

c-field descriptions of nonequilibrium polariton fluids

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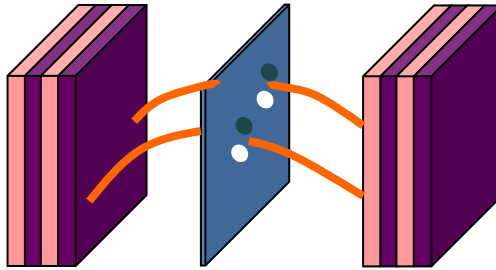
Iacopo Carusotto, Vincenzo Savona



polariton characteristics

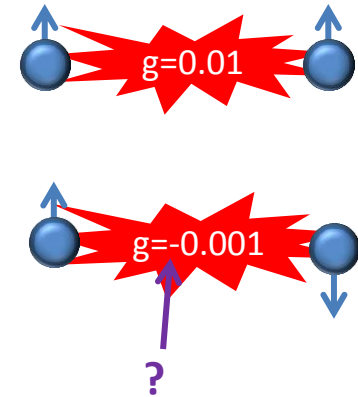
Dispersion

2D / 1D / 0D



1 meV
inhomogeneous broadening
(but long life time)

Interactions



- short range ($R \approx 0.01\mu\text{m}$)
- polarization dependent (bi-exciton “feshbach” resonance?)
- weak: $gm/\hbar^2 \ll 1$
 \Rightarrow mean field theory OK

[C. Ciuti et al. PRB 1998]

[Kwong et al. PRB 2007]

[M.W. PRB 2007]

1 eV

$1\mu\text{m}^{-1}$

0.x meV

\Rightarrow Polariton life time 1-10's ps

\Rightarrow nonequilibrium

overview

c-field models
(Gross-Pitaevskii+
&truncated Wigner)



superfluidity



coherence (BEC)

mean field theory (GPE)

$$\hat{\psi} \rightarrow \langle \hat{\psi} \rangle = \psi$$

$$i \frac{\partial}{\partial t} \psi = \underbrace{\left(\varepsilon(-i\nabla) + V_{ext} + g|\psi|^2 \right)}_{\text{usual GPE}} \underbrace{\psi}_{\text{finite life time (dissipation @all scales)}} - i \frac{\gamma}{2} \psi + \underbrace{F_L(x,t)}_{\text{external pumping laser}}$$

Resonant excitation

$$F_L(x,t) = F_L e^{i(k_L x - \omega_L t)}$$

creates a plane wave **coherent** polariton **quantum fluid**

inherited from laser

super?

$$\psi = \sqrt{n_0} e^{i(k_L x - \omega_L t)}$$

↑ ↑ ↑
independently tunable



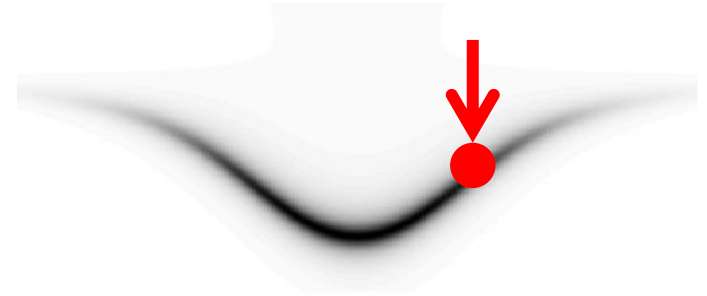
equilibrium condensate

$$\psi = \sqrt{n_0} e^{i(k_L x - \mu t)}$$

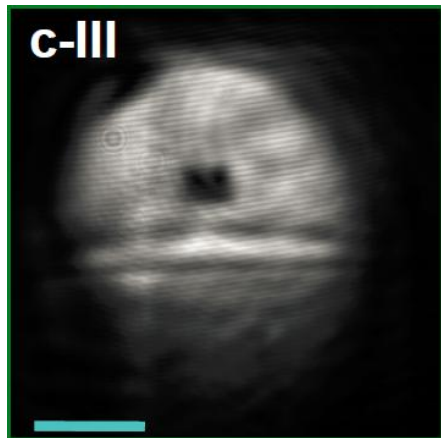
where

$$\mu = \frac{k_L^2}{2m} + gn_0$$

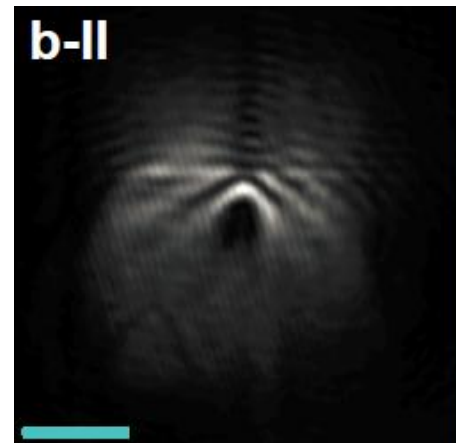
superfluidity



- persistent currents : trivial (driven by laser)
- scattering off weak defects: nontrivial \square Landau criterion



$v < v_c$: flow without scattering



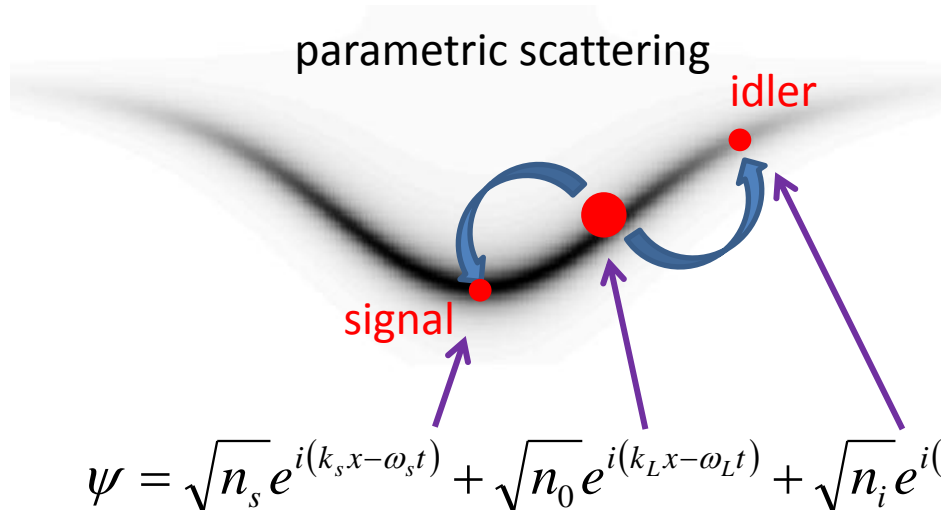
$v > v_c$: Cerenkov radiation

- great similarity with equilibrium GPE

[Carusotto and Ciuti, PRL 2004; Amo et al. Nat. Phys. 2009]

dynamical instability

$$F_L(x, t) = F_L e^{i(k_L x - \omega_L t)} \xrightarrow{\text{always ??}} \psi = \sqrt{n_0} e^{i(k_L x - \omega_L t)}$$

