

Proton scattering on an electron gas

Janos Polonyi

University of Strasbourg

J.P. K. Zazoua, 2012

1. Collision in many-body systems
with target-beam entanglement and final state interactions
2. Closed Time Path method
to handle system-environment entanglement
3. Reduction formulae
for transition probabilities
4. Perturbation expansion
5. $p+e \rightarrow p+e$
 - 5.1 entanglement
 - 5.2 decoherence
 - 5.3 real and virtual excitations
6. Summary

Collision in many-body systems

$e(p_e) + p(p) \rightarrow e(q_e) + p(q)$ on electron gas

Transition amplitude: $\mathcal{A}_n = \langle \Psi_{n, \hat{\mathbf{q}}_e}^{gas} | a_p(\mathbf{q}) S a_p^\dagger(\mathbf{p}) | \Psi_0^{gas} \rangle$.

with or without identifying the state of the participating electron

Transition probability: $P = \sum_n |\mathcal{A}_n|^2$

$$\begin{aligned}
 &= \sum_n \langle \Psi_0^{gas} | a_p(\mathbf{p}) S^\dagger a_p^\dagger(\mathbf{q}) | \Psi_{n, \hat{\mathbf{q}}_e}^{gas} \rangle \langle \Psi_{n, \hat{\mathbf{q}}_e}^{gas} | a_p(\mathbf{q}) S a_p^\dagger(\mathbf{p}) | \Psi_0^{gas} \rangle \\
 &= \langle \Psi_0^{gas} | a_p(\mathbf{p}) S^\dagger a_p^\dagger(\mathbf{q}) \widehat{a_e^\dagger(\mathbf{q}_e) a_e(\mathbf{q}_e)} a_p(\mathbf{q}) S a_p^\dagger(\mathbf{p}) | \Psi_0^{gas} \rangle \\
 &= \text{Tr} \left[\underbrace{a_p^\dagger(\mathbf{q}) a_e^\dagger(\mathbf{q}_e) a_e(\mathbf{q}_e) a_p(\mathbf{q})}_O S \underbrace{a_p^\dagger(\mathbf{p}) | \Psi_0^{gas} \rangle \langle \Psi_0^{gas} | a_p(\mathbf{p}_2)}_{\rho_i} S^\dagger \right] \\
 &= \text{Tr}[OS\rho_i S^\dagger] \quad \leftarrow \quad \text{CTP}
 \end{aligned}$$

1. Collision in many body systems
2. Closed Time Path method ←
3. Reduction formulae
4. Perturbation expansion
5. $p+e \rightarrow p+e$
6. Summary

Closed Time Path method

An incomplete list

- QED (Schwinger 1961)
- relaxation in many-body systems (Kadanoff, Baym 1962)
- density matrix by path integral (Feynman, Vernon 1963, Diosi 1986)
- perturbation expansion for retarded Green-functions (Keldysh, 1964)
- time symmetric quantum mechanics
(Aharonov, Bergmann, Lebowitz 1964)
- thermal field theory (Umezawa 1975)
- consistent histories (Griffith 1984, Gell-Mann, Hartle 1993)
- equations of motion for the expectation value of local operators
(Jordan 1986)
- non-equilibrium processes (Calzetta, Hu 1988)
- classical mechanics (Polonyi 2012)

Closed Time Path

Open systems

Expectation value as opposed to transition amplitude

$$\bar{\mathcal{O}} = \text{Tr}[\mathcal{O}U(t_f, t_i)\rho_i U^\dagger(t_f, t_i)],$$

Generator functional:

$$\begin{aligned} e^{iW[\hat{j}]} &= \text{Tr} T[e^{-i \int dt' [H(t') - j^+(t')x(t')]}] \rho_i T^* e^{i \int dt' [H(t') + j^-(t')x(t')]} \\ &= \int D[\hat{x}] e^{iS[x^+] - iS^*[x^-] + i \int dt \hat{j}(t)\hat{x}(t)} \\ &= \int D[\hat{x}] e^{iS_{CTP}[\hat{x}] + i \int dt \hat{j}(t)\hat{x}(t)} \end{aligned}$$

reduplication of the source $\hat{j} = (j^+, j^-)$ and
the degrees of freedom $\hat{x} = (x^+, x^-)$

