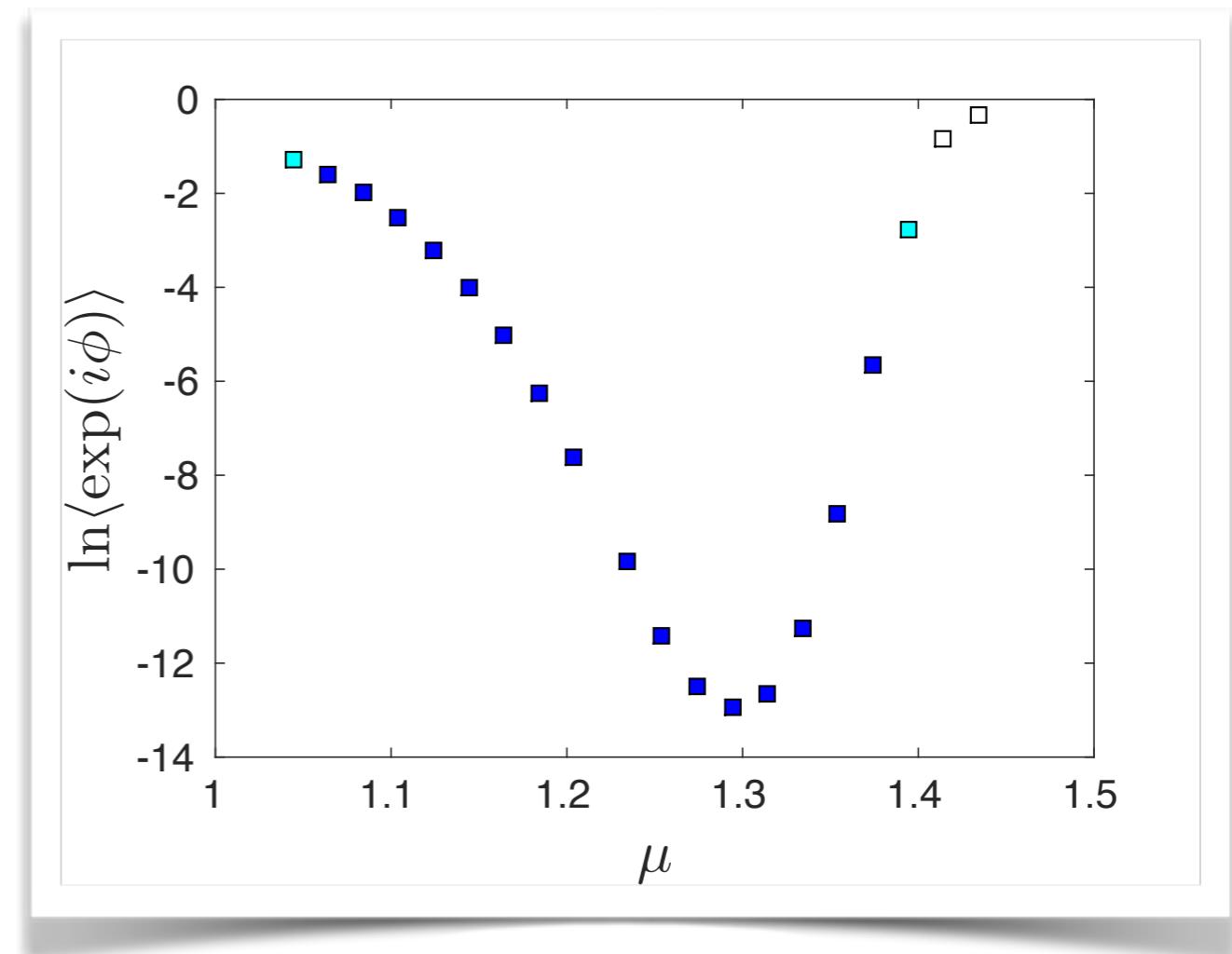


Can the density-of-states method solve sign problems ?



Kurt Langfeld (Plymouth & Liverpool)

Delta Meeting, Heidelberg, 28-30 April 2016

Biagio Lucini (Swansea)

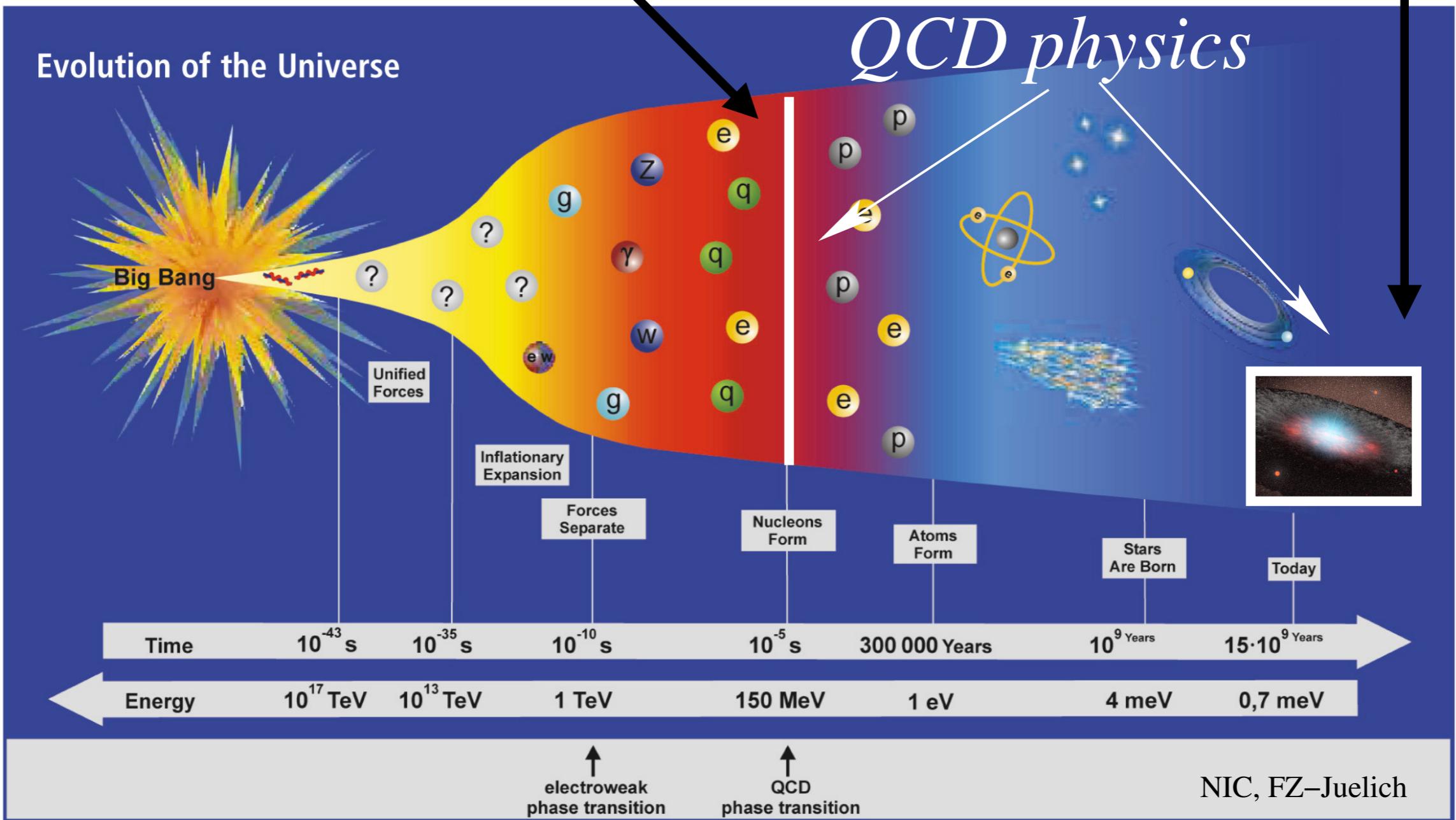
Antonio Rago, Nicolas Garron (Plymouth)

Roberto Pellegrini (Edinburgh)

low density & high T
Standard LGT:



high density & low T
Standard LGT:



The density-of-states method:

Consider the high dimensional integral:

$$Z = \int \mathcal{D}\phi \exp\{\beta S[\phi]\}$$

The density-of-states:

$$\rho(E) = \int \mathcal{D}\phi \delta(E - S[\phi])$$

A 1-dimensional integral:

$$Z = \int dE P(E)$$

Gibbs factor

entropy



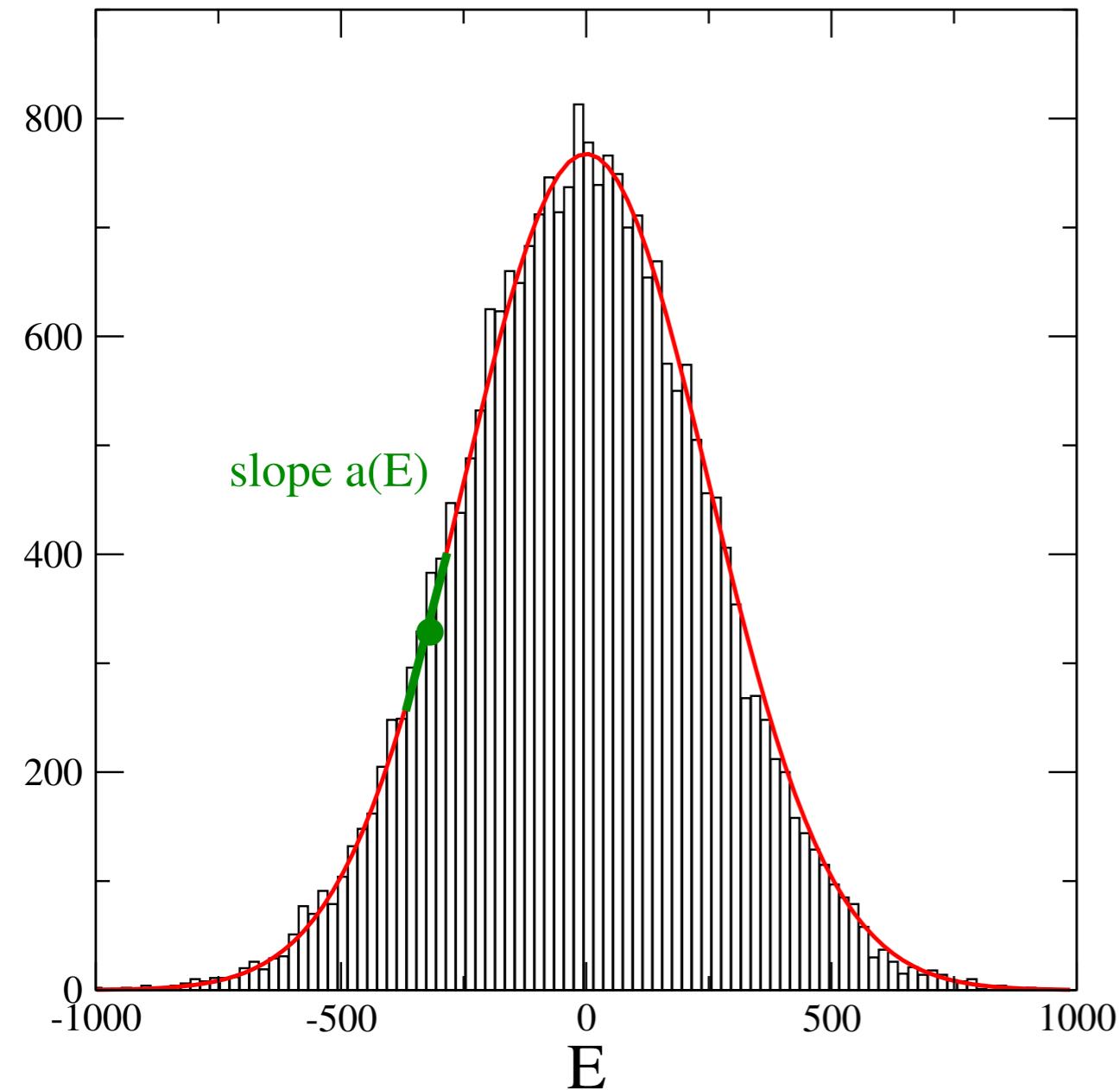
$$P(E) = \rho(E) e^{\beta E}$$

Probabilistic weight

How do I find the density-of-states?

straightforward: Histogram of the action S

The LLR approach to the density-of-states:



[Langfeld, Lucini, Rago, PRL 109 (2012) 111601]

