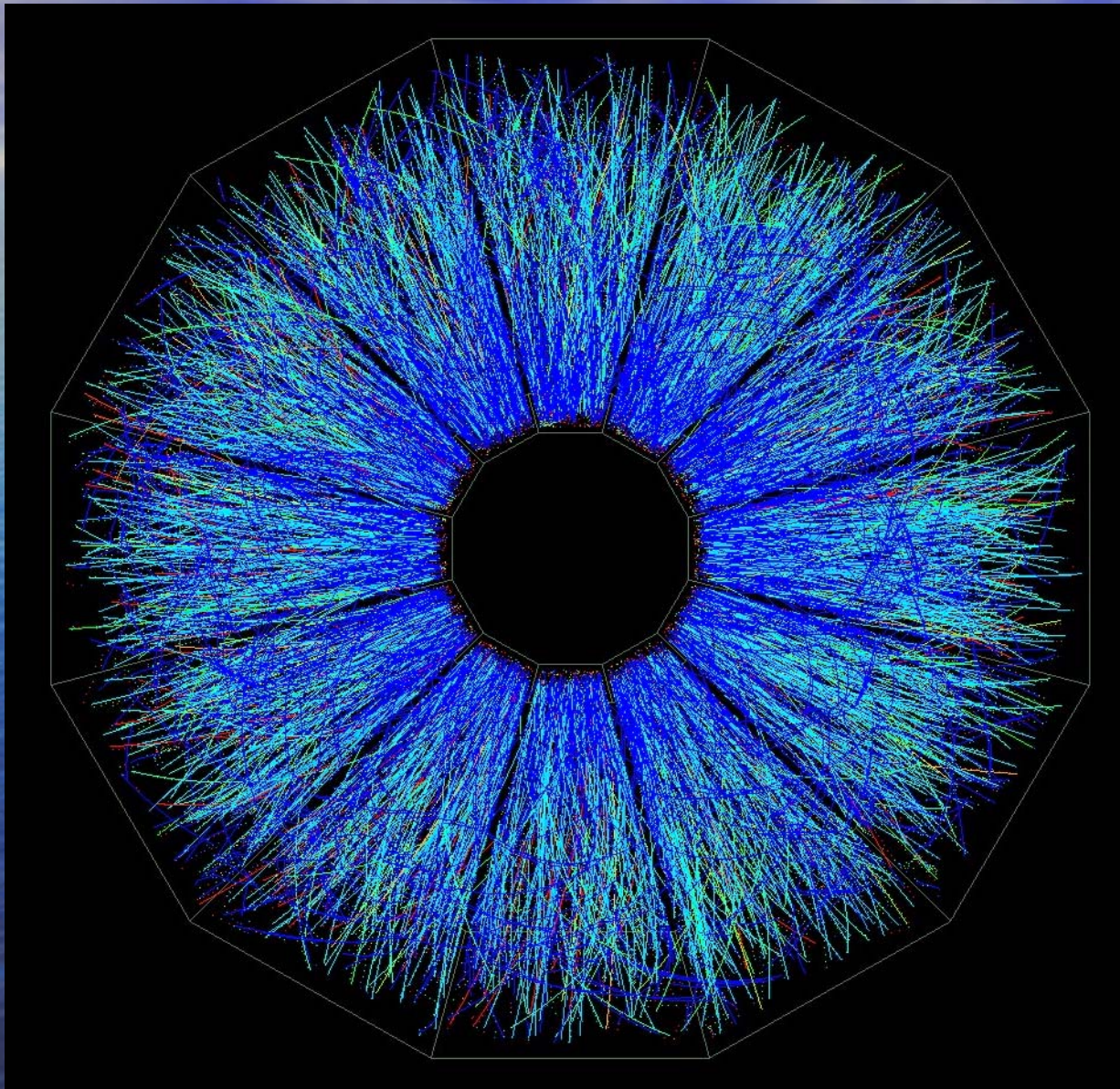


Has the
critical temperature of the
QCD phase transition
been measured ?

Heavy ion collision



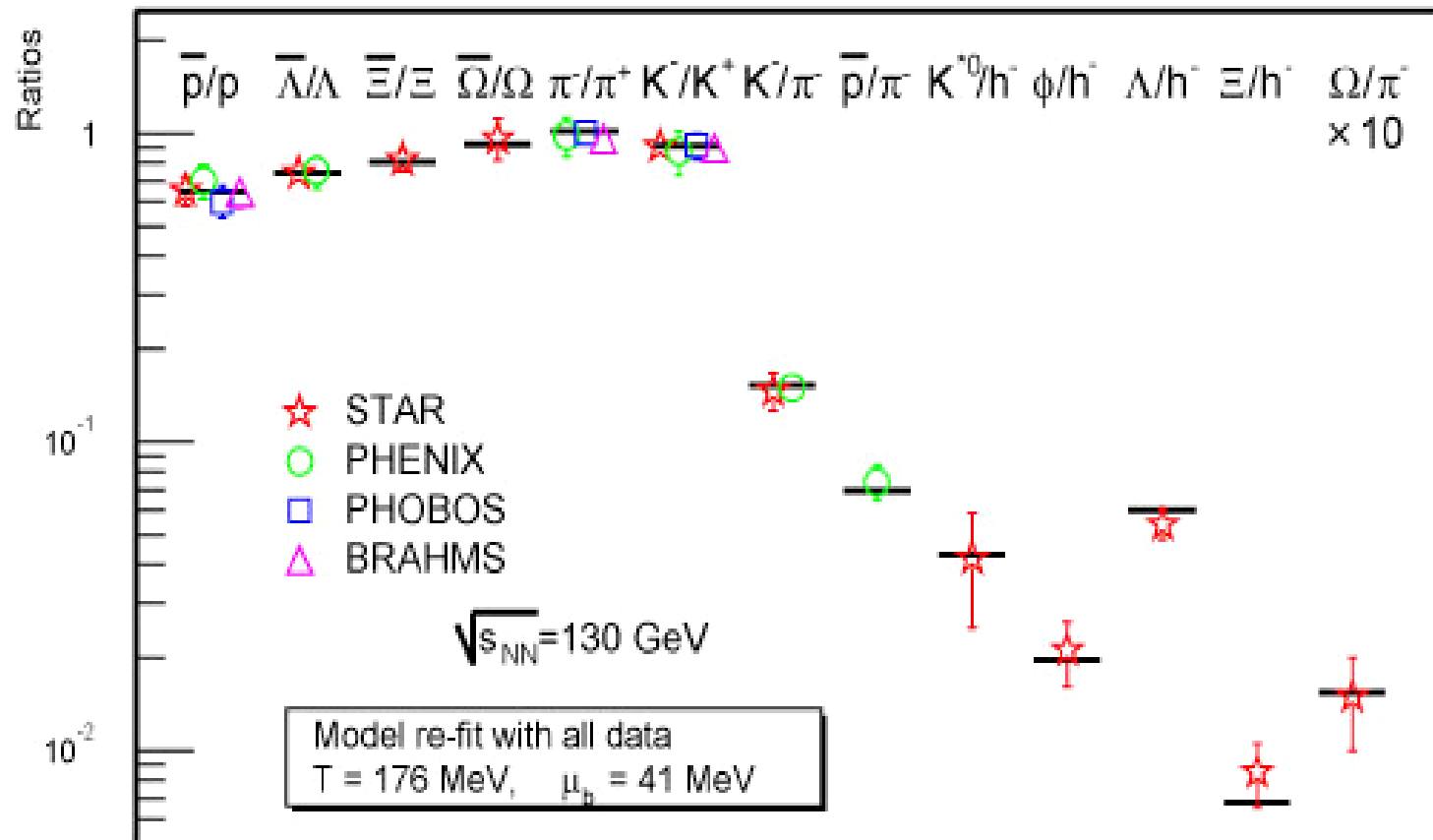
Yes !

$$0.95 T_c < T_{ch} < T_c$$

- not : " I have a model where $T_c \approx T_{ch}$ "
- not : " I use T_c as a free parameter and find that in a model simulation it is close to the lattice value (or T_{ch}) "

$$T_{ch} \approx 176 \text{ MeV} \quad (?)$$

Hadron abundancies



Braun-Munzinger et al., PLB 518 (2001) 41 D. Magestro (updated July 22, 2002)

Has T_c been measured ?

