

The background of the slide is a deep-field astronomical image, likely from the Hubble Space Telescope. It shows a vast field of galaxies and stars against a black background. The galaxies are of various shapes and sizes, including spiral, elliptical, and irregular forms. Some are bright and clear, while others are faint and distant. The stars appear as small, bright points of light, some with prominent diffraction spikes. The overall scene conveys the immense scale and complexity of the universe.

**Expanding Universe or  
shrinking atoms ?**

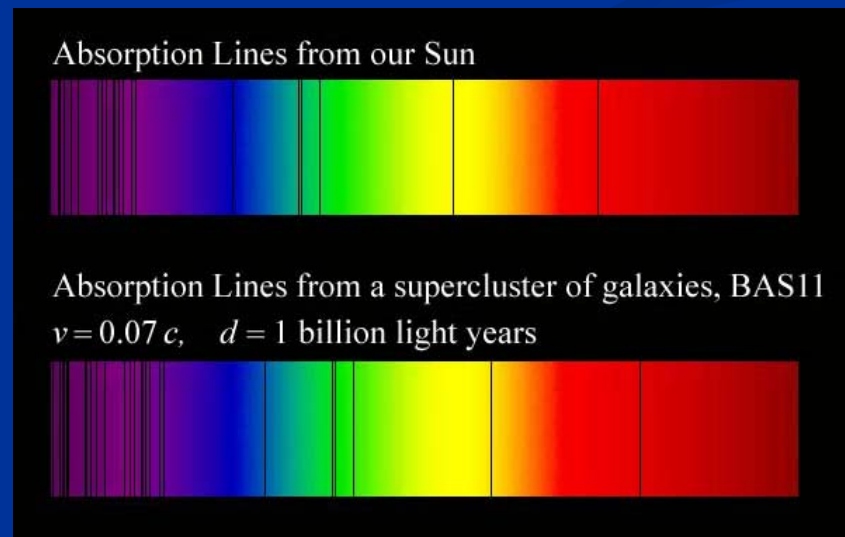
A deep-field astronomical image, likely from the Hubble Space Telescope, showing a vast field of galaxies and distant stars. The galaxies are of various shapes and sizes, including spiral, elliptical, and irregular forms, scattered across a black background. The text "Big bang or freeze ?" is overlaid in the center in a yellow, serif font.

**Big bang or freeze ?**

# Do we know that the Universe expands ?

instead of redshift due to expansion :

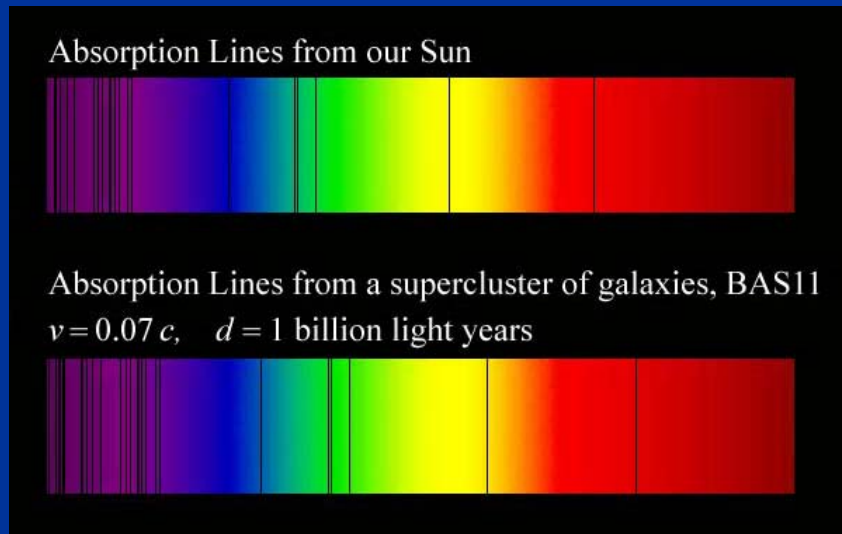
smaller frequencies have been emitted in the past,  
because electron mass was smaller !





# Why do we see redshift of photons emitted in the distant past ?

photons are more red because they have been **emitted** with longer wavelength



frequency  $\sim$  mass

wavelength  $\sim$   
atoms size



# What is increasing ?

Ratio of distance between galaxies  
over size of atoms !

atom size constant : expanding geometry

alternative : shrinking size of atoms

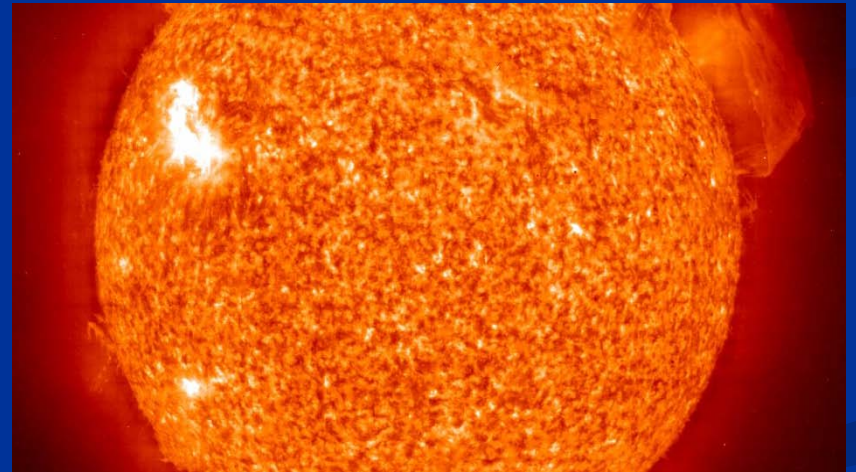
# How can particle masses change with time ?

- Particle masses are proportional to scalar field  $\chi$  .
- Scalar field varies with time.
- Ratios of particle masses are independent of  $\chi$  and therefore remain constant.
- Compatibility with observational bounds on time dependence of particle mass ratios.
- Dimensionless couplings are independent of  $\chi$  .

Do we know that the temperature was higher in the early Universe than now ?

Cosmic microwave radiation , nucleosynthesis

instead of  
higher temperature :  
smaller particle masses

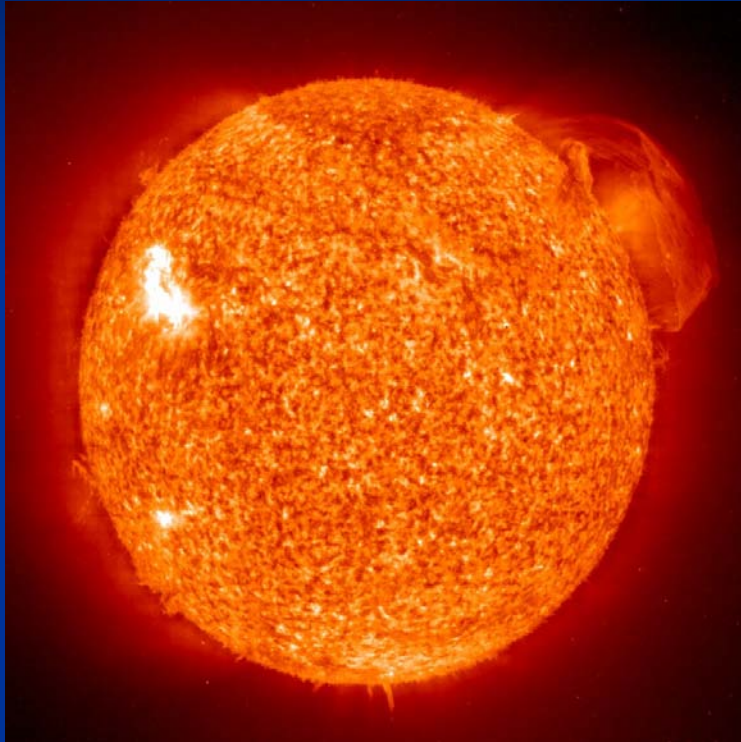




# Hot plasma ?

- Temperature in radiation dominated Universe :  
 $T \sim \chi^{1/2}$  **smaller** than today
- Ratio temperature / particle mass :  
 $T / m_p \sim \chi^{-1/2}$  **larger** than today
- $T/m_p$  counts ! This ratio decreases with time.
- Nucleosynthesis , CMB emission as in standard cosmology !

# Big bang or freeze ?



freeze picture :  
only rods for measurements  
are set differently !



*Big bang is not wrong,*

*but alternative pictures exist !*



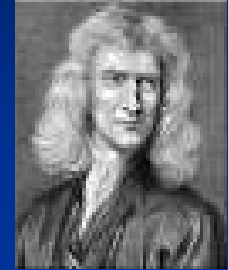
# Field relativity :

## different pictures of cosmology

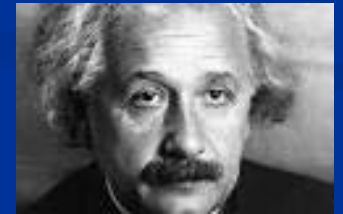
- same physical content can be described by different pictures
- related by field – redefinitions ,  
e.g. Weyl scaling , conformal scaling of metric
- which picture is usefull ?

