Non-perturbative aspects of gauge theories Exercise sheet 13 – Scalars in Quantum Gravity

Lectures: Jan Pawlowski Tutorials: Nicolas Wink Manuel Reichert Institut für Theoretische Physik, Uni Heidelberg

j.pawlowski@thphys.uni-heidelberg.de n.wink@thphys.uni-heidelberg.de reichert@cp3.sdu.dk due date: 04 February 2019

Exercise 21: Scalars in Quantum Gravity

In this exercise we consider N_s scalar fields minimally coupled to quantum gravity. The ansatz for the effective action is given by

$$\Gamma = \Gamma_{\text{gravity}} + \sum_{i=1}^{N_{\text{s}}} \frac{Z_{\varphi^{i}}}{2} \int d^{4}x \sqrt{g} g^{\mu\nu} \partial_{\mu} \varphi^{i} \partial_{\nu} \varphi^{i} \,.$$
(1)

We want to estimate qualitatively the influence of the scalar fields on the fixed-point value of the Newton's coupling and of the cosmological constant. In a fluctuation computation we would obtain this from the scalar contribution to graviton n-point functions. Draw all diagrams that contribute to the graviton two- and three-point function. You should find that all diagrams contain a closed scalar loop and hence all contributions are proportional to $N_{\rm s}$.

For simplicity we compute the scalar contributions to the beta functions from the background field approximation. Compute with heat-kernel methods the trace of right-hand side of the Wetterich equation

$$\operatorname{Tr}\left[\frac{1}{\Gamma_k^{(2)} + R_k}\right]_{\varphi\varphi} \partial_t R_k.$$
(2)

In straight analogy to the exercise from last week, choose the regulator function as

$$R_k = \Gamma_{\varphi\varphi}^{(2)} \cdot r_k \left(\frac{\bar{\Delta}}{k^2}\right) \qquad \text{with} \qquad r_k(x) = \left(\frac{1}{x} - 1\right) \Theta(1 - x) \,. \tag{3}$$

The flow involves contributions from the scalar anomalous dimension $\eta_{\varphi} = -\partial_t \ln Z_{\varphi}$. We could determine those from the flow of scalar two-point function $\partial_t \Gamma_{\varphi\varphi}^{(2)}$, but here we set them simply to zero, $\eta_{\varphi} = 0$, which is indeed a good approximation is most cases. In straight analogy to last week, compare the \sqrt{g} and the $\sqrt{g}R$ terms in order to obtain the contributions to the beta function. You should find

$$\partial_t g_k = \beta_{g,\text{gravity}} + g_k^2 \frac{N_s}{6\pi} \left(1 - \frac{\eta_{\varphi}}{4} \right) ,$$

$$\partial_t \lambda_k = \beta_{\lambda,\text{gravity}} + g_k \frac{N_s}{4\pi} \left(1 - \frac{\eta_{\varphi}}{6} \right) .$$
(4)

In which direction are the fixed point values g_k^* and λ_k^* pushed? How do these changes influence the pure gravity contributions $\beta_{g,\text{gravity}}$ and $\beta_{\lambda,\text{gravity}}$, and how does this backreact on the fixed-point values? (See exercise sheet 12 for the precise form of these contributions.) Make a guess: which effect wins in the end?