

2.3 chiral phase transition

(1) without gluons

lit: Brauer,
diss.

Nambu - Jona-Lasinio model (1 flavour, $N_f=1$)

$$\mathcal{L}[\psi, \bar{\psi}] = \int d^4x \left\{ i \bar{\psi} (\not{\partial} + m) \psi + \frac{\lambda_\psi}{2} \left[(\bar{\psi} \psi)^2 - (\bar{\psi} \gamma_5 \psi)^2 \right] \right\}$$

↑ scalar
↑ pseudo-scalar

non-renormalisable: $[\lambda_\psi] = -2$

Flow of λ_ψ -coupling:

$$\partial_\tau \lambda_\psi = \text{fermion-loop} + \text{ghost-loop}$$

dimensionless coupling: $\hat{\lambda}_\psi = \lambda_\psi \cdot k^2$

$$\dot{\hat{\lambda}}_\psi = \hat{\lambda}_\psi \cdot k^2 + 2 \hat{\lambda}_\psi$$

$$= 2 \hat{\lambda}_\psi + \underbrace{A_k(\tau)}_{<0} \hat{\lambda}^2 + B_k(\tau) \cdot \hat{\lambda}_\phi$$

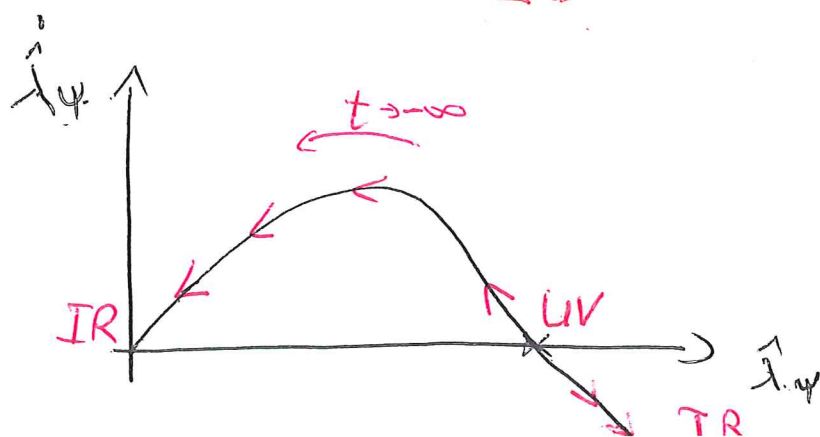


Fig 1

Chiral symmetry breaking:

$S[\psi, \bar{\psi}]$ with $m_q = 0$ is invariant
under chiral rotations

$$\psi \rightarrow e^{i\alpha \gamma_5} \psi, \quad \bar{\psi} \rightarrow \bar{\psi} e^{i\alpha \gamma_5}$$

see p. II-79.

Phases:

$$\langle \bar{\psi} \psi \rangle = 0 \quad \mathcal{N}\text{-symmetric}$$

$$\langle \bar{\psi} \psi \rangle \neq 0 \quad \mathcal{N}\text{-broken}$$

The chiral condensate $\langle \bar{\psi} \psi \rangle$ is an order parameter for spont. chiral symmetry breaking; it is sensitive to chiral symmetry, $\bar{\psi} \psi \xrightarrow{e^{i\alpha \gamma_5}} \bar{\psi} e^{2i\alpha \gamma_5} \psi$ and measures the dynamical quark mass.

Eg. 1-loop:

$$\begin{aligned} \int_x \langle \bar{\psi} \psi \rangle &\simeq -\text{Tr} \langle \psi(x) \bar{\psi}(x) \rangle \simeq -\int \frac{d^4 p}{(2\pi)^4} \text{tr} \frac{1}{\not{p} - im} \\ &\simeq -\int \frac{d^4 p}{(2\pi)^4} \text{tr} \frac{\not{p} + im}{p^2 + m^2} \simeq 4im \int d^4 p \frac{1}{p^2 + m^2} \end{aligned}$$

chiral condensate requires renormalisation

chiral symmetry: $\psi = \begin{pmatrix} \psi_L \\ \psi_R \end{pmatrix}$, $\gamma_5 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

$$\Rightarrow \begin{pmatrix} \psi_L \\ 0 \end{pmatrix} = \frac{1+\gamma_5}{2} \psi, \quad \begin{pmatrix} 0 \\ \psi_R \end{pmatrix} = \frac{1-\gamma_5}{2} \psi$$

chiral transformation:

$$\psi \rightarrow e^{i\alpha\gamma_5} \psi, \quad \bar{\psi} \rightarrow e^{+i\alpha\gamma_5} \bar{\psi}$$

$$\psi_R \rightarrow e^{-i\alpha} \psi_R, \quad \psi_L \rightarrow e^{i\alpha} \psi_L$$

with

$$\{\gamma_5, \gamma_\mu\} = 0$$

$$\{\gamma_5, \gamma_\mu\} = 0 \longrightarrow$$

$$\bar{\psi} = \psi^\dagger \gamma_0 \rightarrow \psi^\dagger e^{-i\alpha\gamma_5} \gamma_0 = \psi^\dagger \gamma_0 e^{i\alpha\gamma_5}$$

