

# Remarks on 2PI formalisms and the FRG

## Functional Renormalization

— from quantum gravity and dark energy to ultracold atoms and condensed matter

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Based on work done in collaboration with

J. Pawłowski and U. Reinosa

A paper to appear... soon !

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# Motivations

Study (controlable ?) non perturbative methods in many-body physics and field theory

## Two exact formulae

Thermodynamic potential as a functional of the propagator

$$\Omega[G] = \frac{1}{2} \text{Tr} \log G^{-1} - \frac{1}{2} \text{Tr} \Sigma G + \Phi[G]$$

Flow of the effective action

$$\partial_\kappa \hat{\Gamma}_\kappa[\phi] = \frac{1}{2} \int_q \partial_\kappa R_\kappa(q) G_\kappa(q, -q; \phi)$$

$$G_\kappa^{-1}(q, -q; \phi) = \Gamma_\kappa^{(2)}(q, -q; \phi) + R_\kappa(q)$$

**These formulae are useful mostly for the approximations that they suggest**

**One can use one formalism to shed light on the other (this talk)**

# Present discussion limited to scalar field theory

(can be generalized)

$$S[\varphi] = \int d^d x \left\{ \frac{1}{2} (\partial\varphi(x))^2 + \frac{m^2}{2} \varphi^2(x) + \frac{\lambda}{4!} \varphi^4(x) \right\}$$

**Some representative recent related works (not limited to scalar field)**

**JPB, J. Pawłowski, U. Reinosa (2010)**

**M. Carrington et al (2014)**

**N. Dupuis (2013)**

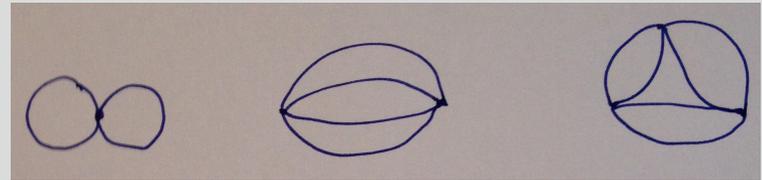
**V. Meden et al (2016)**

# Basics of 2PI formalisms (1)

$$\Omega[G] = \frac{1}{2} \text{Tr} \log G^{-1} - \frac{1}{2} \text{Tr} \Sigma G + \Phi[G]$$

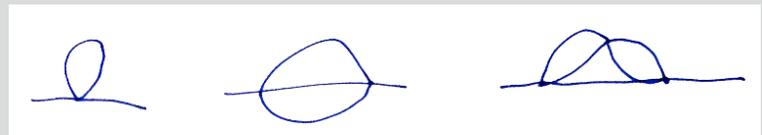
## Luttinger-Ward functional

$$\Phi[G]$$



## Self-energy

$$\Sigma(p) = 2 \frac{\delta \Phi}{\delta G(p)}$$



## Self-consistency condition

$$\Sigma[G] = G^{-1} - G_0^{-1}$$

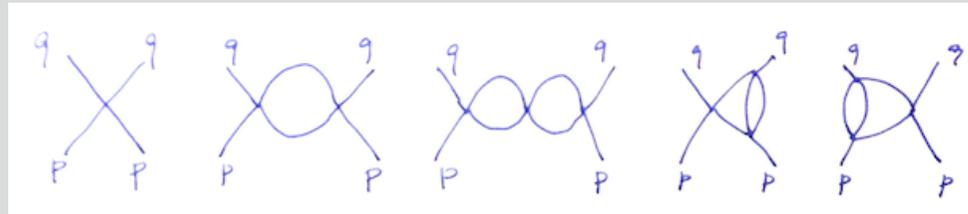
## Stationarity property

$$\left. \frac{\delta \Omega[G]}{\delta G} \right|_{G_0} = 0$$

# Basics of 2PI formalisms (2)

## Irreducible kernel

$$\mathcal{J}(q, p) = 2 \frac{\delta \Sigma(p)}{\delta G(q)} = 4 \frac{\delta^2 \Phi}{\delta G(q) \delta G(p)} = \mathcal{J}(p, q)$$



## Bethe-Salpeter equation

$$\Gamma^{(4)}(q, p) = \mathcal{J}(q, p) - \frac{1}{2} \int_l \Gamma^{(4)}(q, l) G^2(l) \mathcal{J}(l, p)$$