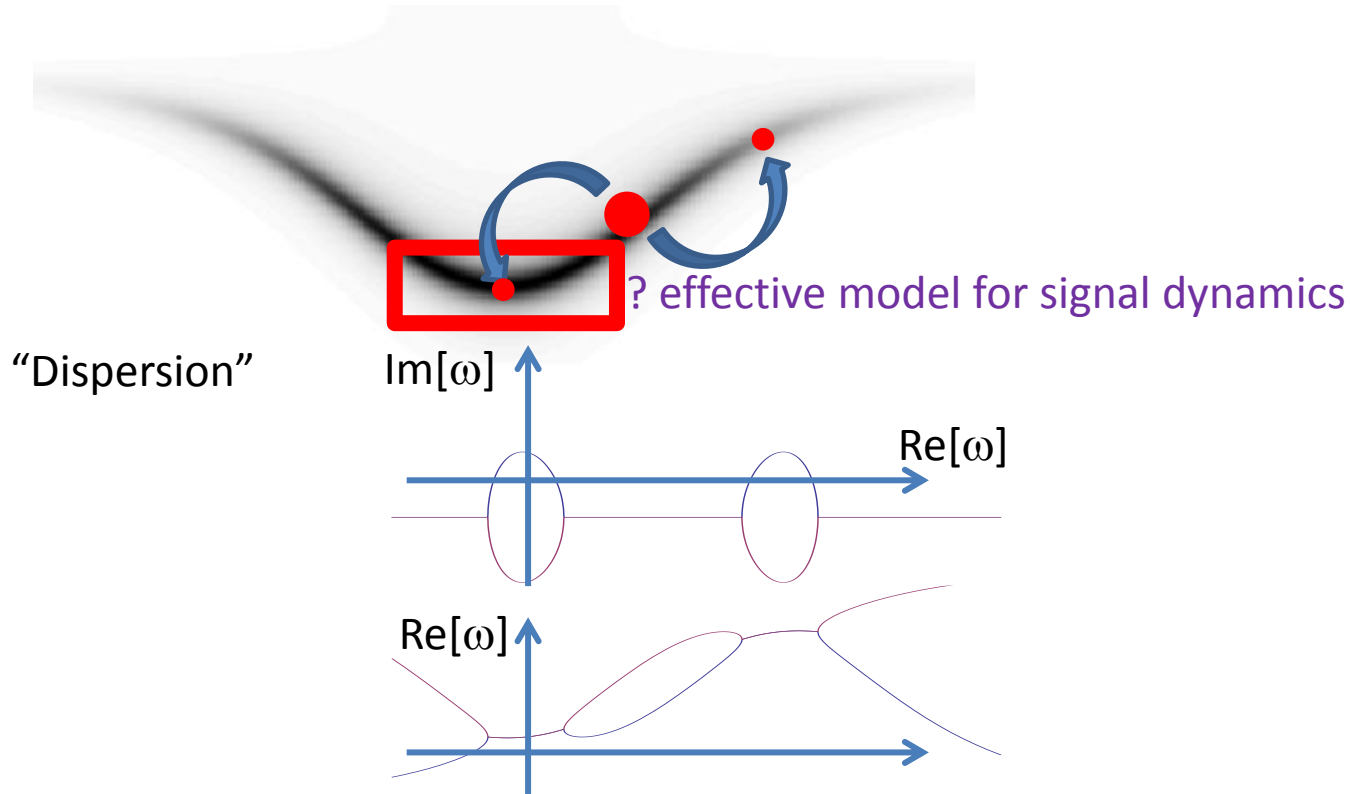


OPO = optical parametric oscillator: occupation of signal/idler in mean field theory

spontaneous U(1) symmetry breaking: $\phi_s + \phi_i = 2\phi_p$ but $\phi_s - \phi_i = \text{free}$

nonequilibrium analog of **BEC** phase transition

dynamics close to bifurcation



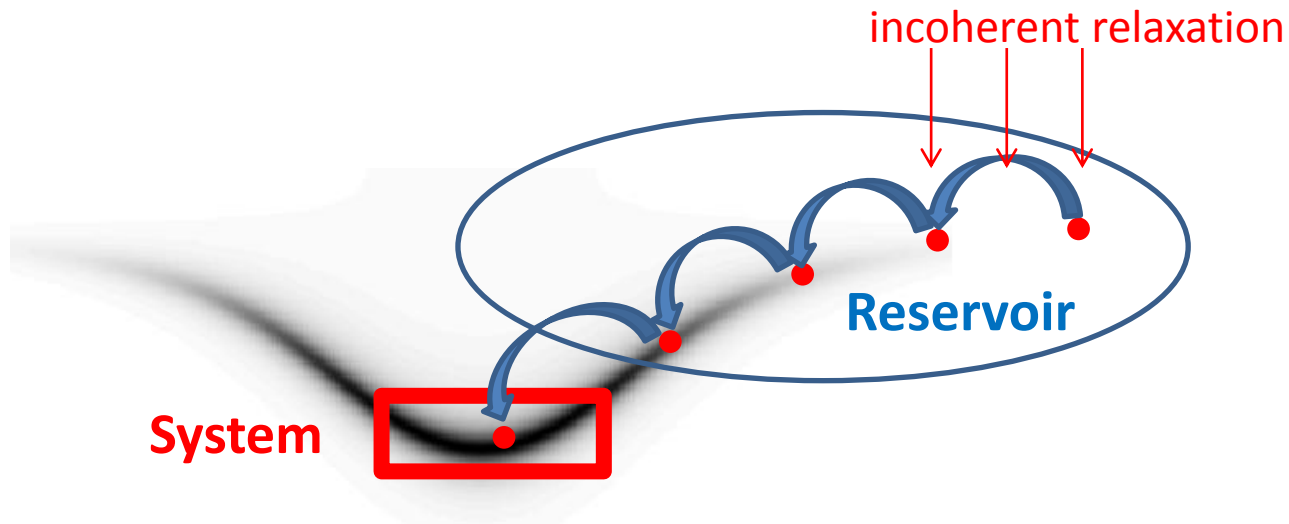
complex Ginzburg Landau-equation (=complex Gross-Pitaevskii Equation)

$$i \frac{\partial}{\partial t} \psi = \left[\underbrace{a + ib + (b + ic) \nabla^2}_{\text{linear: dispersion and gain}} + \underbrace{(-d + ie) |\psi|^2}_{\text{interaction: gain saturation and energy shift}} \right] \psi \quad \mathbf{U(1)} \text{ symmetric}$$

linear: dispersion and gain

interaction: gain saturation and energy shift

“nonresonant” excitation



excitons thermalize with lattice

complicated relaxation process when polaritons are injected at high energy,
but no change in nature of dynamical instability

→ also model with cGLE or with similar phenomenological models

analogy with lasers (Jonathan Keeling)

reservoir model

polariton field dynamics

$$i \frac{\partial}{\partial t} \psi = \left(\varepsilon(-i\nabla) + V_{ext} - i \frac{\gamma - R(n_R)}{2} + g|\psi|^2 + g_R n_R \right) \psi$$