# Relativity of geometry

- Euclid ... Newton : space and time are absolute



- Special relativity : space and time depend on observer
- General relativity : spacetime is influenced by matter ( including radiation )  
geometry is independent of coordinates geometry is observable



- Field relativity : geometry is relative

*Spacetime is a description  
of correlations between “matter”.*

*Different pictures exist.*



# Why should you care about the freeze picture of the Universe ?

*Some aspects are understood easier :*

- Beginning of Universe
- Role of scale symmetry
- Range of impact of quantum gravity

# conclusions

- Big bang singularity is artefact of inappropriate choice of field variables – no physical singularity
- Quantum gravity may be observable in dynamics of present Universe

# variable gravity

*“Newton’s constant is not constant”*



# Variable Gravity

$$\Gamma = \int_x \sqrt{g} \left\{ -\frac{1}{2} \chi^2 R + \mu^2 \chi^2 + \frac{1}{2} (B(\chi/\mu) - 6) \partial^\mu \chi \partial_\mu \chi \right\}$$

quantum effective action,  
variation yields field equations

Einstein gravity :  $\Gamma = \int_x \sqrt{g} \left\{ -\frac{1}{2} M^2 R \right\}$

# Variable Gravity

- Scalar field coupled to gravity
- Effective Planck mass depends on scalar field
- Simple quadratic scalar potential involves intrinsic mass  $\mu$
- Nucleon and electron mass proportional to dynamical Planck mass
- Neutrino mass has different dependence on scalar field

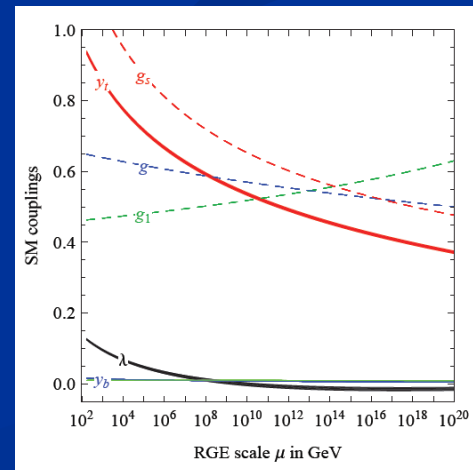
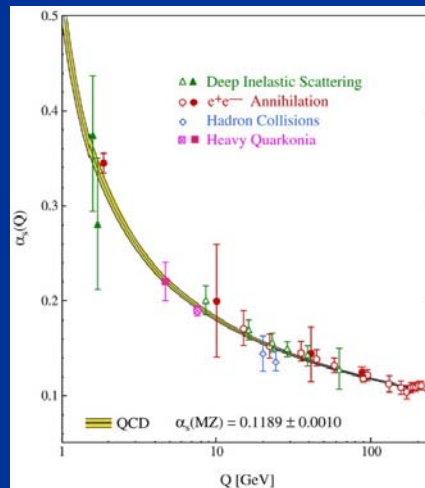
$$\Gamma = \int_x \sqrt{g} \left\{ -\frac{1}{2} \chi^2 R + \mu^2 \chi^2 + \frac{1}{2} (B(\chi/\mu) - 6) \partial^\mu \chi \partial_\mu \chi \right\}$$

# Running coupling

- B varies if intrinsic scale  $\mu$  changes

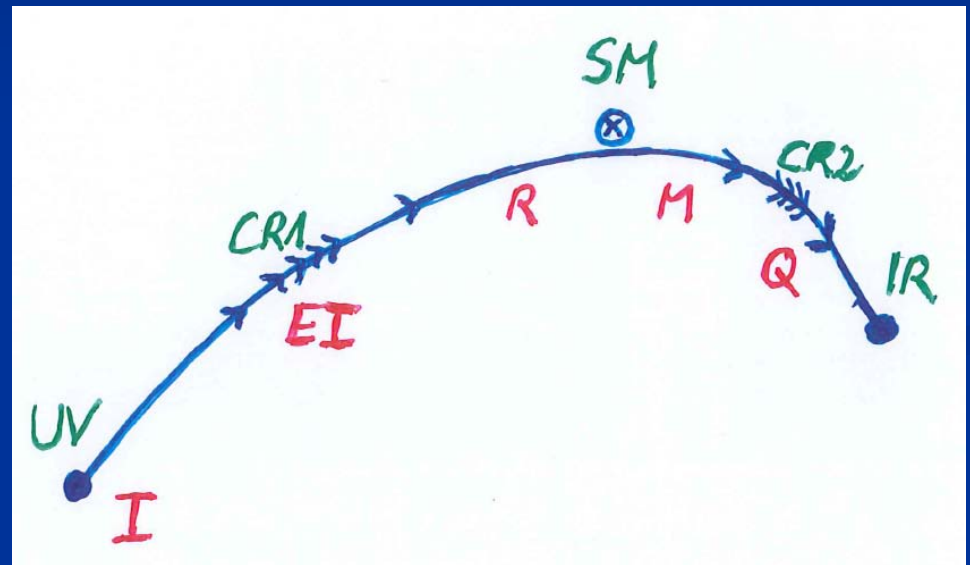
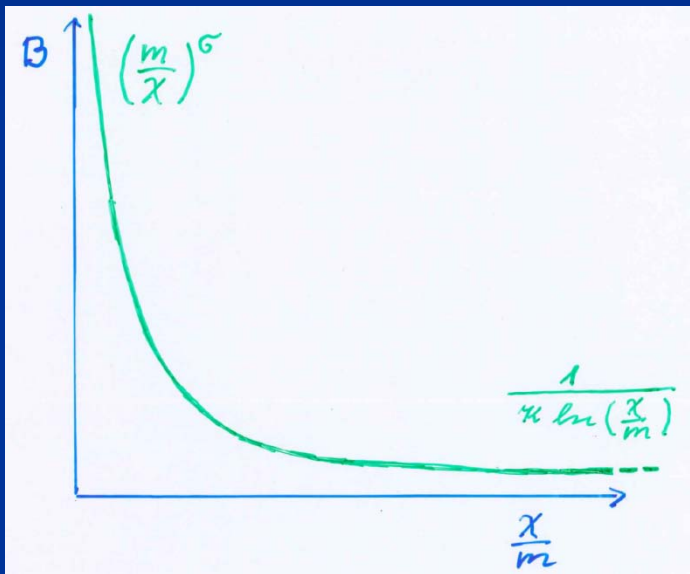
$$\Gamma = \int_x \sqrt{g} \left\{ -\frac{1}{2} \chi^2 R + \mu^2 \chi^2 + \frac{1}{2} (B(\chi/\mu) - 6) \partial^\mu \chi \partial_\mu \chi \right\}$$

- similar to QCD or standard model

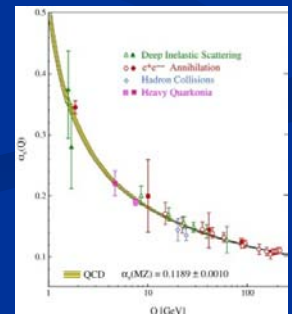


# Kinetic B :

## Crossover between two fixed points



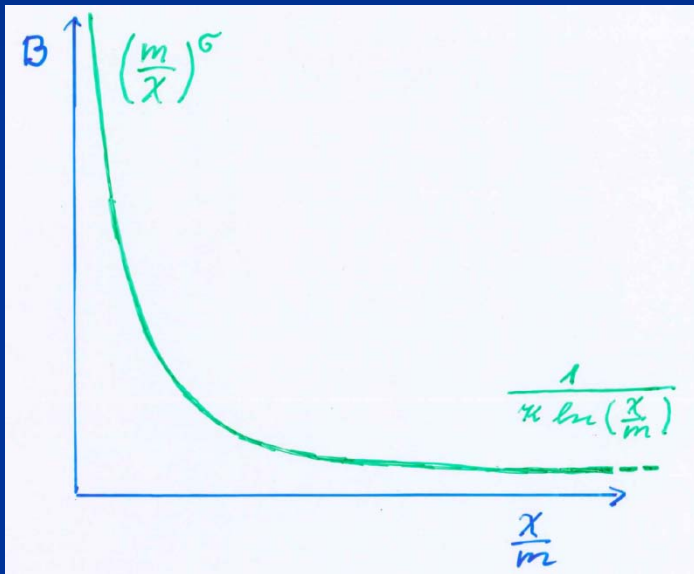
$$\mu \frac{\partial B}{\partial \mu} = \frac{\kappa \sigma B^2}{\sigma + \kappa B}$$





# Kinetic B :

## Crossover between two fixed points



running  
coupling obeys  
flow equation

$$\mu \frac{\partial B}{\partial \mu} = \frac{\kappa \sigma B^2}{\sigma + \kappa B}$$

$$B^{-1} - \frac{\kappa}{\sigma} \ln B = \kappa \left[ \ln \left( \frac{\chi}{\mu} \right) - c_t \right] = \kappa \ln \left( \frac{\chi}{m} \right)$$

$m$  : scale of crossover

can be exponentially larger than intrinsic scale  $\mu$

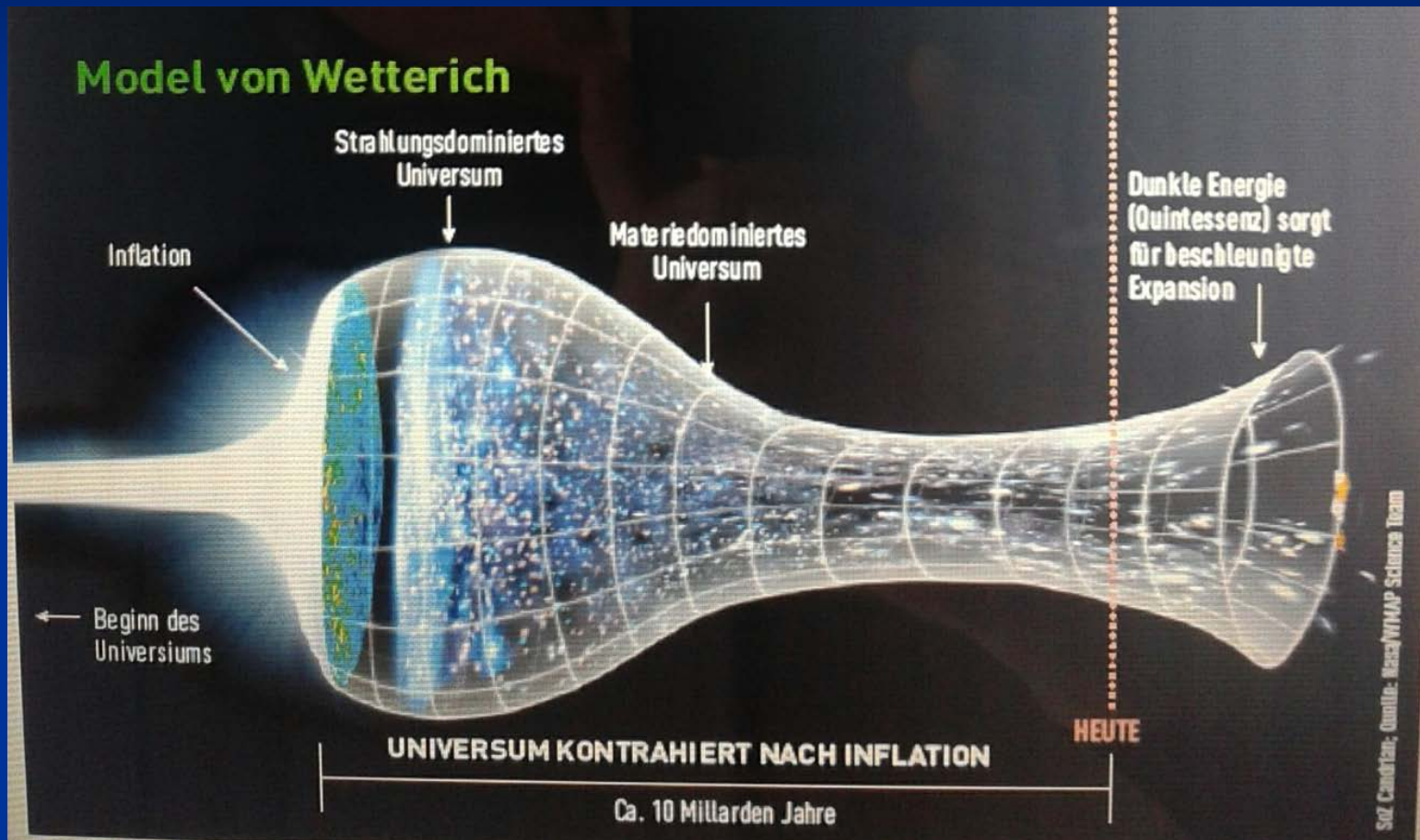
# Four-parameter model

- model has four dimensionless parameters
- three in kinetic :
  - $\sigma \sim 2.5$
  - $\kappa \sim 0.5$
  - $c_t \sim 14$  ( or  $m/\mu$  )
- one parameter for growth rate of neutrino mass over electron mass :  $\gamma \sim 8$
- + standard model particles and dark matter : sufficient for realistic cosmology from inflation to dark energy
- no more free parameters than  $\Lambda$ CDM