I. Calculate the slope at E from a stochastic non-linear equation:

$$a(E) = \frac{d \ln \rho(E)}{dE}$$

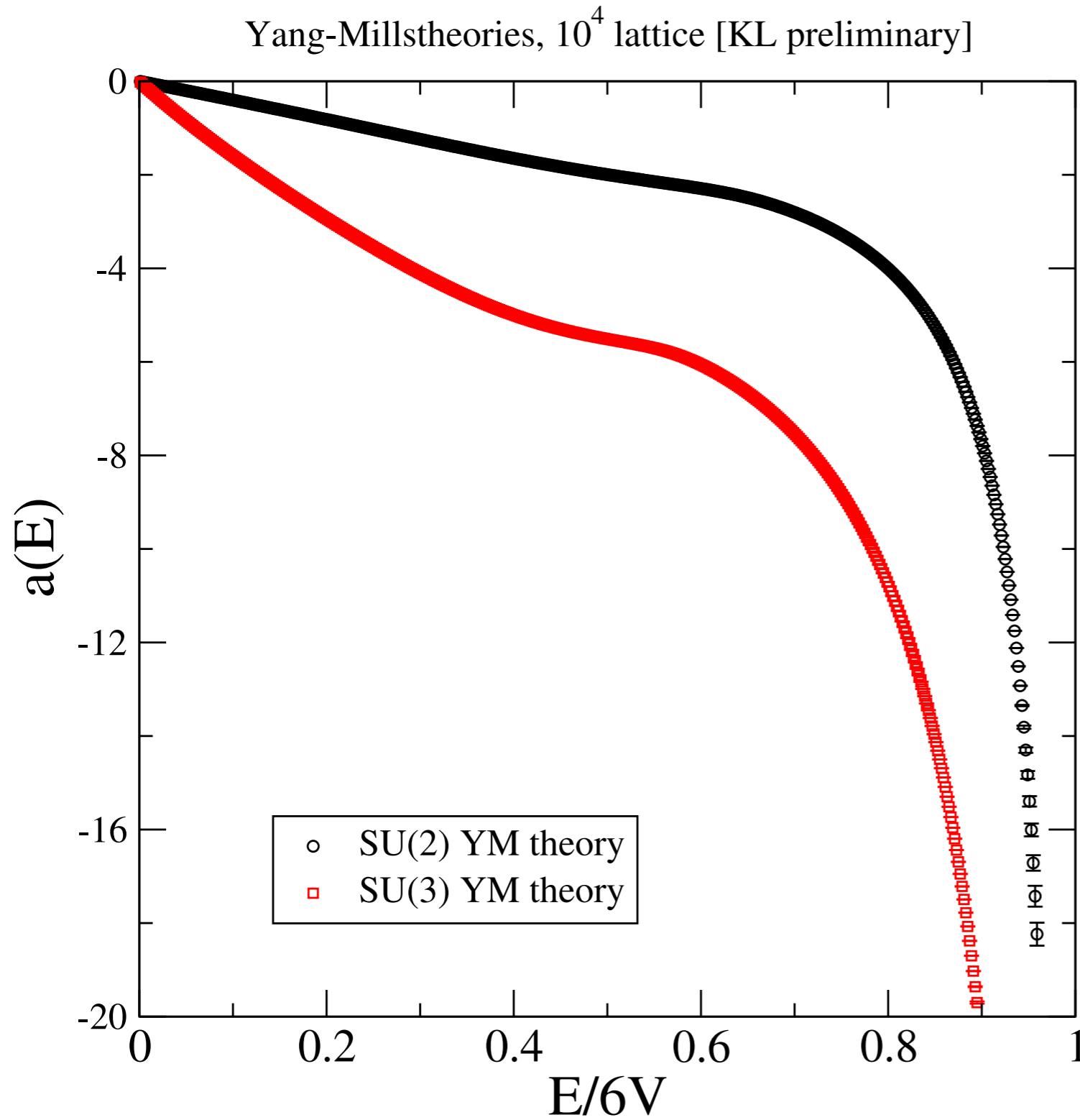
$$\langle\langle S - E \rangle\rangle(a) = 0$$

↑
MC average

2. Reconstruct $\rho(E)$

3. Find: $P(E) = \rho(E) e^{\beta E}$

Showcase: SU(2) and SU(3) Yang-Mills theory

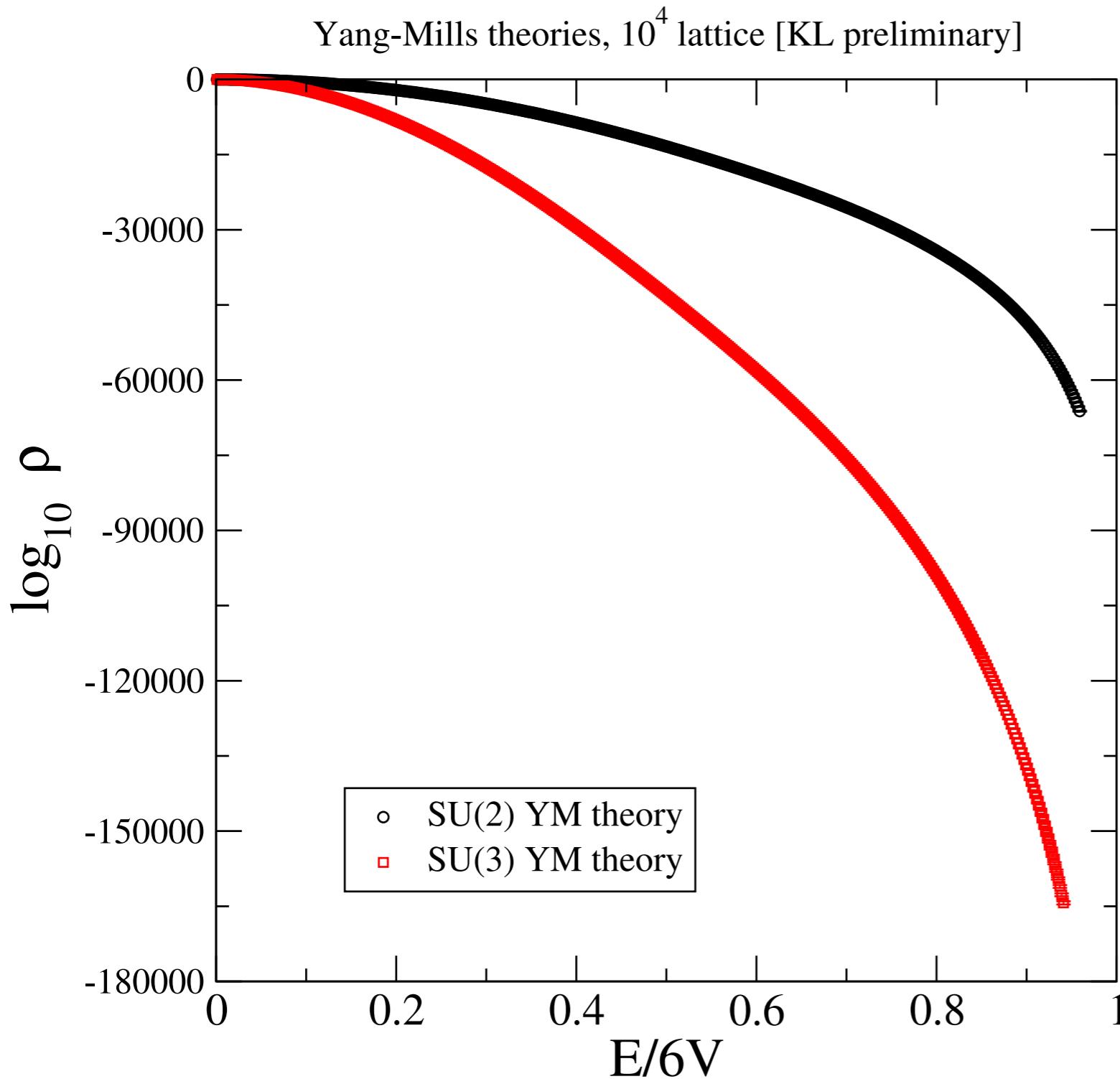


Error bars from bootstrap
(20 independent runs per interval)

features an *exponential error suppression*

[Langfeld, Lucini, Pellegrini, Rago,
arXiv:1509.08391]

Showcase: SU(2) and SU(3) Yang-Mills theory

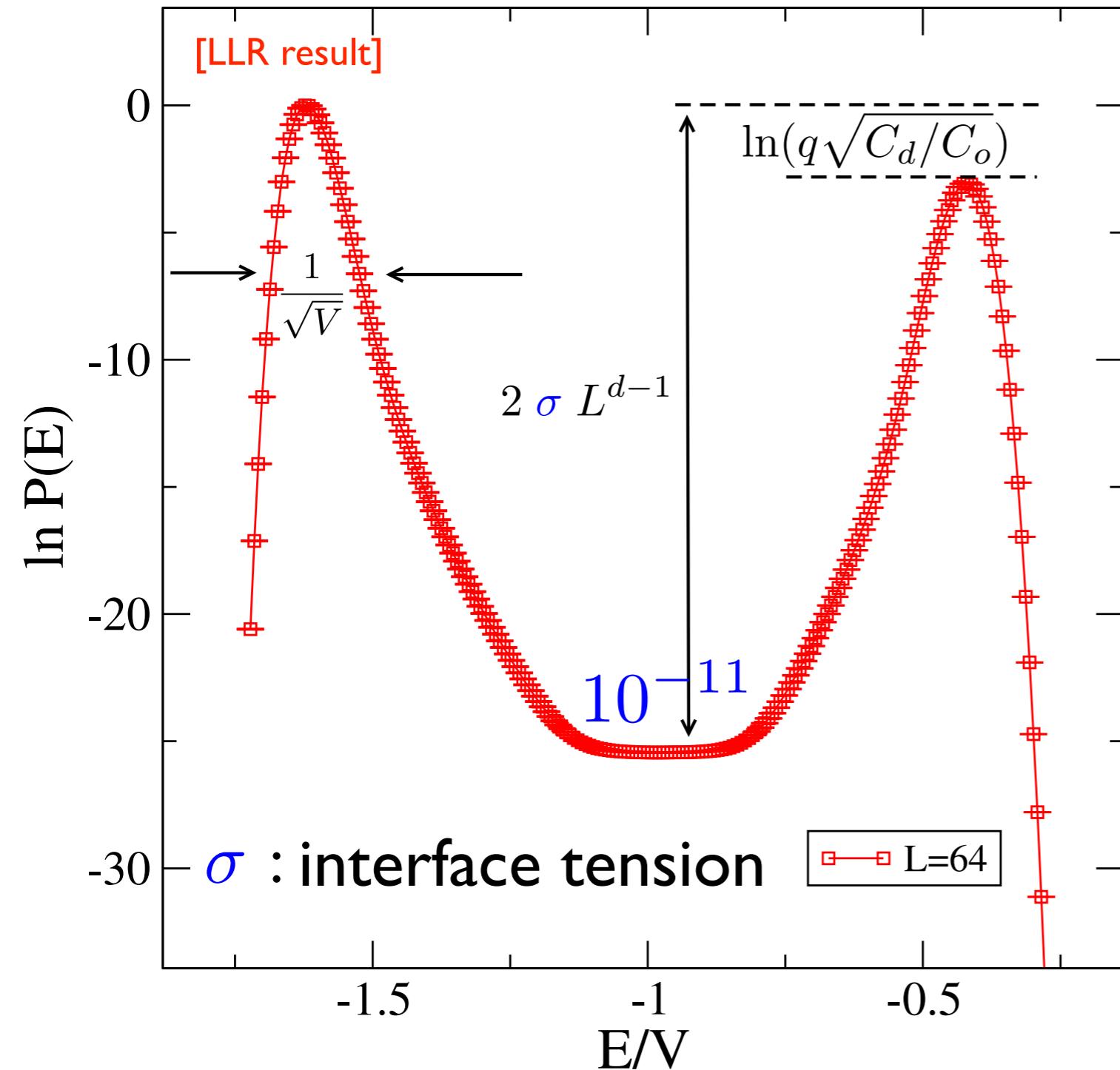


Density of states
over 100,000 orders
of magnitude!

[Gattringer Langfeld, arXiv:1603.09517]

Showcase: q-state Potts model in 2d

$q=20$ Potts model for a L^2 lattice at β_{critical}



LLR solves overlap problems!

Exact solution:

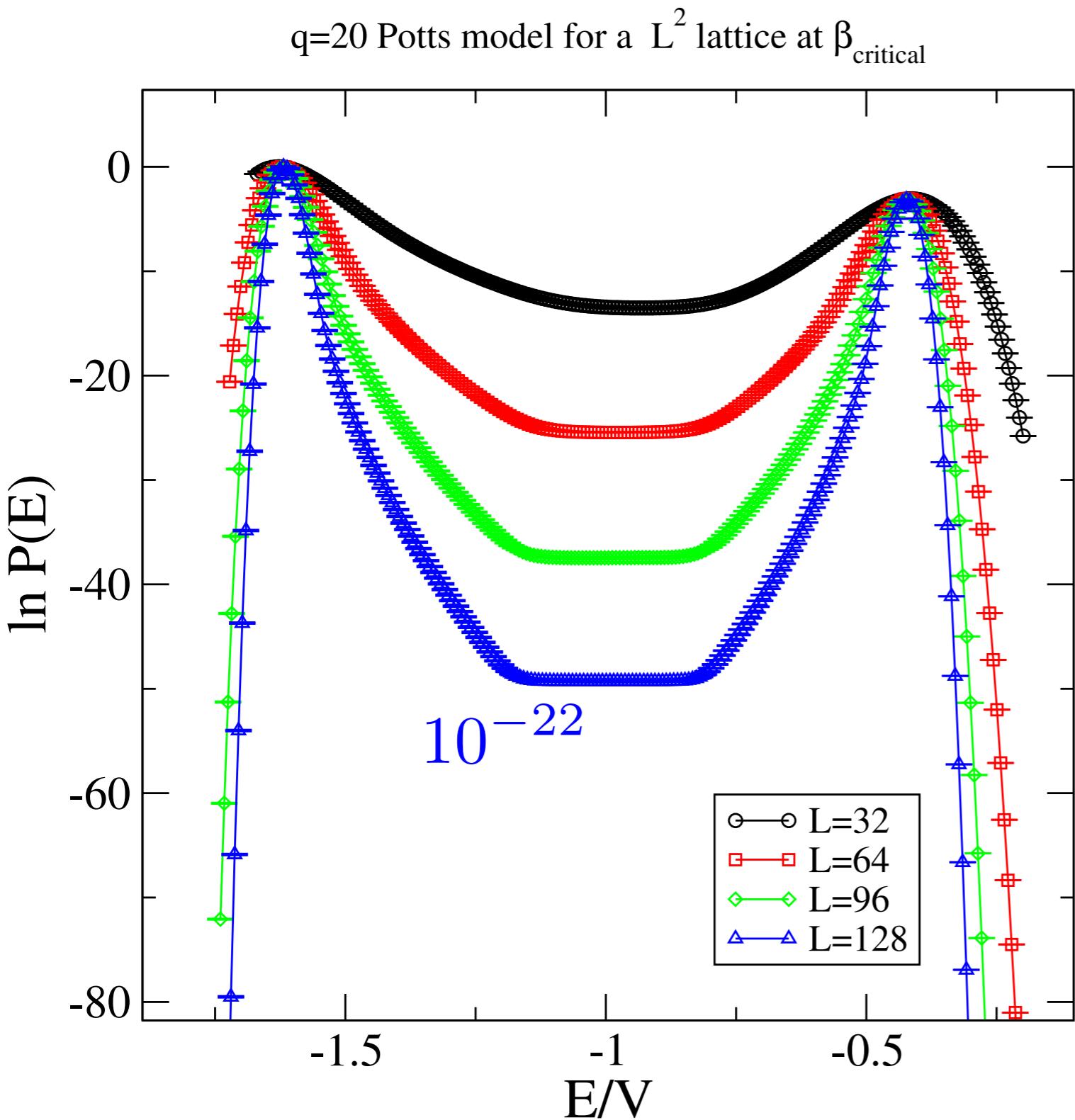
R.J. Baxter, J. Phys. C6 (1973) L445

$$\beta_{\text{critical}} = \frac{1}{2} \ln(1 + \sqrt{q})$$

**First MC $q=20$ simulation:
Multi-canonical approach**

[Berg, Neuhaus, PRL 68 (1992) 9]
[Billoire, Neuhaus, Berg, NPB (1994) 795]

