

Observational constraints in nonlocally modified gravity

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Gravity at the Largest Scales 2015

Heidelberg

Introduction: Theory

- Inspiration

(Arkani-Hamed et al. 2002, Dvali 2006)

$$\mathcal{L}_{\text{proca}} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{m^2}{2}A_\mu A^\mu - A_\mu j^\mu \quad\Leftrightarrow\quad \mathcal{L}_{\text{nl}} = -\frac{1}{4}F_{\mu\nu}\left(1 - \frac{m^2}{\square}\right)F^{\mu\nu} - A_\mu j^\mu$$

where $(\square^{-1}\phi)(x) = \int d^4y G(x,y)\phi(y)$

- Applying the same idea to Fierz-Pauli massive gravity

$$\mathcal{L}_{\text{nl}} = \frac{1}{2}h_{\mu\nu}\left(1 - \frac{m^2}{\square}\right)\mathcal{E}^{\mu\nu\rho\sigma}h_{\rho\sigma} - 2m^2\chi\frac{1}{\square}\partial_\mu\partial_\nu(h^{\mu\nu} - \eta^{\mu\nu}h)$$

→ Obstruction: covariantization $\Rightarrow g^{\mu\nu}R_{\mu\nu} = 0$ "Covariant vDVZ discontinuity"

$$\left[\left(1 - \frac{m^2}{\square_g}\right)G_{\mu\nu}\right]^T = 8\pi GT_{\mu\nu} \quad(\text{Porrati 2002; Jaccard, Maggiore, Mitsou 2013})$$

- ▷ Unviable background cosmology
- ▷ $\square^{-1}R_{\mu\nu} \subset \square^{-1}G_{\mu\nu}$ generates instabilities (Ferreira, Maroto 2013)
- ▷ $g_{\mu\nu}\square^{-1}R \subset \square^{-1}G_{\mu\nu}$ stable (Foffa, Maggiore, Mitsou 2013)

Introduction: Phenomenology

Model RT : $G_{\mu\nu} - m^2 (g_{\mu\nu} \square_{ret}^{-1} R)^T = 8\pi G T_{\mu\nu}$ (MM 2013)

where $(\square^{-1}\phi)(x) = \int d^4y G(x,y)\phi(y)$

- Two models modifying General Relativity nonlocally – in the infrared
 - ▶ m^2 sets a new reference energy scale
→ Nonlocal terms contributes for $\square_g \ll m^2$ and vice versa
- Phenomenological approach
- Interesting cosmology:
 - ▶ FRW background/linear perturbations
 - ▶ Observational constraints and model comparison

Model RR : $S_{RR} = \frac{1}{16\pi G} \int d^4x \sqrt{-g} [R - \frac{1}{2}m^2 R \frac{1}{\square^2} R]$ (MM, Mancarella 2013)

Application to Cosmology

Model RT

$$G_{\mu\nu} - m^2(g_{\mu\nu}\square_{ret}^{-1}R)^T = 8\pi GT_{\mu\nu}$$

Model RR

$$G_{\mu\nu} - m^2K_{\mu\nu}(\square_{ret}^{-1}R, \square_{ret}^{-2}R) = 8\pi GT_{\mu\nu}$$

- Resolution method: Localisation

$$\square V = R \quad \Rightarrow \quad V = \square^{-1}R + V^{(hom)}$$

- ▷ Auxiliary fields with *vanishing initial conditions*
- ▷ They are not genuine (*freely* propagating) degrees of freedom

$$G_{\mu\nu} + m^2 \left[U g_{\mu\nu} - \frac{1}{2} (\nabla_\mu S_\nu + \nabla_\nu S_\mu) \right] = 8\pi GT_{\mu\nu}$$

$$\square_g U = -R, \quad \partial_\mu U = \frac{1}{2} \nabla_\nu (\nabla_\mu S^\nu + \nabla^\nu S_\mu)$$

