


4.2 Perturbation theory

(a) 1-loop: $\partial_t \Gamma_k^{1\text{-loop}} = \frac{1}{2}$ 

$$\Rightarrow \text{propagator } x \text{---} y = \frac{1}{S^{(2)}[\phi] + R_k}(x, y)$$

and hence

$$\partial_t \Gamma_k^{1\text{-loop}}[\phi] = \frac{1}{2} \text{Tr} \frac{1}{S^{(2)}[\phi] + R_k} \partial_t R_k \quad (4.55)$$

$\swarrow \searrow$
 $k\text{-dep}$

$$= \frac{1}{2} \text{Tr} \partial_t \ln(S^{(2)}[\phi] + R_k)$$

$\hat{=}$ (Tr $\partial_t \neq \partial_t$ Tr, strictly speaking)

Integration:

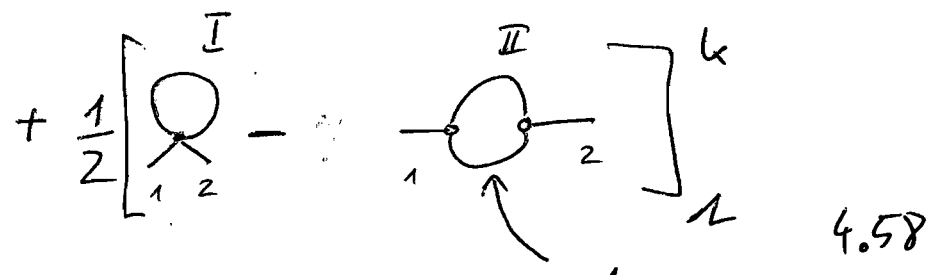
$$\Gamma_k^{1\text{-loop}}[\phi] = \Gamma_{\Lambda}^{1\text{-loop}}[\phi] + \int_{\Lambda}^k \frac{dk'}{k'} \partial_{t'} \Gamma_{k'}^{1\text{-loop}}[\phi]$$

$$\Rightarrow \Gamma_k^{1-loop}[\phi] = \Gamma_k^{1-loop}[\phi] + \frac{1}{2} \text{Tr} \left[\ln(S^{(2)}[\phi] + R_k) - \ln(S^{(2)} + R_k) \right]$$

finite (4.57)

e.g.:

$$\Gamma_k^{1-loop(2)}[\phi] = \Gamma_k^{1-loop(2)}[\phi]$$



$$\Gamma_k^{1-loop}[\sigma](p_1, \dots, p_4) = \frac{1}{2} \left[\text{Diagram I} \right]_k + \Gamma_k^{1-loop(2)}[\sigma](p_1, \dots, p_4)$$

$$\Gamma_k^{1-loop(4)}[\phi] = \frac{\delta^2}{\delta\phi^2} \Gamma_k^{1-loop(2)}[\phi]$$

See pages I-74(a,b)



$$\Rightarrow \Gamma_k^{1-loop(4)}[\sigma](p_1, \dots, p_4) = -\frac{1}{2} \left[\text{Diagram I} \right]_k + \Gamma_k^{1-loop(4)}[\sigma](p_1, \dots, p_4)$$

(p₁, ..., p₄)

Renormalisation:

(1) Γ_k indep. of Λ !

$$\Rightarrow \Lambda \partial_\Lambda \Gamma_k = 0$$

$$= \Lambda \partial_\Lambda \Gamma_\Lambda + \frac{1}{2} \text{Tr} \left[\ln(\delta^{(2)}(p) [A] + R) \right]_\Lambda^k$$

(2) Λ -dep. of Γ_Λ is fixed by Flow (4.59)

\Rightarrow Renormalisation is

(A) adjusting Λ -indep. of $\Gamma_k, k \neq \Lambda$
 \sim regularisation

(B) fixing Λ -indep parts of Γ_Λ
 \sim renormalisation conditions

(3) extends trivially to full flow

Example I: m^2 in ϕ^4 - theory in 4 dim.

