## Quantum gravity prediction

## dynamical dark energy

- for

$$\partial_k \Gamma_k[\varphi] = \frac{1}{2} \text{Tr} \left\{ \left( \Gamma_k^{(2)}[\varphi] + R_k \right)^{-1} \partial_k R_k \right\}$$

#### Quantum gravity

Gravity is field theory. Similar to electrodynamics. Metric field.

Gravity is gauge theory. Similar to QED or QCD. Gauge symmetry: general coordinate transformations (diffeomorphisms)

Quantum gravity: include metric fluctuations in functional integral

#### Metric fluctuations matter



#### Quantum gravity needs method to take them into account

#### Quantum gravity

- Quantum gravity is similar to other quantum field theories
- Difference: metric is tensor, gauge bosons are vectors
- Difference: Einstein gravity is not perturbatively renormalizable
- no small dimensionless coupling constant, effective coupling q<sup>2</sup>/M<sup>2</sup>



Quantum gravity is non-perturbatively renormalizable

Asymptotic safety : non-perturbative renormalizabilty Weinberg, Reuter, ...

Use functional renormalization !

### Flowing couplings

Couplings change with renormalization scale k due to (quantum) fluctuations.

Renormalization scale k : Only fluctuations with momenta larger k are included for (quantum) effective action.

Flow of k to zero : all fluctuations included, IR-limit Flow of k to infinity : UV-limit

k-dependence can differ from momentum dependence

#### Ultraviolet fixed point

- Flow of couplings stops for  $k \rightarrow \infty$
- Self-similarity: no explicit dependence on k
- Theory can be extrapolated to arbitrary short distances
- Completeness

Renormalizability

#### Asymptotic safety Asymptotic freedom





### Asymptotically safe gravity

Ultraviolet fixed point exists for quantum field theory for metric ( or vierbein ).

#### Predictivity

- asymptotic safety is predictive
   few relevant parameters govern flow away from UV-fixed point
- translate to renormalizable couplings in standard model or extensions

#### Prediction of mass of Higgs boson

#### Asymptotic safety of gravity and the Higgs boson mass

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#### Abstract

There are indications that gravity is asymptotically safe. The Standard Model (SM) plus gravity could be valid up to arbitrarily high energies. Supposing that this is indeed the case and assuming that there are no intermediate energy scales between the Fermi and Planck scales we address the question of whether the mass of the Higgs boson  $m_H$  can be predicted. For a positive gravity induced anomalous dimension  $A_{\lambda} > 0$  the running of the quartic scalar self interaction  $\lambda$  at scales beyond the Planck mass is determined by a fixed point at zero. This results in  $m_H = m_{\min} = 126$  GeV, with only a few GeV uncertainty. This prediction is independent of the details of the short distance running and holds for a wide class of extensions of the SM as well.



#### Metric + scalar field

Inflation : add scalar field (inflaton)

 Dynamical dark energy or quintessence: add scalar field (cosmon) Can the scalar potential be predicted by functional renormalization for quantum gravity ?

## Quantum gravity prediction for generic form of potential for singlet scalar



prediction of (approximate) quantum scale symmetry:

dynamical dark energy, generated by scalar field (cosmon)

> C.Wetterich, Nucl.Phys.B302(1988)668, 24.9.87 P.J.E.Peebles, B.Ratra, ApJ.Lett.325(1988)L17, 20.10.87

#### **Quintessential inflation**



Spokoiny, Peebles, Vilenkin, Peloso, Rosati, Dimopoulos, Valle, Giovannini, Brax, Martin, Hossain, Myrzakulov, Sami, Saridakis, de Haro, Salo, Bettoni, Rubio...



### **Scaling solutions**

- At fixed point: all (infinitely many) dimensionless couplings take fixed values Full momentum dependence of graviton propagator (no polynomial expansion) Whole scalar potential is fixed, for arbitrary values of scalar field
- Functional flow equations are needed

#### Scaling solutions are restrictive

 scaling solutions are solutions of non-linear differential equations
 scaling potential needs to extend over whole range of scalar field
 predictivity !

in presence of gravitational fluctuations: scalar effective potential no longer approximated by polynomial Relevant parameters and relevant functions

flow away from scaling solution is dominated by relevant functions

these should not be singular for finite values of scalar field, or diverge exponentially for increasing scalar field

typically only a few, associated to relevant parameters
Predictivity !