$$G_{\mu\nu} - m^2 K_{\mu\nu}(V, S) = 8\pi GT_{\mu\nu}$$

$$\square_g V = R, \quad \square_g S = V$$

- Specialisation to flat FRW

$$ds^2 = -dt^2 + a^2(t)\delta_{ij}dx^i dx^j$$

Background evolution

- Modified Friedmann equations :

$$H^2(t) = 8\pi G \sum_i \bar{\rho}_i(t) + m^2 Y(\{\bar{V}_k\}, H(t))$$

+ auxiliary EoM for $\{\bar{V}_k\}$

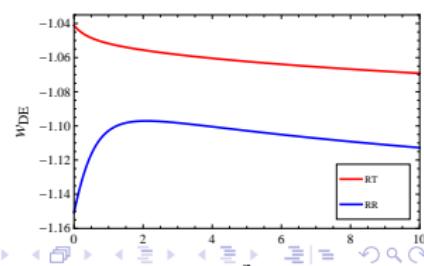
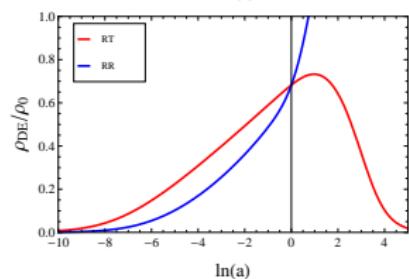
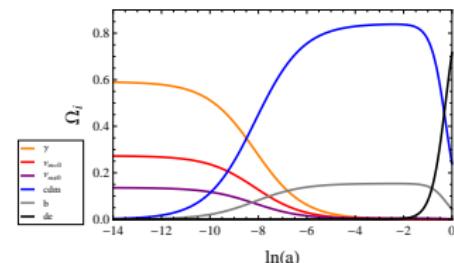
- $m^2 Y \equiv \bar{\rho}_{\text{DE}}(t)$: Dynamical dark energy
 - $\square^{-1} R|_{\text{RD}} = 0$: Late-time effectiveness
 - Flatness today: $m_{\text{RT}} \simeq 0.67 H_0$, $m_{\text{RR}} \simeq 0.28 H_0$
 - From $\dot{\bar{\rho}}_{\text{DE}} = -3H(1 + w_{\text{DE}})\bar{\rho}_{\text{DE}}$

$$\text{Fit : } w(t) = w_0 + (1 - a(t)) w_a$$

RT: $w_0 \approx -1.04$, $w_a \approx -0.02$

RR: $w_0 \approx -1.15$, $w_a \approx 0.08$

→ On the phantom side: $w_{\text{DE}} < -1$

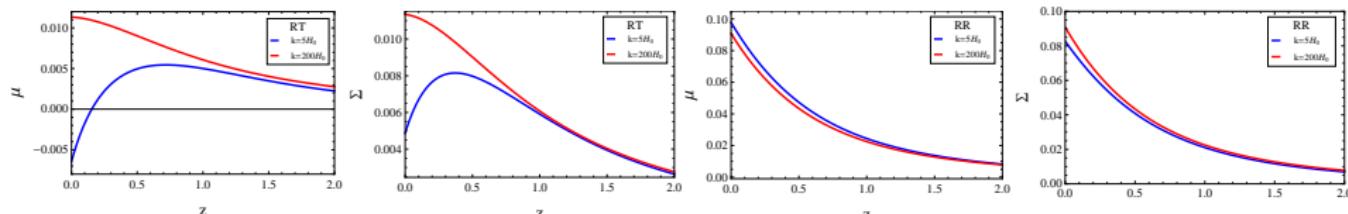


Scalar perturbations and Structure Formation

- Gravitational Ψ and lensing potential $(\Psi - \Phi)$

(YD, Foffa, Khosravi, Kunz, Maggiore 2014)

$$\Psi = [1 + \mu(z, k)] \Psi_{\Lambda\text{CDM}}, \quad (\Psi - \Phi) = [1 + \Sigma(z, k)] (\Psi - \Phi)_{\Lambda\text{CDM}}$$