Closed Time Path

Entanglement

System: ϕ , environment: χ , $S[\phi, \chi] = S_s[\phi] + S_e[\phi, \chi]$

$$\begin{aligned} e^{iW[\hat{J}]} &= \int D[\hat{\phi}] D[\hat{\chi}] e^{iS[\phi^+, \chi^+] - iS[\phi^-, \chi^-] + i\hat{J} \cdot \hat{\phi}} \\ &= \int D[\hat{\phi}] e^{iS_s[\phi^+] - iS_s[\phi^-] - iS_I[\phi^+, \phi^-] + i\hat{J} \cdot \hat{\phi}} \end{aligned}$$

Influence functional (Feynman):

$$e^{iS_I[\phi^+, \phi^-]} = \int D[\hat{\chi}] e^{iS_e[\phi^+, \chi^+] - iS_e[\phi^-, \chi^-]}$$

S_I : system-environment entanglement and mixed system state

1. Collision in many body systems
2. Closed Time Path method
3. Reduction formulae ←
4. Perturbation expansion
5. $p+e \rightarrow p+e$
6. Summary

Reduction formulae

Generator functional

$$\begin{aligned}e^{iW[\hat{J}, \hat{\eta}, \hat{\eta}]} &= \text{Tr} [S(\eta^+, \bar{\eta}^+, j^+) \rho_i S^\dagger(-\eta_-, -\bar{\eta}_-, -j^-)] \\&= \int D[\hat{\psi}] D[\hat{\bar{\psi}}] D[\hat{A}] e^{iS[\hat{A}, \hat{\psi}, \hat{\bar{\psi}}] + i \int dx [\hat{\eta} \hat{\psi} + \hat{\bar{\psi}} \hat{\eta} + \hat{J} \hat{A}]} \\S[\hat{A}, \hat{\psi}, \hat{\bar{\psi}}] &= \int dx \hat{\bar{\psi}} (\hat{G}^{-1} - e \hat{A}) \hat{\psi} + \frac{1}{2} \int dx \hat{A} \hat{D}_0^{-1} \hat{A} \\ \hat{G}^{-1} &= \begin{pmatrix} i\partial - m + i\epsilon & 0 \\ 0 & -\gamma^0 (i\partial m + i\epsilon)^\dagger \gamma^0 \end{pmatrix} + \hat{G}_{BC}^{-1} \\ \hat{D}_{0\mu\nu}^{-1} &= g_{\mu\nu} \begin{pmatrix} \square + i\epsilon & 0 \\ 0 & -\square + i\epsilon \end{pmatrix} + \hat{D}_{BC}^{-1},\end{aligned}$$

Final conditions on the trajectories:

$$\begin{aligned}\psi^+(t_f, \mathbf{x}) &= \psi^-(t_f, \mathbf{x}), \\ \bar{\psi}^+(t_f, \mathbf{x}) &= \bar{\psi}^-(t_f, \mathbf{x}), \\ A_\mu^+(t_f, \mathbf{x}) &= A_\mu^-(t_f, \mathbf{x}),\end{aligned}$$

Reduction formulae

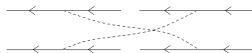
Transition probability

Electron state registered: $e+p \rightarrow e + p$ elastic collision

$$\begin{aligned}
 P = & Z_e^{-2} Z_p^{-2} \int dx_1^+ dx_2^+ dy_1^+ dy_2^+ dx_1^- dx_2^- dy_1^- dy_2^- \\
 & e^{iq_1(y_1^+ - y_1^-) + iq_2(y_2^+ - y_2^-) - ip_1(x_1^+ - x_1^-) - ip_2(x_2^+ - x_2^-)} \\
 & (\bar{u}_{ef} K)_{y_e^+}^+ (\bar{u}_{pf} K)_{y_p^+}^+ (K u_{ef})_{y_e^-}^- (K u_{pf})_{y_p^-}^- \\
 & (K u_{ei})_{x_e^+}^+ (K u_{pi})_{x_p^+}^+ (\bar{u}_{ei} K)_{x_e^-}^- (\bar{u}_{pi} K)_{x_p^-}^- \\
 & \delta^8 W \\
 & \hline
 & \delta\eta_p^-(y_p^-) \delta\eta_e^-(y_e^-) \delta\bar{\eta}_p^-(x_p^-) \delta\bar{\eta}_e^-(x_e^-) \delta\bar{\eta}_e^+(y_e^+) \delta\bar{\eta}_p^+(y_p^+) \delta\eta_e^+(x_e^+) \delta\eta_p^+(x_p^+)
 \end{aligned}$$