- Observation : statistical distribution of hadron species with "chemical freeze out temperature " $T_{ch}=176$ MeV
- T_{ch} cannot be much smaller than T_c : hadronic rates for $T < T_c$ are too small to produce multistrange hadrons ($\Omega, \dots$)
- Only near T_c multiparticle scattering becomes important (collective excitations ...) – proportional to high power of density



$$T_{ch} \approx T_c$$

Exclusion argument

Assume T is a meaningful concept -
complex issue, to be discussed later

$T_{ch} < T_c$: hadrochemical equilibrium

Exclude T_{ch} much smaller than T_c :

say $T_{ch} > 0.95 T_c$

$$0.95 < T_{ch} / T_c < 1$$

Estimate of critical temperature

For $T_{\text{ch}} \approx 176 \text{ MeV}$:

$$0.95 < T_{\text{ch}} / T_c$$

- $176 \text{ MeV} < T_c < 185 \text{ MeV}$

$$0.75 < T_{\text{ch}} / T_c$$

- $176 \text{ MeV} < T_c < 235 \text{ MeV}$

Quantitative issue matters!

needed :

lower bound on

$$T_{ch} / T_c$$

Key argument

- Two particle scattering rates not sufficient to produce Ω
- “multiparticle scattering for Ω -production” : dominant only in immediate vicinity of T_c

Mechanisms for production of multistrange hadrons

Many proposals

- Hadronization
- Quark-hadron equilibrium
- Decay of collective excitation (σ – field)
- Multi-hadron-scattering

Different pictures !

Hadronic picture of Ω - production

Should exist, at least semi-quantitatively, if $T_{\text{ch}} < T_c$
(for $T_{\text{ch}} = T_c$: $T_{\text{ch}} > 0.95 T_c$ is fulfilled anyhow)

e.g. collective excitations $\approx$ multi-hadron-scattering
(not necessarily the best and simplest picture)

multihadron $\rightarrow \Omega + X$ should have sufficient rate

Check of consistency for many models

Necessary if $T_{\text{ch}} \neq T_c$ and temperature is defined

Way to give quantitative bound on T_{ch} / T_c

Rates for multiparticle scattering

2 pions + 3 kaons $\rightarrow \Omega$ + antiproton

$$r(n_{in}, n_{out}) = \bar{n}(T)^{n_{in}} |\mathcal{M}|^2 \phi$$

$$\phi = \prod_{k=1}^{n_{out}} \left(\int \frac{d^3 p_k}{(2\pi)^3 (2E_k)} \right) (2\pi)^4 \delta^4 \left(\sum_k p_k^\mu \right)$$

$$r_\Omega = n_\pi^5 (n_K/n_\pi)^3 |\mathcal{M}|^2 \phi.$$

Very rapid density increase

...in vicinity of critical temperature

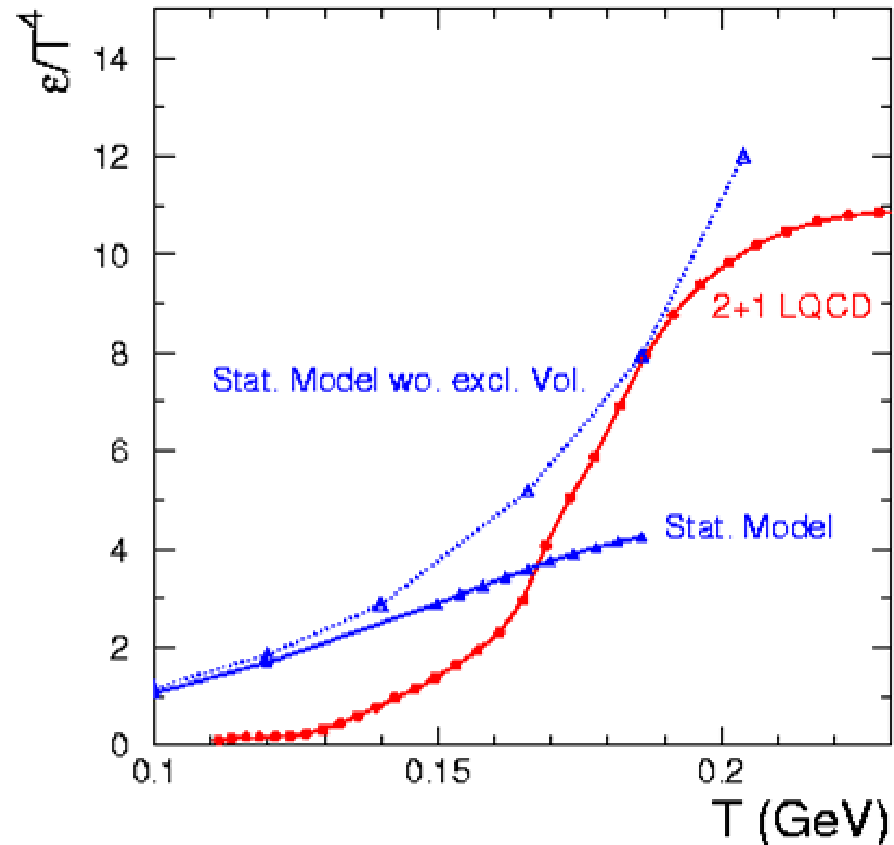
Extremely rapid increase of rate of multiparticle scattering processes

(proportional to very high power of density)

Energy density

Lattice simulations
Karsch et al

even more dramatic
for first order
transition



Phase space

- increases very rapidly with energy and therefore with temperature
- effective dependence of time needed to produce Ω

$$T_{\Omega} \sim T^{-60} !$$

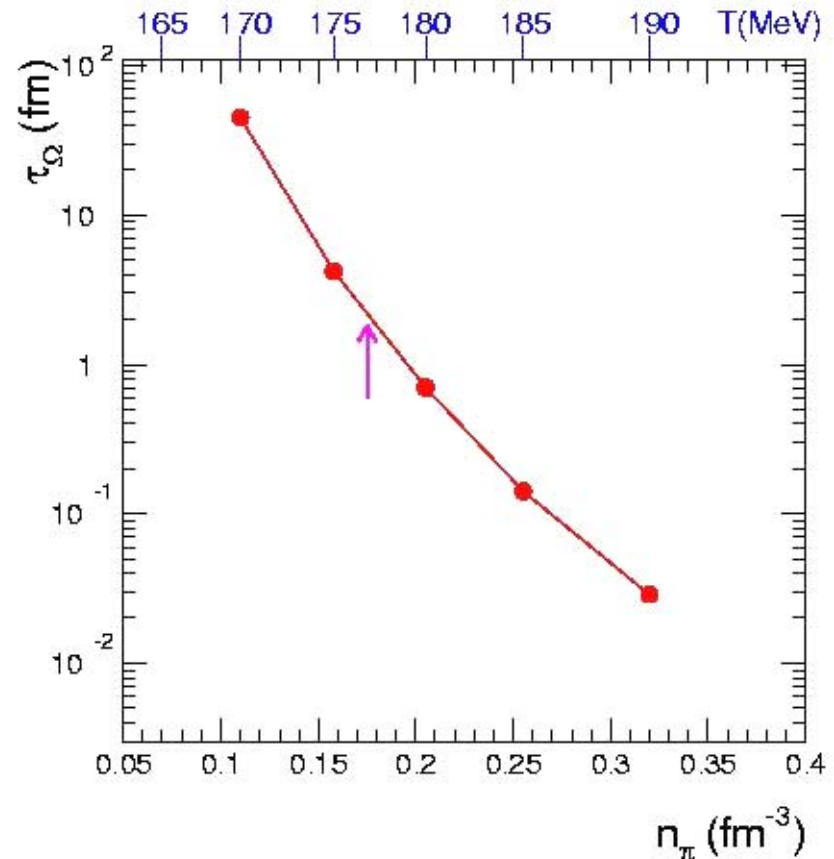
This will even be more dramatic if transition is closer to first order phase transition