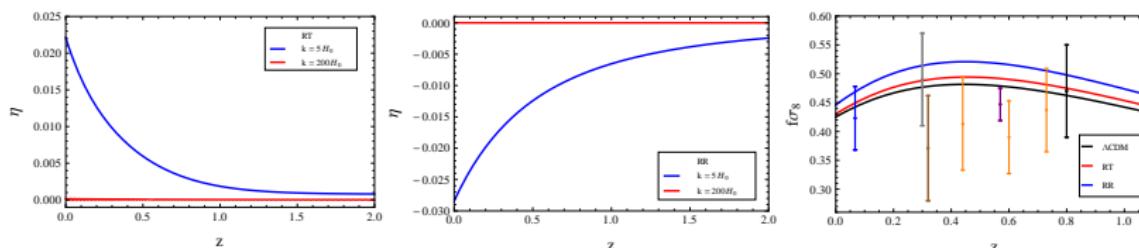
▷ Fit: $\mu(t) = \mu_s a(t)^s$ RT: $\mu_s = 0.01$, $s = 0.8$, RR: $\mu_s = 0.09$, $s = 2$ (EUCLID: $\Delta\mu_s = 0.01$)

- Gravitational slip and RSD

(6dF, SDSS LRG, BOSS LOWZ+CMASS, WiggleZ, VIPERS)

$$\eta = (\Psi + \Phi)/\Phi,$$

$$f \equiv \frac{d \ln D}{d \ln a} \text{ with } D(a) \sim \delta_M(a)$$



▷ Consistency with structure formation

▷ Nonlinear structure formation for RR: N-body simulation

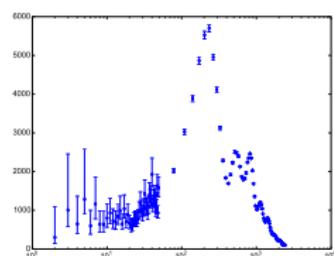
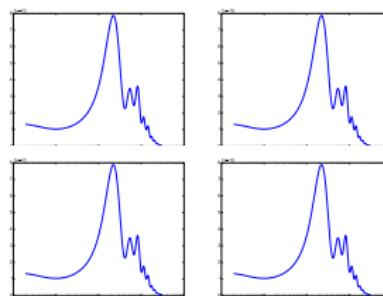
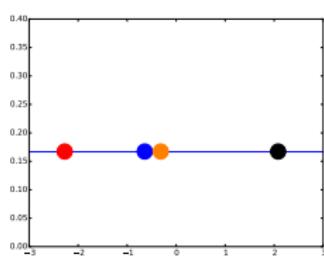
Boltzmann Code and Parameter Inference

- Implementation in CLASS: Computation of CMB and LSS observables
- Observational constraints and model comparison with MONTEPYTHON
(Lesgourges, Audren et al.)
- Cosmological scenario: *Planck* baseline
 - ▷ 6 cosmo parameters varied: $\{\omega_b, \omega_c, H_0, A_s, n_s, z_{reio}\}$
 - ▷ Neutrino: Two massless species $N_{\text{eff}} = 2.03351$, one massive $m_\nu = 0.06\text{eV}$
- Datasets:
 - ▷ CMB: [Planck 2013](#), [Planck 2015](#)
 - ▷ Supernovae: SDSS-II/SNLS3 Joint Light-Curve Analysis (JLA 2014)
 - ▷ BAO: BOSS LOWZ+CMASS DR10&11 ([iso.](#), [aniso.](#)), 6dF and [SDSS MGS](#)
 - ▷ H_0 : HST ($70.6 \pm 3.3, 73.0 \pm 2.4, 73.8 \pm 2.4$)
(YD, Foffa, Kunz, MM, Pettorino, 2014)
(YD, Foffa, Kunz, MM, Pettorino, in prep.)