direct



entanglement

Reduction formulae

Transition probability

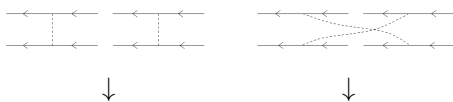
$$\begin{aligned}P &= -\frac{ie^2}{2M^2} [p^\mu q^\nu + q^\mu p^\nu + g^{\mu\nu} (M^2 - pq)] D_{0\mu\nu}^{+-}(r) \\P_d &= \frac{e^4}{4m^2 M^2} \{ 2\Re[(p_1 D^{++} p_2)(q_1 D^{++*} q_2)] + 2\Re[(p_1 D^{++} q_2)(q_1 D^{++*} p_2)] \\&\quad + \text{tr}[D^{++} D^{++*}](m^2 - p_1 q_1)(M^2 - p_2 q_2) \\&\quad + 2(M^2 - p_2 q_2)\Re[p_1 D^{++} D^{++*tr} q_1] \\&\quad + 2(m^2 - p_1 q_1)\Re[p_2 D^{++tr} D^{++*} q_2] \}, \\P_e &= \frac{e^4}{4m^2 M^2} \{ 2\Re[(p_1 D^{+-} p_2)(q_1 D^{+-*} q_2)] + 2\Re[(p_1 D^{+-} q_2)(q_1 D^{+-*} p_2)] \\&\quad + \text{tr}[D^{+-} D^{+-\dagger}](m^2 - p_1 q_1)(M^2 - p_2 q_2) \\&\quad + 2(M^2 - p_2 q_2)\Re[p_1 D^{+-} D^{+-\dagger} q_1] \\&\quad + 2(m^2 - p_1 q_1)\Re[q_2 D^{-+\dagger}(-r)D^{-+}(-r)p_2] \},\end{aligned}$$

1. Collision in many body systems
2. Closed Time Path method
3. Reduction formulae
4. Perturbation expansion ←
5. $p+e \rightarrow p+e$
6. Summary

Perturbation expansion

t-channel resummation

closing electron lines:
 G^{+-} on mass shell



D^{+-}

mass shell \rightarrow vanishing

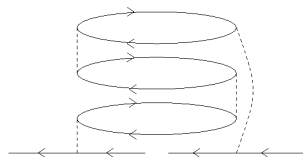
$D^{++}\Pi^{+-}D^{--}$

$D^{+-}\Pi^{-+}D^{+-}$

vanishing

non-vanishing contributions:

$D^{++}\Pi^{+-}D^{--}\Pi^{-+}D^{++}\Pi^{+-}D^{--}$



Schwinger-Dyson resummation: $\hat{D} = \frac{1}{\hat{D}_0^{-1} - \hat{\Pi}}$

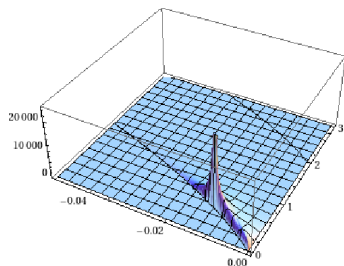
for homogeneous electron gas, $k_F \neq 0$

1. Collision in many body systems
2. Closed Time Path method
3. Reduction formulae
4. Perturbation expansion
5. $p+e \rightarrow p+e$, entanglement, real and virtual excitations ←
6. Summary

$p+e \rightarrow p+e$

Quasi-particles

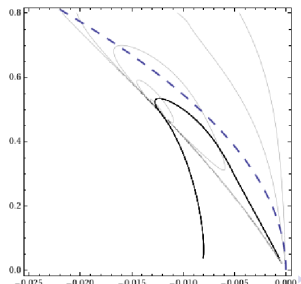
Spectral function, $\Im D_\ell^{+-}(\omega, \mathbf{r})$,
of dressed electric photons
on the plane $(\frac{\omega}{k_F}, \frac{|\mathbf{r}|}{k_F})$



Solid line: Electric dispersion relation

$$\Re D^{+-}(\omega, \mathbf{r}) = 0$$

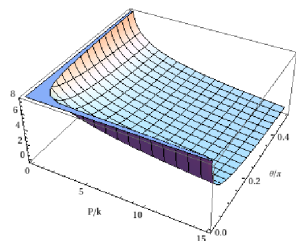
Thin lines: The contour plot of $\Im D_\ell^{+-}$