It follows that

$$(a) \quad i\bar{\psi} \not{\partial} \psi \rightarrow i\bar{\psi} e^{i\alpha\gamma_5} \not{\partial} e^{i\alpha\gamma_5} \psi = \bar{\psi} \underbrace{e^{i\alpha\gamma_5} e^{-i\alpha\gamma_5}}_1 \not{\partial} \psi$$

$$(b) \quad i\bar{\psi} \not{m} \psi \rightarrow i\bar{\psi} e^{2i\alpha\gamma_5} \psi$$

and $(\bar{\psi} \psi)^2 - (\bar{\psi} \gamma_5 \psi)^2 = (\psi_R^\dagger \psi_L + \psi_L^\dagger \psi_R)^2 - (\psi_R^\dagger \psi_L - \psi_L^\dagger \psi_R)^2$

$$\gamma_0 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

chiral rep.

$$\longrightarrow = 2(\psi_R^\dagger \psi_L)(\psi_L^\dagger \psi_R)$$

$$(c) \Rightarrow (\bar{\psi} \psi)^2 - (\bar{\psi} \gamma_5 \psi)^2 \longrightarrow (\bar{\psi} \psi)^2 - (\bar{\psi} \gamma_5 \psi)^2$$

$$\psi_R^\dagger \psi_L \rightarrow \psi_R^\dagger e^{-i\alpha} e^{i\alpha} \psi_L = \psi_R^\dagger \psi_L, \quad \psi_L^\dagger \psi_R \rightarrow \psi_L^\dagger \psi_R$$

Bosonisation: "Introducing the order parameter as a field"

σ -meson: $\sigma \sim \bar{\psi} \psi$

ρ -meson: $\pi \sim \bar{\psi} \gamma_5 \psi$

Hubbard-Stratonovich

$$S_{QM}[\psi, \bar{\psi}, \sigma, \pi] = \int d^4x \left\{ i \bar{\psi} (\not{\partial} + m) \psi + i h \bar{\psi} (\sigma + i \gamma_5 \pi) \psi + \frac{m_\sigma^2}{2} (\sigma^2 + \pi^2) \right.$$

kinetic term $+ \frac{1}{2} \int d^4x \left\{ \sigma (-\partial_\mu^2) \sigma + \pi (-\partial_\mu^2) \pi \right.$

We have $S_{QM}[\psi, \bar{\psi}, \sigma, \pi] \Big|_{EOM(\sigma, \pi)} = S_{NJL}[\psi, \bar{\psi}]$

EOM: $\frac{\partial S_{QM}}{\partial \sigma} = 0 : \quad \sigma = - \frac{i h}{m_\sigma^2} \bar{\psi} \psi$

const. fields

$\frac{\partial S_{QM}}{\partial \pi} = 0 : \quad \pi = \frac{h}{m_\pi^2} \bar{\psi} \gamma_5 \psi$

with

$$\lambda_\psi = \frac{h^2}{m_\sigma^2}$$

Integrate out fermions $\rightarrow$ bosonic theory

with $O(2)$ -symmetry.
 $\stackrel{is}{\sim} U(1)$



Sym. breaking: $\langle \sigma \rangle = 0$ $\mathbb{Z}$ -sym.

