Quantum Field Theory 1 – Problem set 2

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Suggested reading before solving these problems: Chapters 2.3 in the script and/or Chapters 2.3 to 2.4 of *Peskin & Schroeder*.

Problem 1: Commutation relations

For a real scalar field $\phi(x)$ the Lagrangian density is

$$\mathcal{L} = \frac{1}{2} \partial_{\mu} \phi \partial^{\mu} \phi - \frac{1}{2} m^2 \phi^2.$$

The canonical momentum density is $\pi = \partial \mathcal{L} / \partial \dot{\phi} = \dot{\phi}$. The theory is quantized by promoting ϕ and π to operators in the Schrödinger picture with the commutation relations

$$\begin{array}{lll} [\phi({\bm{x}}),\pi({\bm{y}})] &=& i\delta^{(3)}({\bm{x}}-{\bm{y}}), \\ [\phi({\bm{x}}),\phi({\bm{y}})] &=& [\pi({\bm{x}}),\pi({\bm{y}})] = 0 \end{array}$$

Introduce now the operators a and a^{\dagger} by

$$\begin{split} \phi(\boldsymbol{x}) &= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{2\omega_{\boldsymbol{p}}} \left\{ a(\boldsymbol{p}) \, e^{i\boldsymbol{p}\boldsymbol{x}} + a^{\dagger}(\boldsymbol{p}) \, e^{-i\boldsymbol{p}\boldsymbol{x}} \right\}, \\ \pi(\boldsymbol{x}) &= \frac{-i}{2} \int \frac{d^3 p}{(2\pi)^3} \left\{ a(\boldsymbol{p}) \, e^{i\boldsymbol{p}\boldsymbol{x}} - a^{\dagger}(\boldsymbol{p}) \, e^{-i\boldsymbol{p}\boldsymbol{x}} \right\} \end{split}$$

and derive the commutation relations

$$[a(\boldsymbol{p}), a^{\dagger}(\boldsymbol{q})], \quad [a(\boldsymbol{p}), a(\boldsymbol{q})], \quad [a^{\dagger}(\boldsymbol{p}), a^{\dagger}(\boldsymbol{q})].$$

Problem 2: Complex scalar field: quantization

Consider the Lagrangian density for a free complex scalar field

$$\mathcal{L} = \partial_{\mu} \phi^* \partial^{\mu} \phi - m^2 \phi^* \phi.$$

a) Show that the canonical momenta of ϕ and ϕ^* are

$$\pi = \dot{\phi}^*, \quad \pi^* = \dot{\phi}$$

and derive an expression for the Hamiltonian H.

- b) Proceed to quantization by promoting ϕ, ϕ^* and π, π^* to operators ϕ, ϕ^{\dagger} and π, π^{\dagger} (in the Schrödinger picture). What would you postulate as their commutation relations?
- c) Introduce now creation and annihilation operators by writing

$$\begin{split} \phi(\boldsymbol{x}) &= \int \frac{d^3 p}{(2\pi)^3} \frac{1}{2\omega_{\boldsymbol{p}}} \left\{ a(\boldsymbol{p}) \, e^{i\boldsymbol{p}\boldsymbol{x}} + b^{\dagger}(\boldsymbol{p}) \, e^{-i\boldsymbol{p}\boldsymbol{x}} \right\}, \\ \pi(\boldsymbol{x}) &= \int \frac{d^3 p}{(2\pi)^3} \frac{-i}{2} \left\{ -a^{\dagger}(\boldsymbol{p}) \, e^{-i\boldsymbol{p}\boldsymbol{x}} + b(\boldsymbol{p}) \, e^{i\boldsymbol{p}\boldsymbol{x}} \right\}. \end{split}$$

Why do we now need operators b, b^{\dagger} in addition to a, a^{\dagger} ? Convince yourself that the commutation relations

$$\begin{aligned} &[a(\boldsymbol{p}), a^{\dagger}(\boldsymbol{q})] &= [b(\boldsymbol{p}), b^{\dagger}(\boldsymbol{q})] = 2\omega_{\boldsymbol{p}} (2\pi)^{3} \,\delta^{(3)}(\boldsymbol{p} - \boldsymbol{q}), \\ &[a(\boldsymbol{p}), a(\boldsymbol{q})] &= [b(\boldsymbol{p}), b(\boldsymbol{q})] = 0, \\ &[a(\boldsymbol{p}), b(\boldsymbol{q})] &= [a(\boldsymbol{p}), b^{\dagger}(\boldsymbol{q})] = 0 \end{aligned}$$

are consistent with the ones postulated in part b).

d) Show that the Hamiltonian can be written as

$$H = \int \frac{d^3p}{(2\pi)^3} \frac{1}{2} \left\{ a^{\dagger}(\boldsymbol{p})a(\boldsymbol{p}) + b^{\dagger}(\boldsymbol{p})b(\boldsymbol{p}) \right\} + \text{const.}$$

Why is the positive sign in front of the $b^{\dagger}b$ term important? What is the physical interpretation of b and b^{\dagger} ?

e) Switch now to the Heisenberg picture

$$\phi_H(t, \boldsymbol{x}) = e^{iHt} \phi(\boldsymbol{x}) e^{-iHt}.$$

Show that

$$e^{iHt} a(\mathbf{p}) e^{-iHt} = a(\mathbf{p}) e^{-i\omega_{\mathbf{p}}t}, \qquad e^{iHt} a^{\dagger}(\mathbf{p}) e^{-iHt} = a^{\dagger}(\mathbf{p}) e^{i\omega_{\mathbf{p}}t}, e^{iHt} b(\mathbf{p}) e^{-iHt} = b(\mathbf{p}) e^{-i\omega_{\mathbf{p}}t}, \qquad e^{iHt} b^{\dagger}(\mathbf{p}) e^{-iHt} = b^{\dagger}(\mathbf{p}) e^{i\omega_{\mathbf{p}}t},$$
(1)

and therefore

$$\phi_H(t, \boldsymbol{x}) = \int \frac{d^3 p}{(2\pi)^3} \frac{1}{2\omega_{\boldsymbol{p}}} \left\{ a(\boldsymbol{p}) e^{-ipx} + b^{\dagger}(\boldsymbol{p}) e^{ipx} \right\}.$$