# Basics of functional RG

## Flow equation (Wetterich)

$$\partial_\kappa \hat{\Gamma}_\kappa[\phi] = \frac{1}{2} \int_q \partial_\kappa R_\kappa(q) G_\kappa(q, -q; \phi) = \text{Diagram}$$

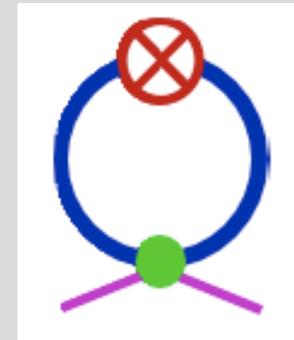
$$G_\kappa^{-1}(q, -q; \phi) = \Gamma_\kappa^{(2)}(q, -q; \phi) + R_\kappa(q)$$

Infinite hierarchy of coupled flow equations for the n-point functions

## Equation for the 2-point function

$$\partial_\kappa \Gamma_\kappa^{(2)}(p) = -\frac{1}{2} \int_q \partial_\kappa R_\kappa(q) G_\kappa^2(q) \Gamma_\kappa^{(4)}(q, p)$$

And so on.....



# The theory in the presence of $R_K(q)$

All formal relations between n-point functions hold for any  $K$

One can then take derivatives w.r.t.  $K$

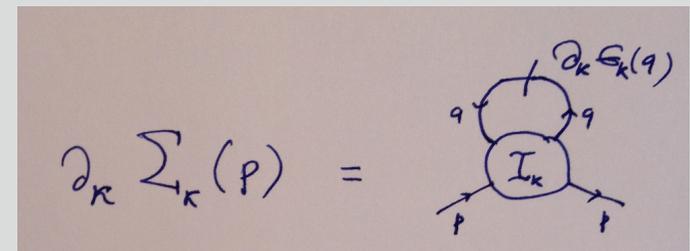
.... thereby obtaining flow equations

Equation for the 2-point function (or self-energy)

$$G_K^{-1}(p) = p^2 + m^2 + \Sigma_K(p) + R_K(p)$$

$$\partial_K \Sigma_K(p) = 2 \int_q \partial_K G_K(q) \frac{\delta^2 \Phi[G]}{\delta G(q) \delta G(p)} \Big|_{G_K} = \frac{1}{2} \int_q \partial_K G_K(q) \mathcal{J}_K(q, p)$$

This is NOT quite the usual flow equation



The diagram shows the derivative of the self-energy,  $\partial_K \Sigma_K(p)$ , represented as a diagrammatic equation. On the left is the expression  $\partial_K \Sigma_K(p)$ . On the right is an equals sign followed by a diagram. The diagram consists of a central circle labeled  $\mathcal{I}_K$ . Two external lines with arrows pointing outwards are labeled  $p$ . Two internal lines with arrows pointing inwards are labeled  $q$ . A vertical line segment connects the two  $q$  lines, and this segment is labeled  $\partial_K \Sigma_K(q)$ .

Solving the Bethe-Salpeter equation to get  $\Gamma^{(4)}(q, p)$

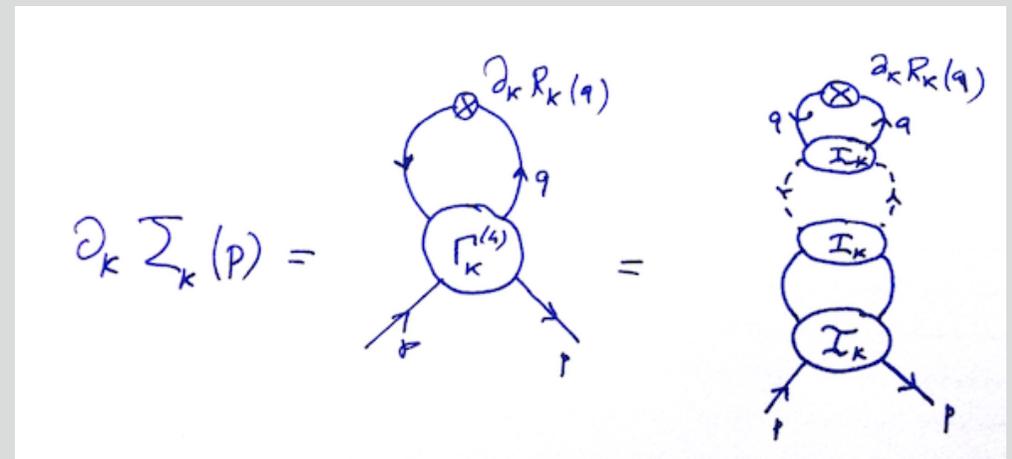
$$\Gamma_K^{(4)}(q, p) = \mathcal{J}_K(q, p) - \frac{1}{2} \int_l \Gamma_K^{(4)}(q, l) G_K^2(l) \mathcal{J}_K(l, p)$$

and using this equation to eliminate  $\mathcal{J}_K(q, p)$

we are left with

$$\partial_K \Sigma_K(p) = -\frac{1}{2} \int_q \partial_K R_K(q) G_K^2(q) \Gamma_K^{(4)}(q, p)$$