$\langle \sigma \rangle \neq 0$ $\mathbb{Z}$ -sym. breaking

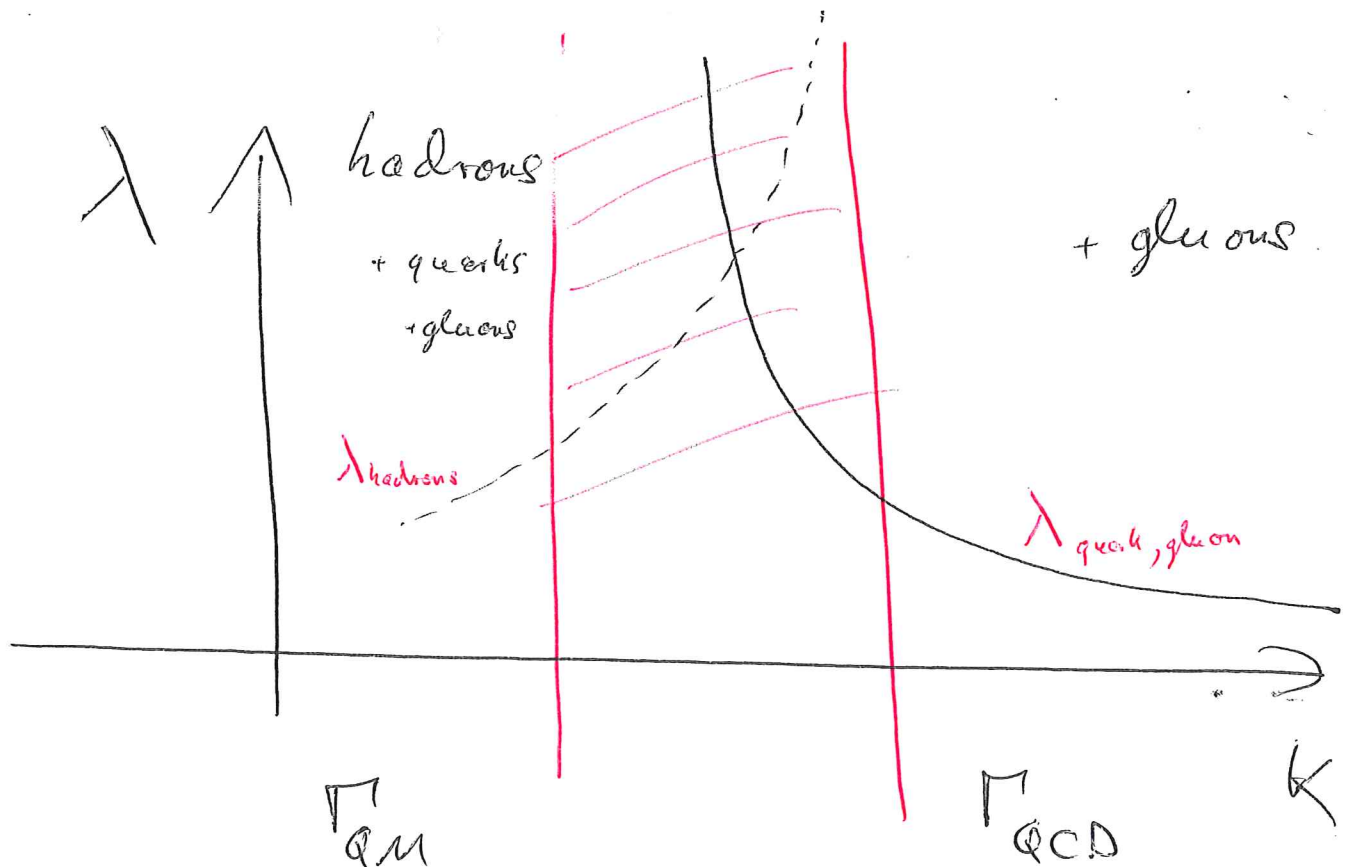
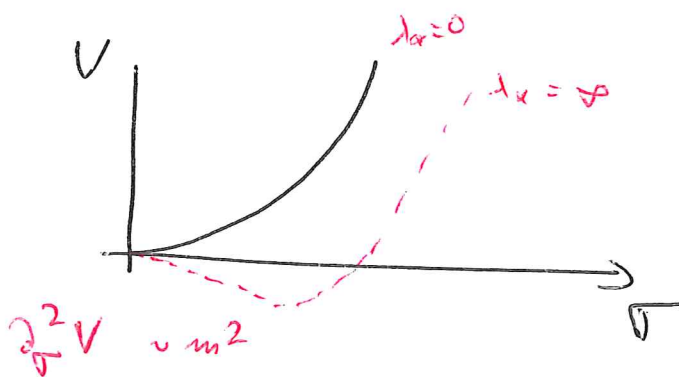
Fixed points :

IR : $\lambda_* = 0$

$\lambda_* = \infty$ $\leftarrow$

UV : $\lambda_* = \lambda_*(A)$

$$m_\sigma^2 \sim \frac{\hbar^2}{\lambda \varphi} \sim \begin{cases} 0 & \lambda_* = \infty \\ \infty & \lambda_* = 0 \end{cases}$$



chiral symmetry breaking:

$$\langle \bar{\psi} \psi \rangle \neq 0 \quad \therefore \sigma = 2\lambda \psi \bar{\psi} \quad \langle \sigma \rangle = -2\lambda \langle \bar{\psi} \psi \rangle$$

$$\pi = 2i\lambda \sigma \bar{\psi} \gamma_5 \psi \quad \langle \pi \rangle = -2i\lambda \langle \bar{\psi} \gamma_5 \psi \rangle$$

(part.) Bosonisation: (Hubbard-Stratonovich)

$$S_{QM}[\psi, \bar{\psi}, \sigma, \pi] = \int d^d x \left\{ \bar{\psi} i \not{\partial} \psi + i h_\sigma \bar{\psi} (\sigma + i \gamma_5 \pi) \psi - \frac{h_\sigma^2}{4\lambda \sigma} (\sigma^2 + \pi^2) \right\}$$