# Cosmological solution

- scalar field  $\chi$  vanishes in the infinite past
- scalar field  $\chi$  diverges in the infinite future

# Strange evolution of Universe



Sonntagszeitung Zürich , Laukenmann

# Model is compatible with present observations

Together with variation of neutrino mass over  
electron mass in present cosmological epoch :  
model is compatible with all present  
observations, including inflation and dark energy

$$\Gamma = \int_x \sqrt{g} \left\{ -\frac{1}{2} \chi^2 R + \mu^2 \chi^2 + \frac{1}{2} (B(\chi/\mu) - 6) \partial^\mu \chi \partial_\mu \chi \right\}$$

$$B^{-1} - \frac{\kappa}{\sigma} \ln B = \kappa \left[ \ln \left( \frac{\chi}{\mu} \right) - c_t \right] = \kappa \ln \left( \frac{\chi}{m} \right)$$

# Einstein frame

- “Weyl scaling” maps variable gravity model to Universe with fixed masses and standard expansion history.
- Exact equivalence of different frames !
- Standard gravity coupled to scalar field.
- Only neutrino masses are growing.



# Einstein frame

Weyl scaling :

$$g'_{\mu\nu} = \frac{\chi^2}{M^2} g_{\mu\nu} , \quad \varphi = \frac{2M}{\alpha} \ln \left( \frac{\chi}{\mu} \right)$$

effective action in Einstein frame :

$$\Gamma = \int_x \sqrt{g'} \left\{ -\frac{1}{2} M^2 R' + V'(\varphi) + \frac{1}{2} k^2(\varphi) \partial^\mu \varphi \partial_\mu \varphi \right\}$$

$$V'(\varphi) = M^4 \exp \left( -\frac{\alpha \varphi}{M} \right)$$

$$k^2 = \frac{\alpha^2 B}{4}$$

# Field relativity

Weyl scaling :

$$g'_{\mu\nu} = \frac{\chi^2}{M^2} g_{\mu\nu}$$

changes geometry,  
not a coordinate transformation

infinite past

# Infinite past : slow inflation

$\sigma = 2$  : field equations

$$\ddot{\chi} + \left( 3H + \frac{1}{2} \frac{\dot{\chi}}{\chi} \right) \dot{\chi} = \frac{2\mu^2 \chi^2}{m}$$

$$H = \sqrt{\frac{\mu^2}{3} + \frac{m\dot{\chi}^2}{6\chi^3}} - \frac{\dot{\chi}}{\chi}$$

approximative  
solution

$$H = \frac{\mu}{\sqrt{3}}, \quad \chi = \frac{3^{\frac{1}{4}} m}{2\sqrt{\mu}} (t_c - t)^{-\frac{1}{2}}$$

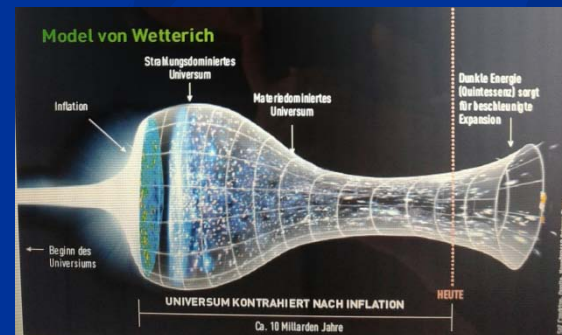
particles become massless in infinite past !

# Eternal Universe

Asymptotic solution in freeze frame :

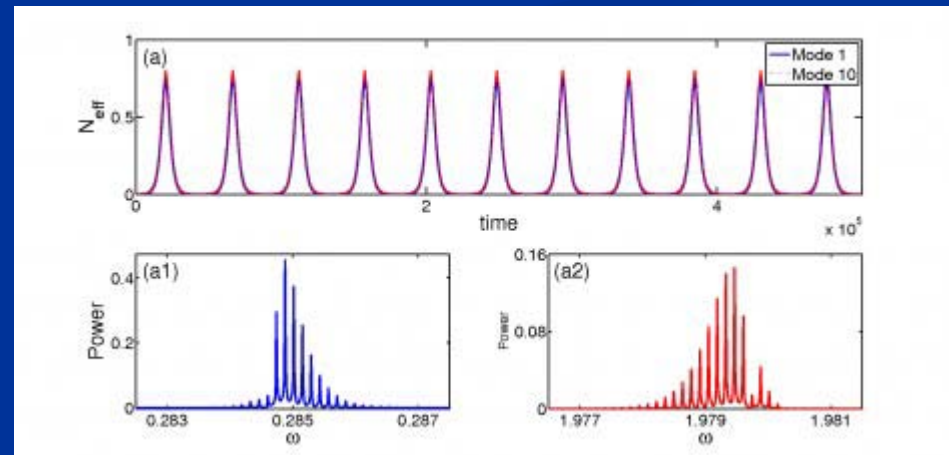
$$H = \frac{\mu}{\sqrt{3}}, \quad \chi = \frac{3^{\frac{1}{4}} m}{2\sqrt{\mu}} (t_c - t)^{-\frac{1}{2}}$$

- solution valid back to the infinite past in physical time
- no singularity
- physical time to infinite past is infinite



# Physical time

count oscillations ....





# Physical time

field equation for scalar field mode

$$(\partial_\eta^2 + 2Ha\partial_\eta + k^2 + a^2m^2)\varphi_k = 0$$

$$\varphi_k = \frac{\tilde{\varphi}_k}{a} \quad \left\{ \partial_\eta^2 + k^2 + a^2 \left( m^2 - \frac{R}{6} \right) \right\} \tilde{\varphi}_k = 0$$

determine **physical time** by counting number of oscillations

$$\tilde{t}_p = n_k$$

$$n_k = \frac{k\eta}{\pi}$$

( m=0 )

# Physical time

- counting : discrete
- invariant under field transformations
- same in all frames

*Big bang singularity  
in Einstein frame is  
field singularity !*

$$g'_{\mu\nu} = \frac{\chi^2}{M^2} g_{\mu\nu} , \quad \varphi = \frac{2M}{\alpha} \ln \left( \frac{\chi}{\mu} \right)$$

choice of frame with constant particle masses is not well suited if physical masses go to zero !

# Inflation

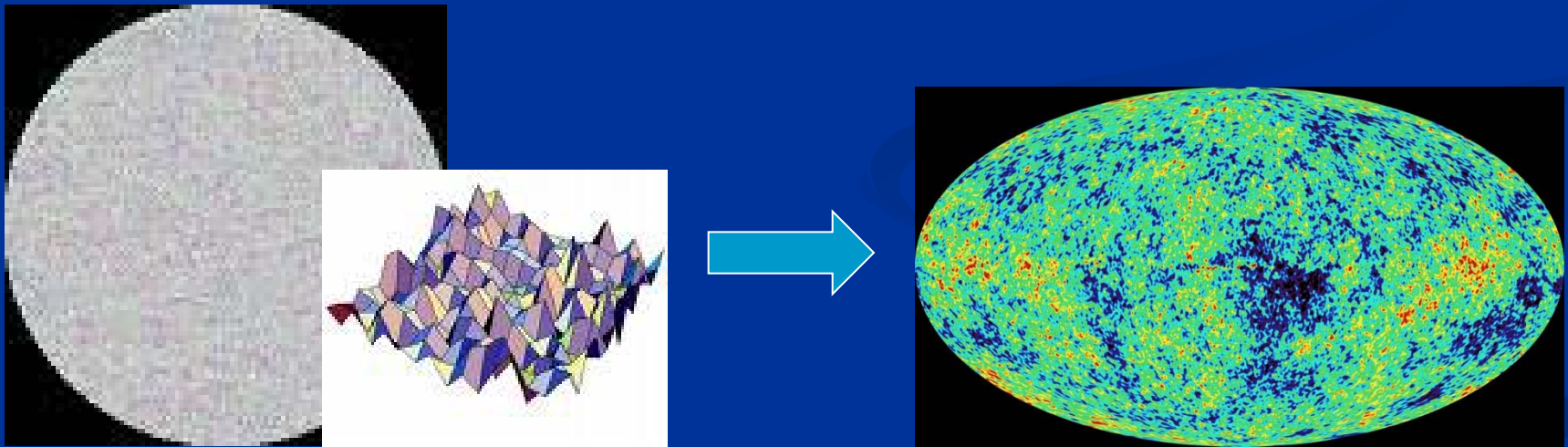
solution for small  $\chi$  : inflationary epoch

kinetial characterized by  
anomalous dimension  $\sigma$

$$B = b \left( \frac{\mu}{\chi} \right)^{\sigma} = \left( \frac{m}{\chi} \right)^{\sigma}$$

# Primordial fluctuations

- inflaton field :  $\chi$
- primordial fluctuations of inflaton become observable in cosmic microwave background