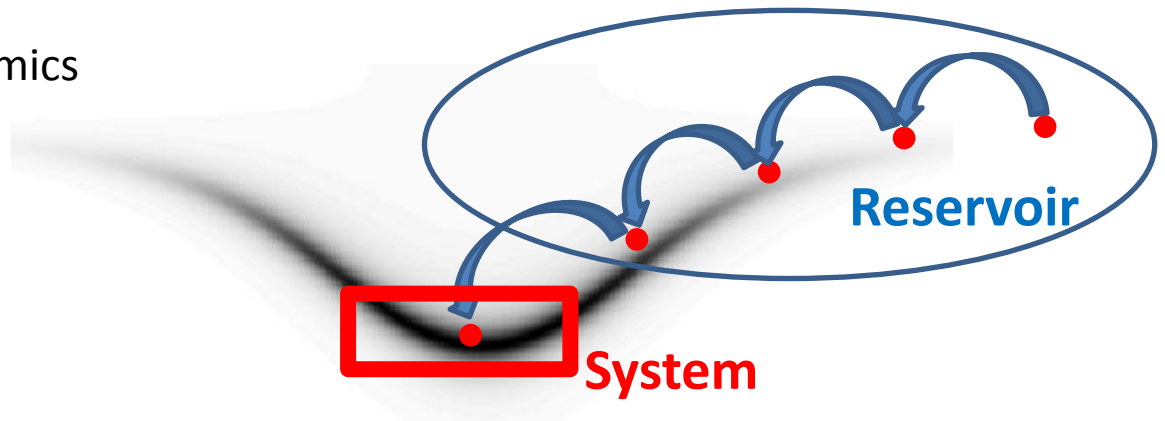
gain

exciton-polariton interaction

coupled to exciton density dynamics

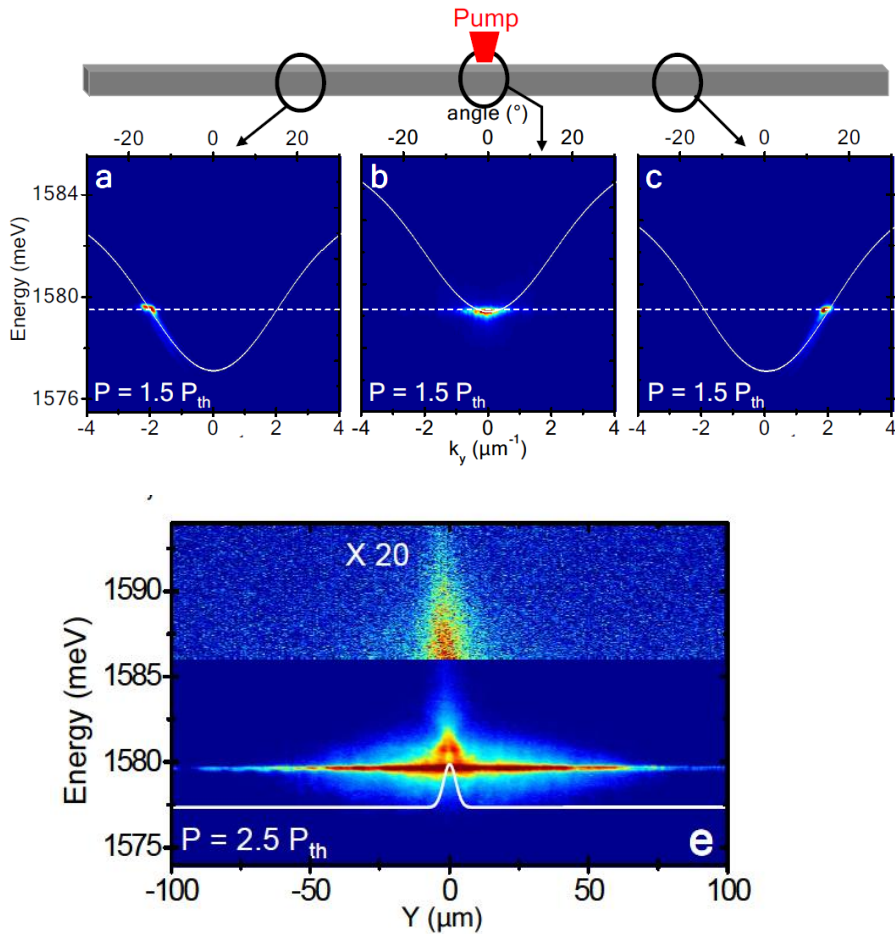
$$\frac{\partial}{\partial t} n_R = P - \gamma_R n_R - R(n_R) |\psi|^2$$

gain saturation

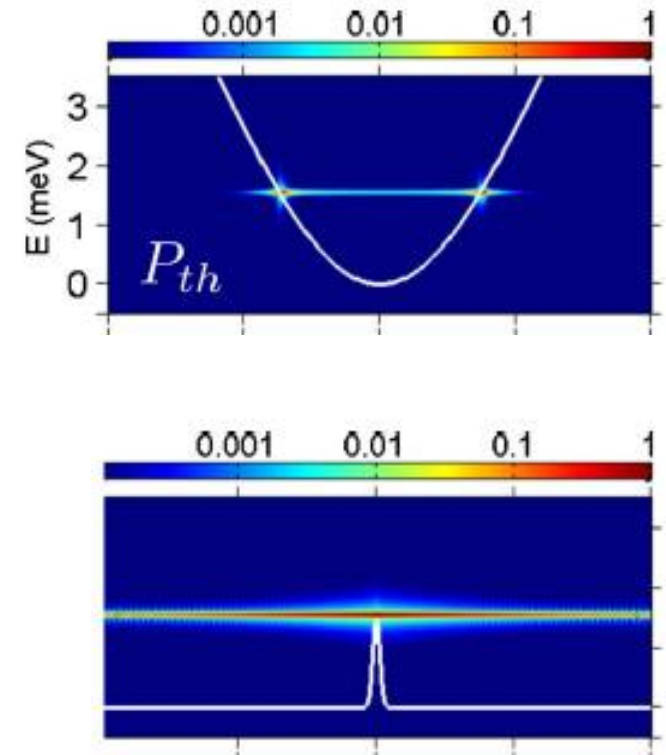


application: accelerated condensate

experiment



theory



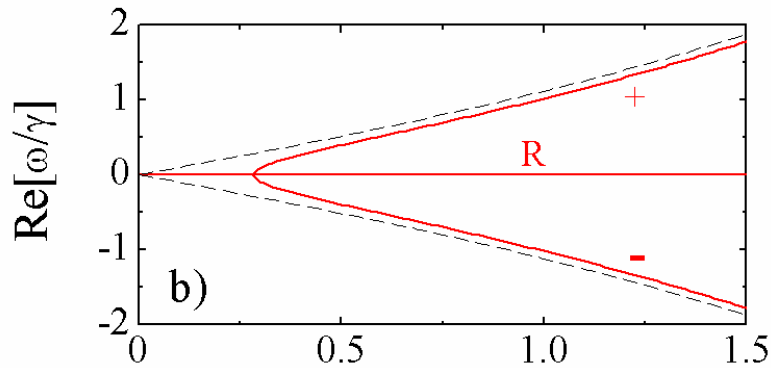
[E. Wertz et al. Nat. Phys. 2010]

[M.W. et al PRB 2010]