+ kinetic term

$$+ \frac{1}{2} \int d^d x \left\{ \sigma (-\partial^2) \sigma + \pi (-\partial^2) \pi \right\}$$

in QCD iso-vector
($\pi^0, \pi^\pm$)

$$EoM: \frac{\partial S}{\partial \sigma} = 0: \quad \sigma = 2\lambda \sigma \bar{\psi} \psi / h_\sigma$$

$$\frac{\partial S}{\partial \pi} = 0: \quad \pi = 2i\lambda \sigma \bar{\psi} \gamma_5 \psi / h_\sigma$$

$\Rightarrow$ on-shell

$$S_{QM} = S_{NJL}$$

explicit
symbol

$$- \frac{1}{4} \lambda (\sigma^2 + \pi^2 - V^2)^2 - c \sigma$$

Integrate-out fermions: (non-) linear σ -model

Flow of λ_5 -coupling: (schematically)

$$2\lambda_5 = \text{diagram 1} + \text{diagram 2} + \text{diagram 3} + \text{diagram 4} + \dots$$

(drop & insertions)

$$\dot{\lambda}_5 = 2\lambda_5 - A_4 \lambda_5^2 + \dots \Rightarrow \text{Fig 1 page}$$

1) Warning: potential overcounting

2) Renormalisation: page II-80a
QCD & (1 flavour)

$$S_{\text{QCD}}[\phi, \psi, \bar{\psi}] = S[\phi] + \int d^d x \bar{\psi} i \not{D} \psi$$

$$\text{with } i \not{D} = i \not{\partial} + i g A$$

$$\text{diagram 1} + \text{diagram 2} + \text{diagram 3} + \dots$$

$$\Rightarrow \Gamma = S_{\text{QCD}} + \int d^d x \lambda_5 [(\bar{\psi} \psi)^2 - (\bar{\psi} \gamma_5 \psi)^2] + \dots$$

Re bosonisation?

Idea:

Gies, Wetzel '01

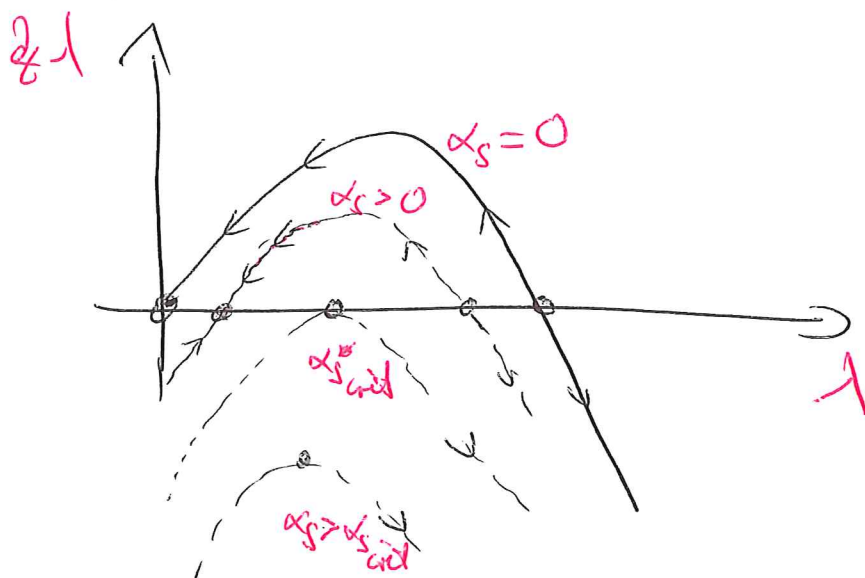
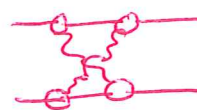
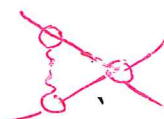
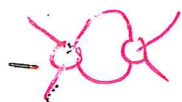
bosonise at each flow-step, e.g.

$$(\lambda_{\sigma_k}, h_{\sigma_k}, m_{\sigma_k}, \lambda_k) \rightarrow (\lambda_{\sigma}, \bar{h}_{\sigma_k}, \bar{m}_{\sigma_k}, \bar{\lambda}_k)$$

- absorb running of λ_{σ} in running of Yukawa coupling h_{σ} , mass m_{σ} , ϕ^4 -coupl. λ .
- moves bound state/condensate fractions in λ_{σ} to bosonic dofs.