$$\partial_{\epsilon} m^2 = \partial_{\bar{\phi}}^2 \partial_{\epsilon} \Gamma_n[\bar{\phi} \text{ const}] \Big|_{\bar{\phi}=0}^{1\text{-loop}} \quad (4.60)$$

$$\partial_{\epsilon} \text{eq. (4.58)} \rightarrow = -\frac{1}{2} \text{diagram} + \text{diagram}$$

or p. 1B a

$$\left. \begin{array}{l} \text{diagram} \\ \text{diagram} \end{array} \right|_{\bar{\phi}=0} = 0$$

$$\left. \begin{array}{l} \text{diagram} = \frac{1}{p^2 + m^2 + R_u} \\ \text{diagram} = \lambda \\ \text{diagram} = \partial_{\epsilon} R_u \end{array} \right\} \partial_{\epsilon} m^2 = -\frac{1}{2} \int \frac{d^4 p}{(2\pi)^4} \frac{1}{(p^2 + m^2 + R_u)^2} \partial_{\epsilon} R_u \quad (4.61)$$

$$\text{with } S[\phi] = \frac{1}{2} \int \frac{d^4 p}{(2\pi)^4} \phi(p) (p^2 + m^2) \phi(-p)$$

$$+ \frac{\lambda}{4!} \int_{p_1, \dots, p_4} \phi(p_1) \dots \phi(p_4) (2\pi)^4 \delta(p_1 + \dots + p_4)$$

$$\text{with } \int_p = \int \frac{d^4 p}{(2\pi)^4} \quad (4.62)$$

Now we rescale all dim. full variables/parameters with $k^{-\#}$, where $\#$ is the mass dimension, i.e. e.g.

$$\hat{p}^2 = p^2/k^2, \quad \hat{m}^2 = m^2/k^2 \quad (4.63)$$

Hence ($\hat{k} = 1$)

$$\partial_{\epsilon} m_u^2 = -k^2/2 \int \frac{d^4 \hat{p}}{(2\pi)^4} \frac{1}{(\hat{p}^2 + \hat{m}^2 + R_1)^2} \partial_{\epsilon} R_1$$

$$= -k^2/2 \int \frac{d^4 \hat{p}}{(2\pi)^4} \frac{1}{(\hat{p}^2 + R_1)^2} \partial_{\epsilon} R_1 + \underbrace{k^2 \mathcal{O}(\hat{m}^2)}_{\mathcal{O}(k^0)}$$

$$= c(R_1) \cdot k^2 + \mathcal{O}(k^0) \quad (4.64)$$

With

$$R_k = (k^2 - p^2) \Theta(k^2 - p^2) \quad (4.65)$$

we get

$$\partial_{\epsilon} m_u^2 = -k^2 \int \frac{d^4 \hat{p}}{(2\pi)^4} \frac{1}{(1 + \hat{m}^2)^2} \Theta(1 - \hat{p}^2)$$

$$= \frac{k^2}{(1 + \hat{m}^2)^2} \frac{\Omega_4}{(2\pi)^4}, \quad \Omega_4 = 2\pi^2 \quad (4.66)$$

see exercise

Example II: β -function in ϕ^4 -theory in 4-dim

127c

$$S[\phi] = \frac{1}{2} \int \frac{d^4 p}{(2\pi)^4} \phi(p) (p^2 + m^2) \phi(-p)$$

$$+ \frac{\lambda}{4!} \int_{p_1, \dots, p_4} \phi(p_1) \dots \phi(p_4) \cdot (2\pi)^4 \delta^{(4)}(p_1 + \dots + p_4)$$

(4.67)

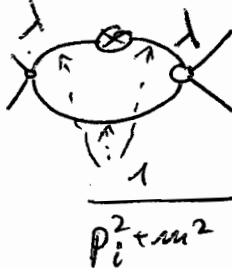
with $\int_p = \int \frac{d^d p}{(2\pi)^d}$, here $d=4$.

