History, Present, Future: Functional RG for unconventional superconductors

Ronny Thomale

Julius-Maximilians Universität Würzburg



"Functional Renormalization - from quantum gravity and dark energy to ultracold atoms and condensed matter"

Unconventional superconductors





Outline

History

- Timeline of materials studied through FRG
- Recent methodological step: multi-band/orbital superconductivity

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- Pnictides: extended d-wave vs. extended s-wave
- Sodium cobaltates: chiral d-wave?

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- Sodium cobaltates: chiral d-wave?

Future

- SrPtAs: Weyl superconductors
- LAO/STO heterostructures

History



Wednesday, March 8, 17



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The favored real space pairing function

 $\Delta_{i,j} \sim \langle c_i^{\dagger} c_j^{\dagger} \rangle$

transforms under irreducible point group representations

C_{4v}	E	<i>C</i> ₂	$2C_4$	$2\sigma_v$	$2\sigma_d$
A_1	1	1	1	1	1
A2	1	1	1	-1	-1
B ₁	1	1	-1	1	-1
<i>B</i> ₂	1	1	-1	-1	1
E	2	-2	0	0	0

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B_1	1	1	-1	1	-1	
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$$\Delta_s(\mathbf{k}) = 1 + \alpha(\cos k_x + \cos k_y)$$

$$\Delta_{d_{x^2-y^2}}(\mathbf{k}) = \cos k_x - \cos k_y + \alpha (\cos 2k_x - \cos 2k_y)$$

Extended s-wave in the pnictides



Extended s-wave in the pnictides



LaFeAsO vs. LaFePO: s-wave anisotropy

Thomale, Platt, Hanke, Bernevig, PRL 106, 187003 (2011)



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Prediction of extended d-wave in pnictides

RAPID COMMUNICATIONS

PHYSICAL REVIEW B 80, 180505(R) (2009)

Functional renormalization-group study of the doping dependence of pairing symmetry in the iron pnictide superconductors

 Ronny Thomale,¹ Christian Platt,² Jiangping Hu,³ Carsten Honerkamp,² and B. Andrei Bernevig⁴
¹Institut für Theorie der Kondensierten Materie, Universität Karlsruhe, D 76128 Karlsruhe, Germany
²Theoretical Physics, University of Würzburg, D-97074 Würzburg, Germany
³Department of Physics, Purdue University, West Lafayette, Indiana 47907, USA
⁴Department of Physics, Princeton University, Princeton, New Jersey 08544, USA (Received 19 October 2009; published 13 November 2009)

We use the functional renormalization group to analyze the phase diagram of a four-band model for the iron pnictides subject to band interactions with certain A_{1g} momentum dependence. We determine the parameter regimes where an extended *s*-wave pairing instability with and without nodes emerges. For electron doping, the parameter regime in which a nodal gap appears is in correspondence to recent predictions [A. Chubukov *et al.*, arXiv:0903.5547 (unpublished)], however, at very low T_c . Upon hole doping, the *s*-wave gap never becomes nodal: above a critical strength of the intraband repulsion, the system favors an exotic extended *d*-wave instability on the enlarged hole pockets. At half filling, we find that a strong momentum dependence of

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Present

D-wave in hole doped iron pnictides?

- moderate hole doping: electron and hole pockets
- strong hole doping: only hole pockets present



FRG: s to d-wave transition in K-doped Ba-122

Thomale, Platt, Hanke, Hu, Bernevig, PRL 107, 117001 (2011).



Evidence for d-wave: thermal transport

Taillefer group, Reid et al. PRL 2012



Evidence against d-wave: Laser ARPES



Raman: extended d-wave dominant subleading order

Maiti et al., PRL 117, 257001 (2016); Böhm et al., in preparation.



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Raman: extended d-wave dominant subleading order

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To be submitted: Raman scattering confirms the prediction of extended d-wave as the dominant subleading d-wave instability in K-doped Ba-122.