Flow of $\hat{\lambda}_4$:

$$\hat{\lambda}_4 = 2 \hat{\lambda}_4 + \underbrace{A_4(\tau)}_{<0} \hat{\lambda}^2 + \underbrace{B_4(\tau)}_{<0} \hat{\lambda} \alpha_s + \underbrace{C_4(\tau)}_{<0} \alpha_s^2 + \dots$$



FRG: $\boxed{\alpha_{s_{\text{crit}}} \sim 0.85}$ ($3=N_f$), $T=0$ Gies, Jäkel '08
 $\equiv N_c$

Landau gauge: $\alpha_s^* \sim 3$

$T > 0$: $\boxed{\alpha_{s_{\text{crit}}}(T) > \alpha_{s_{\text{crit}}}(0)}$

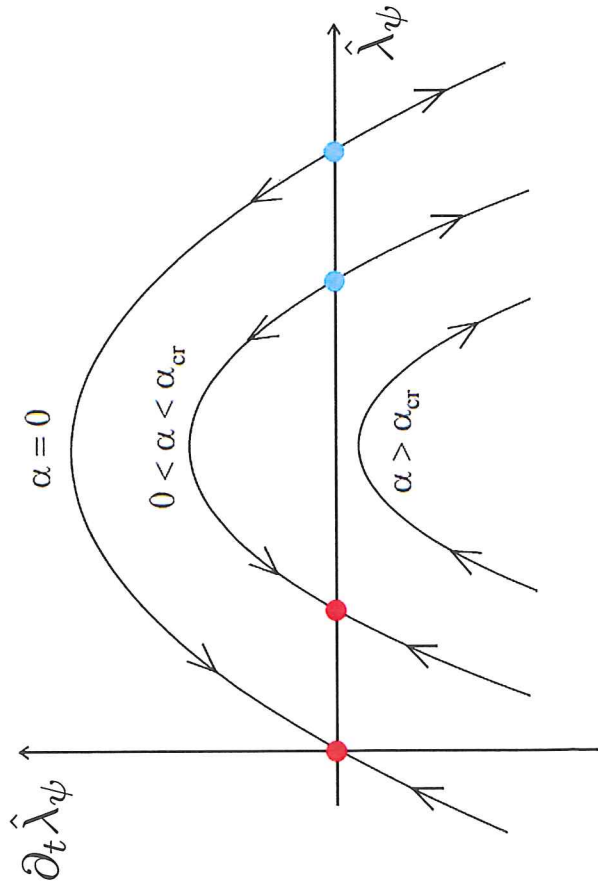
Remark: low energy QCD successfully modelled with
 Quark-Meson model

Chiral symmetry breaking

A glimpse at chiral symmetry breaking in QCD within the FRG

Flow for four-fermion coupling $\hat{\lambda}_\psi = \lambda_\psi k^2$ with infrared scale k

$$k \partial_k \hat{\lambda}_\psi = 2\hat{\lambda}_\psi + A\left(\frac{T}{k}\right)\hat{\lambda}_\psi^2 + B\left(\frac{T}{k}\right)\hat{\lambda}_\psi \alpha_s + C\left(\frac{T}{k}\right)\alpha_s^2 + \dots$$

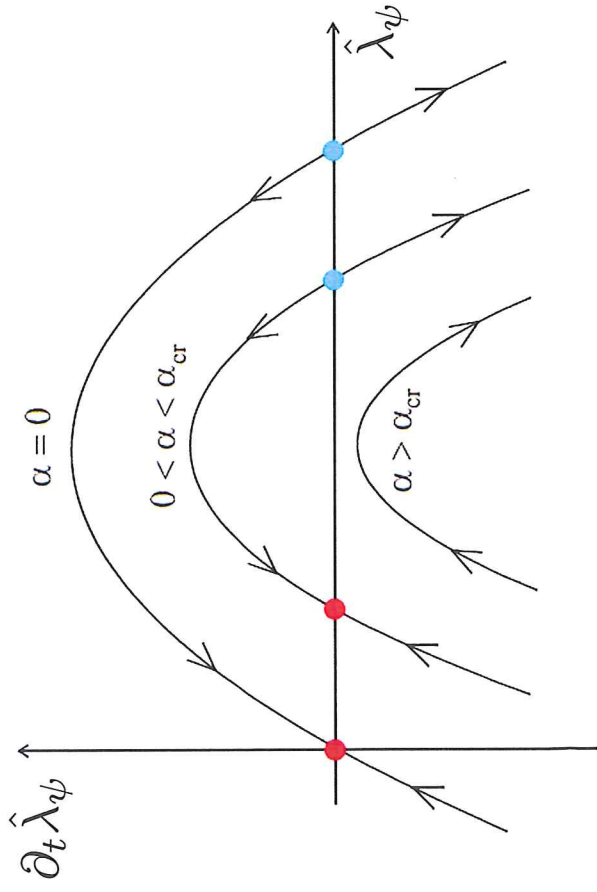
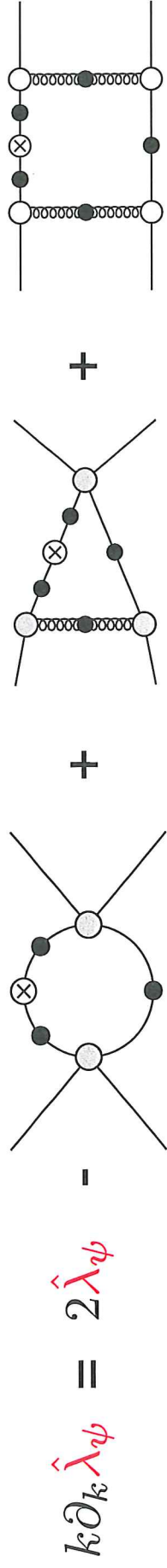


