Solutions to be handed in and class exercises discussed in the tutorials of Week 12 (18/01/08)

Ex. H22: Kronig-Penney model

A one-dimensional periodic potential is given with period l, thus V(x) = V(x+l).

a) Prove Bloch's Theorem: there exists a basis of energy eigenstates, whose wavefunctions assume the form $\psi_{\kappa}(x) = e^{-i\kappa x}u_{\kappa}(x)$ with periodic functions $u_{\kappa}(x) = u_{\kappa}(x+l)$.

Hint: The law of simultaneous diagonalisability of commuting operators holds not only for Hermitian operators. Apply it here to the unitary translation operator T(l) (see Ex. H16) and the Hamiltonian.

Consider now a particle with energy E in the range $0 < E < V_0$, moving in the following potential with period l = a + b:

$$V(x) = \begin{cases} 0, & 0 \le x \le a \\ V_0, & a < x < a + b \end{cases}$$

(periodically continued for all $x \in \mathbb{R}$). This is a simple one-dimensional model for an electron in a solid. The positively charged ions are located on a crystal lattice at a separation of a + b, separated by potential barriers of height V_0 and width b. The energy spectrum shows a band structure, in which certain ranges of energy are allowed and others are forbidden.

b) Show that the following condition holds for the energy:

$$-1 \le \cos(ka)\cosh(qb) + \frac{q^2 - k^2}{2qk}\sin(ka)\sinh(qb) \le 1,$$

where

$$k = \frac{\sqrt{2mE}}{\hbar}, \qquad q = \frac{\sqrt{2m(V_0 - E)}}{\hbar}.$$

c) Consider the limiting case

 $b \to 0, \qquad V_0 \to \infty, \qquad V_0 b = {\rm const.}$

and energies $E \ll V_0$. Show that

$$-1 \le \cos(ka) + \gamma \frac{\sin(ka)}{ka} \le 1,$$

where $\gamma = mabV_0/\hbar^2$. Determine the allowed ranges for ka graphically for $-4\pi \leq ka \leq 4\pi$ and $\gamma = 5$.

(5 points)

Ex. H23: Three-dimensional harmonic oscillator

(5 points)

The potential of the three-dimensional harmonic oscillator is

$$V = \frac{m\omega^2}{2}r^2, \qquad \text{mit } r = |\mathbf{x}|.$$

- a) The corresponding Schrödinger equation separates in Cartesian coordinaten. State the possible energy eigenvalues, and determine the degeneracy of the nth energy eigenstate.
- b) Since we are dealing with a central potential, the Schrödinger equation also separates in spherical polar coordinates, for which the solutions of the angular part are already known to be the spherical harmonic functions $Y_l^m(\theta, \phi)$:

$$\psi(r,\theta,\phi) = R(r) Y_l^m(\theta,\phi).$$

It now remains to determine the radial part R(r). Set up the radial equation, and discuss the limits $r \to 0$ and $r \to \infty$ with the ansatz

$$u(r) = r R(r) = r^{\delta} e^{-\gamma r^2} g(r)$$

(where $g(0) \neq 0$ and |g(r)| does not either grow or fall off too quickly for $r \to \infty$). Determine δ and γ .

- c) Set up an equation to determine g(r).
- d) Give a reason why, with the Ansatz

$$g(r) = \sum_{n} a_n r^n,$$

the power series must terminate at a finite value of n. Determine thereby the energy spectrum and the degeneracies of energies, and compare with a).

Ex. H24: Landau levels

Consider an electron (whose spin is to be neglected) moving in a constant magnetic field pointing in the z-direction, $\mathbf{B} = (0, 0, B)$.

- a) Show that the vector potential can be chosen as $\mathbf{A} = (0, Bx, 0)$.
- b) Show that the Hamiltonian

$$H = \frac{1}{2m} \left(\mathbf{P} + \frac{e}{c} \mathbf{A} \right)^2$$

commutes with both the y- and z-components of the momentum operators. The eigenstates of H can thus be determined with the separation of variables $\psi(x, y, z) = f(x)g(y)h(z)$. State g and h. We wish to choose h(z) such that eigenvalue of P_z is precisely 0.

- c) Derive an eigenvalue equation for f(x), and trace this back to the eigenvalue equation of the simple harmonic oscillator.
- d) Read off the energy spectrum. The energy levels resulting from the equation for f are called *Landau levels*.