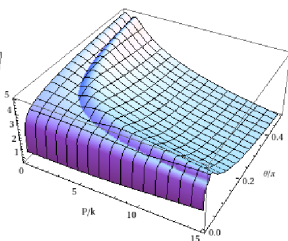
$p+e \rightarrow p+e$

Cross section

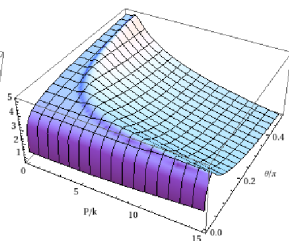
$\frac{d^2\sigma}{dPd\Theta}$ as function of $\frac{P}{k_F}$ and $\frac{\Theta}{\pi}$:



$\log \sigma_{vac}$



$\log \sigma_{direct}$



$\log(\sigma_{direct} + \sigma_{entangl})$

σ_{vac} OR σ_{direct}



factorizable

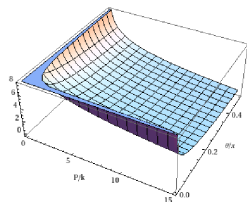
$\sigma_{entangl}$



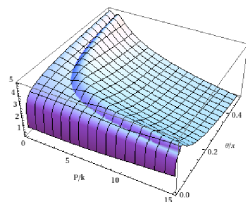
photon-charge
entanglement

$p \dagger e \rightarrow p \dagger e$

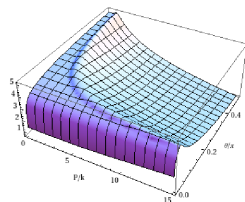
Spectral strengths at $\Theta = 0.5\pi$



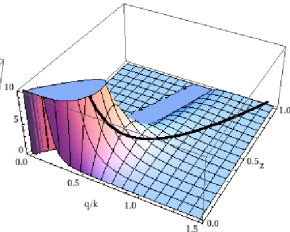
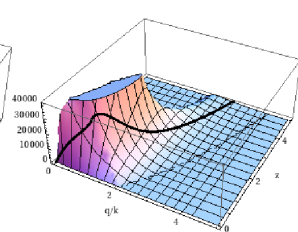
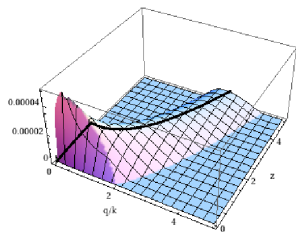
$\log \sigma_{vac}$



$\log \sigma_{direct}$



$\log(\sigma_{direct} + \sigma_{entangl})$



$z = \frac{m\omega}{k_F^2}$ current

electric photon

magnetic photon

spectral functions

$p \leftarrow e \rightarrow p \leftarrow e$

Effective theory of non-relativistic charges

$$Z = \int D[\hat{A}] D[\hat{\mathbf{x}}] e^{iS[\mathbf{x}^+] - iS[\mathbf{x}^-] + \frac{i}{2} \hat{A} \hat{D}_0^{-1} \hat{A} - ij^+ A^+ + ij^- A^-} = \int D[\hat{\mathbf{x}}] e^{iS_{\text{eff}}[\hat{\mathbf{x}}]}$$

$$e^{iW_0[\hat{j}]} = \int D[\hat{A}] e^{\frac{i}{2} \hat{A} \hat{D}_0^{-1} \hat{A} + i\hat{j} \hat{A}} = e^{-\frac{i}{2} \hat{j} \hat{D}_0 \hat{j}}$$

$$S_{\text{eff}}[\hat{\mathbf{x}}] = S[\mathbf{x}^+] - S[\mathbf{x}^-] + W_0[-ej[\mathbf{x}^+], ej[\mathbf{x}^-]]$$

$$\text{Parametrization: } x^\pm = x \pm \frac{x^d}{2}, j^\pm = j \pm \frac{j^d}{2}$$

Real part: expectation values

$$\Re S_{\text{eff}} = S[\mathbf{x}^+] - S[\mathbf{x}^-] - \frac{e^2}{2} [j D^a j^d - j^d D^r j]$$