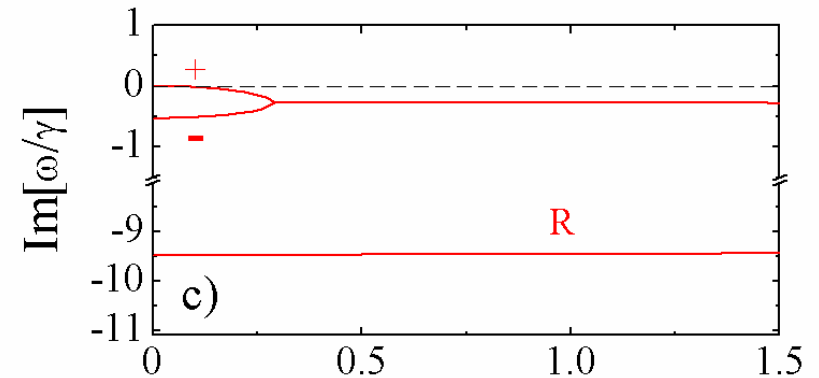
excitation spectrum

excitation spectrum

real part



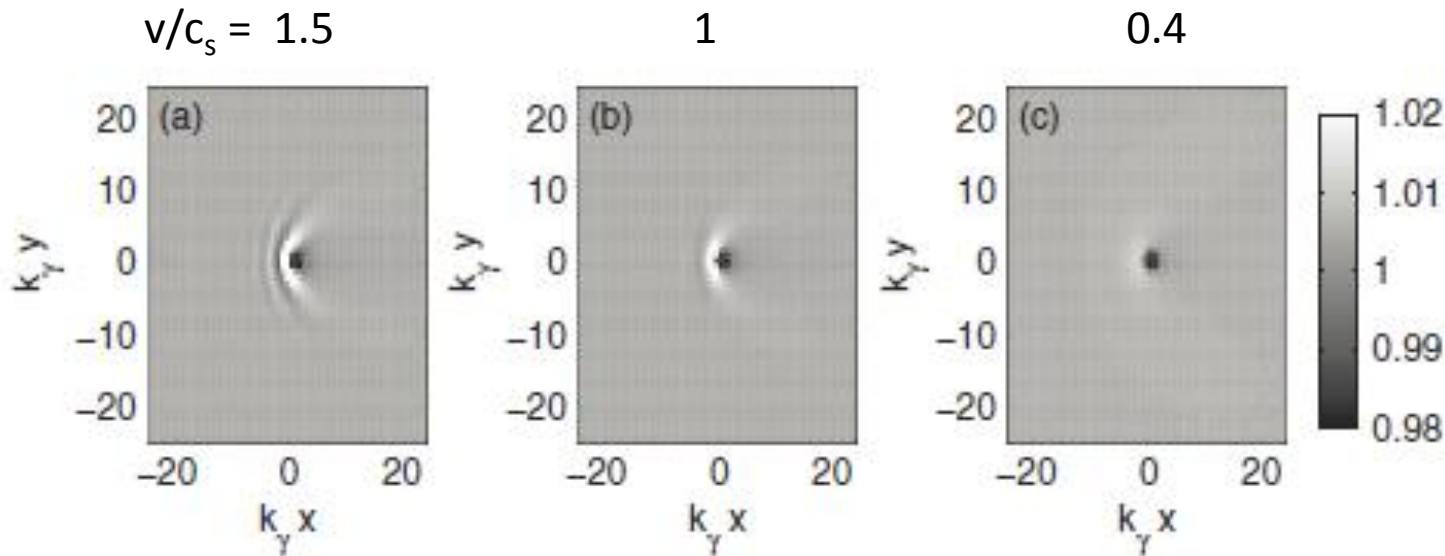
imaginary part



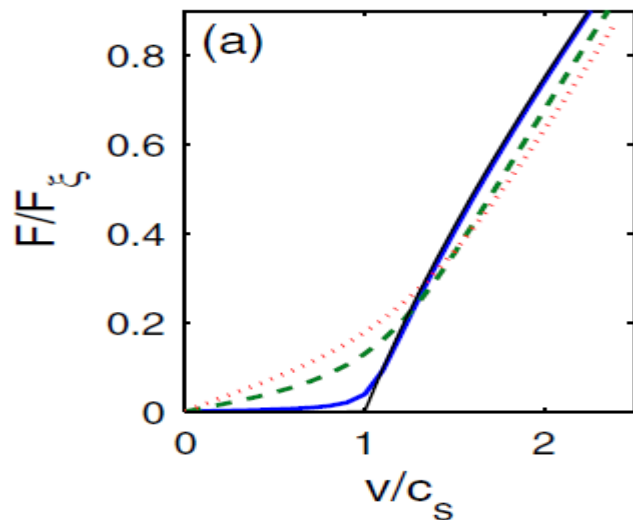
→ zero critical velocity ?

[M.Szymanska et al. PRL 2006)] [M. Wouters and I. Carusotto, PRL 2007]

superfluidity : interaction with a defect



Force on the defect



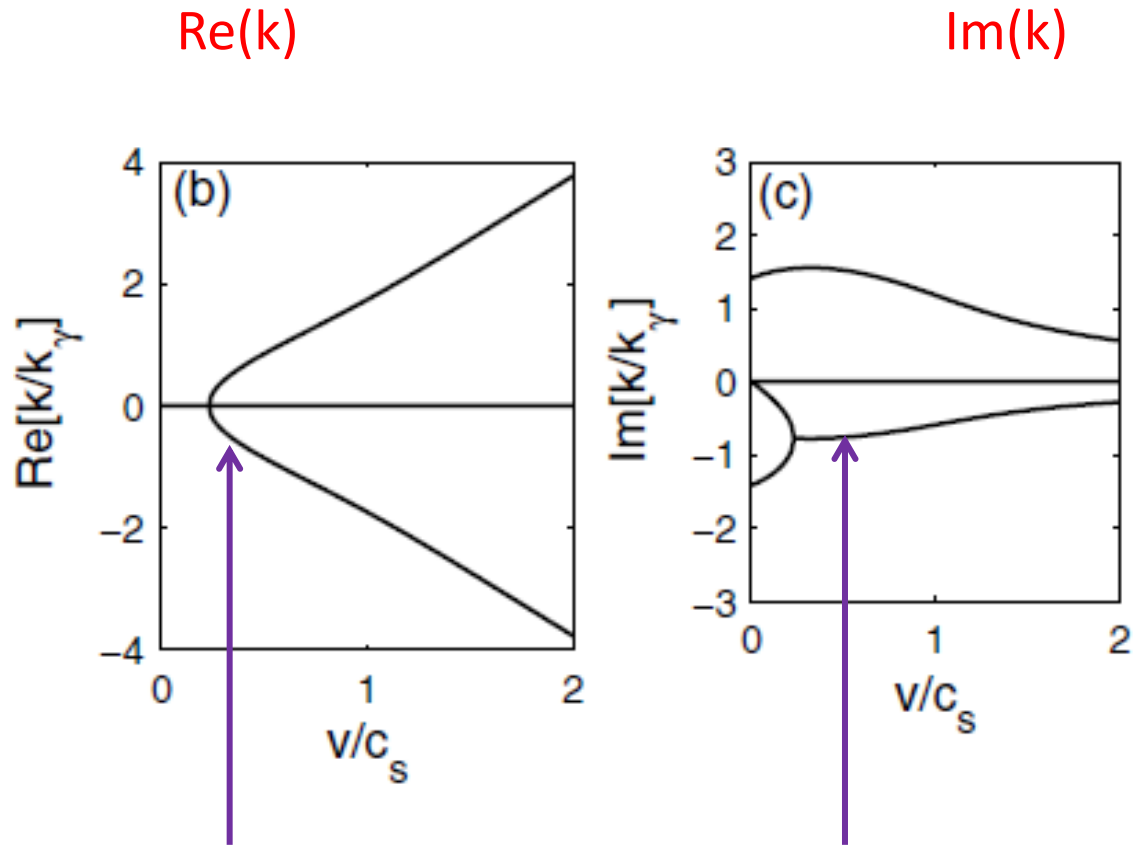
finite superfluid fraction

[J. Keeling, PRL 2011]