Showcase: q-state Potts model in 2d



LLR result:
216 energy intervals
replica method

Showcase: q-state Potts model in 2d

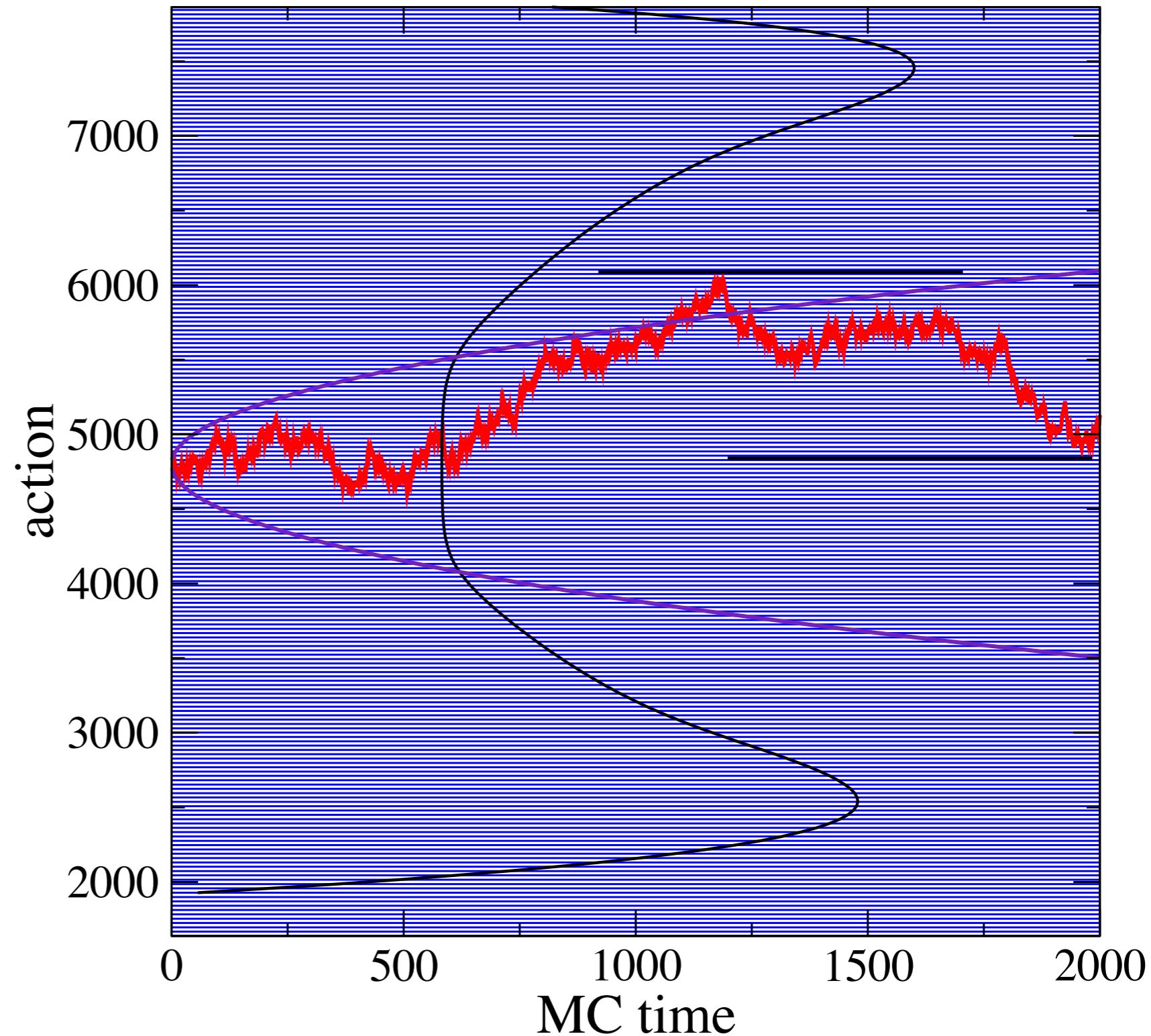
[q=20, L=64]

Tunnelling between
LLR action intervals:

Interval size: 29

bridged 42
intervals within
750 sweeps

$$[\sqrt{750} = 27.38 \dots]$$



How do we simulate dense matter QFT with the density-of-states method?

- Quantum Field Theory:

$$Z = \int \mathcal{D}\phi \exp\{\beta S_R[\phi]\}$$

infinite dimension integral

Monte-Carlo simulation (importance sampling!)

$\exp\{\beta S_R[\phi]\}$: probabilistic weight

- QFT at finite densities:

$$Z = \int \mathcal{D}\phi \exp\{\beta S_R[\phi] + i\mu S_I[\phi]\}$$

μ : chemical potential

S_I : imaginary part of the action

complex!

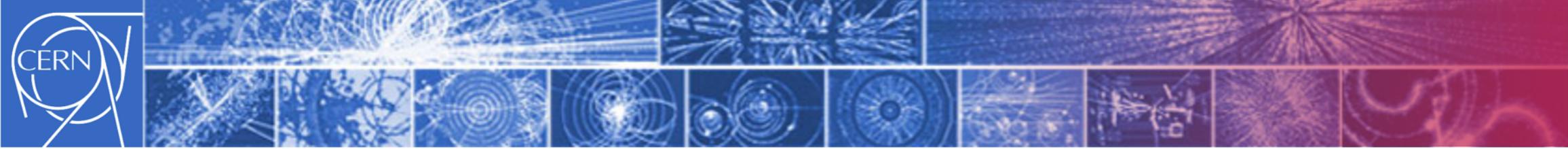
How can we quantify the problem?

- If we drop the imaginary part of the action:

$$Z_{\text{PQ}}(\mu) = \int \mathcal{D}\phi \exp\{S_R[\phi]\}$$

- Define the *overlap* between full and phase quenched theory

$$O(\mu) = \frac{Z(\mu)}{Z_{\text{PQ}}(\mu)} = \langle \exp\{i\mu S_I\} \rangle_{\text{PQ}}$$



Overlap problem:

- Z and Z_{PQ} have different free energy densities

$$O(\mu) = \frac{Z(\mu)}{Z_{PQ}(\mu)} = \exp\{-\Delta f V\} \Rightarrow \text{can be very small} \quad (\Delta f > 0)$$

re-weighting is inefficient!

- Trivially: $Z(\mu) = \frac{Z(\mu)}{Z_{PQ}(\mu)} Z_{PQ}(\mu) = O(\mu) Z_{PQ}(\mu)$

$$\rho(\mu) = \frac{T}{V_3} \frac{\partial}{\partial \mu} O(\mu) + \rho_{PQ}(\mu)$$

standard Monte-Carlo

generically dominant!

The density-of-states approach for complex theories:

- Recall: theory with complex action

$$Z = \int \mathcal{D}\phi \exp\{\beta S_R[\phi] + i\mu S_I[\phi]\}$$

- Define the generalised density-of-states:

$$P_\beta(s) = \int \mathcal{D}\phi \delta(s - S_I[\phi]) \exp\{\beta S_R[\phi]\}$$

Could get it by histogramming

- Partition function emerges from a FT:

$$Z = \int ds P_\beta(s) \exp\{i\mu s\}$$

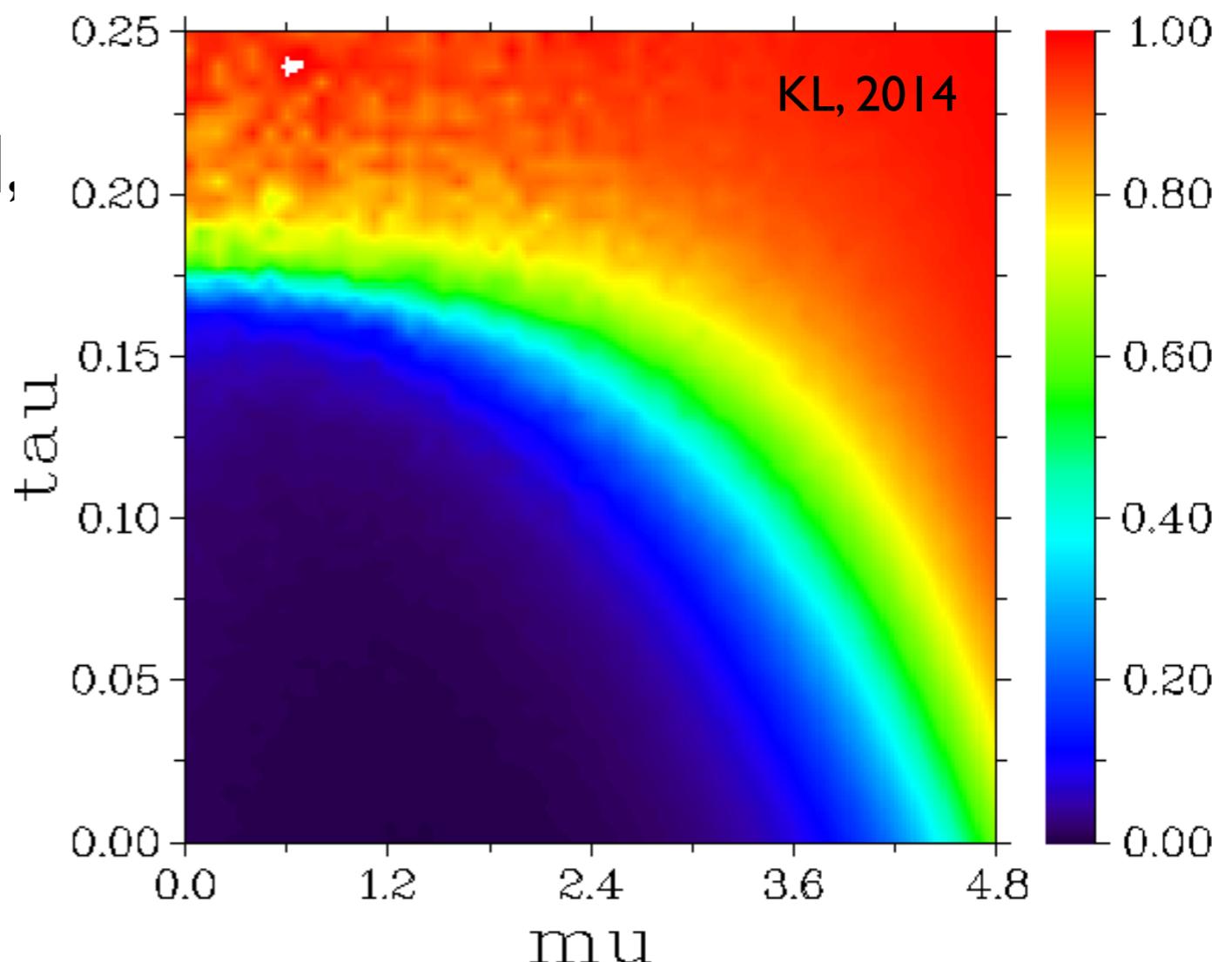
The Z3 showcase:

$$S[z] = \tau \sum_{x,\nu} [z_x z_{x+\nu}^* + cc] + \sum_x [\eta z_x + \bar{\eta} z_x^*]$$

τ : temperature, $\eta = \kappa e^\mu$, $\bar{\eta} = \kappa e^{-\mu}$ $z \in Z_3$

- Solvable: dual theory is real, efficient flux algorithm

[Mercado, Evertz, Gattringer,
PRL 106 (2011) 222001]



What is the scale of the problem?