Production time for Ω

multi-meson
scattering

$\pi + \pi + \pi + K + K \rightarrow$
 $\Omega + p$

strong
dependence on
pion density



P.Braun-Munzinger, J.Stachel, CW

enough time for Ω - production

at $T=176$ MeV :

$$\tau_{\Omega} \sim 2.3 \text{ fm}$$

consistency !

extremely rapid change

lowering T by 5 MeV below critical temperature :

rate of Ω – production decreases by
factor 10

This restricts chemical freeze out to close vicinity
of critical temperature

$$0.95 < T_{\text{ch}} / T_c < 1$$

Relevant time scale in hadronic phase

rates needed for equilibration of Ω and kaons:

$$\bar{r}_j = \frac{\dot{N}_j}{V} = \dot{n}_j + n_j \dot{V}/V.$$

$$\left| \frac{\bar{r}_\Omega}{n_\Omega} - \frac{\bar{r}_K}{n_K} \right| = \frac{\ln F_{\Omega K}}{\tau_T} \frac{T_{\text{ch}}}{\Delta T} = (1.10 - 0.55)/\text{fm}$$

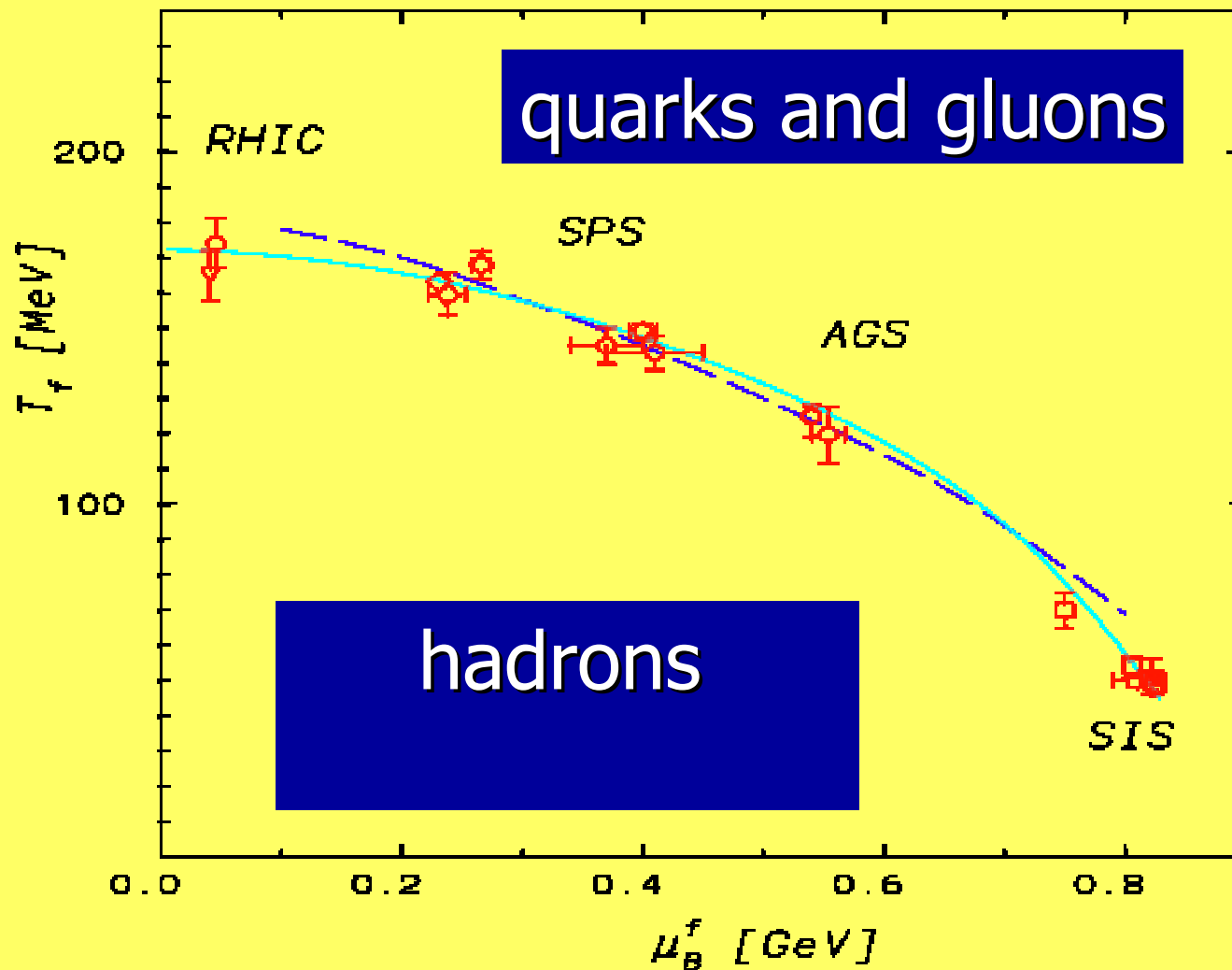
$$\begin{aligned} \Delta T &= 5 \text{ MeV}, \\ F_{\Omega K} &= 1.13, \\ \tau_T &= 8 \text{ fm} \end{aligned}$$

two –particle – scattering :

$$\left| \frac{\bar{r}_\Omega}{n_\Omega} - \frac{\bar{r}_K}{n_K} \right| = (0.02 - 0.2)/\text{fm}$$


$$T_{ch} \approx T_c$$

Phase diagram

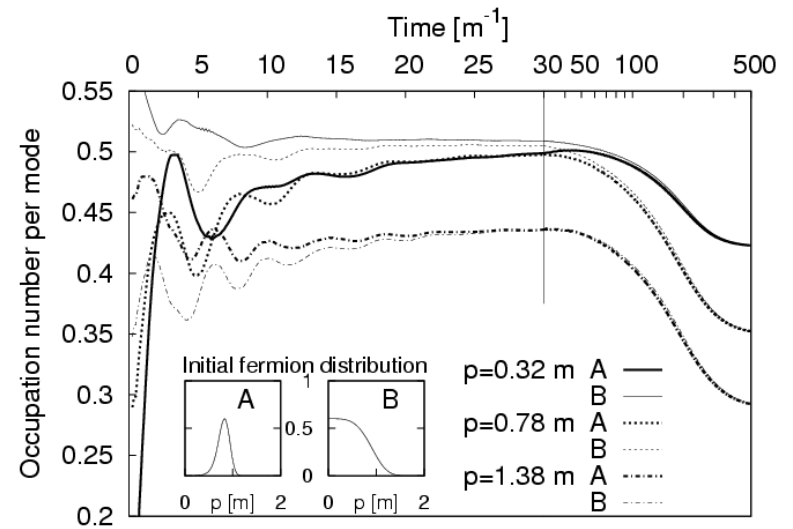
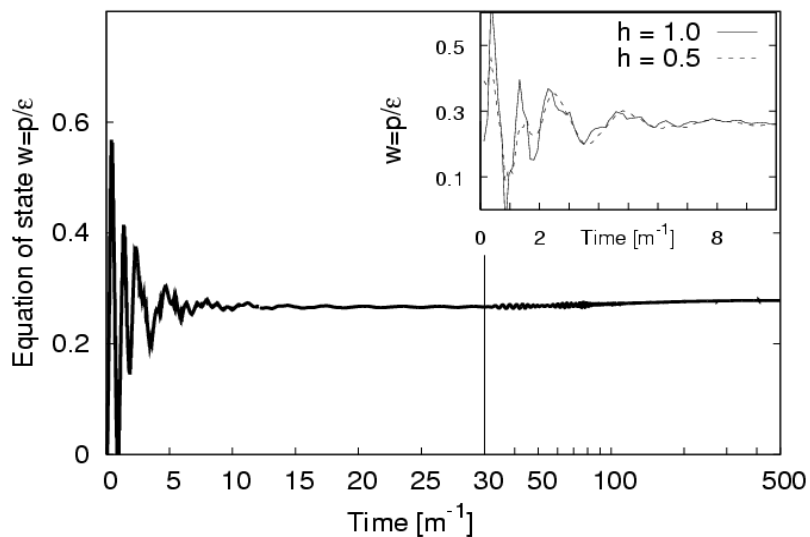