Inserting into $\partial_z \Gamma^{(4)}|_{\phi=0} = \partial_z \lambda = \underline{\underline{\dot{\lambda}}}$

1-loop: $\Gamma^{(2)} = S^{(2)} = (p^2 + m^2) \delta^{(4)}(p+q)$

$\Gamma^{(4)} = S^{(4)} = \lambda (2\pi)^4 \delta^{(4)}(p_1 + \dots + p_4)$

(4.68)

E-146 $\Rightarrow \dot{\Gamma}^{(4)} = 3 \cdot$ 

$$= 3 \lambda^2 \int \frac{d^4 p}{(2\pi)^4} \left(\frac{1}{p^2 + m^2 + R_U} \dot{R}_U(p^2) \frac{1}{p^2 + m^2} \right) \cdot \frac{1}{p^2 + m^2 + R_U}$$

(4.69)

$$\dot{\Gamma}(u) = \dot{\lambda} = 3\lambda^2 \cdot \int \frac{d\Omega_4}{(2\pi)^4} \cdot \frac{1}{2} \int_0^\infty dp^2 p^2 \left(\frac{1}{p^2 + m^2 + R_u(p^2)} \right)^3 R_u(p^2) \quad 127d$$

Introduce $x = p^2/u^2$ (4.70)

$$R_u(p^2) = p^2 r(x)$$

$$\Rightarrow R_u(p^2) = -2p^2 x r'(x) = -2k^2 x r'(x) \quad (4.71)$$

$$\Rightarrow \dot{\lambda} = -\frac{2}{2} 3\lambda^2 \frac{\Omega_4}{(2\pi)^4} \int_0^\infty dx x^3 \left(\frac{1}{x(1+r) + m^2} \right)^3 r'$$

$$= -3\lambda^2 \frac{\Omega_4}{(2\pi)^4} \int_0^\infty dx \left(\frac{1}{1+r + m^2/x} \right)^3 r'$$

$$= \frac{3}{2} \lambda^2 \frac{\Omega_4}{(2\pi)^4} \int_0^\infty dx \left[\frac{d}{dx} \frac{1}{(1+r + m^2/x)^2} - \frac{m^2}{x^2} \frac{1}{(1+r + m^2/x)^3} \right]$$

$$= \frac{3}{2} \lambda^2 \frac{\Omega_4}{(2\pi)^4} + m^2 \text{-corrections} \quad (4.72)$$