- Indicative result: $P_\beta(s) = \exp\left\{-\frac{s^2}{V}\right\}$

action

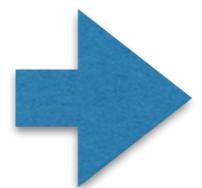
volume

$$Z = \int ds e^{-s^2/V} \exp\{i \mu s\} \propto \exp\left\{-\frac{\mu^2}{4} V\right\}$$

↑ statistical errors

exponentially small

Need exponential error suppression over the whole action range



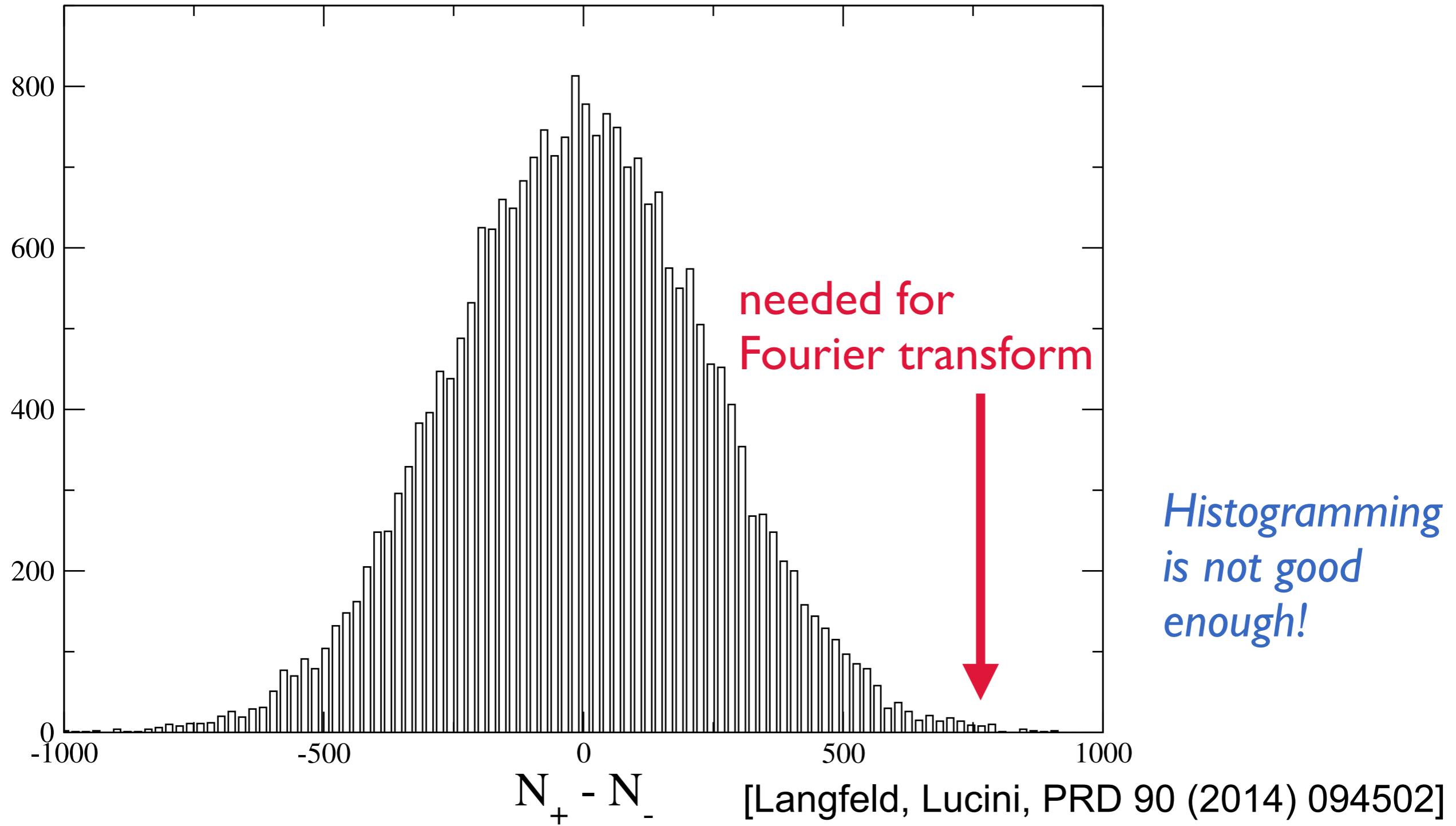
LLR approach:

[Langfeld, Lucini, PRD 90 (2014) 094502]

[Langfeld, Lucini, Rago, PRL 109 (2012) 111601]

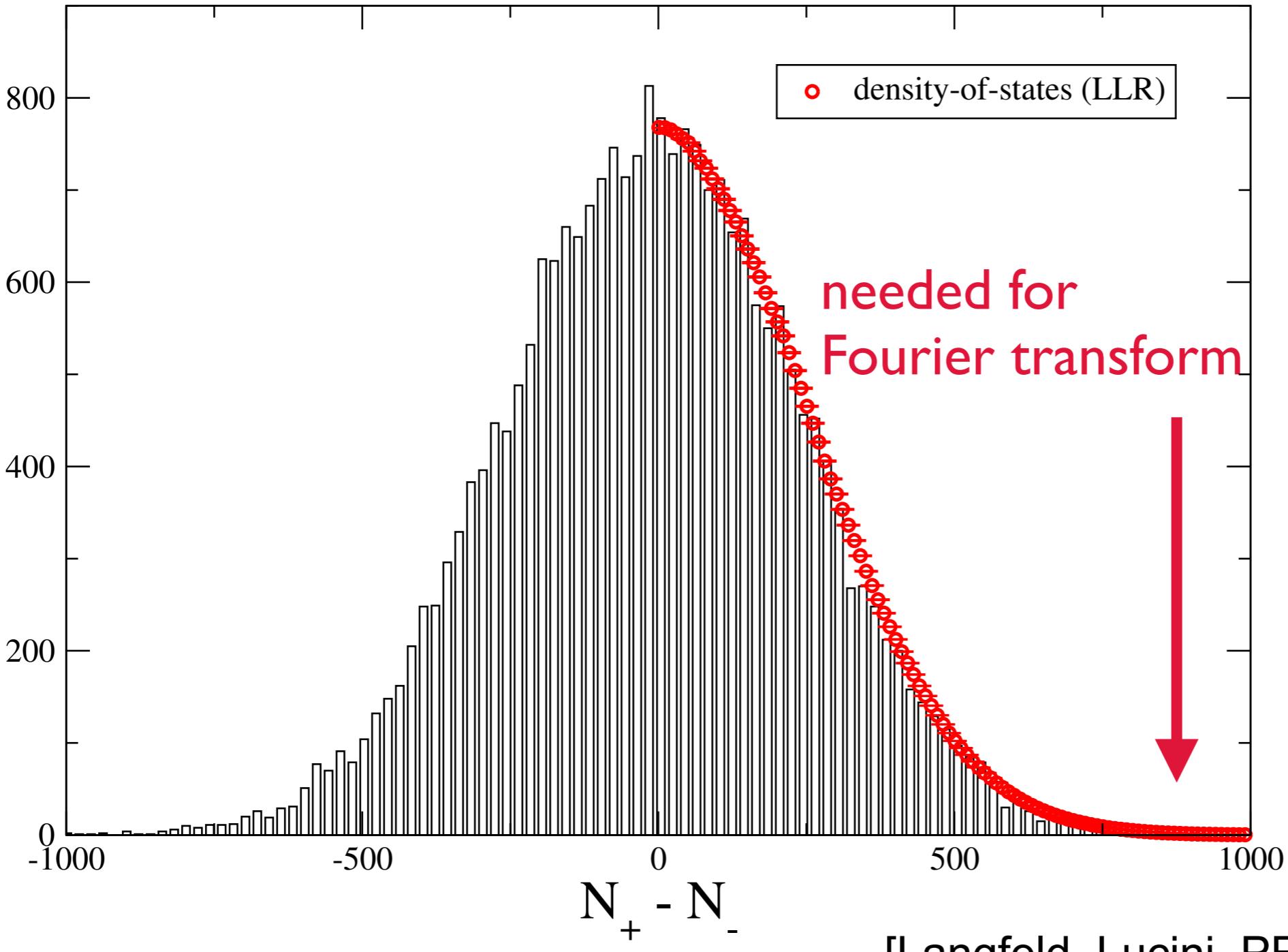
What do we find for $P(s)$?

Polyakov spin model: 24^3 tau=0.17 (*2), kappa=0.05

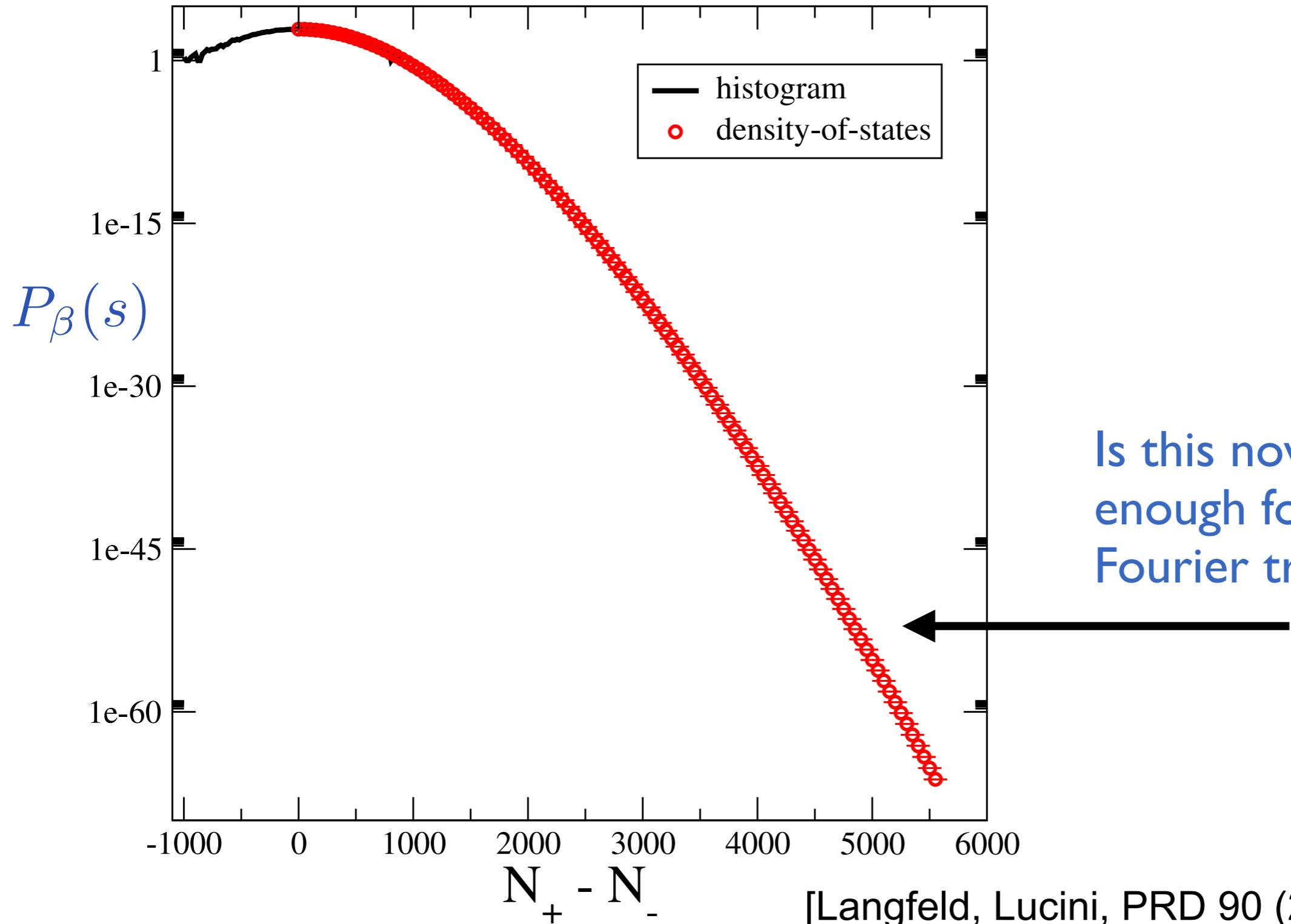


Let's go back to $P(\text{imaginary part})$ of the $Z3$ theory:

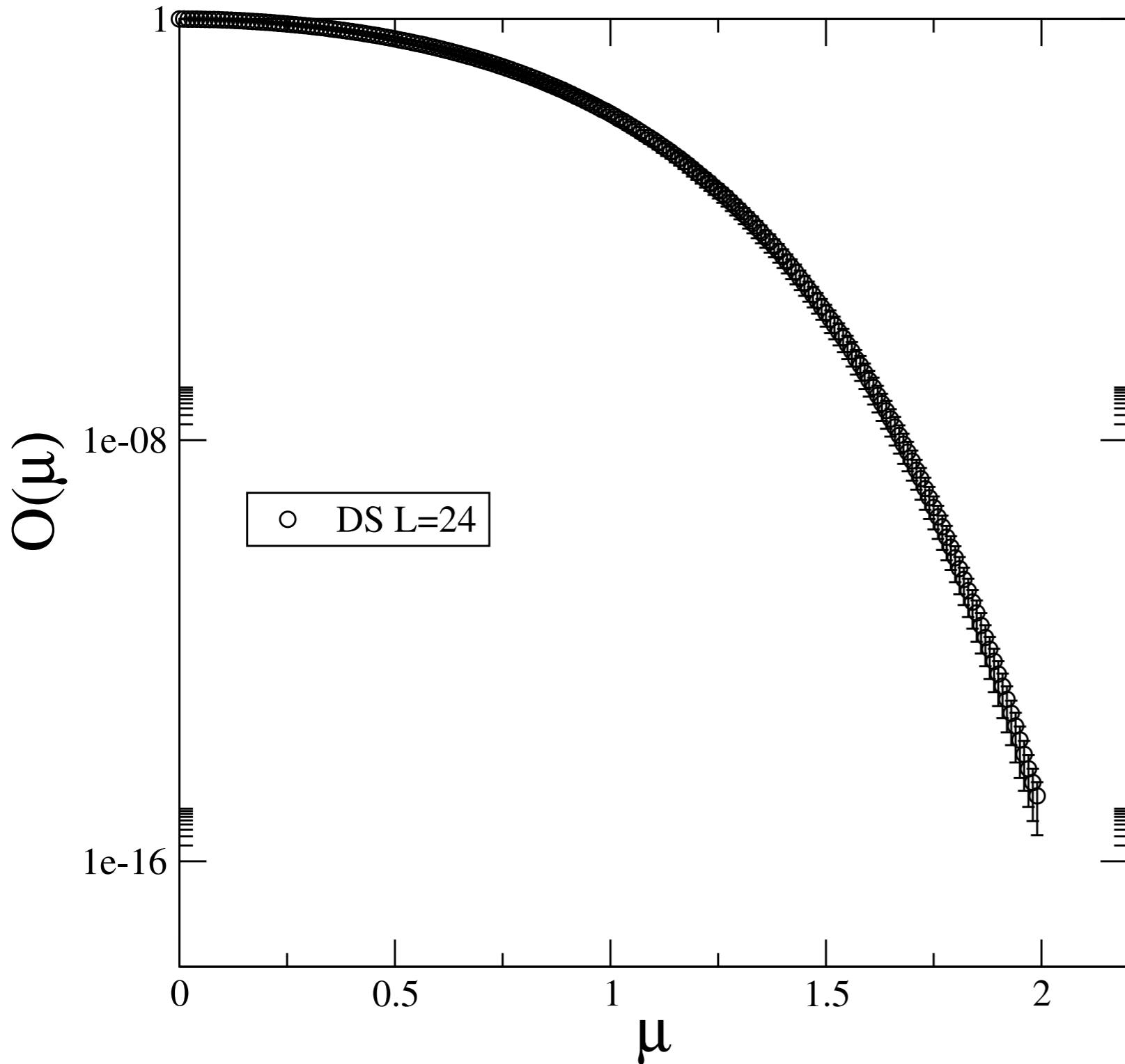
Polyakov spin model: 24^3 tau=0.17 (*2), kappa=0.05