Is temperature defined ?

Does comparison with
equilibrium critical temperature
make sense ?

Prethermalization

J. Berges, Sz. Borsanyi, CW



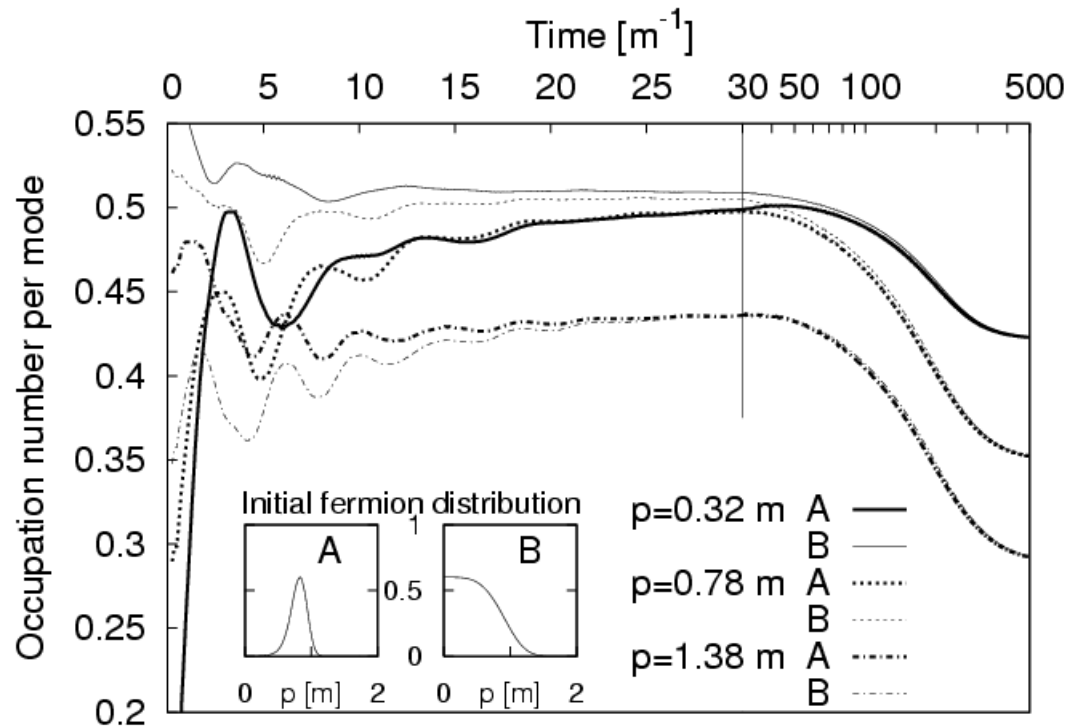
Vastly different time scales

for “thermalization” of different quantities

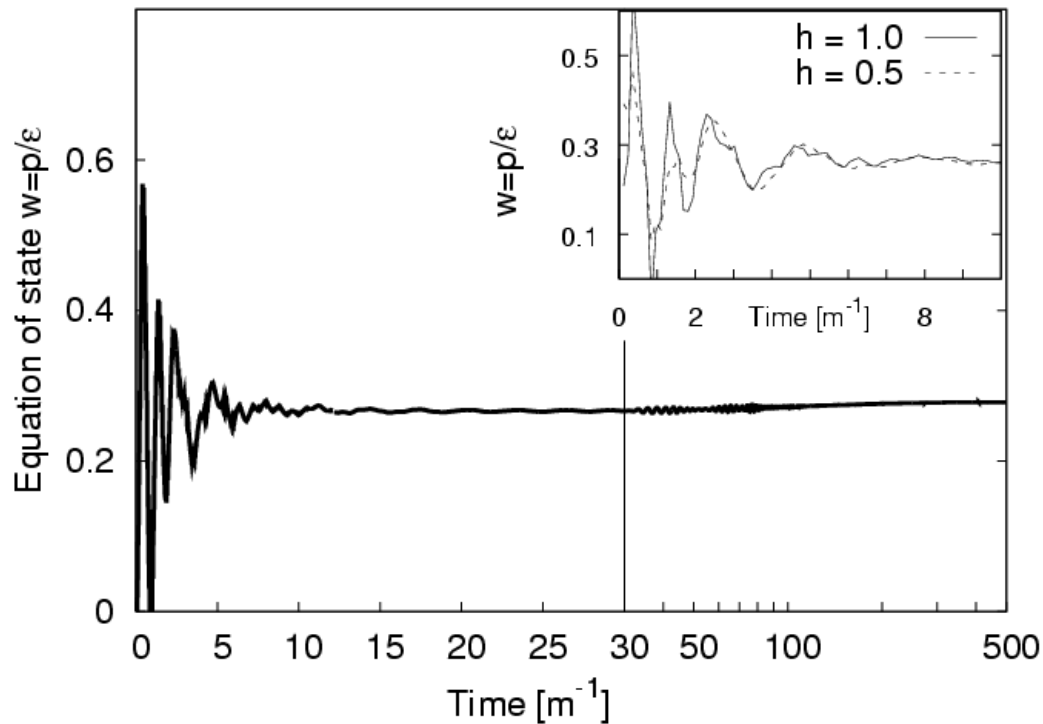
here : scalar with mass m coupled to fermions
(linear quark-meson-model)

method : two particle irreducible non-
equilibrium effective action (J.Berges et al)