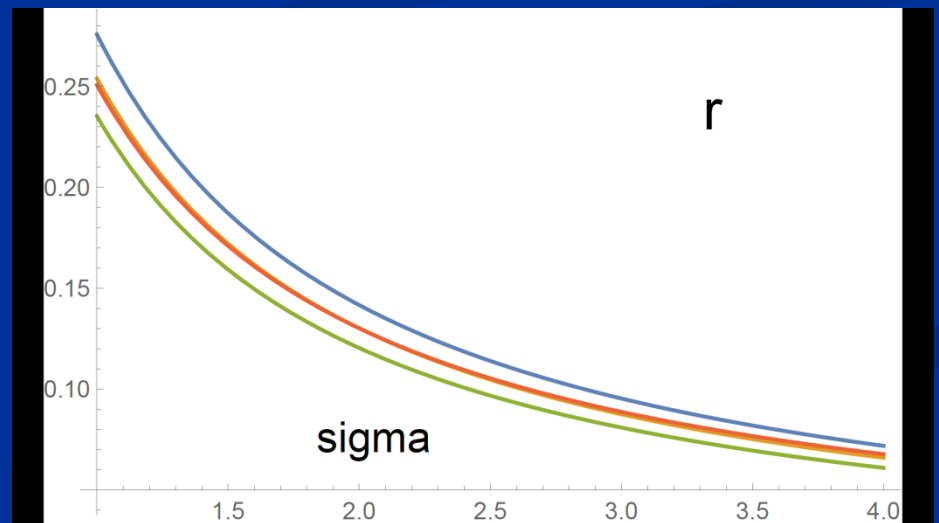
# Anomalous dimension determines spectrum of primordial fluctuations

$$r = \frac{0.26}{\sigma}$$

$$n = 1 - \frac{0.065}{\sigma} \cdot \left(1 + \frac{\sigma - 2}{4}\right)$$

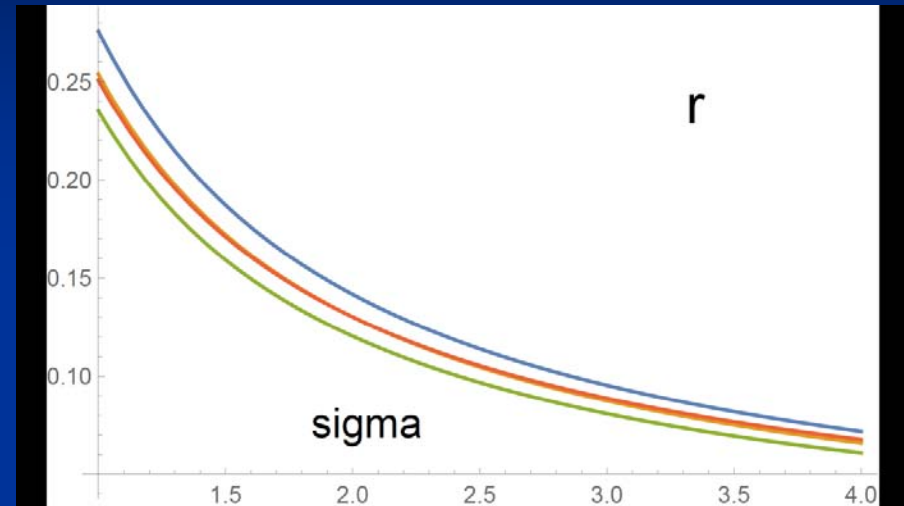
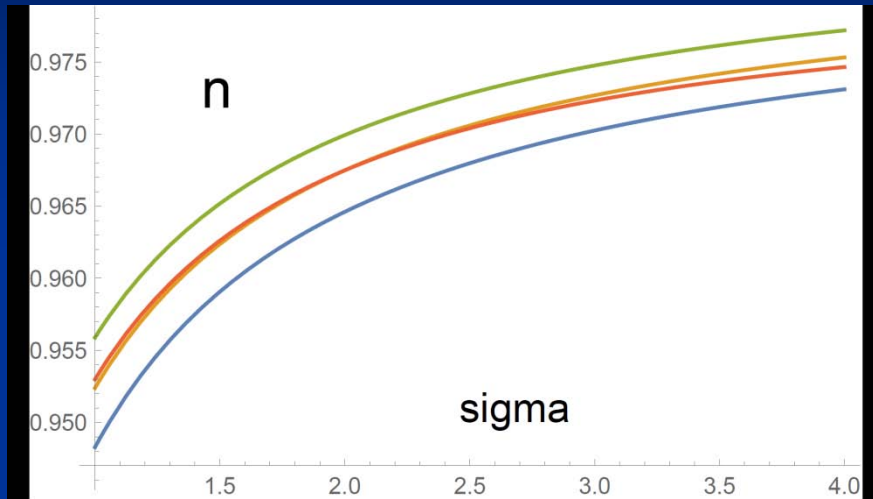
spectral index  $n$

tensor amplitude  $r$





# relation between n and r



$$r = 8.19 ( 1 - n ) - 0.1365$$

# Amplitude of density fluctuations

small because of logarithmic running  
near UV fixed point !

$$\mathcal{A} = \frac{(N+3)^3}{4} e^{-2c_t}$$

$$c_t = \ln \left( \frac{m}{\mu} \right) = 14.1 \quad \sigma=1$$

$$\frac{m}{\mu} = \frac{(N+3)^{\frac{3}{2}}}{2\sqrt{\mathcal{A}}} = 1.32 \cdot 10^6 \left( \frac{N}{60} \right)^{\frac{3}{2}}$$

$$B^{-1} - \frac{\kappa}{\sigma} \ln B = \kappa \left[ \ln \left( \frac{\chi}{\mu} \right) - c_t \right] = \kappa \ln \left( \frac{\chi}{m} \right)$$

N : number of e – foldings at horizon crossing

no small parameter for  
dark energy

# Four-parameter model

- model has four dimensionless parameters
- three in kinetic :
  - $\sigma \sim 2.5$
  - $\kappa \sim 0.5$
  - $c_t \sim 14$  ( or  $m/\mu$  )
- one parameter for growth rate of neutrino mass over electron mass :  $\gamma \sim 8$
- + standard model particles and dark matter : sufficient for realistic cosmology from inflation to dark energy
- no more free parameters than  $\Lambda$ CDM

# No tiny dimensionless parameters ( except gauge hierarchy )

- one mass scale  $\mu = 2 \cdot 10^{-33} \text{ eV}$
- one time scale  $\mu^{-1} = 10^{10} \text{ yr}$
- Planck mass does not appear as parameter
- Planck mass grows large dynamically

# Slow Universe

Asymptotic solution in  
freeze frame :

$$H = \frac{\mu}{\sqrt{3}} , \quad \chi = \frac{3^{\frac{1}{4}} m}{2\sqrt{\mu}} (t_c - t)^{-\frac{1}{2}}$$

$$\mu = 2 \cdot 10^{-33} \text{ eV}$$

Expansion or shrinking always slow ,  
characteristic time scale of the order of the age of the  
Universe :  $t_{\text{ch}} \sim \mu^{-1} \sim 10 \text{ billion years} !$

Hubble parameter of the order of **present** Hubble  
parameter for all times , including inflation and big bang !  
Slow increase of particle masses !



# asymptotically vanishing cosmological „constant“

- What matters : Ratio of potential divided by fourth power of Planck mass

$$\frac{V}{\chi^4} = \frac{\mu^2 \chi^2}{\chi^4} = \frac{\mu^2}{\chi^2}$$

- vanishes for  $\chi \rightarrow \infty$  !

# Einstein frame

Weyl scaling :

$$g'_{\mu\nu} = \frac{\chi^2}{M^2} g_{\mu\nu} , \quad \varphi = \frac{2M}{\alpha} \ln \left( \frac{\chi}{\mu} \right)$$

effective action in Einstein frame :

$$\Gamma = \int_x \sqrt{g'} \left\{ -\frac{1}{2} M^2 R' + V'(\varphi) + \frac{1}{2} k^2(\varphi) \partial^\mu \varphi \partial_\mu \varphi \right\}$$

$$V'(\varphi) = M^4 \exp \left( -\frac{\alpha \varphi}{M} \right)$$

$$k^2 = \frac{\alpha^2 B}{4}$$

# small dimensionless number ?

- needs two intrinsic mass scales
- standard approach :  $V$  and  $M$  ( cosmological constant and Planck mass )
- variable gravity : Planck mass moving to infinity , with fixed  $V$  → ratio vanishes asymptotically !

# Quintessence

Dynamical dark energy ,  
generated by scalar field (cosmon )

C.Wetterich,Nucl.Phys.B302(1988)668, 24.9.87

P.J.E.Peebles,B.Ratra,ApJ.Lett.325(1988)L17, 20.10.87

**Prediction :**

**homogeneous dark energy  
influences recent cosmology**

**- of same order as dark matter -**

Original models do not fit the present observations  
.... modifications  
( different growth of neutrino mass )

# Cosmon inflation

Unified picture of inflation and  
dynamical dark energy

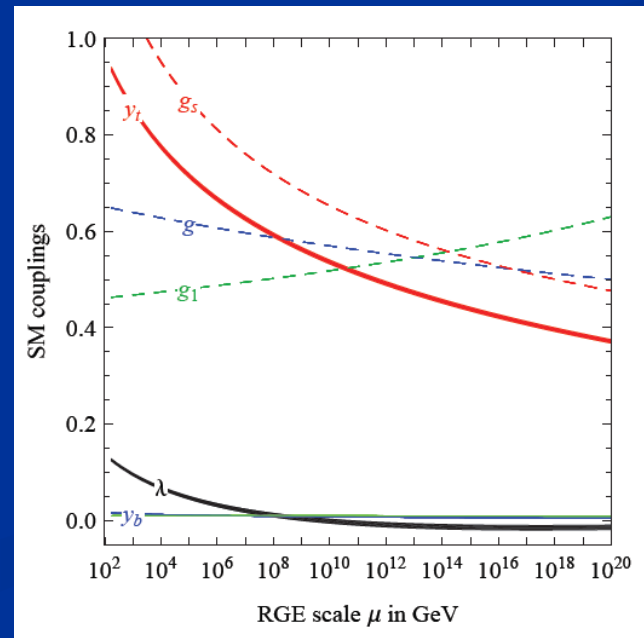
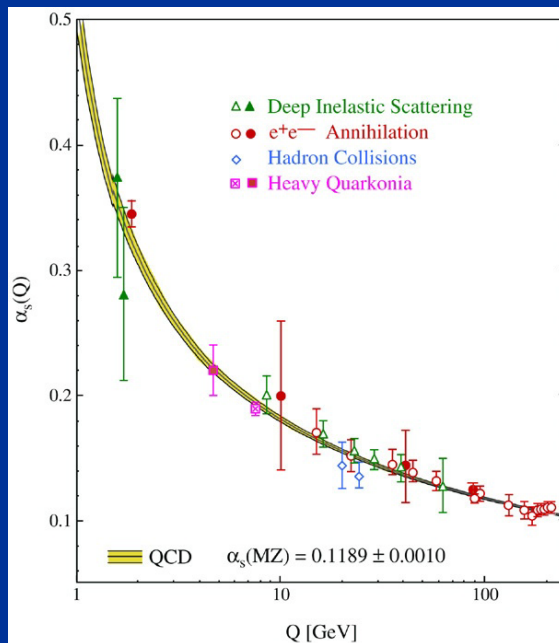
Cosmon and inflaton are the same  
scalar field