Let's go back to $P(\text{imaginary part})$ of the Z_3 theory:



Numerical results: Z3 gauge theory

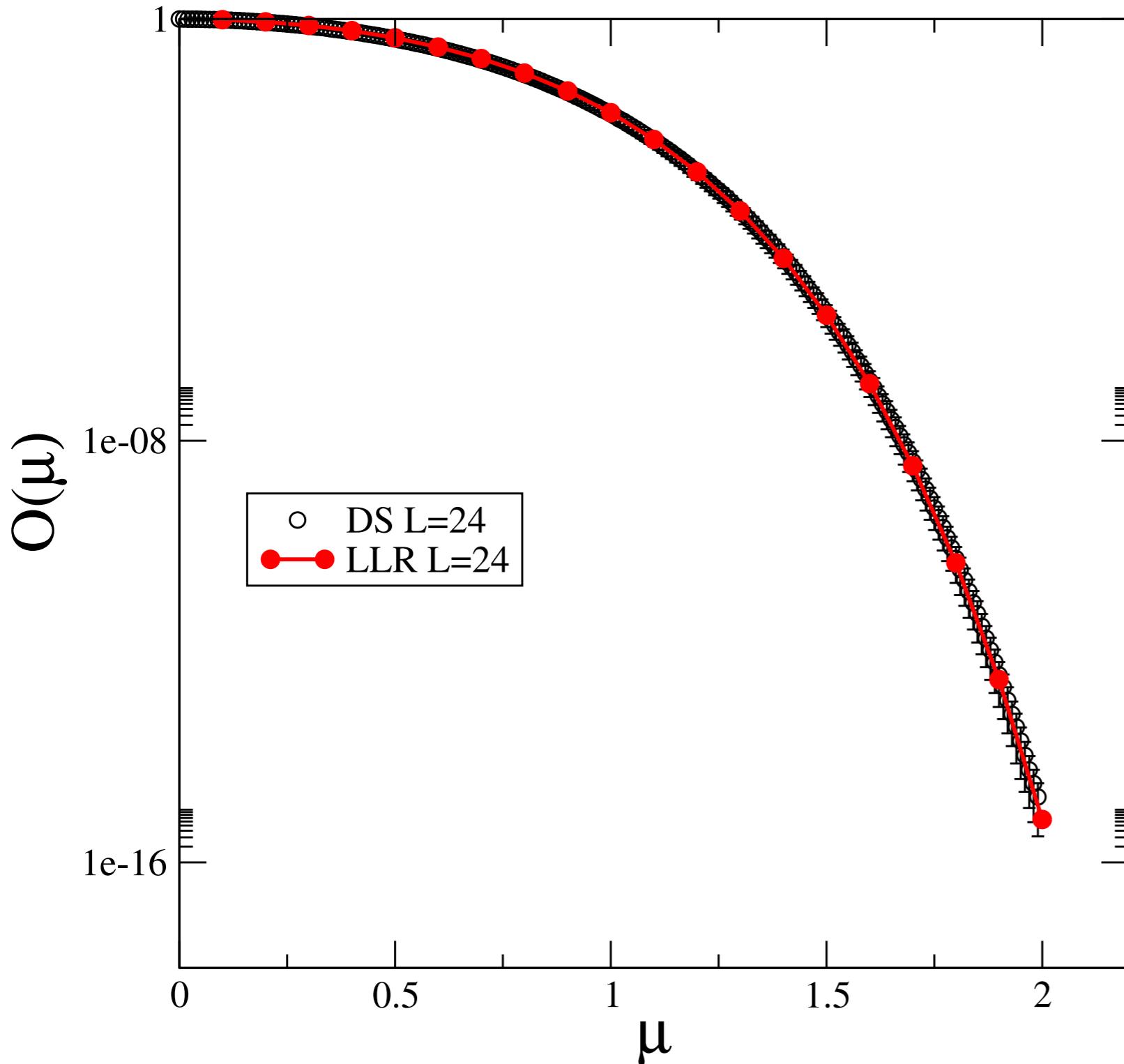


$$O(\mu) = \langle e^{i\varphi} \rangle_{\text{PQ}}$$

Result from the
(real) dual theory:
“snake algorithm”



Numerical results: Z3 gauge theory



$$O(\mu) = \langle e^{i\varphi} \rangle_{\text{PQ}}$$

Result from the
(real) dual theory:
“snake algorithm”



Excellent agreement!

*First “head on”
solution of a sign
problem!*

[Langfeld, Lucini, PRD9 (2014) 094502]

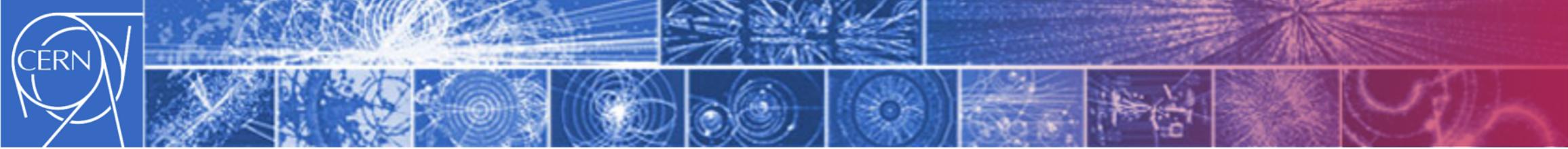
Anatomy of a sign problem: Heavy-Dense QCD (HDQCD)

Starting point $Z(\mu) = \int \mathcal{D}U_\mu \exp\{\beta S_{\text{YM}}[U]\} \text{Det}M(\mu)$

QCD:

Limit quark mass m, μ large, $\mu/m \rightarrow$ finite

[Bender, Hashimoto, Karsch, Linke, Nakamura, Plewnia,
Nucl. Phys. Proc. Suppl. 26 (1992) 323]



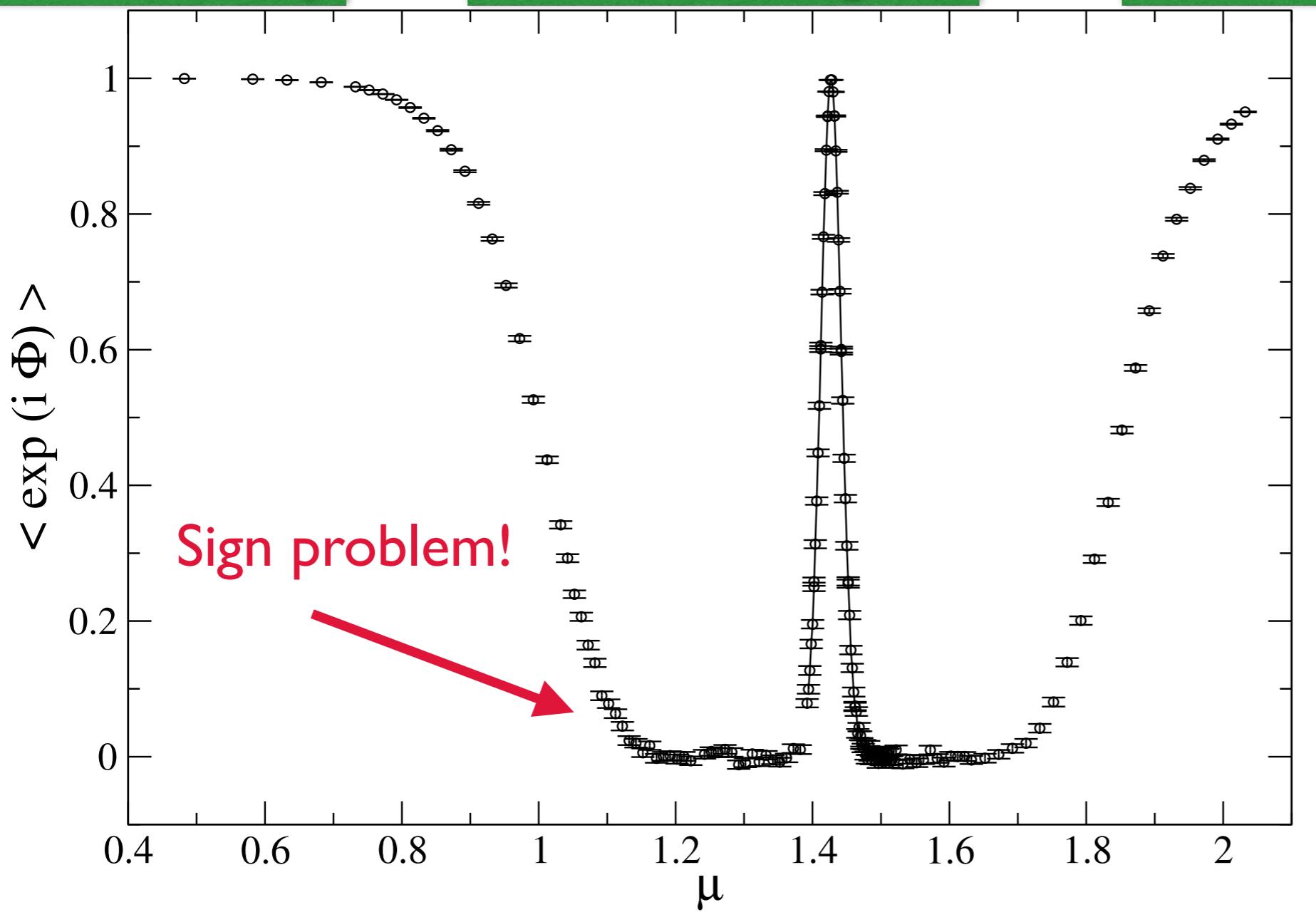
Heavy-dense QCD

[Garron, Langfeld, et al, in preparation]

low density

“half-filling”

saturation



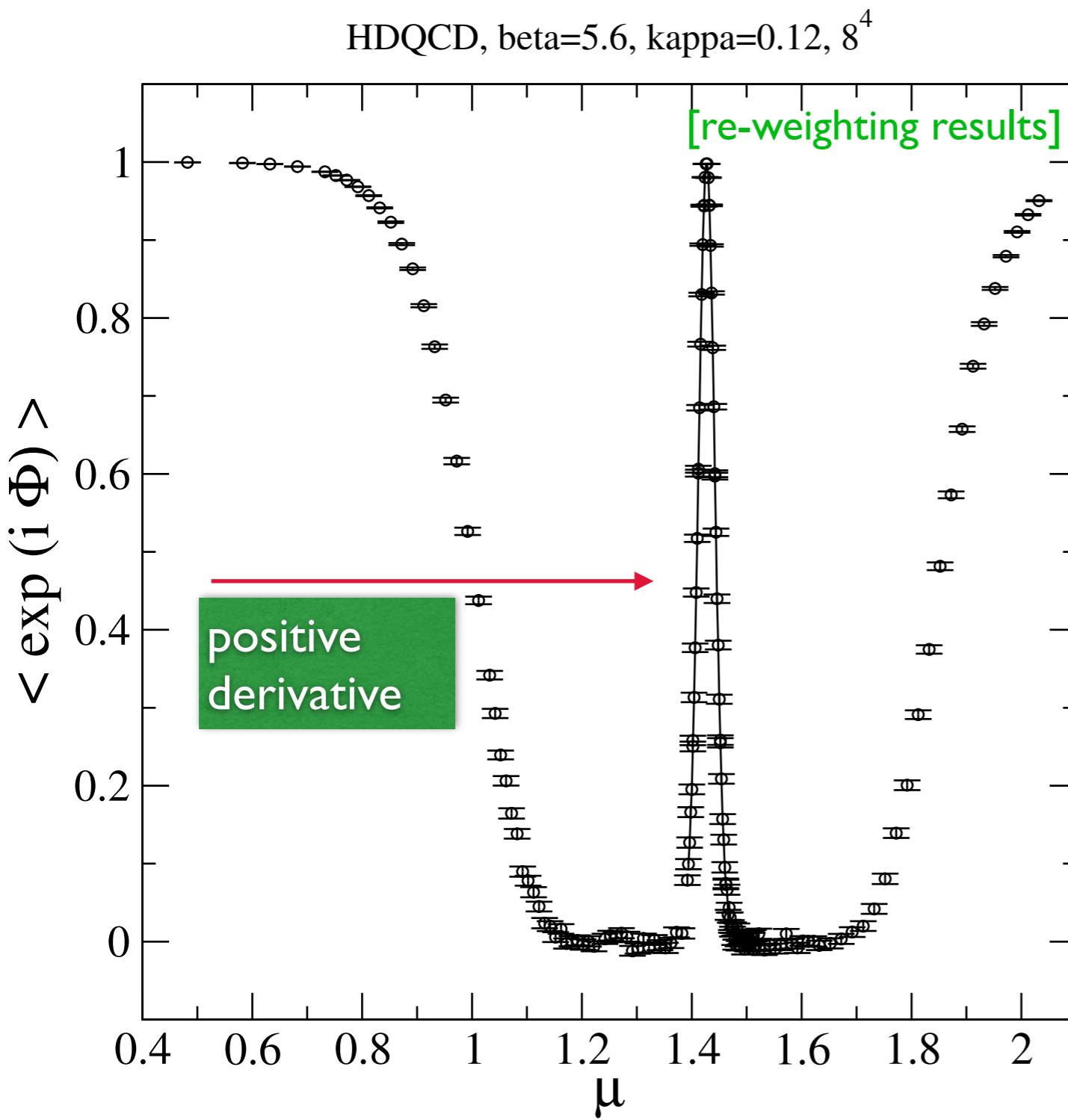
see also [Rindlisbacher, de Forcrand, JHEP 1602 (2016) 051]

“half-filling”
“ $\mu = m$ ”
(physical)

Is QCD
real at
threshold
 $\mu = m_B/3$?

