
Quantum Field Theory 1 – Problem set 2

Lectures: Jörg Jäckel

Jan Pawłowski

Tutorials: Malo Tarpin

Institut für Theoretische Physik, Uni Heidelberg

J.Jaeckel@thphys.uni-heidelberg.de

J.Pawlowski@thphys.uni-heidelberg.de

M.Tarpin@thphys.uni-heidelberg.de

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Suggested reading before solving these problems: Chapters 2.3 in the script and/or Chapters 2.3 to 2.4 of *Peskin & Schroeder*.

Problem 1: Commutation relations

For a real scalar field $\phi(x)$ the Lagrangian density is

$$\mathcal{L} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m^2 \phi^2.$$

The canonical momentum density is $\pi = \partial \mathcal{L} / \partial \dot{\phi} = \dot{\phi}$. The theory is quantized by promoting ϕ and π to operators in the Schrödinger picture with the commutation relations

$$\begin{aligned} [\phi(\mathbf{x}), \pi(\mathbf{y})] &= i\delta^{(3)}(\mathbf{x} - \mathbf{y}), \\ [\phi(\mathbf{x}), \phi(\mathbf{y})] &= [\pi(\mathbf{x}), \pi(\mathbf{y})] = 0. \end{aligned}$$

Introduce now the operators a and a^\dagger by

$$\begin{aligned} \phi(\mathbf{x}) &= \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2\omega_{\mathbf{p}}}} \{ a(\mathbf{p}) e^{i\mathbf{p}\mathbf{x}} + a^\dagger(\mathbf{p}) e^{-i\mathbf{p}\mathbf{x}} \}, \\ \pi(\mathbf{x}) &= -i \int \frac{d^3p}{(2\pi)^3} \sqrt{\frac{\omega_{\mathbf{p}}}{2}} \{ a(\mathbf{p}) e^{i\mathbf{p}\mathbf{x}} - a^\dagger(\mathbf{p}) e^{-i\mathbf{p}\mathbf{x}} \} \end{aligned}$$

and derive the commutation relations

$$[a(\mathbf{p}), a^\dagger(\mathbf{q})], \quad [a(\mathbf{p}), a(\mathbf{q})], \quad [a^\dagger(\mathbf{p}), a^\dagger(\mathbf{q})].$$

Problem 2: Complex scalar field: quantization

Consider the Lagrangian density for a free complex scalar field

$$\mathcal{L} = \partial_\mu \phi^* \partial^\mu \phi - m^2 \phi^* \phi.$$

a) Show that the canonical momenta of ϕ and ϕ^* are

$$\pi = \dot{\phi}^*, \quad \pi^* = \dot{\phi}.$$

and derive an expression for the Hamiltonian H .

- b) Proceed to quantization by promoting ϕ, ϕ^* and π, π^* to operators ϕ, ϕ^\dagger and π, π^\dagger (in the Schrödinger picture). What would you postulate as their commutation relations?
- c) Introduce now creation and annihilation operators by writing

$$\begin{aligned}\phi(\mathbf{x}) &= \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2\omega_{\mathbf{p}}}} \{a(\mathbf{p}) e^{i\mathbf{p}\mathbf{x}} + b^\dagger(\mathbf{p}) e^{-i\mathbf{p}\mathbf{x}}\}, \\ \pi(\mathbf{x}) &= -i \int \frac{d^3p}{(2\pi)^3} \sqrt{\frac{\omega_{\mathbf{p}}}{2}} \{-a^\dagger(\mathbf{p}) e^{-i\mathbf{p}\mathbf{x}} + b(\mathbf{p}) e^{i\mathbf{p}\mathbf{x}}\}.\end{aligned}$$

Why do we now need operators b, b^\dagger in addition to a, a^\dagger ? Convince yourself that the commutation relations

$$\begin{aligned}[a(\mathbf{p}), a^\dagger(\mathbf{q})] &= [b(\mathbf{p}), b^\dagger(\mathbf{q})] = (2\pi)^3 \delta^{(3)}(\mathbf{p} - \mathbf{q}), \\ [a(\mathbf{p}), a(\mathbf{q})] &= [b(\mathbf{p}), b(\mathbf{q})] = 0, \\ [a(\mathbf{p}), b(\mathbf{q})] &= [a(\mathbf{p}), b^\dagger(\mathbf{q})] = 0\end{aligned}$$

are consistent with the ones postulated in part b).

- d) Show that the Hamiltonian can be written as

$$H = \int \frac{d^3p}{(2\pi)^3} \omega_{\mathbf{p}} \{a^\dagger(\mathbf{p})a(\mathbf{p}) + b^\dagger(\mathbf{p})b(\mathbf{p})\} + \text{const.}$$

Why is the positive sign in front of the $b^\dagger b$ term important? What is the physical interpretation of b and b^\dagger ?

- e) Switch now to the Heisenberg picture

$$\phi_H(t, \mathbf{x}) = e^{iHt} \phi(\mathbf{x}) e^{-iHt}.$$

Show that

$$\begin{aligned}e^{iHt} a(\mathbf{p}) e^{-iHt} &= a(\mathbf{p}) e^{-i\omega_{\mathbf{p}}t}, & e^{iHt} a^\dagger(\mathbf{p}) e^{-iHt} &= a^\dagger(\mathbf{p}) e^{i\omega_{\mathbf{p}}t}, \\ e^{iHt} b(\mathbf{p}) e^{-iHt} &= b(\mathbf{p}) e^{-i\omega_{\mathbf{p}}t}, & e^{iHt} b^\dagger(\mathbf{p}) e^{-iHt} &= b^\dagger(\mathbf{p}) e^{i\omega_{\mathbf{p}}t},\end{aligned}\tag{1}$$

and therefore

$$\phi_H(t, \mathbf{x}) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2\omega_{\mathbf{p}}}} \{a(\mathbf{p}) e^{-i\mathbf{p}\mathbf{x}} + b^\dagger(\mathbf{p}) e^{i\mathbf{p}\mathbf{x}}\}.$$