quantum gravity with  
scalar field –  
the role of scale symmetry



# fluctuations induce running couplings

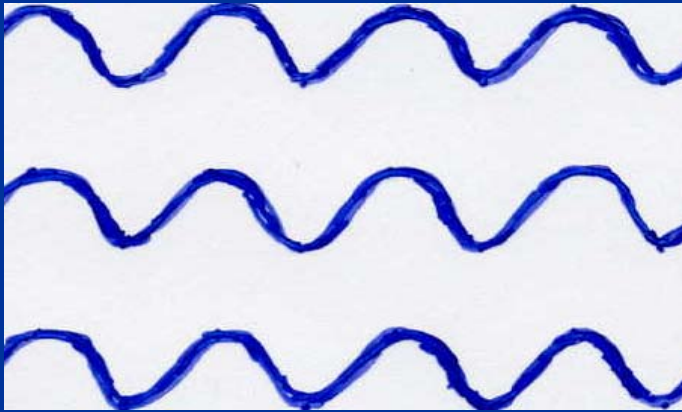
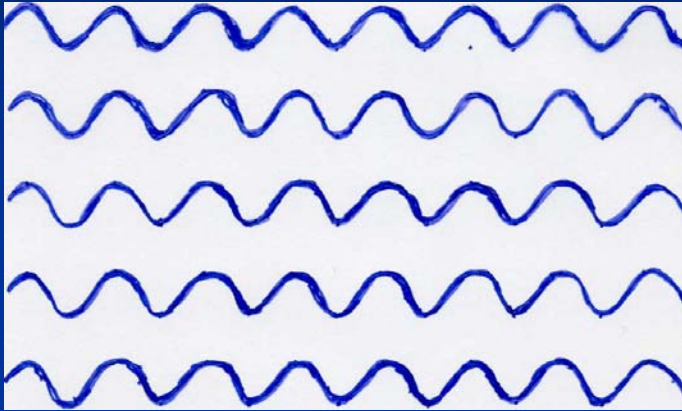
- violation of scale symmetry
- well known in QCD or standard model



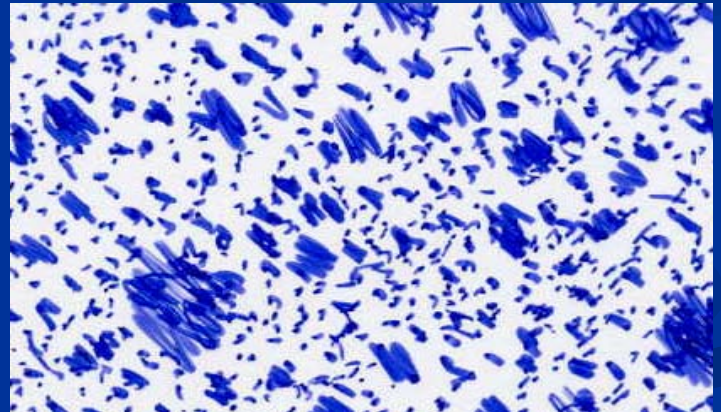
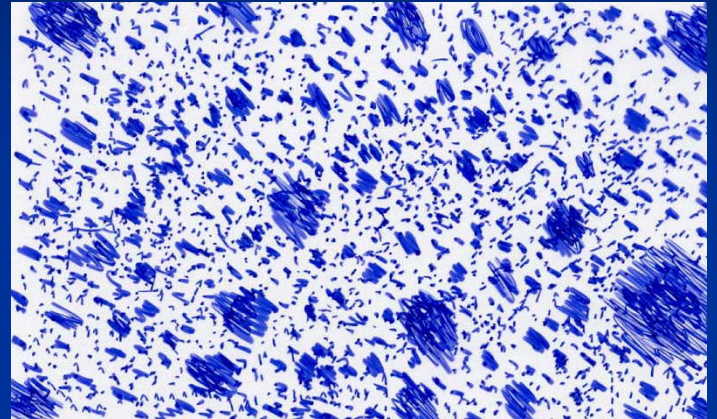
# Quantum scale symmetry

- quantum fluctuations violate scale symmetry
- running dimensionless couplings
- at fixed points , scale symmetry is exact !