Thermal equilibration : occupation numbers



Prethermalization equation of state p/ε

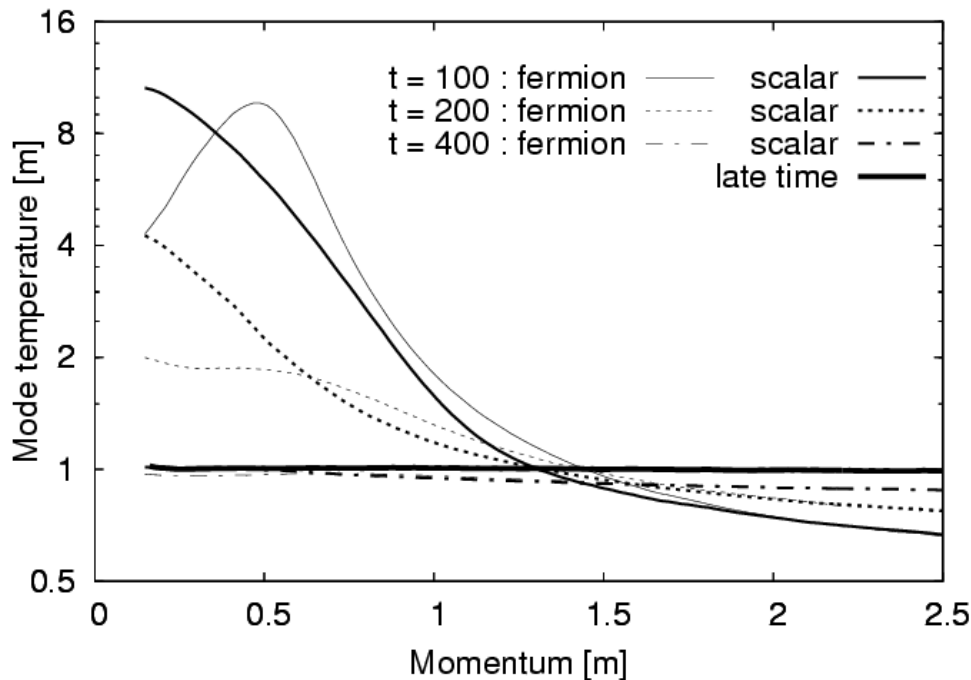


similar for kinetic temperature

A wide-angle photograph of a calm, deep blue ocean stretching to the horizon. The sky is a clear, vibrant blue with some wispy white clouds near the horizon. On the far left, a bright rainbow is visible, its colors blending into the blue of the sky and sea. The text "different 'temperatures'" is overlaid in the lower-middle part of the image.

different "temperatures"

Mode temperature



$$n_p(t) \stackrel{!}{=} \frac{1}{\exp [\omega_p(t)/T_p(t)] \pm 1}$$

$\omega_p^{(f,s)}(t)$ determined by peak of spectral function

n_p : occupation number
for momentum p

late time:

**Bose-Einstein or
Fermi-Dirac distribution**

Global kinetic temperature T_{kin}

Practical definition:

- association of temperature with average kinetic energy per d.o.f.

$$T_{\text{kin}}(t) = E_{\text{kin}}(t)/c_{\text{eq}}$$

- $c_{\text{eq}} = E_{\text{kin,eq}}/T_{\text{eq}}$ is given solely in terms of equilibrium quantities
(E.g. relativistic plasma: $E_{\text{kin}}/N = \epsilon/n = \alpha T$)

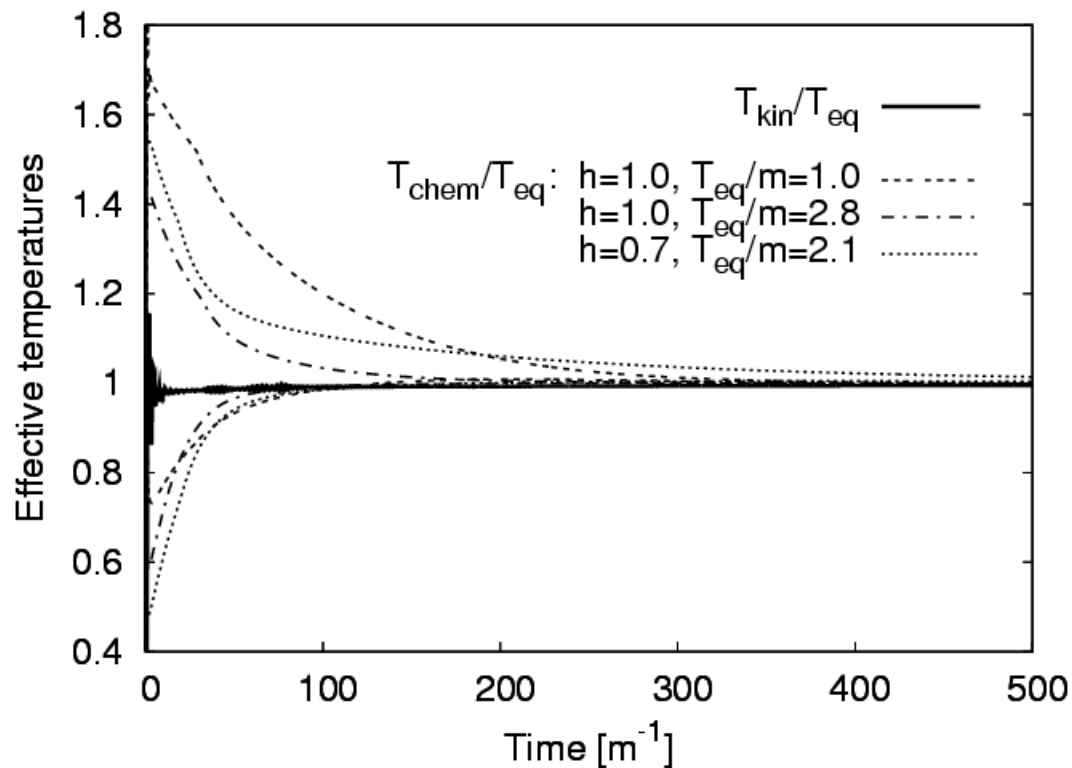
$$\text{Kinetic equilibration: } T_{\text{kin}}(t) = T_{\text{eq}}$$

Consider also *chemical temperatures* $T_{\text{ch}}^{(f,s)}$ from integrated number density of each species, $n^{(f,s)}(t) = g^{(f,s)} \int d^3p/(2\pi)^3 n_p^{(f,s)}(t)$:

$$n(t) \stackrel{!}{=} \frac{g}{2\pi^2} \int_0^\infty dp p^2 [\exp(\omega_p(t)/T_{\text{ch}}(t)) \pm 1]^{-1}$$

$$\text{Chemical equilibration: } T_{\text{ch}}^{(f)}(t) = T_{\text{ch}}^{(s)}(t)$$

Kinetic equilibration before chemical equilibration



Once a temperature becomes stationary it takes the value of the equilibrium temperature.

Once chemical equilibration has been reached the chemical temperature equals the kinetic temperature and can be associated with the overall equilibrium temperature.

Comparison of chemical freeze out temperature with critical temperature of phase transition makes sense

A possible source of error :
temperature-dependent particle masses

Chiral order parameter σ depends on T

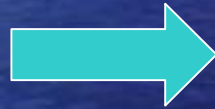
$$M_j(T) = h_j(T, \mu) \sigma(T, \mu)$$

$$\frac{\sigma(T_{\text{ch}}, \mu)}{T_{\text{ch}}} = \frac{\sigma(0, 0)}{T_{\text{obs}}}.$$