[M. W. and I. Carusotto, PRL 2010]

excitation spectrum II

Static defect excites modes for which $\omega(k) = 0$

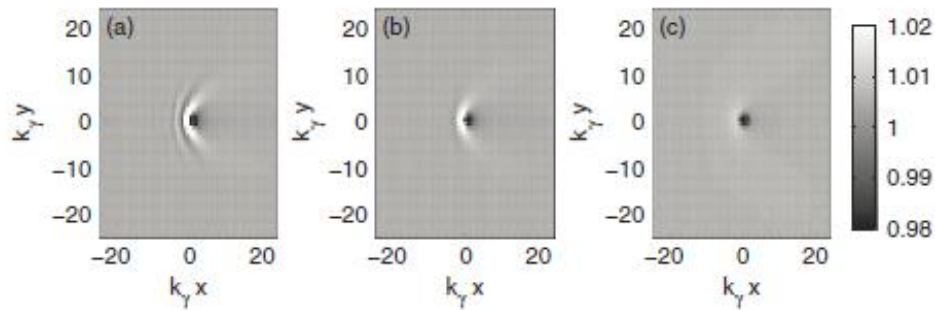


Condensate perturbation oscillation

Spatial decay rate

metastable flows

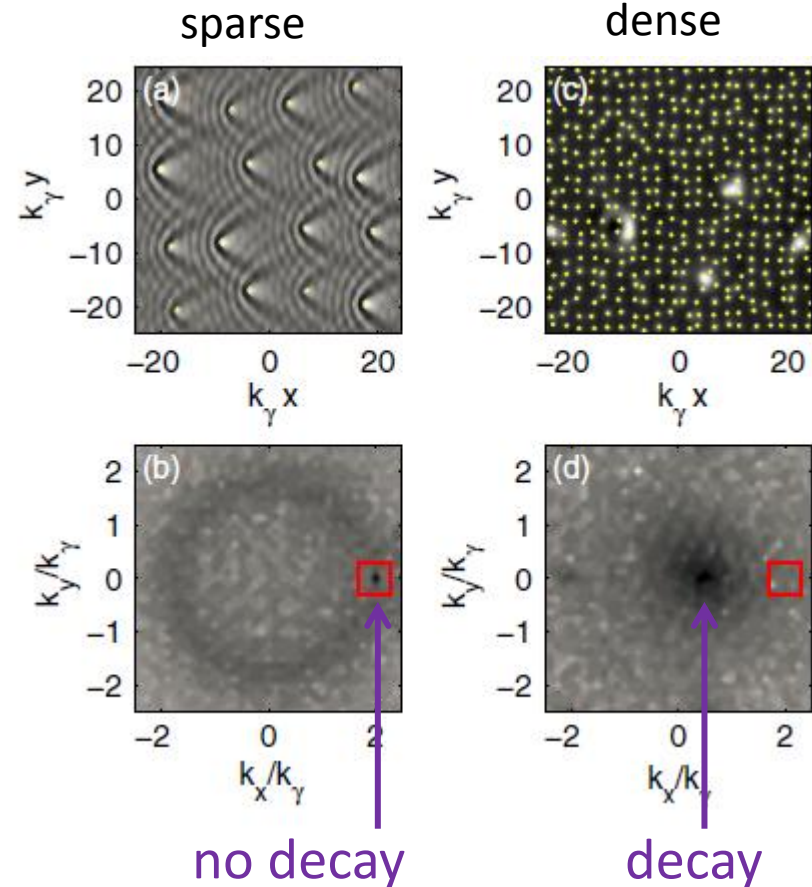
single defect



perturbation of the wave function,
force on defect,
but **no decay** of superflow

Damping of excitations enhances
stability of superflow!

many defects



long range order

- BEC: U(1) broken state has ODLRO:

$$\langle \hat{\psi}^+(x) \hat{\psi}(x') \rangle \rightarrow n_c \quad \text{for} \quad |x - x'| \rightarrow \infty$$

- What about the OPO ?

description of fluctuation needed \rightarrow go beyond mean field theory

- out of equilibrium \rightarrow no $\exp[-\beta H]$ Boltzmann distribution

- dynamical theory of open system :

$$\frac{\partial}{\partial t} \rho = \underbrace{-i[\rho, H]}_{\text{Hamiltonian evolution}} + \underbrace{L[\rho]}_{\text{Lindblad: losses, incoherent gain from reservoir}} \quad \leftarrow \text{simplification needed!!}$$

Hamiltonian evolution

Lindblad: losses, incoherent gain from reservoir

quantum \rightarrow c-field

- map quantum dissipative dynamics to classical stochastic process by means of the Wigner (quasi-) probability distribution $P_W(\varphi)$

= Truncated Winger approximation

- Observables: e.g. 1-body density matrix

$$\frac{1}{2} \langle \hat{\psi}^+(x) \hat{\psi}(x') + \hat{\psi}(x') \hat{\psi}^+(x) \rangle = \int [d^2 \varphi(x)] P_W(\varphi) \varphi^*(x) \varphi(x')$$

- $P_W(\varphi)$ is sampled by functions that follow the stochastic process

$$d\varphi(x) = \underbrace{\left\{ \left[\varepsilon(-i\nabla) + g|\varphi(x)|^2 - i\frac{\gamma}{2} \right] \varphi(x) + F_L \right\}}_{\text{GPE}} dt + \underbrace{\sqrt{\frac{4\gamma}{\Delta V}} dW(x,t)}_{\text{noise, related to losses}}$$

$$\langle dW^*(x,t) dW(x',t') \rangle = 2 dt \delta_{x,x'} \delta_{t,t'}$$

relation to Boltzmann eqn.

for homogeneous system: $\langle \hat{\psi}^+(x) \hat{\psi}(x') \rangle = \int n(k) e^{ik(x-x')} dk$

→ information on LRO in momentum distribution, that can be computed with Boltzmann

“RPA” on GPE **without noise** → Boltzmann equation without “spontaneous scattering”

$$\frac{d}{dt} n(k) = \sum_{k_i} \delta(\varepsilon_1 + \varepsilon_2 - \varepsilon_3 - \varepsilon_4) [-n_1 n_2 n_3 - n_1 n_2 n_4 + n_3 n_4 n_1 + n_3 n_4 n_2]$$

“RPA” on GPE **with noise** → Boltzmann equation with spontaneous scattering + ...

