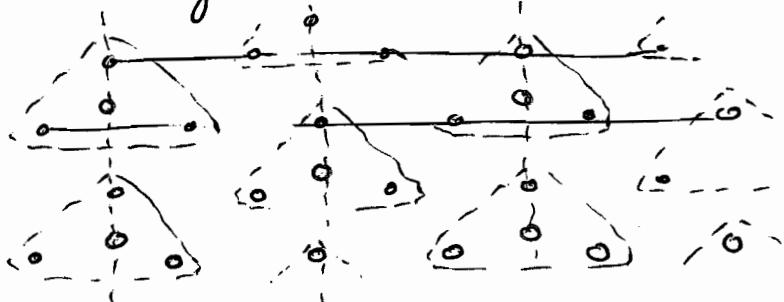


3.3 Approximation by cumulants

We consider the Ising model in two dim.

on a triangular lattice



This allows to define a simple, non-linear, RG-transformation

$$S'_\alpha = \text{sign} [S_\alpha^{(1)} + S_\alpha^{(2)} + S_\alpha^{(3)}] = f_\alpha(S_i) \quad (3.87)$$

Evidently, eq. (3.87) maps spin variables $S_i = \pm 1$ into spin variables $S_\alpha = \pm 1$.

The scaling factor is $s = \sqrt{3}$ but the map does not preserve the angles!

The new hamiltonian reads

$$e^{-\Omega - H'[S']} = \sum_{\{S\}} \prod_{\alpha} \delta[S'_\alpha - f_\alpha(S_i)] e^{-H[S]} \quad (3.88)$$

The δ -fcts. in eq. (3.88) fix the spin configs.

in each block to obey $f_\alpha(s_i) = s'_\alpha$. Hence the

hamiltonian H separates naturally into

interaction between blocks

$$H = H_0 + V \quad \begin{matrix} \uparrow \\ \text{interaction inside blocks} \end{matrix} \quad (3.89)$$

where H_0 stands for interactions in a given block

and V mediates interactions between blocks.

We use eq. (3.89) to reorder eq. (3.88) as

$$\boxed{e^{-\Omega - H'[\{s'\}]}} = \frac{\sum_{\{s\}} e^{-H_0 \prod_\alpha \delta[s'_\alpha - f_\alpha(s)]} \sum_{\{s\}} e^{-H_0 - V \prod_\alpha \delta[s'_\alpha - f_\alpha(s)]}}{\left[\sum_{\{s\}} e^{-H_0 \prod_\alpha \delta[s'_\alpha - f_\alpha(s)]} \right]} \quad (3.90)$$

The normalisation/prefactor $\sum e^{-H[\{s\}]}$ is

straightforward to compute: one has to sum over all configs $\{s\}$ that satisfy the constraint

with the weight factor $\sum_{\{s^{(i)}\}} \delta[s'_\alpha - f_\alpha(s_i)] e^{-H_0} = Z_\alpha$.

The remaining term is the expectation value of the operator e^{-V} in the theory with hamiltonian H_0 (with restricted configs. due to $\delta f = J$). We write

$$\langle \sigma \rangle_o = \frac{\sum_{\{S\}} \sigma[\{S\}] e^{-H_0[\{S\}]} \prod_{\alpha} \delta[S_{\alpha}' - f_{\alpha}(S)]}{\sum_{\{S\}} e^{-H_0[\{S\}]} \prod_{\alpha} \delta[S_{\alpha}' - f_{\alpha}(S)]} \quad (3.91)$$

These considerations lead to

$$e^{-\Omega - H'[\{S'\}]} = Z_o \cdot \langle e^{-V} \rangle_o \quad (3.92)$$

with

$$Z_o = \sum_{\{S\}} e^{-H_0[\{S\}]} \prod_{\alpha} \delta[S_{\alpha}' - f_{\alpha}(S)]$$

Note that $\langle \sigma \rangle_o$ rather is $\langle \sigma \rangle_o[\{S'\}]$. Even though Z_o is simple to compute, $\langle e^{-V} \rangle_o$ is not. Indeed e^{-V} couples all blocks (V^n mediates couplings of n neighbouring blocks). For the sake of simplicity, let us expand $\langle e^{-V} \rangle$ in powers of V .

