Solutions to be handed in in the tutorials of Week 10 (21st December

## Ex. H16: Translation operator

The translation operator  $T(\mathbf{a})$  is defined by

$$T(\mathbf{a}) = \exp\left(-\frac{i}{\hbar}\mathbf{a}\cdot\mathbf{P}\right),$$

where  $\mathbf{a} \in \mathbb{R}^3$  and  $\mathbf{P} = (P_1, P_2, P_3)$  is the momentum operator in three dimensions.

- a) Show that  $T(\mathbf{a})T(\mathbf{b}) = T(\mathbf{a} + \mathbf{b})$ .
- b) Show that, given the position operator  $\mathbf{X} = (X_1, X_2, X_3)$ , the following holds:

$$T(\mathbf{a})^{\dagger} \mathbf{X} T(\mathbf{a}) = \mathbf{X} + \mathbf{a} \mathbb{1}.$$

## Ex. H17: Optical theorem

Prove the *optical theorem* of scattering theory:

$$\sigma_{\rm tot} = \frac{4\pi}{k} \operatorname{Im} f(\theta, \phi) \bigg|_{\cos \theta = 1}.$$

Here  $\sigma_{\text{tot}}$  is the total (integrated) cross-section, k is the wavenumber and f is the scattering amplitude.

Instructions: Integrate the continuity equation for the probability current over a sphere of radius r around the scattering centre, and consider the limit  $r \to \infty$ . Note furthermore that  $f(\theta, \phi)$  is not dependent on  $\phi$  for  $\cos \theta = 1$ .

## Ex. H18: Partial wave expansion

Consider the scattering of an incident plane wave travelling in the positive z-direction on a spherically symmetric, short-ranged potential. The aim of the exercise is to expand the scattering amplitude  $f(\theta)$  in spherical harmonics (due to the spherical symmetry f is not dependent on  $\phi$ ) and to represent it in the following form:

$$f(\theta) = \sum_{l=0}^{\infty} (2l+1) f_l P_l(\cos \theta).$$

(3 points)

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(5 \text{ points})
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(8 points)

Here  $P_l$  is the *l*th Legendre polynomial given by  $P_l = P_l^0$  with the associated Legendre function  $P_l^m$  already familiar from Ex. P4. The coefficient  $f_l$  (the so-called *partial wave amplitude*) may be calculated from the phase shift  $\delta_l$  between the incident plane wave and the outgoing spherical wave.

a) Start by considering the expansion of a general  $\phi$ -independent solution of the Schrödinger equation:

$$\psi(r,\theta) = \sum_{l=0}^{\infty} c_l \frac{u_l(r)}{r} Y_l^0(\theta) \equiv \sum_{l=0}^{\infty} (2l+1) a_l \frac{u_l(r)}{r} P_l(\cos\theta).$$

Show that, for  $r \to \infty$ , the radial part of the wavefunction  $u_l(r)$  can be written in the following form:

$$u_l(r) = \frac{1}{k} \sin\left(kr - \frac{l\pi}{2} + \delta_l\right).$$

Here the value of  $\delta_l$  depends on l, and we have chosen to split off the constant  $\frac{l\pi}{2}$  from each of the terms (without loss of generality).

b) Show that  $\delta_l = 0$  for an incident plane wave travelling in the positive zdirection with wavenumber k, und determine the corresponding coefficients  $a_l$ .

*Hint:* The following identity holds

$$e^{ikr\cos\theta} = \sum_{l=0}^{\infty} i^l (2l+1) j_l(kr) P_l(\cos\theta)$$

with the spherical Bessel functions  $j_l(x)$ , that are related to the usual Bessel functions by  $j_l(x) = \sqrt{\pi/2x} J_{l+1/2}(x)$ . The asymptotic behaviour of  $J_{\alpha}$  is

$$J_{\alpha}(x) \approx \sqrt{\frac{2}{\pi x}} \cos\left(x - \frac{\alpha \pi}{2} - \frac{\pi}{4}\right) \qquad (x \to \infty).$$

c) Show that, for the complete scattering solution, the partial wave amplitudes in the decomposition  $\bigcirc$  are given by

$$f_l = \frac{1}{2ik} \left( e^{2i\delta_l} - 1 \right).$$

d) Show that the total cross-section  $\sigma_{tot}$  can be described as follows:

$$\sigma_{\rm tot} = \frac{4\pi}{k^2} \sum_{l=0}^{\infty} (2l+1) \sin^2 \delta_l.$$

*Hint:* Use the orthogonality of the Legendre polynomials:

$$\int_{-1}^{1} dx P_l(x) P_{l'}(x) = \frac{2\delta_{l,l'}}{2l+1}$$