$$\frac{d}{dt} n(k) = I_B + \frac{\pi g^2}{2} \sum_{k_i} \delta(\varepsilon_1 + \varepsilon_2 - \varepsilon_3 - \varepsilon_4) [-n_1 - n_2 + n_3 + n_4]$$

full usual Boltzmann ↑

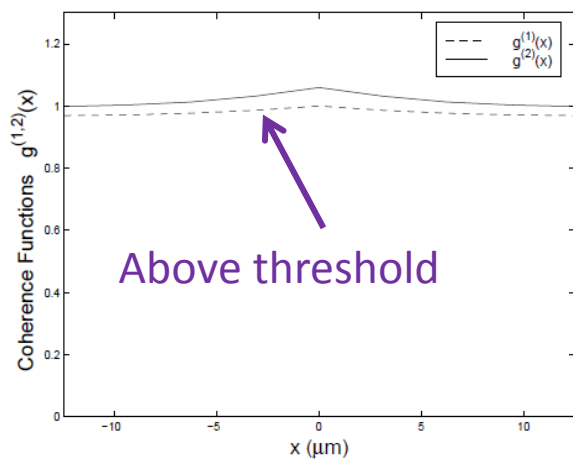
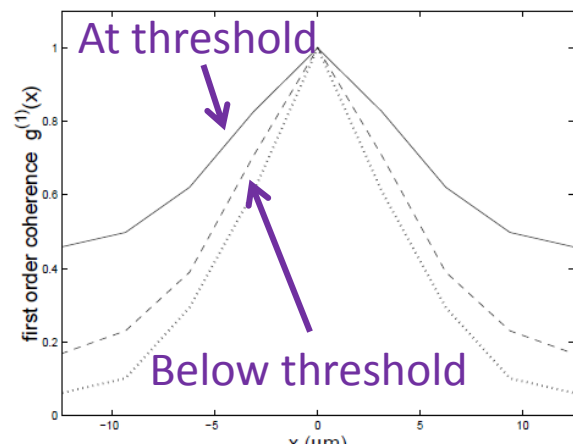
↑ spurious, should be small with respect to I_B or γn

satisfied if $n k_{\max}^2 \gg 1$ or $\gamma \gg g k_{\max}^2$

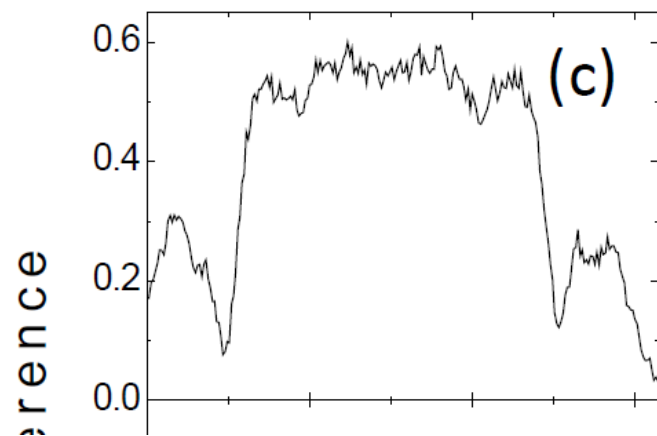
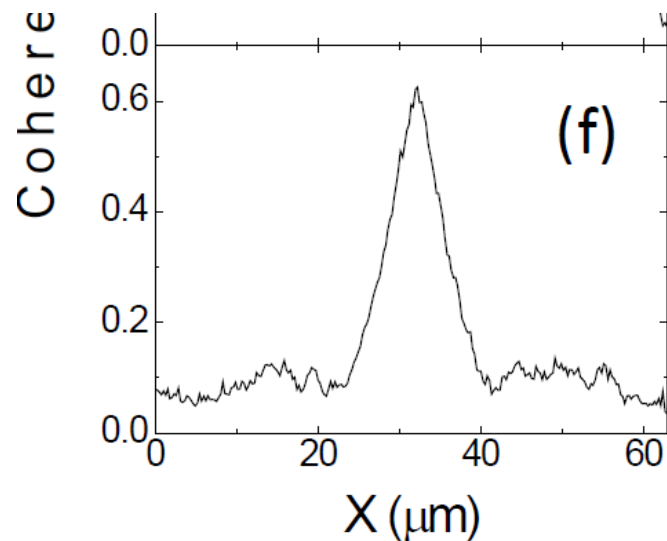
see e.g. [Y. Kagan in “Bose-Einstein condensation”, Griffin, Snook, Stringari eds. 1994.]

OPO: coherence across threshold

theory



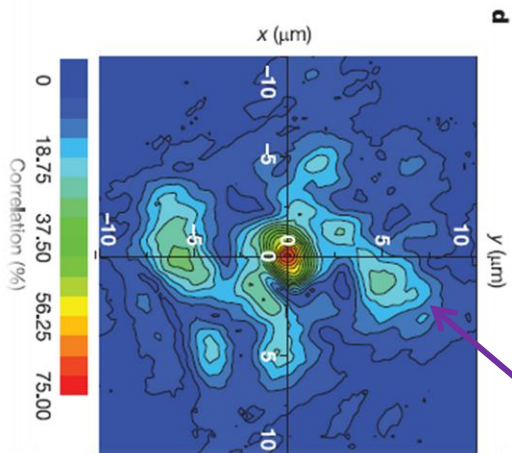
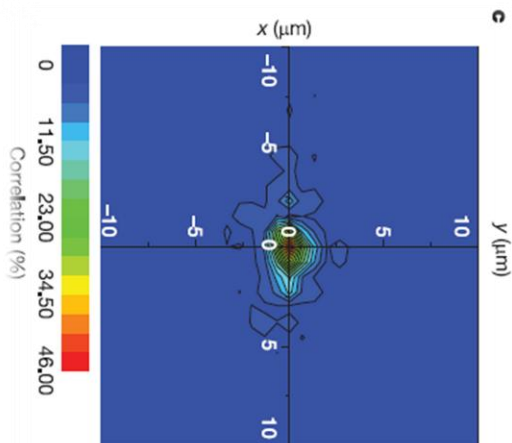
experiment



Nonresonant excitation: coherence

experiment

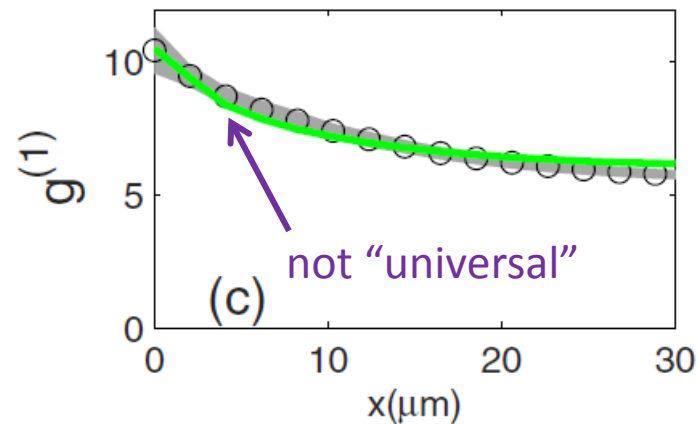
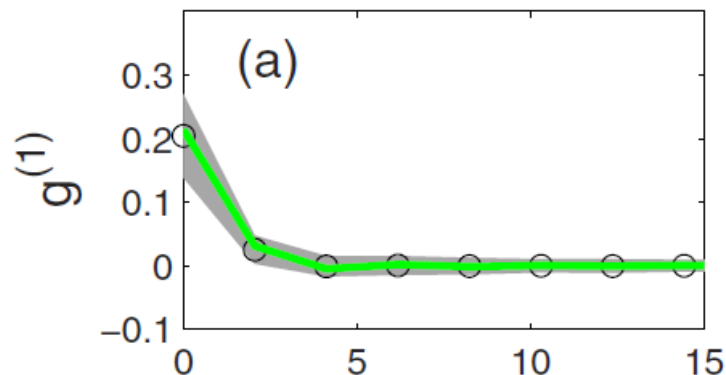
$$g^{(1)}(x, -x)$$



disorder

theory

(reservoir+Wigner)