To heat and we use that

$$\begin{aligned}
 \ln \langle e^{-v} \rangle &= \ln \left(\langle e^{-(V-\langle v \rangle)} \rangle e^{-\langle v \rangle} \right) \\
 &= -\langle V \rangle + \ln \left\langle 1 - (V - \langle v \rangle) + \frac{1}{2} (V - \langle v \rangle)^2 \right. \\
 &\quad \left. - \frac{1}{3!} (V - \langle v \rangle)^3 + \dots \right\rangle \\
 &= -\langle V \rangle - (V - \langle v \rangle)^2 + O((V - \langle v \rangle)^3)
 \end{aligned} \tag{3.100}$$

As indicated above, the comp. of the higher terms gets increasingly involved. Here, we already drop the second term, to wit

$$\langle e^{-V} \rangle = e^{-\langle V \rangle_0} \tag{3.101}$$

which is, again, a mean-field approximation.

It is left to compute $\langle V \rangle_0$ & Z_0 :

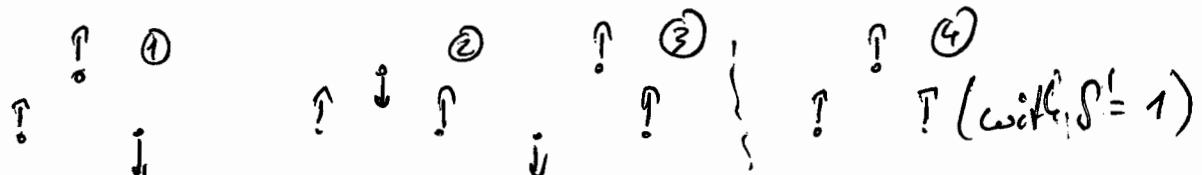
- (i) Z_0 factorises as H_0 only has interactions inside blocks by definition.

$$Z_0 = \prod Z_{\text{block}} [S_\alpha] \tag{3.102}$$

However, Z_{block} is independent of S_i^z as it is quadratic in the spins S_i^z :

$$Z_{\text{block}} = e^{3K} + 3e^{-K} \stackrel{\beta J}{=} (3.103)$$

with



Hence

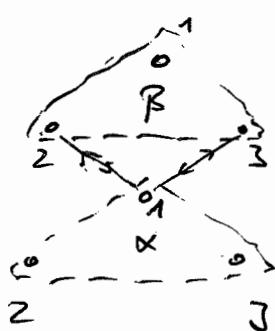
$$Z_0 = Z_{\text{block}}^{\frac{N}{3}} = (e^{3u} + 3e^{-u})^{\frac{N}{3}} \quad (3.104)$$

of Blocks

(ii) Computation of $\langle V \rangle_0$ where

$$V_{\alpha\beta} = -K S_\alpha^{(1)} (S_\beta^{(2)} + S_\beta^{(3)}) \quad (3.105)$$

between 2 blocks α, β with



As H_0 does not mediate interactions between different blocks, we have

$$\langle S_\alpha^{(i)} S_\beta^{(j)} \rangle = \langle S_\alpha^{(i)} \rangle \langle S_\beta^{(j)} \rangle \quad (3.106)$$

and we only have to evaluate the expectation values of single spins S_i for a given $S_{\alpha/\beta}^1$ with the Hamiltonian H_0 . This leads to

$$\boxed{\langle S_{\alpha}^{(1)} \rangle_o = \frac{1}{Z_{\text{block}}} \sum_{\{S_{\alpha}^{(1)}\}} S_{\alpha}^{(1)} e^{K[S_{\alpha}^{(1)} S_{\alpha}^{(2)} + S_{\alpha}^{(1)} S_{\alpha}^{(3)} + S_{\alpha}^{(2)} S_{\alpha}^{(3)}]}$$

