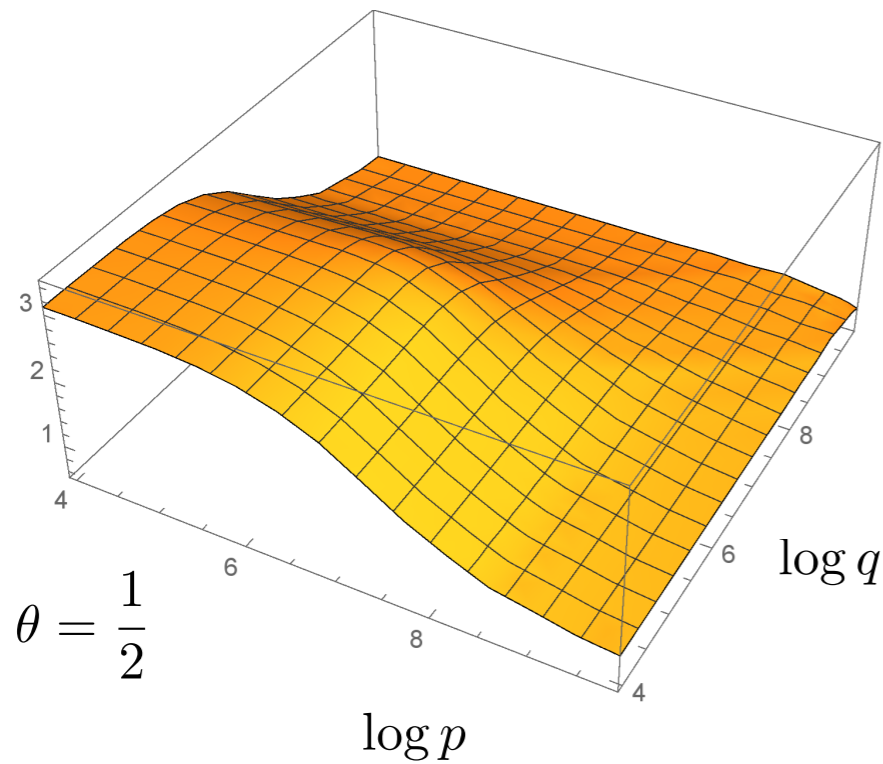
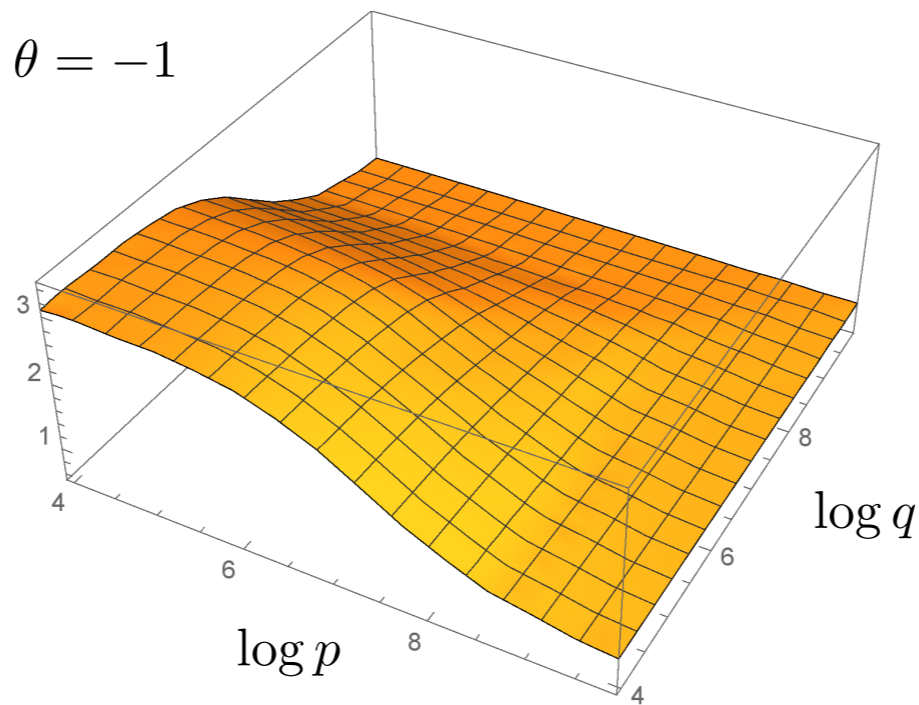
# Quark-gluon vertex

$$\theta = \frac{p \cdot q}{\sqrt{p^2 q^2}}$$

**p,q in MeV**

$\lambda_1(p, q)$

$\theta = -1$



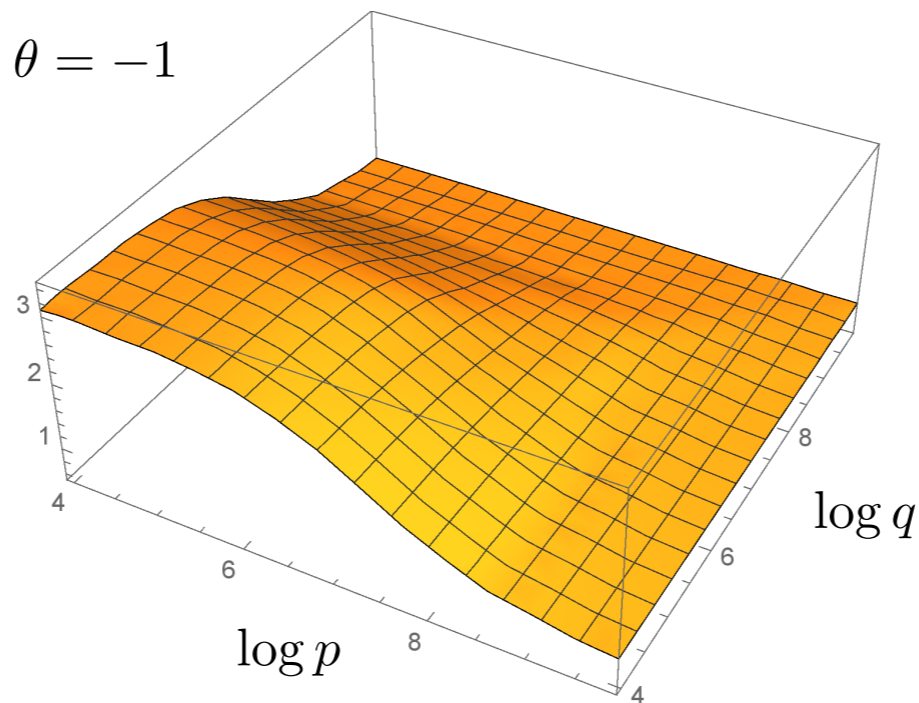
**Aiming at apparent convergence**

# Quark-gluon vertex

$$\theta = \frac{p \cdot q}{\sqrt{p^2 q^2}}$$

**p,q in MeV**

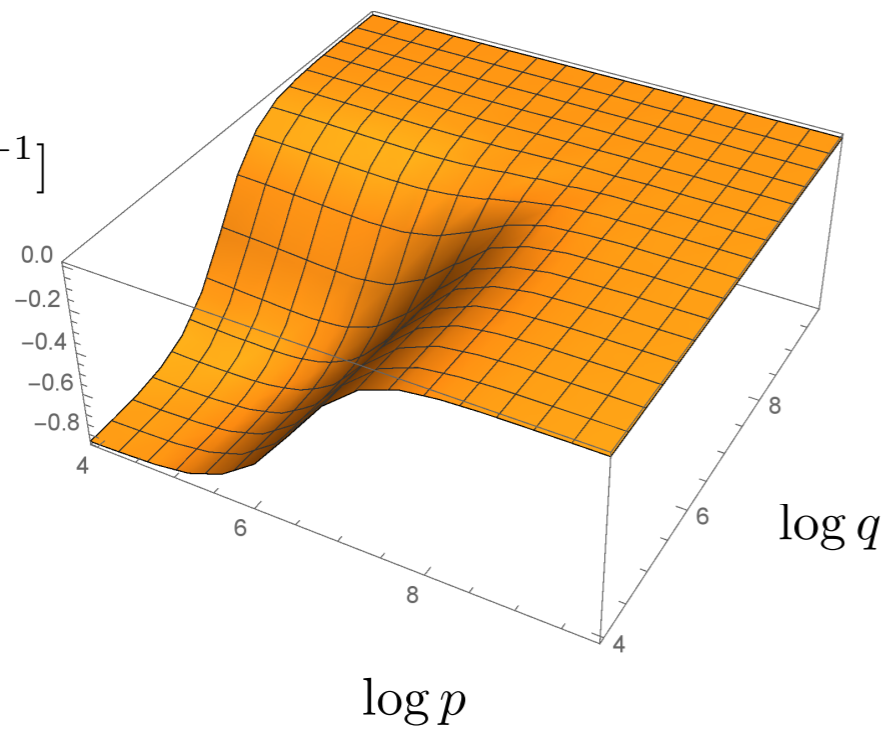
$$\theta = -1$$



$$\lambda_1(p, q)$$

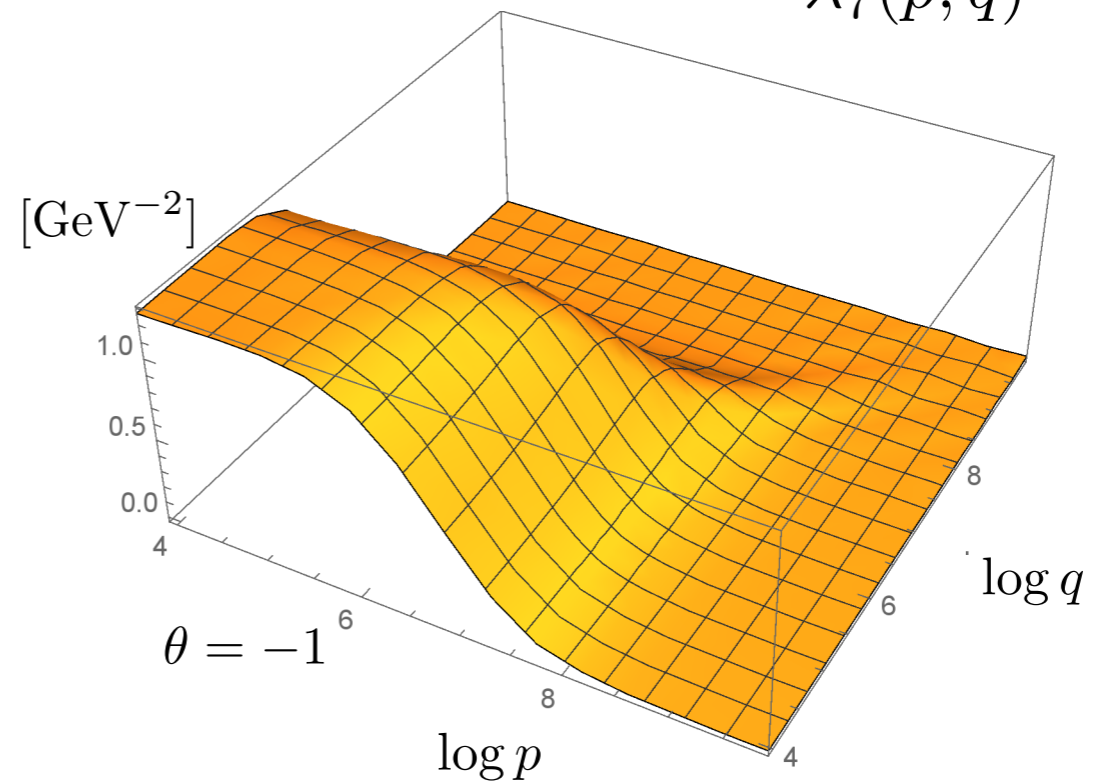
$$\lambda_4(p, q)$$

[GeV<sup>-1</sup>]



$$\lambda_7(p, q)$$

[GeV<sup>-2</sup>]



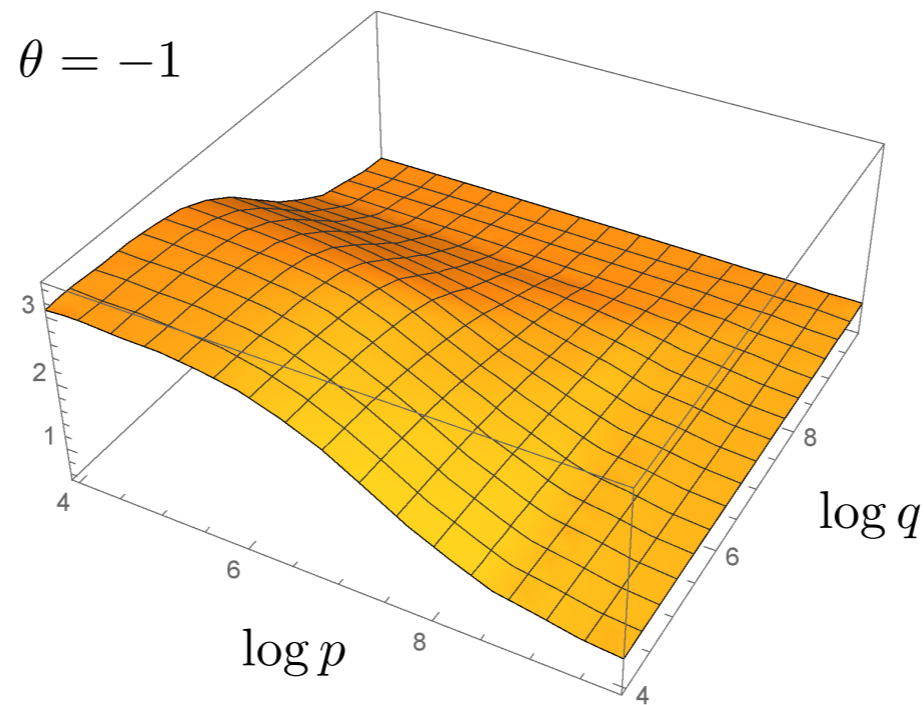
**Aiming at apparent convergence**



# QCD: Quark-gluon vertex

$$\theta = \frac{p \cdot q}{\sqrt{p^2 q^2}}$$

**p,q in MeV**



$$\lambda_1(p, q)$$

## up-to-date 1st principles works:

### FunMethods:

Williams, EPJ A51 (2015) 57  
Sanchis-Alepuz, Williams, PLB 749 (2015) 592  
Williams, Fischer, Heupel, PRD 93 (2016) 034026

Aguilar, Binosi, Ibanez, Papavassiliou, PRD 89 (2014) 065027  
Binosi, Chang, Papavassiliou, Qin, Roberts, PRD 95 (2017) 031501  
Aguilar, Cardona, Ferreira, Papavassiliou, arXiv:1610.06158

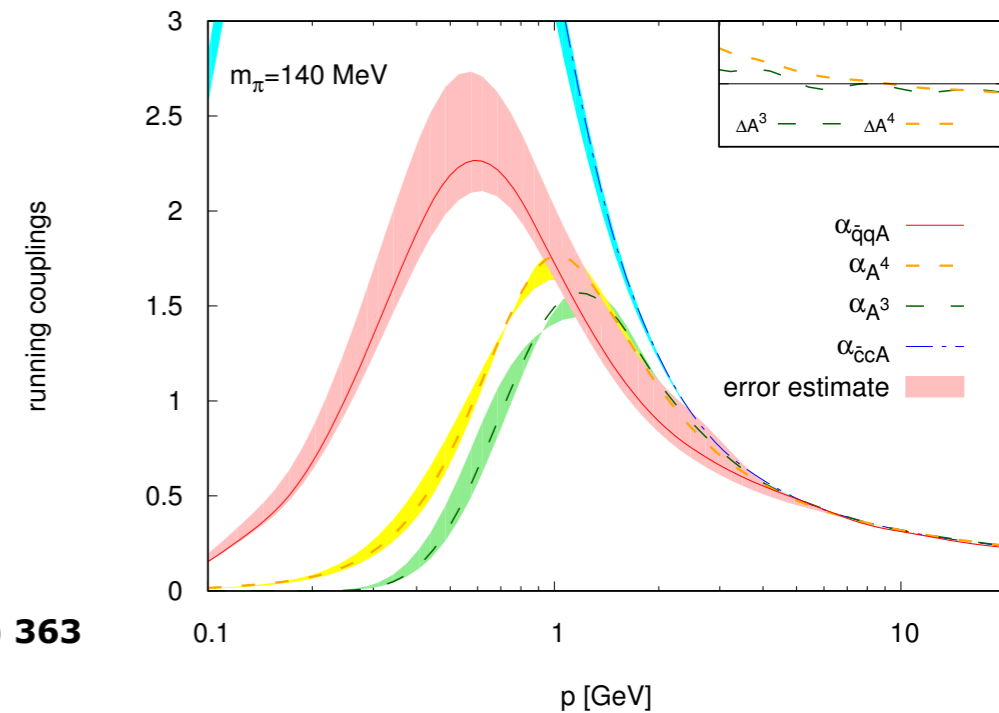
Mitter, JMP, Strodthoff, PRD 91 (2015) 054035

Pelaez, Tissier, Wschebor, PRD 92 (2015) 045012

Eichmann, Sanchis-Alepuz, Williams, Alkofer, Fischer, PPNP 91 (2016) 1

**lattice:** Oliveira, Kizilersü, Silva, Skullerud, Sternbeck, Williams, APP Suppl. 9 (2016) 363

## Beware of BRST



**Aiming at apparent convergence**

# **(III) Phase structure of QCD and dynamics**

---

- **Yang-Mills theory at finite temperature**
  - Order parameter potential for confinement
  - Correlation functions at finite temperature
  - Polyakov loop from functional methods
  
- **Application to the phase structure of QCD and dynamics\***
  - QCD-assisted hydrodynamics\*
  - QCD-assisted transport\*
  - QCD at imaginary chemical potential\*

# **Yang-Mills theory at finite temperature**

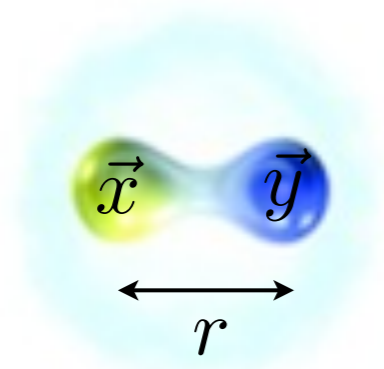


# **Order parameter potential for Confinement**

# Confinement

Free energy  $F_{q\bar{q}}$  of a quark - antiquark pair

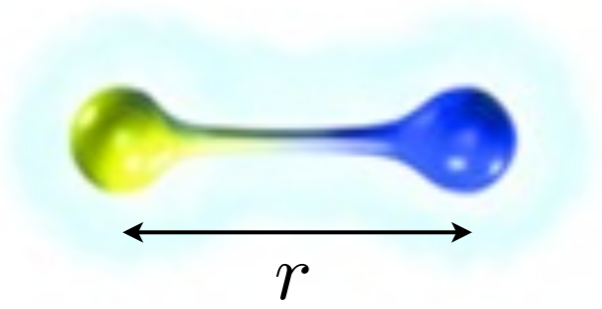
**Reminder**



$$F_{q\bar{q}} \simeq -\frac{1}{r}$$

**Order parameter**  $\sim \langle q \rangle'$

$$\Phi = e^{-\frac{1}{2T} F_{q\bar{q}}(\infty)}$$

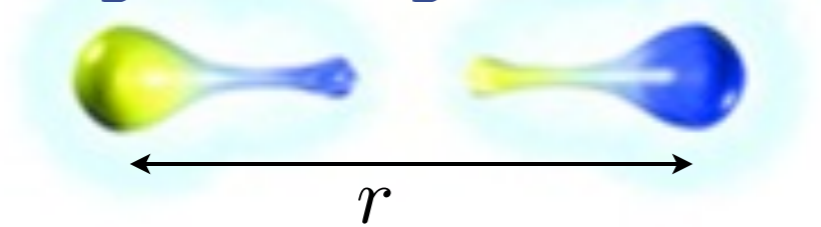


$$F_{q\bar{q}} \simeq \sigma r$$

• **Confinement**  $\Phi = 0$

• **Deconfinement**  $\Phi \neq 0$

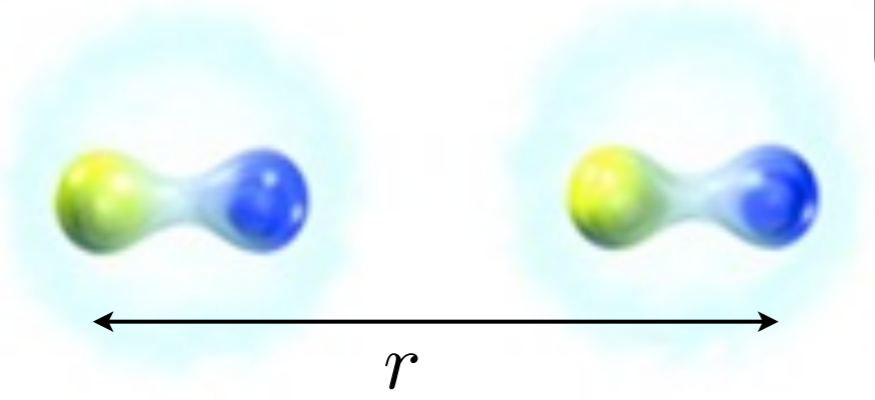
string breaking at  $r \approx 1\text{fm}$



$$F_{q\bar{q}} \simeq \text{const.}$$

**Polyakov loop**

$$\Phi = \frac{1}{3} \langle \text{Tr } \mathcal{P} \exp \{ ig \int_0^{1/T} dx_0 A_0 \} \rangle$$



# Confinement

## Order parameters

**Polyakov loop operator**

$$L[A_0] = \frac{1}{N_c} \text{tr} \mathcal{P} e^{ig \int_0^{1/T} dt A_0}$$

$$\Phi = \langle L[A_0] \rangle$$

**order parameter**

$L[\langle A_0 \rangle]$  **order parameter**

$$L[\langle A_0 \rangle] = 0 \longleftrightarrow \langle L[A_0] \rangle = 0$$
$$L[\langle A_0 \rangle] \geq \langle L[A_0] \rangle$$

**Braun, Gies, JMP '07**  
**Marhauser, JMP '08**

**up to lattice renormalisation**

$\langle A_0 \rangle$  **order parameter**

$$\left. \frac{\partial V[A_0]}{\partial A_0} \right|_{A_0 = \langle A_0 \rangle} = 0$$

$$V[A_0] = \frac{1}{\beta \text{Vol}_3} \Gamma[A_0]$$

**constant backgrounds**

**background Landau gauge**

# Confinement

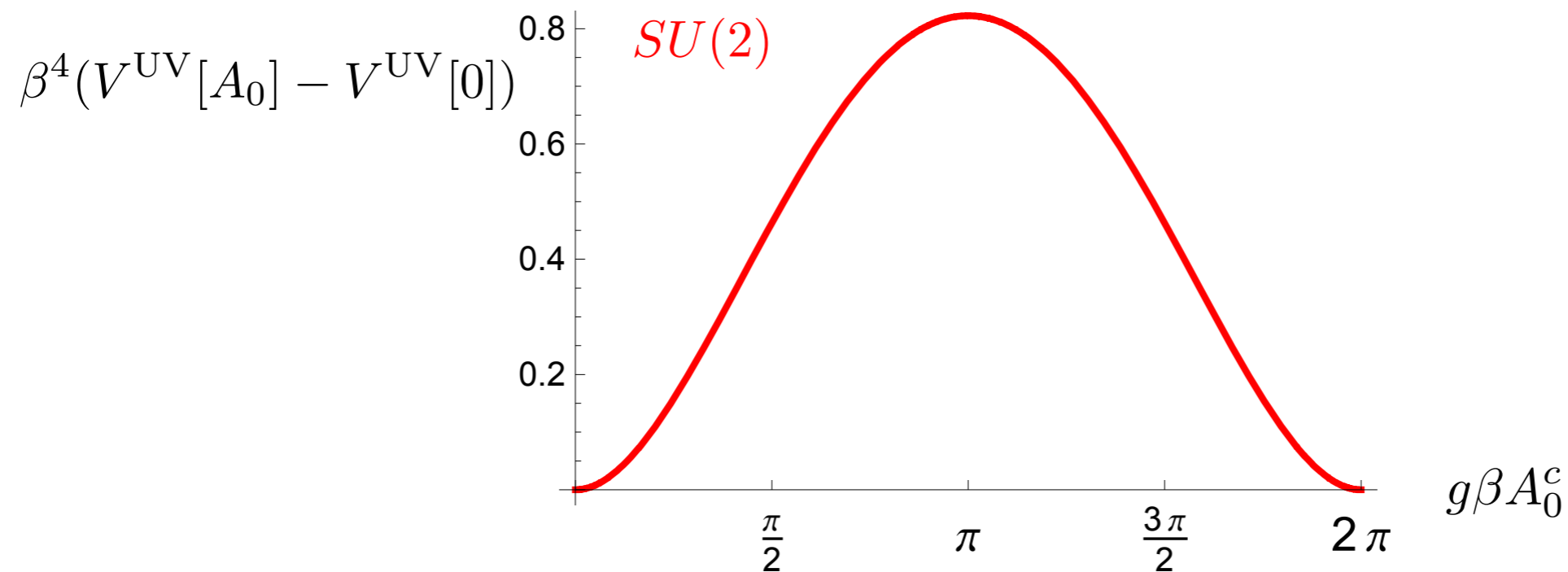
## Effective Polyakov loop potential

**One-loop**

$$V^{\text{UV}}[A_0] = \frac{1}{2\Omega} \text{Tr} \log S_{AA}^{(2)}[A_0] - \frac{1}{\Omega} \text{Tr} \log S_{C\bar{C}}^{(2)}[A_0]$$

**free energy**

**Gross, Pisarski, Yaffe '81  
Weiss '81**



$$SU(2) : \Phi[A_0] = \cos \frac{1}{2} \beta g A_0^c \quad \text{with} \quad A_0 = A_0^c \frac{\sigma_3}{2}$$

## Non-perturbative effective potential

$$V[A_0] = -\frac{1}{2} \text{Tr} \log \langle AA \rangle [A_0] + O(\partial_t \langle AA \rangle) + \text{Tr} \log \langle C\bar{C} \rangle [A_0] + O(\partial_t \langle C\bar{C} \rangle)$$

**free energy**

# Confinement

## Effective Polyakov loop potential

### Non-perturbative effective potential

$$V[A_0] \simeq -\frac{1}{2} \text{Tr} \log \langle AA \rangle [A_0] + O(\partial_t \langle AA \rangle) + \text{Tr} \log \langle C \bar{C} \rangle [A_0] + O(\partial_t \langle C \bar{C} \rangle)$$

**free energy**

**flow**

$$\frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k = \frac{1}{2} \text{Tr} \partial_t \log(\Gamma_k^{(2)}[\phi] + R_k) - \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t \Gamma_k^{(2)}[\phi]$$

# Confinement

## Effective Polyakov loop potential

### Non-perturbative effective potential

$$V[A_0] \simeq -\frac{1}{2} \text{Tr} \log \langle AA \rangle [A_0] + O(\partial_t \langle AA \rangle) + \text{Tr} \log \langle C \bar{C} \rangle [A_0] + O(\partial_t \langle C \bar{C} \rangle)$$

free energy

flow

$$\frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k = \frac{1}{2} \text{Tr} \partial_t \log(\Gamma_k^{(2)}[\phi] + R_k) - \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t \Gamma_k^{(2)}[\phi]$$

## Propagators

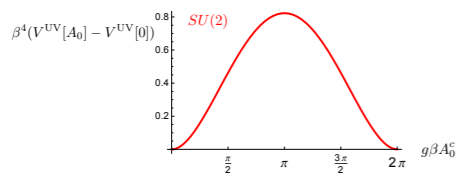
$$\langle AA \rangle [A_0] \simeq \frac{1}{-D_\mu^2(A_0)} \frac{1}{Z[-D_\mu^2(A_0)]}$$

## Integrals & sums

$$\text{Tr} f[-D_\mu^2(A_0)] = \sum_{\vec{p}, \pm} f[(2\pi T)^2 (n \pm \varphi)^2 + \vec{p}^2] + \varphi - \text{indep. terms}$$

$$g A_0 = \frac{\varphi}{2\pi T} \tau_{\text{ad}}^3$$

## One-loop result



$$\beta^4 V^{UV} [A_0] = -2 * 3 \left( \frac{\pi^2}{90} - \frac{2\pi^2}{3} \tilde{\varphi}^2 (1 - \tilde{\varphi})^2 \right)$$

$$\tilde{\varphi} = \varphi \pmod{1}$$

# Confinement

## Effective Polyakov loop potential

### Non-perturbative effective potential

$$V[A_0] \simeq -\frac{1}{2} \text{Tr} \log \langle AA \rangle [A_0] + O(\partial_t \langle AA \rangle) + \text{Tr} \log \langle C \bar{C} \rangle [A_0] + O(\partial_t \langle C \bar{C} \rangle)$$

free energy

flow

$$\frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k = \frac{1}{2} \text{Tr} \partial_t \log(\Gamma_k^{(2)}[\phi] + R_k) - \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t \Gamma_k^{(2)}[\phi]$$

## Propagators

$$\langle AA \rangle [A_0] \simeq \frac{1}{-D_\mu^2(A_0)} \frac{1}{Z[-D_\mu^2(A_0)]}$$

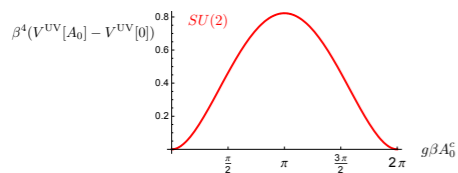
## Integrals & sums

$$\text{Tr} f[-D_\mu^2(A_0)] = \sum_{\vec{p}, \pm} f[(2\pi T)^2 (n \pm \varphi)^2 + \vec{p}^2] + \varphi - \text{indep. terms}$$

$$\beta^4 p_{\text{SB}}$$

$$g A_0 = \frac{\varphi}{2\pi T} \tau_{\text{ad}}^3$$

## One-loop result



$$\beta^4 V^{UV} [A_0] = -2 * 3 \left( \frac{\pi^2}{90} - \frac{2\pi^2}{3} \tilde{\varphi}^2 (1 - \tilde{\varphi})^2 \right)$$

$$\tilde{\varphi} = \varphi \pmod{1}$$

# Confinement

## Effective Polyakov loop potential

### Non-perturbative effective potential

$$V[A_0] \simeq -\frac{1}{2} \text{Tr} \log \langle AA \rangle [A_0] + O(\partial_t \langle AA \rangle) + \text{Tr} \log \langle C \bar{C} \rangle [A_0] + O(\partial_t \langle C \bar{C} \rangle)$$

free energy

flow

$$\frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t R_k = \frac{1}{2} \text{Tr} \partial_t \log(\Gamma_k^{(2)}[\phi] + R_k) - \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)}[\phi] + R_k} \partial_t \Gamma_k^{(2)}[\phi]$$

## Propagators

$$\langle AA \rangle [A_0] \simeq \frac{1}{-D_\mu^2(A_0)} \frac{1}{Z[-D_\mu^2(A_0)]}$$

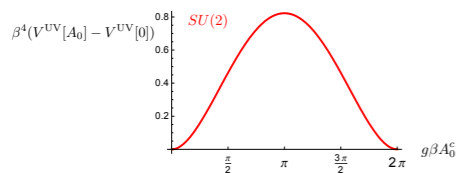
## Integrals & sums

$$\text{Tr} f[-D_\mu^2(A_0)] = \sum_{\vec{p}, \pm} f[(2\pi T)^2(n \pm \varphi)^2 + \vec{p}^2] + \varphi - \text{indep. terms}$$

$$N_c^2 - 1$$

$$g A_0 = \frac{\varphi}{2\pi T} \tau_{\text{ad}}^3$$

## One-loop result



$$\beta^4 V^{UV} [A_0] = -2 * 3 \left( \frac{\pi^2}{90} - \frac{2\pi^2}{3} \tilde{\varphi}^2 (1 - \tilde{\varphi})^2 \right)$$

$$\tilde{\varphi} = \varphi \pmod{1}$$



# Confinement

## Effective Polyakov loop potential

### Non-perturbative effective potential

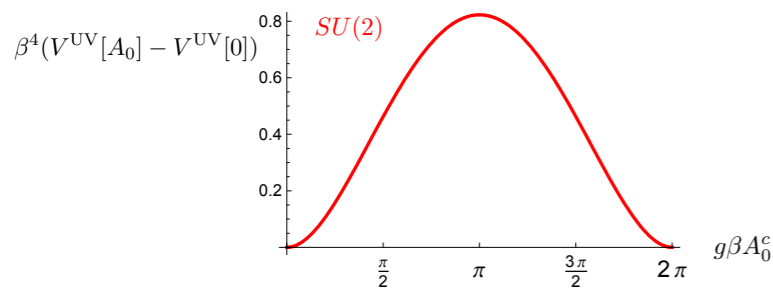
$$V[A_0] \simeq -\frac{1}{2} \text{Tr} \log \langle AA \rangle [A_0] + O(\partial_t \langle AA \rangle) + \text{Tr} \log \langle C\bar{C} \rangle [A_0] + O(\partial_t \langle C\bar{C} \rangle)$$

free energy

### Confinement criterion

Braun, Gies, JMP '07

Fister, JMP '13



$$\beta^4 V^{UV}[A_0] = 2 * 3 \left( \frac{\pi^2}{90} - \frac{2\pi^2}{3} \tilde{\varphi}^2 (1 - \tilde{\varphi})^2 \right)$$

**2** = 2 transversal physical polarisations + 1 transversal (zero mode) + 1 longitudinal - 2 ghosts

# Confinement

## Effective Polyakov loop potential

### Non-perturbative effective potential

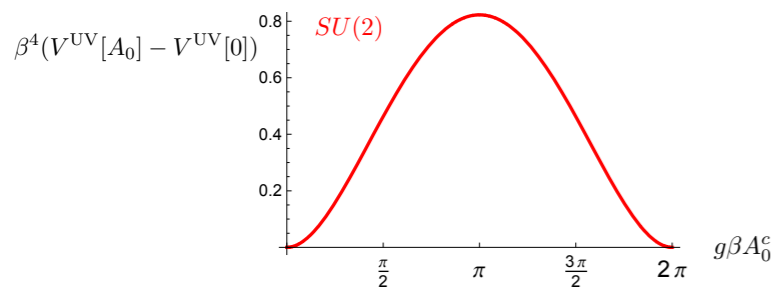
$$V[A_0] \simeq -\frac{1}{2} \text{Tr} \log \langle AA \rangle [A_0] + O(\partial_t \langle AA \rangle) + \text{Tr} \log \langle C\bar{C} \rangle [A_0] + O(\partial_t \langle C\bar{C} \rangle)$$

free energy

### Confinement criterion

Braun, Gies, JMP '07

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$$\beta^4 V^{UV} [A_0] = -2 * 3 \left( \frac{\pi^2}{90} - \frac{2\pi^2}{3} \tilde{\varphi}^2 (1 - \tilde{\varphi})^2 \right)$$

**2** = 2 transversal physical polarisations + 1 transversal (zero mode) + 1 longitudinal - 2 ghosts

**Glueon contribution deconfines**

**Ghost contribution confines**

**Confinement**  $\longleftrightarrow$  **suppression of the gluon relative to the ghost**

# **Correlation functions at finite temperature**

# YM-theory: gluonic correlation functions

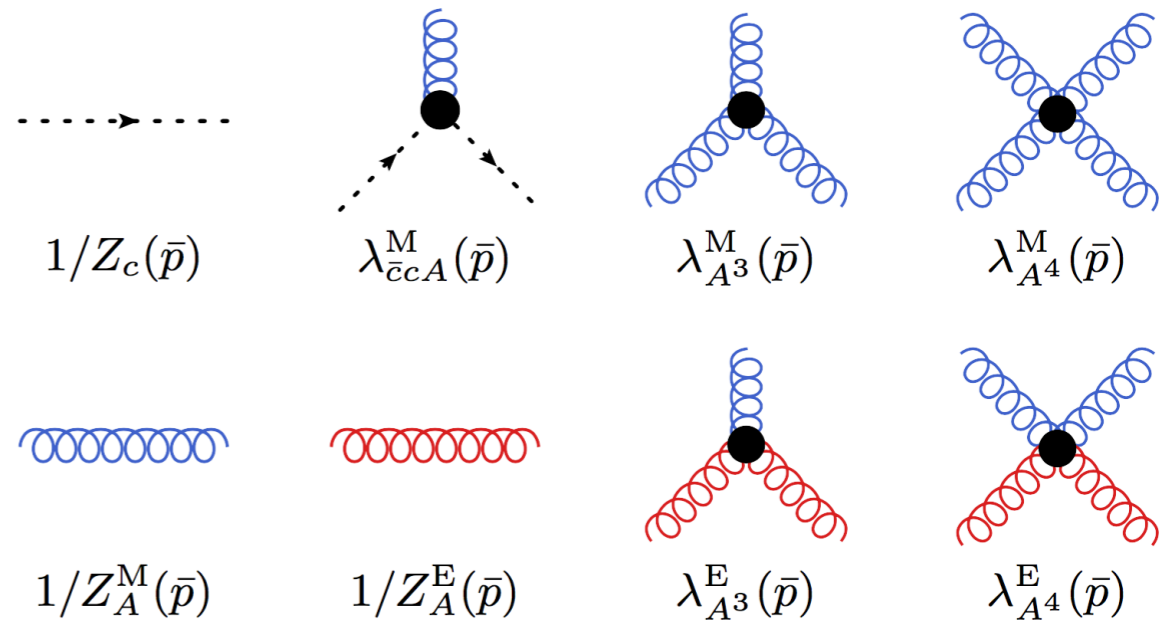
$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---}$$

$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} - 2 \text{---} \text{---} \text{---} + \frac{1}{2} \text{---} \text{---} \text{---}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$



**Aiming at apparent convergence**

# YM-theory: gluonic correlation functions

$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---}$$

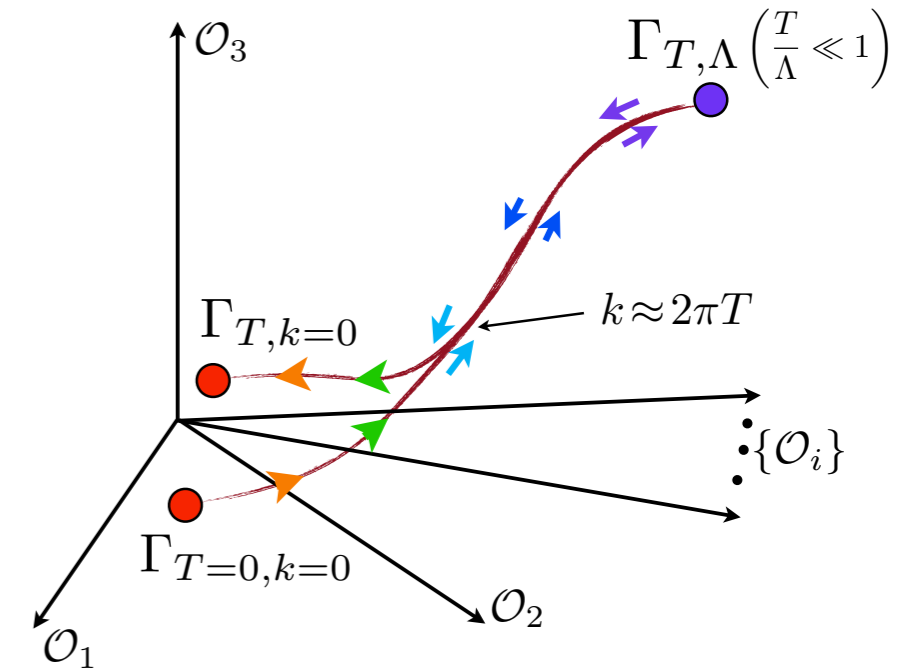
$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} - 2 \text{---} \text{---} \text{---} + \frac{1}{2} \text{---} \text{---} \text{---}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