$$T_c = 176^{+5}_{-18} \text{ MeV}.$$

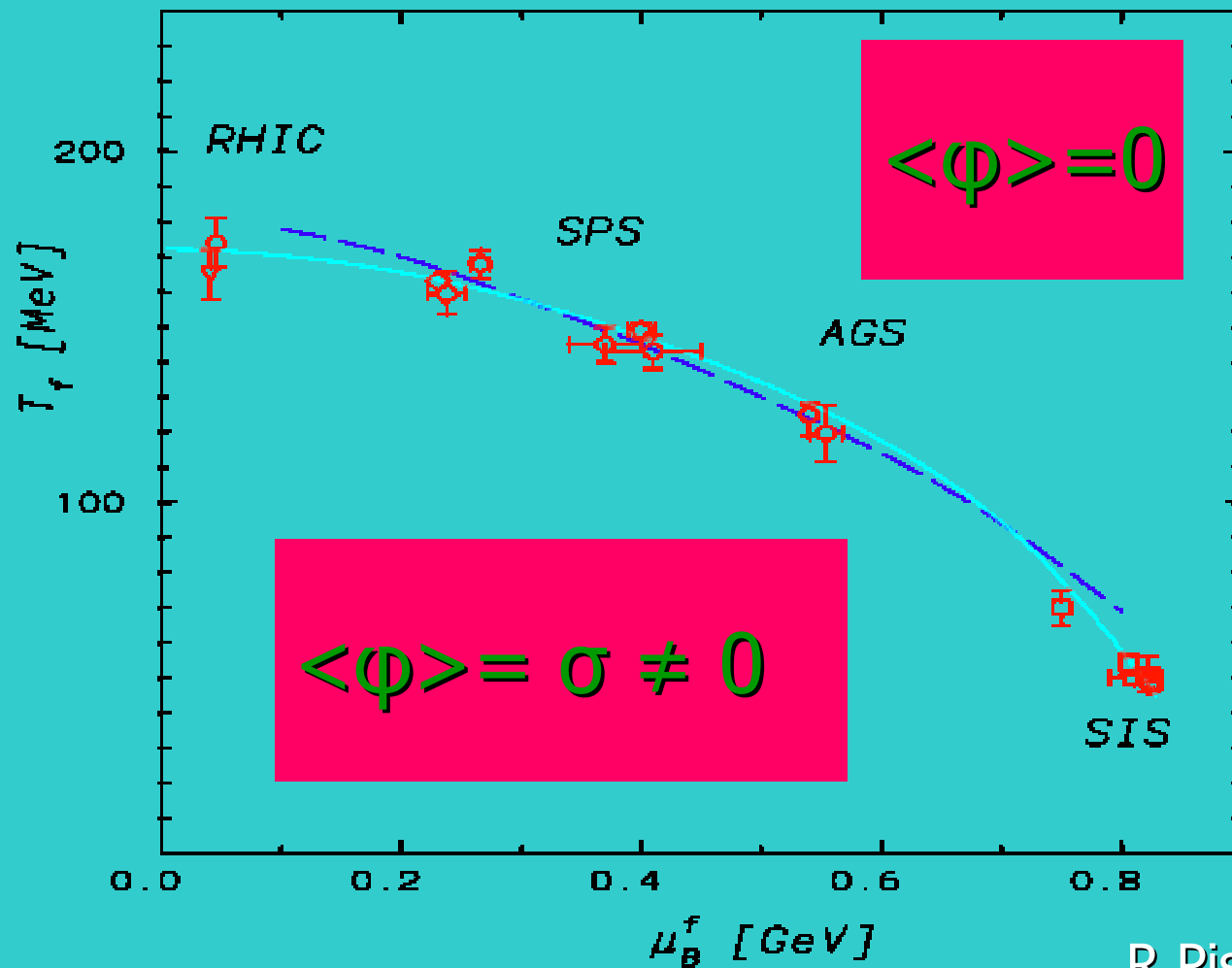
chemical
freeze out
measures
 T/m !

uncertainty in $m(T)$



uncertainty in critical temperature

Phase diagram



Chiral symmetry restoration at high temperature

Low T

SSB

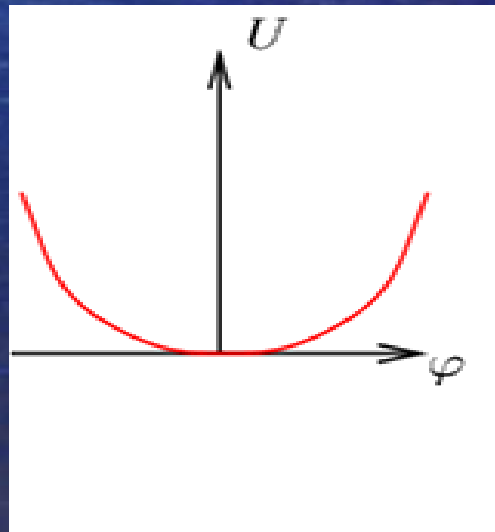
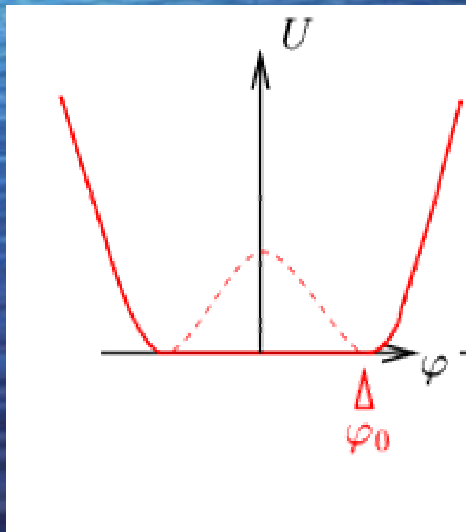
$$\langle \varphi \rangle = \varphi_0 \neq 0$$

High T

SYM

$$\langle \varphi \rangle = 0$$

at high T :
less order
more symmetry

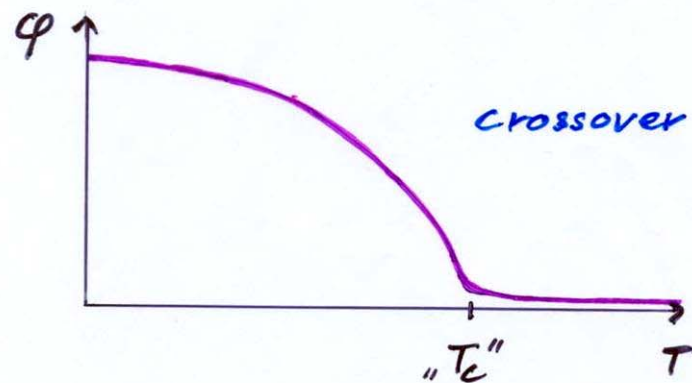
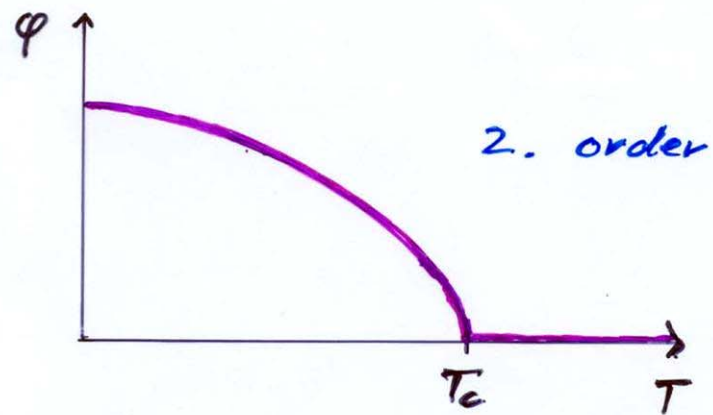
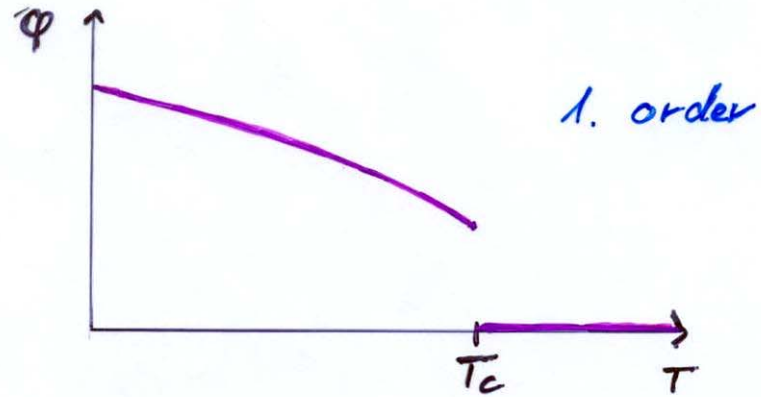