Observational constraints and parameter inference

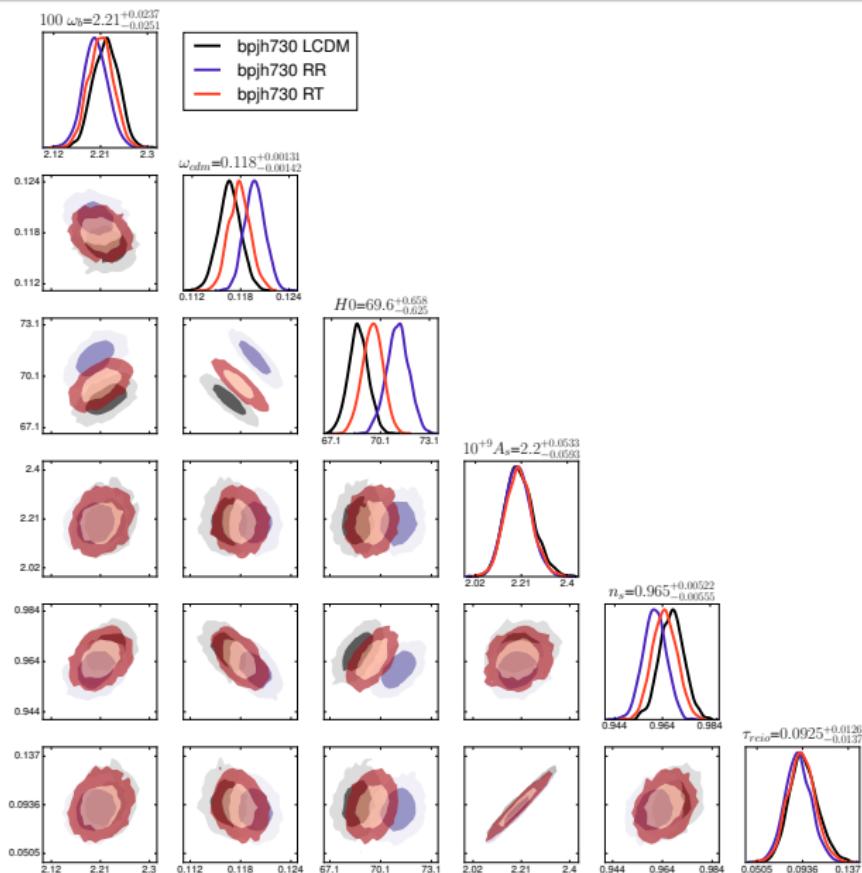
- Bayesian inference:

- Observed datasets: Planck 2013/2015, JLA, BAO, HST, etc
- Statistical models: Λ CDM, RT and RR with $\{\omega_b, \omega_c, H_0, A_s, n_s, z_{reio}\}$
- Parameter estimation: Update our degree of belief through observations



- Minimum χ^2 estimation:

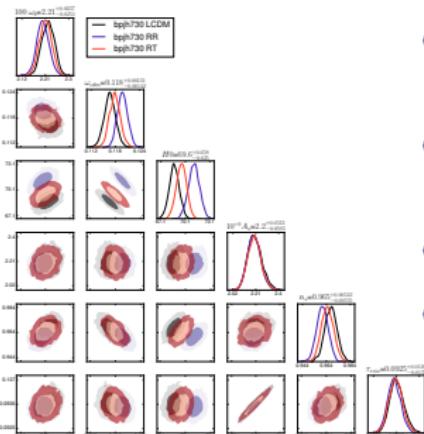
$$\chi^2 = \sum_{\text{dataset}} \chi_{\text{dataset}}^2 \quad \text{with} \quad \chi_{\text{dataset}}^2 = \sum_i \frac{(\theta_{\text{theo}}^i - \theta_{\text{obs}}^i)^2}{(\sigma_{\text{obs}}^i)^2}$$