# Scale symmetry

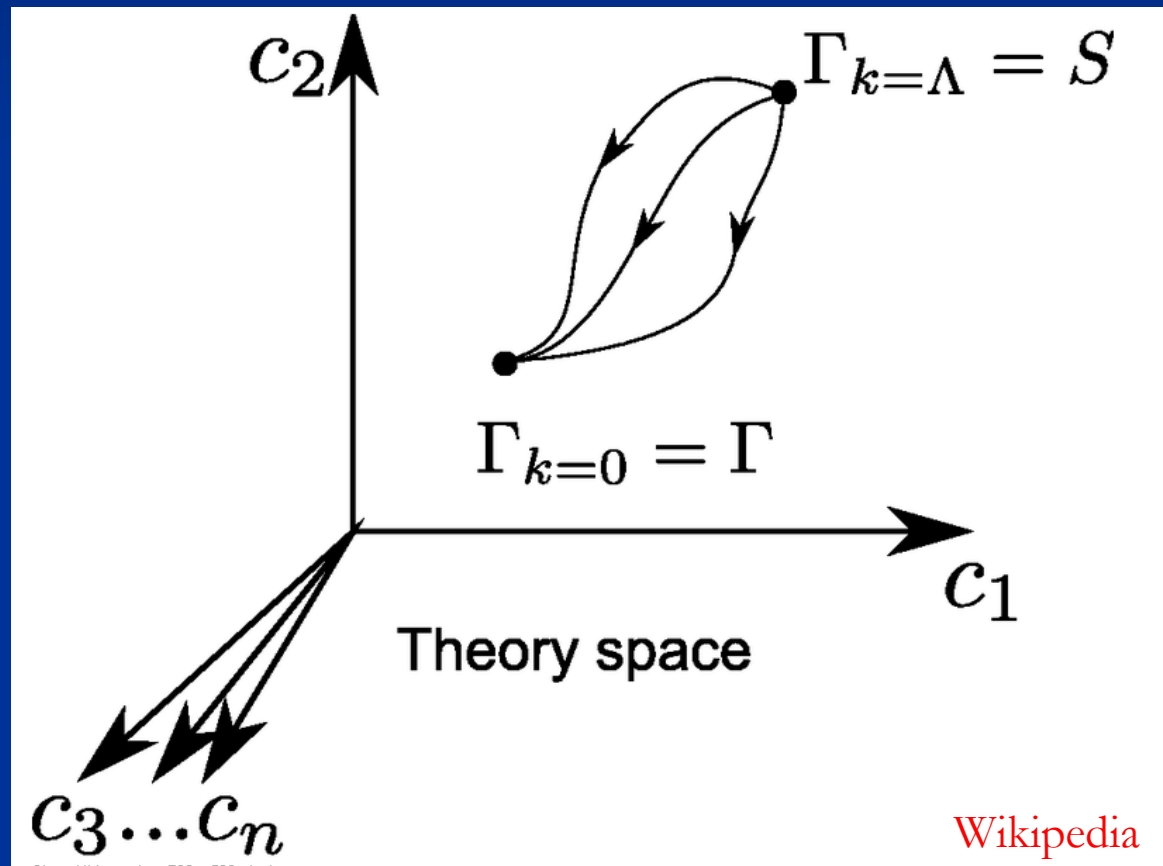


no scale symmetry

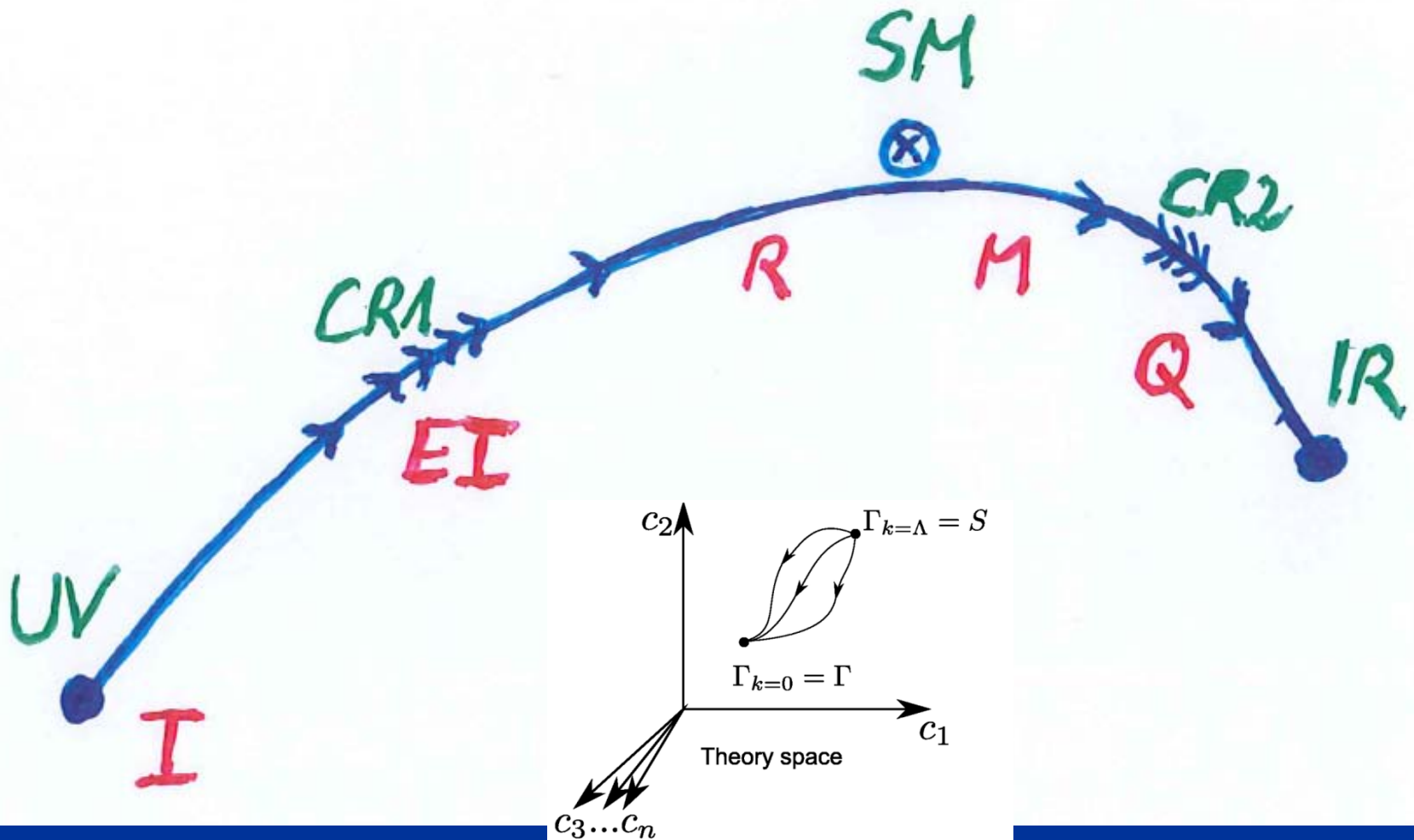


scale symmetry

# functional renormalization : flowing action

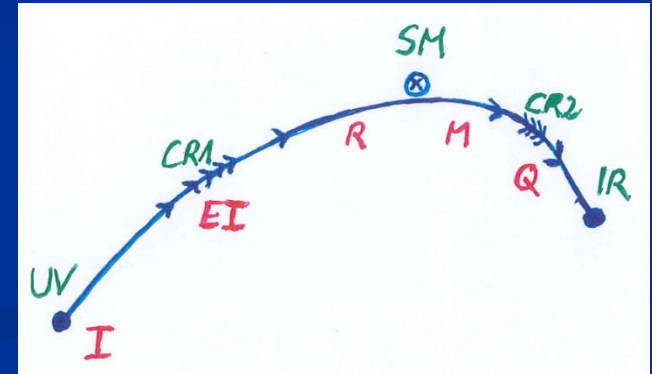


# Crossover in quantum gravity



# Origin of mass

- UV fixed point : scale symmetry unbroken  
all particles are massless
- IR fixed point :  
scale symmetry spontaneously broken,  
massive particles , massless dilaton
- crossover : explicit mass scale  $\mu$  important
- approximate SM fixed point : approximate scale symmetry  
spontaneously broken, massive particles , almost massless  
cosmon, tiny cosmon potential





# Spontaneous breaking of scale symmetry

- expectation value of scalar field breaks scale symmetry spontaneously
- massive particles are compatible with scale symmetry
- in presence of massive particles : sign of exact scale symmetry is exactly **massless Goldstone boson** – the dilaton



# Approximate scale symmetry near fixed points

- UV : approximate scale invariance of primordial fluctuation spectrum from inflation
- IR : cosmon is pseudo Goldstone boson of spontaneously broken scale symmetry, tiny mass, responsible for dynamical Dark Energy

# Asymptotic safety

if UV fixed point exists :

*quantum gravity is  
non-perturbatively renormalizable !*

S. Weinberg , M. Reuter

# a prediction...

## Asymptotic safety of gravity and the Higgs boson mass

Mikhail Shaposhnikov

*Institut de Théorie des Phénomènes Physiques, École Polytechnique Fédérale de Lausanne, CH-1015 Lausanne, Switzerland*

Christof Wetterich

*Institut für Theoretische Physik, Universität Heidelberg, Philosophenweg 16, D-69120 Heidelberg, Germany*

12 January 2010

### Abstract

There are indications that gravity is asymptotically safe. The Standard Model (SM) plus gravity could be valid up to arbitrarily high energies. Supposing that this is indeed the case and assuming that there are no intermediate energy scales between the Fermi and Planck scales we address the question of whether the mass of the Higgs boson  $m_H$  can be predicted. For a positive gravity induced anomalous dimension  $A_\lambda > 0$  the running of the quartic scalar self interaction  $\lambda$  at scales beyond the Planck mass is determined by a fixed point at zero. This results in  $m_H = m_{\min} = 126$  GeV, with only a few GeV uncertainty. This prediction is independent of the details of the short distance running and holds for a wide class of extensions of the SM as well.

s in  $m_H = m_{\min} = 126$  GeV, with o

# IR fixed point in quantum gravity

## Dilaton Quantum Gravity

T. Henz, J. M. Pawłowski, A. Rodigast, and C. Wetterich

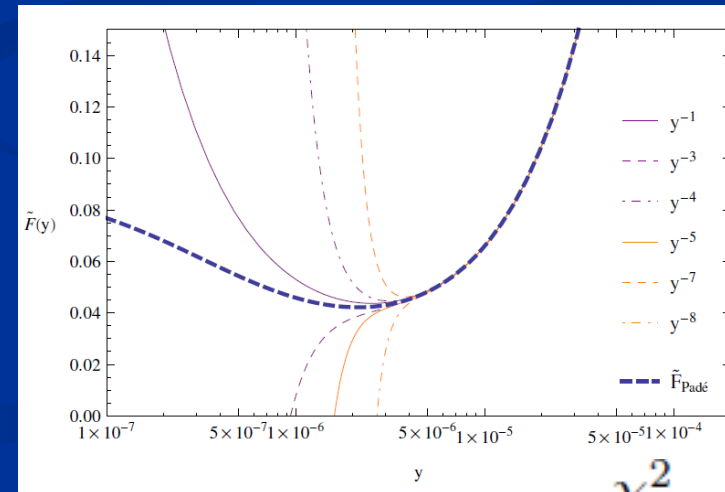
First positive indication from  
functional renormalization  
flow with truncation :

$$\Gamma_k = \int d^4x \sqrt{g} \left( V_k(\chi^2) - \frac{1}{2} F_k(\chi^2) R + \frac{1}{2} g^{\mu\nu} \partial_\mu \chi \partial_\nu \chi \right)$$

fixed point effective action :

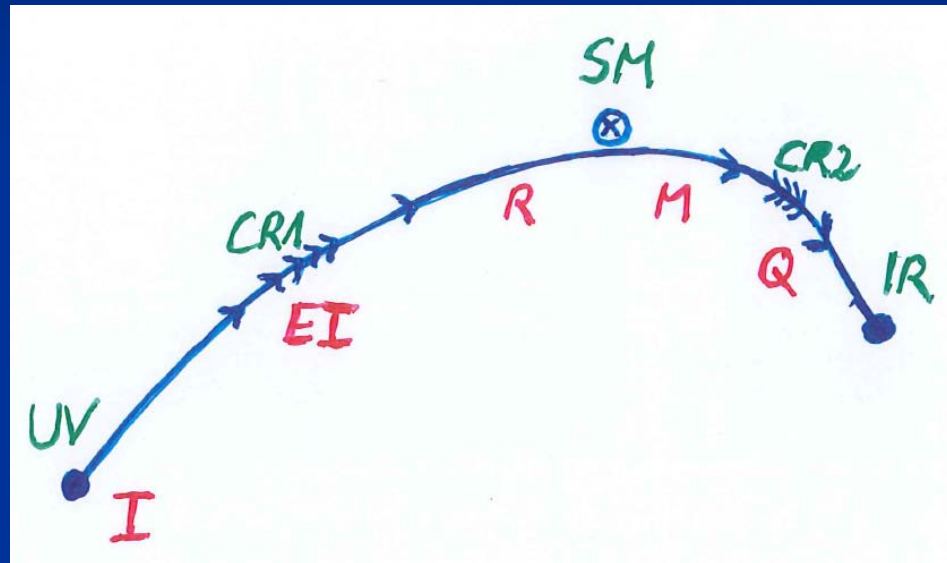
$$\Gamma = \int d^4x \sqrt{g} \left( \frac{1}{2} g^{\mu\nu} \partial_\mu \chi \partial_\nu \chi - \frac{1}{2} \xi \chi^2 R \right)$$

large field  
behavior of  $F$



$$y = \frac{\chi^2}{k^2}$$

# Possible consequences of crossover in quantum gravity

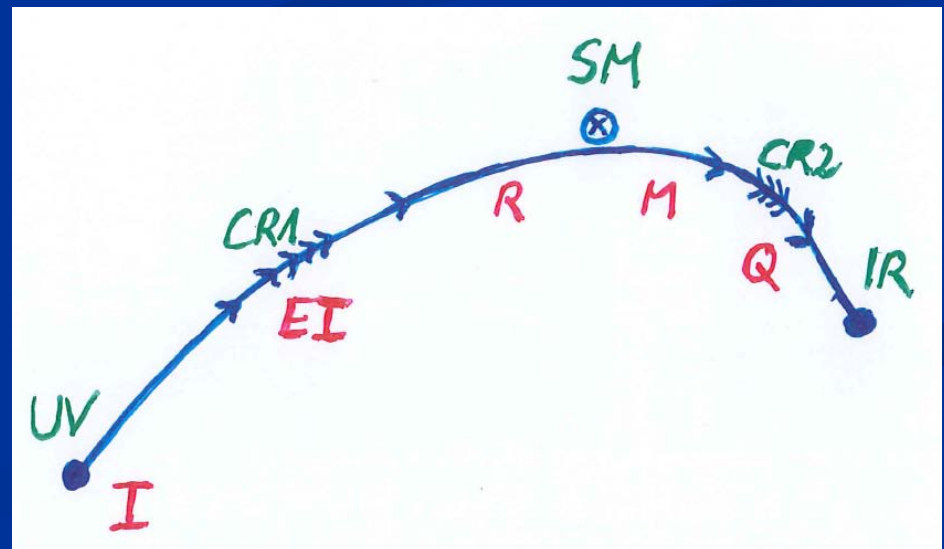


Realistic model for inflation and dark energy  
with single scalar field

# Cosmological solution : crossover from UV to IR fixed point

- Dimensionless functions as  $B$   
depend only on ratio  $\mu/\chi$ .
- IR:  $\mu \rightarrow 0$  ,  $\chi \rightarrow \infty$
- UV:  $\mu \rightarrow \infty$  ,  $\chi \rightarrow 0$

**Cosmology makes  
crossover between  
fixed points by  
variation of  $\chi$ .**



# renormalization flow and cosmological evolution

- renormalization flow as function of  $\mu$

is mapped by dimensionless functions to

- field dependence of effective action on scalar field  $\chi$

translates by solution of field equation to

- dependence of cosmology on time  $t$  or  $\eta$



# Simplicity

simple description of **all** cosmological epochs

natural incorporation of Dark Energy :

- inflation
- Early Dark Energy
- present Dark Energy dominated epoch

# conclusions

Quantum gravity may be observable in  
dynamics of present Universe

Fixed points and scale symmetry crucial

Big bang singularity is artefact  
of inappropriate choice of field variables –  
no physical singularity