[Kasprzak et al. Nature 2006]

[M.W. and V. Savona, PRB 2009]
[M.H. Szymanska et al PRB 2008]

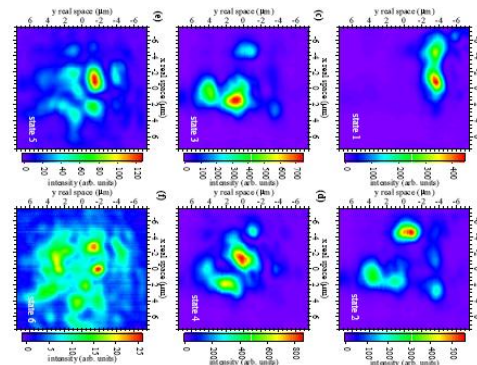
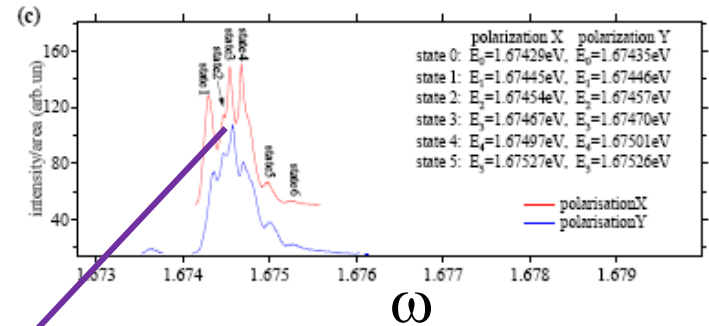
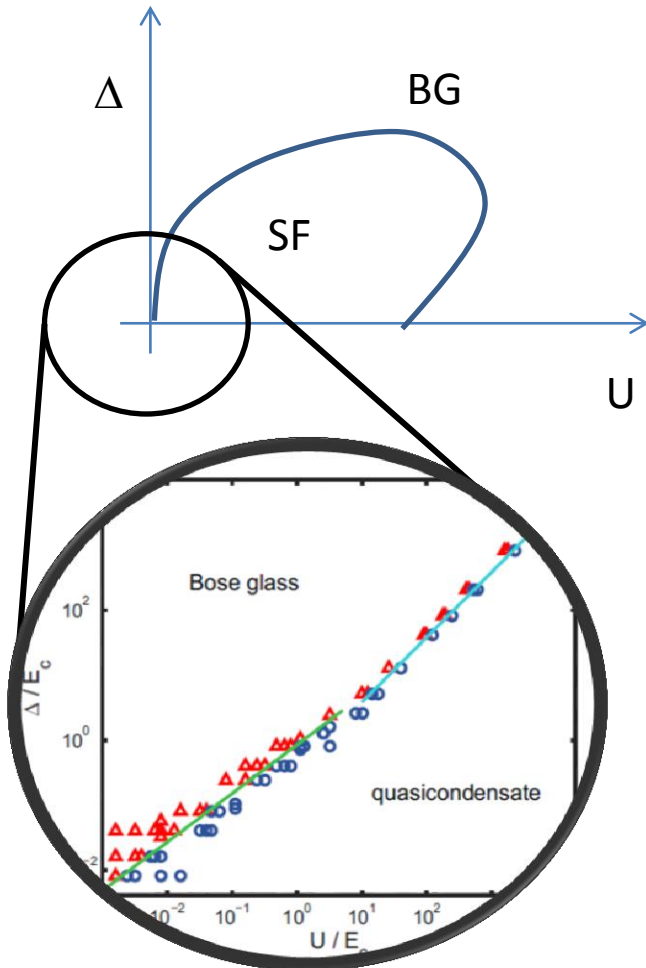
Disorder: Bose Glass?

equilibrium disordered Bose gas: theory

superfluidity: qualitatively survives

→ what with Bose glass ??

nonequilibrium: experiment



remember:

$$\omega \Leftrightarrow \mu$$

⇒ not a single chemical potential!

[Fisher et al. PRB 1989]

[Fontanesi et al. PRL 2009] see also [Altmann et al. PRB 2010]

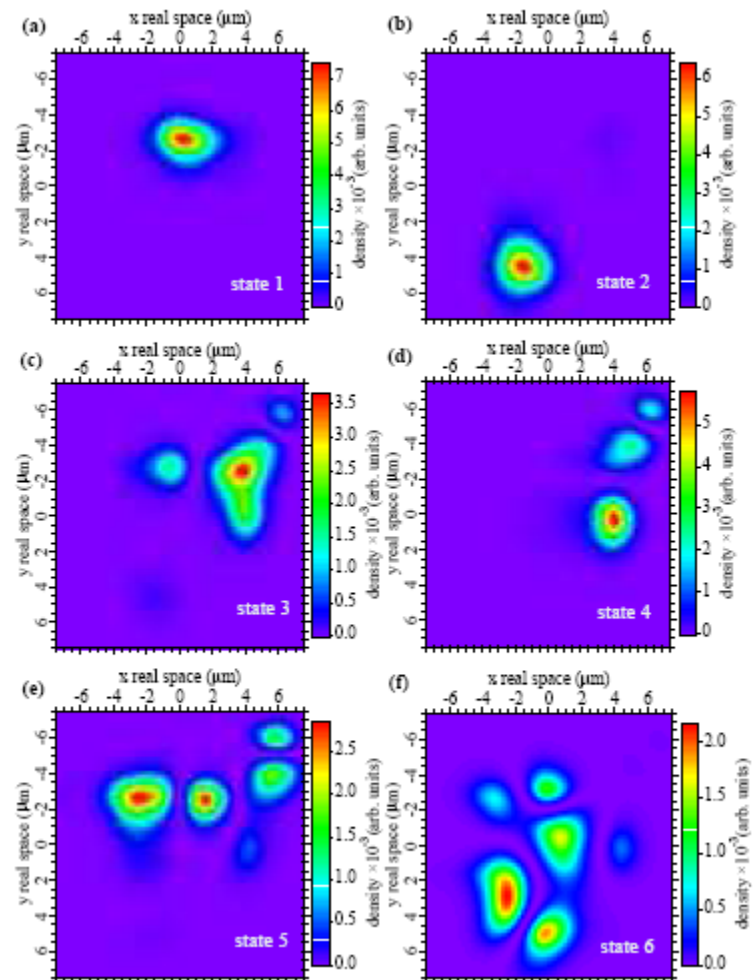
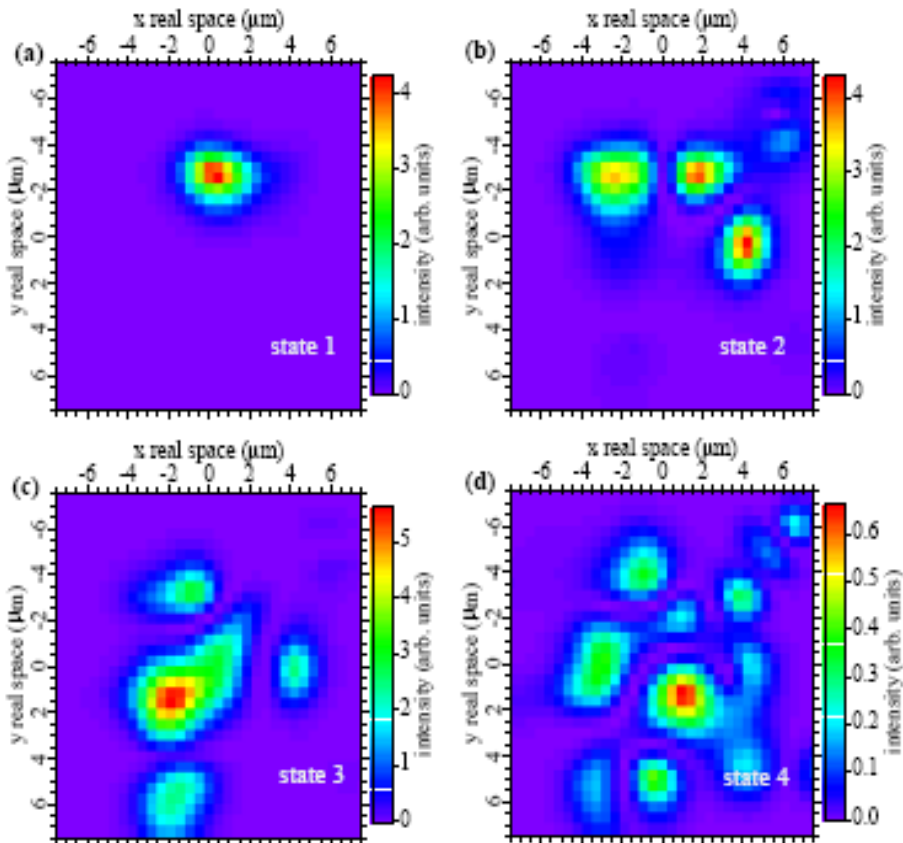
[Krizhanovskii et al, PRB 2009]

