

Problem 29: Bogoliubov transformation

- (a) The bosonic operators \hat{a} , \hat{a}^\dagger satisfy the canonical commutation relations $[\hat{a}, \hat{a}^\dagger] = 1$, $[\hat{a}, \hat{a}] = 0$. Define a new pair of operators by the transformation

$$\hat{b} = u\hat{a} + v\hat{a}^\dagger, \quad \hat{b}^\dagger = u^*\hat{a}^\dagger + v^*\hat{a} \quad (1)$$

with $u, v \in \mathbb{C}$. Under which condition is this transformation canonical, *i.e.*, $[\hat{b}, \hat{b}^\dagger] = 1$? Find a parametrization for u and v in terms of real numbers.

- (b) The Bogoliubov Hamiltonian for spin-1/2 electrons ($\sigma = \uparrow, \downarrow$) with isotropic dispersion relation $\xi_{\mathbf{k}}$ and real energy gap $\Delta \in \mathbb{R}$ is defined as

$$\mathcal{H} = \sum_{\mathbf{k}\sigma} \xi_{\mathbf{k}} c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}\sigma} - \sum_{\mathbf{k}} (\Delta c_{-\mathbf{k}\downarrow} c_{\mathbf{k}\uparrow} + \Delta c_{\mathbf{k}\uparrow}^\dagger c_{-\mathbf{k}\downarrow}^\dagger). \quad (2)$$

Show that this Hamiltonian can be brought into the matrix form

$$\mathcal{H} = \sum_{\mathbf{k}} \begin{pmatrix} c_{\mathbf{k}\uparrow}^\dagger & c_{-\mathbf{k}\downarrow} \end{pmatrix} \begin{pmatrix} \xi_{\mathbf{k}} & -\Delta \\ -\Delta & -\xi_{\mathbf{k}} \end{pmatrix} \begin{pmatrix} c_{\mathbf{k}\uparrow} \\ c_{-\mathbf{k}\downarrow}^\dagger \end{pmatrix} + \sum_{\mathbf{k}} \xi_{\mathbf{k}}. \quad (3)$$

Diagonalize \mathcal{H} with the help of the Bogoliubov transformation

$$\begin{pmatrix} c_{\mathbf{k}\uparrow} \\ c_{-\mathbf{k}\downarrow}^\dagger \end{pmatrix} = \begin{pmatrix} \cos \theta_{\mathbf{k}} & \sin \theta_{\mathbf{k}} \\ \sin \theta_{\mathbf{k}} & -\cos \theta_{\mathbf{k}} \end{pmatrix} \begin{pmatrix} \gamma_{\mathbf{k}\uparrow} \\ \gamma_{-\mathbf{k}\downarrow}^\dagger \end{pmatrix}; \quad (4)$$

determine $\theta_{\mathbf{k}}$ such that the anomalous terms with $\gamma\gamma$ and $\gamma^\dagger\gamma^\dagger$ vanish. Sketch the spectrum of fermionic quasiparticles $\gamma_{\mathbf{k}\sigma}$ as a function of \mathbf{k} and of $\xi_{\mathbf{k}}$.

Problem 30: Energy of the superconducting ground state

Compute the condensation energy in the superconducting ground state ($\Delta > 0$) relative to the normal state ($\Delta = 0$, $E_k = |\xi_k|$),

$$\Delta E = \langle \mathcal{H} - \mu N \rangle_S - \langle \mathcal{H} - \mu N \rangle_N, \quad (5)$$

expressed in terms of the gap Δ , the volume and the electronic density of states ν_0 , in the limit of weak coupling $\nu_0 g \ll 1$. One can, *e.g.*, compute the expectation value of $\mathcal{H} - \mu N$ with the BCS wave function $|\text{BCS}\rangle$, or (recommended) evaluate the term for the ground state energy in the Bogoliubov Hamiltonian. Assume that the attractive interaction is constant in a shell of width $\hbar\omega_D$ around the Fermi surface (as in the lecture), and use particle-hole symmetry within this shell, $-\sum_{k < k_F} \xi_k = \sum_{k > k_F} \xi_k$.

Problem 31: Hubbard-Stratonovich transformation

The partition function of an electron-phonon system can be written as the path integral

$$\mathcal{Z} = \int D\psi^* D\psi \int D\phi^* D\phi \exp \left\{ -S_{\text{el}}[\psi^*, \psi] - S_{\text{ph}}[\phi^*, \phi] - S_{\text{el-ph}}[\psi^*, \psi, \phi^*, \phi] \right\} \quad (6)$$

with Grassmann fields ψ^* , ψ representing the fermions and complex fields ϕ^* , ϕ for the phonons. The action for the phonons S_{ph} and for the electron-phonon interaction $S_{\text{el-ph}}$ have the form

$$S_{\text{ph}}[\phi^*, \phi] = \sum_q \phi_q^* (-i\omega_m + \omega_q) \phi_q \quad (7)$$

$$S_{\text{el-ph}}[\psi^*, \psi, \phi^*, \phi] = \sum_q M_{\mathbf{q}} \rho_{\mathbf{q}} (\phi_{\mathbf{q}} + \phi_{-\mathbf{q}}^*) \quad (8)$$

with the electronic density operator $\rho_{\mathbf{q}} = \sum_k \psi_{k+\mathbf{q}}^* \psi_k$, electron-phonon coupling $M_{\mathbf{q}}$, and the usual kinetic term $S_{\text{el}}[\psi^*, \psi]$ for the electrons. The multi-indices $q = (i\omega_m, \mathbf{q})$ are a shorthand notation for frequency and momentum indices, and the sum $\sum_q = \sum_{\mathbf{q}} \beta^{-1} \sum_m$ includes a Matsubara sum as well as a momentum sum.

Perform a Gaussian integration over the phonon fields ϕ^* , ϕ in \mathcal{Z} and derive an effective action $S_{\text{eff}}[\psi^*, \psi]$ for the electrons alone (neglect a constant term which does not depend on ψ^* and ψ):

$$\mathcal{Z} = \int D\psi^* D\psi e^{-S_{\text{eff}}[\psi^*, \psi]} \quad (9)$$

$$e^{-S_{\text{eff}}[\psi^*, \psi]} = \int D\phi^* D\phi \exp \left\{ -S_{\text{el}}[\psi^*, \psi] - S_{\text{ph}}[\phi^*, \phi] - S_{\text{el-ph}}[\psi^*, \psi, \phi^*, \phi] \right\} \quad (10)$$

The effective action contains an interaction term between the electrons; compare this to the result of the Fröhlich transformation given in the lecture.