$$T=0: \alpha_s \sim 1$$

$$T>0: \alpha_s(T) > \alpha_s(0)$$

Dynamical hadronisation

Flow for four-fermion coupling $\hat{\lambda}_\psi = \lambda_\psi k^2$ with infrared scale k

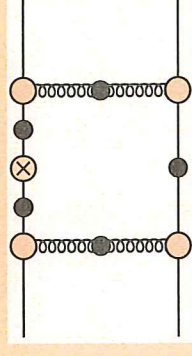
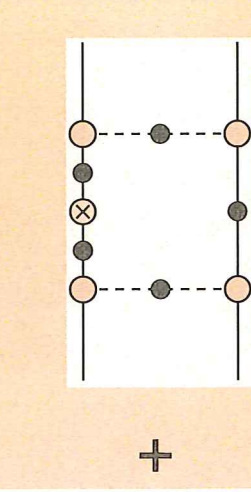
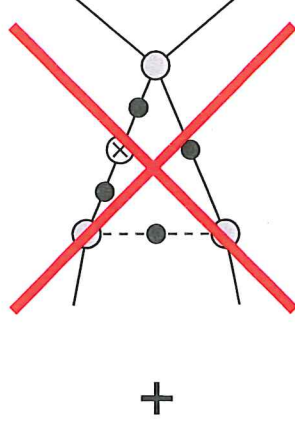
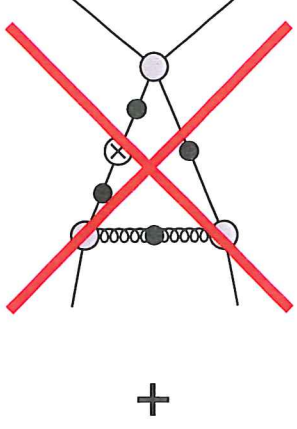
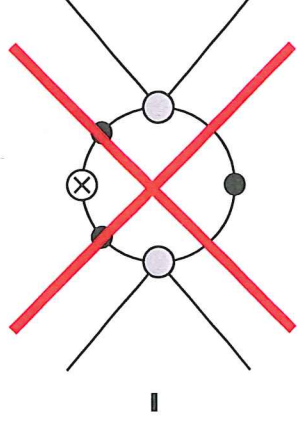
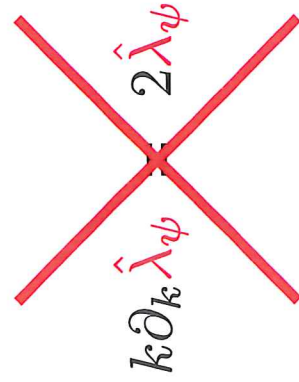


h^2
 $\partial_t \frac{h^2}{m^2}$ - terms

+ ...

