



# Shear Viscosity from Black Holes

Kai Schweda  
Physikalisches Institut  
University of Heidelberg



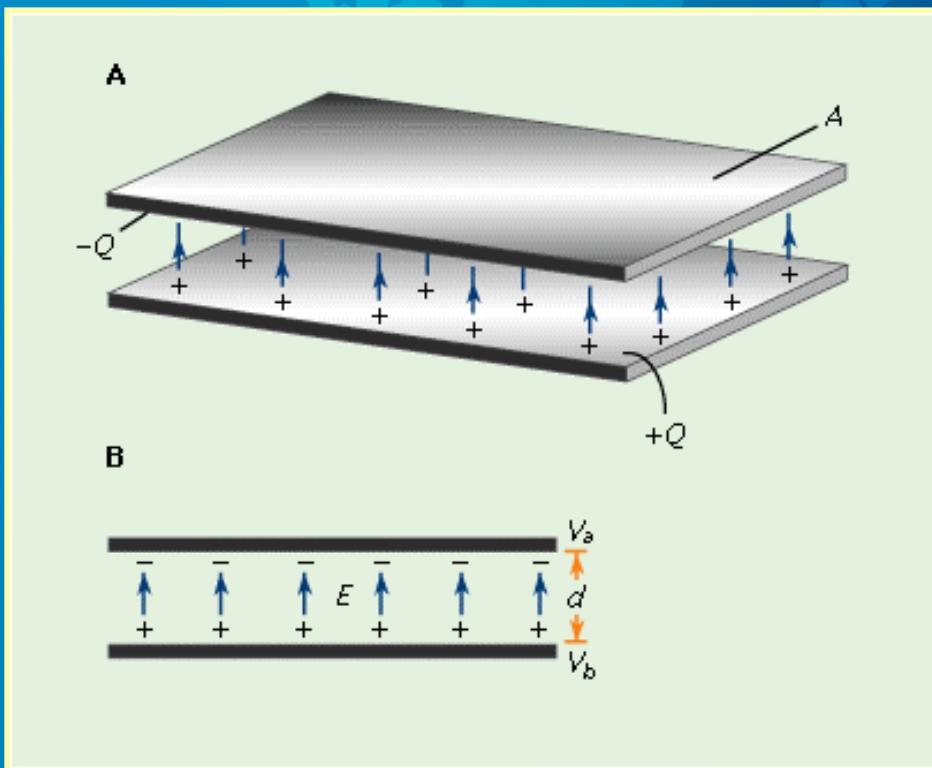
EMMI seminar, PI, Heidelberg

14 Jun 2010

# General Considerations

- Strong coupling  $\rightarrow$  quantum effects large
- Use AdS/CFT correspondence
- Holographic duality: relate string theory of higher dimension to 4-d gauge theory on the boundary
- Limit of strong coupling: string theory  $\rightarrow$  classical gravity (GR)

# Parallel Plate Capacitor



Source: <http://www.britannica.com>

- Bulk: 3-d space between plates
- Fluctuations of the field in the bulk induce fluctuations of electric charges on the surface (boundary)
- Correlations of surface charges correlated to bulk field

# Black Hole



Source: <http://media.photobucket.com>

- Black hole, mass  $M$
- Temp.  $T = \frac{\hbar c^3}{8\pi G M k_B}$
- Entropy  
 $S = A/4 \cdot (k_B c^3 / G \hbar)$   
A: area of horizon of boundary
- Physics of the interior region projected onto boundary: hologram

# Holographic Principle

- Conjectured by 't Hooft
- Quantum gravity in  $(d+1)$  dimensions  $\Leftrightarrow$  equivalent theory living on  $d$ -dimensional boundary  
 $\Rightarrow$  holographic dual

# AdS/CFT Correspondence

- Fields that propagate in the bulk have well defined values at asymptotic infinity (boundary)
- Asymptotic values behave like field and coupling at the boundary
- Anti-de Sitter spacetime: negative curvature
- Holographic duals are sometimes gauge theories
- E.g.  $\text{AdS}_5 \Leftrightarrow N=4$  Super Yang-Mills

# **AdS<sub>5</sub>×S<sub>5</sub> Geometry**

- AdS<sub>5</sub>: 5 dimensional Anti-de-Sitter space
- Infinitesimal line element

$$ds^2 = \frac{r^2}{L^2}(-dt^2 + dx^2) + \frac{L^2}{r^2}dr^2 + L^2d\Omega_5^2$$

S<sub>5</sub>: 5 dimensional sphere, neglect

- r: radial coordinate
- R = const.: 3+1 dim. flat Minkowski space
- R →∞: boundary
- L: curvature radius