Observational constraints and parameter inference

Param	Planck			BAO+Planck+JLA		
	Λ CDM	$g_{\mu\nu}\square^{-1}R$	$R\square^{-2}R$	Λ CDM	$g_{\mu\nu}\square^{-1}R$	$R\square^{-2}R$
ω_c	$0.1194^{+0.0027}_{-0.0026}$	$0.1195^{+0.0026}_{-0.0028}$	$0.1191^{+0.0027}_{-0.0028}$	$0.1175^{+0.0015}_{-0.0014}$	$0.1188^{+0.0014}_{-0.0014}$	$0.1204^{+0.0014}_{-0.0013}$
H_0	$67.56^{+1.2}_{-1.3}$	$68.95^{+1.3}_{-1.3}$	$71.67^{+1.5}_{-1.5}$	$68.43^{+0.61}_{-0.69}$	$69.3^{+0.68}_{-0.66}$	$70.94^{+0.74}_{-0.7}$
$\Delta\chi^2_{\min}$	9801.7	9801.3	9800.1	10485.5	10485.0	10488.7

	BAO+Planck+JLA+ $H_0 = 73.0 \pm 2.4$		
Param	Λ CDM	$g_{\mu\nu}\square^{-1}R$	$R\square^{-2}R$
ω_c	$0.117^{+0.0014}_{-0.0014}$	$0.1182^{+0.0013}_{-0.0014}$	$0.1201^{+0.0013}_{-0.0013}$
H_0	$68.72^{+0.61}_{-0.63}$	$69.60^{+0.66}_{-0.63}$	$71.14^{+0.72}_{-0.69}$
$\Delta\chi^2_{\min}$	10488.9	10487.3	10489.3

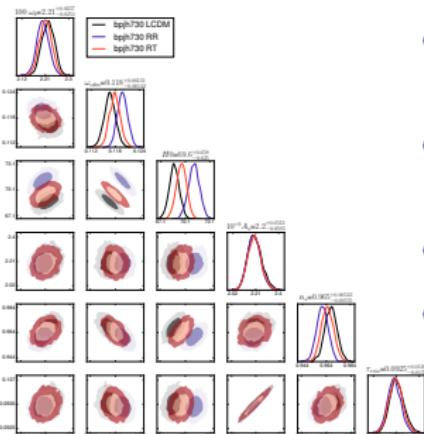


- Few parameters with $\gtrsim 1\sigma$ deviation from Λ CDM
→ Nonlocal models prefer a bigger H_0
- Nonlocal vs Λ CDM: Overall $|\Delta\chi^2| \lesssim 2$
→ Mostly statistically equivalent to Λ CDM
- Planck: RR fits slightly better C_l^{TT} at low- l
- BAO+Planck+JLA: RR creates a Planck-JLA 1σ -tension