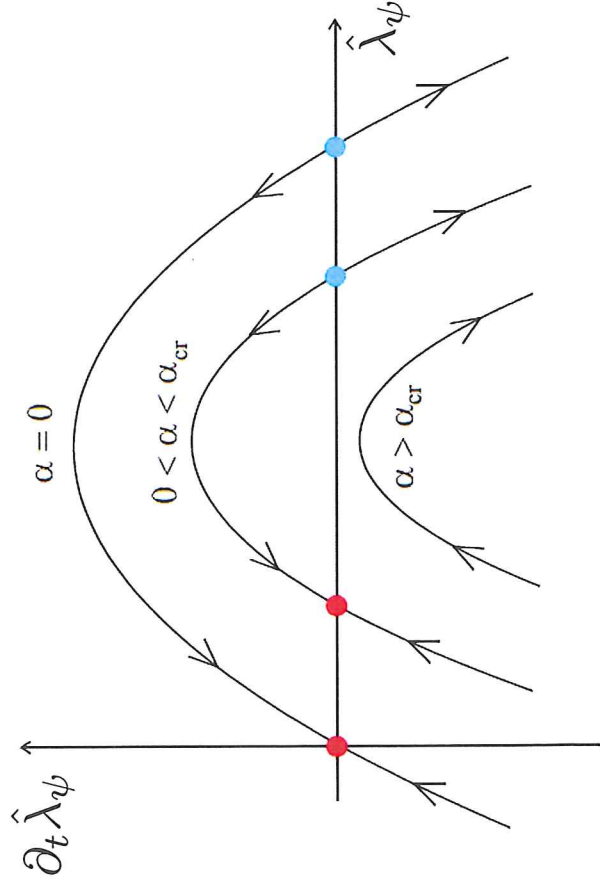
Dynamical hadronisation

Full bosonisation $\lambda_\psi = 0$



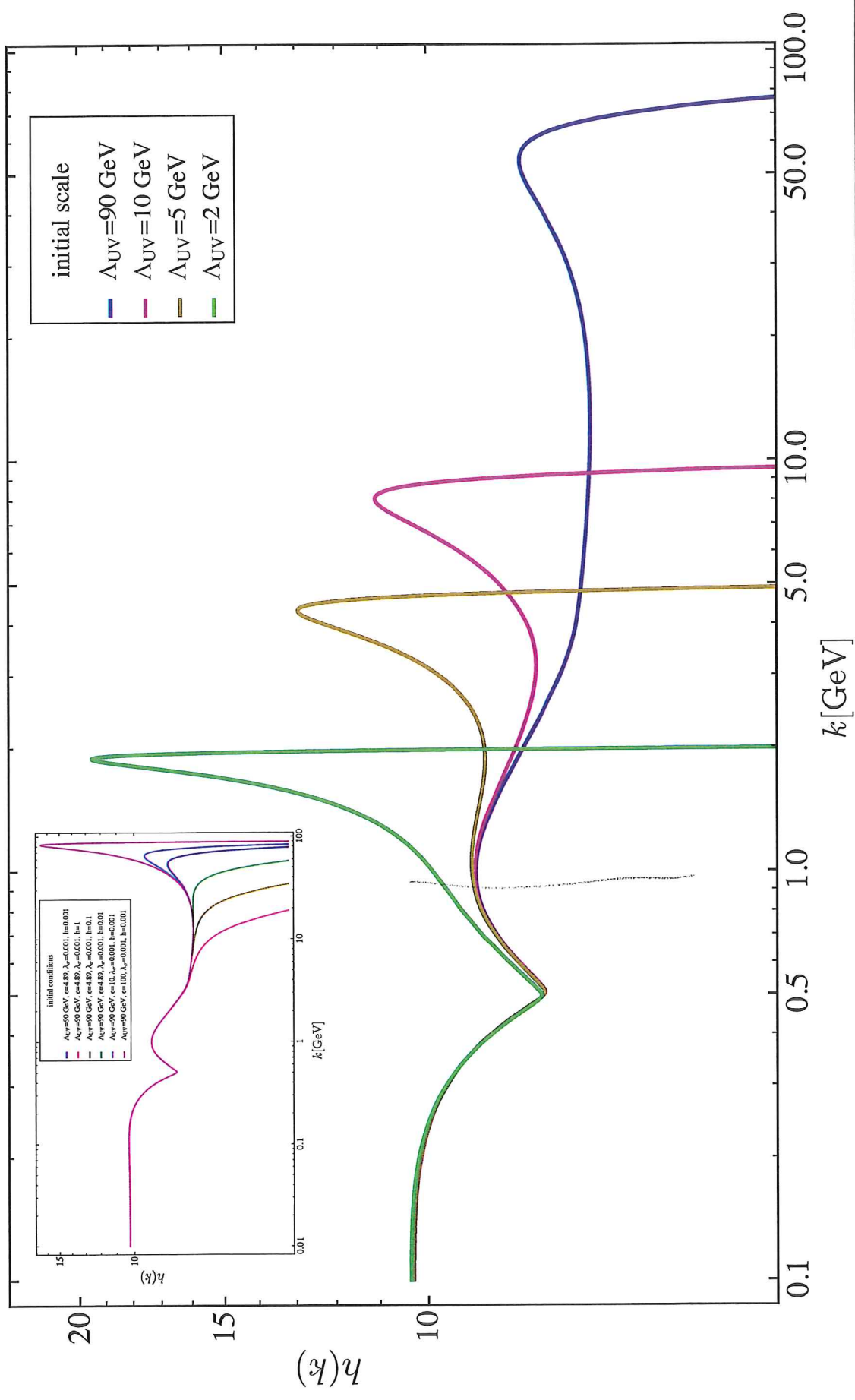
$$+ \partial_t^2 \frac{\hbar^2}{m^2} \text{ - terms}$$