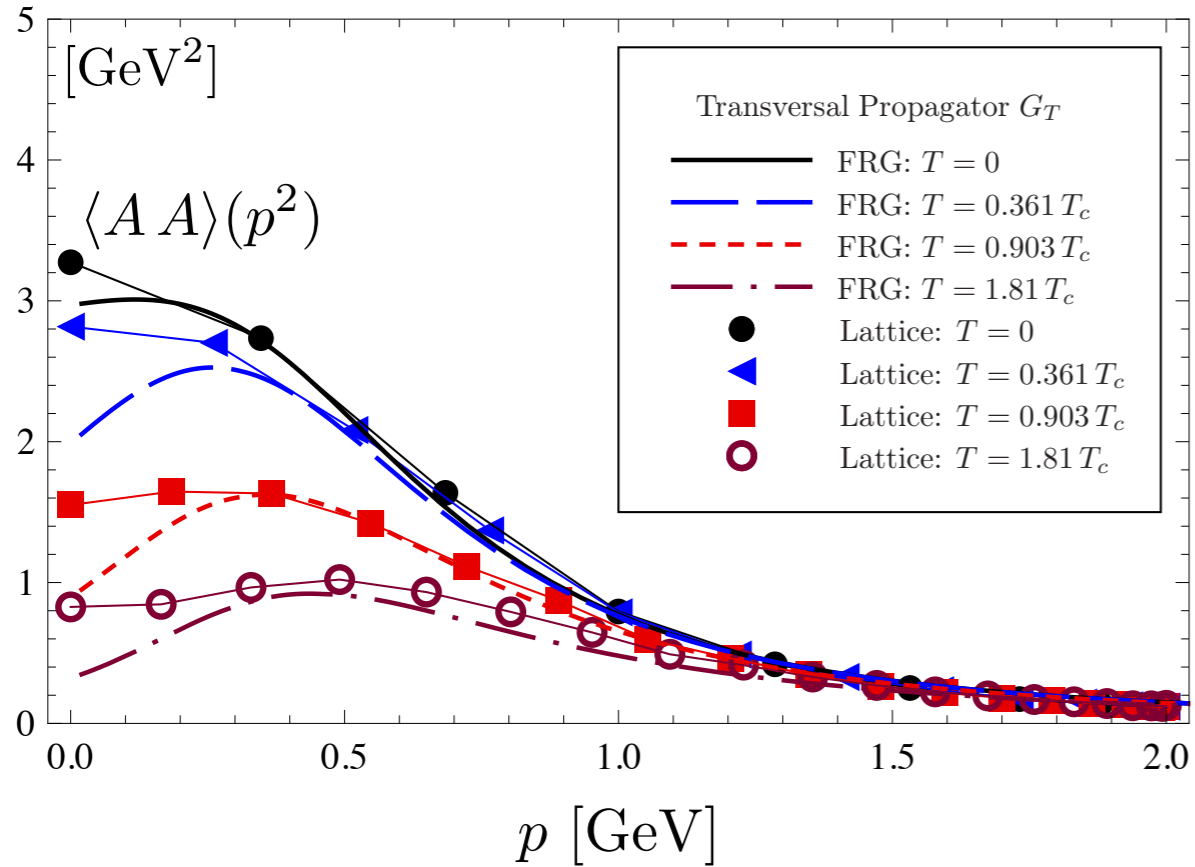
## Thermal flows



**Aiming at apparent convergence**

# Euclidean gluon propagator at finite T

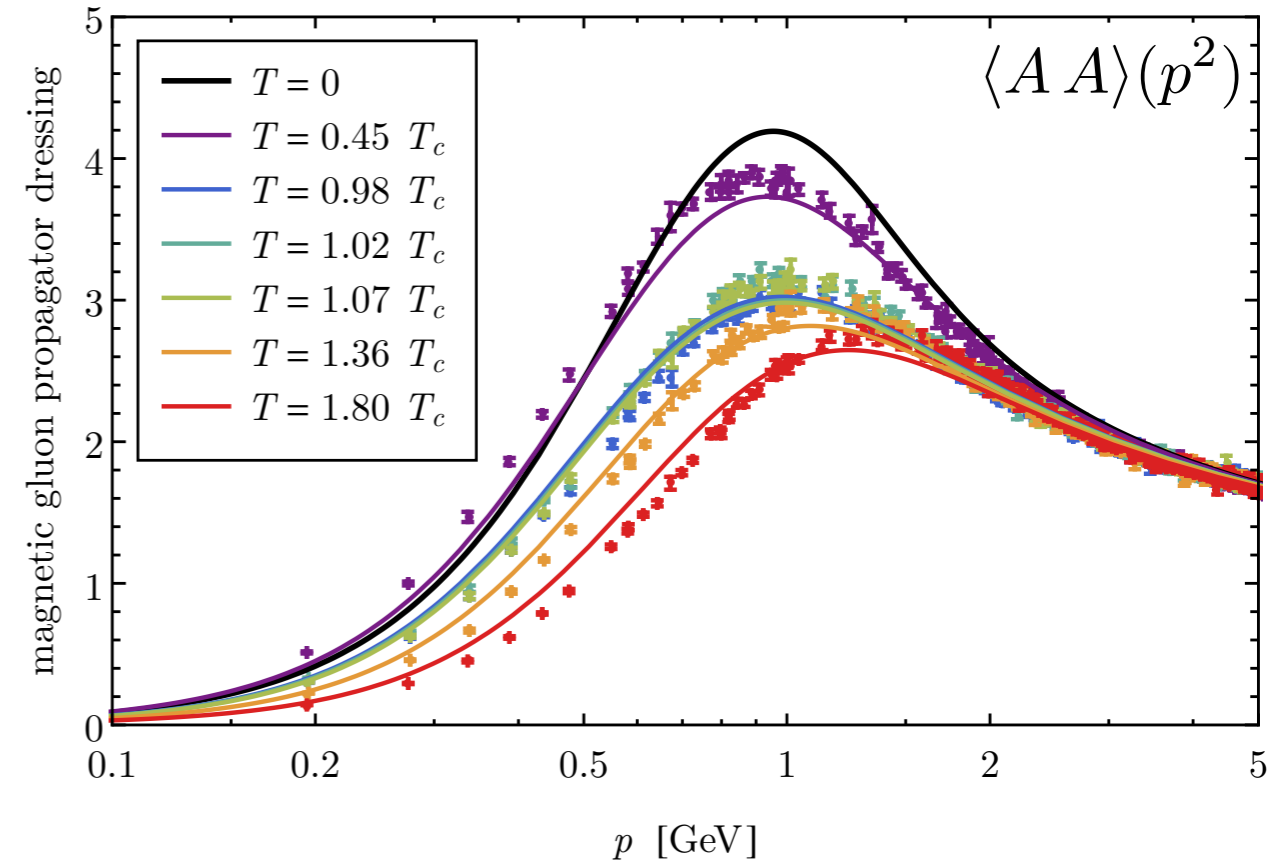
## chromo-magnetic propagator



Fister, JMP, arXiv:1112.5440

Lattice: Maas, JMP, Smekal, Spielmann, PRD 85 (2012) 034037

CF model: Reinoso, Serreau, Tissier, Tresmontant, PRD 95 (2017) 045014



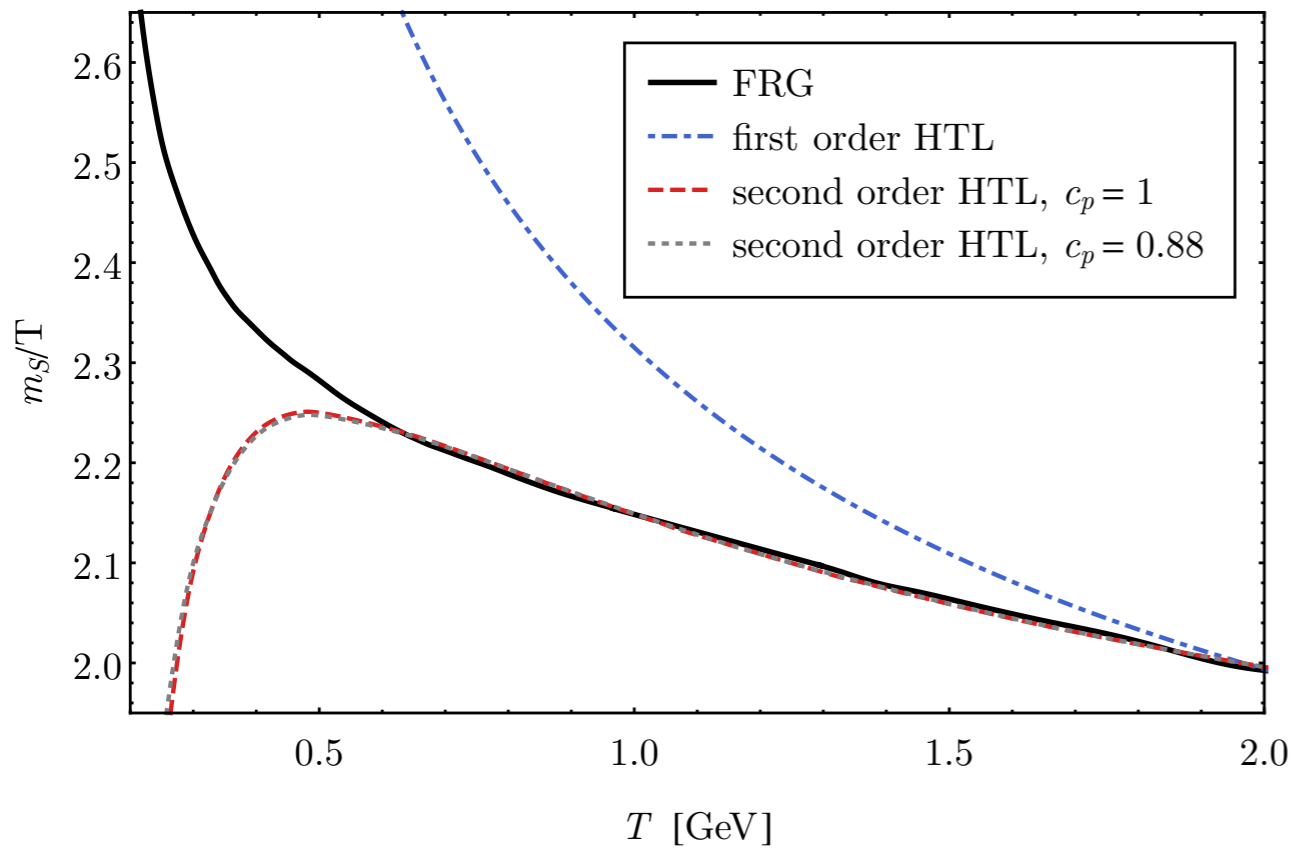
Lattice: Silva, Oliviera, Bicudo, Cardoso, PRD89 (2014) 7, 074503

**Aiming at apparent convergence**

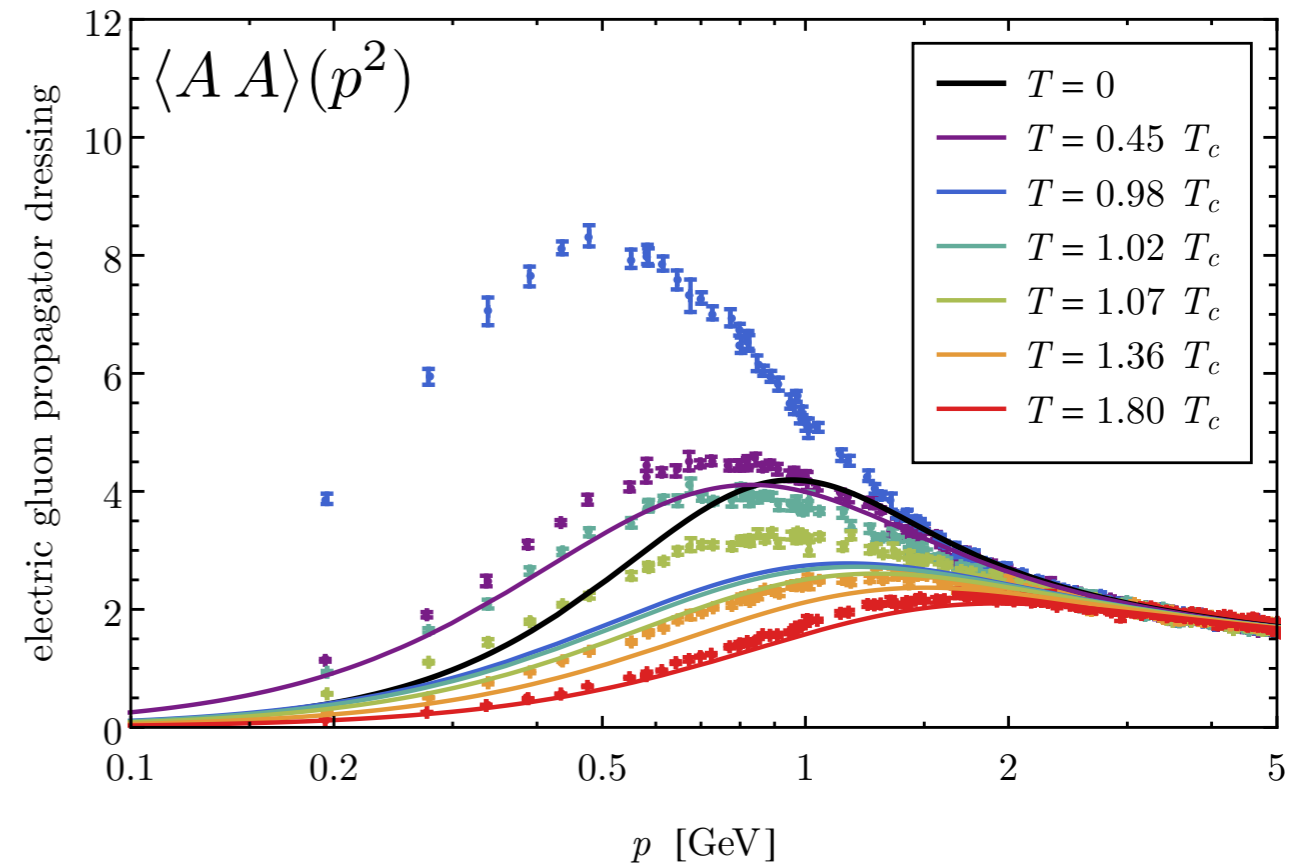
Cyrol, Fister, Mitter, JMP, Strodthoff, PRD 97 (2018) 5, 054015

# Euclidean gluon propagator at finite T

Debye mass (chromo-electric)



chromo-electric propagator

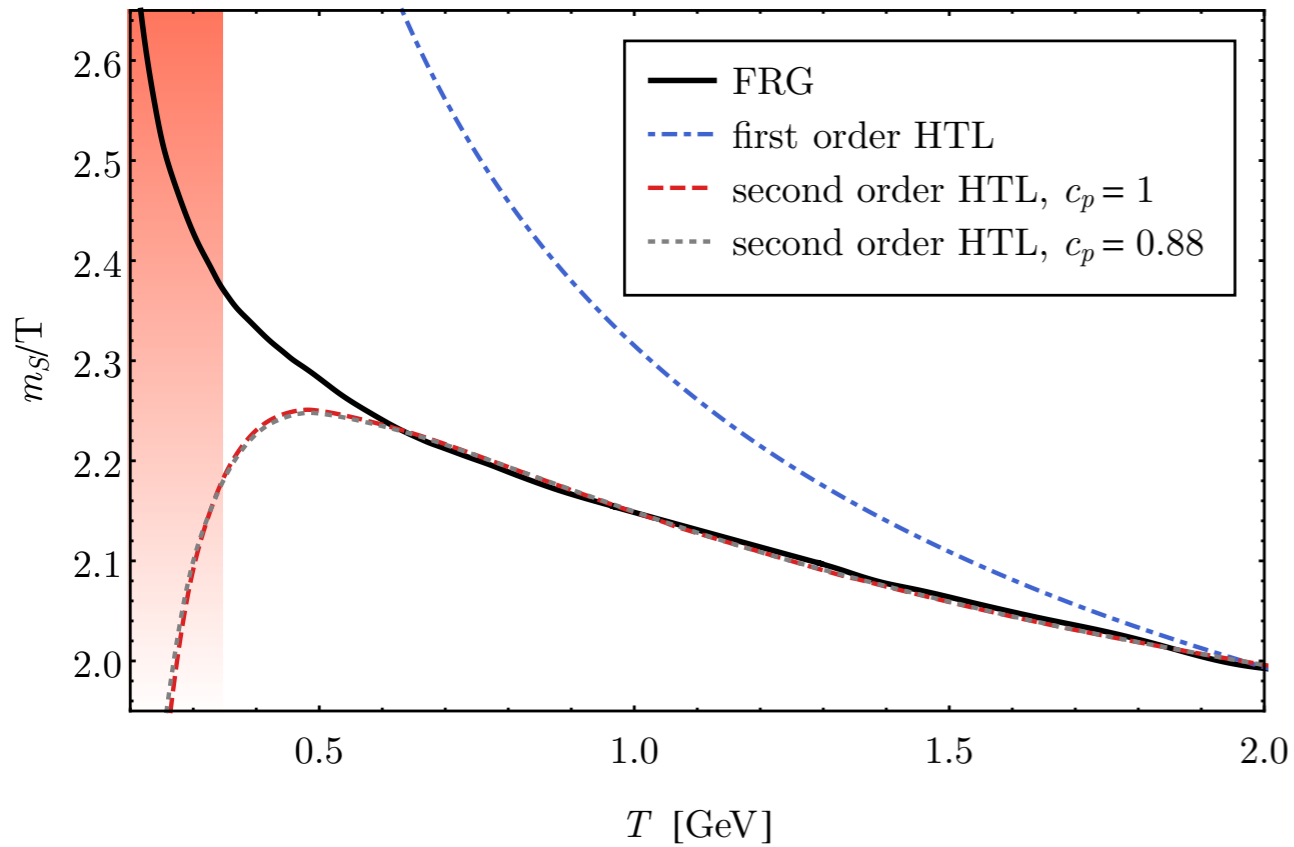


Lattice: Silva, Oliveira, Bicudo, Cardoso, PRD89 (2014) 7, 074503

Cyrol, Fister, Mitter, JMP, Strodthoff, PRD 97 (2018) 5, 054015

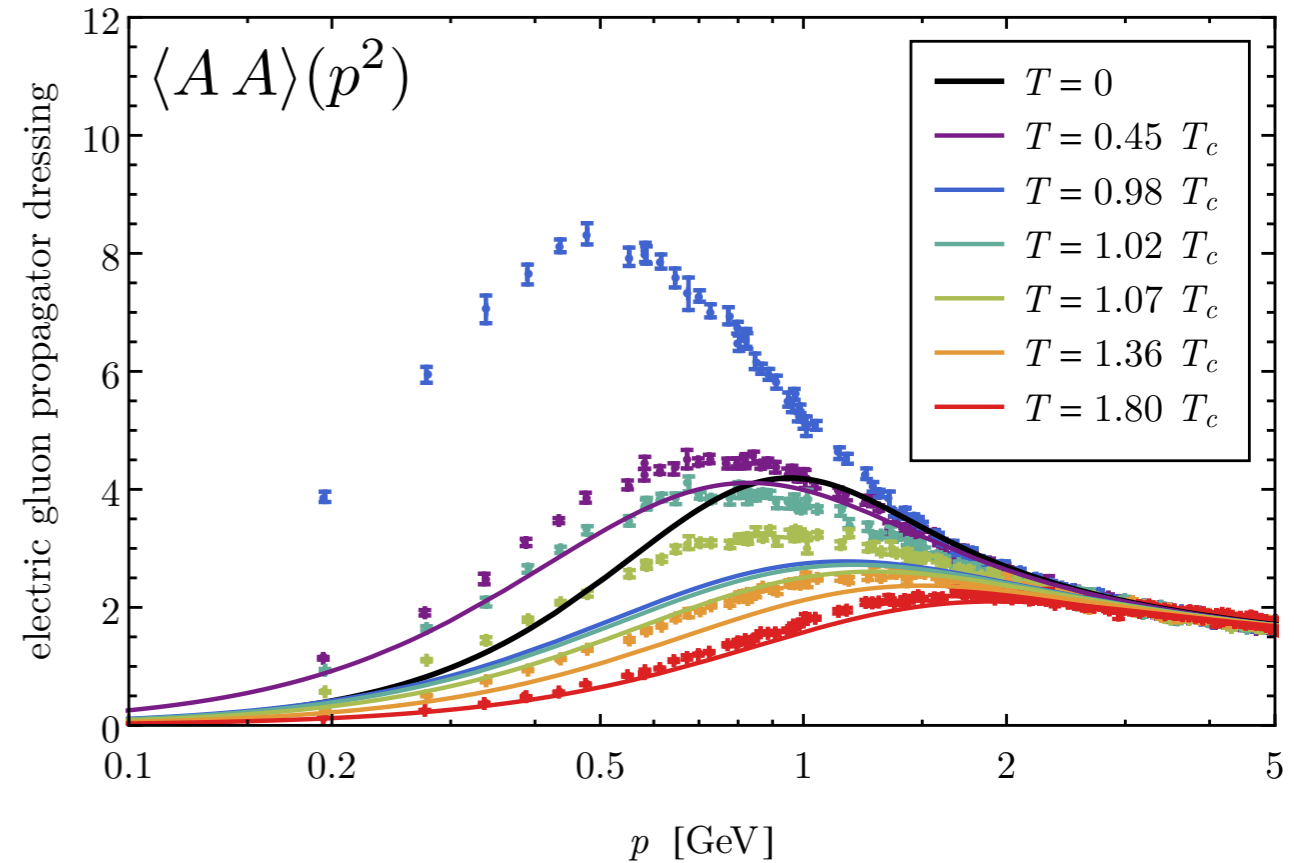
# Euclidean gluon propagator at finite T

Debye mass (chromo-electric)



$$\langle A_0 \rangle \neq 0$$

chromo-electric propagator



Lattice: Silva, Oliveira, Bicudo, Cardoso, PRD89 (2014) 7, 074503

Cyrol, Fister, Mitter, JMP, Strodthoff, PRD 97 (2018) 5, 054015



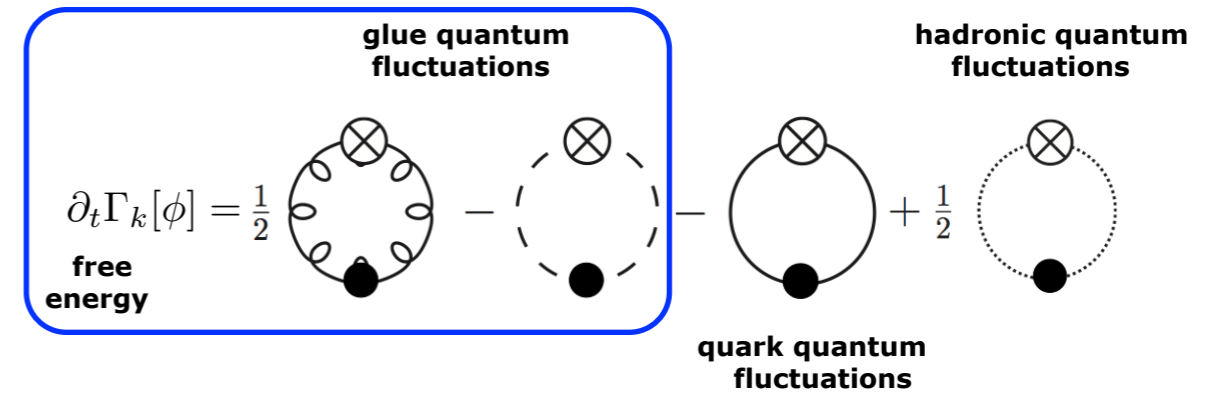
# **Polyakov loop from functional methods**

# Confinement

**FRG:** Braun, Gies, JMP, PLB 684 (2010) 262

**FRG, DSE, 2PI:** Fister, JMP, PRD 88 (2013) 045010

$$L[A_0] = \frac{1}{N_c} \text{tr} \mathcal{P} e^{i g \int_0^\beta \mathbf{A}_0(\mathbf{x})}$$

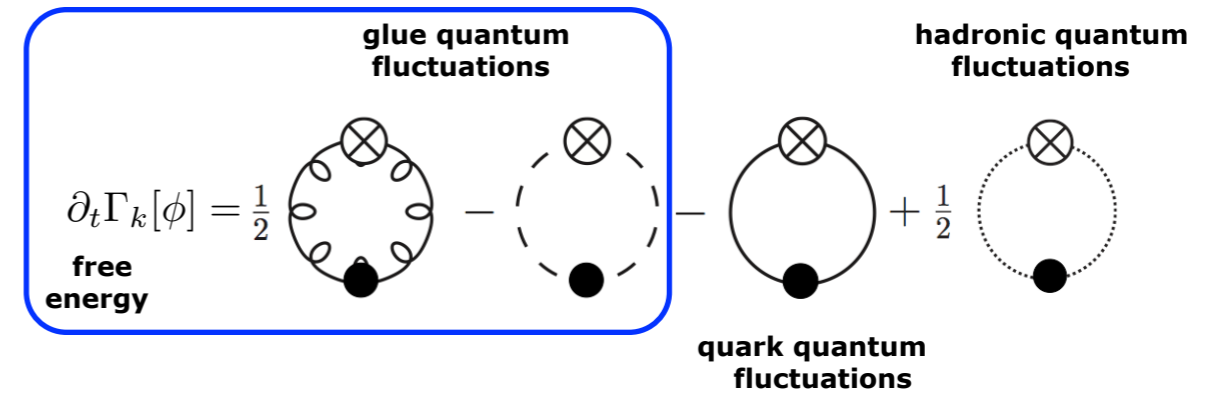


# Confinement

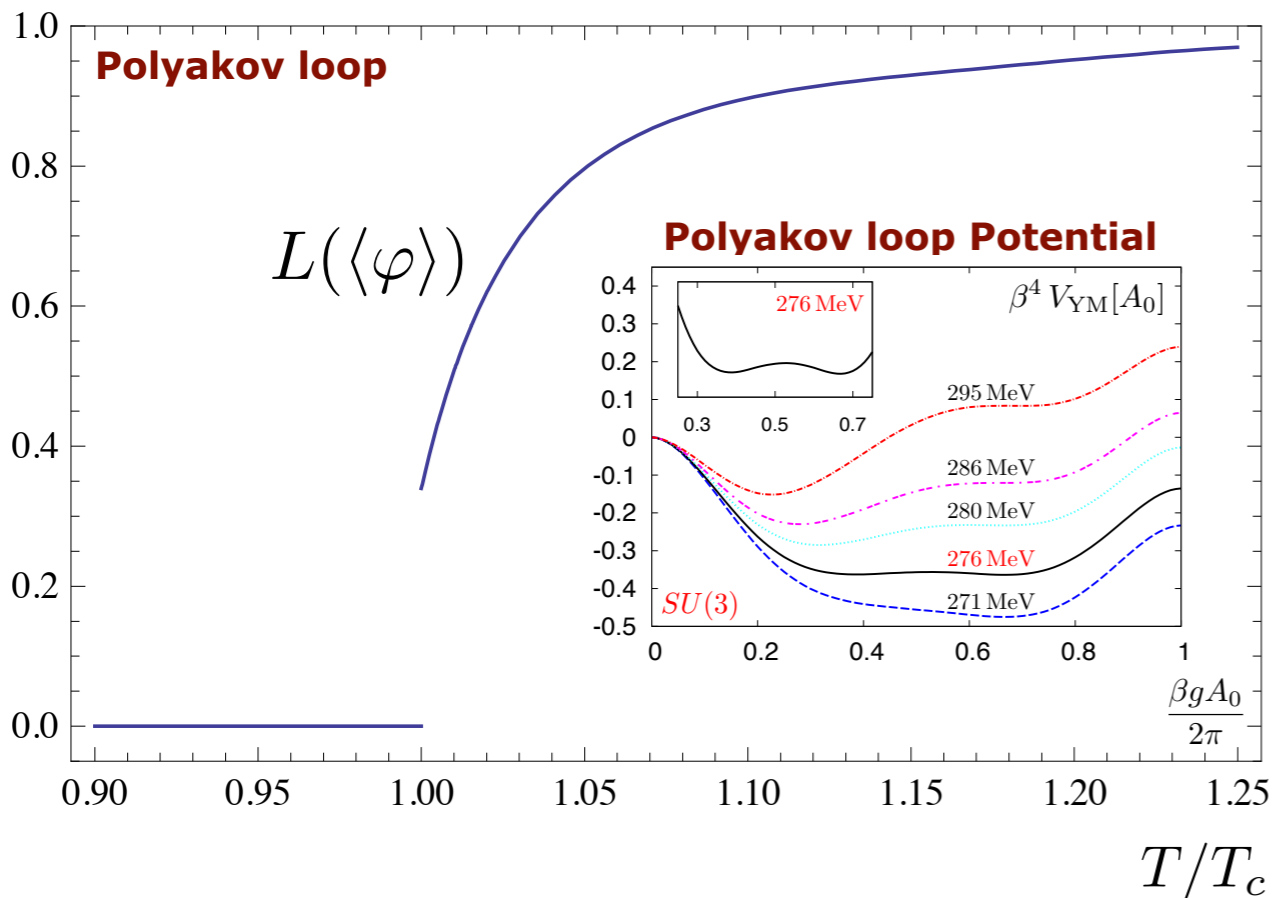
**FRG: Braun, Gies, JMP, PLB 684 (2010) 262**

**FRG, DSE, 2PI: Fister, JMP, PRD 88 (2013) 045010**

$$L[A_0] = \frac{1}{N_c} \text{tr} \mathcal{P} e^{i g \int_0^\beta A_0(x)}$$



$$\mathcal{P} e^{i g \int_0^\beta A_0(x)} = e^{i\varphi}$$



$$T_c/\sqrt{\sigma} = 0.658 \pm 0.023$$

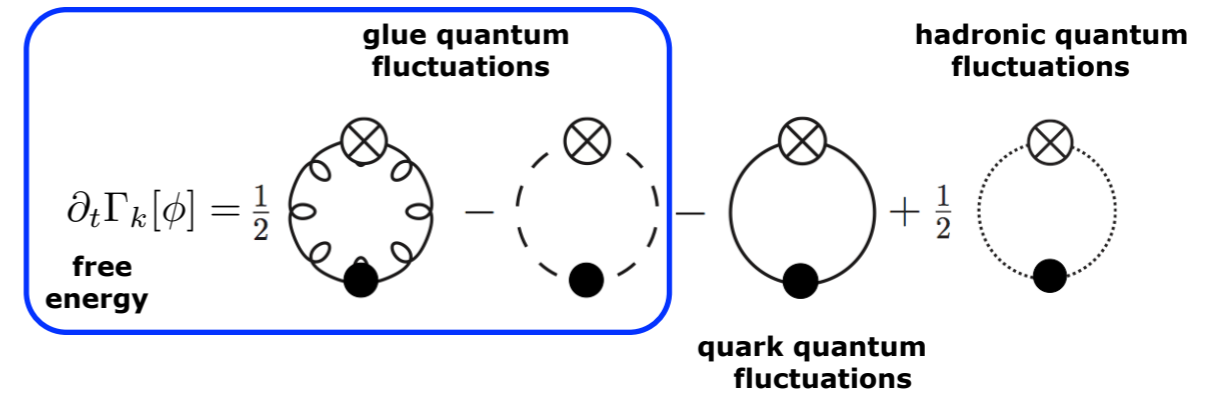
$$\text{lattice : } T_c/\sqrt{\sigma} = 0.646$$

# Confinement

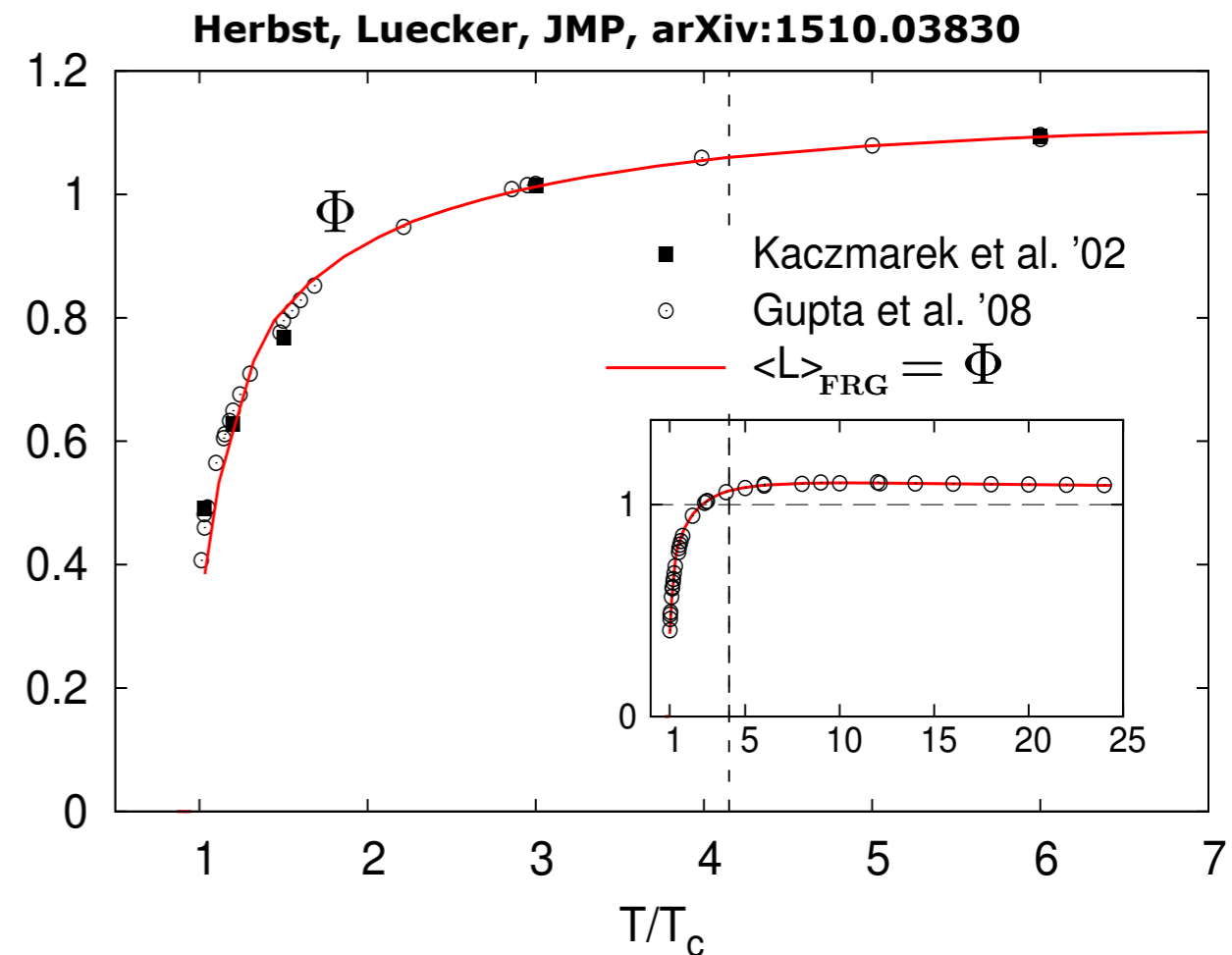
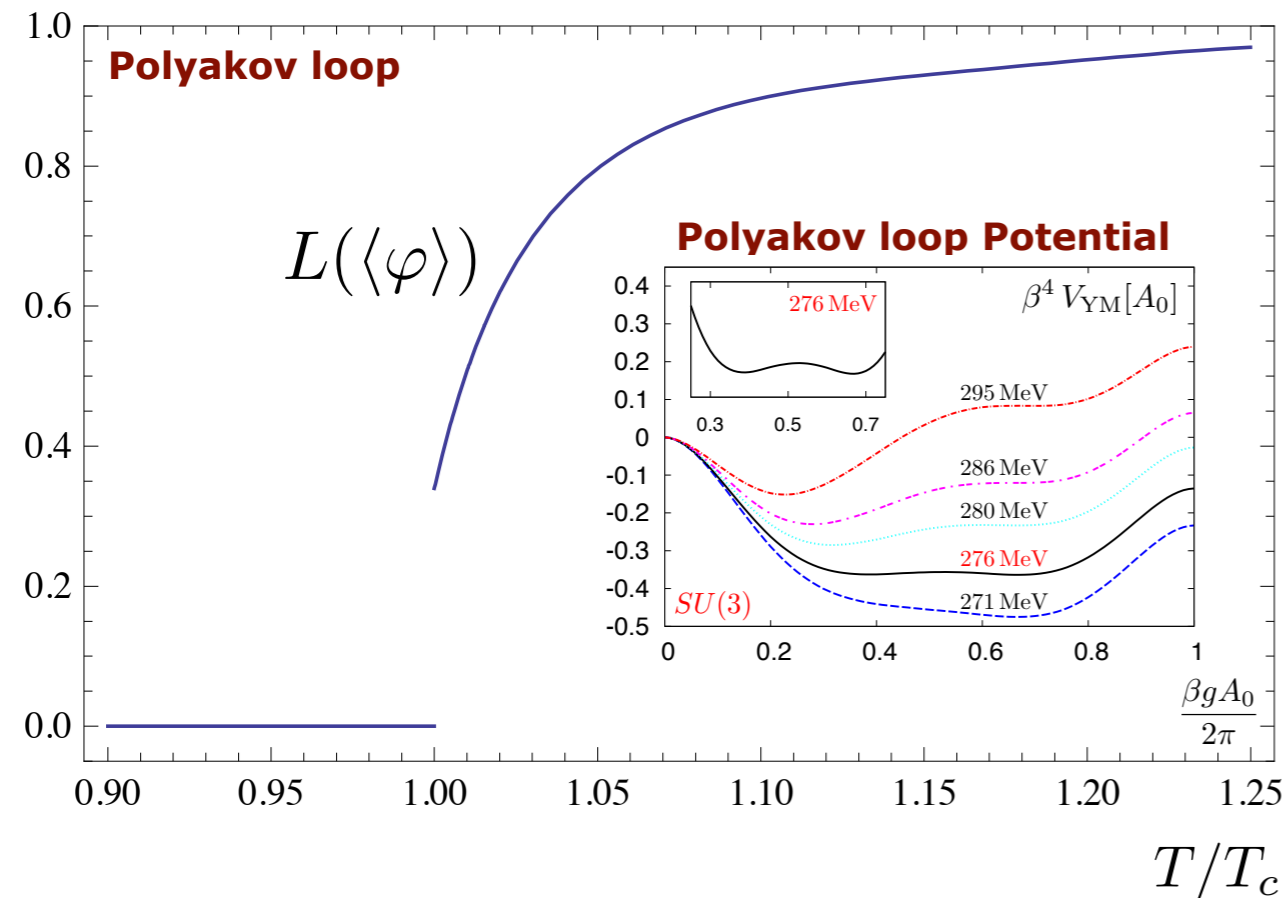
**FRG: Braun, Gies, JMP, PLB 684 (2010) 262**

**FRG, DSE, 2PI: Fister, JMP, PRD 88 (2013) 045010**

$$L[A_0] = \frac{1}{N_c} \text{tr} \mathcal{P} e^{i g \int_0^\beta A_0(x)}$$



$$\mathcal{P} e^{i g \int_0^\beta A_0(x)} = e^{i\varphi}$$



# Confinement

Herbst, Luecker, JMP, arXiv:1510.03830

Flow equation for the Polyakov loop expectation value

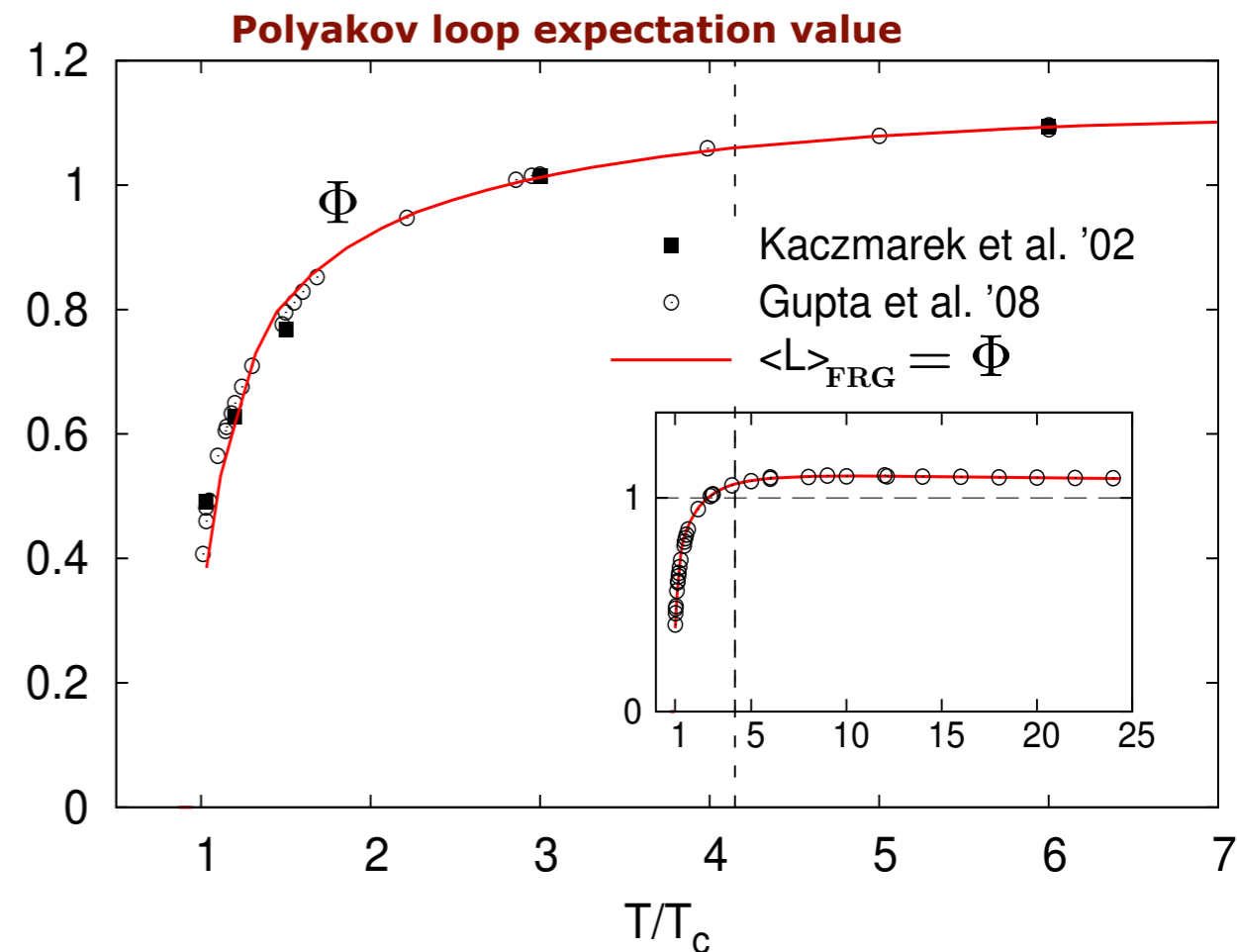
$$\partial_t \langle L[A_0] \rangle = - \frac{1}{2} \left[ \frac{\delta^2 \langle L[A_0] \rangle}{\delta A^2} - \frac{\delta^2 \langle L[A_0] \rangle}{\delta c \delta \bar{c}} \right]$$

