

# 2-, 3- and 4-point functions in 2, 3 and 4 dimensions



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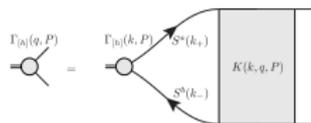
666. WE-Heraeus-Seminar - From correlation functions to QCD phenomenology  
Bad Honnef, Germany

April 3, 2018



# Hadronic bound states

Bound state equations:



Ingredients:

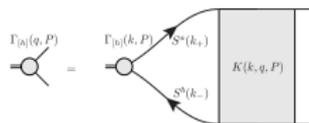
- Interaction kernel  $K$

- Quark propagator  $S$

$$S(p)^{-1} = S_0(p)^{-1} + \gamma_\mu \int \frac{D_{\mu\nu}(p-q)}{S(q)} \Gamma_\nu(p, q)$$

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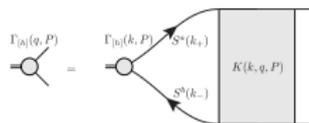
Approaches:

- Phenomenological:  
Model interactions

- From first principles:  
Piecing together the pieces

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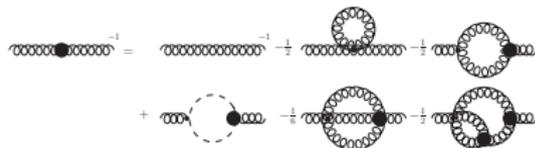
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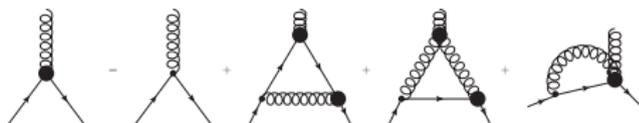
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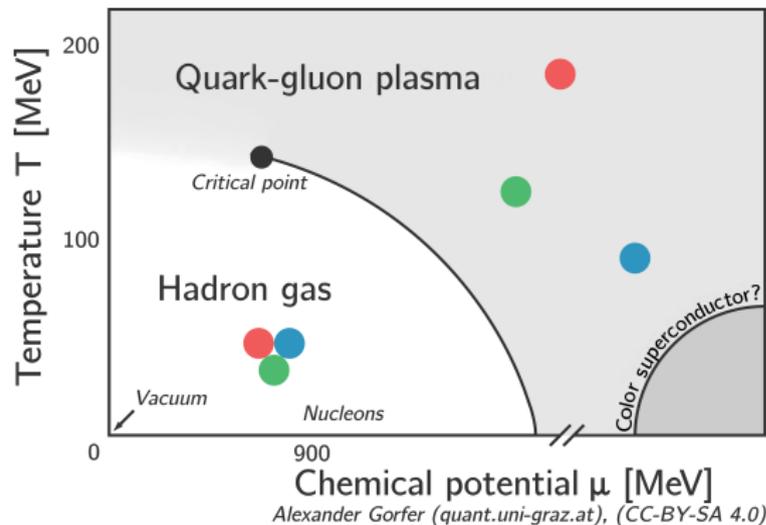


→ Couples to infinity of equations.

# QCD phase diagram

## Questions:

- Phases and transitions between them, critical point
- Experimental signatures



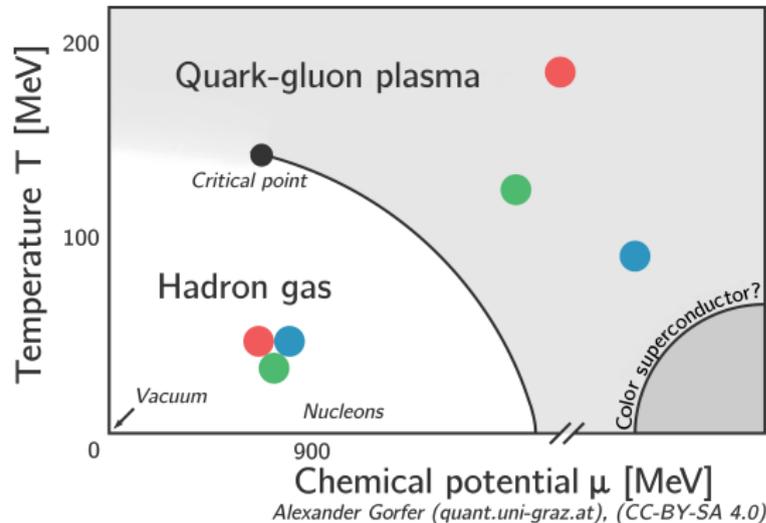
## Theoretical challenges:

- Model description
- Mathematical, e.g., complex action for lattice QCD
- Complexity, e.g., truncations of function eqs.
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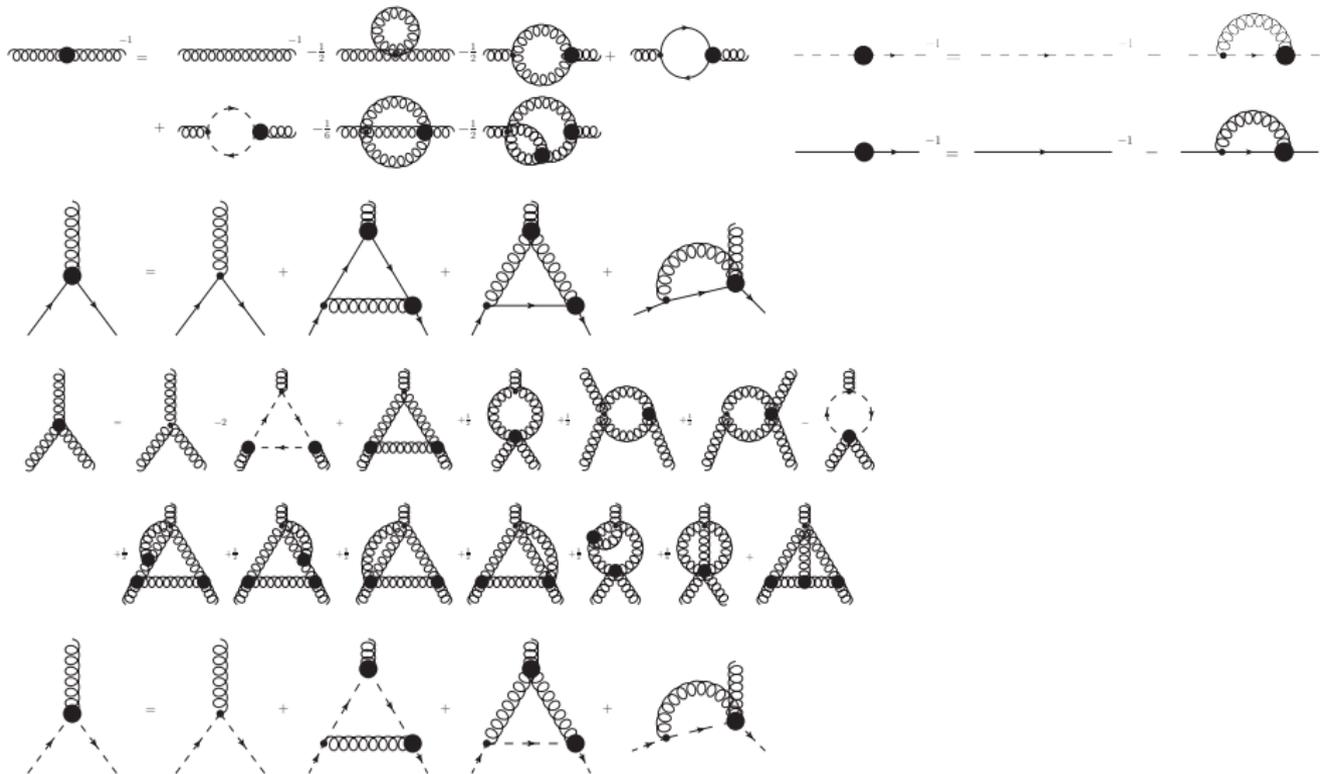
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- How to realize **resummation**?  
higher loop contributions?
- Equivalence between different functional methods?  
FRG, DSEs, nPI, Hamiltonian approach