+



Dynamical hadronisation

Full bosonisation $\hat{\lambda}_\psi = 0$



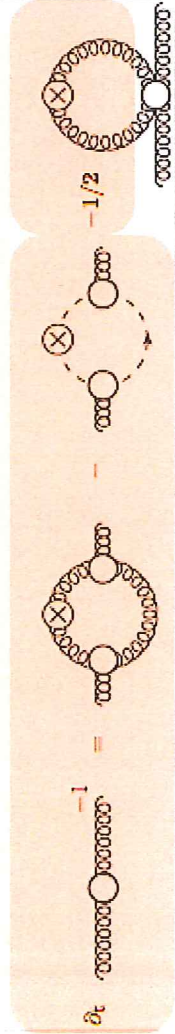
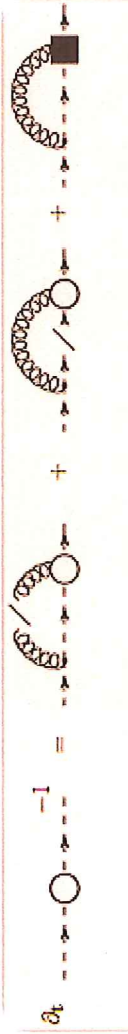
Approximation scheme



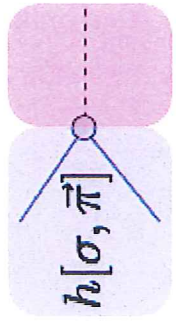
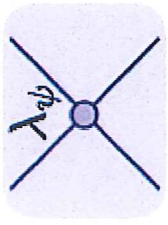
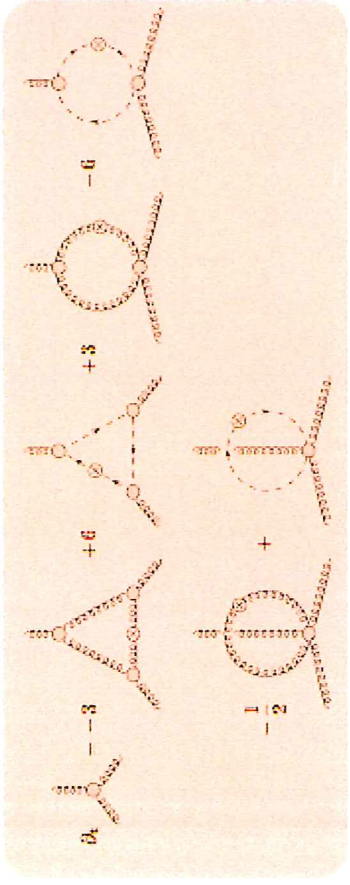
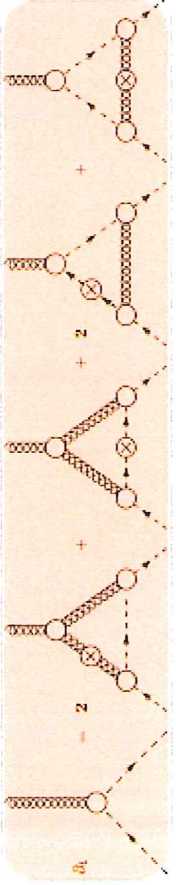
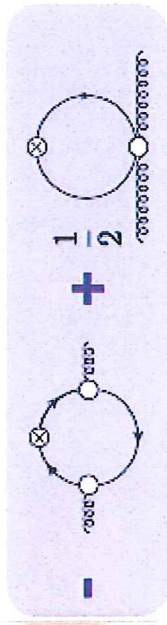
Yang-Mills

DSE

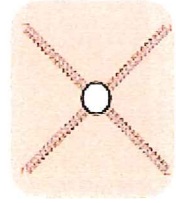
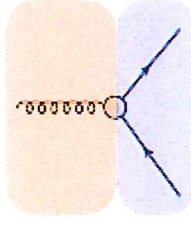
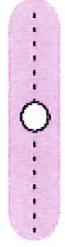
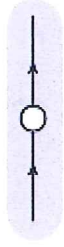
Matter



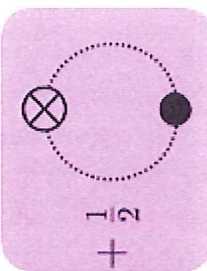
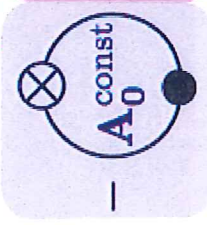
2PI-resummation



+matter-
contributions



$V_{\text{eff}}[\sigma, \vec{\pi}; A_0]$



Bonus: dynamical creation of hadrons

II-82

Challenges:

(1) QCD flow at all scales

(2) finite μ

$\Rightarrow$ bound state spectrum of QCD

(3) Non-equilibrium