Flow equation for composite operators

JMP, AP 322 (2007) 2831

Igarashi, Itoh, Sonoda, PTP Suppl. 181 (2010) 1

Pagani, PRD 94 (2016) 045001



# Confinement

Herbst, Luecker, JMP, arXiv:1510.03830

## Flow equation for the Polyakov loop expectation value

$$\partial_t \langle L[A_0] \rangle = - \frac{1}{2} \left( \frac{\delta^2 \langle L[A_0] \rangle}{\delta A^2} - \frac{\delta^2 \langle L[A_0] \rangle}{\delta c \delta \bar{c}} \right)$$

## Flow equation for composite operators

JMP, AP 322 (2007) 2831

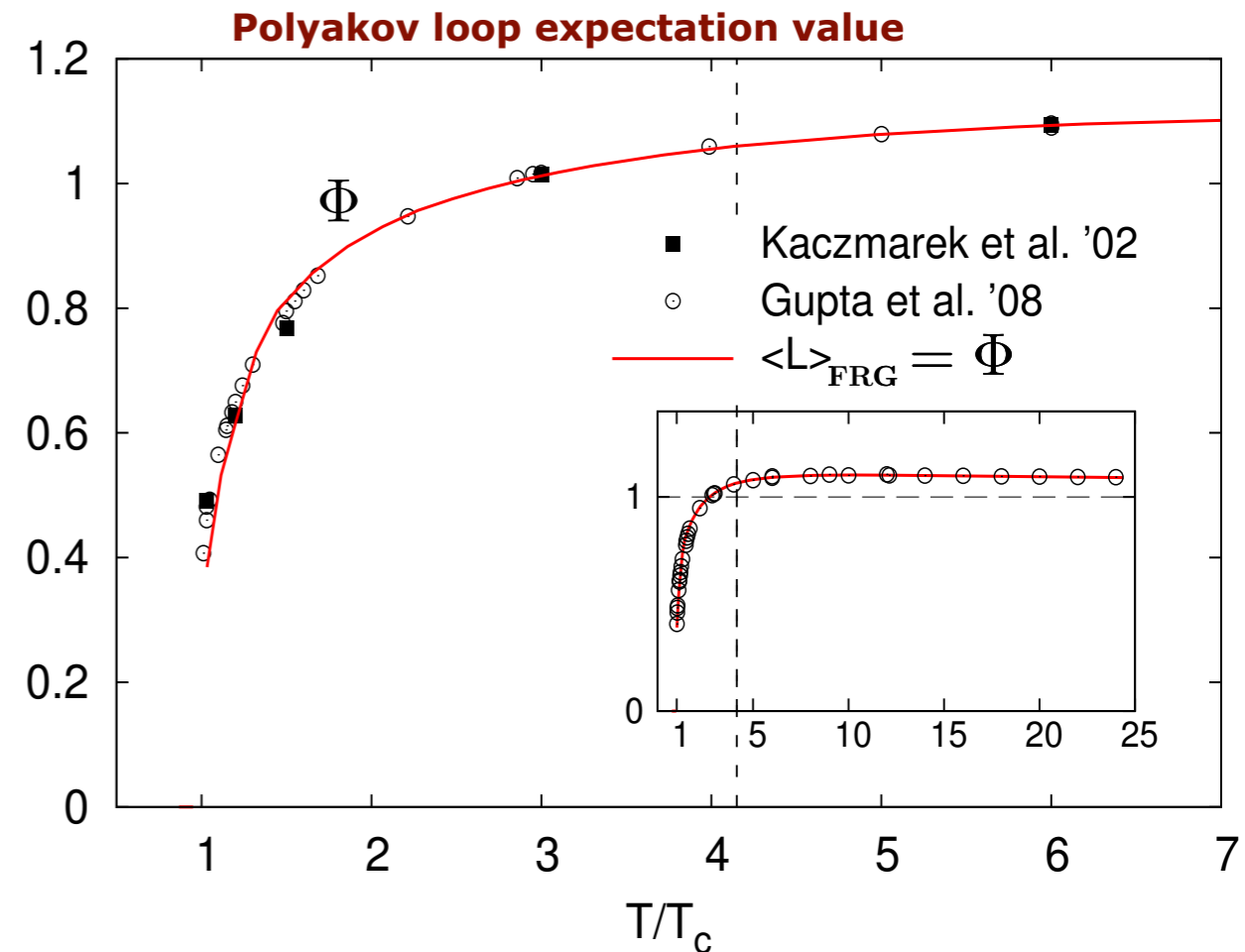
Igarashi, Itoh, Sonoda, PTP Suppl. 181 (2010) 1

Pagani, PRD 94 (2016) 045001

## Parameterisation

$$\langle L[A_0] \rangle = Z_L[\bar{A}, \phi] \cdot L[A_0]$$

with  $\phi = (a_\mu, c, \bar{c})$



# Confinement

Herbst, Luecker, JMP, arXiv:1510.03830

## Flow equation for the Polyakov loop expectation value

$$\partial_t \langle L[A_0] \rangle = - \frac{1}{2} \left[ \frac{\delta^2 \langle L[A_0] \rangle}{\delta A^2} - \frac{\delta^2 \langle L[A_0] \rangle}{\delta c \delta \bar{c}} \right]$$

## Flow equation for composite operators

JMP, AP 322 (2007) 2831

Igarashi, Itoh, Sonoda, PTP Suppl. 181 (2010) 1

Pagani, PRD 94 (2016) 045001

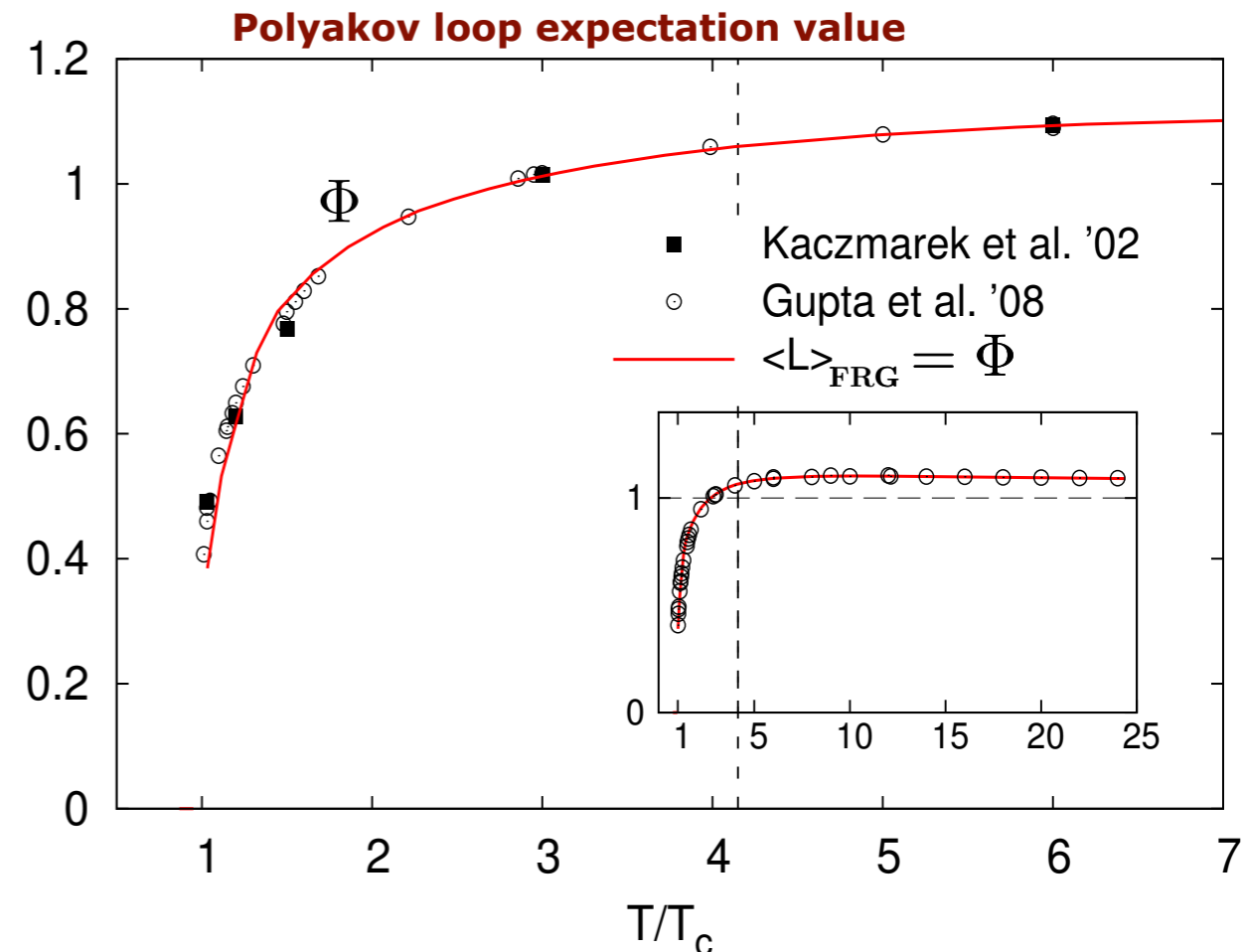
## Parameterisation

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with  $\phi = (a_\mu, c, \bar{c})$

## Flow for Polyakov loop wave function

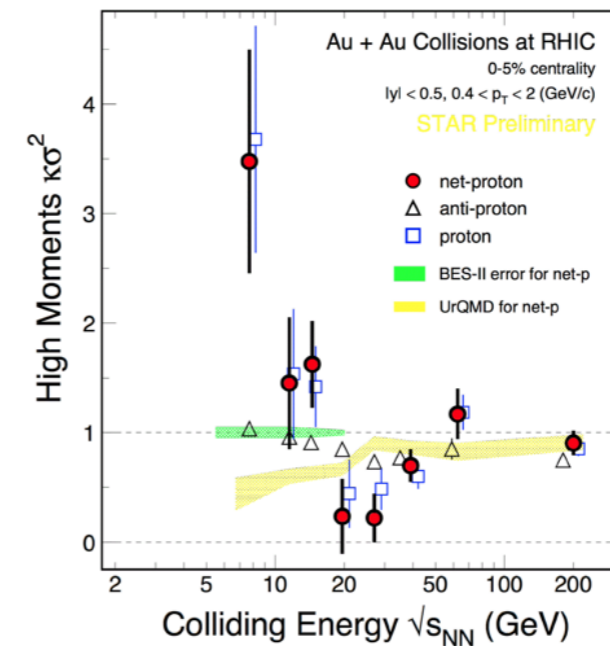
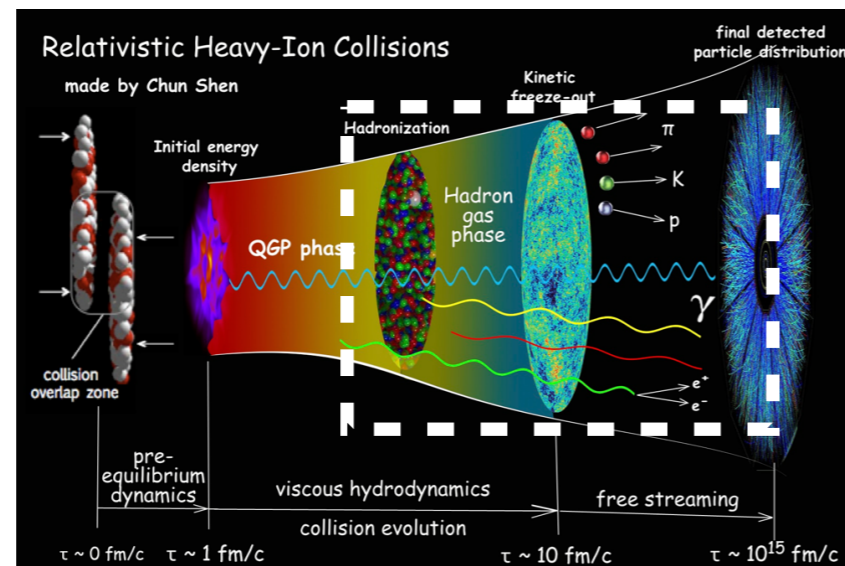
$$\partial_t Z_L[\bar{A}, \phi] = \text{Flow}_{Z_L}[\bar{A}; Z_L, G_A, G_c, L[A_0]]$$



# **Phase structure of QCD and dynamics**



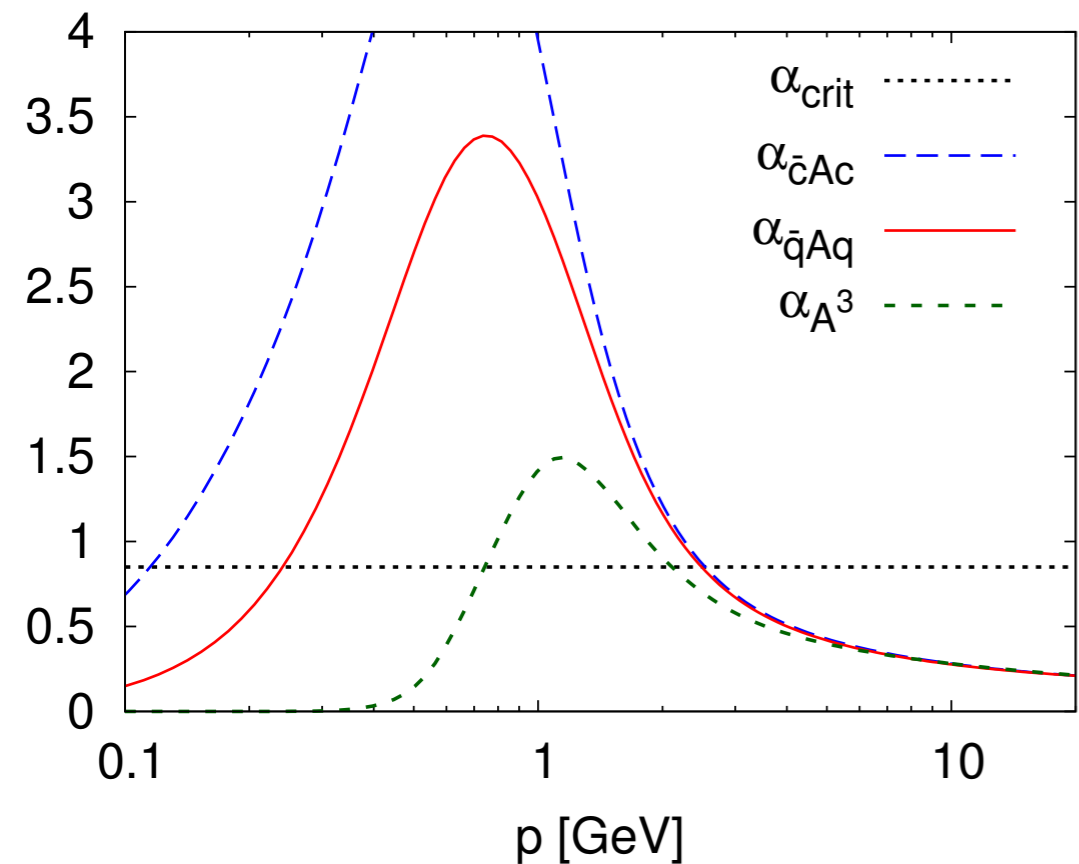
# QCD-assisted transport



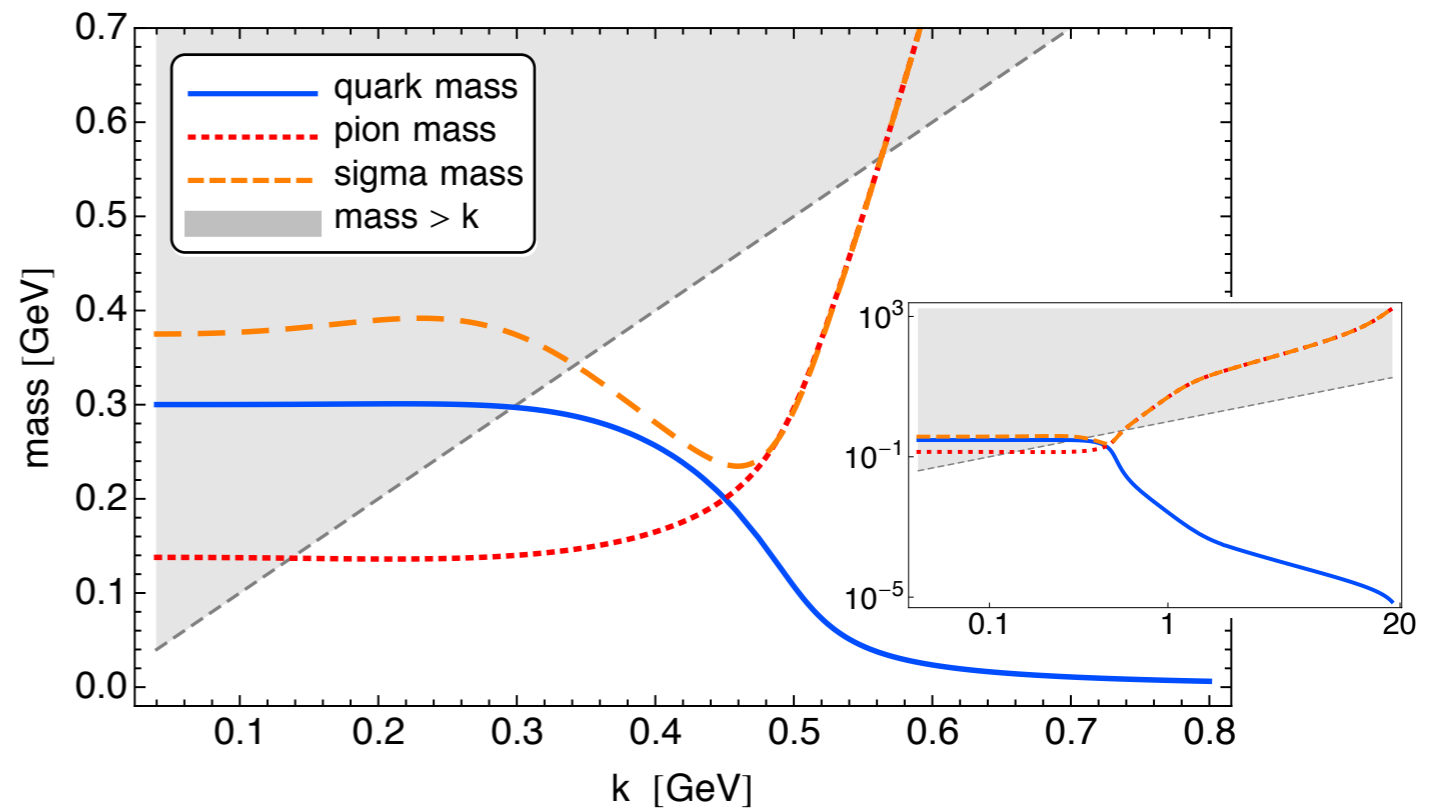
# On the unreasonable effectiveness of low energy effective theories

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \left( \text{diagram 1} - \text{diagram 2} - \text{diagram 3} + \frac{1}{2} \text{diagram 4} \right)$$

**Sequential decoupling of gluon, quark, sigma, pion fluctuations**



Mitter, JMP, Strodthoff, PRD 91 (2015) 054035



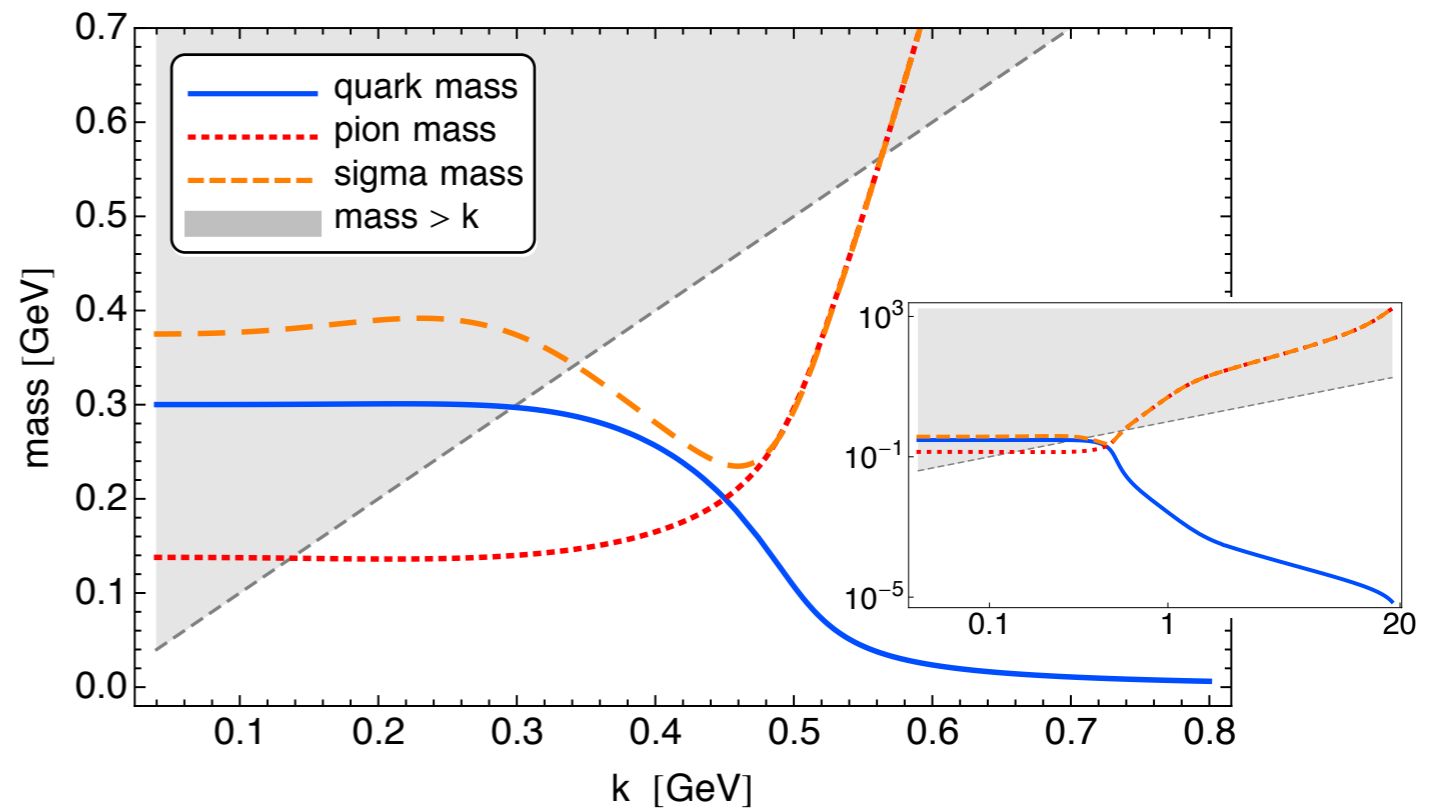
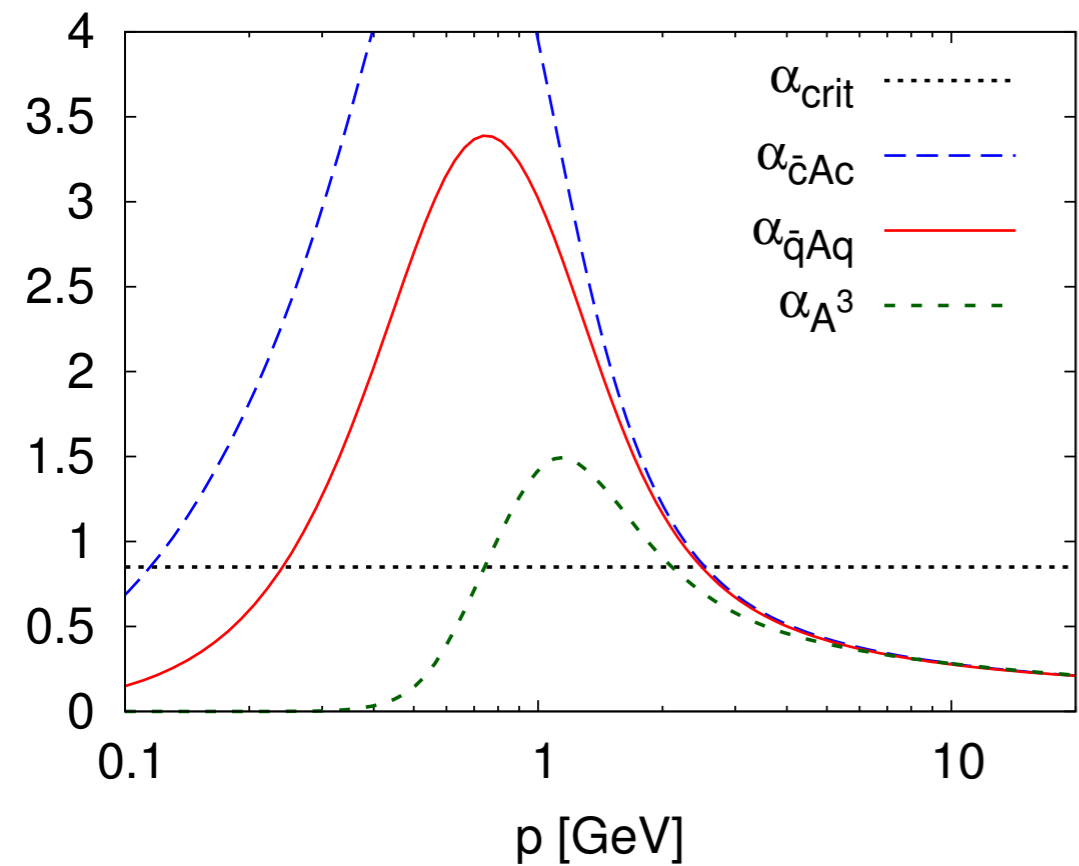
Braun, Fister, Haas, JMP, Rennecke, PRD 94 (2016) 034016

Rennecke, PRD 92 (2015) 076012

# On the unreasonable effectiveness of low energy effective theories

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{[diagram 1]} - \text{[diagram 2]} - \text{[diagram 3]} + \frac{1}{2} \text{[diagram 4]}$$

**Sequential decoupling of gluon, quark, sigma, pion fluctuations**



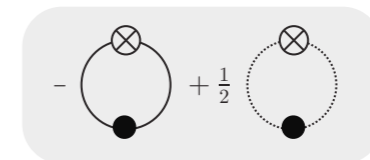
**PQM-model**



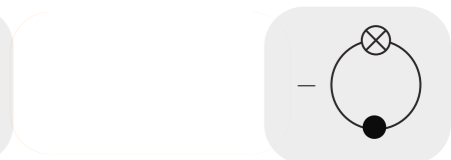
**PNJL-model**



**QM-model**

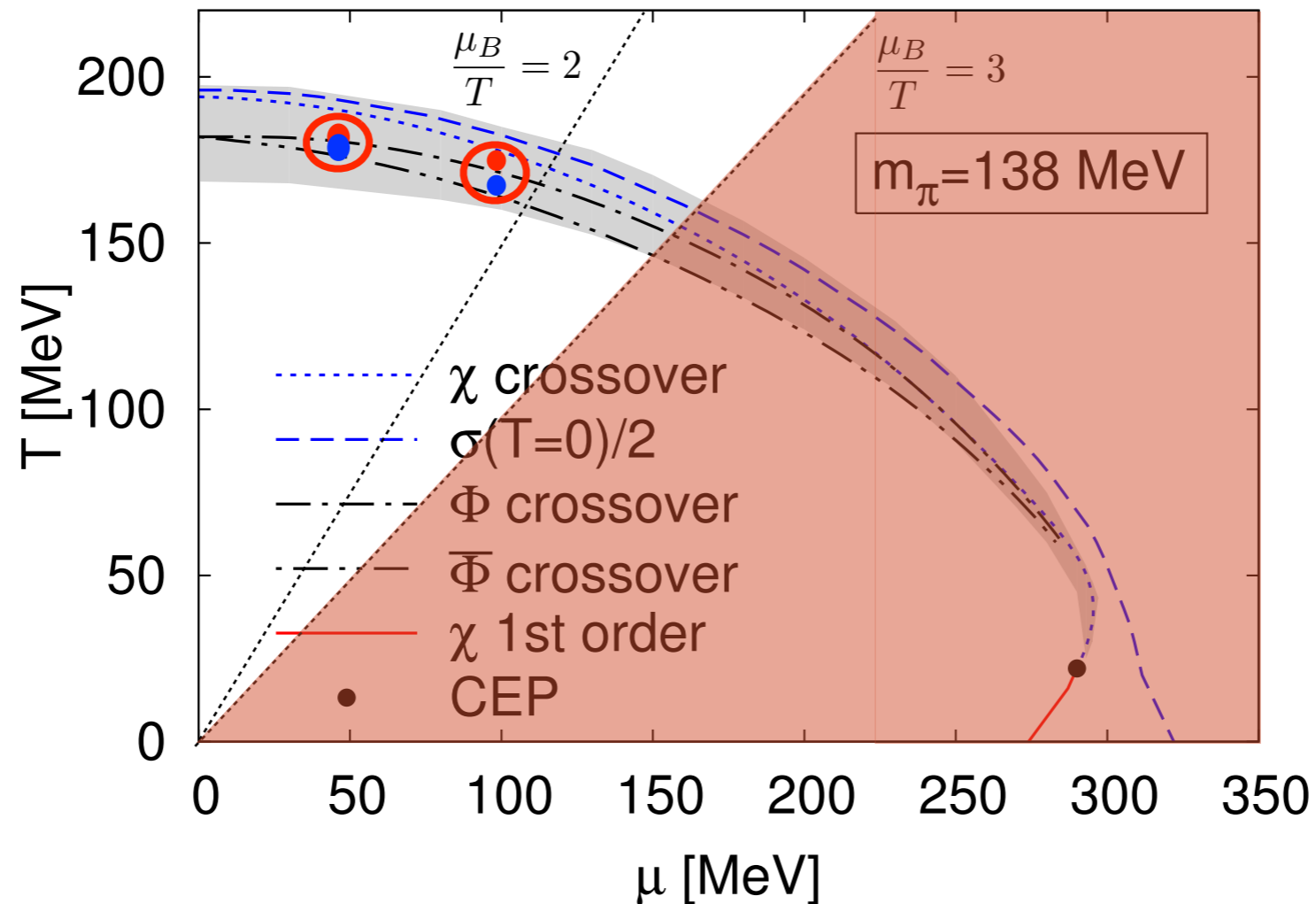


**NJL-model**



# QCD at finite density

## Phase diagram of the QCD-enhanced PQM



Herbst, JMP, Schaefer, PLB 696 (2011) 58-67  
PRD 88 (2013) 1, 014007



FRG QCD results at finite density

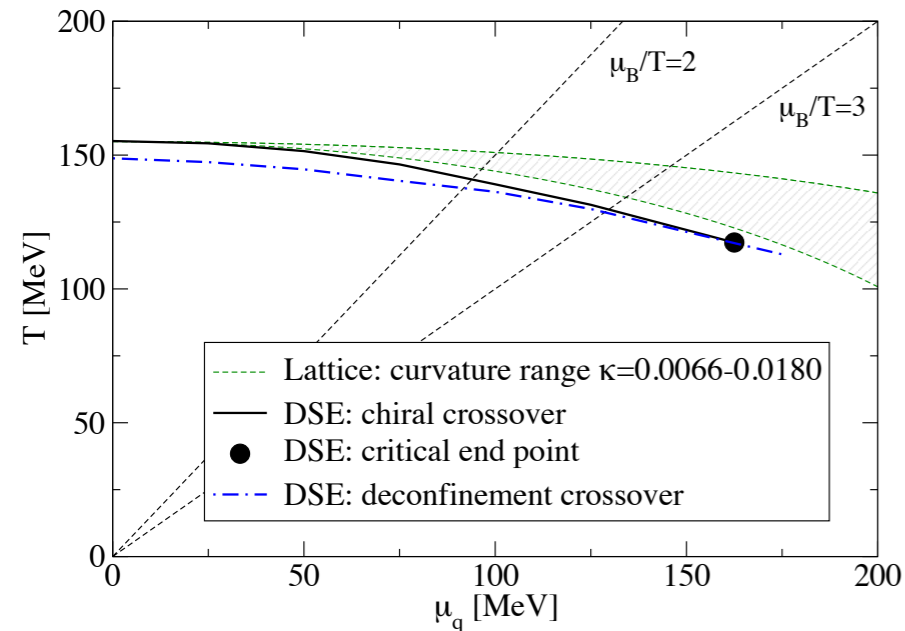
Haas, Braun, JMP '09, unpublished

Extension of FRG QCD results at imaginary chemical potential

Braun, Haas, Marhauser, JMP, PRL 106 (2011) 022002

# Phase structure at finite density

Phase diagram of 2+1 flavor QCD



Kaczmarek et al. '11  
 Endrodi, Fodor, Katz, Szabo '11  
 Cea, Cosmai, Papa '14

Fischer, Fister, Luecker, JMP, PLB732 (2014)

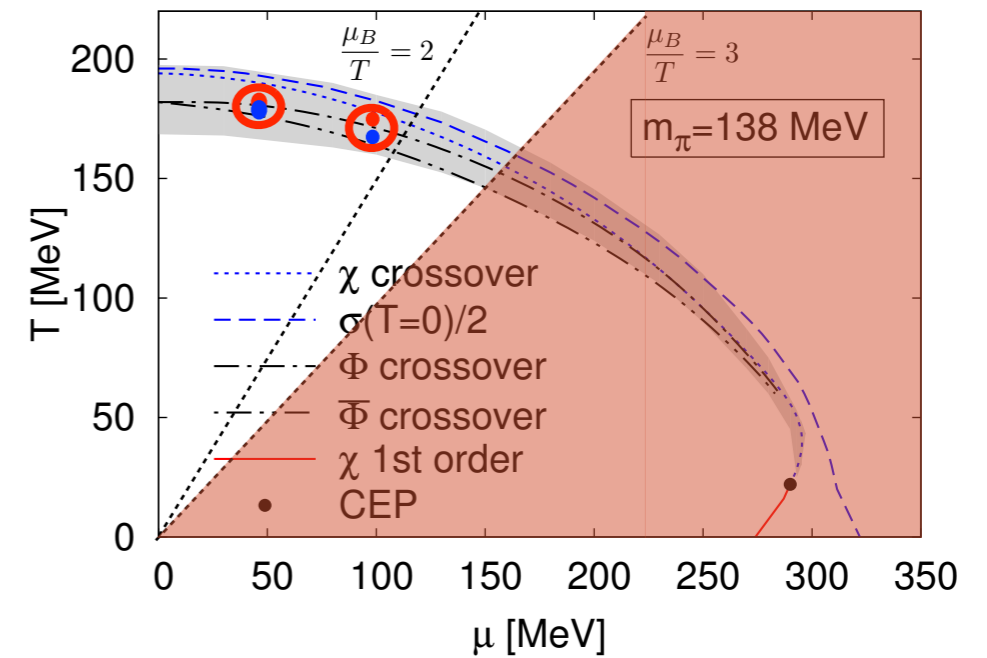
Fischer, Luecker, Welzbacher, PRD 90 (2014) 034022

Eichmann, Fischer, Welzbacher, PRD 93 (2014) 034013

## Chiral phase structure

Qin, Chang, Chen, Liu, Roberts, PRL 106 (2011) 172301

Phase diagram of QCD-enhanced 2-flavor PQM-model



Herbst, JMP, Schaefer, PLB 696 (2011) 58-67  
 PRD 88 (2013) 1, 014007

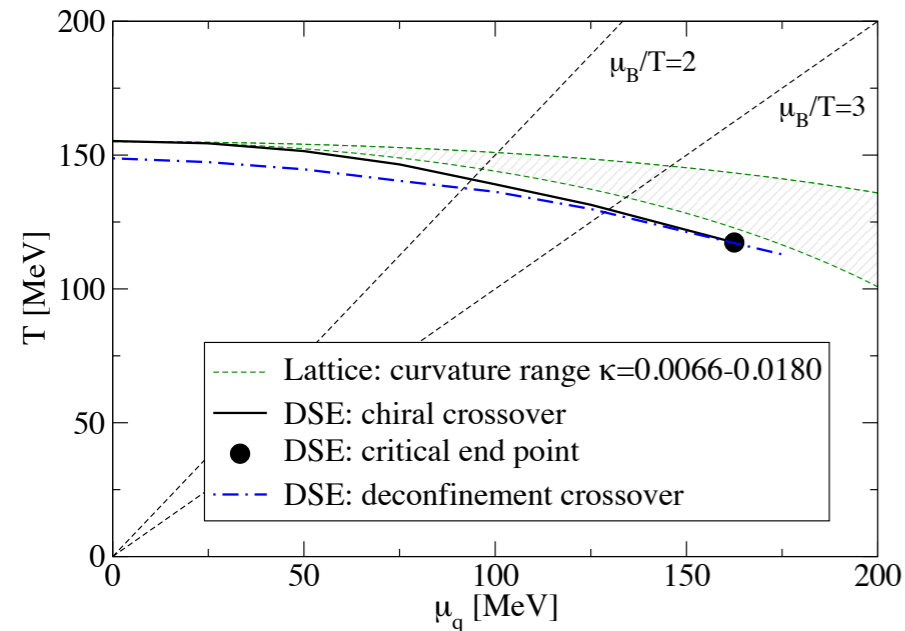


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Fischer, Fister, Luecker, JMP, PLB732 (2014)

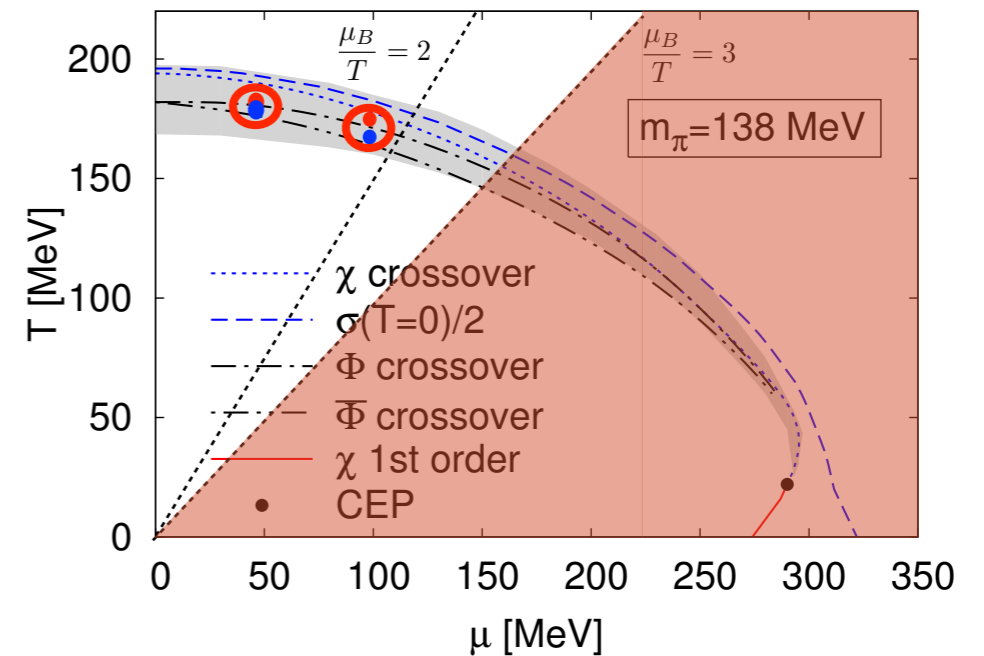
Fischer, Luecker, Welzbacher, PRD 90 (2014) 034022

