PROBLEMS FOR QUANTUM FIELD THEORY 1 2. Sheet

Suggested reading before solving these problems: Chapters 2.3 in the script and/or Chapters 2.3 to 2.4 of *Peskin & Schroeder*.

PROBLEM 1: Commutation relations For a real scalar field $\phi(x)$ the Lagrangian density is

$$\mathcal{L} = \frac{1}{2} \partial_{\mu} \phi \partial^{\mu} \phi - \frac{1}{2} m^2 \phi^2.$$

The canonical momentum density is $\pi = \partial \mathcal{L} / \partial \dot{\phi} = \dot{\phi}$. The theory is quantized by promoting ϕ and π to operators in the Schrödinger picture with the commutation relations

$$\begin{aligned} [\phi(\vec{x}), \pi(\vec{y})] &= i\delta^{(3)}(\vec{x} - \vec{y}), \\ [\phi(\vec{x}), \phi(\vec{y})] &= [\pi(\vec{x}), \pi(\vec{y})] = 0. \end{aligned}$$

Introduce now the operators a and a^{\dagger} by

$$\begin{split} \phi(\vec{x}) &= \int \frac{d^3p}{(2\pi)^3} \frac{1}{2\omega_{\vec{p}}} \left\{ a(\vec{p}) e^{i\vec{p}\vec{x}} + a^{\dagger}(\vec{p}) e^{-i\vec{p}\vec{x}} \right\}, \\ \pi(\vec{x}) &= \frac{-i}{2} \int \frac{d^3p}{(2\pi)^3} \left\{ a(\vec{p}) e^{i\vec{p}\vec{y}} - a^{\dagger}(\vec{p}) e^{-i\vec{p}\vec{x}} \right\} \end{split}$$

and derive the commutation relations

$$[a(\vec{p}), a^{\dagger}(\vec{q})], \quad [a(\vec{p}), a(\vec{q})], \quad [a^{\dagger}(\vec{p}), a^{\dagger}(\vec{q})].$$

PROBLEM 2: Complex scalar field: quantization

Consider the Lagrangian density for a free complex scalar field

$$\mathcal{L} = \partial_{\mu} \phi^* \partial^{\mu} \phi - m^2 \phi^* \phi.$$

a) Show that the canonical momenta of ϕ and ϕ^* are

$$\pi = \dot{\phi}^*, \quad \pi^* = \dot{\phi}$$

and derive an expression for the Hamiltonian H.

b) Proceed to quantization by promoting ϕ, ϕ^* and π, π^* to operators ϕ, ϕ^{\dagger} and π, π^{\dagger} (in the Schrödinger picture). What would you postulate as their commutation relations?

c) Introduce now creation and annihilation operators by writing

$$\begin{split} \phi(\vec{x}) &= \int \frac{d^3p}{(2\pi)^3} \frac{1}{2\omega_{\vec{p}}} \left\{ a(\vec{p}) e^{i\vec{p}\vec{x}} + b^{\dagger}(\vec{p}) e^{-i\vec{p}\vec{x}} \right\}, \\ \pi(\vec{x}) &= \int \frac{d^3p}{(2\pi)^3} \frac{-i}{2} \left\{ -a^{\dagger}(\vec{p}) e^{-i\vec{p}\vec{x}} + b(\vec{p}) e^{i\vec{p}\vec{x}} \right\}. \end{split}$$

Why do we now need operators b, b^{\dagger} in addition to a, a^{\dagger} ? Convince yourself that the commutation relations

$$\begin{bmatrix} a(\vec{p}), a^{\dagger}(\vec{q}) \end{bmatrix} = \begin{bmatrix} b(\vec{p}), b^{\dagger}(\vec{q}) \end{bmatrix} = 2\omega_{\vec{p}} (2\pi)^{3} \, \delta^{(3)}(\vec{p} - \vec{q}), \\ \begin{bmatrix} a(\vec{p}), a(\vec{q}) \end{bmatrix} = \begin{bmatrix} b(\vec{p}), b(\vec{q}) \end{bmatrix} = 0, \\ \begin{bmatrix} a(\vec{p}), b(\vec{q}) \end{bmatrix} = \begin{bmatrix} a(\vec{p}), b^{\dagger}(\vec{q}) \end{bmatrix} = 0$$

are consistent with the ones postulated in part b).

d) Show that the Hamiltonian can be written as

$$H = \int \frac{d^3p}{(2\pi)^3} \frac{1}{2} \left\{ a^{\dagger}(\vec{p})a(\vec{p}) + b^{\dagger}(\vec{p})b(\vec{p}) \right\} + \text{const.}$$

Why is the positive sign in front of the $b^{\dagger}b$ term important? What is the physical interpretation of b and b^{\dagger} ?

e) Switch now to the Heisenberg picture

$$\phi_H(t, \vec{x}) = e^{iHt} \phi(\vec{x}) e^{-iHt}$$

Show that

$$e^{iHt} a(\vec{p}) e^{-iHt} = a(\vec{p}) e^{-i\omega_{\vec{p}}t}, \qquad e^{iHt} a^{\dagger}(\vec{p}) e^{-iHt} = a^{\dagger}(\vec{p}) e^{i\omega_{\vec{p}}t}, e^{iHt} b(\vec{p}) e^{-iHt} = b(\vec{p}) e^{-i\omega_{\vec{p}}t}, \qquad e^{iHt} b^{\dagger}(\vec{p}) e^{-iHt} = b^{\dagger}(\vec{p}) e^{i\omega_{\vec{p}}t},$$
(1)

and therefore

$$\phi_H(t,\vec{x}) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{2\omega_{\vec{p}}} \left\{ a(\vec{p}) e^{-ipx} + b^{\dagger}(\vec{p}) e^{ipx} \right\}.$$

PROBLEM 3: "String quantization"

Consider the Lagrangian of the clamped string derived in problem 1 on sheet 1

$$L = \int_0^R dx \left\{ \frac{1}{2} \sigma \left(\frac{\partial y}{\partial t} \right)^2 - \frac{1}{2} \tau \left(\frac{\partial y}{\partial x} \right)^2 \right\}$$
$$= \sum_{n=1}^\infty \left\{ \frac{1}{2} \sigma \dot{q}_n^2 - \frac{1}{2} \tau \left(\frac{n\pi}{R} \right)^2 q_n^2 \right\}.$$

- a) How would you quantize that string? What commutation relations would you postulate for the "field operator" y(x) and its canonical momentum $(\pi_y)(x)$ or for the operators q_n and their canonical momenta $(\pi_q)_n$?
- b) Express the Schrödinger picture operators q_n and its canonical momentum in terms of creation and annihilation operators a_n, a_n^{\dagger} . (Use a normalization as in the lecture.) Can you guess the form of the Hamiltonian H expressed in terms of these operators?