# Dyson-Schwinger equations



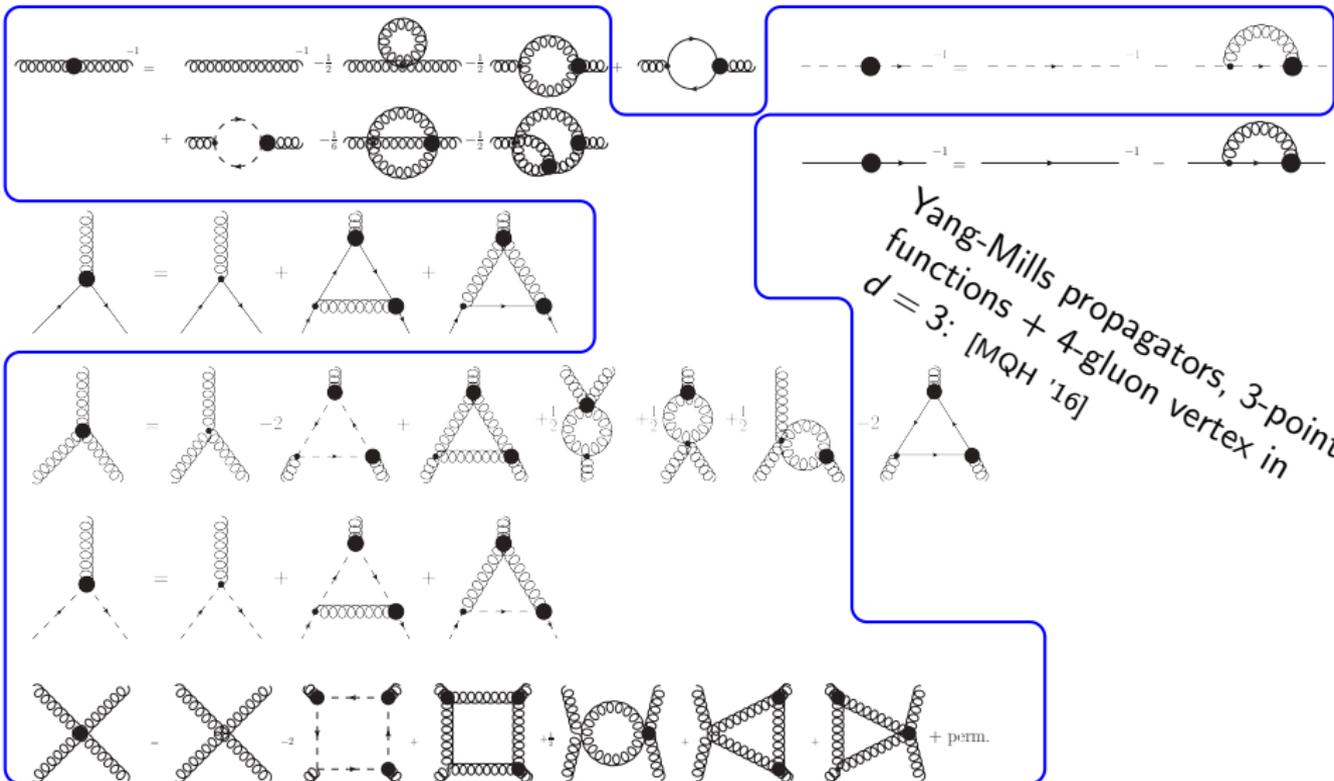
# Coupled systems of Dyson-Schwinger equations

$$\begin{aligned}
 \text{Quark Propagator}^{-1} &= \text{Quark Propagator}^{-1} + \frac{1}{2} \text{Loop (Ghost)}^{-1} + \frac{1}{2} \text{Loop (Quark)}^{-1} + \text{Loop (Ghost-Quark)}^{-1} \\
 \text{Ghost Propagator}^{-1} &= \text{Ghost Propagator}^{-1} + \text{Loop (Ghost)}^{-1} + \text{Loop (Quark)}^{-1}
 \end{aligned}$$

$$\begin{aligned}
 \text{3-point (Quark-Ghost)} &= \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \text{Diagram 4} \\
 \text{3-point (Quark-Quark)} &= \text{Diagram 1} - 2 \text{Diagram 2} + \text{Diagram 3} + \frac{1}{2} \text{Diagram 4} + \frac{1}{2} \text{Diagram 5} + \frac{1}{2} \text{Diagram 6} - 2 \text{Diagram 7} \\
 \text{3-point (Quark-Ghost)} &= \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \text{Diagram 4}
 \end{aligned}$$

quark propagator + 3-point functions: [Williams, Fischer, Heupel '15]  $\rightarrow$  application to bound states