Eichmann, Fischer, Welzbacher, PRD 93 (2014) 034013

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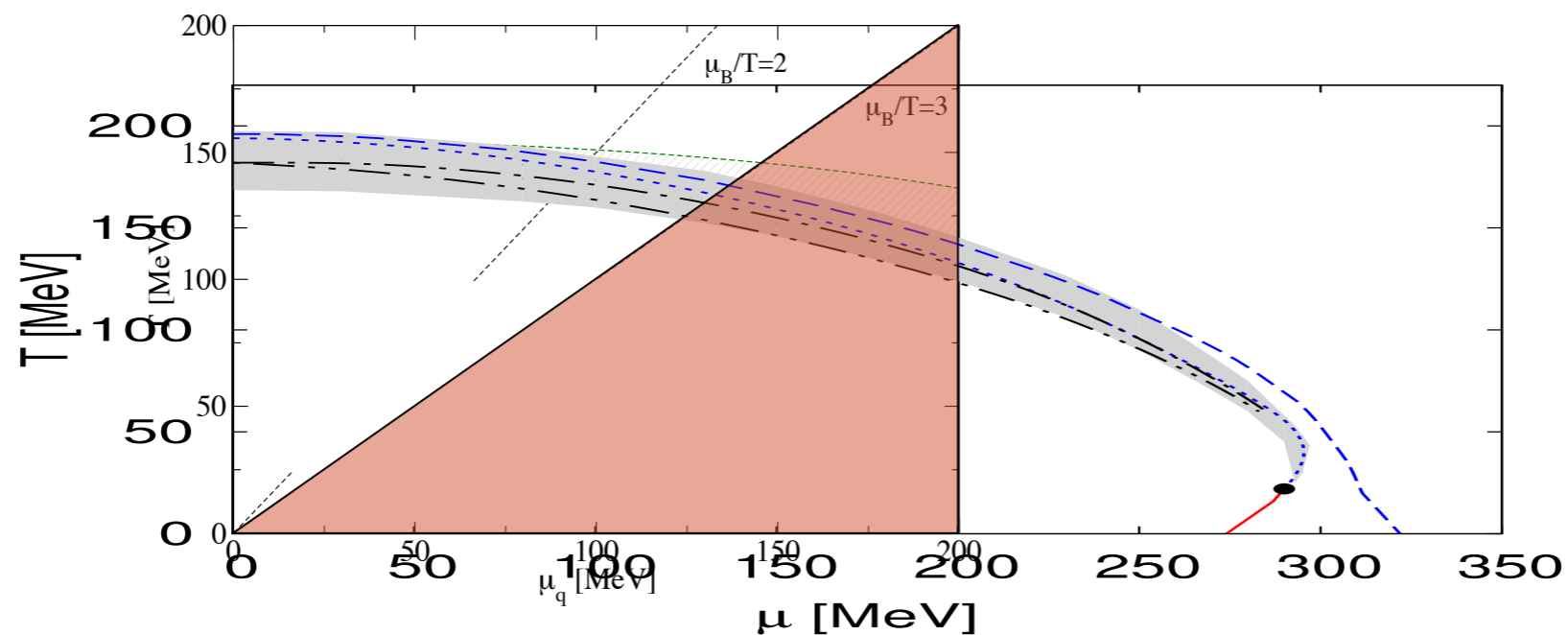
Herbst, JMP, Schaefer, PLB 696 (2011) 58-67  
PRD 88 (2013) 1, 014007



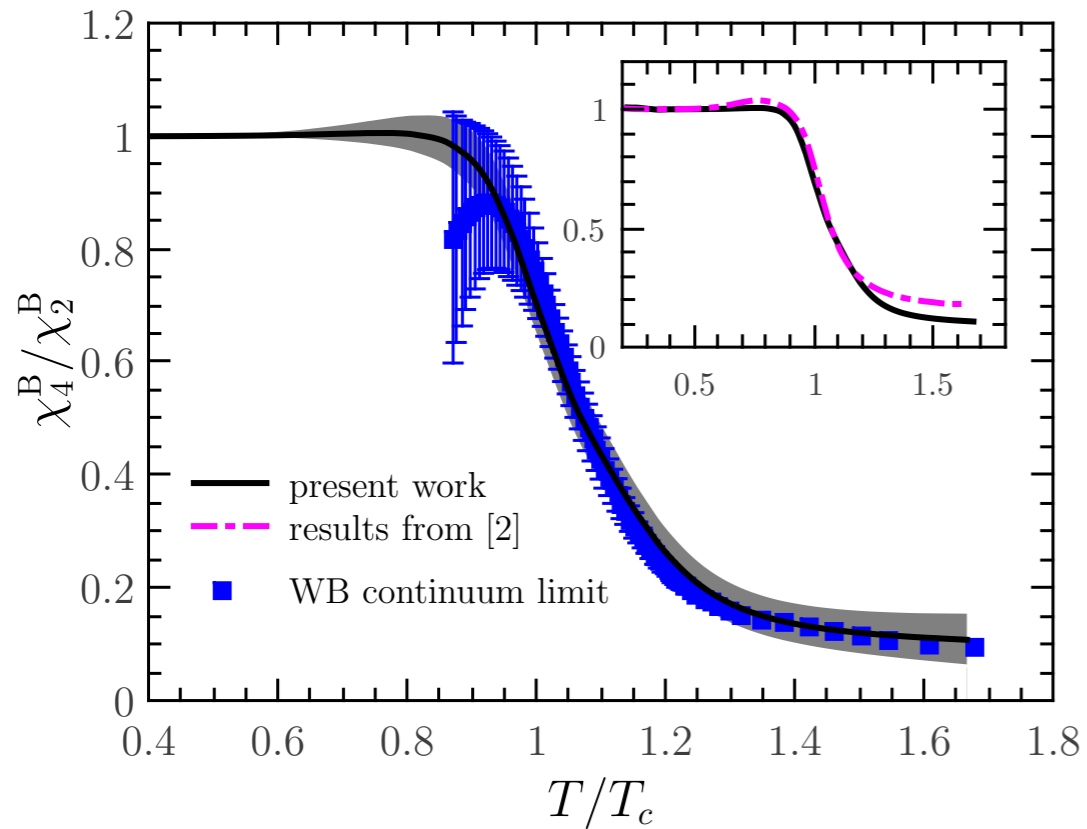
FRG QCD results at finite density

Haas, Braun, JMP '09, unpublished

Comparison with 2 flavor vs 2+1 flavor scale matching of  $T_c$



# Fluctuations as a measure of confinement



[2] Fu, JMP, PRD 92 (2015) 116006

Karsch, Schaefer, Wagner, Wambach, PLB 698 (2011) 256

Friman, Karsch, Redlich, Skokov, EPJ C71 (2011) 1694

Schaefer, Wagner, PRD 85 (2012) 034027

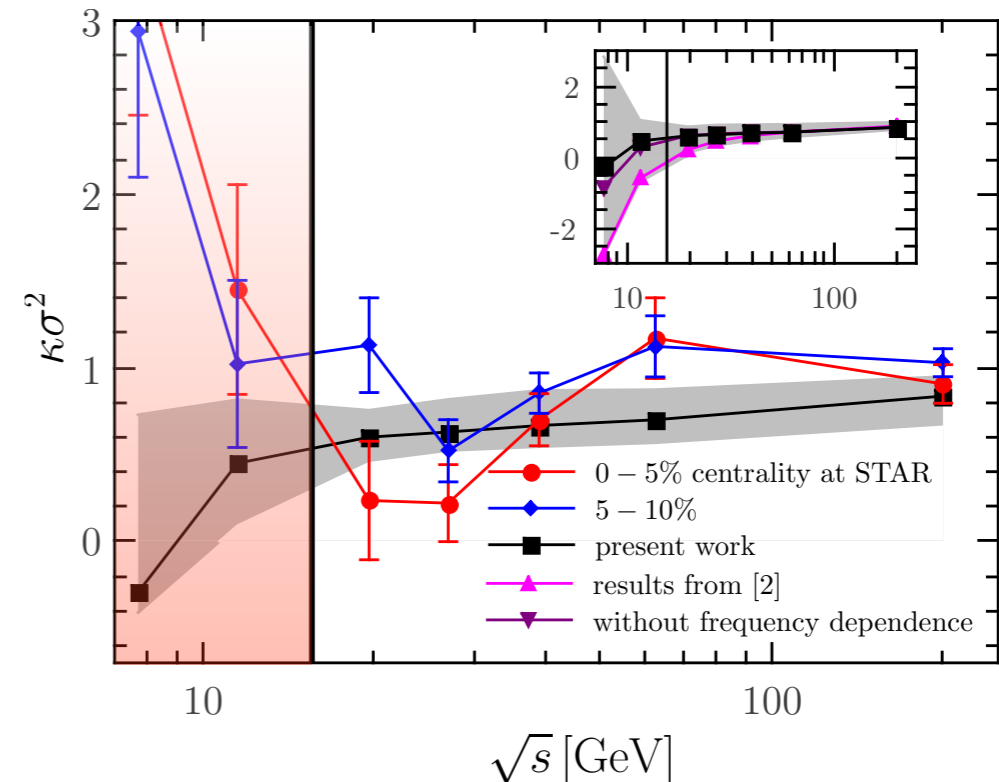
Skokov, Friman, Redlich, PRC 88 (2013) 034911

Almasi, Friman, Redlich, Nucl.Phys. A956 (2016) 356-359

$$\chi_n^B = \frac{\partial^n}{\partial(\mu_B/T)^n} \frac{p}{T^4}$$

Skewness, Kurtosis

$$\sigma^2 = VT^3 \chi_2^B \quad S = \frac{\chi_3^B}{\chi_2^B \sigma} \quad \kappa = \frac{\chi_4^B}{\chi_2^B \sigma^2}$$



[2] Fu, JMP, PRD 93 (2016) 091501

Fu, JMP, Schaefer, Rennecke, PRD 94 (2016) 116020

# Transport approach to QCD

Blum, Jiang, Mitter, Nahrgang, JMP, Rennecke, Wink

**Time evolution of the critical (scalar)  $\sigma$  mode**

$$\frac{\delta\Gamma}{\delta\sigma} = \xi$$

**quantum equation of motion**

**noise field**

**Extension of mean-field version**

**Nahrgang, Leupold, Herold, Bleicher PRC84 (2011)**

**see also**

**Stephanov, Rajagopal, Shuryak PRL81 (1998)**

**Mukherjee, Venugopalan, Yin PRC92 (2015)**

**Herold, Nahrgang, Yan, Kobdaj PRC93 (2016)**

**Nahrgang, Bluhm, Schäfer, Bass arXiv:1804.05728**



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**noise field**

**Input from equilibrium low energy effective action of QCD**

$$\text{Re } \Gamma_{\sigma}^{(2)}(\omega, \vec{p})$$

**kinetic term**

$$\text{Im } \Gamma_{\sigma}^{(2)}(\omega, \vec{p})$$

**diffusion term  $\eta \partial_t \sigma$**

$$U(\sigma)$$

**effective potential**

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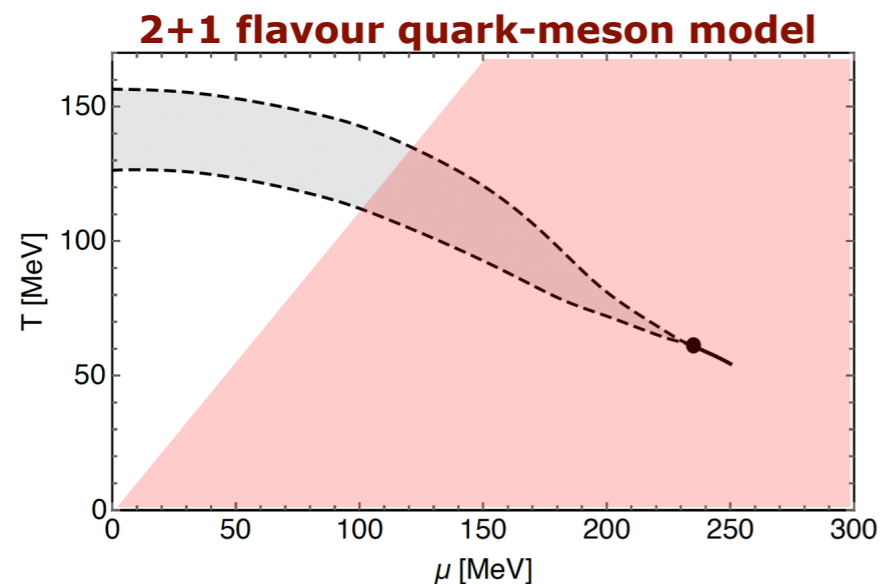
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effective potential

Phase structure of low energy QCD



Schaefer, Rennecke, PRD 96 (2017) 1, 016009

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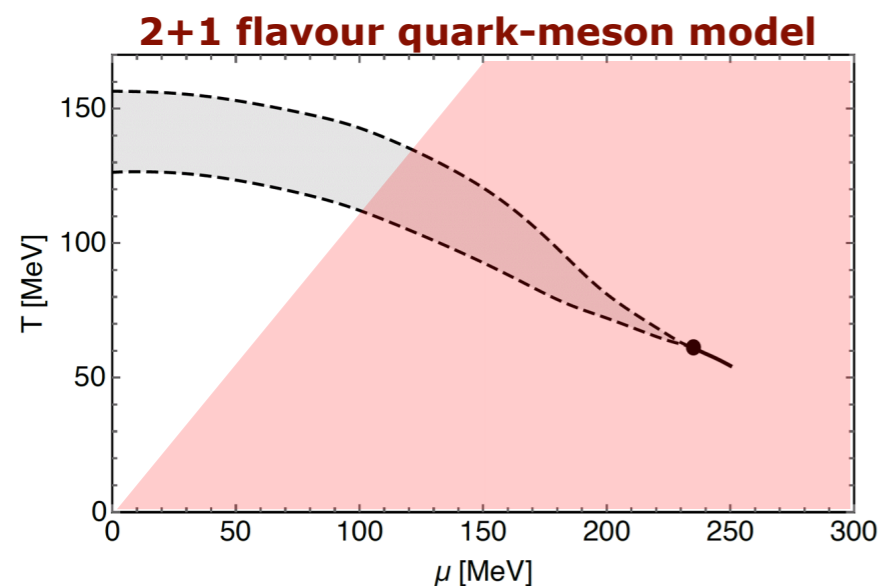
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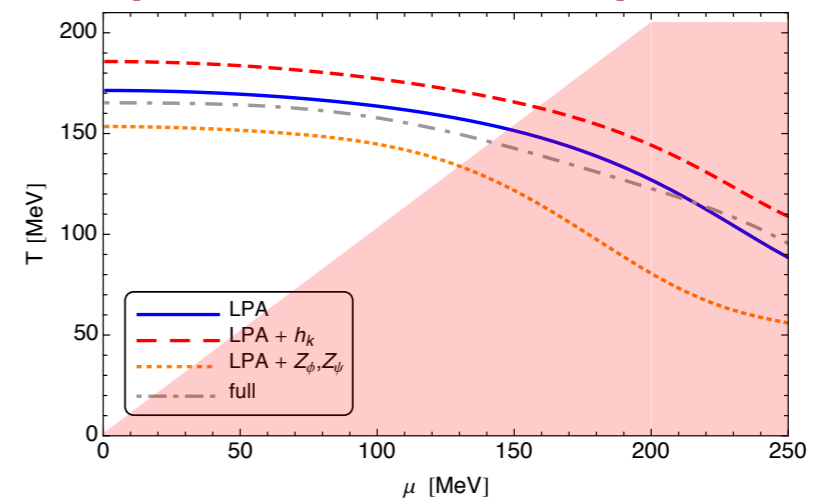
effective potential

Phase structure of low energy QCD



Schaefer, Rennecke, PRD 96 (2017) 1, 016009

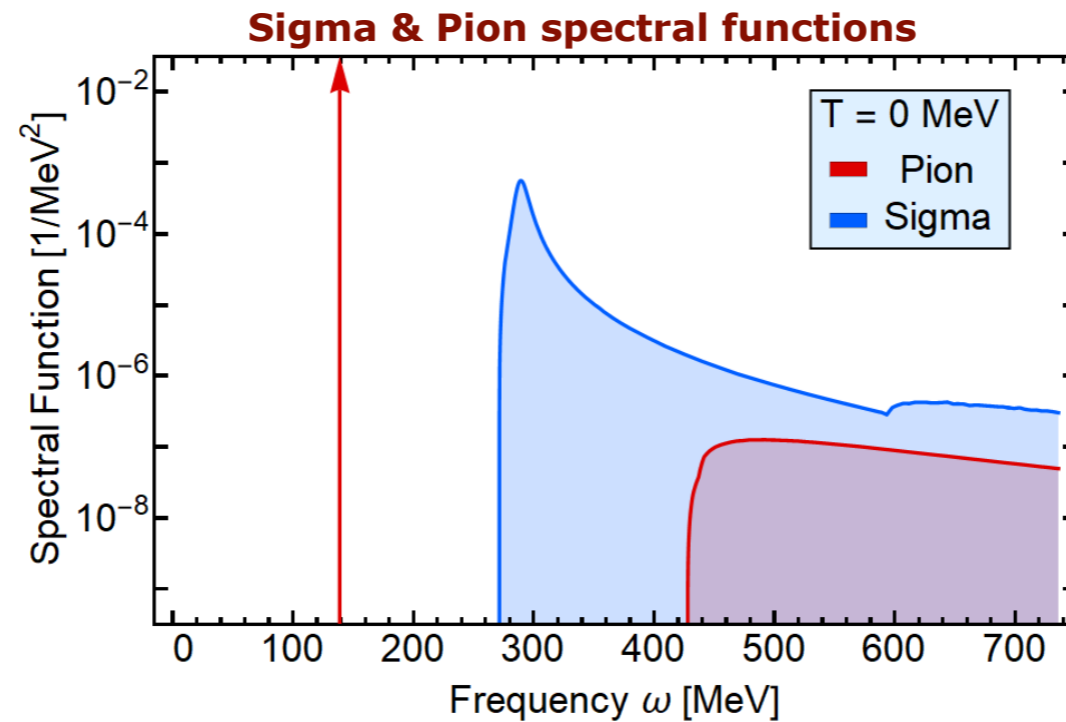
Comparison of truncations (2 flavours)



JMP, Rennecke, PRD 90 (2014) 7, 076002

# Pion & sigma spectral functions

Show case in linear sigma model



**JMP, Strodthoff, Wink, arXiv:1711.07444**

**Real-time FRG computations, e.g.**

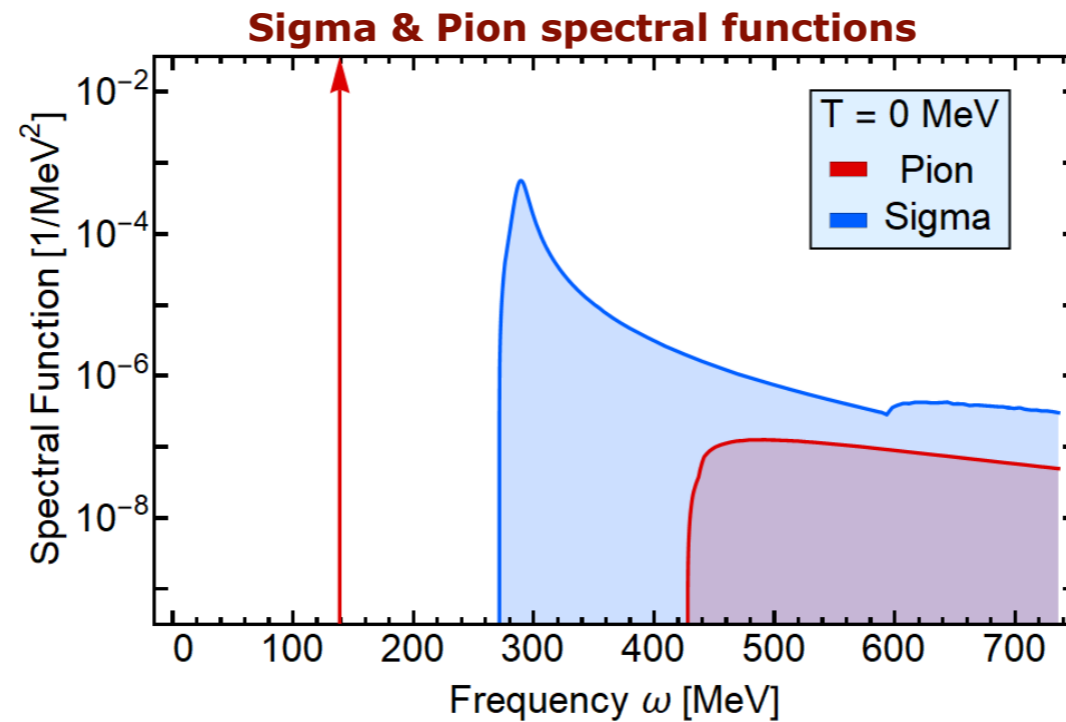
**Flörchinger JHEP 1205 (2012) 021**

**Kamikado, Strodthoff, von Smekal, Wambach, EPJC 74 (2014) 2806**

**JMP, Strodthoff, PRD 92 (2015) 094009**

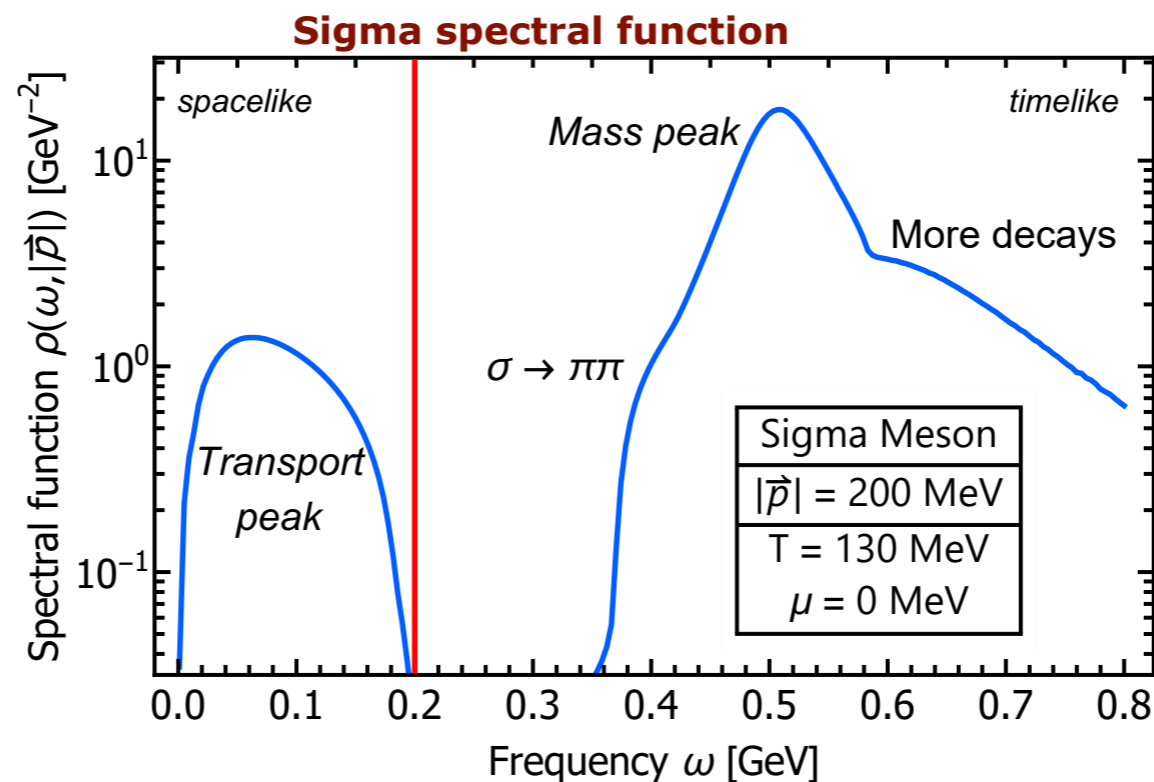
# Pion & sigma spectral functions

Show case in linear sigma model



JMP, Strodthoff, Wink, arXiv:1711.07444

2+1 flavour quark-meson model sigma spectral function

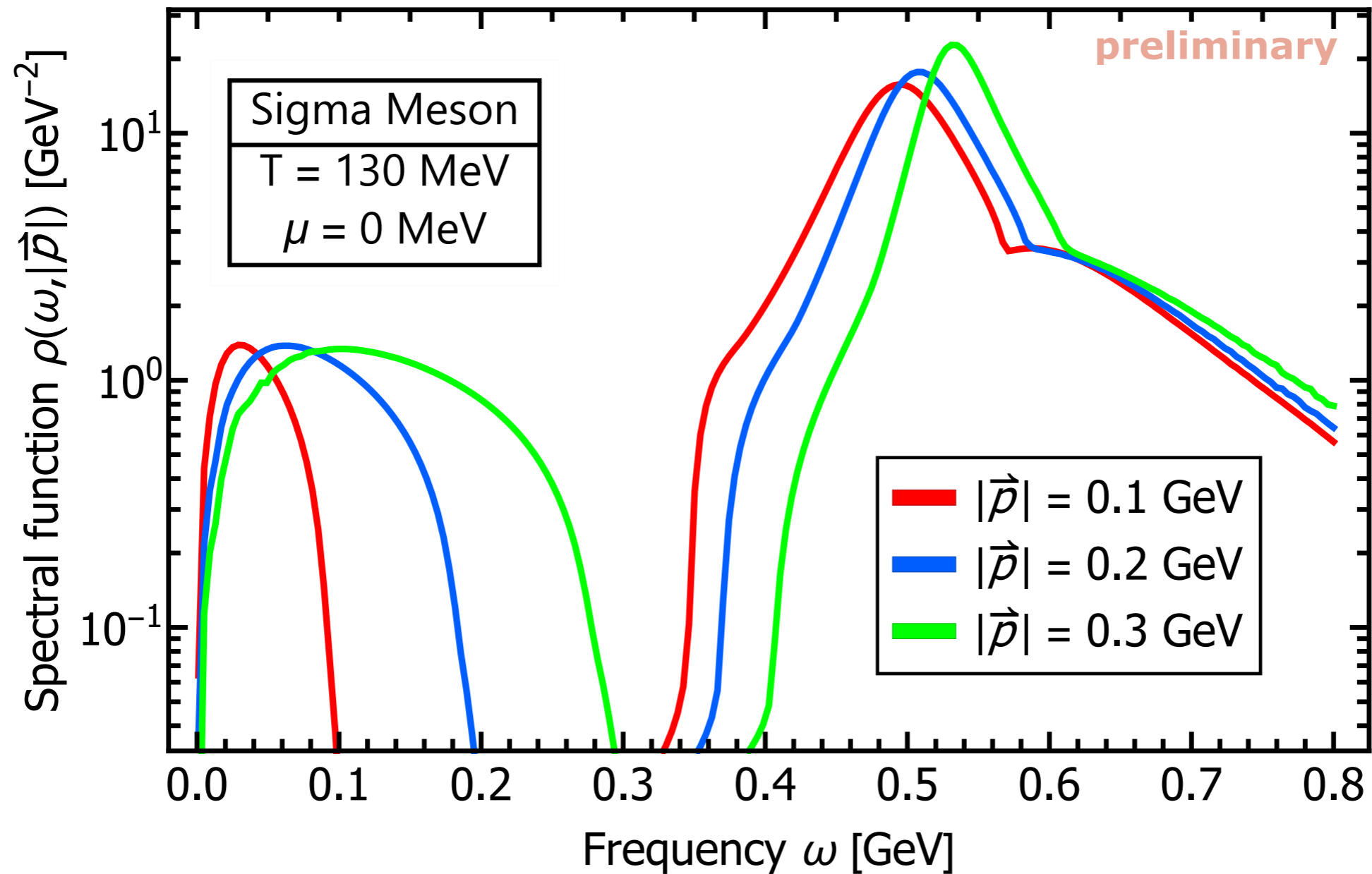


preliminary

JMP, Rennecke, Wink, in prep

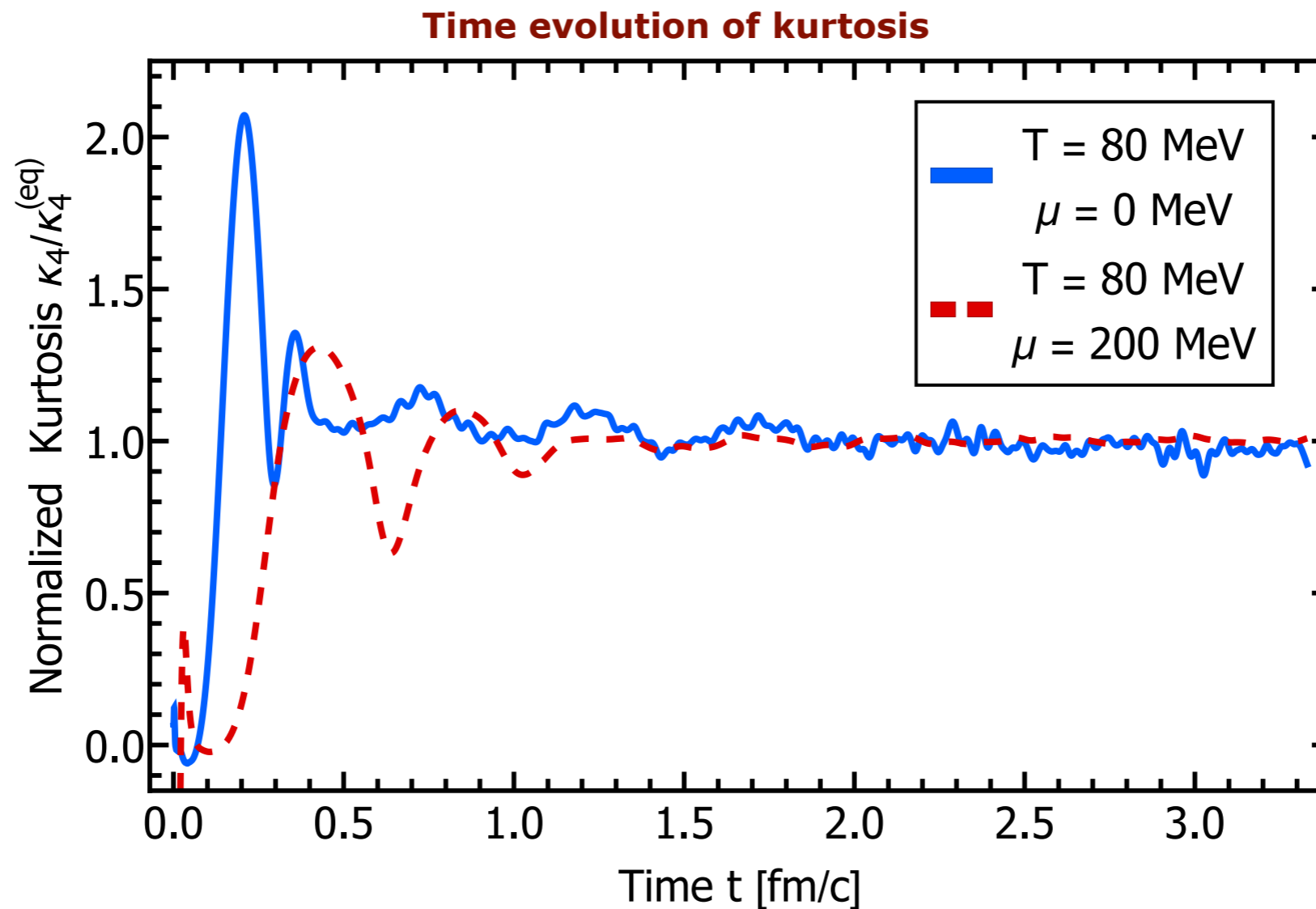
# Pion & sigma spectral functions

2+1 flavour quark-meson model sigma spectral function



# Time evolution of cumulants

Blum, Jiang, Nahrgang, JMP, Rennecke, Wink, in prep



**$n$ th central moment of the sigma field:  $\chi_n$**

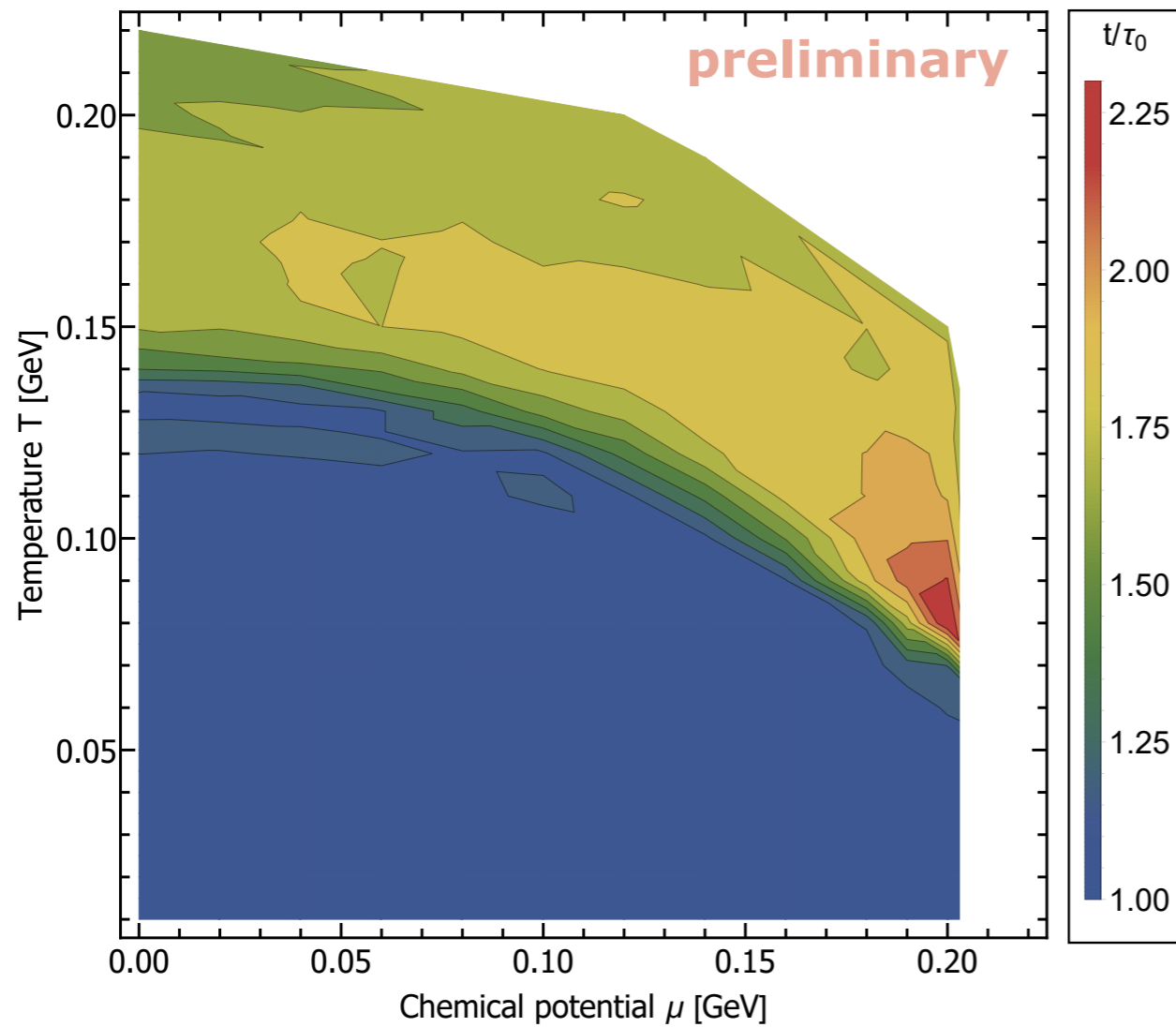
$$\chi_2 = \langle (\sigma - \langle \sigma \rangle)^2 \rangle$$

**kurtosis:**  $\kappa = \frac{\chi_4}{\chi_2^2} - 3$

# Equilibration time phase structure

Blum, Jiang, Nahrgang, JMP, Rennecke, Wink, in prep

## Equilibration time of sigma-kurtosis



**n**th central moment of the sigma field:  $\chi_n$

**variance:**  $\chi_2 = \langle (\sigma - \langle \sigma \rangle)^2 \rangle$

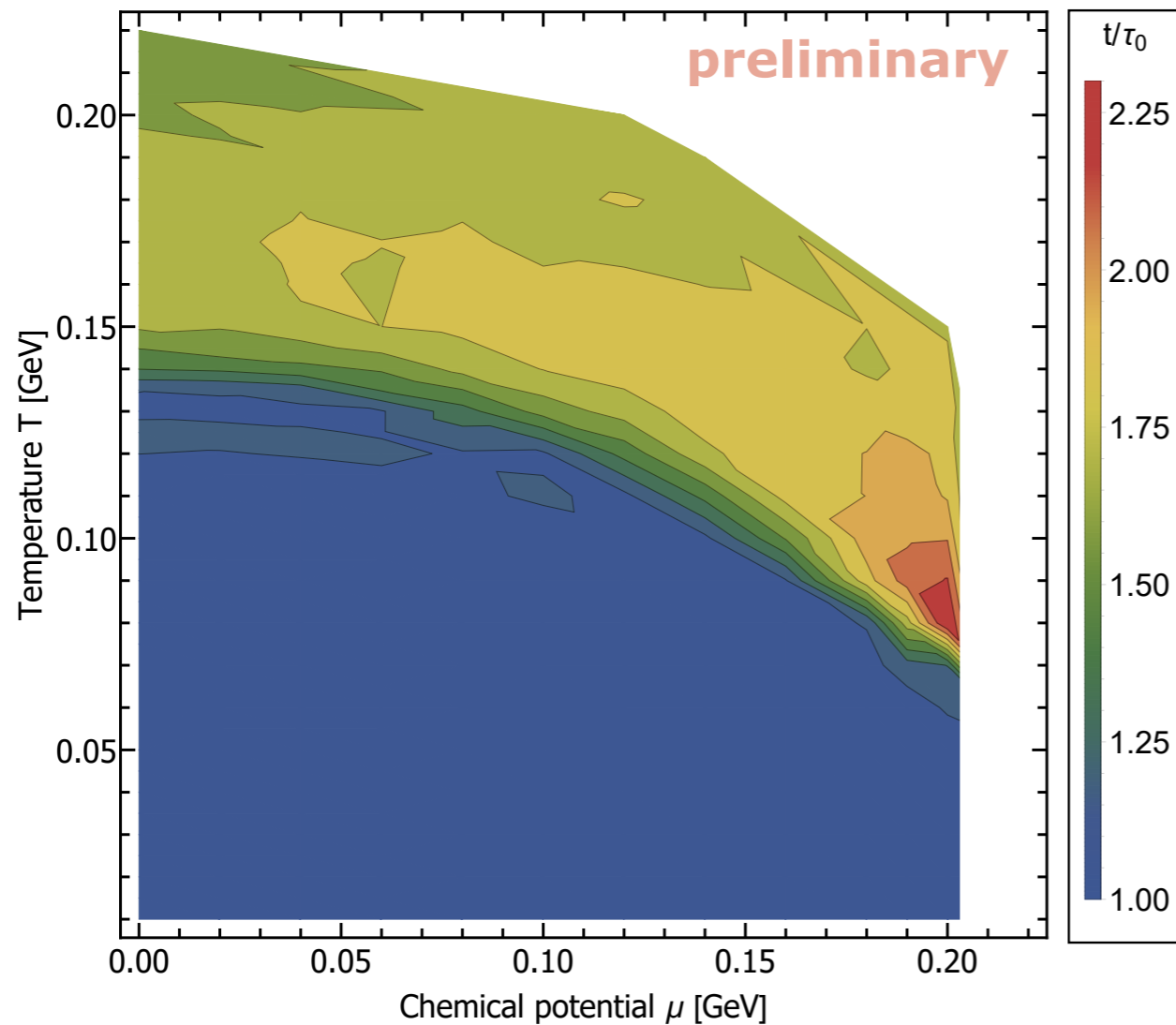
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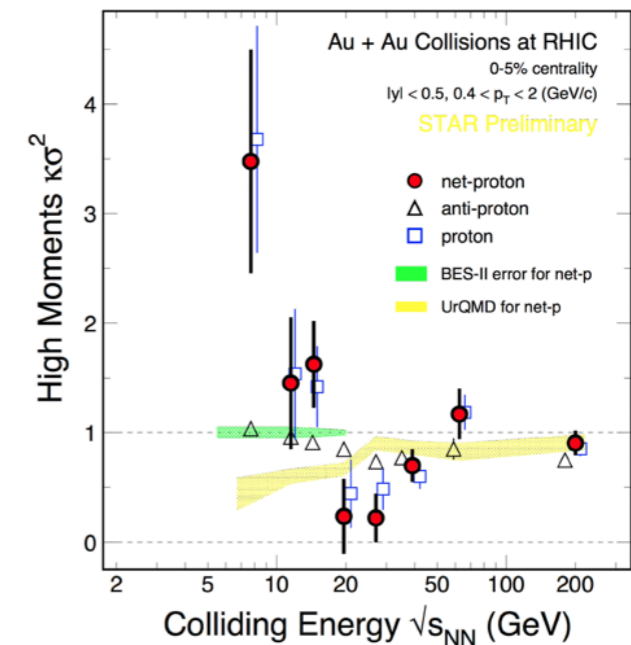
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Blum, Jiang, Nahrgang, JMP, Rennecke, Wink, in prep