**The exact flow equation  
for the 2-point function**



# A possible truncation scheme (1)

**Truncate the Luttinger-Ward functional (keeping selected skeletons)**

**Obtain the kernel** 
$$\mathcal{J}(q, p) = 4 \frac{\delta^2 \Phi}{\delta G(q) \delta G(p)}$$

**Then solve the coupled equations**

$$\Gamma_K^{(4)}(q, p) = \mathcal{J}_K(q, p) - \frac{1}{2} \int_l \Gamma_K^{(4)}(q, l) G_K^2(l) \mathcal{J}_K(l, p)$$

$$\partial_K \Sigma_K(p) = -\frac{1}{2} \int_q \partial_K R_K(q) G_K^2(q) \Gamma_K^{(4)}(q, p)$$

**NB. i) The solution is independent of the choice of the "regulator"**

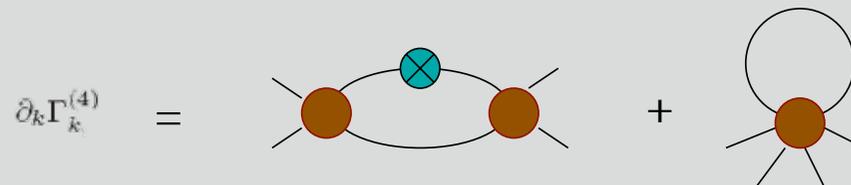
**ii) Not only a truncation of fRG, but an alternative to solving the 2PI equations**

# A possible truncation scheme (2)

Instead of solving the Bethe-Salpeter eqn., write a flow equation for the 4-point function

$$\begin{aligned}
 \partial_\kappa \Gamma_\kappa^{(4)}(p, q) = \partial_\kappa \mathcal{I}_\kappa(p, q) & - \frac{1}{2} \int_l \Gamma_\kappa^{(4)}(p, l) \partial_\kappa G_\kappa^2(l) \Gamma_\kappa^{(4)}(l, q) \\
 & - \frac{1}{2} \int_l \partial_\kappa \mathcal{I}_\kappa(p, l) G_\kappa^2(l) \Gamma_\kappa^{(4)}(l, q) \\
 & - \frac{1}{2} \int_l \Gamma_\kappa^{(4)}(p, l) G_\kappa^2(l) \partial_\kappa \mathcal{I}_\kappa(l, q) \\
 & + \frac{1}{4} \int_l \int_s \Gamma_\kappa^{(4)}(p, l) G_\kappa^2(l) \partial_\kappa \mathcal{I}_\kappa(l, s) G_\kappa^2(s) \Gamma_\kappa^{(4)}(s, q)
 \end{aligned}$$

**NB.** This equation is NOT the "usual" flow equation for the 4-point function



# Renormalization issues

**Not a priori obvious that the integrals are finite**

$$\begin{aligned}\partial_\kappa \Gamma_\kappa^{(4)}(p, q) &= \partial_\kappa \mathcal{I}_\kappa(p, q) - \frac{1}{2} \int_l \Gamma_\kappa^{(4)}(p, l) \partial_\kappa G_\kappa^2(l) \Gamma_\kappa^{(4)}(l, q) \\ &- \frac{1}{2} \int_l \partial_\kappa \mathcal{I}_\kappa(p, l) G_\kappa^2(l) \Gamma_\kappa^{(4)}(l, q) \\ &- \frac{1}{2} \int_l \Gamma_\kappa^{(4)}(p, l) G_\kappa^2(l) \partial_\kappa \mathcal{I}_\kappa(l, q) \\ &+ \frac{1}{4} \int_l \int_s \Gamma_\kappa^{(4)}(p, l) G_\kappa^2(l) \partial_\kappa \mathcal{I}_\kappa(l, s) G_\kappa^2(s) \Gamma_\kappa^{(4)}(s, q)\end{aligned}$$