Heavy-dense QCD

Inverse Silver Blaze feature



Recall:

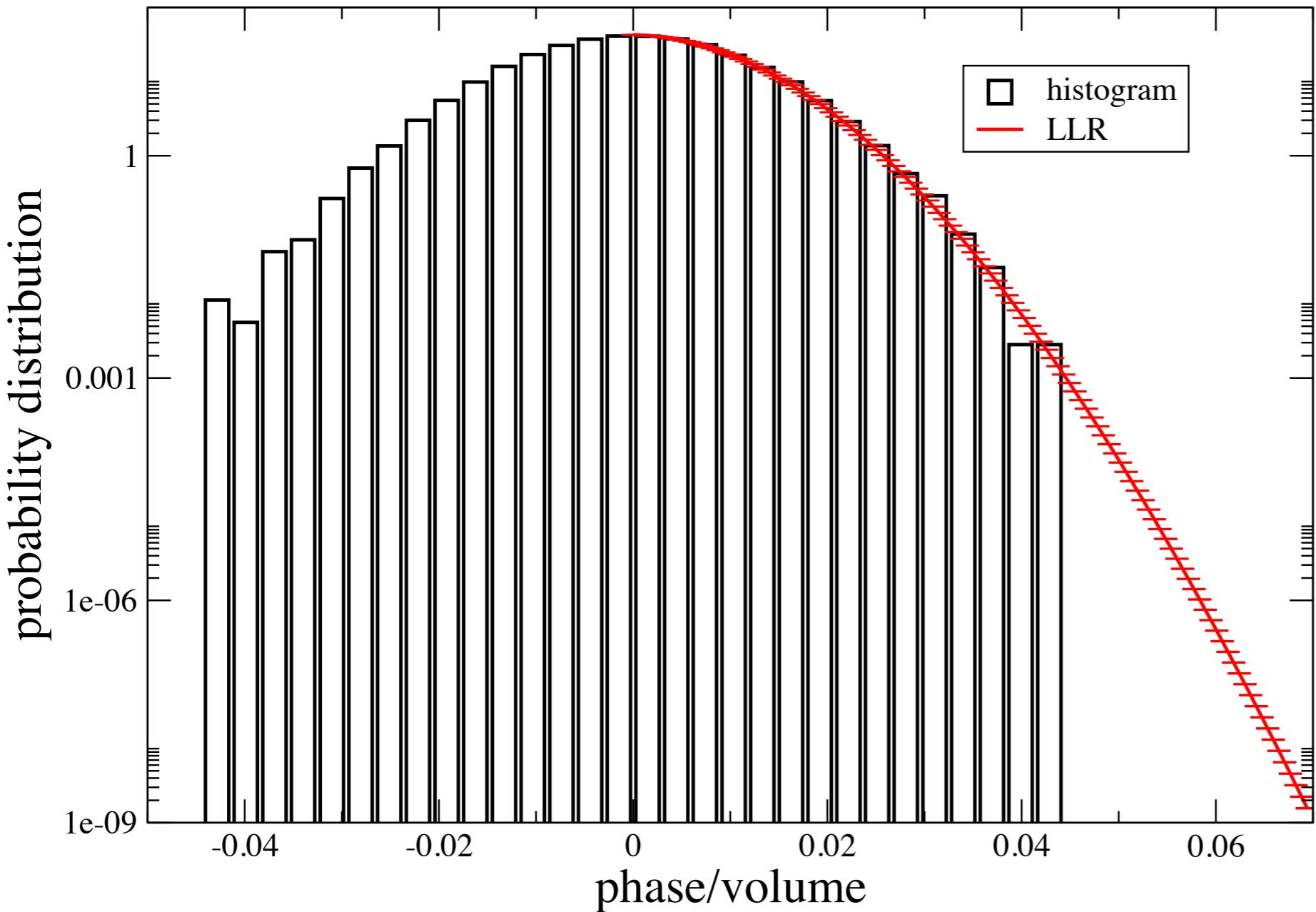
$$\rho(\mu) = \frac{T}{V_3} \frac{\partial}{\partial \mu} O(\mu) + \rho_{\text{PQ}}(\mu)$$

Phase quenching
underestimates the
true density!

What can LLR do for you?

$$\mu = 1.3321$$

SU(3) 8^4 , kappa=0.12, beta=5.8, mu = 1.3321



$P_\beta(s)$
LLR
exponential
error
suppression !

Challenge:

How do we carry out a Fourier transform the result of which is 10^{-14} and the integrand of order $\mathcal{O}(1)$ is only known numerically?

Fit a Polynomial: $\ln P(s) = \sum_{i \text{ even}}^p c_i s^i$, for $p = 2, 4, 6, 8..$

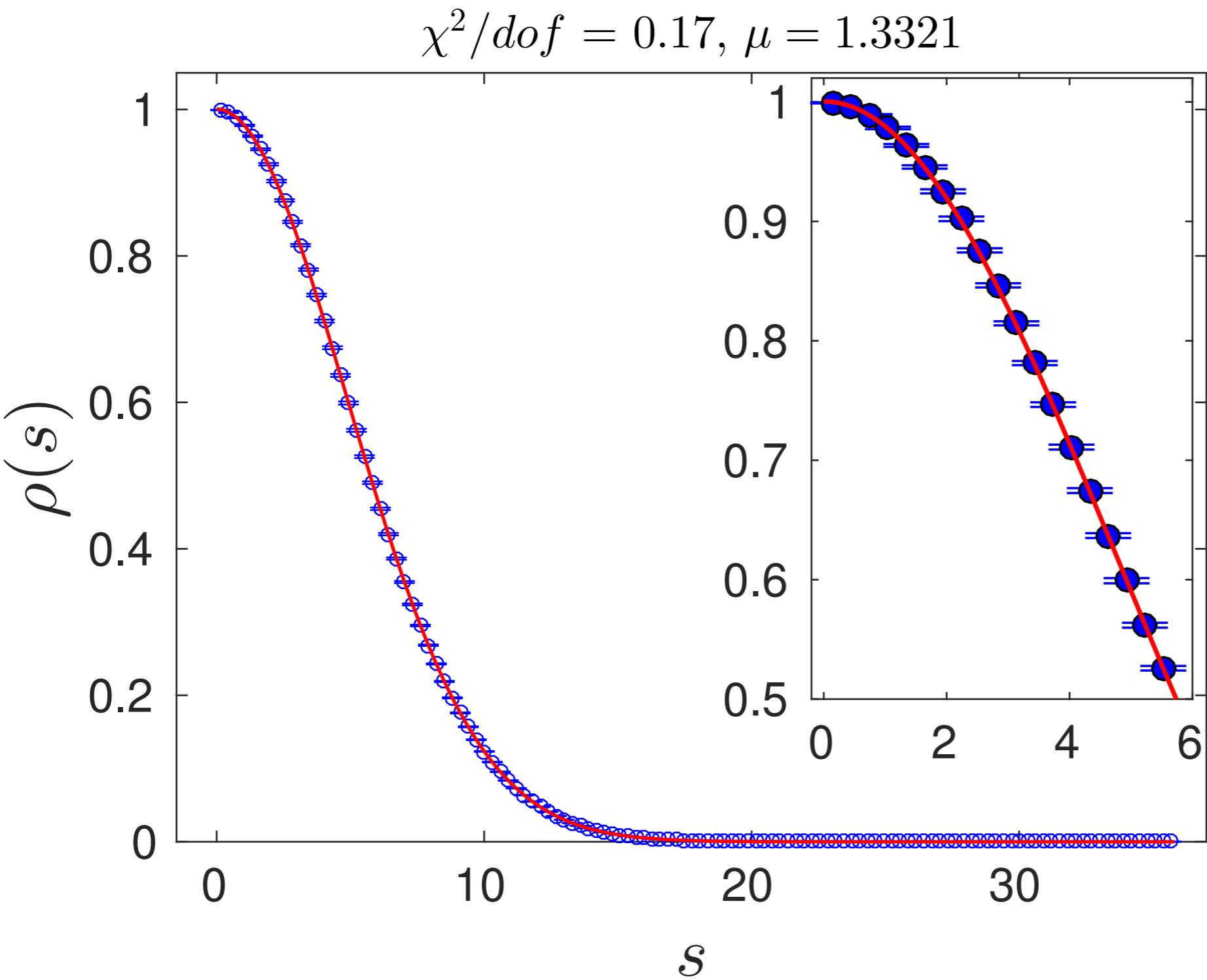
Calculate the Fourier transform semi-analytically

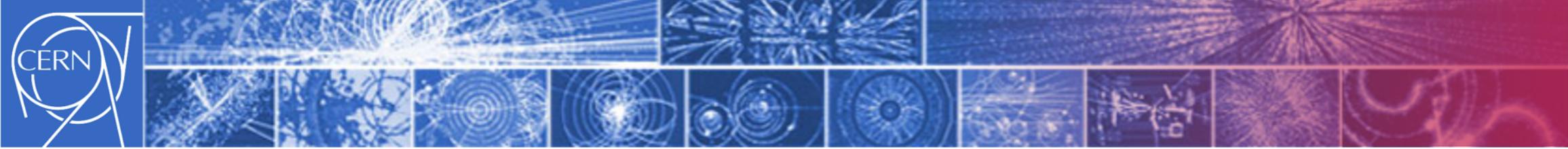
Compressed Sensing:

$\ln P(s) \sim 1000$ data points $\Rightarrow c_i \sim 20$ coefficients

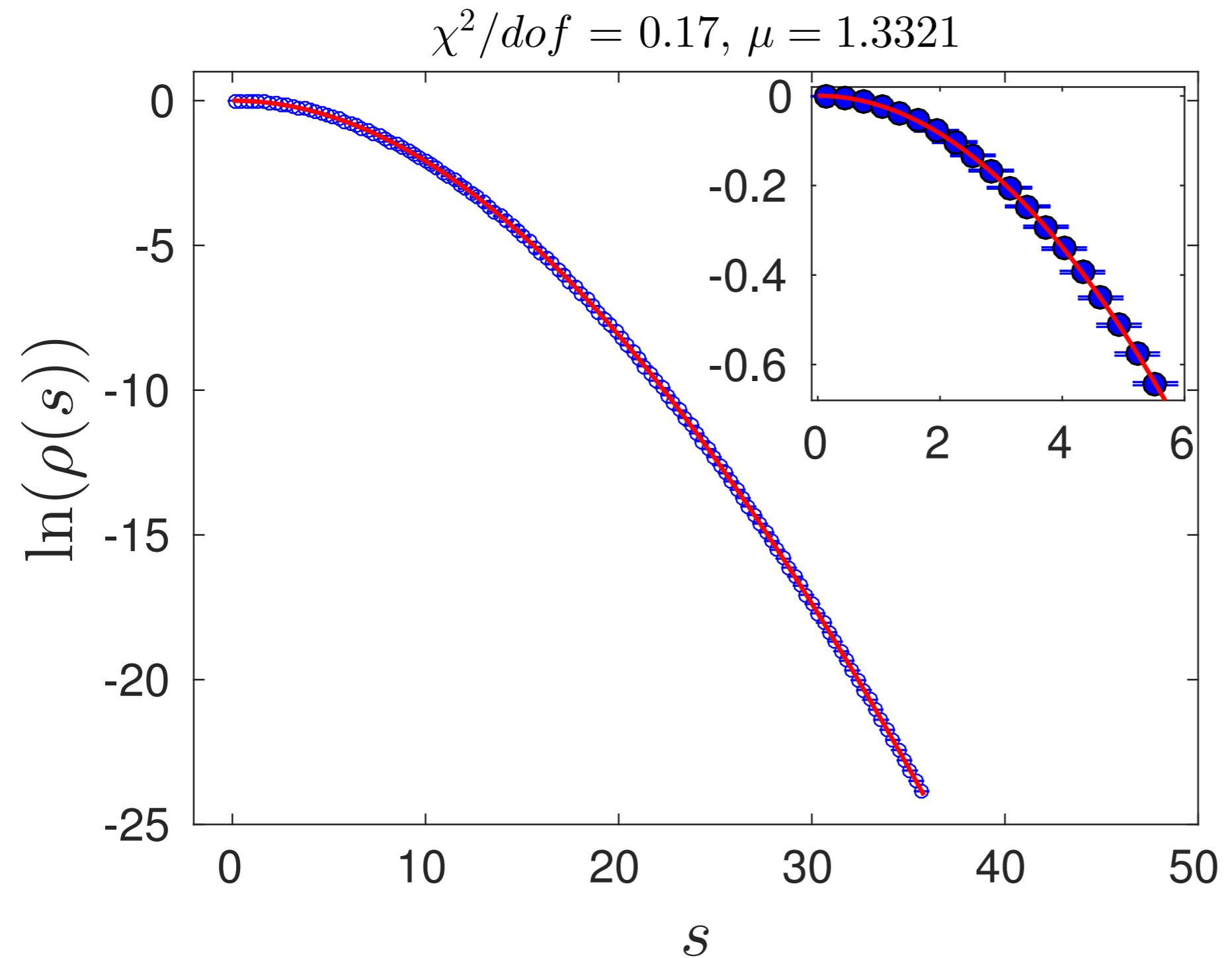
$$\chi^2/\text{dof} = \mathcal{O}(1)$$

Works very well!