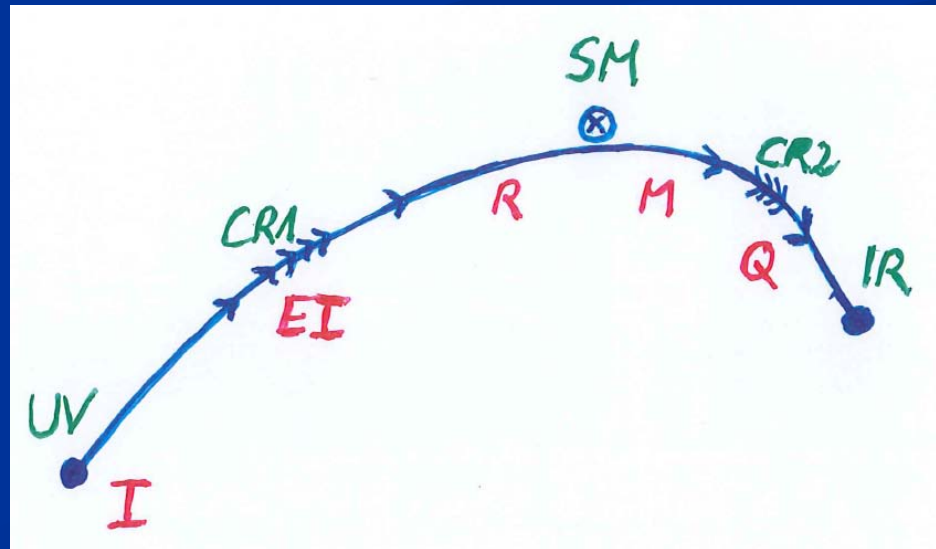
## conclusions (2)

- crossover in quantum gravity is reflected in crossover in cosmology
- quantum gravity becomes testable by cosmology
- quantum gravity plays a role not only for primordial cosmology
- crossover scenario explains different cosmological epochs
- simple model is compatible with present observations
- no more parameters than  $\Lambda$ CDM : tests possible

# Growing neutrino masses and quintessence

# Second stage of crossover

- from SM to IR
- in sector Beyond Standard Model
- affects neutrino masses first ( seesaw or cascade mechanism )



# Varying particle masses at onset of second crossover

- All particle masses **except for neutrinos** are proportional to  $\chi$ .
- Ratios of particle masses remain constant.
- Compatibility with observational bounds on time dependence of particle mass ratios.
- Neutrino masses show stronger increase with  $\chi$ , such that **ratio neutrino mass over electron mass grows**.

# connection between dark energy and neutrino properties

$$[\rho_h(t_0)]^{\frac{1}{4}} = 1.27 \left( \frac{\gamma m_\nu(t_0)}{eV} \right)^{\frac{1}{4}} 10^{-3} eV$$

L.Amendola,  
M.Baldi, ...

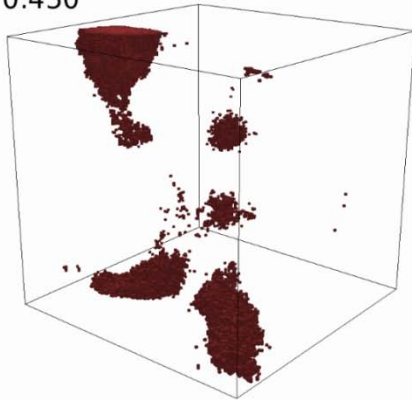
present dark energy density given by neutrino mass

present equation  
of state given by  
neutrino mass !

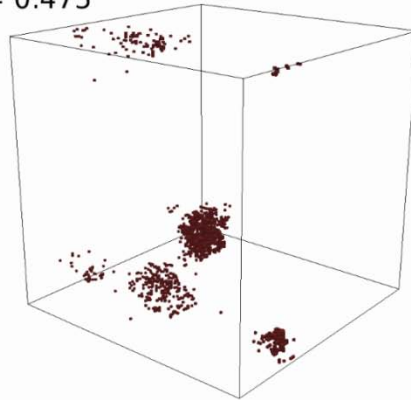
$$w_0 \approx -1 + \frac{m_\nu(t_0)}{12eV}$$

# Oscillating neutrino lumps

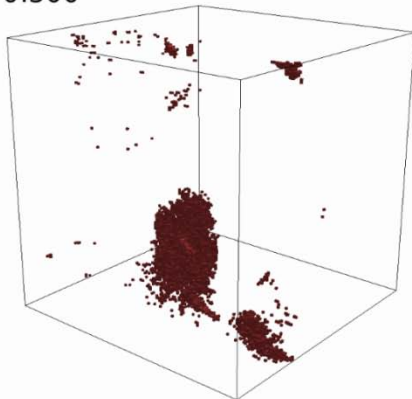
$a = 0.450$



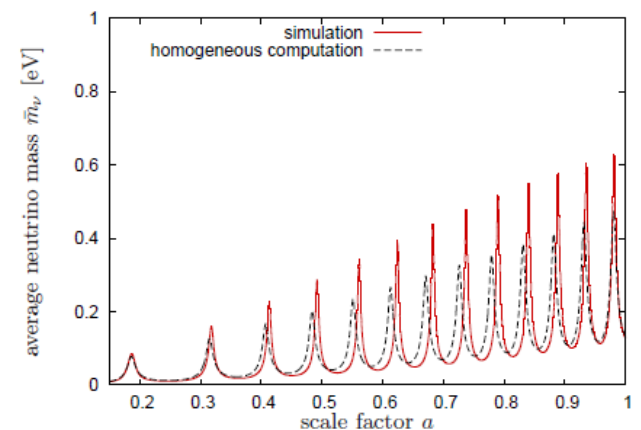
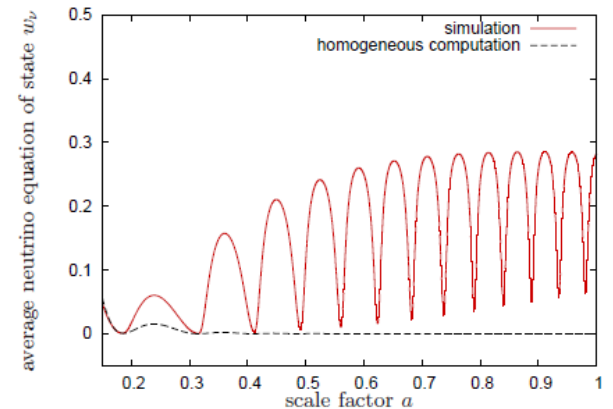
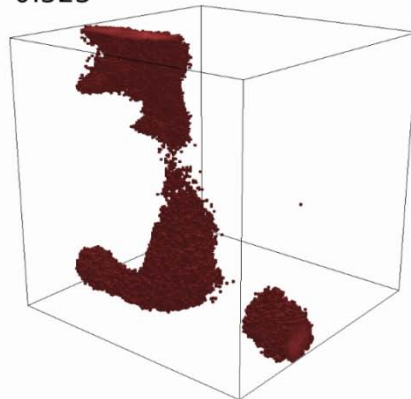
$a = 0.475$



$a = 0.500$



$a = 0.525$

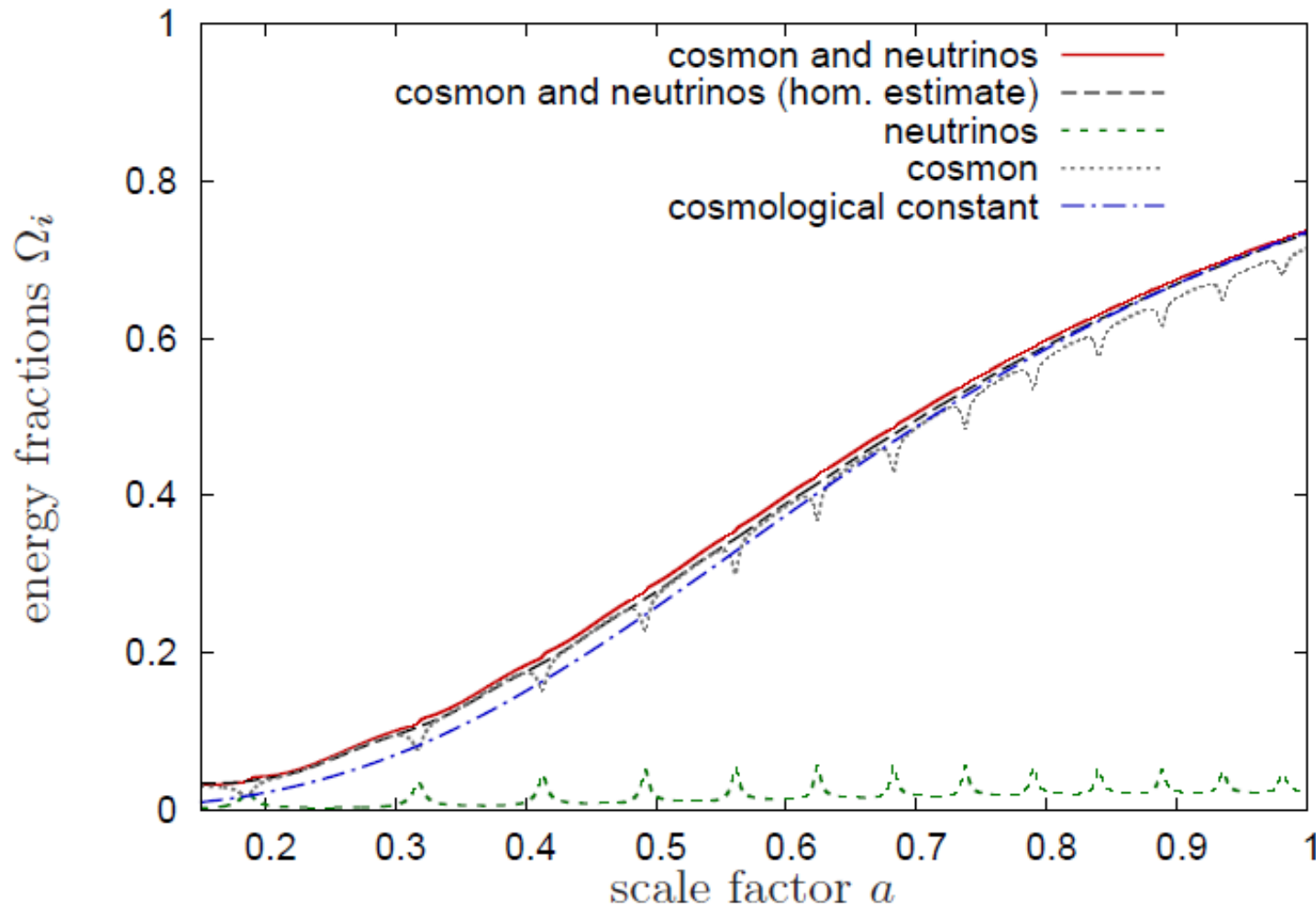


Y.Ayaita, M.Weber,...