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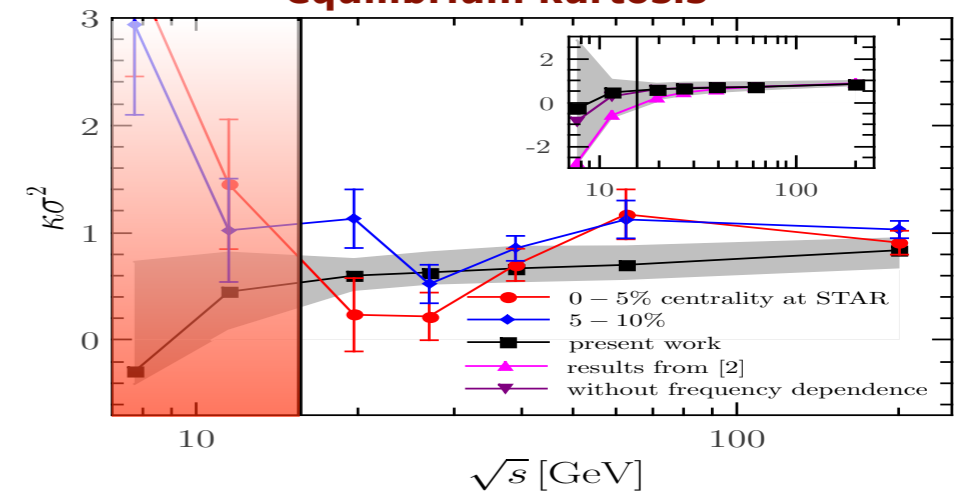


## kurtosis of baryon number fluctuations



Luo, Cu, NST 28 (2017)

## equilibrium kurtosis



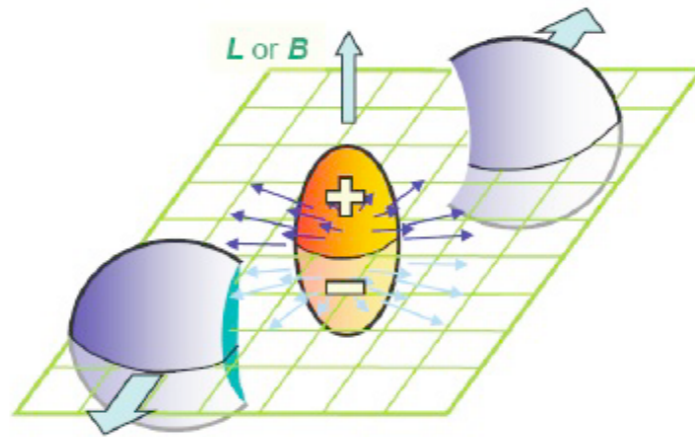
Fu, JMP, Schaefer, Rennecke, PRD 94 (2016) 11, 116020

$n$ th central moment of the sigma field:  $\chi_n$

variance:  $\chi_2 = \langle (\sigma - \langle \sigma \rangle)^2 \rangle$

kurtosis:  $\kappa = \frac{\chi_4}{\chi_2^2} - 3$

# QCD-assisted hydrodynamics



# QCD-assisted hydrodynamics

Dubla, Masciocchi, JMP, Schenke, Shen, Stachel, arXiv:1805.02985

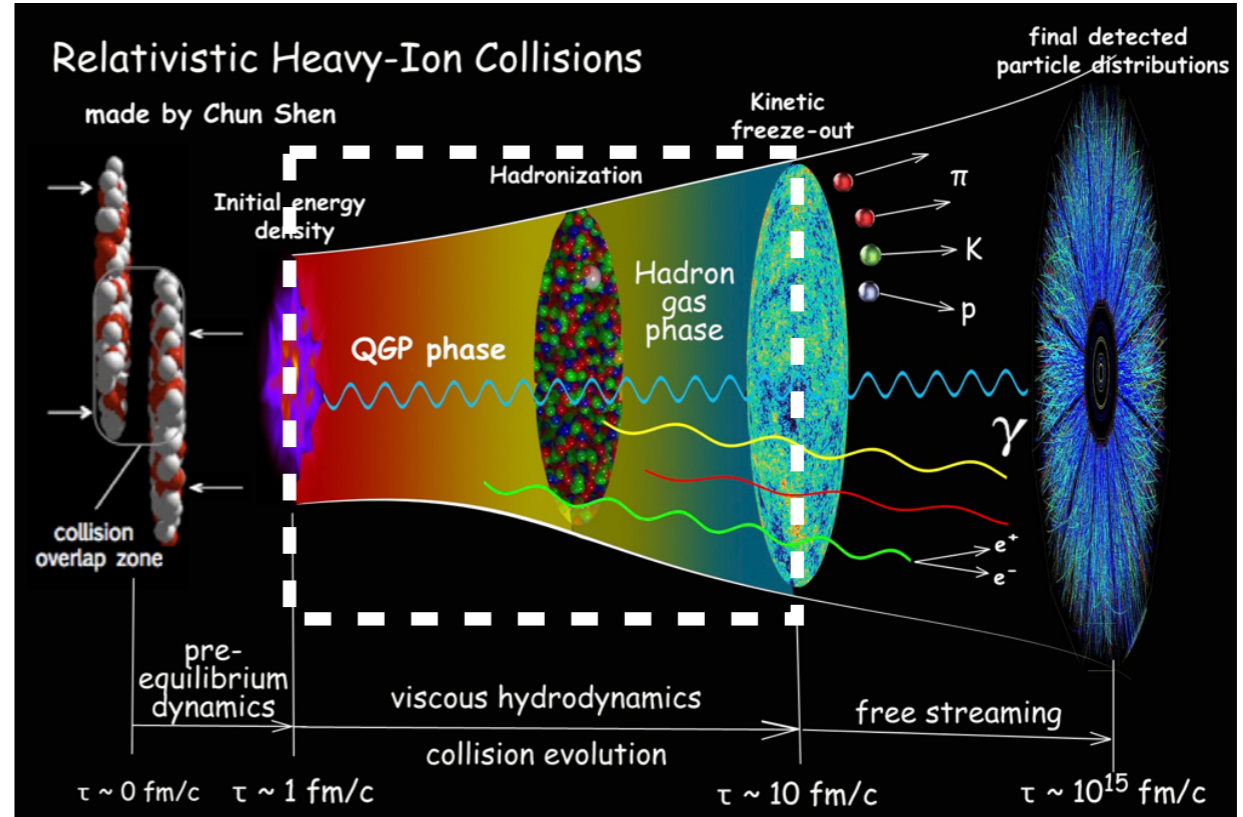
## HIC 'phases'

Far from equilibrium initial phase

Kinetic phase

Hydrodynamical phase

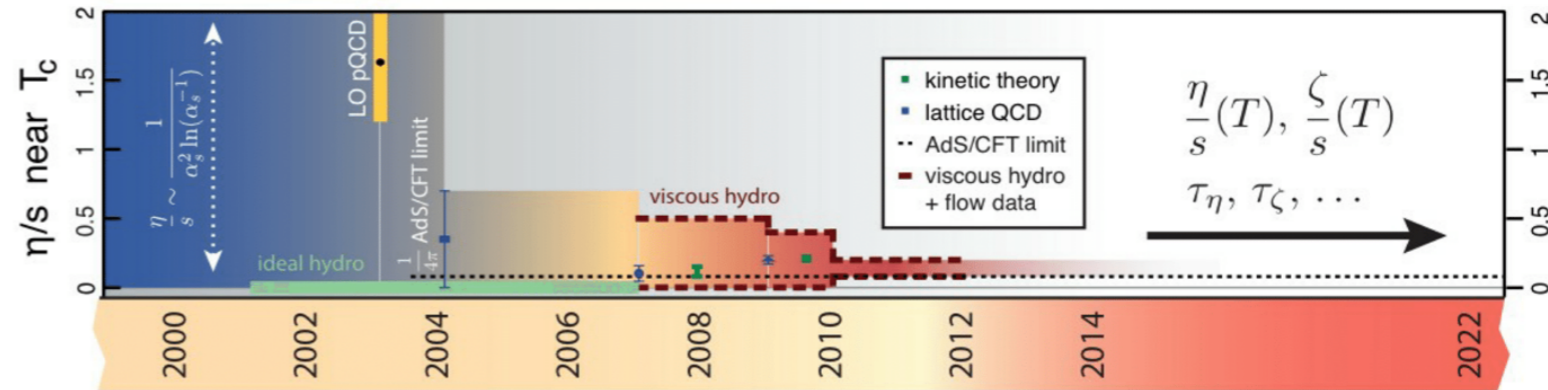
Hadronisation & freeze out



## QCD-assisted transport

Hydro with QCD transport coefficients

Equilibrium transport coefficients



'Steady-state' hydro

Constraints for the other phases

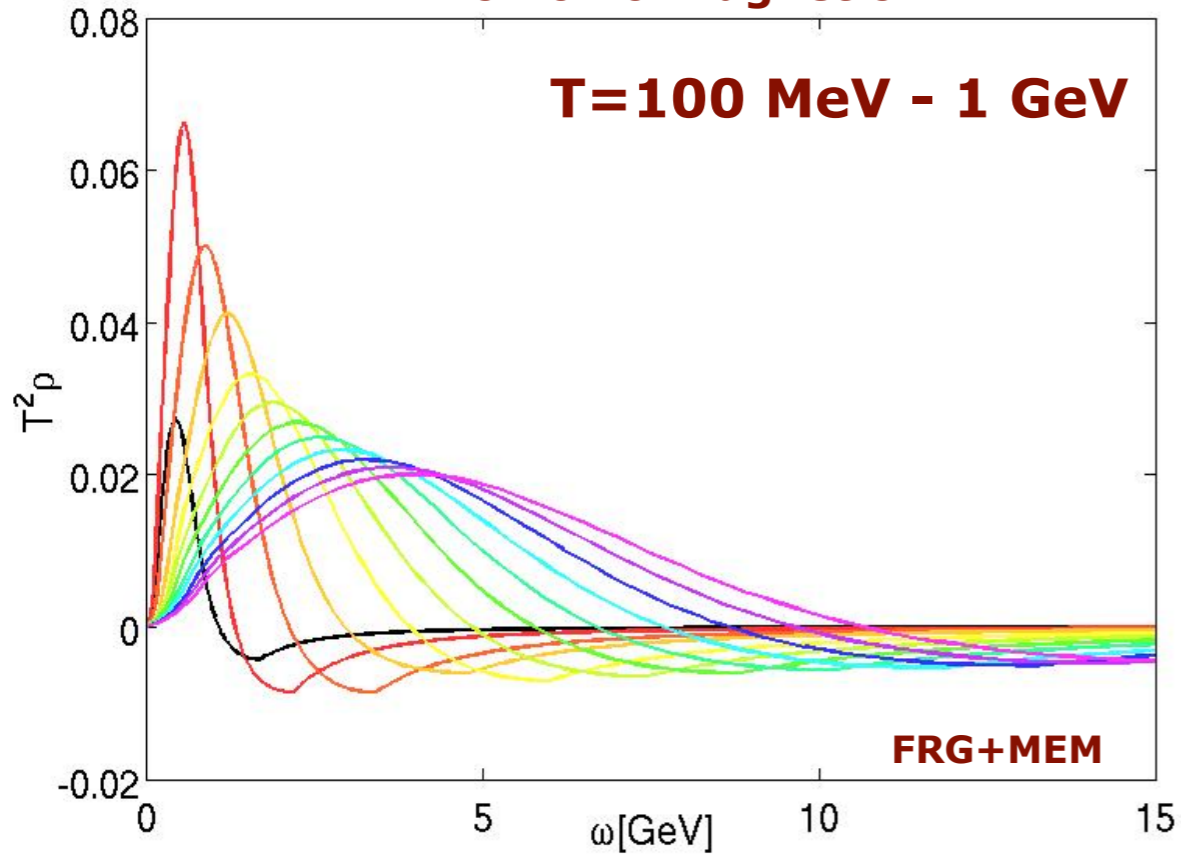
$$\pi^{\mu\nu} = \eta \left( \nabla^\mu u^\nu + \nabla^\nu u^\mu - \frac{2}{3} \Delta^{\mu\nu} \nabla_\alpha u^\alpha \right) - \frac{4}{3} \tau_\pi \pi^{\mu\nu} \partial_\alpha u^\alpha - \tau_\pi \Delta_\alpha^\mu \Delta_\beta^\nu u^\sigma \partial_\sigma \pi^{\alpha\beta},$$

# Single particle spectral functions

$$\rho(p) = 2 \operatorname{Im} \langle A A \rangle_{\text{ret}}(p)$$

# Single particle spectral functions

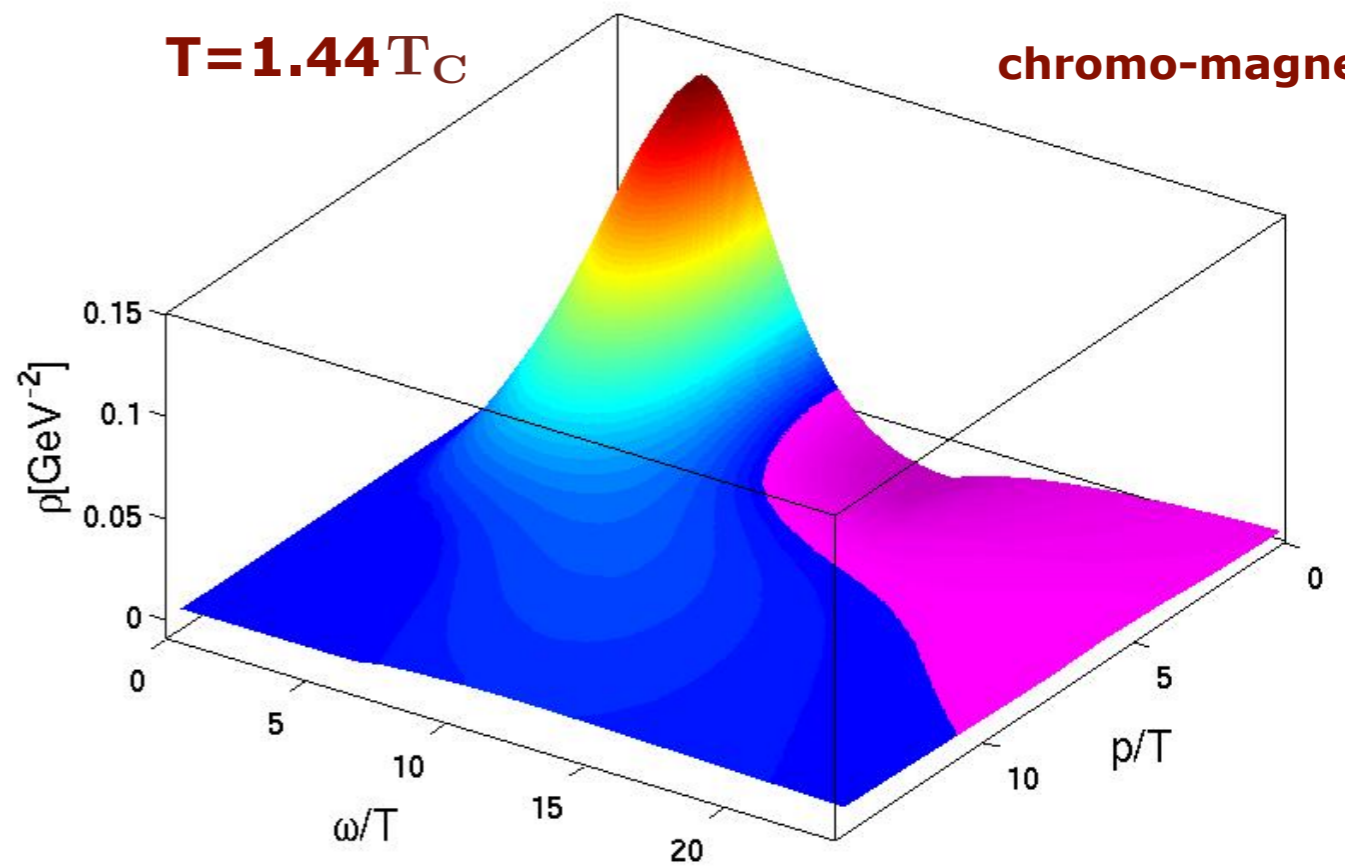
chromo-magnetic



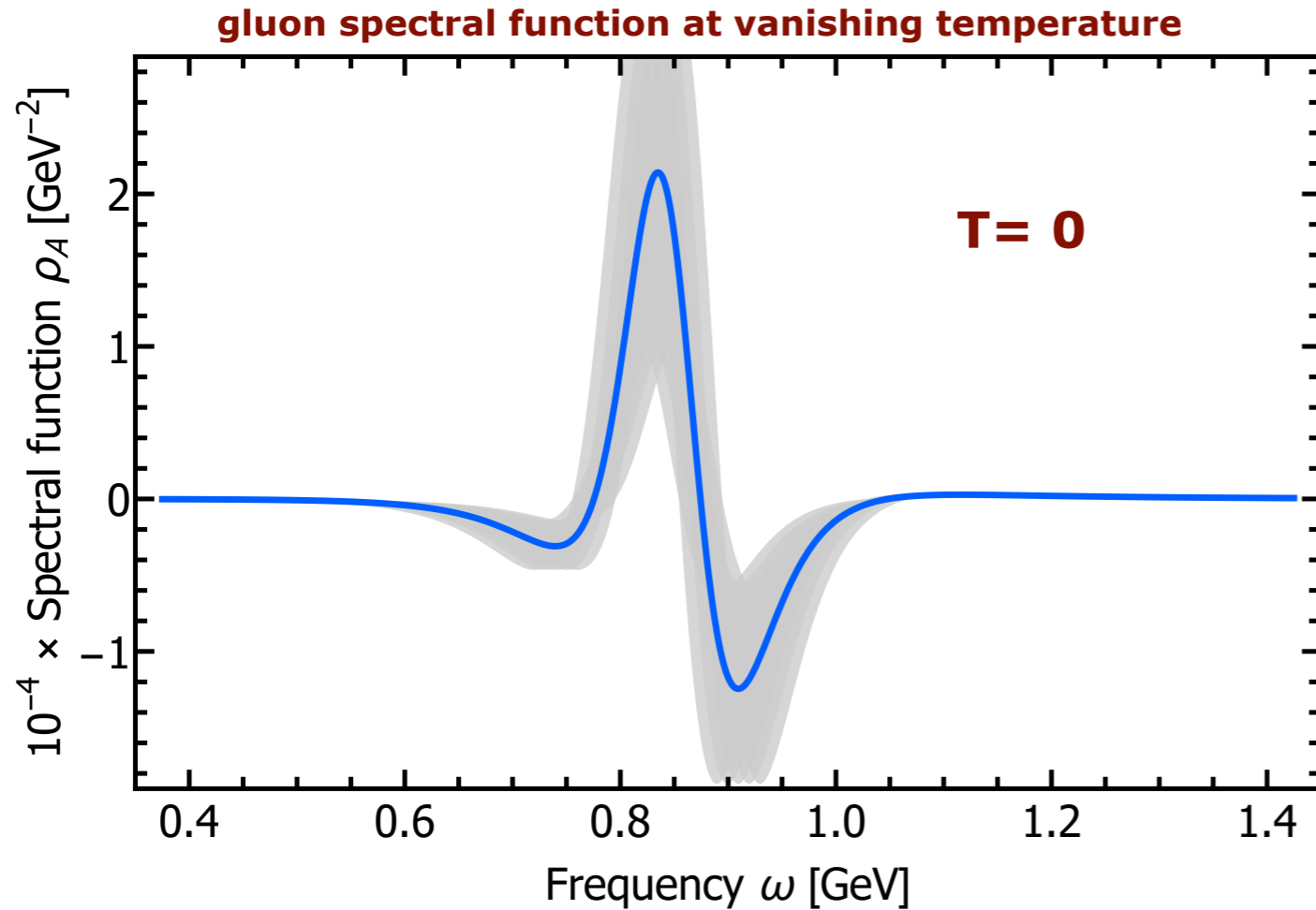
$$\rho(p) = 2 \text{Im} \langle A A \rangle_{\text{ret}}(p)$$

**T=1.44 T<sub>C</sub>**

chromo-magnetic



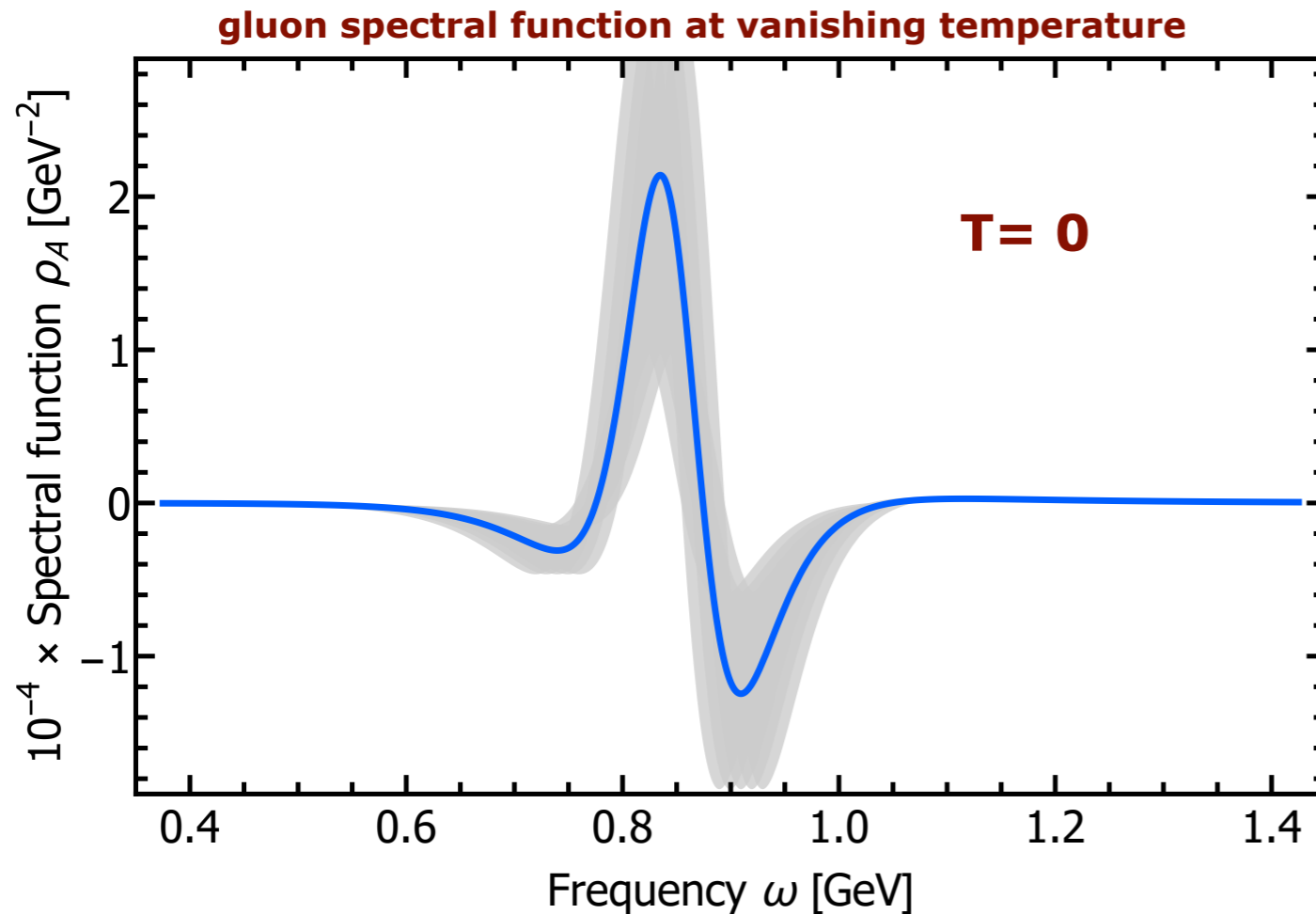
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novel analytic IR (& UV) behaviour and qualitatively refined reconstruction

# Single particle spectral functions



$$\rho(p) = 2 \text{Im} \langle A A \rangle_{\text{ret}}(p)$$

**novel analytic IR (& UV) behaviour and qualitatively refined reconstruction**

'Those are my methods (principles), and if you don't like them...well, I have others'

direct computation

Groucho Marx

**Real-time FRG: wait 5 minutes**

Cyrol, JMP, Rothkopf, Wink, arXiv:1804.00945

# Transport coefficients

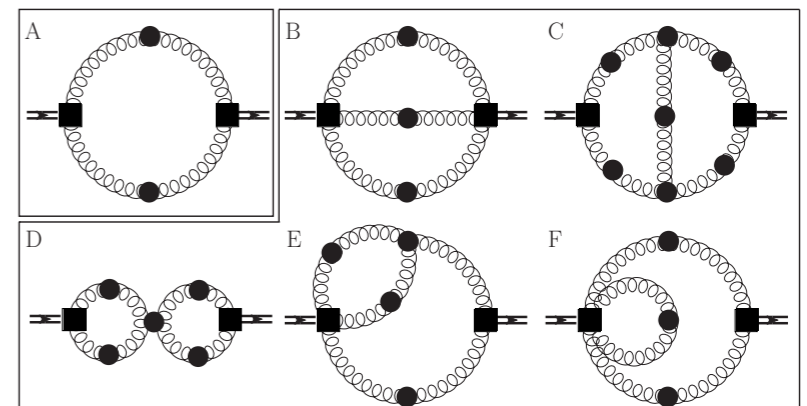
## viscosity over entropy ratio in Yang-Mills theory

### Kubo relation

$$\eta = \frac{1}{20} \left. \frac{d}{d\omega} \right|_{\omega=0} \rho_{\pi\pi}(\omega, 0)$$

### '3-loop' exact functional relation for $\rho_{\pi\pi}$

#### 1 & 2-loop terms



Haas, Fister, JMP, PRD 90 (2014) 091501

Christiansen, Haas, JMP, Strodthoff, PRL 115 (2015) 112002



# Transport coefficients

## viscosity over entropy ratio in Yang-Mills theory

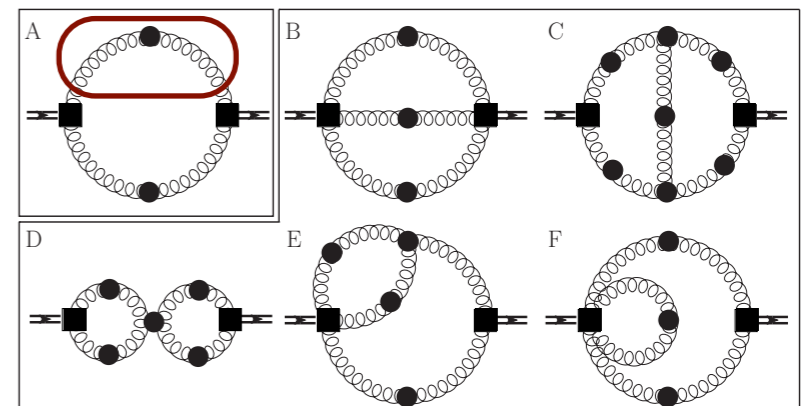
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$$\eta = \frac{1}{20} \left. \frac{d}{d\omega} \right|_{\omega=0} \rho_{\pi\pi}(\omega, 0)$$

Gluon spectral function

### '3-loop' exact functional relation for $\rho_{\pi\pi}$

#### 1 & 2-loop terms



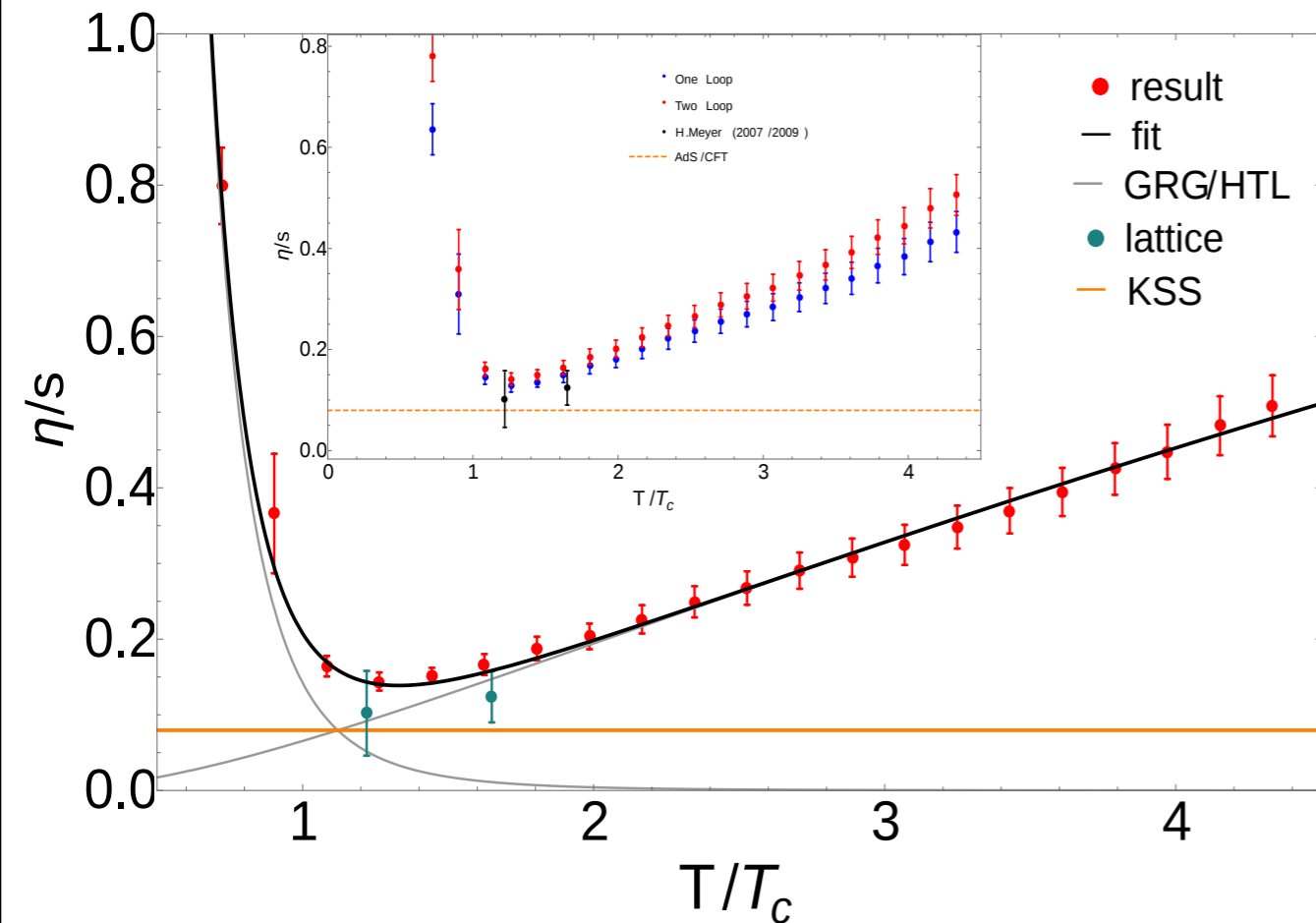
Haas, Fister, JMP, PRD 90 (2014) 091501

Christiansen, Haas, JMP, Strodthoff, PRL 115 (2015) 112002

# Transport coefficients

## viscosity over entropy ratio in Yang-Mills theory

### Yang-Mills viscosity over entropy ratio

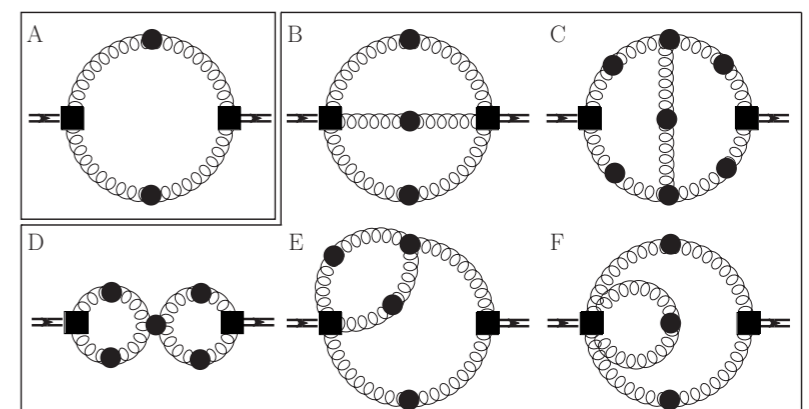


### Kubo relation

$$\eta = \frac{1}{20} \left. \frac{d}{d\omega} \right|_{\omega=0} \rho_{\pi\pi}(\omega, 0)$$

### '3-loop' exact functional relation for $\rho_{\pi\pi}$

#### 1 & 2-loop terms



recent lattice results: Astrakhantsev, Braguta, Kotov, JHEP 1704 (2017) 101  
arXiv:1804.02382

**Aiming at apparent convergence**

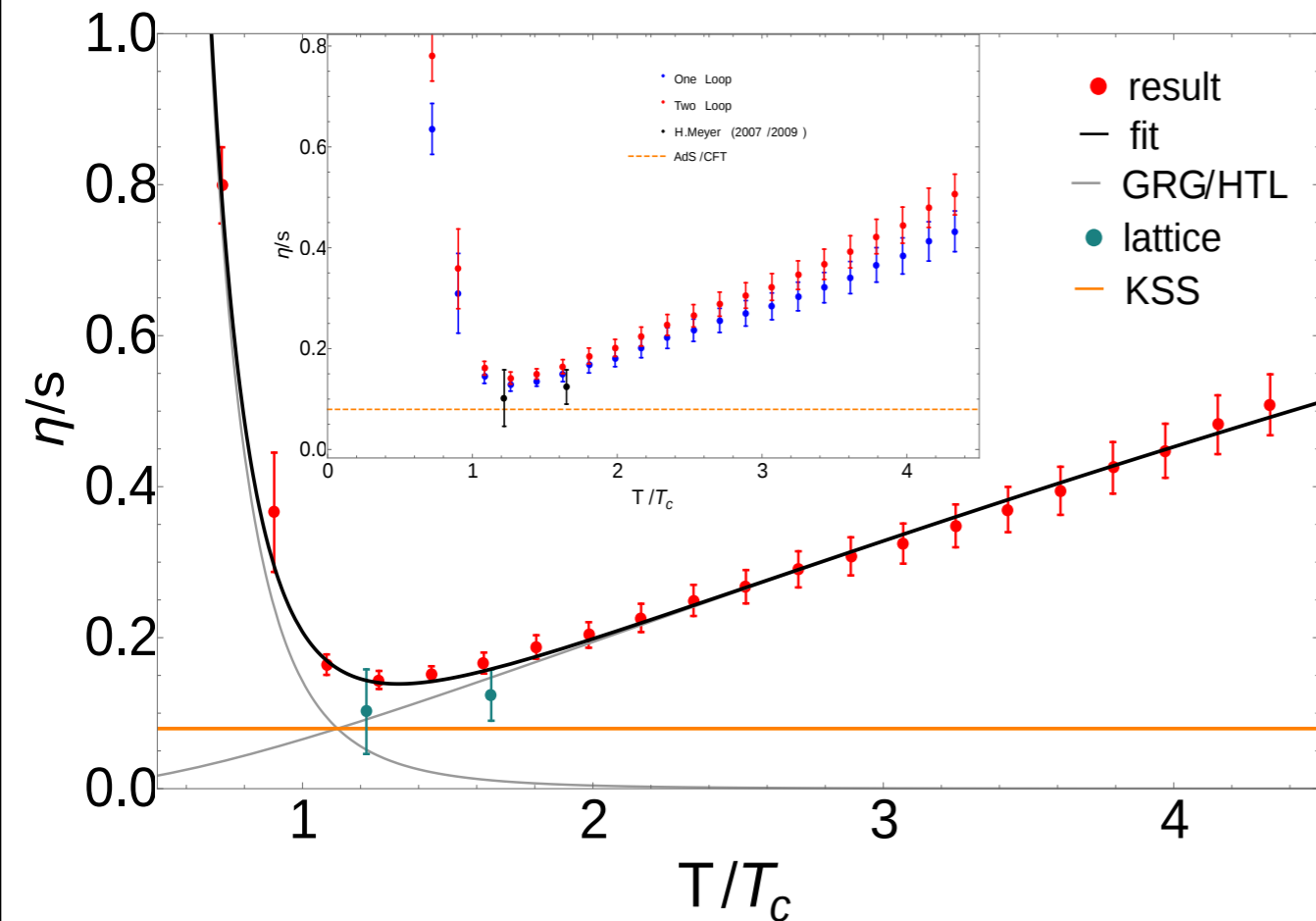
Haas, Fister, JMP, PRD 90 (2014) 091501

Christiansen, Haas, JMP, Strodthoff, PRL 115 (2015) 112002

# Transport coefficients

## QCD - estimate for viscosity over entropy ratio

### viscosity over entropy ratio



$$\gamma_{\text{grg}} \approx 5$$

$$\gamma_{\text{qgp}} \approx 1.6$$

**pure glue**

$$a_{\text{qgp}} \approx 0.15$$

$$a_{\text{hrg}} \approx 0.14$$

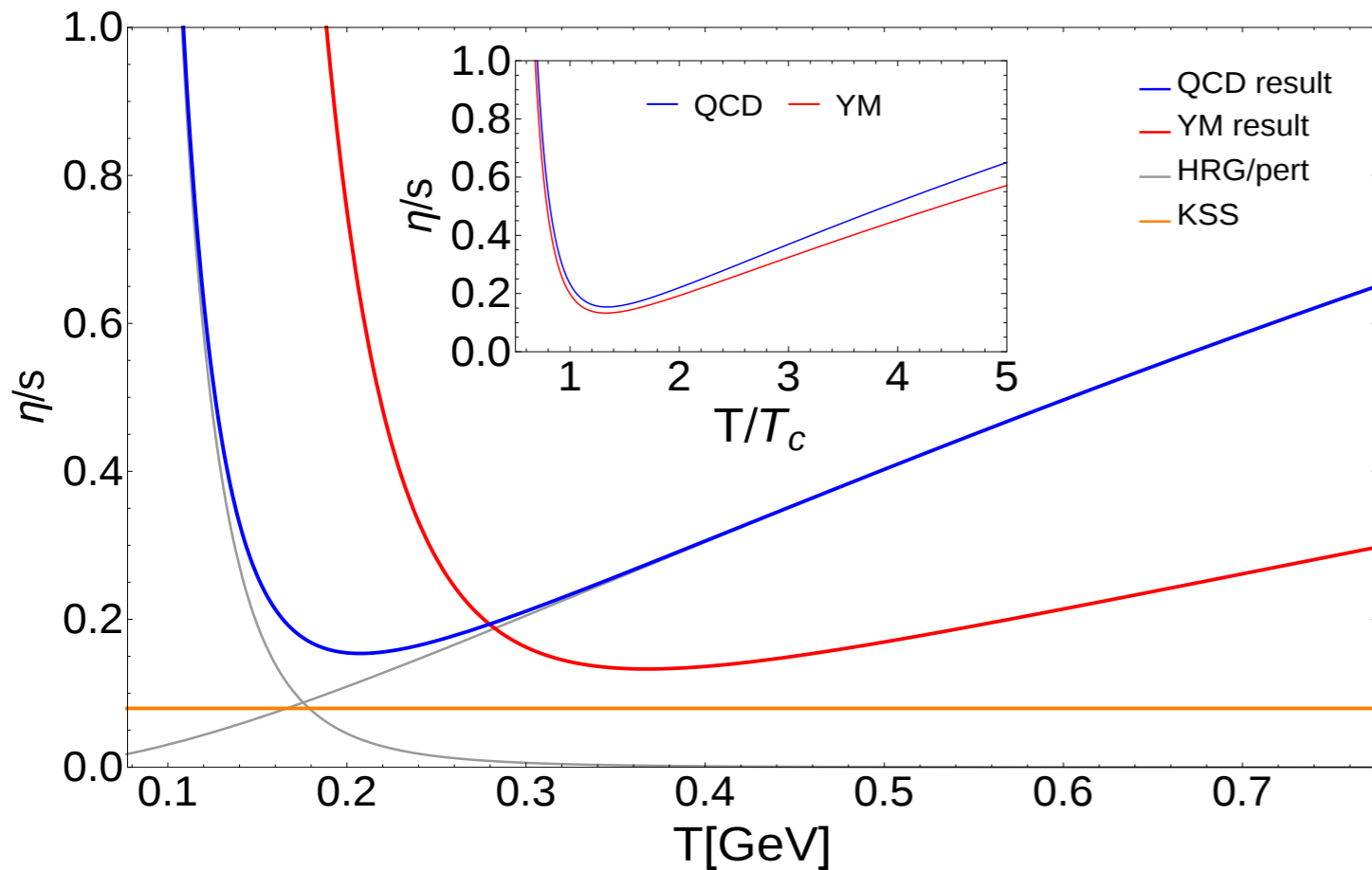
$$c \approx 0.66$$

$$\frac{\eta}{s}(T) = \frac{a_{\text{qgp}}}{\alpha_s^{\gamma_{\text{qgp}}}(cT/T_c)} + \frac{a_{\text{grg}}}{(T/T_c)^{\gamma_{\text{grg}}}}$$