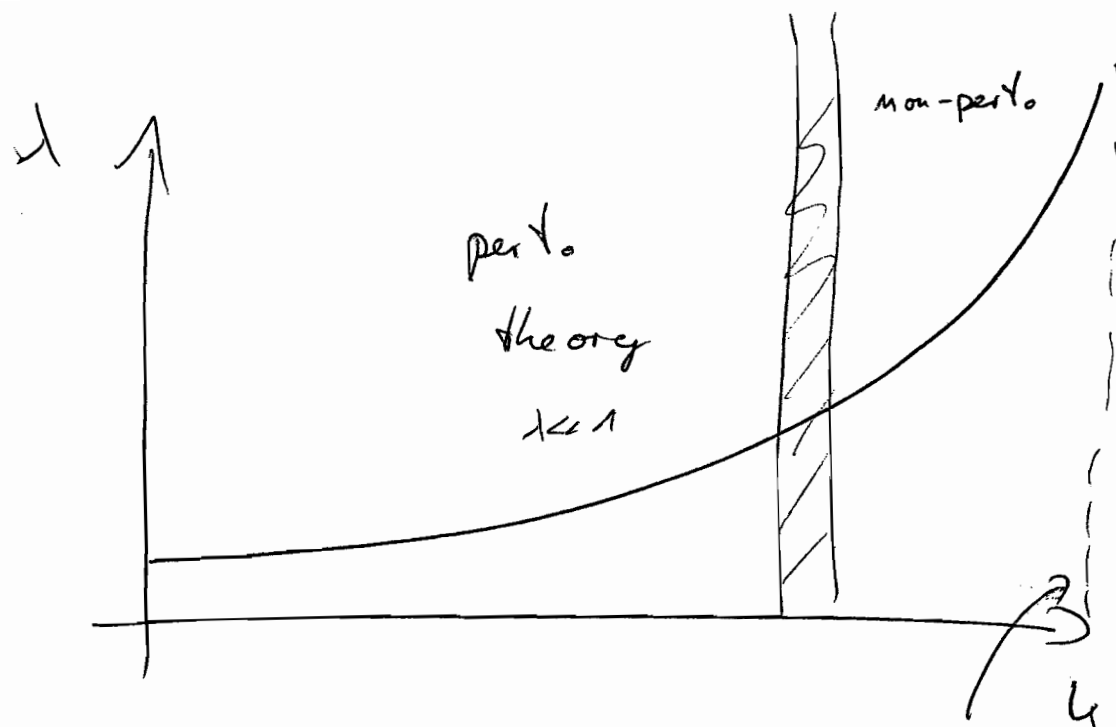
$$\Omega_4 = 2\pi^2$$

$$m^2 = m_0^2/u^2 + \lambda^2 \text{-terms}$$

$$\Rightarrow \boxed{\partial_\pm \lambda \approx \frac{3}{16\pi^2} \lambda^2}$$

$$\Rightarrow \lambda(t) = \frac{\lambda_0}{1 + \frac{3}{16\pi^2} (t-t_0) \lambda_0} \quad (4.73)$$

Landau pole



Landau-pole

triviality: demand finite λ for $t \rightarrow \infty$

- ϕ^4 -theory valid at all scales
- no Landau pole

$$\Rightarrow \lambda_{\text{phys}} = \lambda_{k=0} \stackrel{!}{=} 0 \quad (4.74)$$

[as $t \rightarrow \infty$]

(b) 2-loop :

$$\partial_t \Gamma_k^{2\text{-loop}} = \frac{1}{2} \text{diagram} \quad \text{with } x \oplus y = \frac{1}{\Gamma_k^{1\text{-loop}}(s)} [dJ + R_k]$$

$$= \frac{1}{2} \text{diagram} - \frac{1}{2} \text{diagram} \quad \Delta \Gamma_k^{(2)}$$

↑
1-loop

with $\Delta \Gamma_k^{(2)} = \Gamma_k^{1\text{-loop}}(s) - \int^{(2)}$

$$= \frac{1}{2} \left[\text{diagram} - \text{diagram} \right]_k + \left(\Gamma_k^{1\text{-loop}}(s) - \int^{(2)} \right)$$

$$= \frac{1}{2} \left[\text{diagram} - \text{diagram} \right]$$

$$= \frac{1}{2} (\text{diagram}^k - \text{diagram}^k) - \frac{1}{2} (\text{diagram} - \text{diagram})$$

It follows

$$\partial_t \Gamma_k^{2\text{-loop}} = \frac{1}{2} \text{diagram} - \frac{1}{4} \text{diagram} + \frac{1}{4} \text{diagram}$$

$$\Rightarrow \partial_t \Gamma_k^{2-loop} = \partial_t(1-loop) - \frac{1}{4} \left[\text{diagram 1} - \text{diagram 2} \right]$$

$$+ \frac{1}{4} \left[\text{diagram 3} - \text{diagram 4} \right]$$

$$= \partial_t(1-loop) + \partial_t \left\{ \frac{1}{8} \text{diagram 1} - \frac{1}{4} \text{diagram 2} \right.$$

$$\left. - \frac{1}{12} \text{diagram 3} + \frac{1}{4} \text{diagram 4} \right\}$$

$$\Rightarrow \Gamma_k^{2-loop} = \Gamma_n^{2-loop} + \int_n^k \frac{d\ell}{\ell} \partial_{t'} \Gamma_{\ell}$$

$$= S_{ce} + (1-loop)_{ren}$$

$$+ \frac{1}{8} \text{diagram 1} - \frac{1}{4} \text{diagram 2}$$

$$- \frac{1}{12} \text{diagram 3} + \frac{1}{4} \text{diagram 4} + \left(\Gamma_n^{2-loop} - S - \Gamma_n^{1-loop} \right)$$

$$\Rightarrow \Delta \Gamma_k^{2-loop(2)} \Rightarrow 3-loop$$