**Standard lore in fRg: things become "simple" at the "cutoff scale"**

$$\Gamma_\kappa^{(2)}(p) \equiv Z_\kappa p^2 + m_\kappa^2, \quad Z_\kappa \sim \ln \kappa, \quad m_\kappa^2 \sim \kappa^2 \quad \Gamma_\kappa^{(4)} \sim \ln \kappa, \quad \Gamma_\kappa^{(n>4)} \sim \kappa^{4-n}$$

**One expects of course similar features in the 2PI truncation...**

**...but working out the "details" turned out to be tricky**

# Divergences, and subdivergences.....

Consider the loop expansion of the 4-point function

$$\Gamma^{(0)} = \mathcal{I}^{(0)} = \text{X}$$

$$\Gamma^{(1)} = \mathcal{I}^{(1)} + \mathcal{I}^{(0)} G^2 \Gamma^{(0)} = \text{X} \text{ with loop} + \text{X with bubble}$$

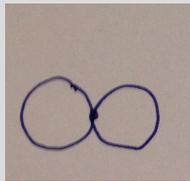
$$\Gamma^{(2)} = \mathcal{I}^{(2)} + \mathcal{I}^{(1)} G^2 \Gamma^{(0)} + \mathcal{I}^{(0)} G^2 \Gamma^{(1)}$$

$$= \text{X with 2 loops} + \text{X with subdivergence} + \text{X with subdivergence} + \text{X with subdivergence} + \text{X with subdivergence} + \text{X with subdivergence}$$

In addition to counterterms needed to renormalise the kernel  $\mathcal{I}$

an infinite number of counterterms are needed to renormalise the BS equation....

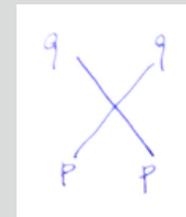
# A simple example (1)



$$\Phi[G]$$



$$\Sigma(p) = 2 \frac{\delta\Phi}{\delta G(p)}$$



$$\mathcal{J}_K(q, p)$$

## Standard 2PI renormalization

### Gap equation

$$\bar{\Sigma} = \delta m^2 + \frac{\lambda + \delta\lambda_{(0)}^{\text{BS}}}{2} \int_q G(q) \quad G(q) = \frac{1}{q^2 + m^2 + \bar{\Sigma}}$$

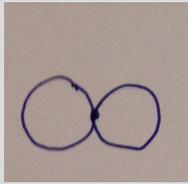
### BS equation

$$\frac{1}{\Gamma_{(0)}^{\text{BS}}} = \frac{1}{\lambda + \lambda_{(0)}^{\text{BS}}} + \frac{1}{2} \int_q G^2(q)$$

### Counterterms

$$\delta\lambda_{(0)}^{\text{BS}} = \lambda \frac{a\lambda}{1 - a\lambda}, \quad a \equiv \frac{1}{2} \int_q G^2(q) \quad \frac{\delta m^2}{m^2} = \frac{a\lambda}{1 - a\lambda} = \frac{\delta\lambda}{\lambda}$$

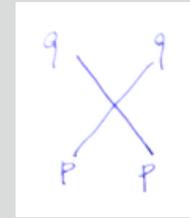
# A simple example (2)



$$\Phi[G]$$



$$\Sigma(p) = 2 \frac{\delta \Phi}{\delta G(p)}$$



$$\mathcal{J}_\kappa(q, p)$$

The two equations to be solved

$$\partial_\kappa m_\kappa^2 = -\frac{1}{2} \Gamma_\kappa \int_q (\partial_\kappa R_\kappa) G_\kappa^2(q) \quad \partial_\kappa \Gamma_\kappa = -\frac{1}{2} \Gamma_\kappa^2 \int_q \partial_\kappa G_\kappa^2(q)$$

**Solution**

$$m_\kappa^2 = m^2 = m_\Lambda^2 + \frac{\Gamma_\Lambda}{2} \int_q \left\{ G_\kappa(q) - G_\Lambda(q) + (m_\kappa^2 - m_\Lambda^2) G_\Lambda^2(q) \right\}$$

**Elimination of "subdivergences " is automatically taken care of by the coupled flow equations**

# Conclusions

- Two non perturbative methods were compared
- Approximation schemes exist where they completely match
- The comparison help to clarify some renormalisation issues in non perturbative schemes, such as 2PI
- Truncating the fRG flow equations with 2PI relations may be useful in some applications