# Transport coefficients

## QCD - estimate for viscosity over entropy ratio

### viscosity over entropy ratio



$$a_{\text{qgp}} \approx 0.2$$

$$a_{\text{hrg}} \approx 0.16$$

$$c \approx 0.79$$

**QCD**

$$\gamma_{\text{grg}} \approx 5$$

$$\gamma_{\text{qgp}} \approx 1.6$$

**pure glue**

$$a_{\text{qgp}} \approx 0.15$$

$$a_{\text{hrg}} \approx 0.14$$

$$c \approx 0.66$$

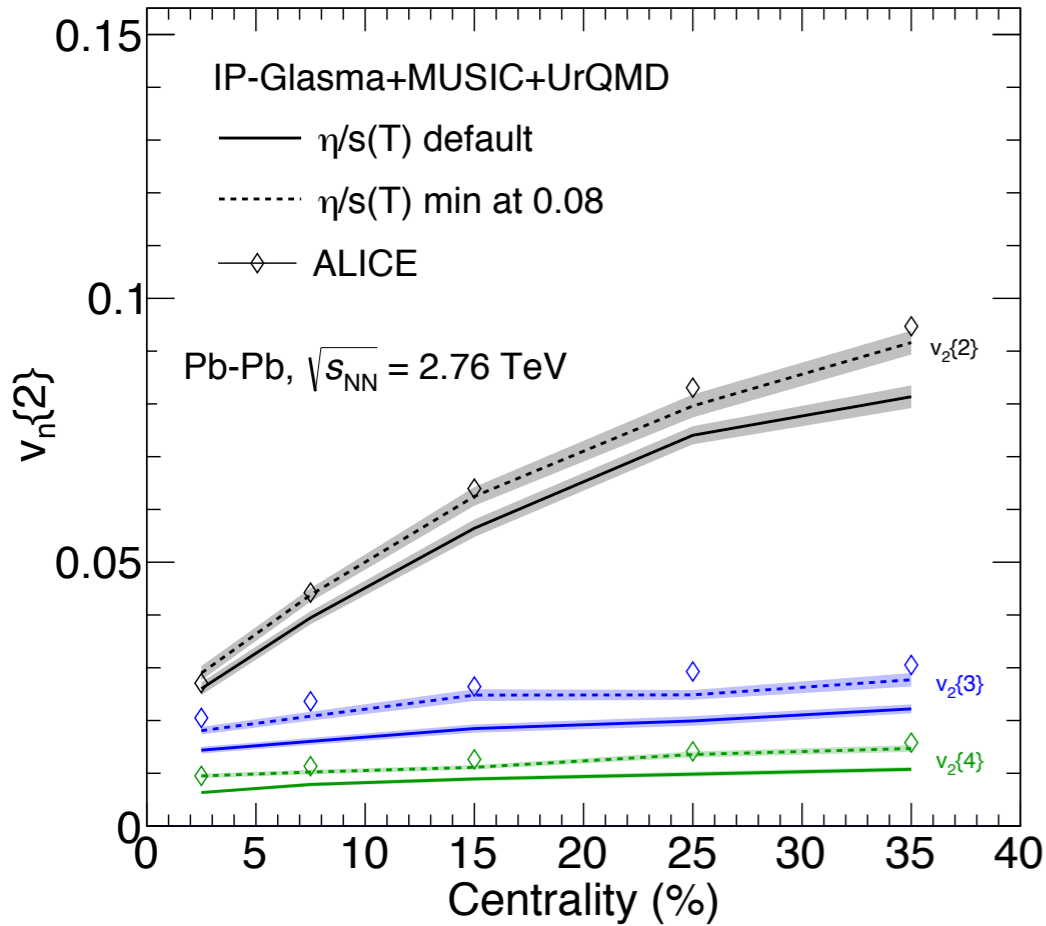
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# QCD-assisted hydrodynamics

Dubla, Masciocchi, JMP, Schenke, Shen, Stachel, arXiv:1805.02985

## IP-Glasma - MUSIC - UrQMD

### $v_n$ as function of centrality

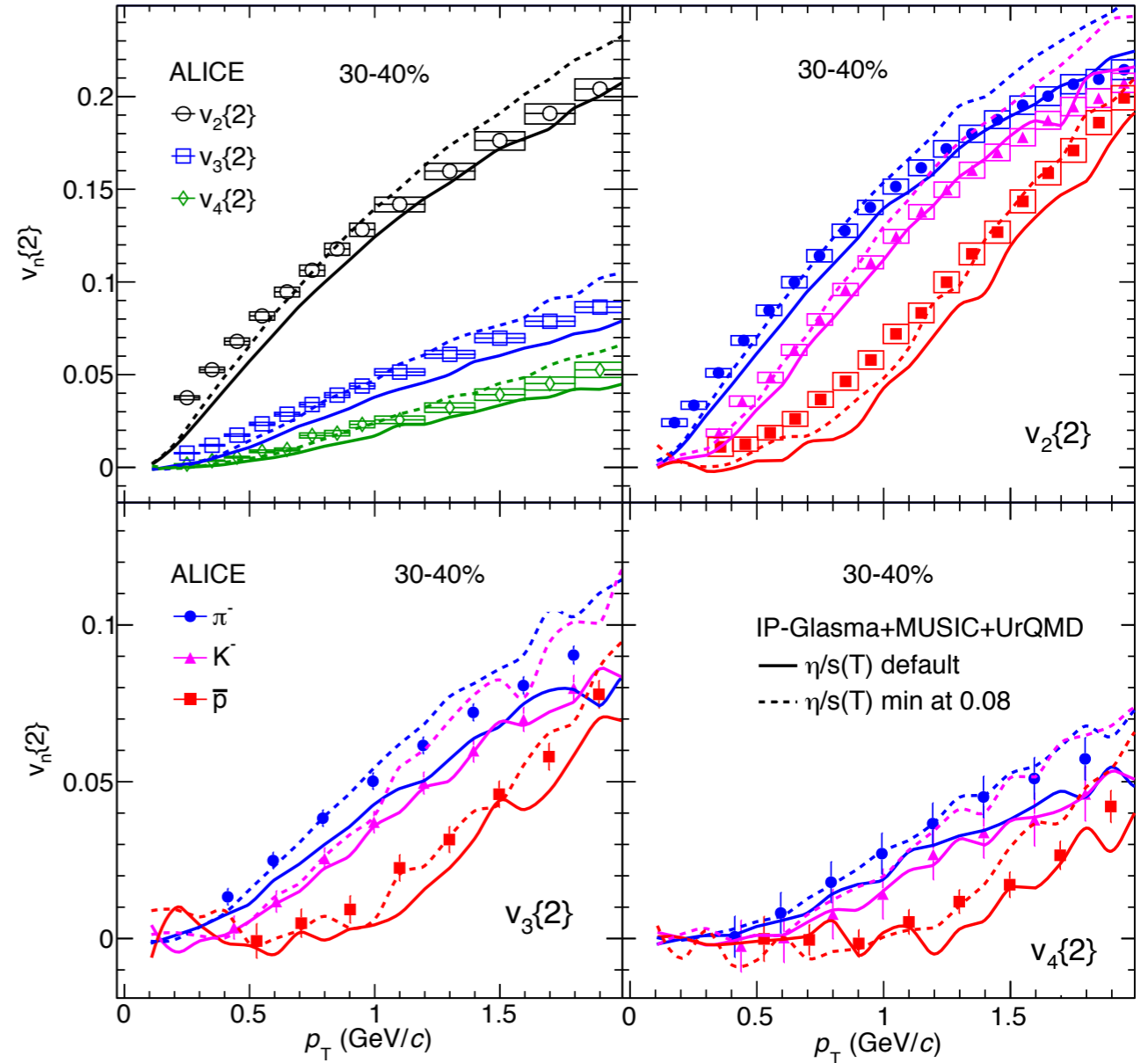


### Test of systematic error

$$\eta/s(T) \rightarrow \eta/s(T) + d$$

$$d \in [-0.06, 0]$$

### $v_n$ as function of $p_T$



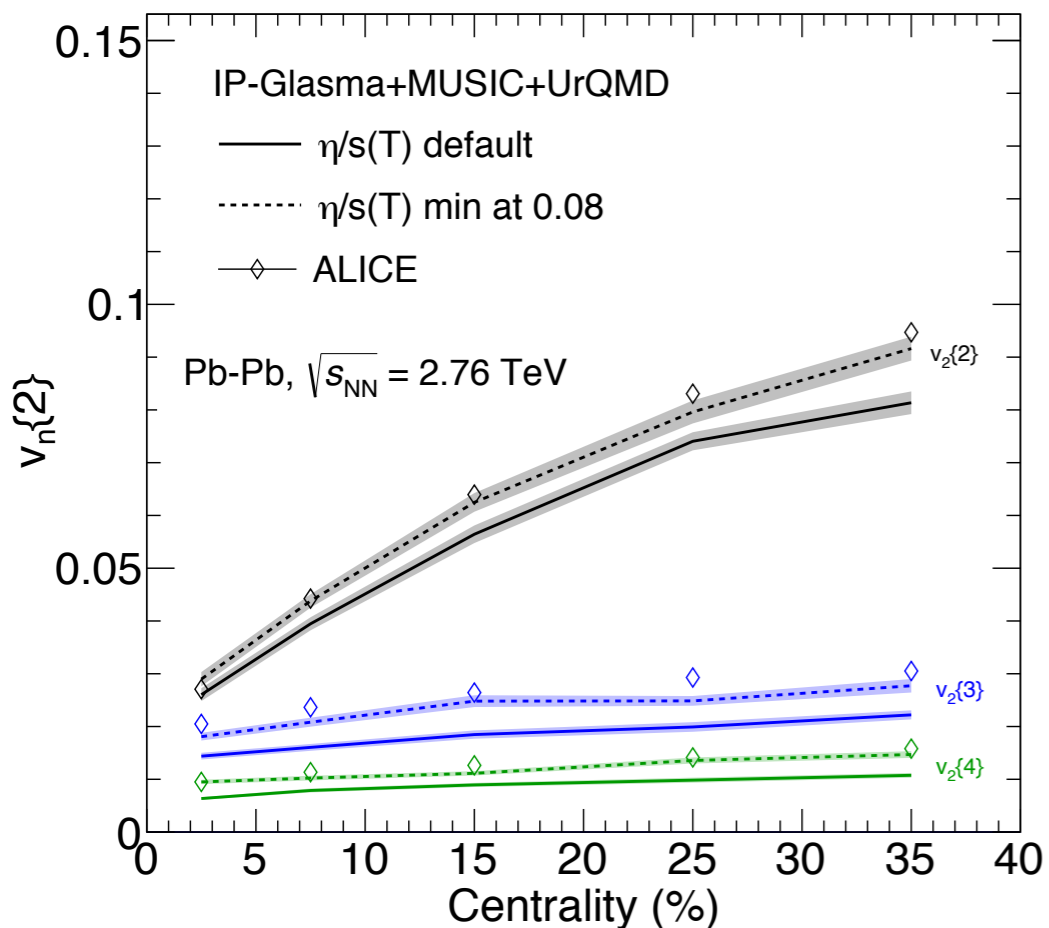
$$\frac{dN}{d(\varphi - \Psi_R)} = \frac{N_0}{2\pi} \left( 1 + 2 \sum_n v_n \cos[n(\varphi - \Psi_R)] \right)$$

# QCD-assisted hydrodynamics

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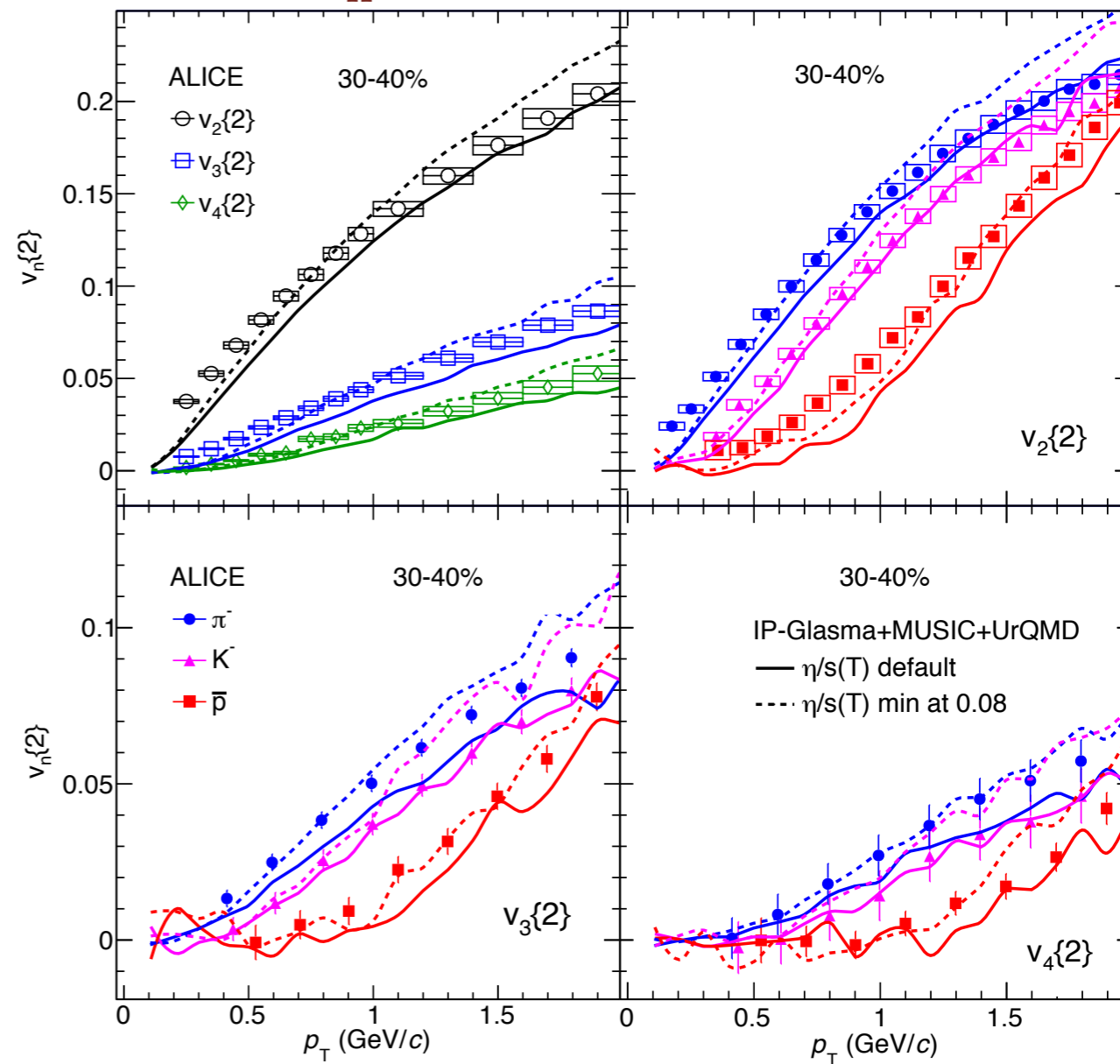
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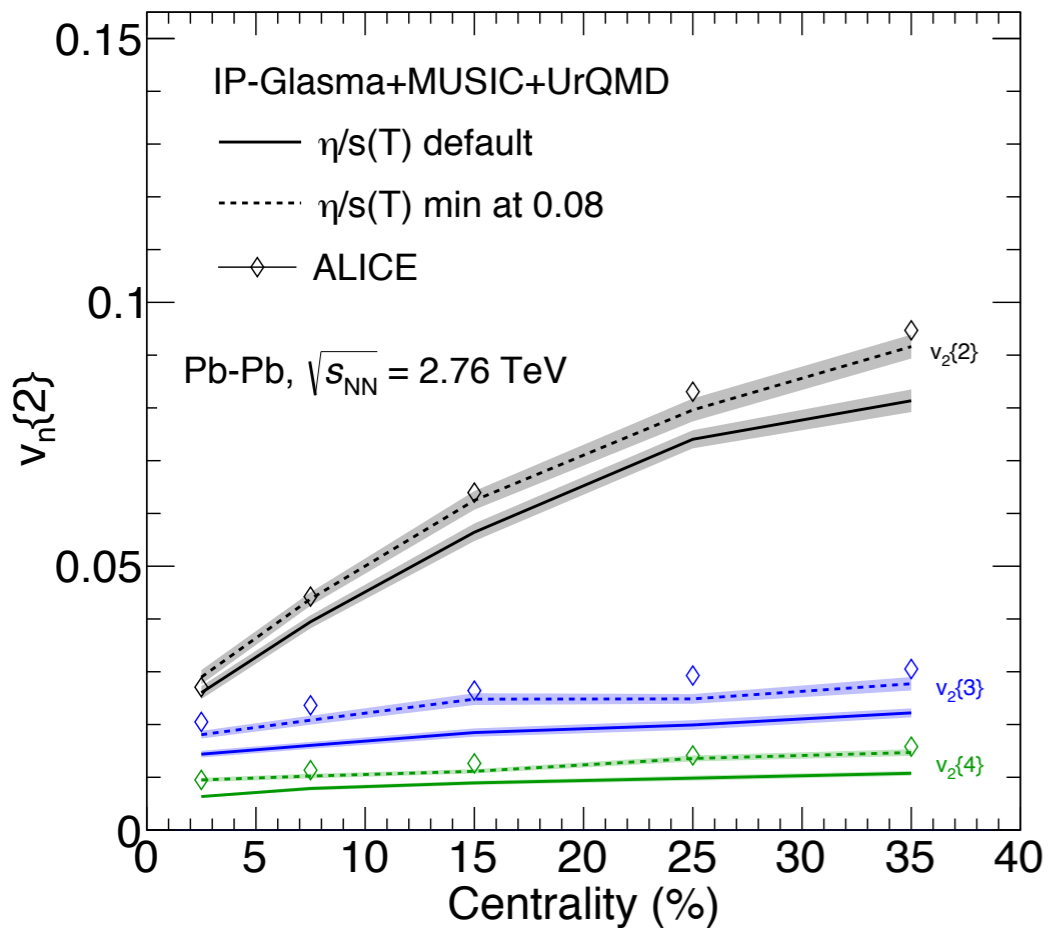


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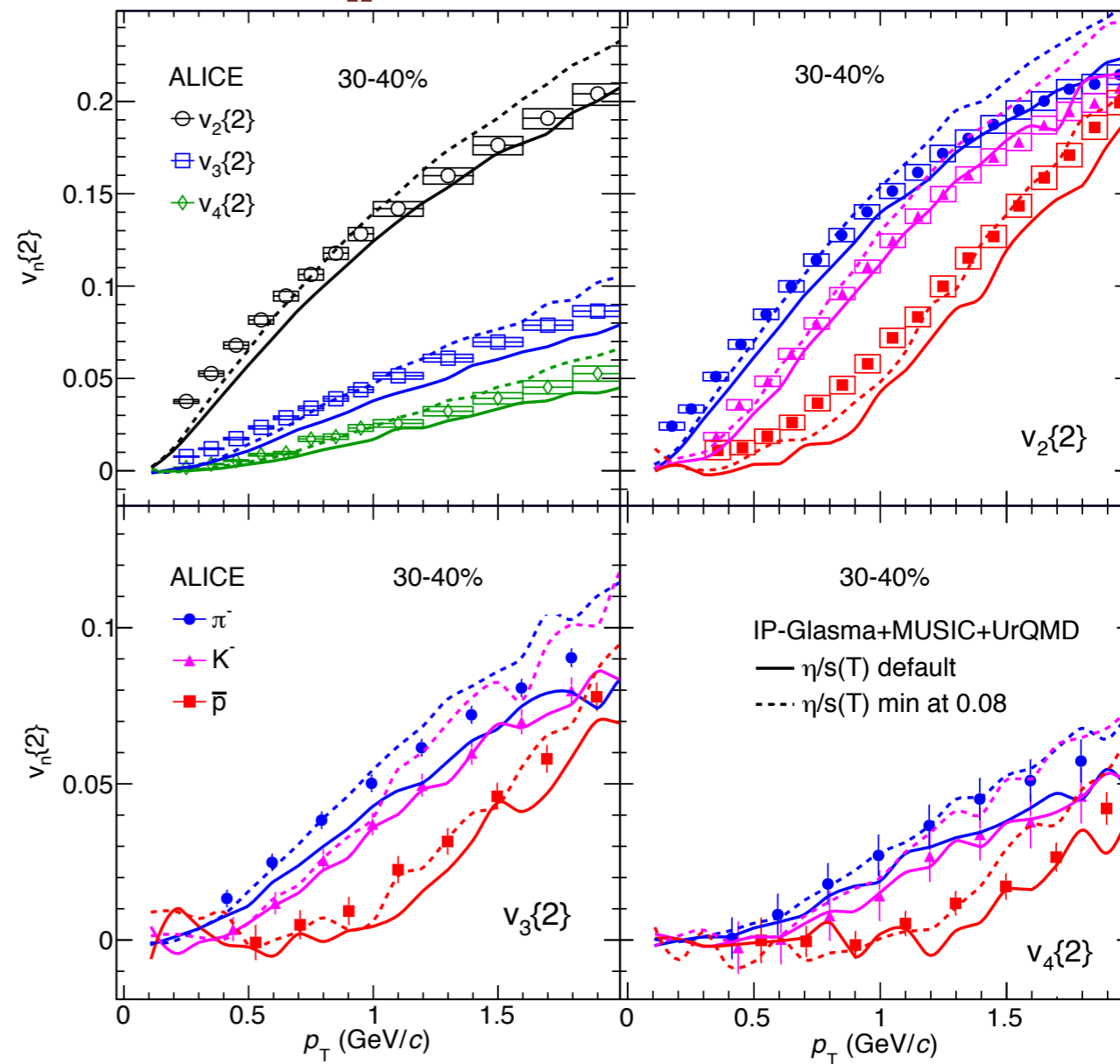
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Initial state fluctuations?

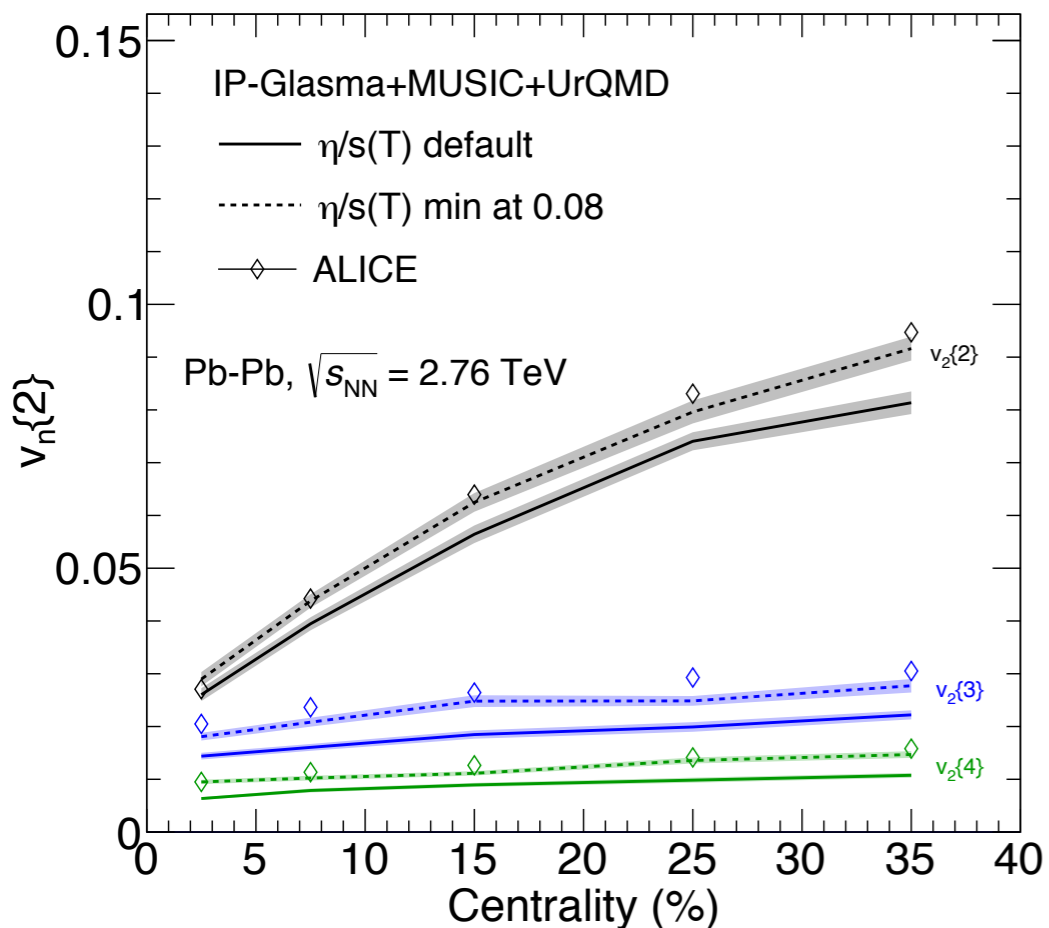
Kinetic phase?

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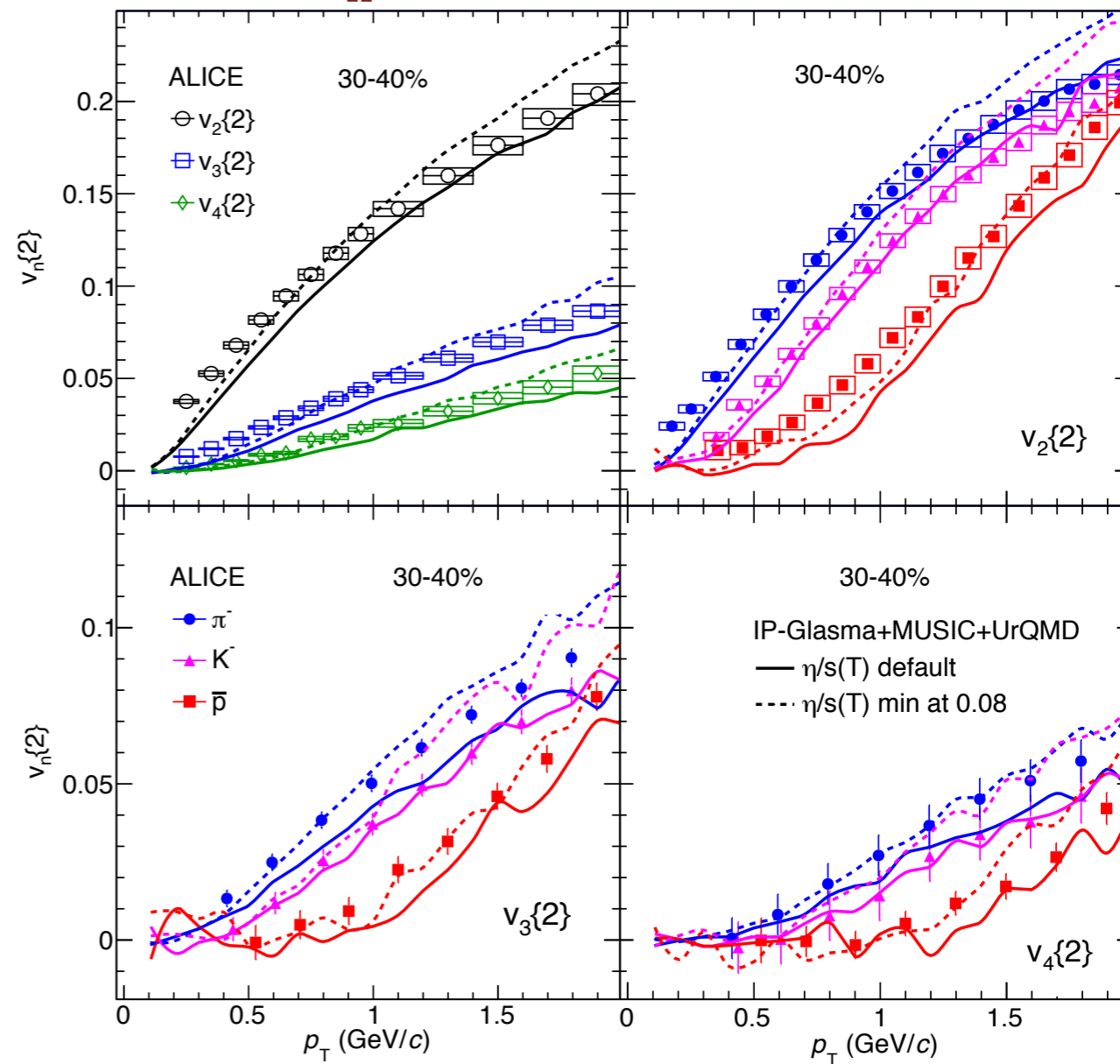
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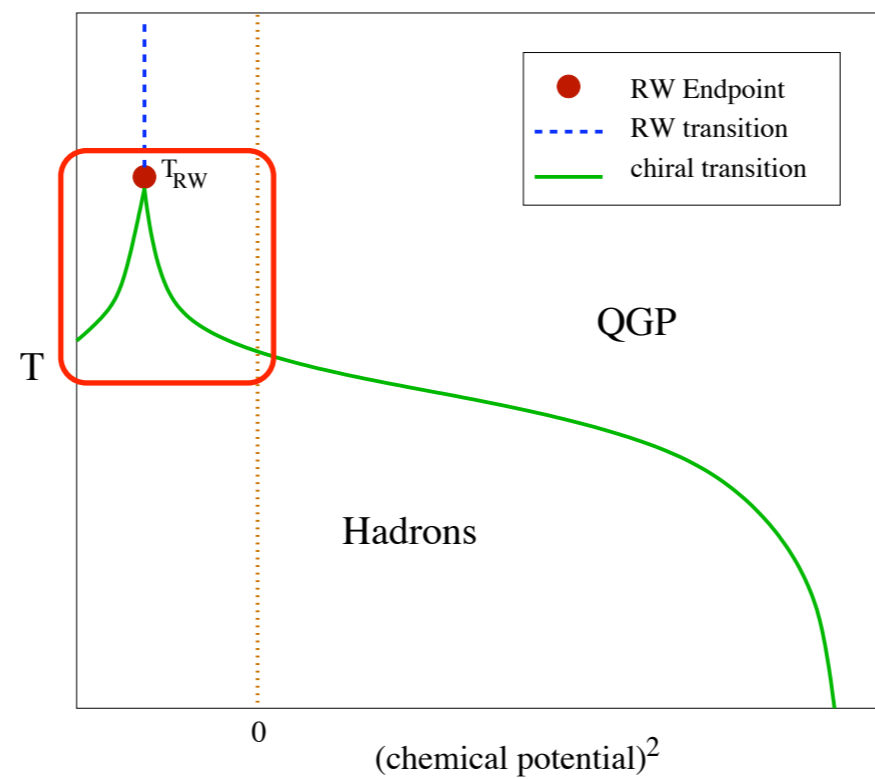
'Steady-state' hydro?

Kinetic phase?

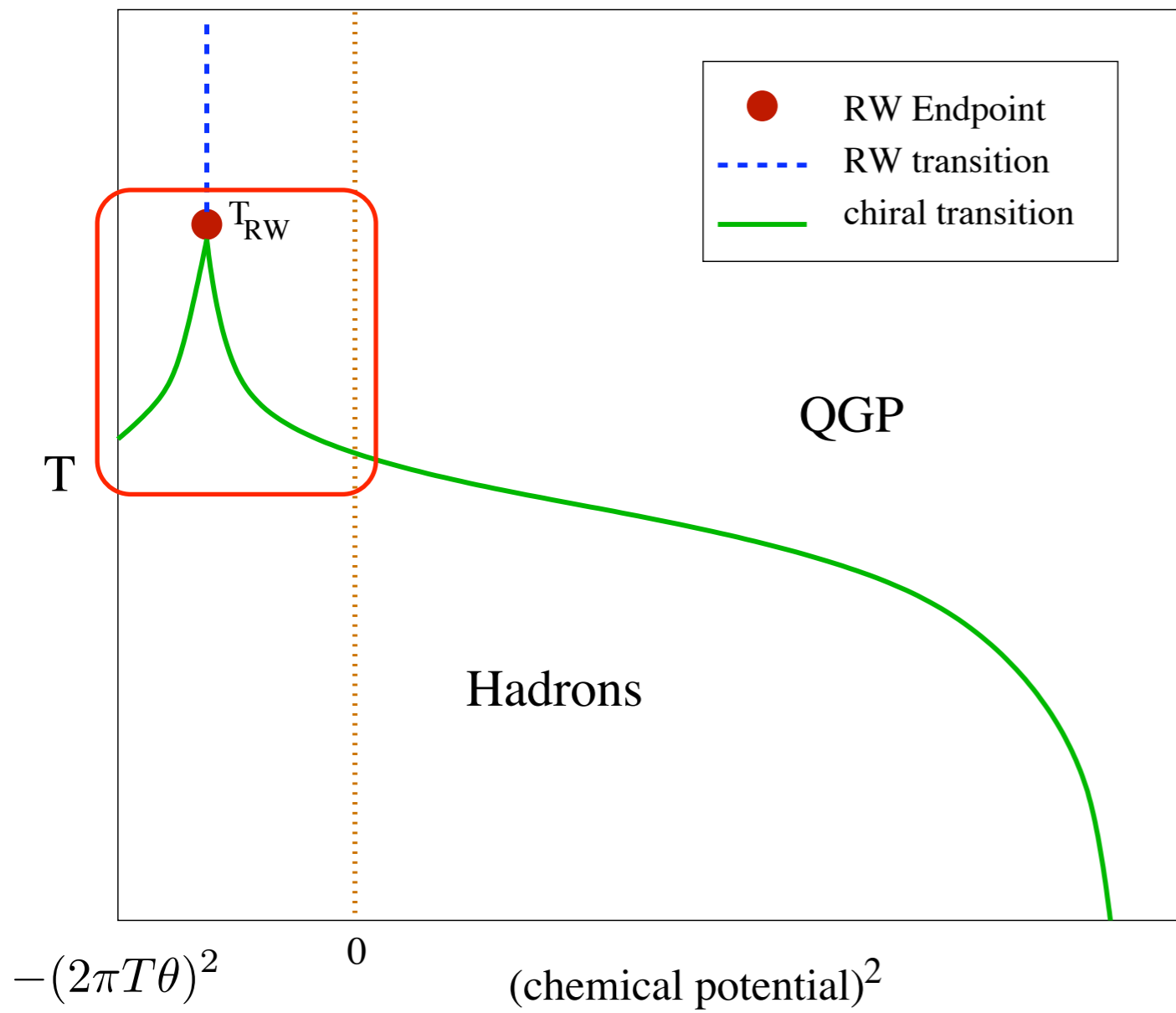
Hadronisation & freeze-out?