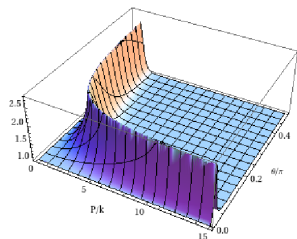
Imaginary part: decoherence, consistence

$$\Im S_{\text{eff}} = -\frac{e^2}{2} j^d D^i j^d$$

$p \dagger e \rightarrow p \dagger e$

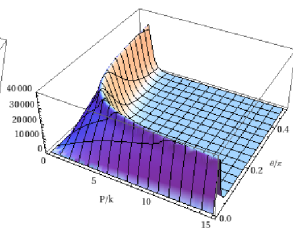
Entanglement and decoherence

$(\frac{P}{K_F}, \frac{\Theta}{\pi})$ plane:



$$\frac{\sigma_{direct} + \sigma_{entangl}}{\sigma_{direct}}$$

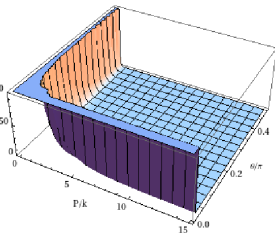
weight of entanglement



$$D_\ell^i$$

inverse width of $\rho(x + \frac{x^d}{2}, x - \frac{x^d}{2})$ in x^d

$$|e^{iS_{eff}}| = e^{-\frac{e^2}{2} j^d D^i j^d}$$



$$D_t^i$$

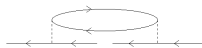
$$p + e \rightarrow p + e$$

Real vs. virtual excitations

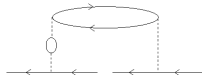
Bloch, Nordsieck, Kinoshita, Lee, Nauenberg:

Cancellation of collinear divergences between real and virtual soft photons
(need of degenerate perturbation expansion)

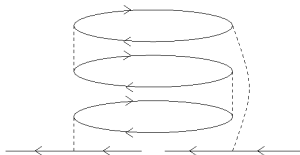
Leading order:



Virtual excitations:
within S or S^\dagger



Real excitations:
at t_f



Is there a partial cancellation between real and virtual degenerate excitations?

$p \dagger e \rightarrow p \dagger e$

Real vs. virtual excitations

$$\hat{\Pi} = \hat{\Pi}_d + \hat{\Pi}_{nd}, \quad \hat{D}_d = \frac{1}{\hat{D}_0^{-1} - \hat{\Pi}_d} = \begin{pmatrix} D_F & 0 \\ 0 & -D_F^\dagger \end{pmatrix}$$

$$\hat{D} = \frac{1}{\hat{D}_d^{-1} - \hat{\Pi}_{nd}}$$

$$\hat{D}^{+-} = \sum_{n=0}^{\infty} [\hat{D}_d (\hat{\Pi}_{nd} \hat{D}_d)^{2n+1}]^{+-} = \hat{D}_d \hat{\Pi}_{nd} \hat{D}_d + \sum_{n=1}^{\infty} [\hat{D}_d (\hat{\Pi}_{nd} \hat{D}_d)^{2n+1}]^{+-}$$



virtual

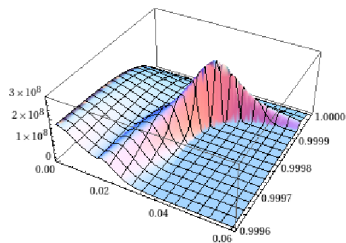


real

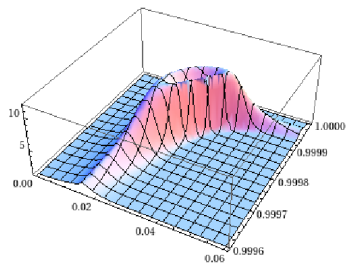
$p+e \rightarrow p+e$

Real vs. virtual excitations

$p \rightarrow p$ inelastic collision: in the plane $(\frac{p}{M}, \frac{q}{p})$ for $\theta = 0.4 \cdot 10^{-4} \pi$



σ



$\frac{\sigma_{virt}}{\sigma}$

strong cancellation between
soft real and virtual excitations

Summary

1. Open systems \leftrightarrow QFT remade in CTP
2. $e + p \rightarrow e + p$
 - 2.1 Solution of the CTP Schwinger-Dyson equation
 - 2.2 Relation between quasi-particles, entanglement and decoherence
 - 2.3 Reduction formulae for transition probability is unitarity protected
3. (Partial) Cancellations among degenerate excitations
An extension of the Kinoshita-Lee-Nauenberg theorem?