Ayaita, Baldi, Fuehrer,  
Puchwein,...



# Evolution of dark energy similar to $\Lambda$ CDM



# Compatibility with observations and possible tests

- Realistic inflation model
- Almost same prediction for radiation, matter, and Dark Energy domination as  $\Lambda$ CDM
- Presence of small fraction of Early Dark Energy
- Large neutrino lumps

# conclusions (3)

- Variable gravity cosmologies can give a simple and realistic description of Universe
- Compatible with tests of equivalence principle and bounds on variation of fundamental couplings if nucleon and electron masses are proportional to variable Planck mass
- Cosmon dependence of ratio neutrino mass/ electron mass can explain why Universe makes a transition to Dark Energy domination **now**
- **characteristic signal : neutrino lumps**

The background is a solid dark blue. On the right side, there are several overlapping, wavy, light blue lines that create a sense of movement or depth, resembling stylized waves or smoke.

end

# Primordial flat frame

$$\Gamma = \int_x \sqrt{g} \left\{ -\frac{1}{2} \chi^2 R + \bar{\lambda} \chi^4 \ln \left( \frac{\bar{m}}{\chi} \right) + \left[ \ln^{-1} \left( \frac{\bar{m}}{\chi} \right) - 3 \right] \partial^\mu \chi \partial_\mu \chi \right\}$$

$$a = a_\infty \exp \left\{ -\frac{\tilde{c}_H}{\ln \left( \frac{\bar{m}}{\chi} \right)} \right\}$$

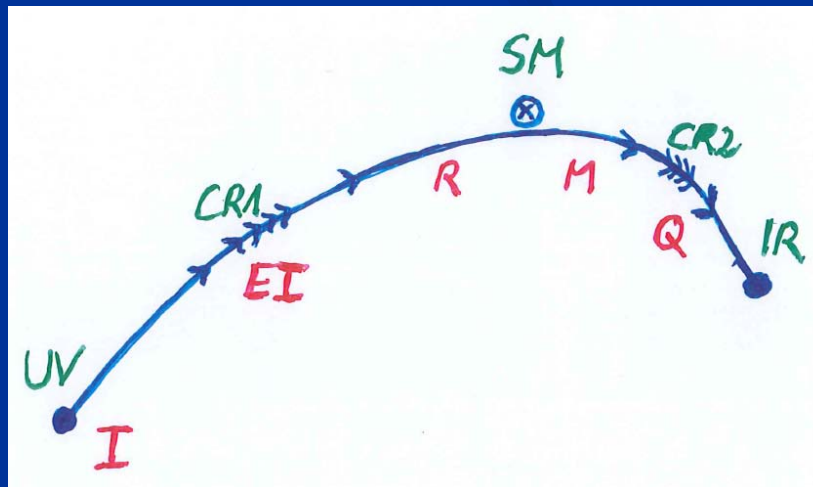
- Minkowski space in infinite past
- absence of any singularity
- geodesic completeness

# First step of crossover ends inflation

- induced by crossover in B

$$B^{-1} - \frac{\kappa}{\sigma} \ln B = \kappa \left[ \ln \left( \frac{\chi}{\mu} \right) - c_t \right] = \kappa \ln \left( \frac{\chi}{m} \right)$$

- after crossover B changes only very slowly



# Scaling solutions near SM fixed point

( approximation for constant B )

$$H = b\mu \ , \ \chi = \chi_0 \exp(c\mu t).$$

Different scaling solutions for  
radiation domination and  
matter domination

# Radiation domination

$$c = \frac{2}{\sqrt{K+6}}$$

$$b = -\frac{c}{2}$$

**Universe  
shrinks !**

$$T_{00} = \rho = \bar{\rho} \mu^2 \chi^2$$

$$\bar{\rho}_r = -3 \frac{K+5}{K+6}$$

$$\mathbf{K = B - 6}$$

solution exists for  $B < 1$  or  $K < -5$

$$S = \int_x \sqrt{g} \left\{ -\frac{1}{2} \chi^2 R + \frac{1}{2} K(\chi) \partial^\mu \chi \partial_\mu \chi + V(\chi) \right\}$$

$$H = b\mu, \quad \chi = \chi_0 \exp(c\mu t)$$



# Varying particle masses near SM fixed point

- All particle masses are proportional to  $\chi$ .  
( scale symmetry )
- Ratios of particle masses remain constant.
- Compatibility with observational bounds on time dependence of particle mass ratios.

# Scaling of particle masses

mass of electron or nucleon is proportional to variable Planck mass  $\chi$  !

effective potential for Higgs doublet  $h$

$$\tilde{V}_h = \frac{1}{2} \lambda_h (\tilde{h}^\dagger \tilde{h} - \epsilon_h \chi^2)^2.$$

# cosmon coupling to matter

$$K(\ddot{\chi} + 3H\dot{\chi}) + \frac{1}{2} \frac{\partial K}{\partial \chi} \dot{\chi}^2 = -\frac{\partial V}{\partial \chi} + \frac{1}{2} \frac{\partial F}{\partial \chi} R + q_\chi$$

$$q_\chi = -(\rho - 3p)/\chi$$

$$F = \chi^2$$

# Matter domination

$$c = \sqrt{\frac{2}{K+6}},$$

$$b = -\frac{1}{3}\sqrt{\frac{2}{K+6}} = -\frac{1}{3}c,$$

**Universe shrinks !**

$$T_{00} = \rho = \bar{\rho}\mu^2\chi^2,$$

solution exists for

$$B < 4/3, \quad K < -14/3$$

$$\bar{\rho}_m = -\frac{2(3K+14)}{3(K+6)}$$

$$K = B - 6$$

# Early Dark Energy

Energy density in radiation increases ,  
proportional to cosmon potential

$$T_{00} = \rho = \bar{\rho} \mu^2 \chi^2, \quad V(\chi) = \mu^2 \chi^2,$$

fraction in early dark energy

$$\Omega_h = \frac{\rho_h}{\rho_r + \rho_h} = \frac{nB(\chi)}{4}$$

or m

observation requires  **$B < 0.02$**  ( at CMB emission )

# Dark Energy domination

neutrino masses scale  
differently from electron mass

$$\left. \frac{\partial \ln m_\nu}{\partial \ln \chi} \right|_{\text{today}} = 2\tilde{\gamma} + 1$$



$$m_\nu = \bar{c}_\nu \chi^{2\tilde{\gamma}+1}$$

$$\chi q_\chi = -(2\tilde{\gamma} + 1)(\rho_\nu - 3p_\nu)$$

new scaling solution. not yet reached.  
at present : transition period

$$\frac{\rho_\nu}{\chi^2} = \bar{\rho}_\nu \mu^2$$

$$b = \frac{1}{3}(2\tilde{\gamma} - 1)c$$

# Infrared fixed point

■  $\mu \rightarrow 0$

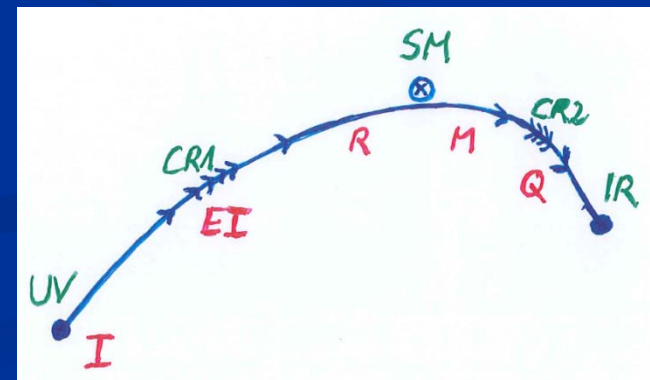
■  $B \rightarrow 0$

$$\mu \partial_\mu B = \kappa B^2 \quad \text{for} \quad B \rightarrow 0$$

$$\Gamma = \int_x \sqrt{g} \left\{ -\frac{1}{2} \chi^2 R + \mu^2 \chi^2 + \frac{1}{2} (B(\chi/\mu) - 6) \partial^\mu \chi \partial_\mu \chi \right\}$$

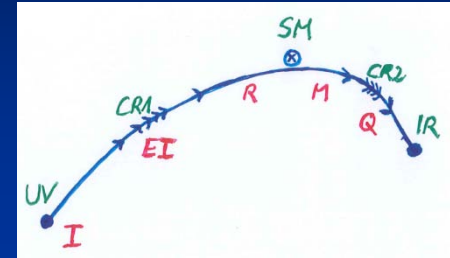
■ no intrinsic mass scale

■ scale symmetry



# Ultraviolet fixed point

■  $\mu \rightarrow \infty$



■ kinetic diverges

$$B = b \left( \frac{\mu}{\chi} \right)^{\sigma} = \left( \frac{m}{\chi} \right)^{\sigma}$$

■ scale symmetry with anomalous dimension  $\sigma$

$$g_{\mu\nu} \rightarrow \alpha^2 g_{\mu\nu} , \quad \chi \rightarrow \alpha^{-\frac{2}{2-\sigma}} \chi$$



# Renormalized field at UV fixed point

$$\chi_R = b^{\frac{1}{2}} \left(1 - \frac{\sigma}{2}\right)^{-1} \mu^{\frac{\sigma}{2}} \chi^{1-\frac{\sigma}{2}}$$

$$1 < \sigma$$

$$\Gamma_{UV} = \int_x \sqrt{g} \left\{ \frac{1}{2} \partial^\mu \chi_R \partial_\mu \chi_R - \frac{1}{2} C R^2 + D R^{\mu\nu} R_{\mu\nu} \right\}$$

no mass  
scale

$$\Delta\Gamma_{UV} = \int_x \sqrt{g} E \left( \mu^2 - \frac{R}{2} \right) \mu^{-\frac{2\sigma}{2-\sigma}} \chi_R^{\frac{4}{2-\sigma}},$$

deviation from  
fixed point  
vanishes for

$$E = b^{-\frac{2}{2-\sigma}} \left(1 - \frac{\sigma}{2}\right)^{\frac{4}{2-\sigma}}$$

$$\mu \rightarrow \infty$$