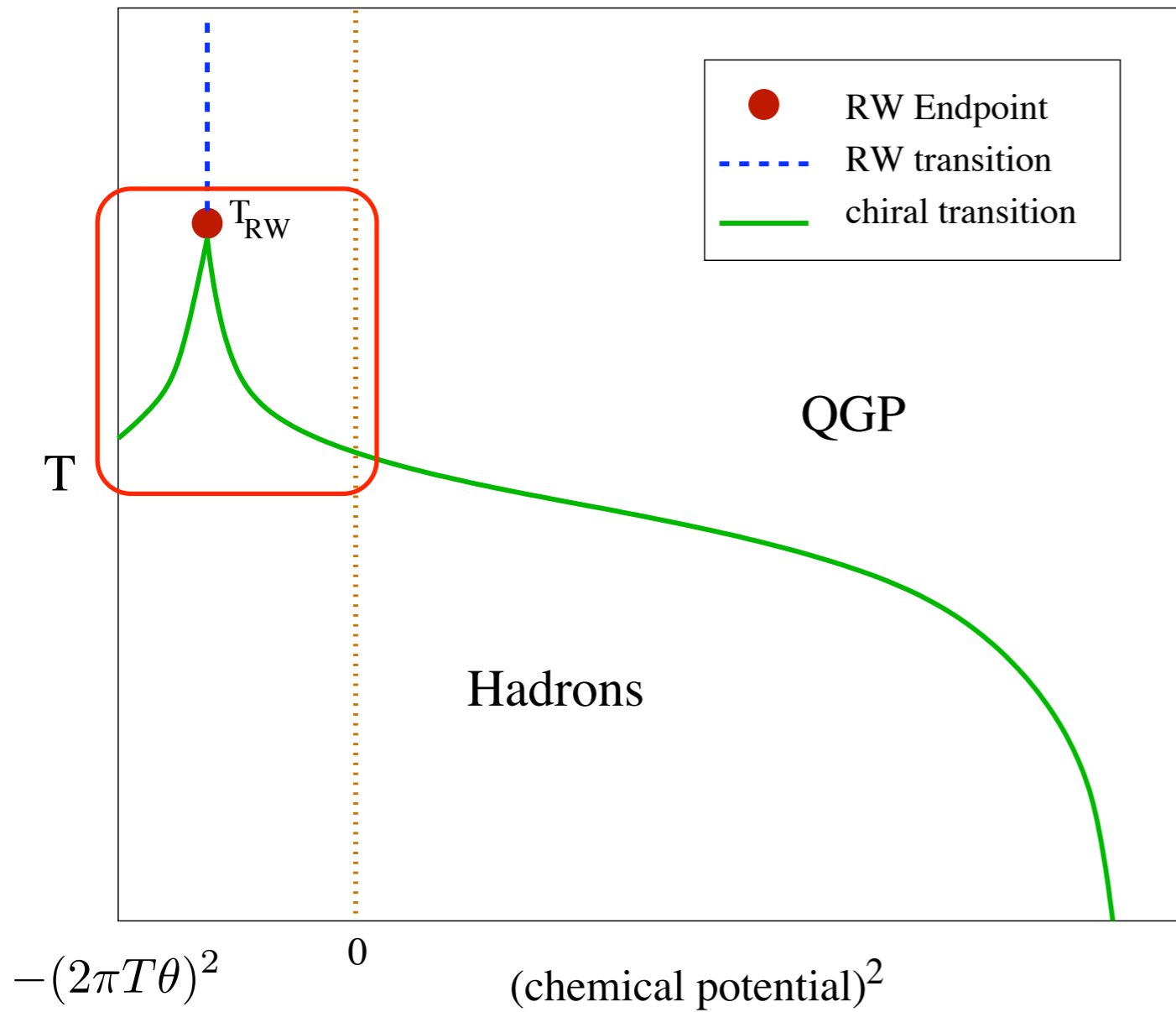
# QCD at imaginary chemical potential



# Imaginary chemical potential



# Imaginary chemical potential



## Dirac term

$$\int_x \bar{q} \cdot (i\not{D} + i m_\psi + i\mu\gamma_0) \cdot q$$

$$\mu = 2\pi T\theta i$$

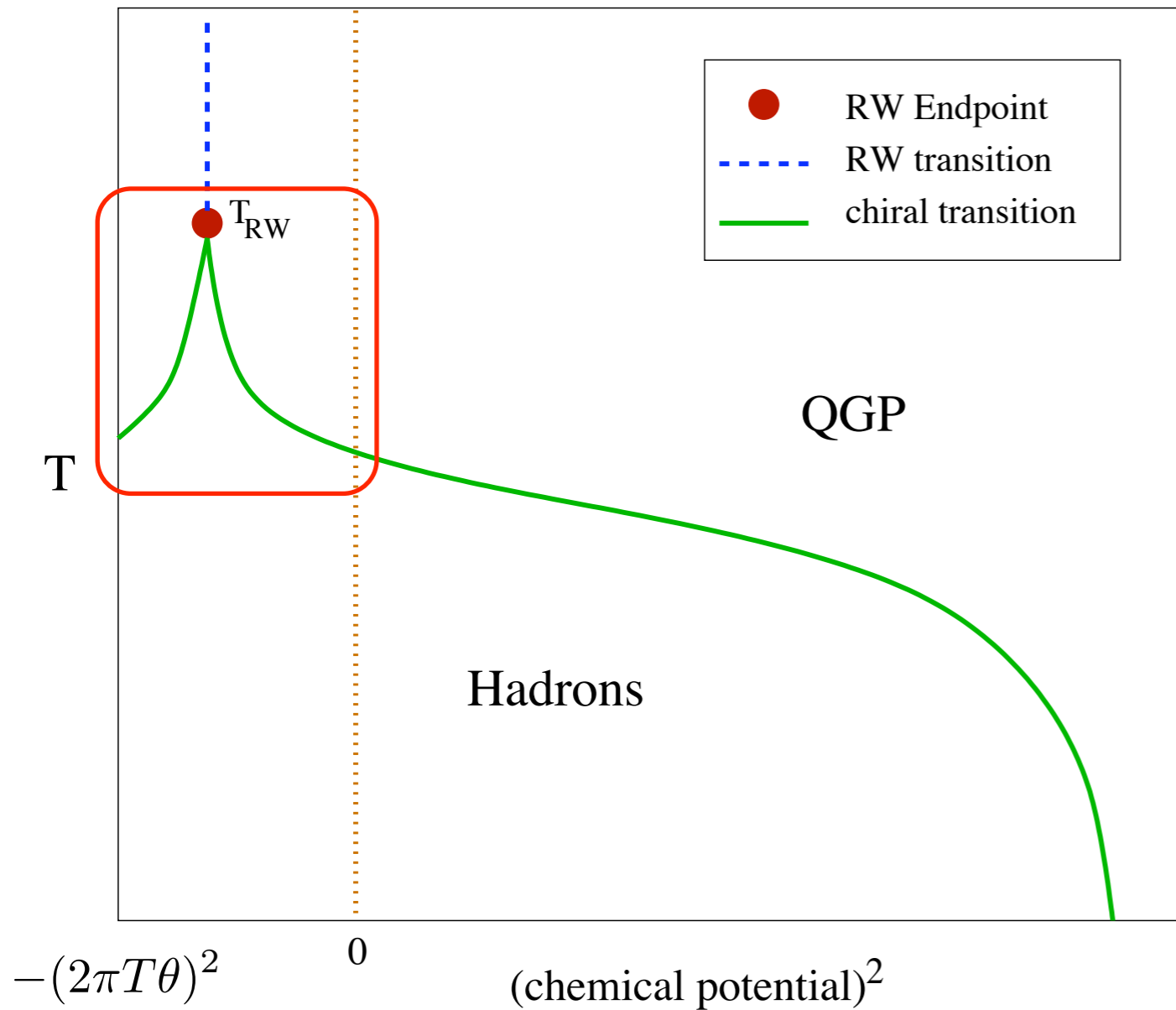
$$\int_x \bar{q}_\theta \cdot (i\not{D} + i m_\psi) \cdot q_\theta$$

$$q_\theta(t, \vec{x}) = e^{2\pi T\theta i t} q(t, x)$$

## Periodicity

$$q_\theta(t + \beta, \vec{x}) = -e^{2\pi\theta i} q_\theta(t, x)$$

# Imaginary chemical potential



## Dirac term

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$$q_\theta(t + \beta, \vec{x}) = -e^{2\pi\theta i} q_\theta(t, x)$$

## Roberge-Weiss symmetry

$$Z_\theta = Z_{\theta+1/3}$$

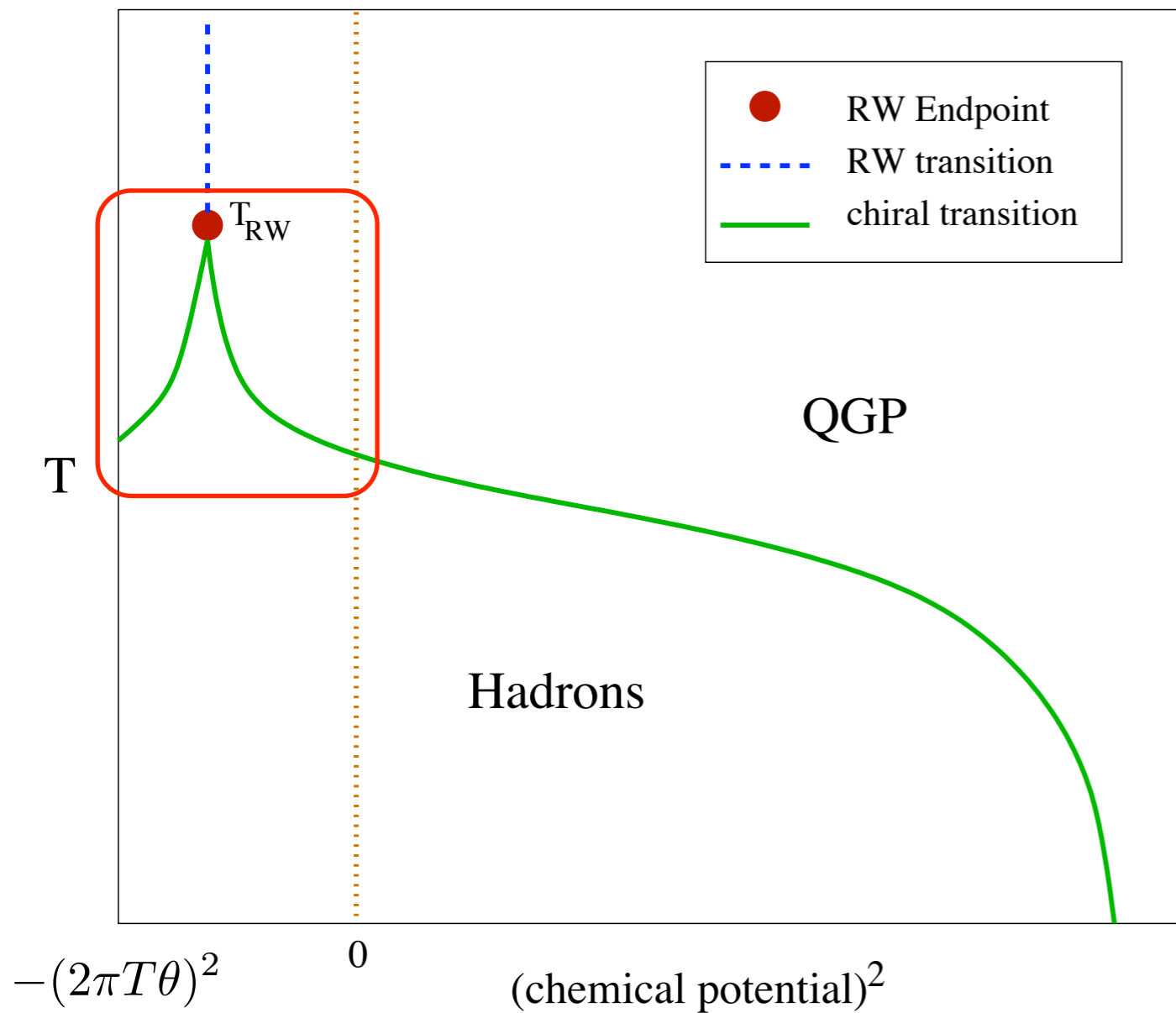
Partition function

via a center transformation

$$e^{\frac{2}{3}\pi i} \mathbb{1} \in \text{center}[SU(3)]$$

gauge field insensitive to center transformations

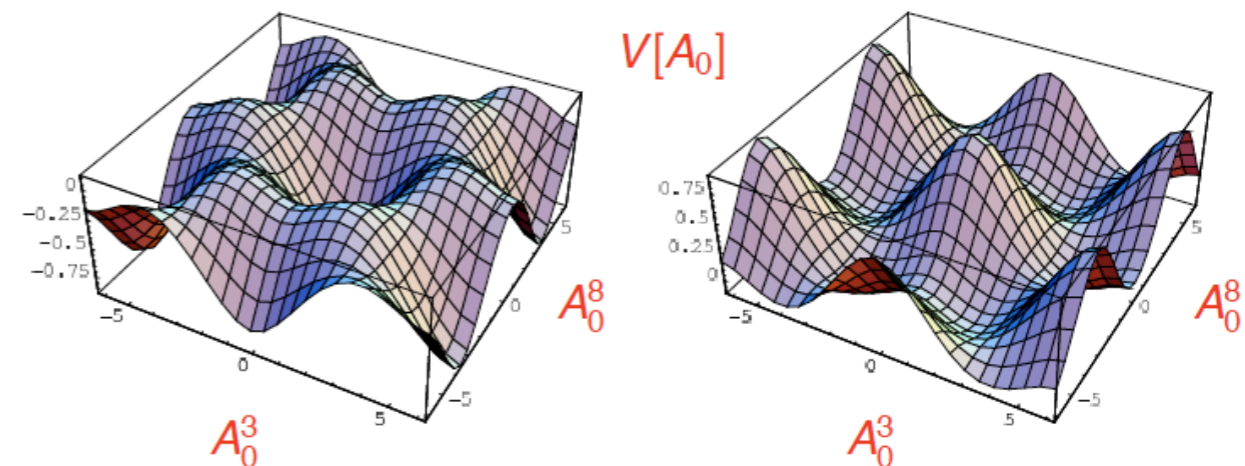
# Imaginary chemical potential



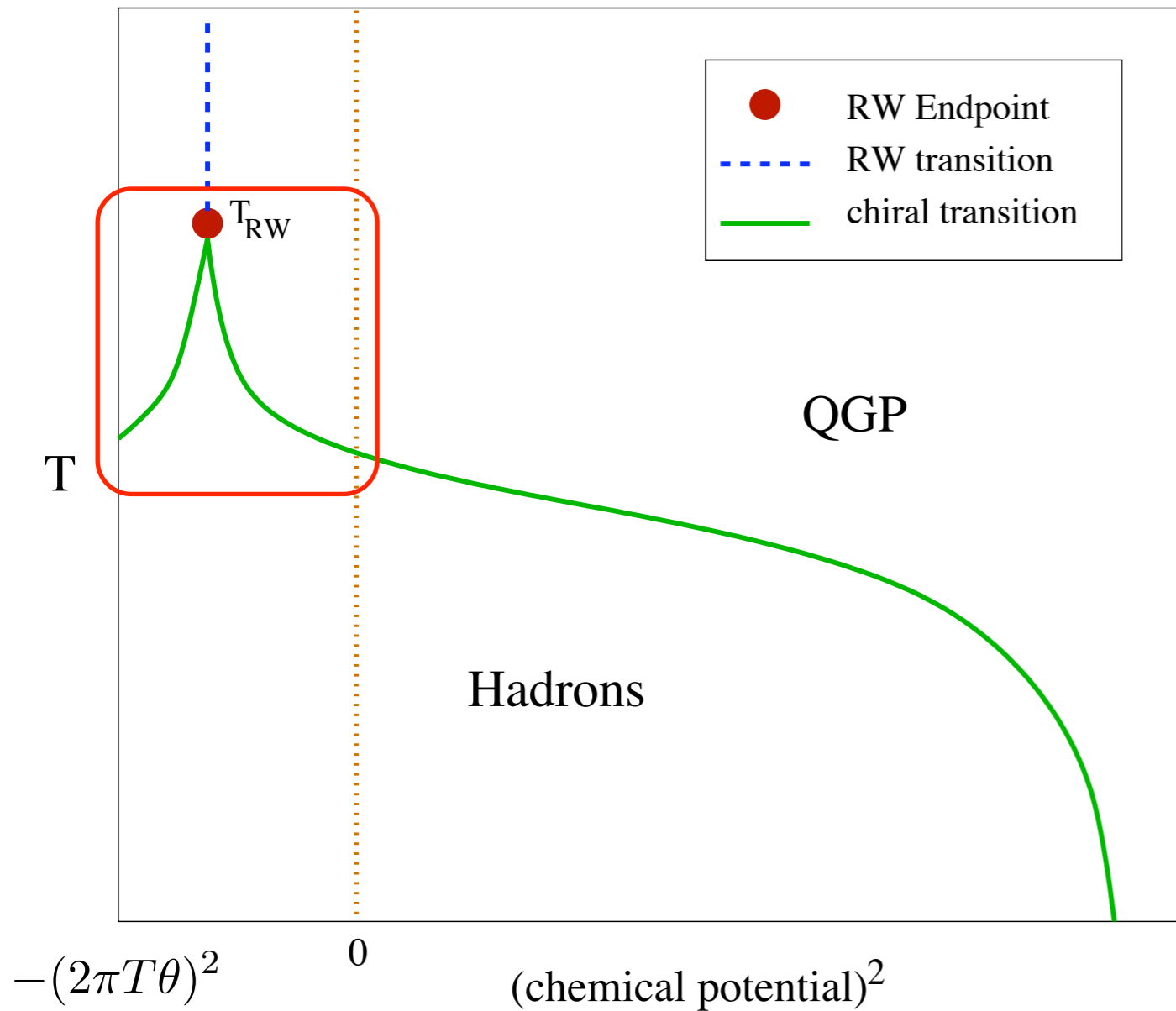
## Periodicity

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## Polyakov loop potential



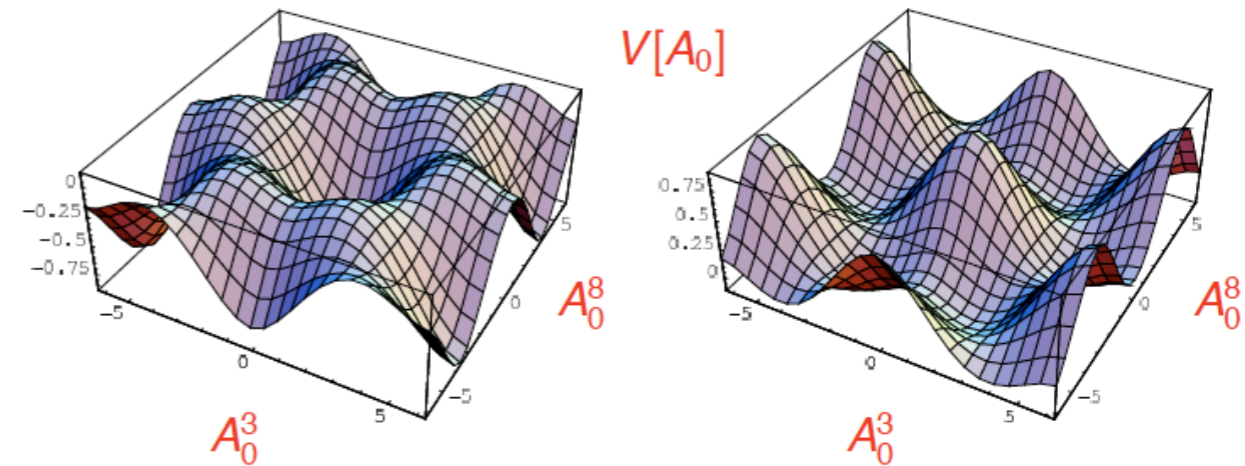
# Imaginary chemical potential



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# Imaginary chemical potential

confinement order parameters

---

$$q_\theta(t + \beta, \vec{x}) = -e^{2\pi\theta i} q_\theta(t, x)$$

Center-sensitive observables

$$\mathcal{O}_\theta = \langle O[q_\theta] \rangle_\theta$$

$$\mathcal{O}_\theta = \langle O[q_\theta] \rangle_{\theta=0}$$

at imaginary chemical potential

at vanishing chemical potential

Dual order parameters

$$\tilde{\mathcal{O}} = \int_0^1 d\theta \mathcal{O}_\theta e^{-2\pi i\theta}$$

$$\tilde{\mathcal{O}} \xrightarrow{z} z\tilde{\mathcal{O}}$$

$$z = 1, e^{\frac{2}{3}\pi i}, e^{\frac{4}{3}\pi i}$$

# Imaginary chemical potential

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vanishing chemical potential

Gattringer '06

Synatschke, Wipf, Wozar '07

Bruckmann, Hagen, Bilgici, Gattringer '08

Lattice

FunMethods

Fischer '09

Fischer, Maas, Müller '10



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imaginary chemical potential

average over diff. theories

$$\tilde{\mathcal{O}}[\langle A_0 \rangle_{\theta=0}]$$

breaking of RW-symmetry

Braun, Haas, Marhauser, JMP '09

Lattice

FunMethods

vanishing chemical potential

Gattringer '06

Synatschke, Wipf, Wozar '07

Bruckmann, Hagen, Bilgici, Gattringer '08

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at imaginary chemical potential

at vanishing chemical potential

FRG

DSE

FRG

DSE

1-loop

1-loop

dual quark propagator

2-loop

3-loop

1-loop

—

dual pressure

—

—

1-loop

1-loop

dual density

2-loop

3-loop

1-loop

1-loop

dual susceptibilities

2-loop

3-loop

dual pressure = -T dual density

dual susceptibility = T dual density

# Imaginary chemical potential

## confinement order parameters

$$\tilde{\mathcal{O}} = \int_0^1 d\theta \mathcal{O}_\theta e^{-2\pi i\theta}$$

at imaginary chemical potential

at vanishing chemical potential

FRG

DSE

FRG

DSE

1-loop

1-loop

dual quark propagator

2-loop

3-loop

1-loop

—

dual pressure

—

—

1-loop

1-loop

dual density

2-loop

3-loop

1-loop

1-loop

dual susceptibilities

2-loop

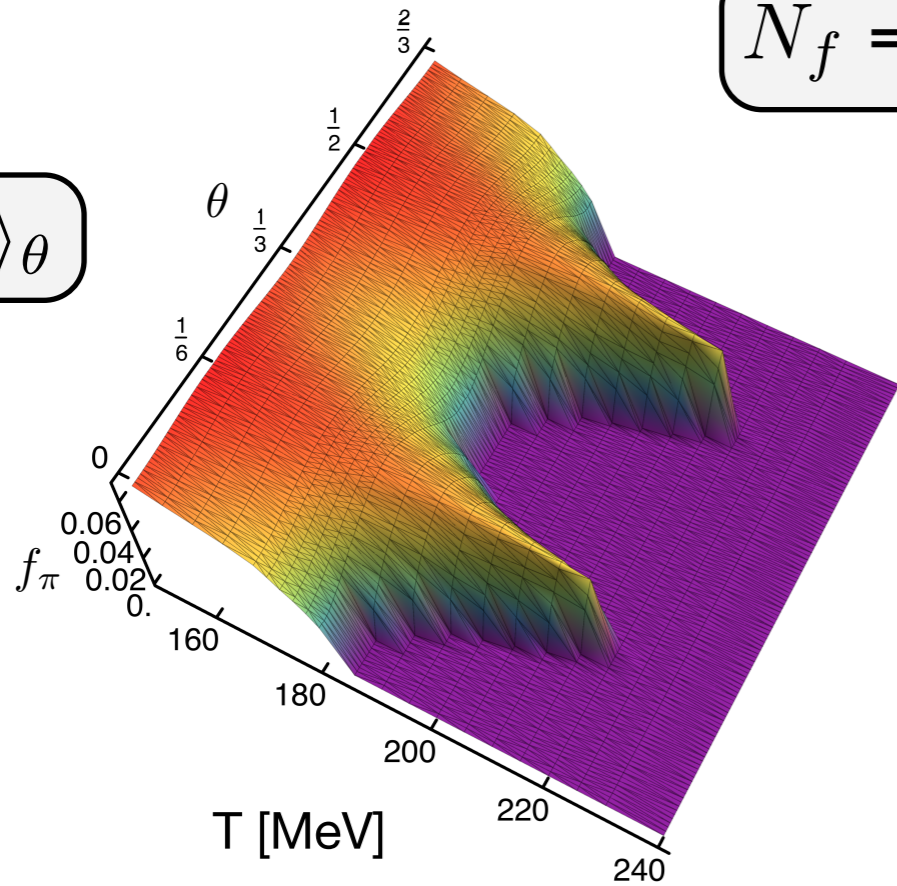
3-loop

$$\frac{1}{2\pi T i} \int_0^1 d\theta (\partial_\theta \mathcal{O}_\theta) e^{-2\pi i\theta} = \frac{1}{T} \int_0^1 d\theta \mathcal{O}_\theta e^{-2\pi i\theta}$$

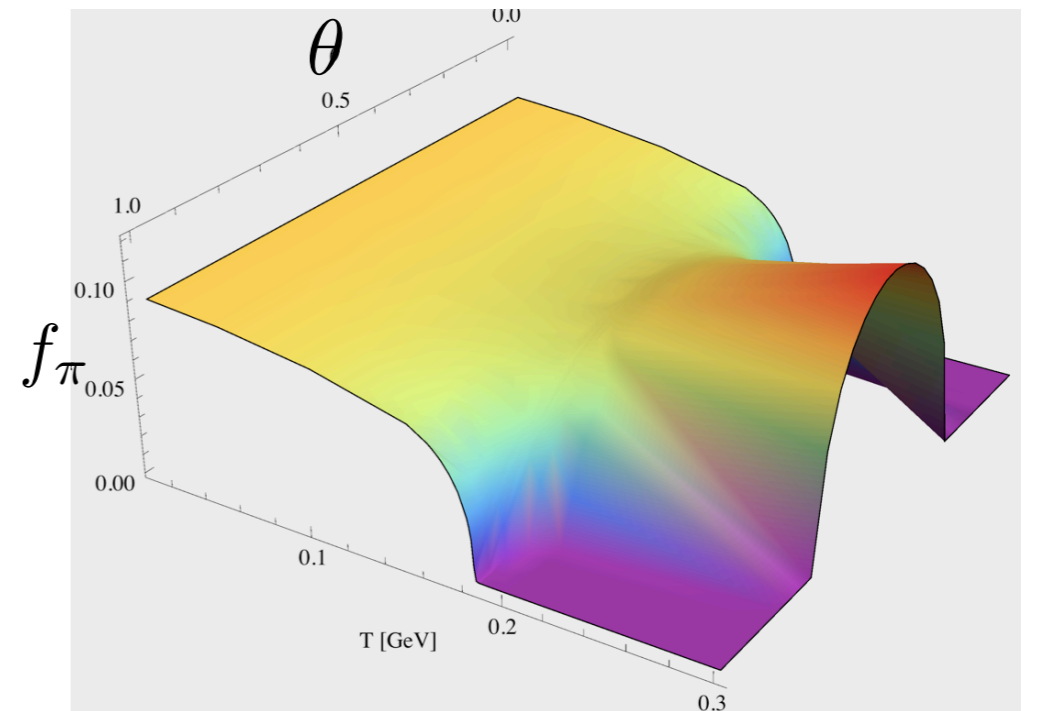
# Imaginary chemical potential

$N_f = 2$  & chiral limit

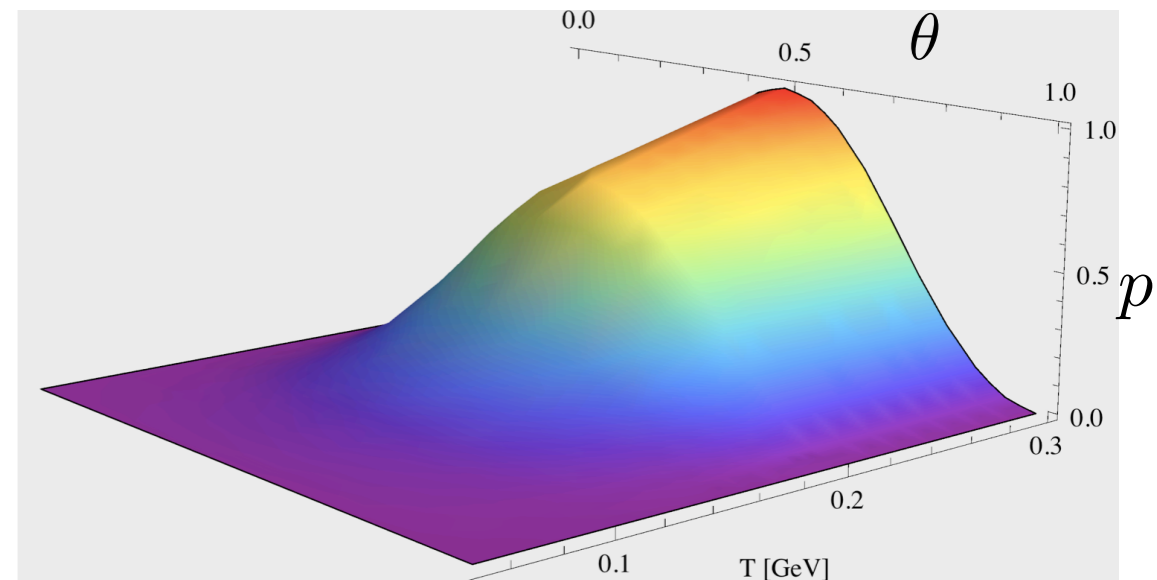
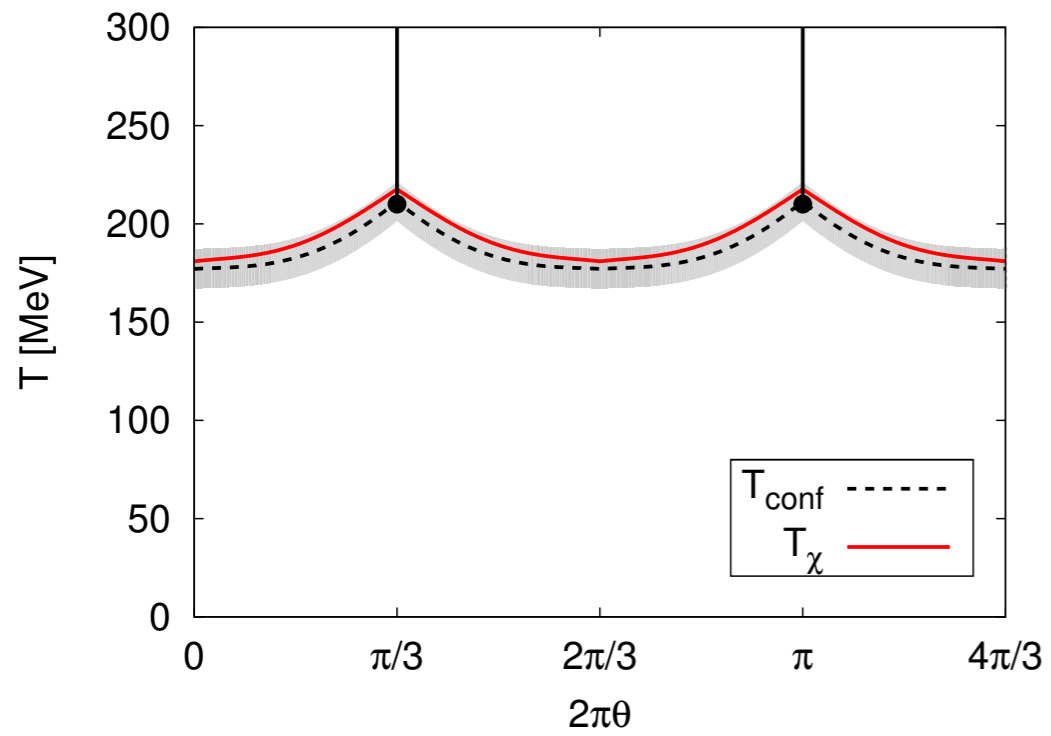
$\langle A_0 \rangle_\theta$



Braun, Haas, Marhauser, JMP '09

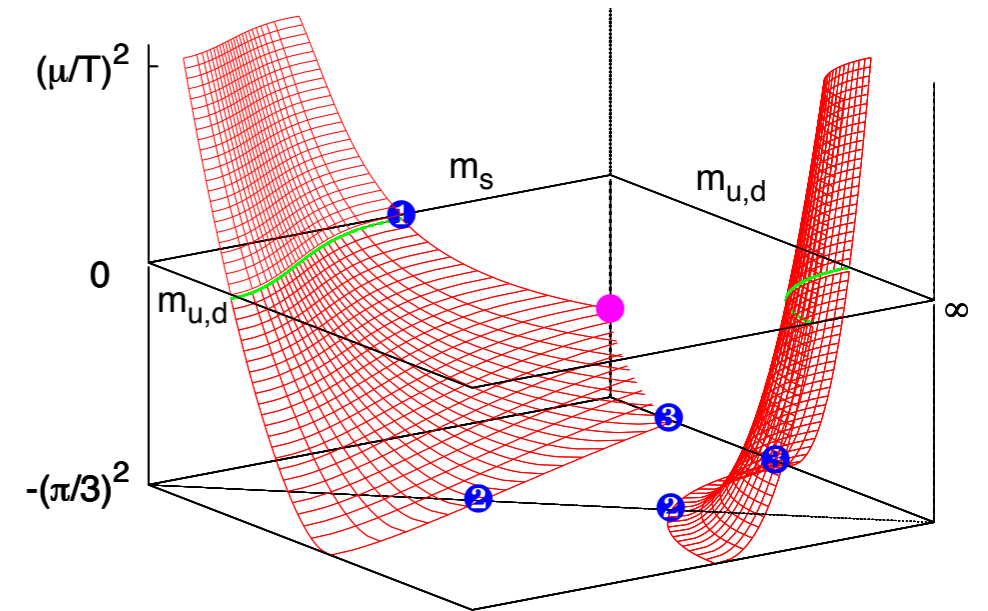
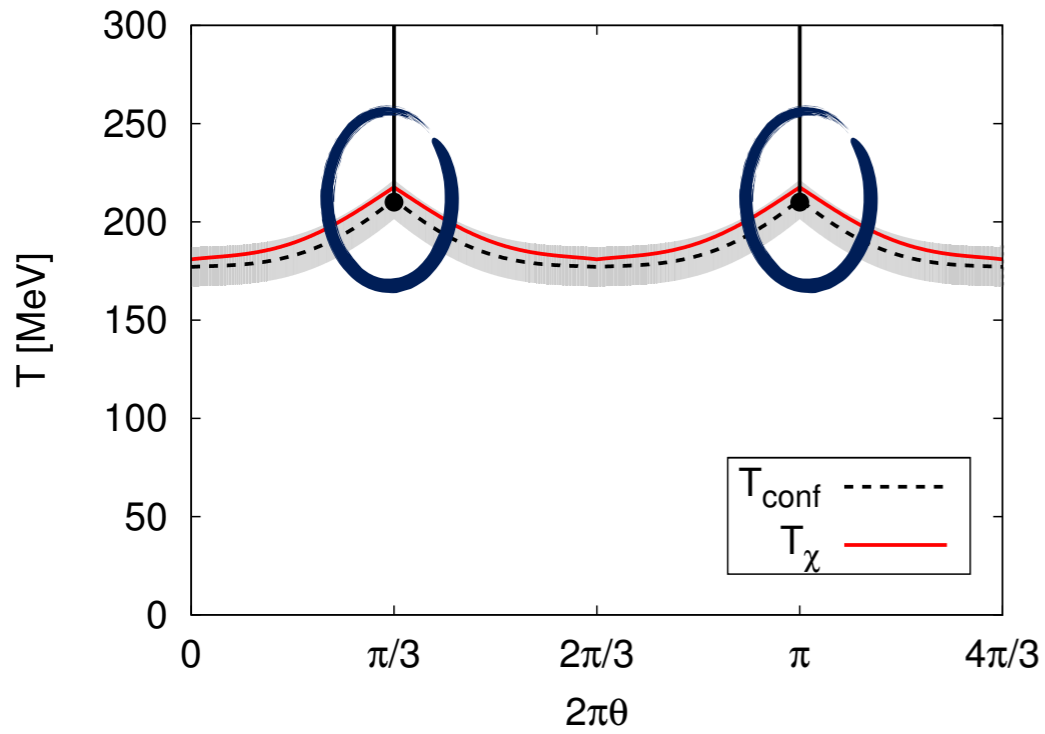
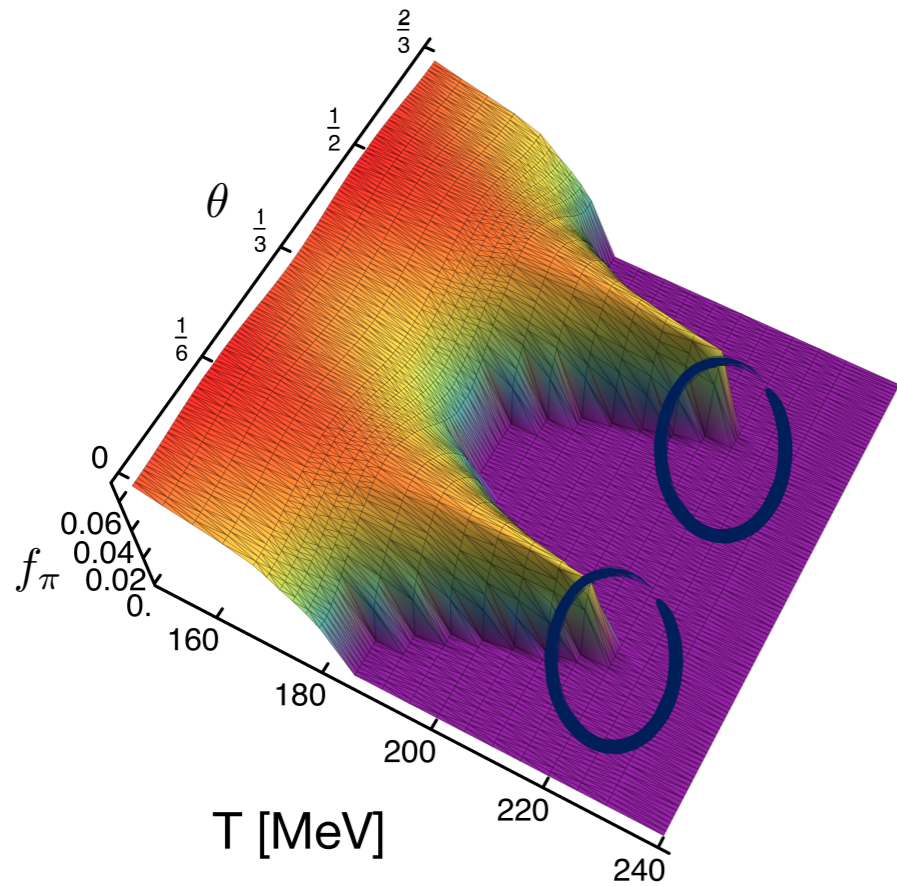


$\langle A_0 \rangle_{\theta=0}$



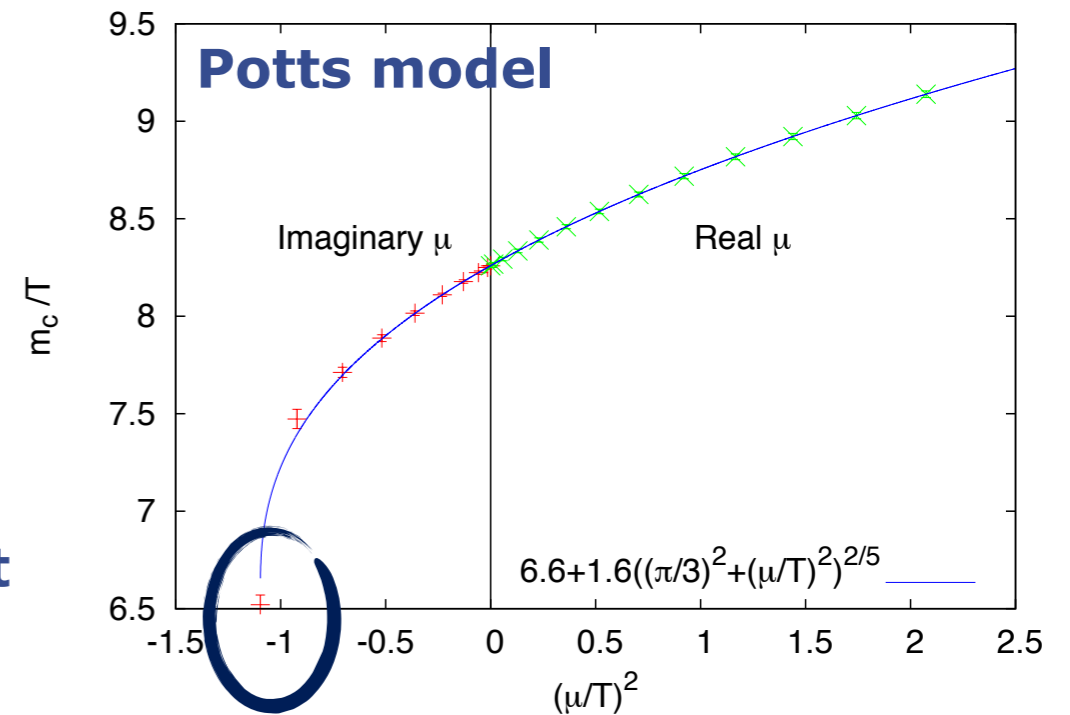
# Imaginary chemical potential

## Nature of the RW endpoint



O. Philipsen '11

RW endpoint

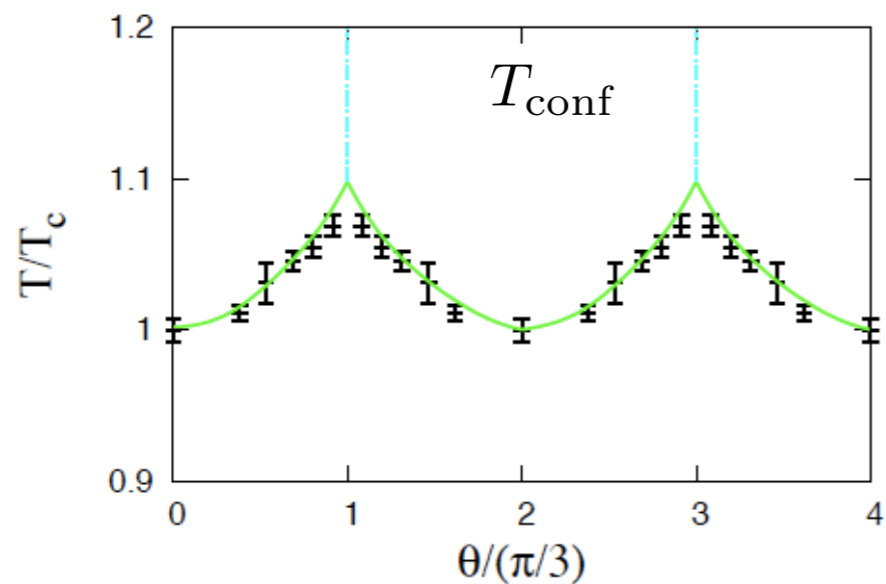
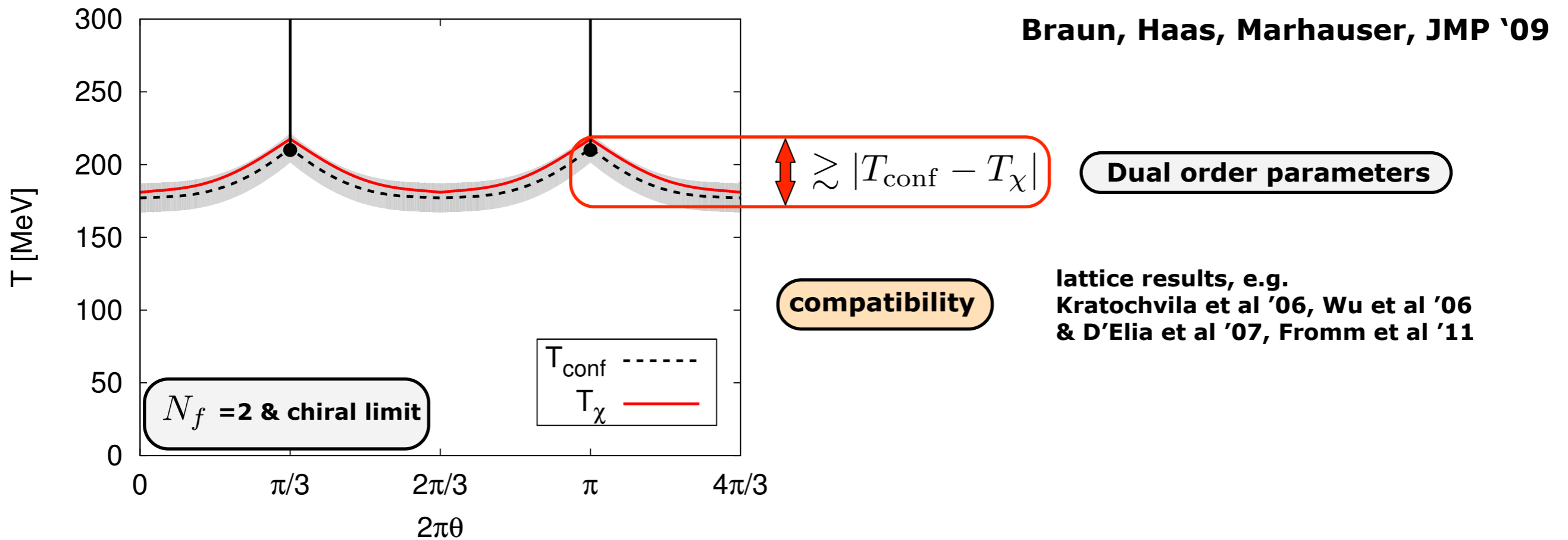


Nature of RW endpoint  
lattice: D'Elia, Sanfilippo '09  
de Forcrand, Philipsen '10

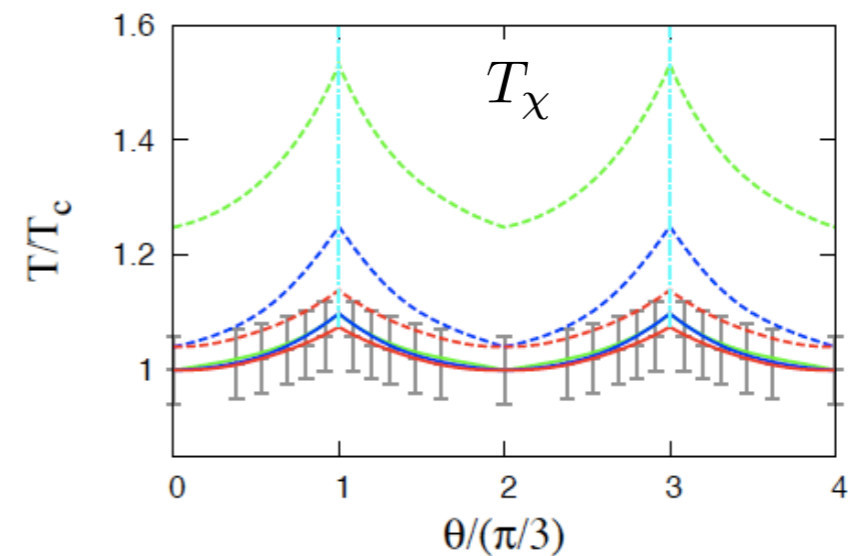
PNJL: Sakai et al '10,  
Morita et al '11

...