examples:
magnets, crystals

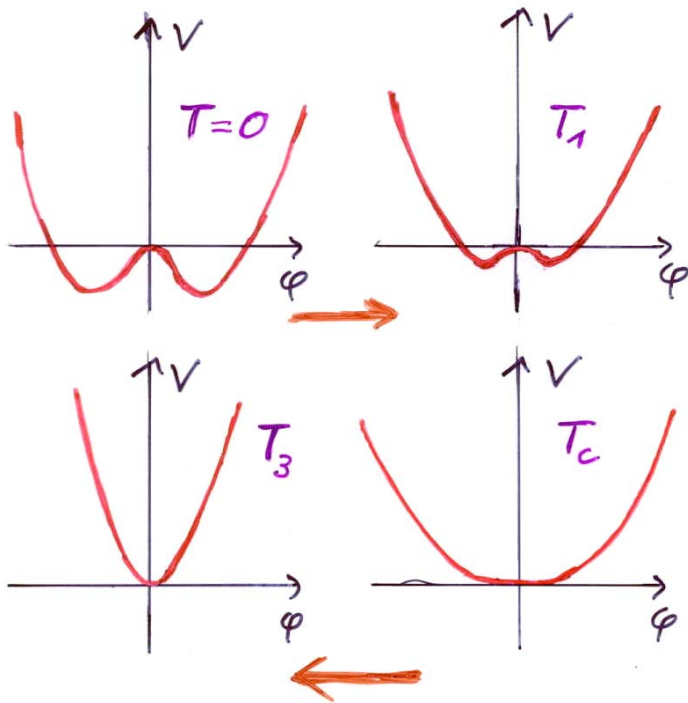


Order of the phase transition is
crucial ingredient for experiments
(heavy ion collisions)
and cosmological phase transition

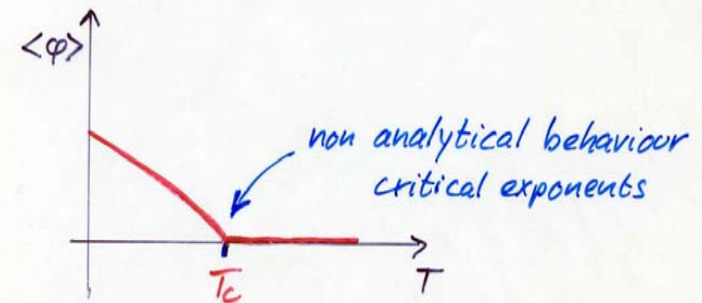
Order of the phase transition



Second order phase transition



$$V = -\mu_0^2 \varphi^\dagger \varphi + c T^2 \varphi^\dagger \varphi + \frac{\lambda}{2} (\varphi^\dagger \varphi)^2$$



second order phase transition

for T only somewhat below T_c :

the order parameter σ is expected to deviate substantially from its vacuum value

This seems to be disfavored by observation of chemical freeze out !

Ratios of particle masses and chemical freeze out

at chemical freeze out :

- ratios of hadron masses seem to be close to vacuum values
- nucleon and meson masses have different characteristic dependence on σ
- $m_{\text{nucleon}} \sim \sigma$, $m_{\pi} \sim \sigma^{-1/2}$
- $\Delta\sigma/\sigma < 0.1$ (conservative)

systematic uncertainty :