#### Dominant relevant mass scale

- Flow away from fixed point induces intrinsic mass scales by dimensional transmutation
- Intrinsic violation of quantum scale symmetry
- Can be in different steps ( example Planck mass, QCD scale )
- Dominant relevant mass scale is the largest intrinsic scale

#### Field dependent mass scales

Scenario: dominant relevant mass scale much below electron mass (say 10<sup>-3</sup> eV in appropriate units) Planck mass and particle masses are given by cosmological (expectation) value of scalar field Spontaneous breaking of scale symmetry Scaling solution very good approximation for all momentum scales larger than dominant relevant mass scale

#### Scale symmetric standard model

**Replace all mass scales by scalar field**  $\chi$ 

Fujii, Zee, CW

(2) Strong gauge coupling, normalized at  $\mu = \chi$ , is independent of  $\chi$ 

$$g(\chi) = \bar{g}$$
  $\land A_{\text{QCD}} = \chi \exp\left(-\frac{1}{b_0 \bar{g}^2}\right)$   $b_0 = \frac{1}{16\pi^2} \left(22 - \frac{4}{3}N_f\right)$ 

(3) Similar for all dimensionless couplings

Quantum effective action for standard model does not involve intrinsic mass or length

**Quantum scale symmetry** CW'87, Shaposhnikov et al

#### Cosmon

- Spontaneously broken scale symmetry induces a Goldstone boson
- Massless dilaton
- Intrinsic mass scale from dominant relevant mass scale: Pseudo-Goldstone boson acquires a tiny mass
- Cosmon

Naturally very light scalar particle !



#### Dilaton quantum gravity

functional renormalization for quantum gravity coupled to a scalar field

Henz, Pawlowski, Rodigast, Yamada, Reichert, Eichhorn, Pauly, Laporte, Pereira, Saueressig, Wang, Knorr, ...

for low order polynomial expansion of potential : Percacci, Narain, ...

# Derivative expansion of effective action

$$\Gamma = \int_{\chi} \sqrt{g} \left\{ -\frac{1}{2} F(\chi) R + \frac{1}{2} K(\chi) \partial^{\mu} \chi \partial_{\mu} \chi + U(\chi) \right\}$$

#### variable gravity

#### Flow equation for scalar potential

$$u = \frac{U}{k^4}, \quad w = \frac{F}{2k^2} \qquad v = \frac{2U}{Fk^2} = \frac{u}{w}$$

$$ho = \chi^2/2$$
,  $ilde{
ho} = 
ho/k^2$ 

$$\partial_t u = \beta_u = -4u + 2\tilde{\rho}\,\partial_{\tilde{\rho}} u + 4c_{\rm U}$$

$$c_{\rm U} = \frac{1}{96\pi^2} \left( \frac{5}{1-v} + \frac{1}{1-v/4} \right) + b_{\rm U}$$

$$b_{\rm U}=\frac{N-4}{128\pi^2}$$

$$N = N_{\rm S} + 2N_{\rm V} - 2N_F$$

#### Differential equation for scaling solution

$$2\tilde{\rho}\frac{\partial u}{\partial\tilde{\rho}} = 4u - \frac{1}{24\pi^2}\left(\frac{5}{1-u/w} + \frac{1}{1-u/4w}\right) - 4b_{\mathrm{U}}$$

Differential equation for scaling solution for effective Planck mass

$$\Gamma = \int_{\chi} \sqrt{g} \left\{ -\frac{1}{2} F(\chi) R + \frac{1}{2} K(\chi) \partial^{\mu} \chi \partial_{\mu} \chi + U(\chi) \right\} \qquad u = \frac{U}{k^4}, \quad w = \frac{F}{2k^2}$$

$$2\tilde{\rho}\,\partial_{\tilde{\rho}}w = 2\left(w - c_M\right)$$

$$c_M = \frac{25}{128\pi^2} \left( 1 - \frac{u}{w} \right)^{-1} + \frac{1}{192\pi^2} \left[ -N_{\rm S} \left( \frac{1}{1+u'} + \frac{3w'}{(1+u')^2} \right) + 2N_V \left( \frac{3}{1+g^2\tilde{\rho}} - 1 \right) - \frac{N_F}{1+y^2\tilde{\rho}} + \frac{43}{6} \right]$$

M.Yamada,

# Scaling potential for particles of standard model



u : dimensionless scalar potential u= U/k<sup>4</sup>

x : logarithm of scalar field value

$$2\tilde{\rho}\frac{\partial u}{\partial\tilde{\rho}} = 4u - \frac{1}{24\pi^2} \left(\frac{5}{1 - u/w} + \frac{1}{1 - u/4w}\right) - 4b_{\mathrm{U}}$$