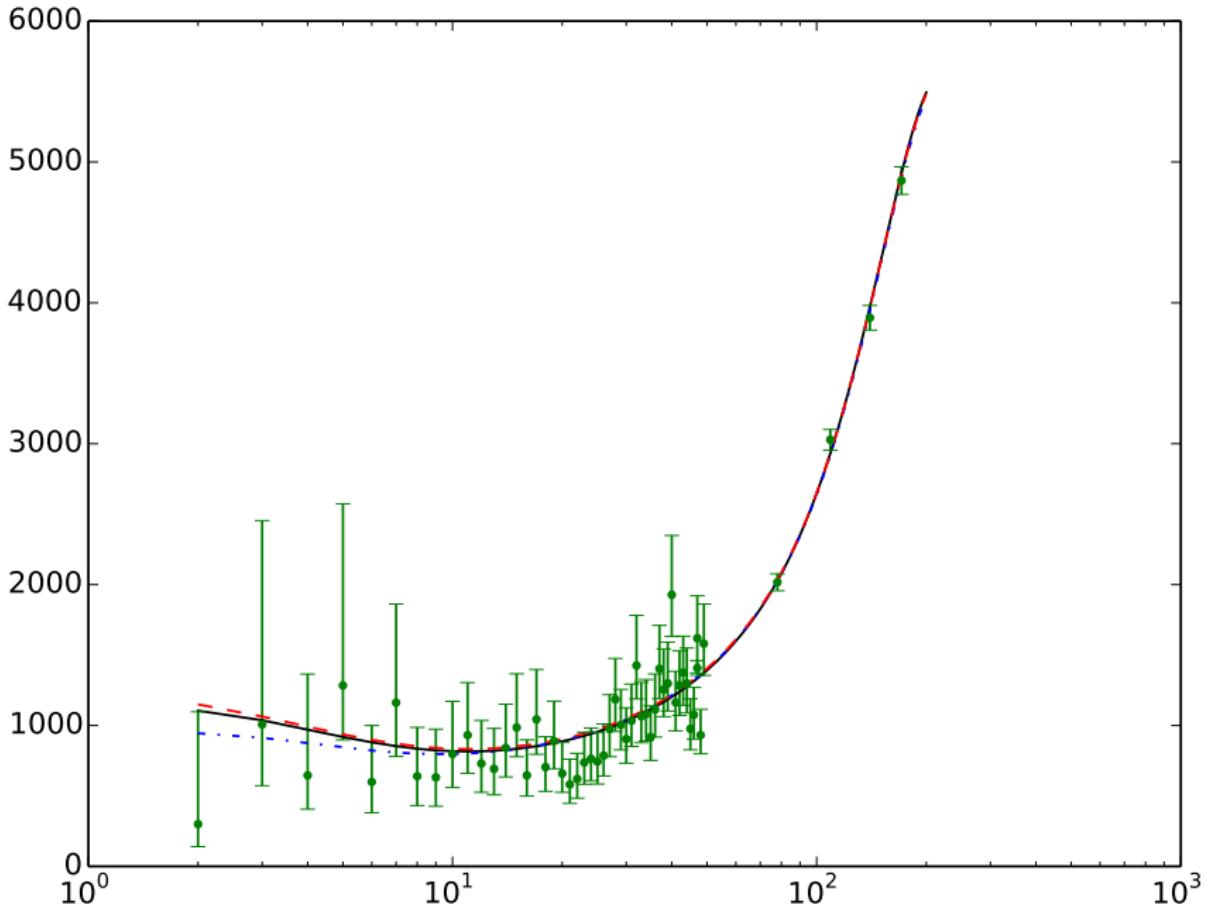
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$\Delta\chi^2_{\min}$	1.6	1.2	0	0.5	0	3.7

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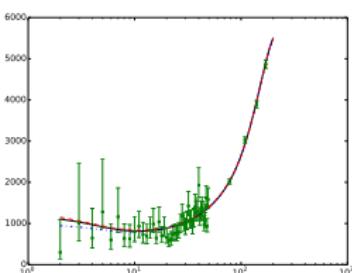
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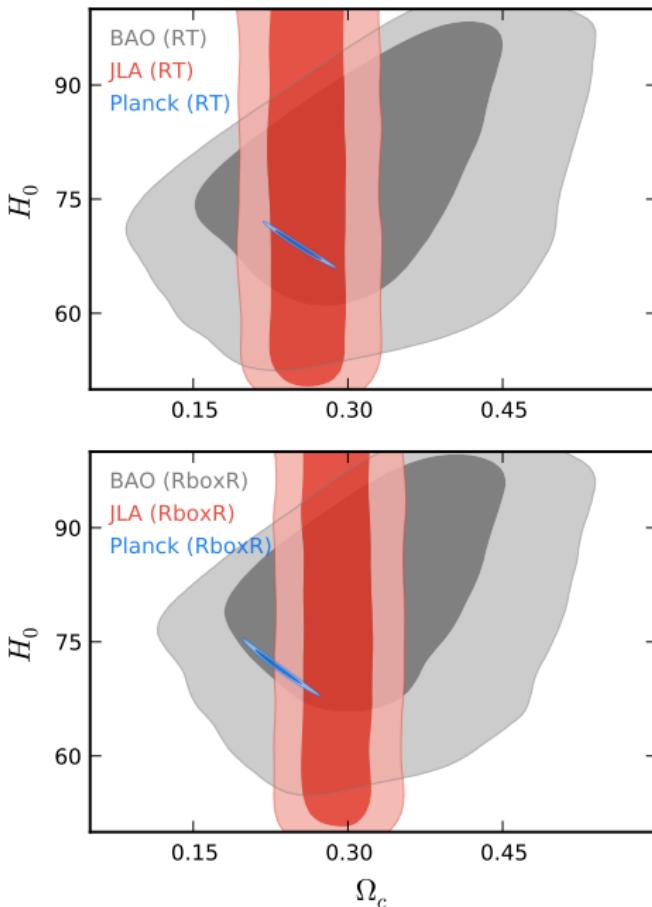
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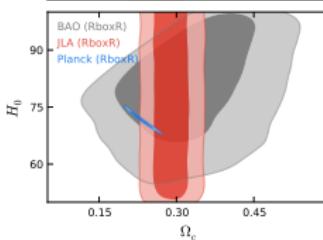
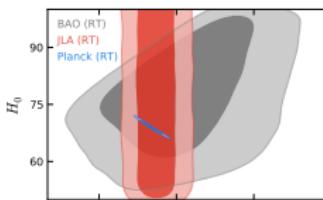
$$\Delta\omega_c = \Delta(\Omega_c h^2) \sim 1\%$$