$$= \frac{1}{Z_{\text{block}}} [(e^{3U} + e^{-K}) D_{+1S_{\alpha}^1} - (e^{3U} + e^{-K}) D_{-1S_{\alpha}^1}]$$

where we have used that for $S_{\alpha}^1 = 1$: (3.107)

$$\begin{array}{cccc} 1 & 2 & 3 & \\ \uparrow & \uparrow & \uparrow & \\ S_{\alpha}^{(1)} = 1 & \downarrow & & \\ + e^{3U} & & & \end{array} \quad \begin{array}{ccc} 12 & 3 & \\ \uparrow \uparrow & \downarrow & \\ \downarrow & -K & \\ + e^{-K} & & \end{array} \quad \begin{array}{ccc} 12 & 3 & \\ \uparrow \downarrow & \uparrow & \\ \downarrow & -K & \\ + e^{-K} & & \end{array} \quad \begin{array}{ccc} 12 & 3 & \\ \downarrow \uparrow \uparrow & & \\ \downarrow & -K & \\ - e^{-K} & & \end{array}$$

$$= e^{3U} + e^{-K}$$

and similarly for $S_{\alpha}^1 = -1$: $-(e^{3U} + e^{-K})$.

In other words

$$\boxed{\langle S_{\alpha}^{(1)} \rangle_o [S_{\alpha}^1] = \frac{1}{e^{3U} + e^{-K}} (e^{3U} + e^{-3K}) S_{\alpha}^1} \quad (3.108)$$

With eq. (3.109) and eq. (3.108) we get

$$\boxed{\begin{aligned} \langle V_{\alpha\beta} \rangle_0 &= 2K \left(\frac{e^{3U} + e^{-U}}{e^{3U} + 3e^{-U}} \right)^2 S_\alpha' S_\beta' \\ &= U' S_\alpha' S_\beta' \end{aligned}} \quad (3.109)$$

with $U' = 2K \left(\frac{e^{3U} + e^{-U}}{e^{3U} + 3e^{-U}} \right)^2$

This leads to the final blocked Hamiltonian

$$\boxed{e^{-\Omega - H'[S']}} = e^{N/3 \ln(e^{3U} + 3e^{-U} + U' \sum_{\langle\alpha\beta\rangle} S_\alpha' S_\beta')} \quad (3.110)$$

The Hamiltonian in eq. (3.110) allows for a simple fixed point analysis. The FP is given by

$$\boxed{U_* = 2U_* \left(\frac{e^{4U_*} + 1}{e^{4U_*} + 3} \right)^2} \quad (3.111)$$

which can be solved algebraically.

We get

$$e^{K_*} = \underbrace{1 + 2\sqrt{2}}_x \quad (3.112)$$

and

$$K_* \approx 0.336 \quad (\text{exact } 0.275)$$

As we have only one coupling there is only one scaling exponent b :

$$s^b = \left. \frac{\partial U'}{\partial U} \right|_{U_*} = 2 \left(\frac{x+1}{x+3} \right)^2 + \frac{32 K_* x}{(x+3)^2} \left(\frac{x+1}{x+3} \right) \quad (3.113)$$

eq.(3.50), p. 83

With $s = \sqrt{3}$ and $s^b \approx 1.634$ we conclude

$$\nu = 1/b = 1 / \frac{\ln 1.634}{\ln \sqrt{3}} = 1.118 \quad (3.114)$$

(exact 1)

We summarise these results in a table

| | exact (ourself) | 1 st order cumulant | MF |
|-------|-----------------|--------------------------------|-------|
| K_* | 0.275 | 0.336.. | $1/6$ |
| ν | 1 | 1.118.. | $1/2$ |