

density matrix  $\rho = \rho^\dagger \geq 0, \text{Tr} \rho = 1$ .

special case: pure state  
( $\rightarrow$  a vector in Hilbert space)

$$\rho_\psi = P_\psi = |\psi\rangle\langle\psi|$$

$$\rightarrow \rho_\psi^2 = \rho_\psi \quad (\text{projection})$$

in general,  $\rho^2 < \rho$  if  $\rho$  is not pure.

DM useful whenever  $\mathcal{H} = \mathcal{H}_1 \otimes \mathcal{H}_2$   
regardless of the dimensions of  $\mathcal{H}_1$  and  $\mathcal{H}_2$ .

Today:  $\mathcal{H}_1 = \mathcal{H}_2 \cong \mathbb{C}^2$



The density matrix for measurements on qubit 1 is

$$\rho_1 = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

(maximal entropy)

e.g. for  $|\psi\rangle \quad A = A^\dagger \in M_2(\mathbb{C})$

$$\begin{aligned} \langle\psi| A \otimes 1 |\psi\rangle &= \frac{1}{2} (\langle 01| - \langle 10|) (A \otimes 1) (|01\rangle - |10\rangle) \\ &= \frac{1}{2} (\langle 01| A \otimes 1 |0\rangle - \langle 01| A \otimes 1 |1\rangle - \langle 10| A \otimes 1 |0\rangle \\ &\quad + \langle 10| A \otimes 1 |1\rangle) \\ &= \frac{1}{2} (\langle 01| A |0\rangle + \langle 11| A |1\rangle) = \frac{1}{2} (A_{00} + A_{11}) \end{aligned}$$

$$\begin{aligned} \langle 01| A \otimes 1 |10\rangle - \langle 01| A |1\rangle \cdot \langle 10| \otimes \langle 1| \rangle &= 0 \\ \langle 10| \otimes \langle 1| \rangle \cdot \langle 10| \otimes |0\rangle &= \frac{1}{2} \text{Tr}(A) = \text{Tr}(A \cdot \frac{1}{2} 1) \end{aligned}$$

$|\psi\rangle$  has the following properties.

$$(\vec{\sigma} \otimes 1) |\psi\rangle = -(1 \otimes \vec{\sigma}) |\psi\rangle$$

$\vec{\sigma} = (\sigma_x, \sigma_y, \sigma_z)$  vector of Pauli matrices, and  $|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, |1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$   
more generally

$$\langle\psi| \vec{a} \cdot \vec{\sigma} \otimes \vec{b} \cdot \vec{\sigma} |\psi\rangle = -\vec{a} \cdot \vec{b}, \quad \forall \vec{a}, \vec{b} \in \mathbb{R}^3.$$

Def. Bell states

$$|\Phi^+\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$

$$|\Phi^-\rangle = \frac{1}{\sqrt{2}} (|00\rangle - |11\rangle)$$

$$|\Psi^+\rangle = \frac{1}{\sqrt{2}} (|01\rangle + |10\rangle)$$

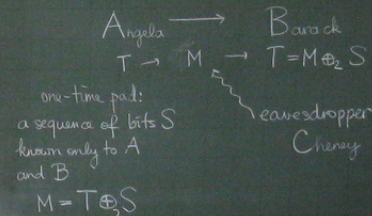
$$|\Psi^-\rangle = \frac{1}{\sqrt{2}} (|01\rangle - |10\rangle) \quad \text{singlet state}$$

These vectors are orthonormal in  $\mathbb{C}^2 \otimes \mathbb{C}^2$   
(exercise!)

They are all entangled states.

## M. Cryptography and Teleportation

M.1 Cryptography, here: key exchange



one-time pad:  
a sequence of bits S  
known only to A  
and B

Cheryl

Protocol: set up a source of particle pairs  
in the  $|\psi\rangle$  state.

A and B perform measurements of  $\sigma_1$  and  $\sigma_3$   
(1 and 3 are chosen randomly with prob  $\frac{1}{2}$ )

A and B make public which of the particles they measured  
 $\sigma_1$  or  $\sigma_3$ , but do not reveal the result of the measurement.

A and B select the subsequence where they measured  
the same  $\sigma_i$ .

A and B pick another subsequence randomly,  
to check if someone listened.  
The remaining sequence is M.

C has messed with the particle pairs so that they are now in the state

$$|\Psi\rangle = |00\rangle \otimes |e_{00}\rangle + |01\rangle \otimes |e_{01}\rangle + |10\rangle \otimes |e_{10}\rangle + |11\rangle \otimes |e_{11}\rangle$$

Verification step (\*): A and B check the property

$$\begin{aligned} \sigma_x \otimes \sigma_x |\Psi\rangle &= -|\Psi\rangle \\ \sigma_y \otimes \sigma_y |\Psi\rangle &= -|\Psi\rangle \\ (\sigma_z \otimes \sigma_z) |\Psi\rangle &= |0\rangle|e_{00}\rangle - |0\rangle|e_{01}\rangle - |1\rangle|e_{10}\rangle + |1\rangle|e_{11}\rangle \\ &= -|\Psi\rangle = -(|0\rangle|e_{00}\rangle - |0\rangle|e_{01}\rangle - |1\rangle|e_{10}\rangle + |1\rangle|e_{11}\rangle) \end{aligned}$$