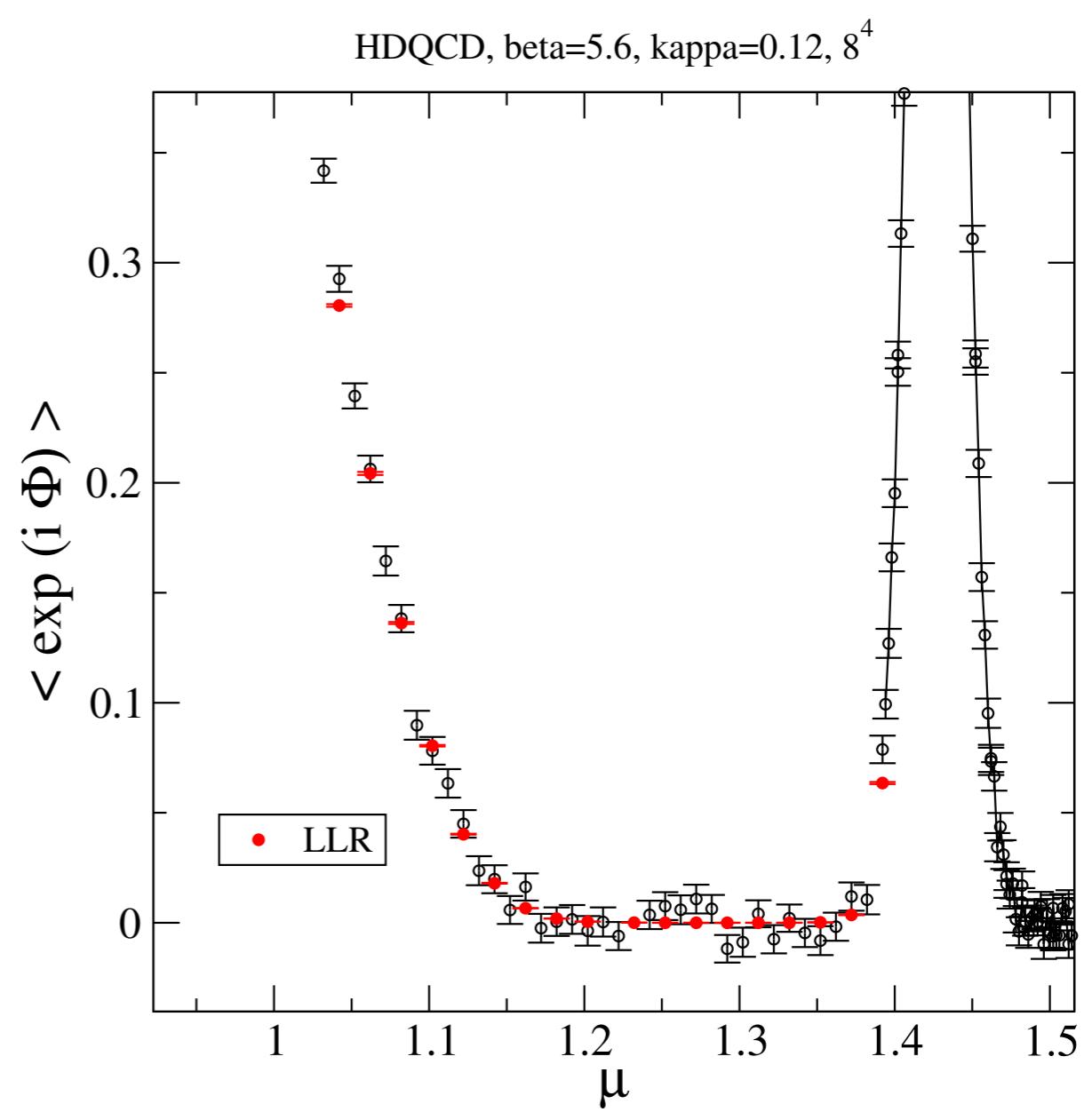
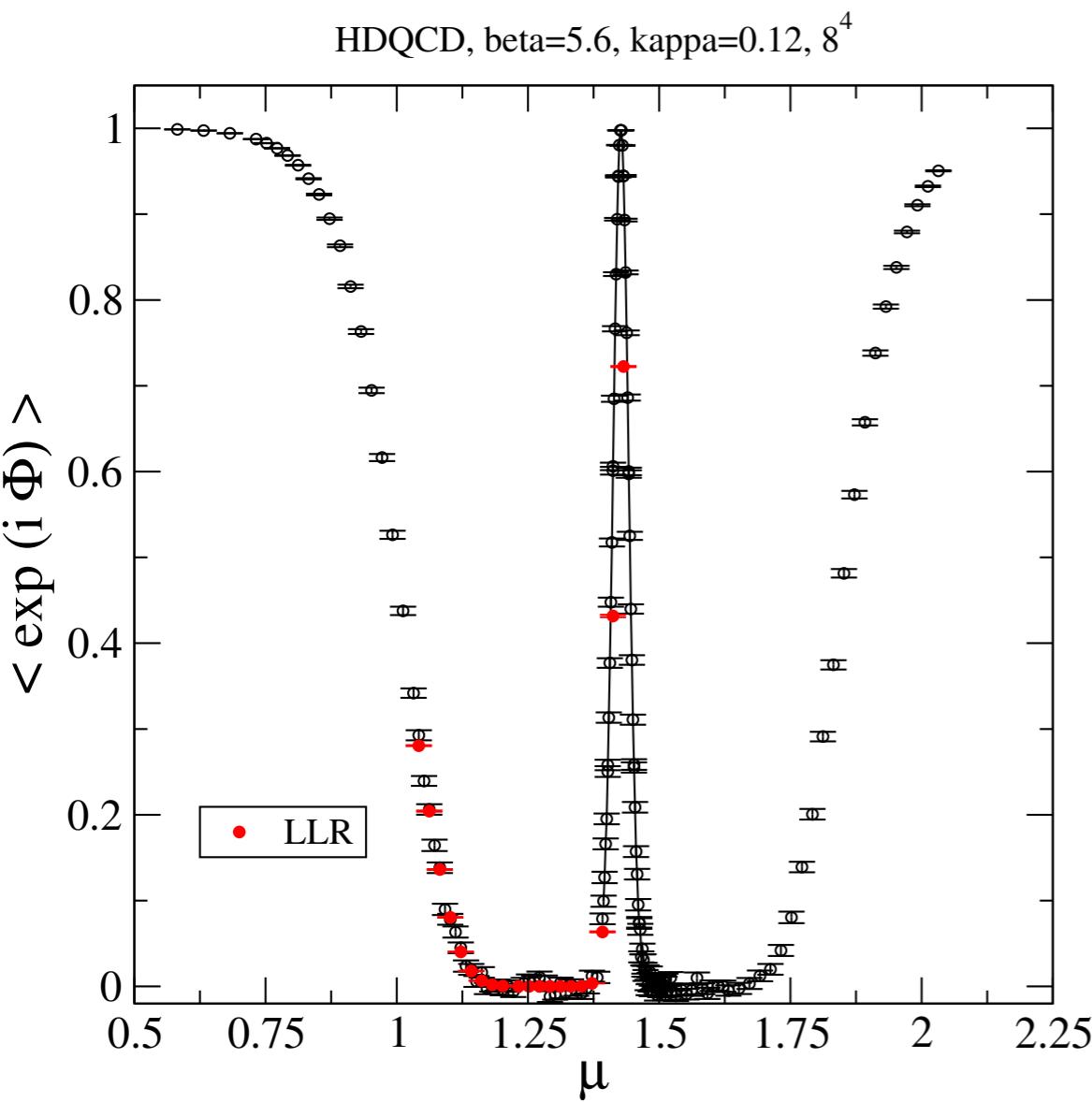




Works very well!

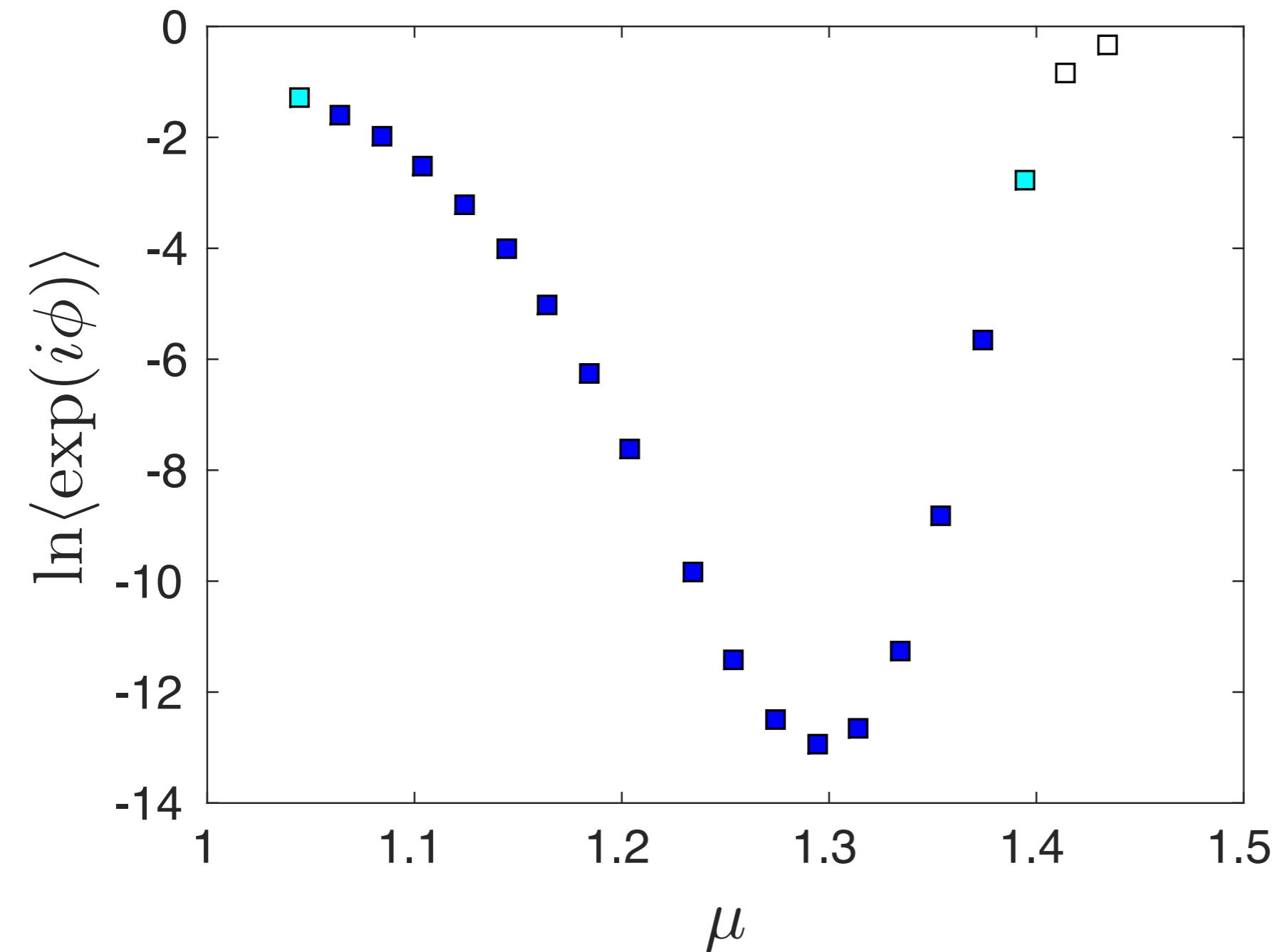


What can LLR do for you?



[Garron, Langfeld, et al, in preparation]

What can LLR do for you?

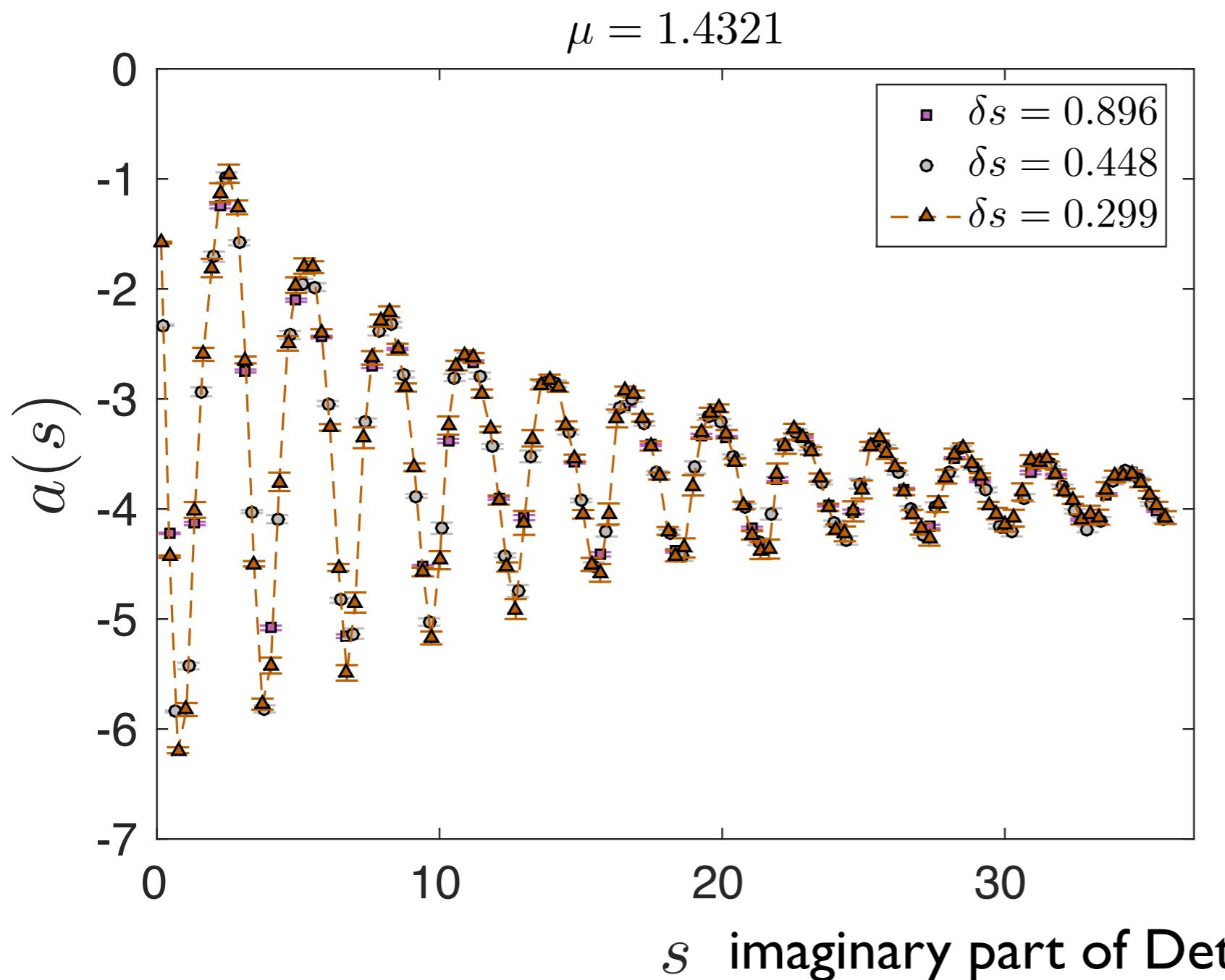


Our final
result for the
Overlap
in HQCD

[Garron, Langfeld, et al, in preparation]

And some surprises...