### Generic form of scaling potential

Interpolates between two plateaus
Scalar potential = field dependent "cosmological constant"
Effectively massless particles contribute to flow
Different numbers of massless particles in different regions of field space

# Coefficient of curvature scalar in standard model



x : logarithm of scalar field value

w : dimensionless field dependent squared Planck

mass

 $u = \frac{U}{k^4}, \quad w = \frac{F}{2k^2}$ 

non-minimal coupling of scalar field to gravity:  $F = \xi \chi^2 R$ 

### Approximate scaling solution

I flat potential: u constant

non-minimal scalar- gravity coupling:

for large scalar field w increases proportional  $\chi^2$ 



looks natural no small parameter no tuning





#### Scaling solution in Einstein frame



Weyl transformation for variable gravity

$$g_{\mu\nu} = (M^2/F)g'_{\mu\nu} \quad \varphi = 4M\ln(\chi/k)$$

$$\Gamma = \int_{\chi} \sqrt{g} \left\{ -\frac{1}{2} F(\chi) R + \frac{1}{2} K(\chi) \partial^{\mu} \chi \partial_{\mu} \chi + U(\chi) \right\}$$

$$\Gamma = \int_{\chi} \sqrt{g} \left\{ -\frac{M^2}{2} R' + \frac{1}{2} Z(\varphi) \partial^{\mu} \varphi \partial_{\mu} \varphi + V(\varphi) \right\}$$

$$V(\varphi) = \frac{UM^4}{F^2}$$

$$Z(\varphi) = \frac{1}{16} \left\{ \frac{\chi^2 K}{F} + \frac{3}{2} \left( \frac{\partial \ln F}{\partial \ln \chi} \right)^2 \right\}$$

### **Scaling solution**

$$U = u_0 k^4$$

$$F = 2w_0k^2 + \xi\chi^2$$



For low energy standard model :

$$u_{\infty} = \frac{7}{256\pi^2}$$

#### Scaling solution in Einstein frame

$$V = \frac{u_0 M^4}{\left(2w_0 + \xi \exp\left(\frac{\varphi}{2M}\right)\right)^2}$$
$$= \frac{u_0 M^4}{\xi^2} \left[1 + \frac{2w_0}{\xi} \exp\left(-\frac{\varphi}{2M}\right)\right]^{-2} \exp\left(-\frac{\varphi}{M}\right)$$



#### Mass scales in Einstein frame

Renormalization scale k is no longer present Planck mass M not intrinsic: introduced only by change of variables

$$V = \frac{u_0 M^4}{\left(2w_0 + \xi \exp\left(\frac{\varphi}{2M}\right)\right)^2}$$
$$= \frac{u_0 M^4}{\xi^2} \left[1 + \frac{2w_0}{\xi} \exp\left(-\frac{\varphi}{2M}\right)\right]^{-2} \exp\left(-\frac{\varphi}{M}\right)$$



## Asymptotic solution of cosmological constant problem

$$V = \frac{u_0 M^4}{\left(2w_0 + \xi \exp\left(\frac{\varphi}{2M}\right)\right)^2}$$
$$= \frac{u_0 M^4}{\xi^2} \left[1 + \frac{2w_0}{\xi} \exp\left(-\frac{\varphi}{2M}\right)\right]^{-2} \exp\left(-\frac{\varphi}{M}\right)$$



no tiny parameter !

$$V(\varphi) = \frac{UM^4}{F^2}$$

#### Predictions for primordial cosmic fluctuations in inflationary cosmology

Depend on form of kinetial K

$$\Gamma = \int_{\chi} \sqrt{g} \left\{ -\frac{1}{2} F(\chi) R + \frac{1}{2} K(\chi) \partial^{\mu} \chi \partial_{\mu} \chi + U(\chi) \right\}$$

 so far form of K assumed : realistic cosmology possible for suitable K
 K needs to be computed!

#### **Cosmological solution**

scalar field χ vanishes in the infinite past
 scalar field χ diverges in the infinite future



J.Rubio,...

#### Conclusions

- Fluctuations of the metric matter
- They influence the behavior of scalar potentials for all field values
- Quantum gravity relevant for early cosmology and late cosmology
- Quantum scale symmetry is central ingredient for understanding cosmology
- Fundamental scale invariance is highly predictive
- Understanding of quantum gravity fluctuations still in beginning stage

#### Conclusion

Fixed points of quantum gravity with associated quantum scale symmetry, scaling solutions and relevant parameters are crucial for understanding the evolution of our Universe