# **AdS<sub>5</sub>×S<sub>5</sub> Geometry, cont'ed**

$$ds^2 = \frac{r^2}{L^2}(-dt^2 + dx^2) + \frac{L^2}{r^2}dr^2$$

- Require  $L \gg l_s$ , (classical approx.)
- 't Hooft coupling:  $\lambda = g_{\text{YM}}^2 N_c$
- $(L/l_s)^4 = \lambda$
- Classical approx. works at strong coupling

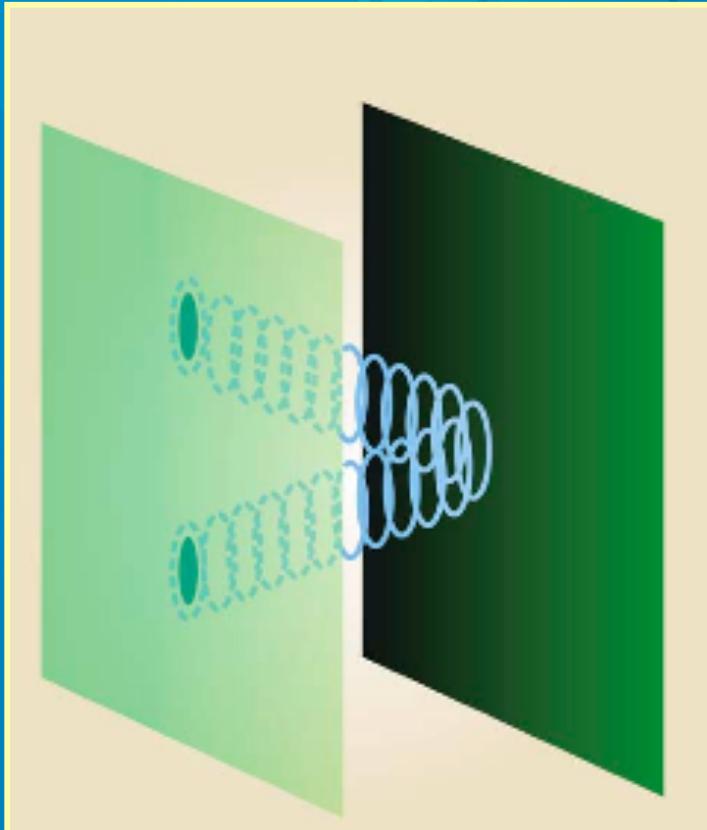
# **AdS<sub>5</sub>×S<sub>5</sub> Geometry, cont'ed**

- Rewrite for AdS<sub>5</sub> black hole

$$ds^2 = \frac{(\pi TL)^2}{u}(-(1-u^2)dt^2 + dx^2) + \frac{L^2}{4u^2(1-u^2)}du^2$$

- $u = (r_0/r)$ ,  $r_0$ : Schwarzschild (horizon) radius
- Horizon at  $u = 1$
- Boundary limit:  $u = \varepsilon$ , then  $\varepsilon \rightarrow 0$

# Ask the AdS/CFT Dictionary...



Source: Physics Today, p29, May 2010

- $\eta$  from  $T^{\mu\nu}$  (Kubo's formula)
- $T^{\mu\nu}$  corresponds to graviton  $h^{\mu\nu}$
- Graviton is disturbance in  $g_{\mu\nu}$
- Graviton at boundary propagates in the bulk and is scattered back
- Cross section  $\propto$  surface A
- Entropy  $s \propto$  surface A
- $\eta/s$  does not depend on A

# KSS bound on $\eta/s$

$$\frac{\eta}{s} \geq \frac{1}{4\pi} \cdot \frac{\hbar}{k_B} \left\{ 1 + \frac{15\zeta(3)}{\lambda^{3/2}} + \dots \right\}$$

Classical implementation bound

from string theory

$$\frac{\eta}{s} = \frac{1}{4\pi} \cdot \frac{\hbar}{k_B} \left( 1 - \frac{1}{2N_C} \right)$$

Potentially lower bound from SU(2)

## Some remarks

- Relativistic fluid, but bound does not depend on speed of light
- $N=4$  Super Yang-Mills is **not** QCD
- $N_c = 3$ , not large
- No confinement
- Quarks are massless
- However, details might not matter too much, system driven by temperature and degrees of freedom

# References

- T. Schaefer and D. Teaney,  
Rep. Prog. Phys. 72 (2009), 126001, p22-26.
- P.K. Kovtun, D.T. Son, and A.O. Starinets,  
Phys. Rev. Lett. 94 (2005) 111601.
- J.M. Maldacena, TASI 2003 Lectures on AdS/CFT  
correspondence, arXiv:hep-th/0309246.
- C.V. Johnson and P. Steinberg,  
Physics Today (May 2010) 29.