...close to onset $\mu \lesssim m$

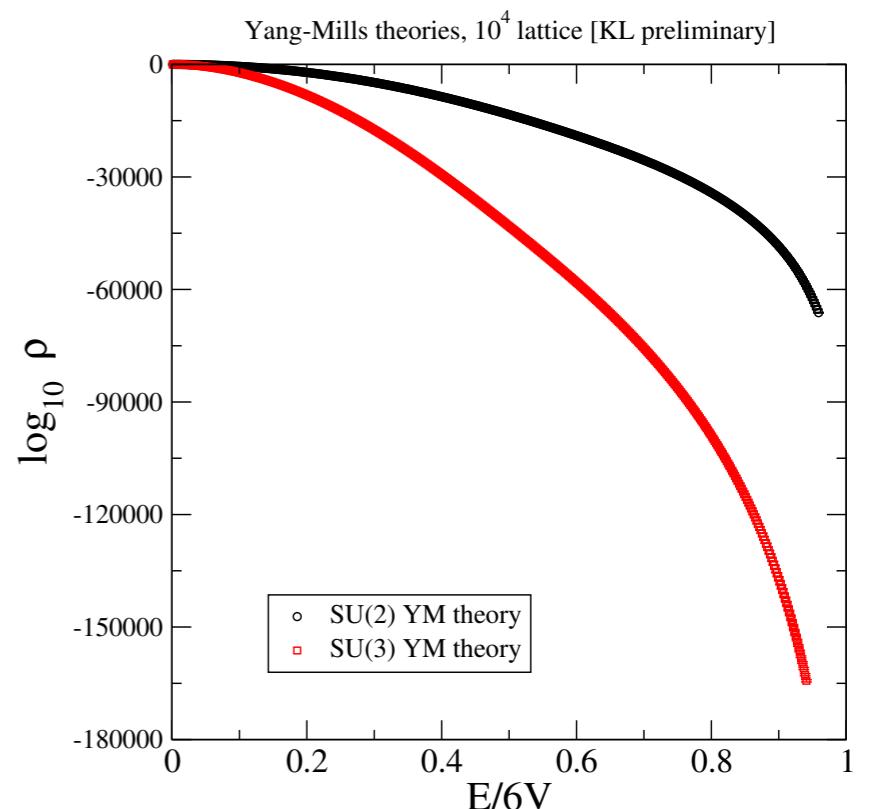
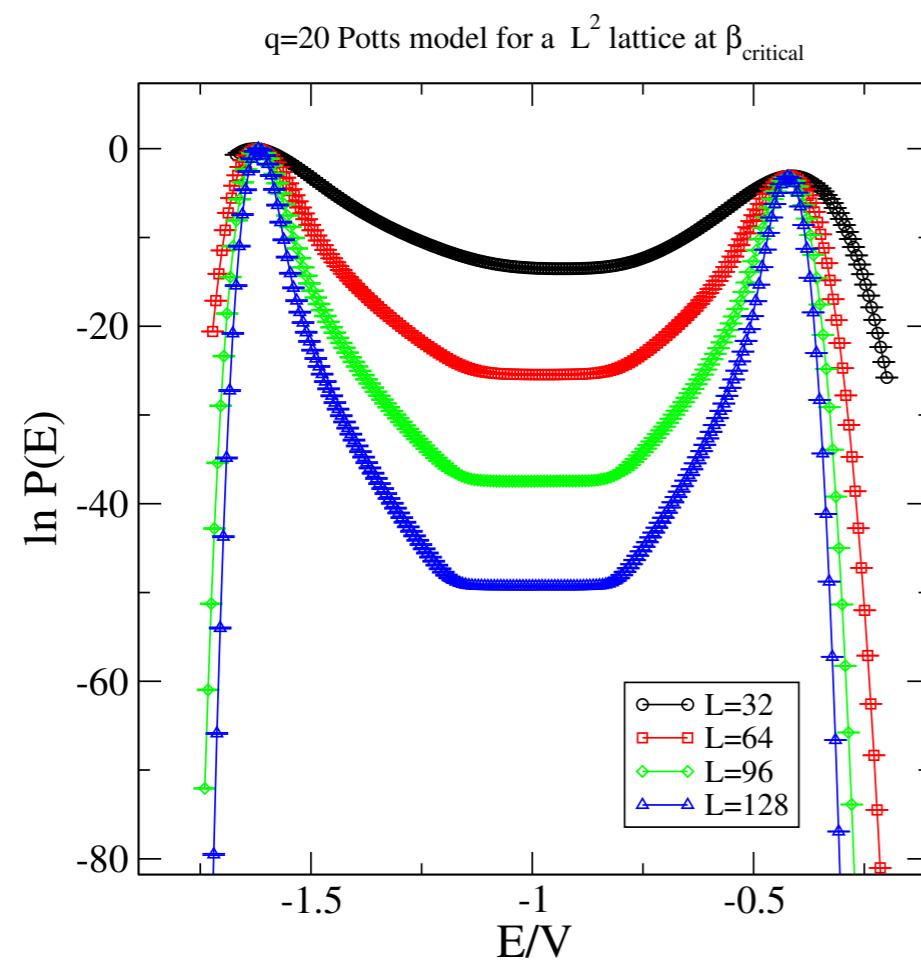


[Garron, Langfeld, et al, in preparation]

Summary:

What is the LLR approach?

Calculates the probability distribution of
(the imaginary part of) the action with
exponential error suppression



Non-Markovian process (does not
rely on importance sampling)

⇒ solves overlap problems

Summary:

Can solve *strong sign problems*:

Z3 gauge theory at finite densities
HD QCD

Open questions:

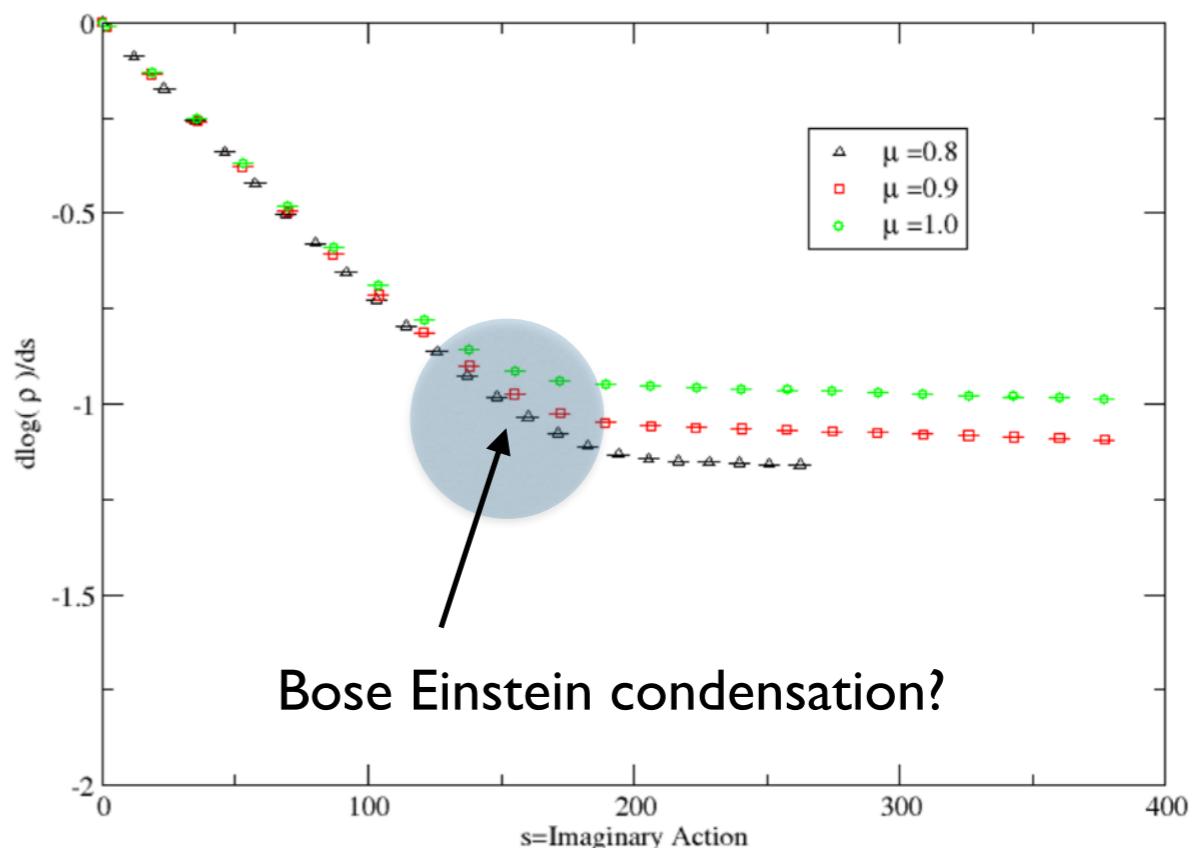
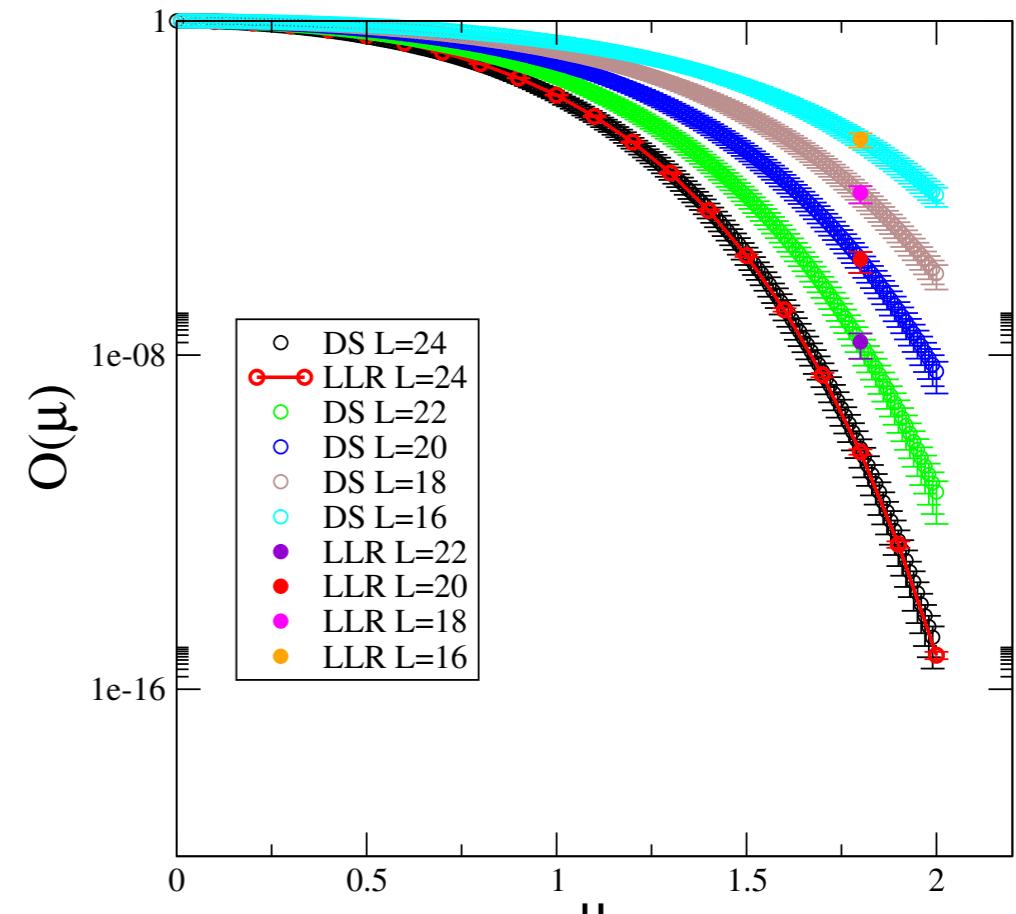
Hinges on the ability to “compress information”

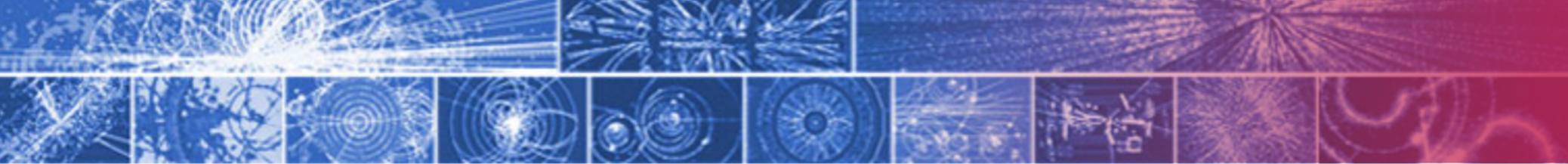
$\ln P(s) \sim 1000$ data points

$\Rightarrow c_i \sim 24$ coefficients

difficult for ϕ^4 theory!

[Bongiovanni, Langfeld, Lucini, Pellegrini, Rago, Lattice 2015]





Conclusion:

Promising first results for finite density QFT!

More new ideas for “compressing information” is needed (FT)

Thank you!