simulated states

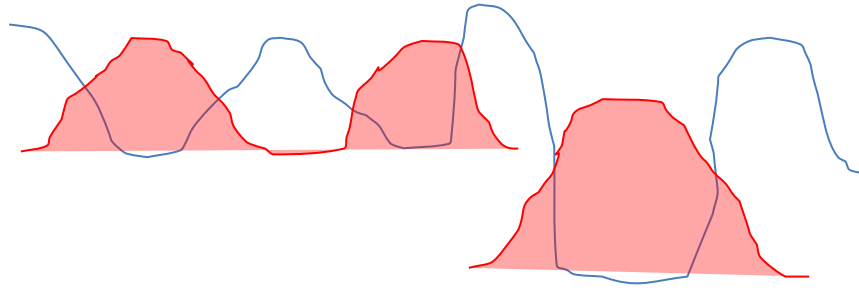
condensate states



linear eigenstates



physical picture



- gain due to excitons should be saturated everywhere
- lowest energies: eigenstates, higher energies: new states (mode locking)
- single state develops when $U \sim \Delta$:
same scaling as in equilibrium for disorder with long correlation length
[Malpuech et al PRL 2006]
- but very different physics: no fluctuations needed for the “incoherent phase” (cf. random lasers)

summary

microcavity polaritons challenge us to revisit
well know phenomena (superfluidity, BEC, Bose glass,)
out of equilibrium

GPE+ is a good tool to address these questions ,
especially when inhomogeneity is important

two-reservoir model

$$\frac{\partial}{\partial t} n_R = P - \gamma_R n_R - R(n_R) |\psi|^2$$

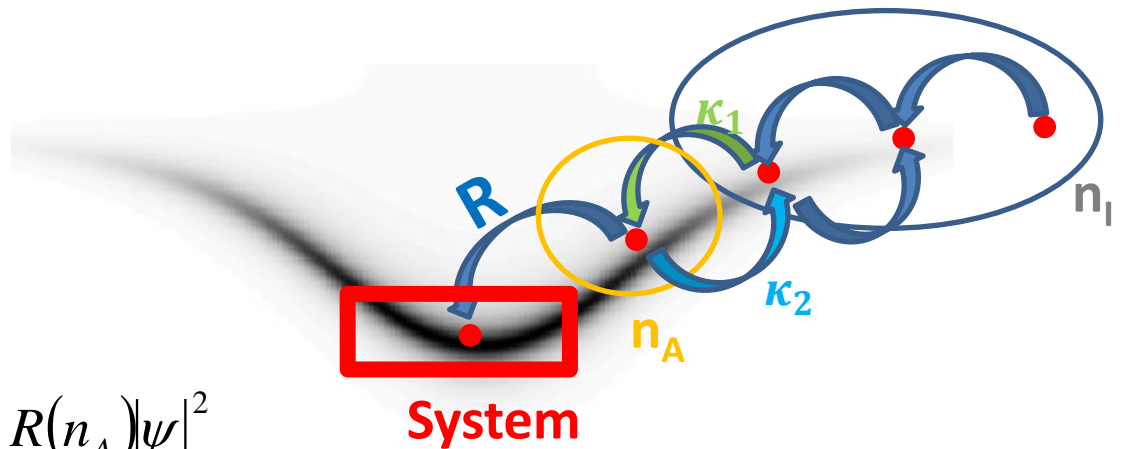
single reservoir time scale

but exciton life time much longer than gain saturation relaxation time scale
 → not good for pulsed excitation

$$\frac{\partial}{\partial t} n_I = P - \gamma_R n_I - \kappa_1 n_A + \kappa_2 n_I$$

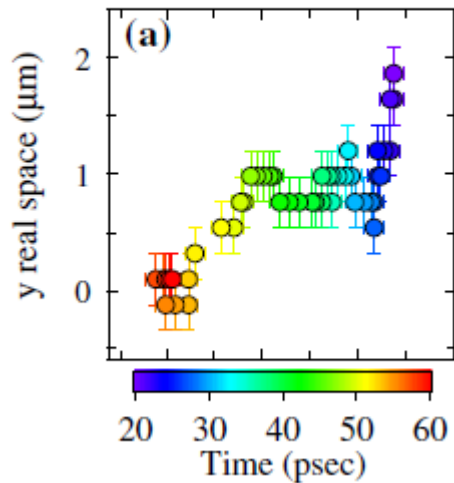
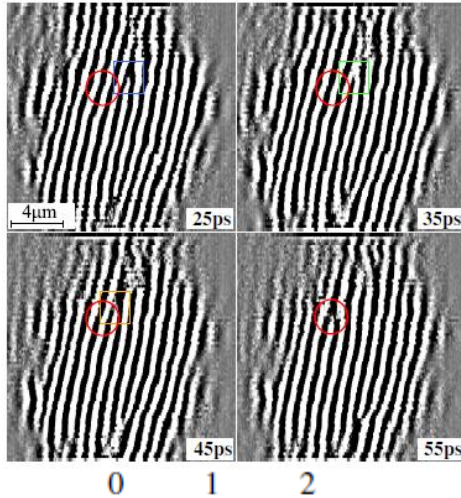
exciton life time gain time scale

$$\frac{\partial}{\partial t} n_A = -\gamma_R n_A + \kappa_1 n_I - \kappa_2 n_A - R(n_A) |\psi|^2$$



application: vortex dynamics

experiment



simulation

vortex positions averaged interferogram