The vectors adding up to  $|\Psi\rangle$  are orthogonal  $\Rightarrow |e_{00}\rangle = 0$  and  $|e_{11}\rangle = 0$

$$\Rightarrow |\Psi\rangle = |10\rangle|e_{10}\rangle + |11\rangle|e_{11}\rangle$$

$$\sigma_x |0\rangle = |1\rangle \quad \sigma_x |1\rangle = |0\rangle$$

$$(\sigma_x \otimes \sigma_x) |\Psi\rangle = |11\rangle|e_{10}\rangle + |10\rangle|e_{11}\rangle$$

$$\Rightarrow |\Psi\rangle = -|10\rangle|e_{11}\rangle - |11\rangle|e_{10}\rangle$$

$$\Rightarrow |e_{11}\rangle = -|e_{10}\rangle$$

$$\Rightarrow |\Psi\rangle = (|10\rangle - |11\rangle) \otimes |e_{10}\rangle = (*)$$

(\*)  $\Rightarrow$  C cannot change the bit sequence while fulfilling both the  $\sigma_x$  and  $\sigma_z$  singlet check

C cannot learn anything either, because:

$$|\psi\rangle \otimes |u\rangle \xrightarrow{\text{unitary}} |\psi\rangle \otimes |v\rangle \quad (\langle u|u\rangle = 1)$$

$$|\phi\rangle \otimes |u\rangle \xrightarrow{\text{unitary}} |\phi\rangle \otimes |v'\rangle$$

if  $\langle \phi|\psi\rangle \neq 0$  then  $|v\rangle = |v'\rangle$

$$\begin{aligned} \langle \phi|\otimes\langle u|(|\psi\rangle \otimes |u\rangle) &= \langle \phi|\psi\rangle \cdot \langle u|u\rangle = \langle \phi|\psi\rangle \\ &= \langle \phi|\otimes\langle v'|(|\psi\rangle \otimes |v\rangle) = \langle \phi|\psi\rangle \cdot \langle v'|v\rangle \end{aligned}$$

$\Rightarrow \langle v'|v\rangle = 1$   
 $\Rightarrow |v'\rangle = |v\rangle$   
because  $\|v'\| = \|v\| = 1$

## M.2 Teleportation.

How to send an unknown quantum state  $|\psi\rangle$  from A to B?

Suppose, A and B have a  $|\phi^\pm\rangle$  pair in common

$$\frac{1}{\sqrt{2}}(|00\rangle \pm |11\rangle)$$

$$|\phi^\pm\rangle_{AB} \otimes |\psi\rangle_C$$

A measures his part in the Bell basis  $|\phi^\pm\rangle, |\psi^\pm\rangle$  and sends the result to B. Depending on A's result, B applies the following operator to his qubit!

$$\begin{aligned} |\phi^+\rangle &\rightarrow I \\ |\phi^-\rangle &\rightarrow \sigma_z \\ |\psi^+\rangle &\rightarrow \sigma_x \\ |\psi^-\rangle &\rightarrow \sigma_x \sigma_z \end{aligned}$$

Then B's qubit is in the state  $|\psi\rangle$

$$|\psi\rangle = a|0\rangle + b|1\rangle$$

$$\begin{aligned} |\psi\rangle_C |\phi^+\rangle_{AB} &= (a|0\rangle_C + b|1\rangle_C) \frac{1}{\sqrt{2}} (|00\rangle_{AB} + |11\rangle_{AB}) \\ &= \frac{1}{\sqrt{2}} [a|00\rangle_{CB} + a|01\rangle_{CB} + b|10\rangle_{CB} + b|11\rangle_{CB}] \\ &= \frac{1}{\sqrt{2}} [a|0\rangle_C \otimes |0\rangle_B + a|0\rangle_C \otimes |1\rangle_B + b|1\rangle_C \otimes |0\rangle_B + b|1\rangle_C \otimes |1\rangle_B] \end{aligned}$$

Use that

$$|00\rangle = \frac{1}{\sqrt{2}} (|\phi^+\rangle + |\phi^-\rangle) \quad |11\rangle = \frac{1}{\sqrt{2}} (|\phi^+\rangle - |\phi^-\rangle)$$

$$|01\rangle = \frac{1}{\sqrt{2}} (|\psi^+\rangle + |\psi^-\rangle) \quad |10\rangle = \frac{1}{\sqrt{2}} (|\psi^+\rangle - |\psi^-\rangle)$$

$$\begin{aligned} |\psi\rangle_C |\phi^+\rangle_{AB} &= |\phi^+\rangle_A \cdot \frac{1}{2} (a|0\rangle_B + b|1\rangle_B) \\ &\quad + |\phi^-\rangle_A \cdot \frac{1}{2} (a|0\rangle_B - b|1\rangle_B) \\ &\quad + |\psi^+\rangle_A \cdot \frac{1}{2} (b|0\rangle_B + a|1\rangle_B) \\ &\quad + |\psi^-\rangle_A \cdot \frac{1}{2} (-b|0\rangle_B + a|1\rangle_B) \\ &= |\phi^+\rangle_A \otimes \frac{1}{2} |\psi\rangle_B + |\phi^-\rangle_A \otimes \frac{1}{2} \sigma_z |\psi\rangle_B \\ &\quad + |\psi^+\rangle_A \otimes \frac{1}{2} \sigma_x |\psi\rangle_B + |\psi^-\rangle_A \otimes \frac{1}{2} (\sigma_x \sigma_z) |\psi\rangle_B \end{aligned}$$