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Bayesian model selection

- Computation of the Bayes factor: done by considering the nested models

$$G_{\mu\nu} - m^2(g_{\mu\nu}\square_{\text{ret}}^{-1}R)^T - g_{\mu\nu}\Lambda = 8\pi GT_{\mu\nu}$$

$$\mathcal{L} = \frac{1}{16\pi G} [R - 2\Lambda - m^2 R \square^{-2} R] + \mathcal{L}_m$$

with cosmological parameter space $\{\omega_b, H_0, A_s, n_s, z_{\text{reio}}, \Omega_\Lambda, \Omega_{de}\}$

→ Non-informative priors are flat on Ω_Λ and Ω_{de}

- Three statistical models in each case: $\mathcal{M}_{\Lambda+de}$, \mathcal{M}_Λ , \mathcal{M}_{de}
- Bayes theorem

$$P(\theta|d, \mathcal{M}) = \frac{P(d, \mathcal{M}|\theta)P(\theta|\mathcal{M})}{P(d, \mathcal{M})}$$

- Savage-Dickey density ratio:

$$B_{\Lambda/(\Lambda+de)} = \frac{P(d, \mathcal{M}_\Lambda)}{P(d, \mathcal{M}_{\Lambda+de})} = \left. \frac{P(\Omega_{de}|d, \mathcal{M}_{\Lambda+de})}{P(\Omega_{de}|\mathcal{M}_{\Lambda+de})} \right|_{\Omega_{de}=0}$$

→ Model Λ (dis)favored with betting odds $B_{\Lambda/(\Lambda+de)} : 1$ wrt $\Lambda + de$

Conclusion

$$G_{\mu\nu} - m^2(g_{\mu\nu}\square_{\text{ret}}^{-1}R)^T = 8\pi GT_{\mu\nu}$$

$$\mathcal{L} = \frac{1}{16\pi G} [R - m^2 R \square^{-2} R] + \mathcal{L}_m$$

- Two observationally viable models of gravity
(JCAP 1504 (2015) 04, 044, arXiv:1411.7692)
- Phenomenological side
 - ▶ Well behaved dynamical dark energy
 - ▶ Same number of free parameters than Λ CDM
 - ▶ Fit the data as well as Λ CDM
- Provide observationally consistent alternatives to Λ CDM
- Theoretical side: Effective models/terms
 - ▶ Suggest effects/mechanisms for dynamical dark energy generation
 - ▶ Dimensional transmutation, conformal anomaly
(Maggiore 2015)