$$\Delta\sigma/\sigma = \Delta T_c/T_c$$

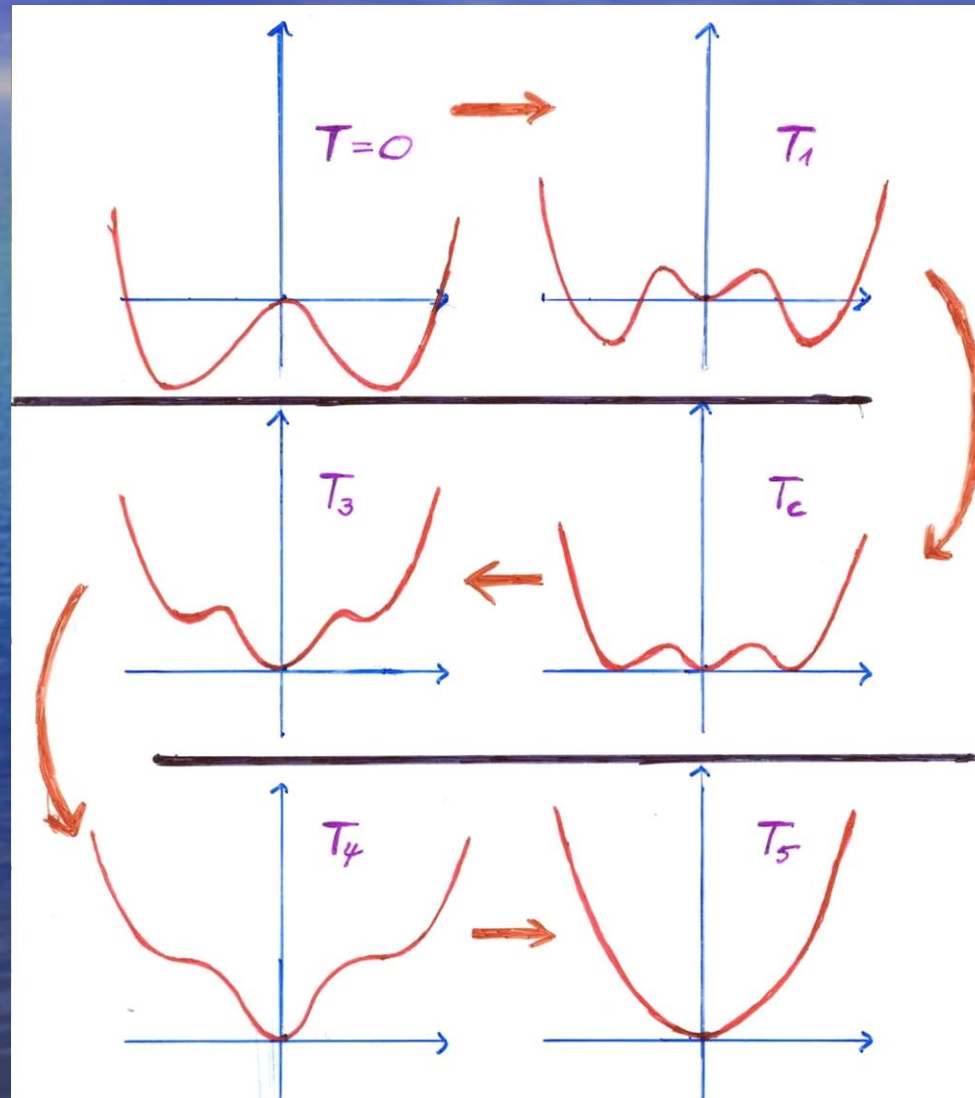
$\Delta\sigma$ is negative

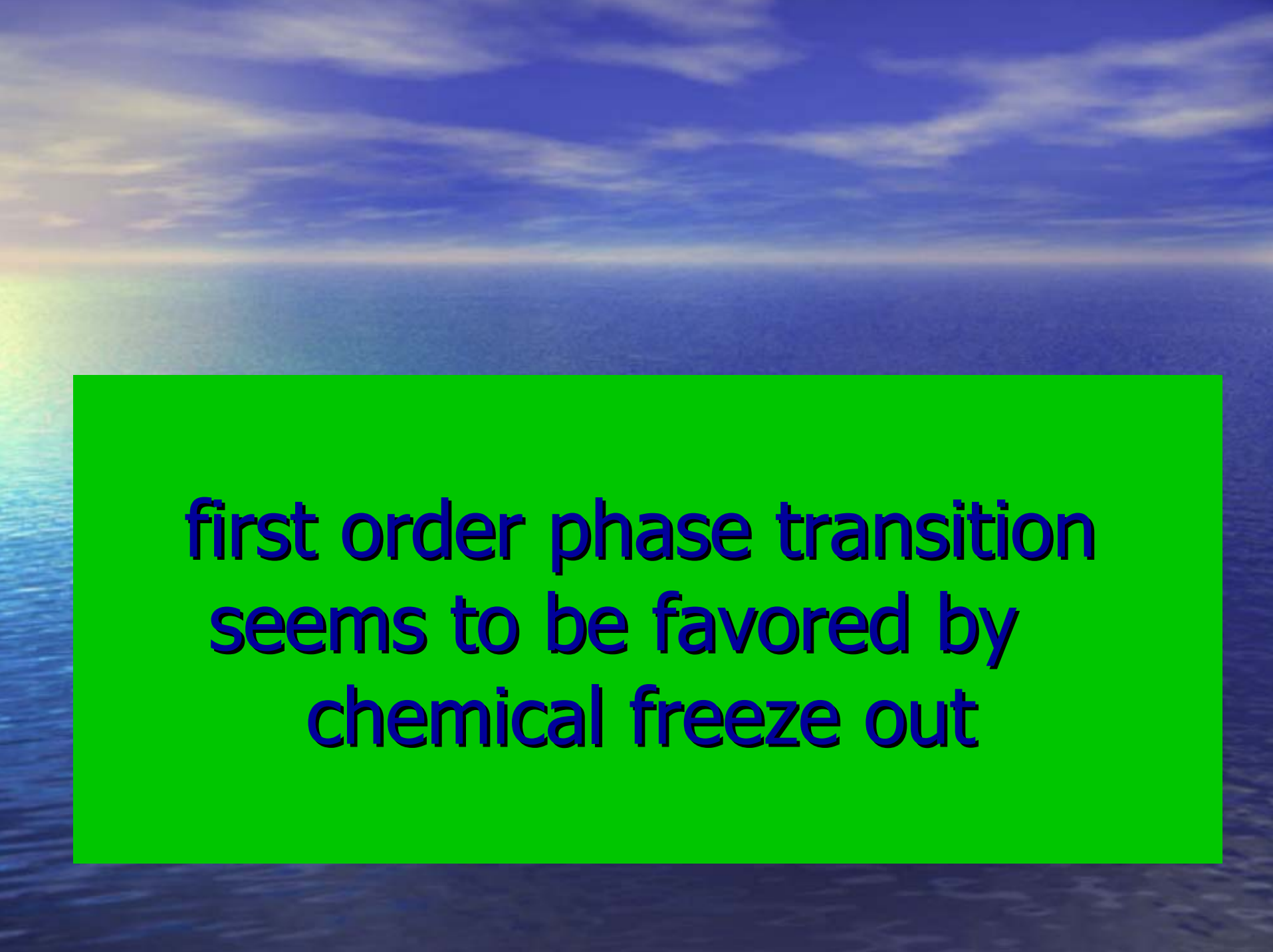
$$M_j(T) = h_j(T, \mu) \sigma(T, \mu)$$

$$\frac{\sigma(T_{\text{ch}}, \mu)}{T_{\text{ch}}} = \frac{\sigma(0, 0)}{T_{\text{obs}}}.$$

$$T_c = 176^{+5}_{-18} \text{ MeV}.$$

First order phase transition



The background of the slide is a photograph of a sunset or sunrise over a vast body of water. The sky is a deep blue with wispy white clouds, and the horizon line is visible in the middle. The water in the foreground is dark blue with gentle ripples. A bright green rectangular box is centered on the slide, containing text in a blue, bold, sans-serif font.

first order phase transition
seems to be favored by
chemical freeze out

Lattice results

e.g. Karsch,Laermann,Peikert

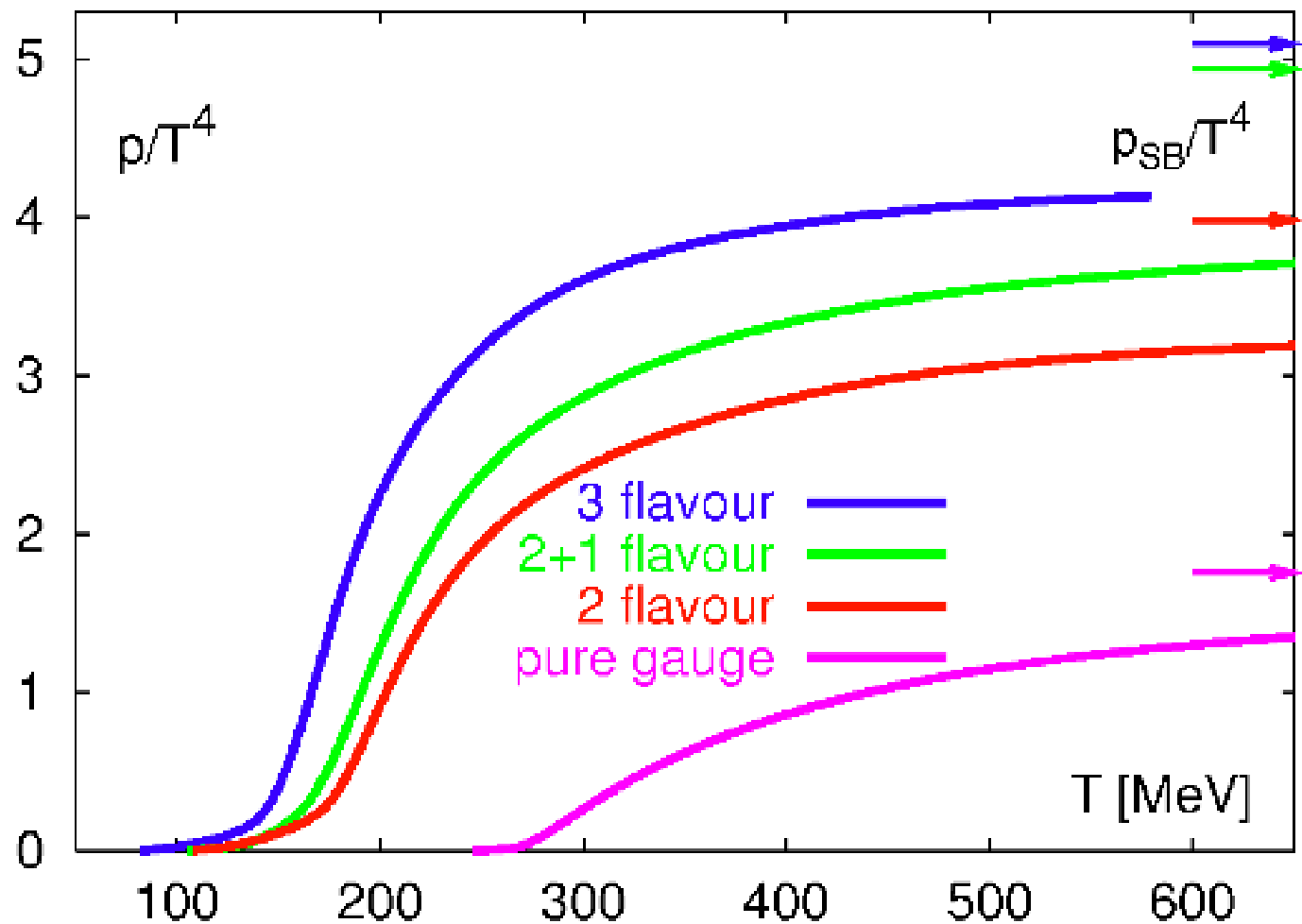
Critical temperature in chiral limit :

$$N_f = 3 : T_c = (154 \pm 8) \text{ MeV}$$

$$N_f = 2 : T_c = (173 \pm 8) \text{ MeV}$$

Chiral symmetry restoration and deconfinement at
same T_c

pressure



realistic QCD

- precise lattice results not yet available for first order transition vs. crossover
- also uncertainties in determination of critical temperature (chiral limit ...)
- extension to nonvanishing baryon number only for QCD with relatively heavy quarks

conclusion

- experimental determination of critical temperature may be more precise than lattice results
- error estimate becomes crucial



end