# Imaginary chemical potential



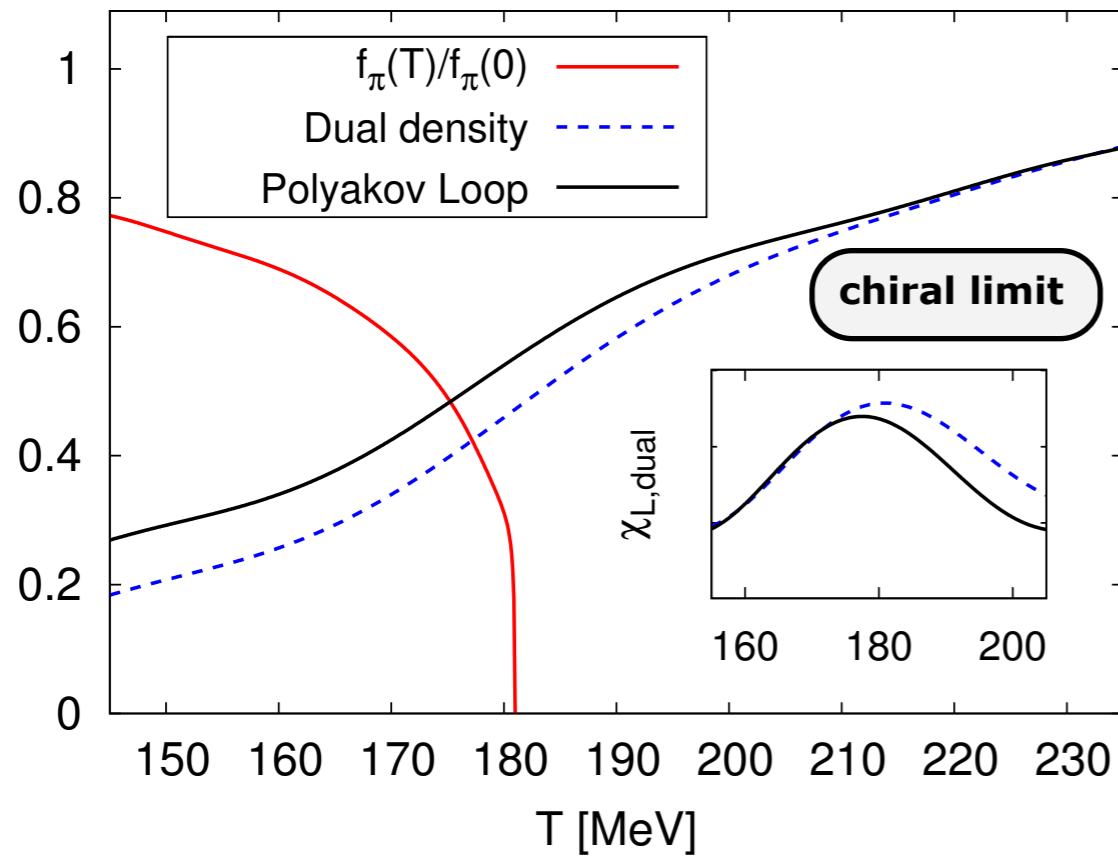
PNJL: Sakai et al '09



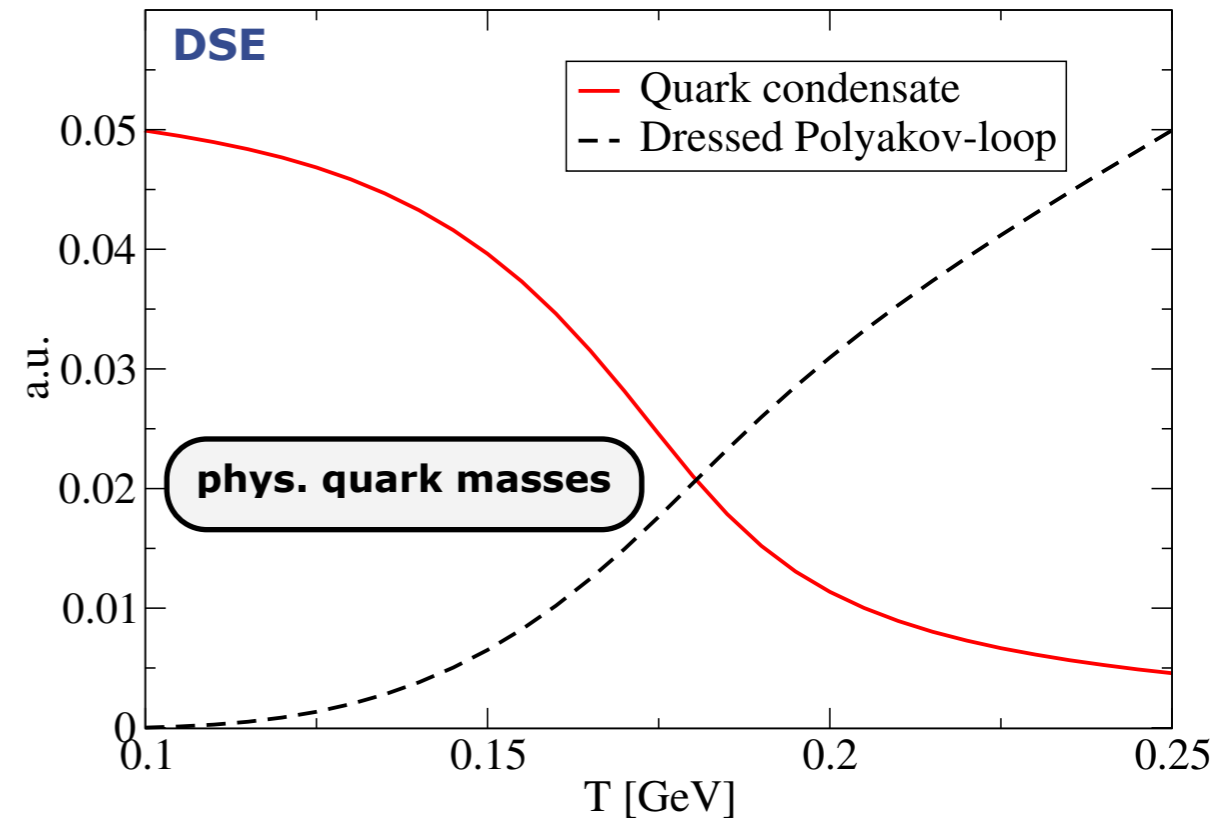
adjust 8-fermi interaction

# Full dynamical QCD: $N_f = 2$ & chiral limit

## Phase structure

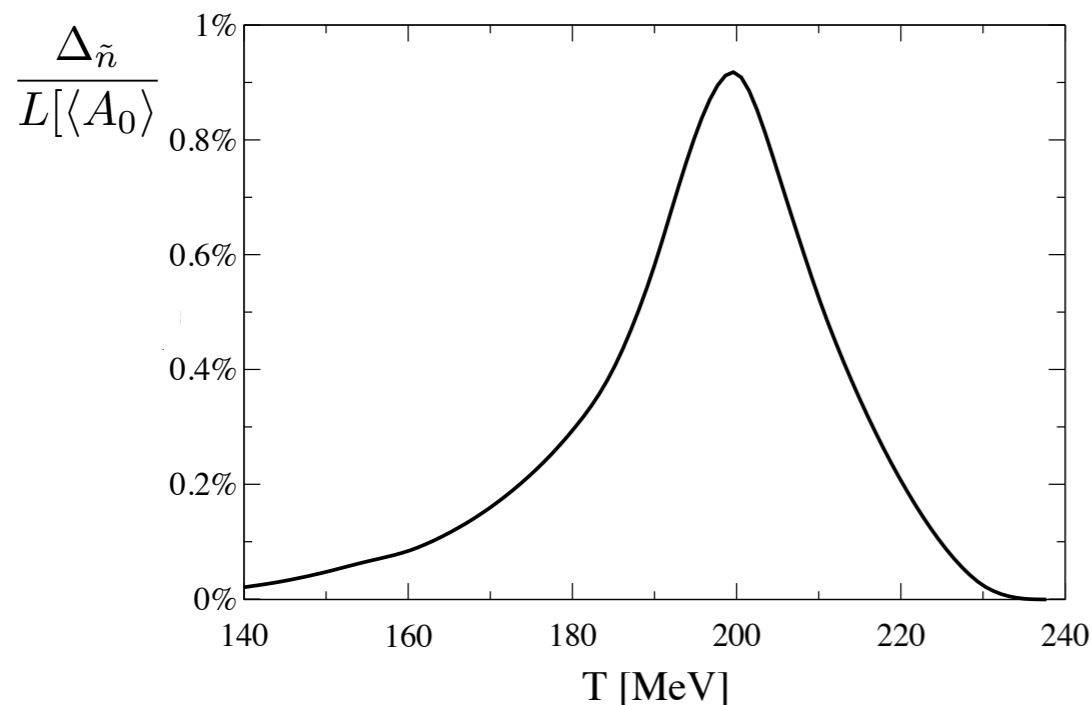


Braun, Haas, Marhauser, JMP '09



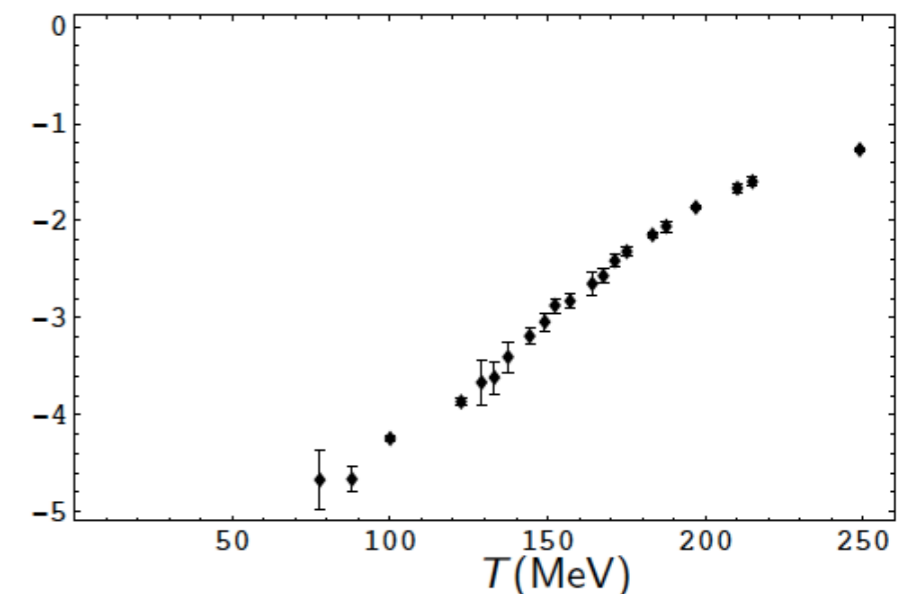
Fischer, Lücker, Müller '11

## factorisation property of dual density



$$\Delta_{\tilde{n}} = \frac{\tilde{n}[\langle A_0 \rangle]}{\tilde{n}[0]} - L[\langle A_0 \rangle]$$

## Log of dual condensate, $m=60$ MeV



Zhang, Bruckmann, Gattringer, Fodor, Szabo '10

# Confinement

## Effective Polyakov loop potential

Learning by diffusion

### Non-perturbative effective potential

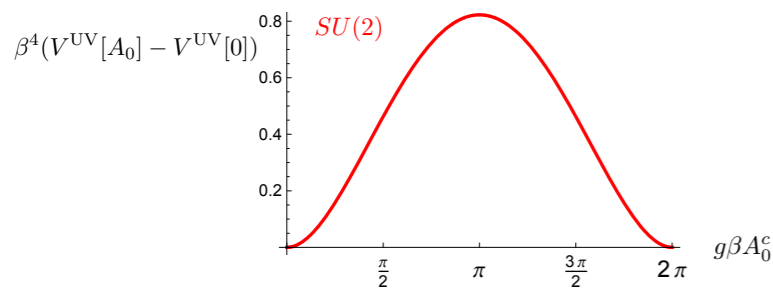
$$V[A_0] \simeq -\frac{1}{2} \text{Tr} \log \langle AA \rangle [A_0] + O(\partial_t \langle AA \rangle) + \text{Tr} \log \langle C\bar{C} \rangle [A_0] + O(\partial_t \langle C\bar{C} \rangle)$$

free energy

### Confinement criterion

Braun, Gies, JMP '07

Fister, JMP '13



$$\beta^4 V^{UV}[A_0] = -2 * 3 \left( \frac{\pi^2}{90} - \frac{2\pi^2}{3} \tilde{\varphi}^2 (1 - \tilde{\varphi})^2 \right)$$

**2** = 2 transversal physical polarisations + 1 transversal (zero mode) + 1 longitudinal - 2 ghosts



# Confinement

## Effective Polyakov loop potential

Learning by diffusion

### Non-perturbative effective potential

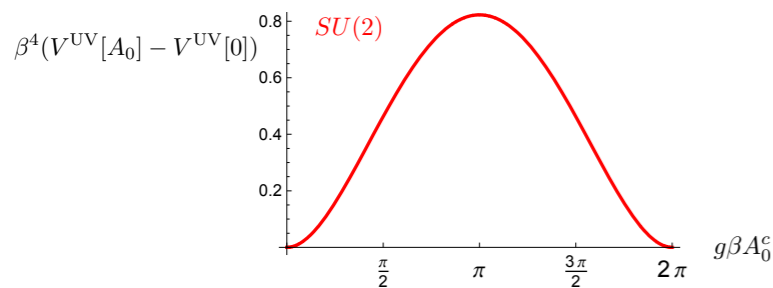
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**Glueon contribution deconfines**

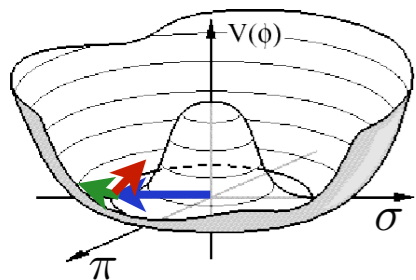
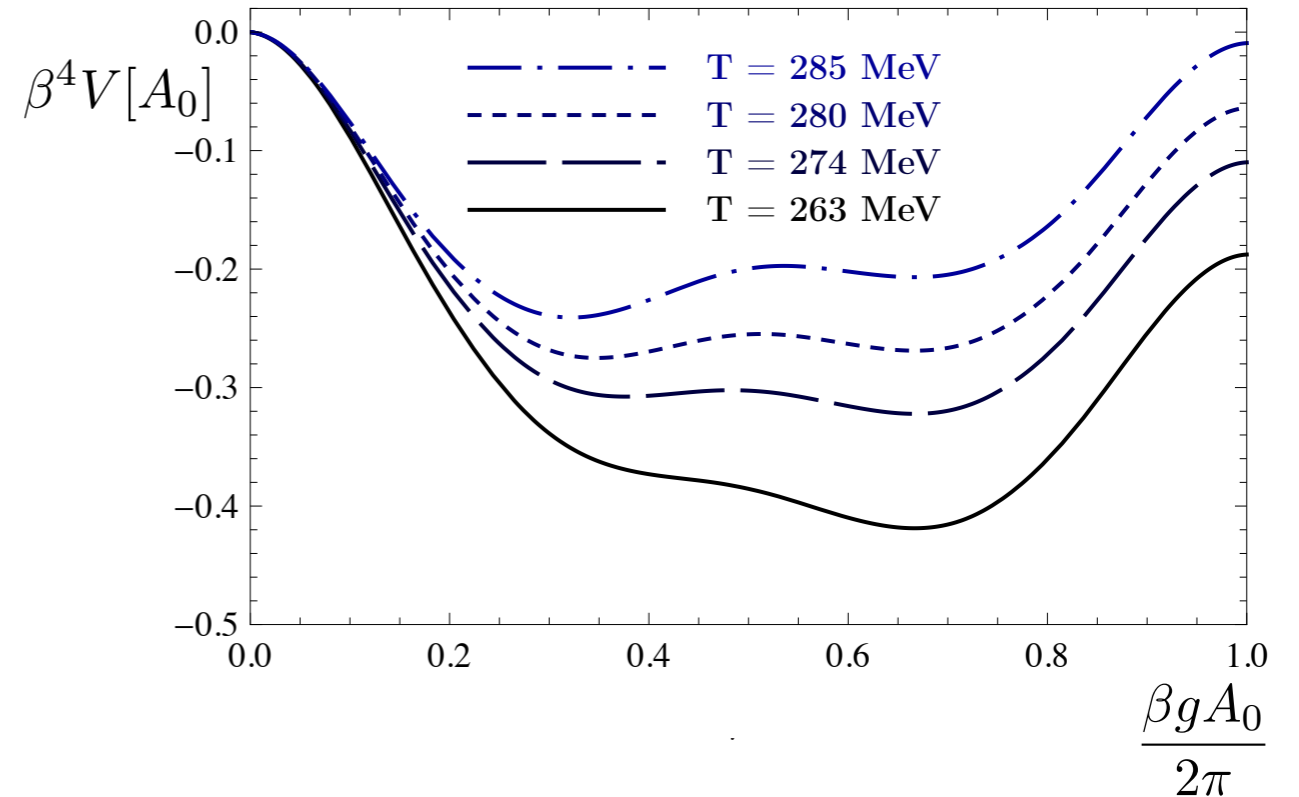
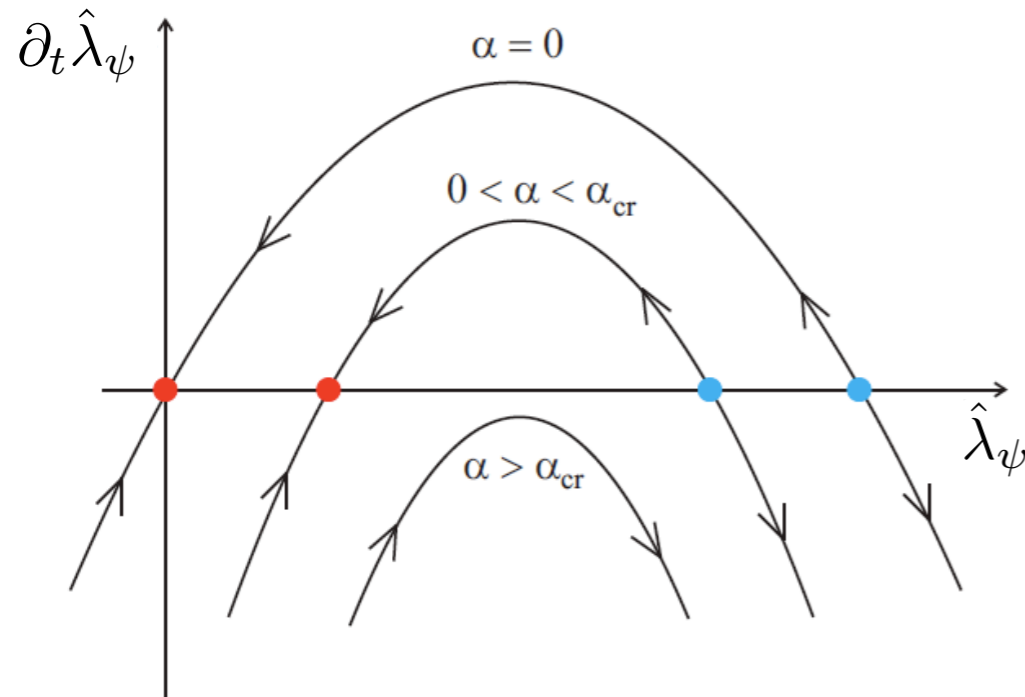
**Ghost contribution confines**

**Confinement**  $\longleftrightarrow$  **suppression of the gluon relative to the ghost**

# Full dynamical QCD: $N_f = 2$ & chiral limit

## Phase structure

Reminder

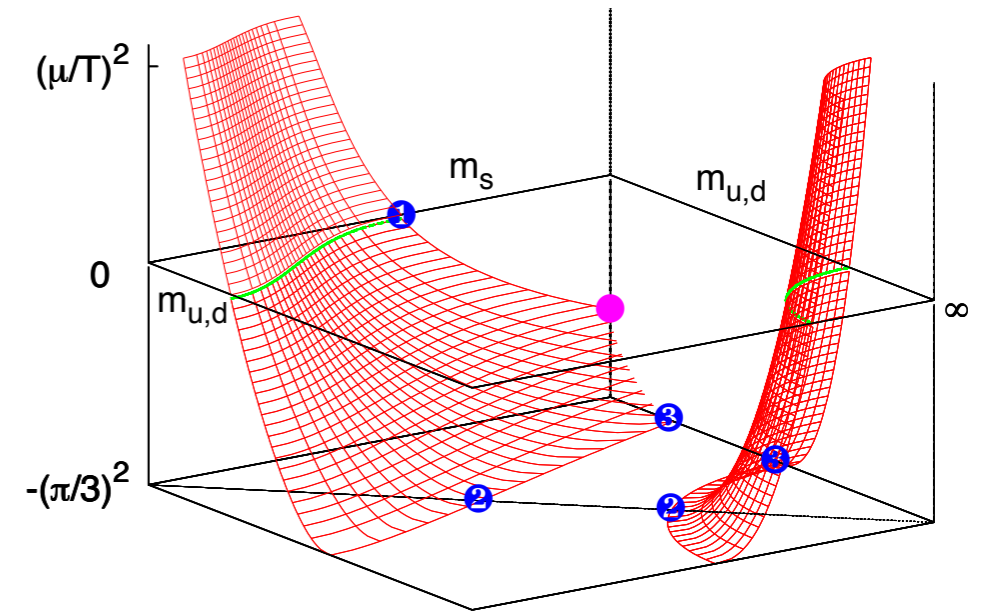
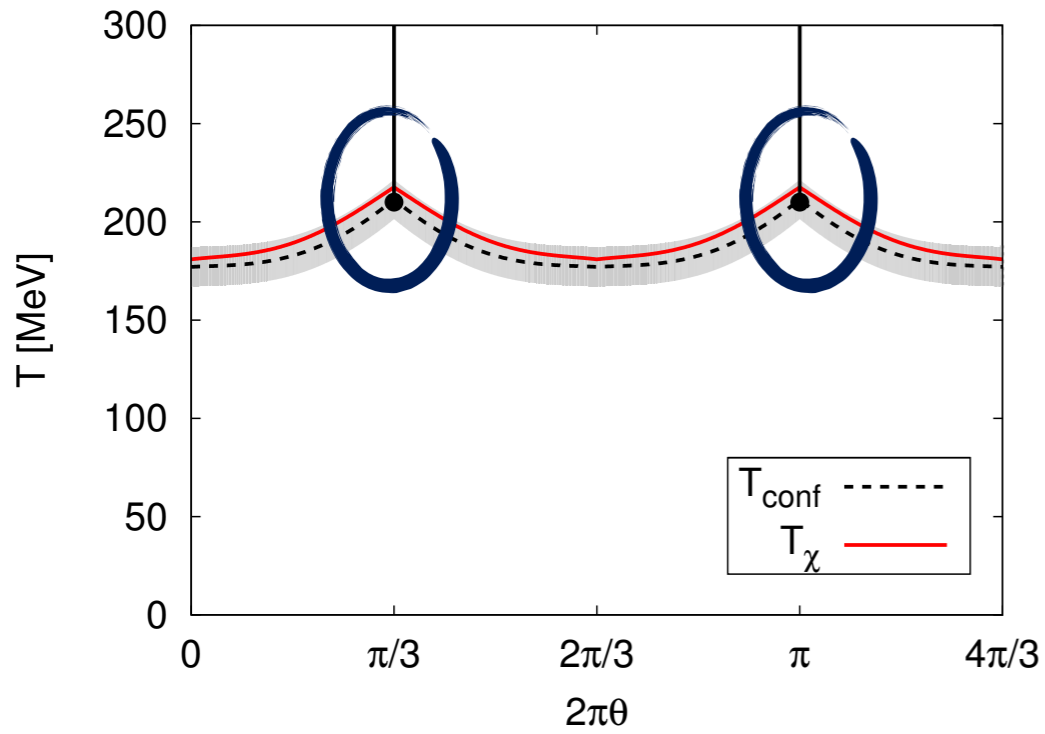
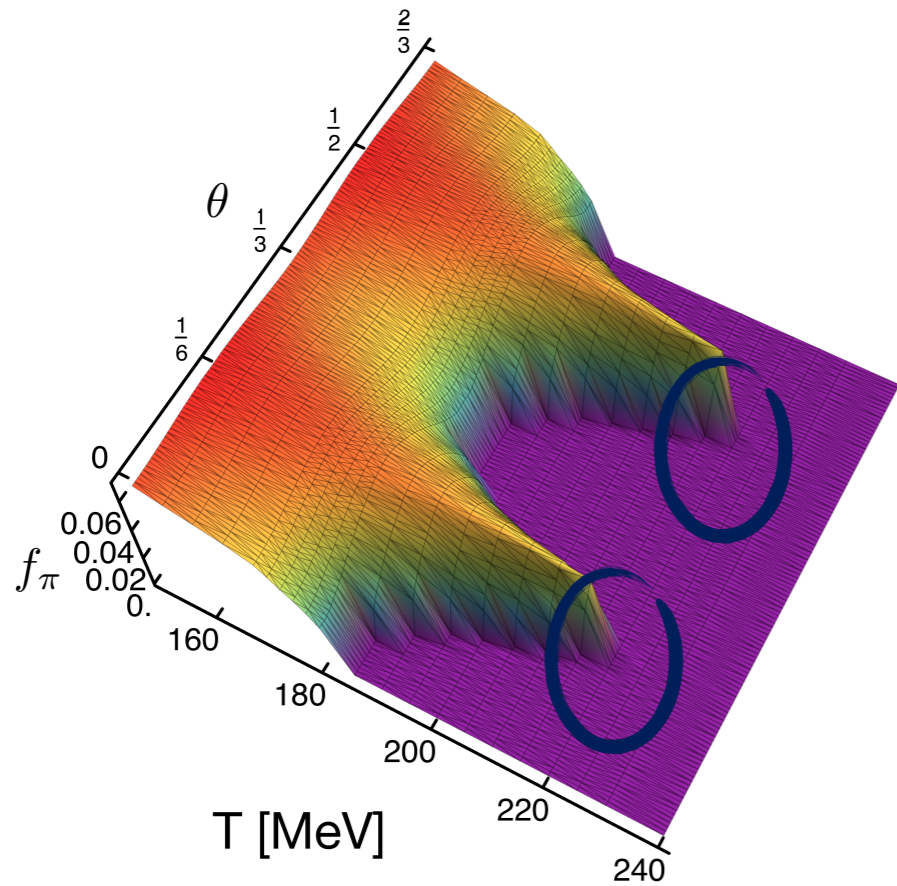


chiral symmetry breaking  $\longleftrightarrow \alpha_s > \alpha_{s,cr}$

Confinement  $\longleftrightarrow$  suppression of the gluon relative to the ghost

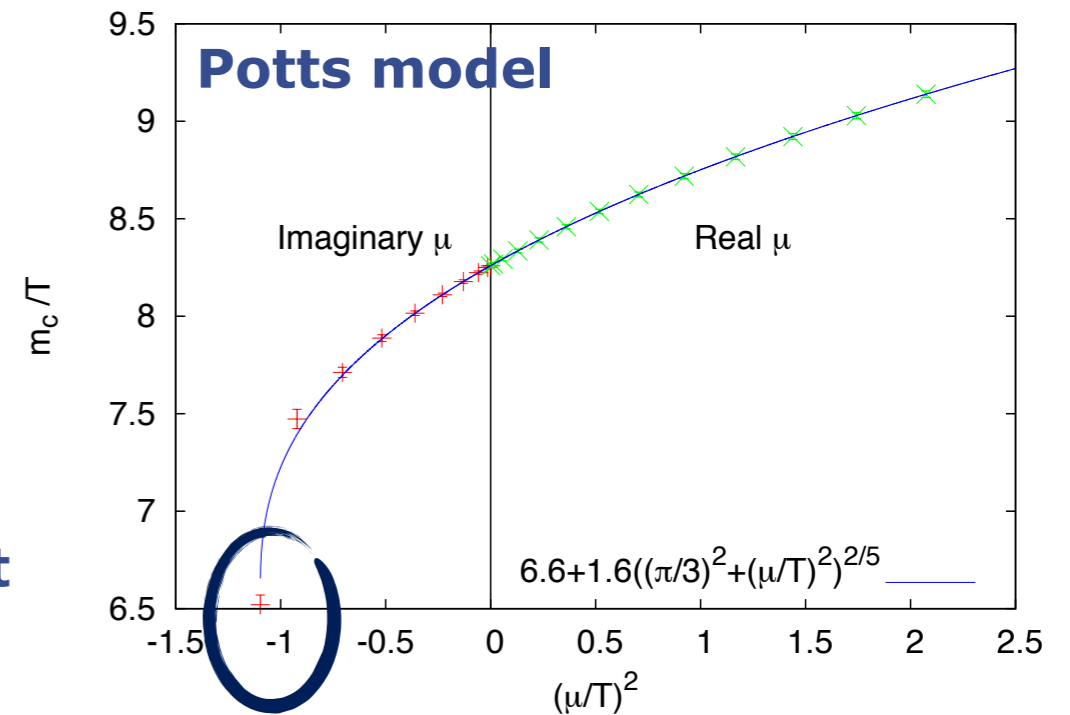
# Imaginary chemical potential

## Nature of the RW endpoint



O. Philipsen '11

RW endpoint



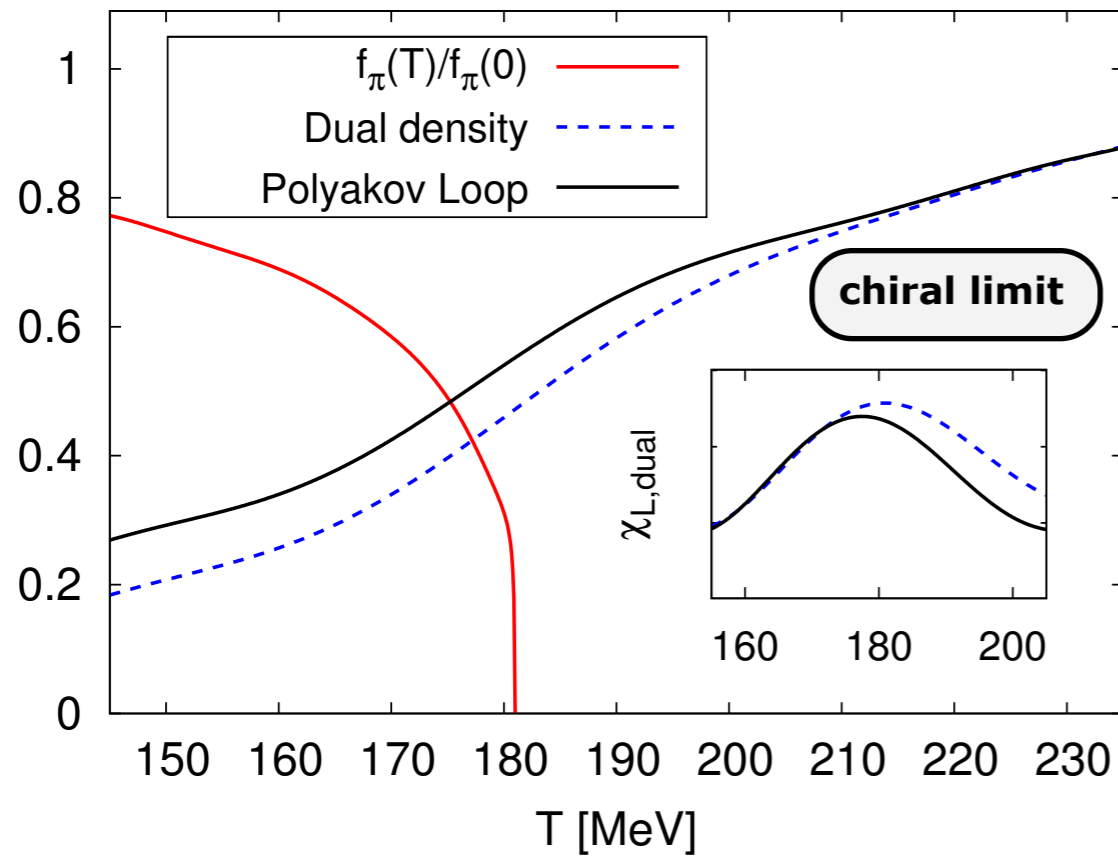
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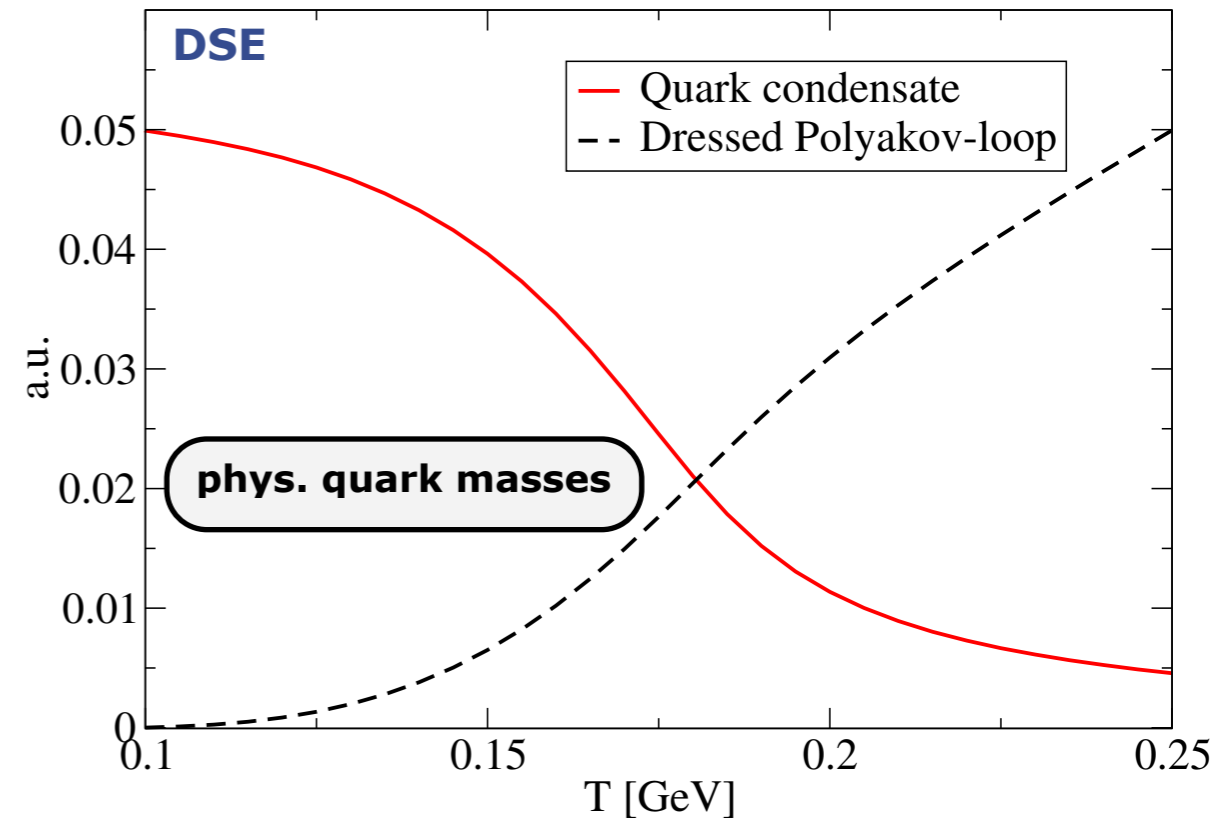
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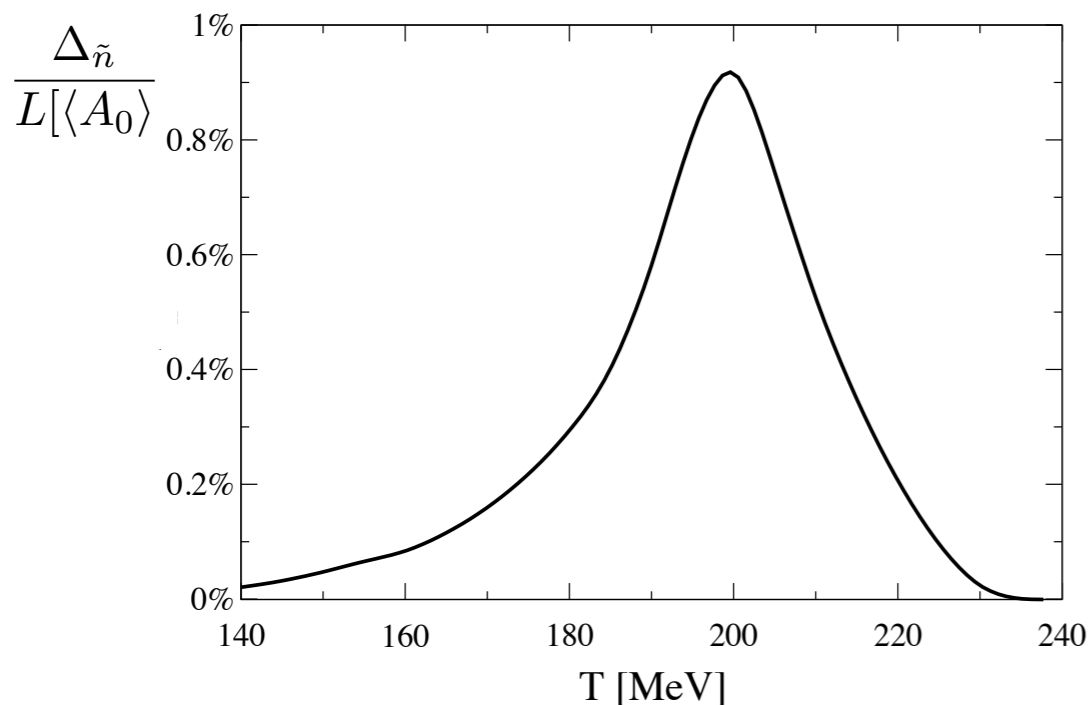


Braun, Haas, Marhauser, JMP '09



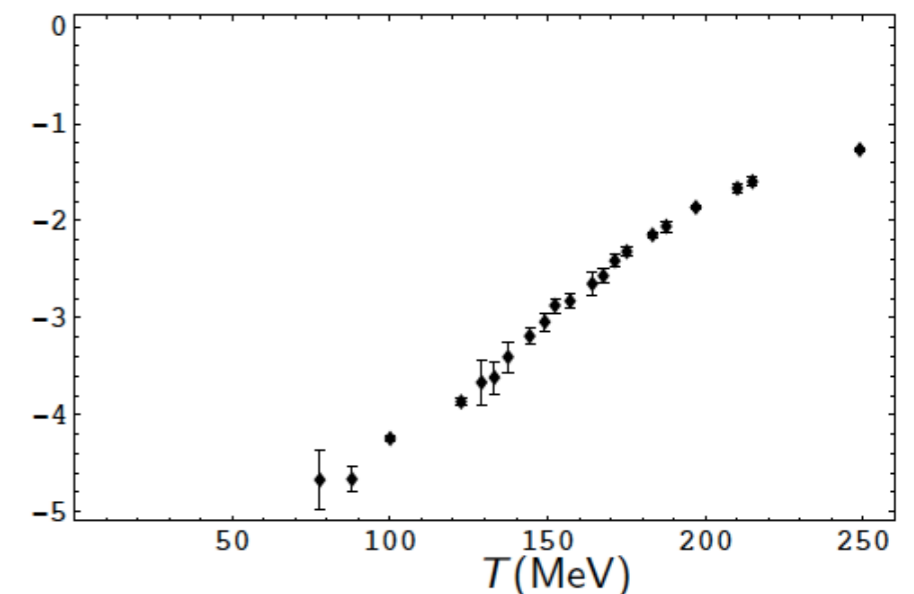
Fischer, Lücker, Müller '11

## factorisation property of dual density



$$\Delta_{\tilde{n}} = \frac{\tilde{n}[\langle A_0 \rangle]}{\tilde{n}[0]} - L[\langle A_0 \rangle]$$

## Log of dual condensate, $m=60$ MeV



Zhang, Bruckmann, Gattringer, Fodor, Szabo '10