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Chiral d-wave superconductivity

$$H_{d+id} = \sum_{\mathbf{k}} (c_{\mathbf{k}\uparrow}^{\dagger}, c_{-\mathbf{k}\downarrow}) \begin{pmatrix} \xi_{\mathbf{k}} & \Delta_{\mathbf{k}} \\ \Delta_{\mathbf{k}}^{*} & -\xi_{\mathbf{k}} \end{pmatrix} \begin{pmatrix} c_{\mathbf{k}\uparrow} \\ c_{-\mathbf{k}\downarrow}^{\dagger} \end{pmatrix}$$

 $\Delta_{\mathbf{k}} = \cos(k_x) - \cos(k_y) + i\sin(k_x)\sin(k_y) = |\Delta_{\mathbf{k}}|e^{i\varphi_{\mathbf{k}}}$ $\xi_{\mathbf{k}} = \cos(k_x) + \cos(k_y)$







Takada et al., Nature 422, 53-55 (2003)

Superconductivity in twodimensional CoO₂ layers

Kazunori Takada*‡, Hiroya Sakurai†, Eiji Takayama-Muromachi†, Fujio Izumi*, Ruben A. Dilanian* & Takayoshi Sasaki*‡

* Advanced Materials Laboratory, National Institute for Materials Science,

Tsukuba, Ibaraki 305-0044, Japan

† Superconducting Materials Center, National Institute for Materials Science,

0.020

0.018

0.016

0.014

Tsukuba, Ibaraki 305-0044, Japan

‡ CREST, Japan Science and Technology Corporation





Cobaltate phase diagram

The phase diagram hosts a plethora of phases interpolating between the parent insulating compound (x=0) and Na-doped insulating limit (x=1)



Cobaltate phase diagram



x=0.1

x=0.2

x=0.3



Cobaltate phase diagram



Evidence at x=0.3: Singlet, close-to-nodal superconductor



Hexagonal d-wave superconductivity

C_{6v}	E	<i>C</i> ₂	$2C_3$	2 <i>C</i> ₆	$3\sigma_v$	$3\sigma_d$
A_1	1	1	1	1	1	1
<i>A</i> ₂	1	1	1	1	-1	-1
B_1	1	-1	1	-1	1	-1
<i>B</i> ₂	1	-1	1	-1	-1	1
E_1	2	-2	-1	1	0	0
<i>E</i> ₂	2	2	-1	-1	0	0

 d_{x2-y2}/d_{xy}

Hexagonal d-wave superconductivity

C_{6v}	E	<i>C</i> ₂	$2C_3$	$2C_6$	$3\sigma_v$	$3\sigma_d$
A_1	1	1	1	1	1	1
A ₂	1	1	1	1	-1	-1
B_1	1	-1	1	-1	1	-1
<i>B</i> ₂	1	-1	1	-1	-1	1
E_1	2	-2	-1	1	0	0
<i>E</i> ₂	2	2	-1	-1	0	0



d_{x2-y2}/d_{xy}

The d-wave hexagonal point group representation is two-dimensional, implying a degeneracy at the instability level.

FRG analysis: anisotropic d+id phase

The multi-orbital Hubbard model at $x \sim 0.3$ yields a d+id-wave superconductor with a strongly anisotropic gap function.



FRG analysis: anisotropic d+id phase



FRG analysis: anisotropic d+id phase











Future

SrPtAs - a Weyl superconductor?

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PHYSICAL REVIEW B 89, 020509(R) (2014)

Mark H. Fischer,^{1,2} Titus Neupert,^{3,4} Christian Platt,⁵ Andreas P. Schnyder,⁶ Werner Hanke,⁵ Jun Goryo,⁷ Ronny Thomale,⁵ and Manfred Sigrist⁴



Bogoliubov spectrum



SrPtAs: First prediction of a Weyl superconductor



LAO/STO - tight binding setup



LAO/STO - tight binding setup



New methodological step: Spin-orbit FRG with complete double group implementation.

LAO/STO - s-wave vs. d-wave



Research team











