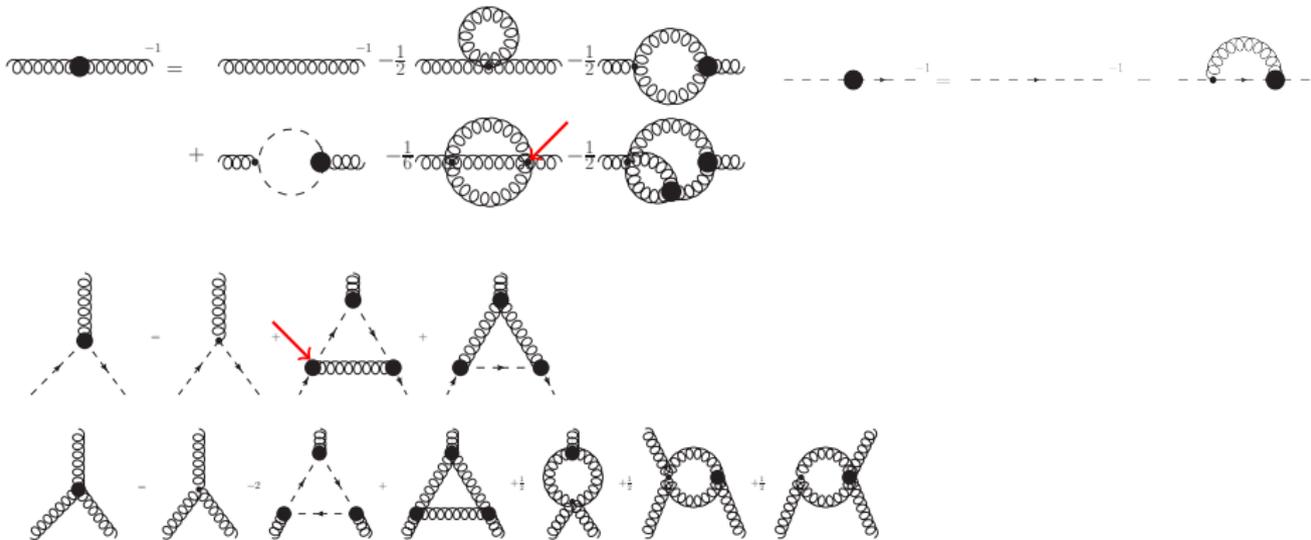
# Coupled systems of Dyson-Schwinger equations





# 3PI system of equations

Three-loop expansion of PI effective action [Berges '04]:



# UV behavior of the gluon propagator

Resummed **one-loop** order: anomalous dimension  $\gamma = -13/22$

$$\left(1 + \frac{\alpha(s)11N_c}{12\pi} \ln \frac{p^2}{s}\right)^\gamma$$

One-loop **anomalous dimension**

Origin in resummation of higher order diagrams.

However, one-loop truncation discards some terms.

$$\begin{aligned}
 & \text{Gluon line with vertex}^{-1} = \text{Gluon line with vertex}^{-1} - \frac{1}{2} \text{Gluon line with gluon loop}^{-1} - \frac{1}{2} \text{Gluon line with ghost loop}^{-1} + \text{Gluon line with ghost loop} + \text{Gluon line with ghost loop and gluon line}
 \end{aligned}$$

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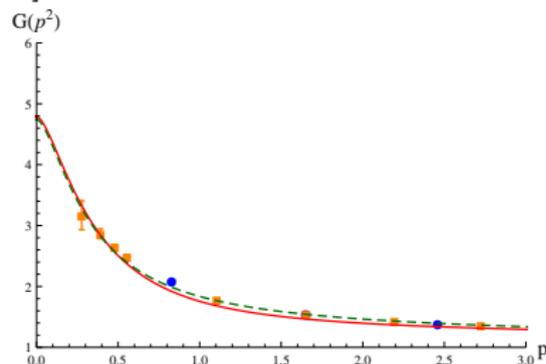
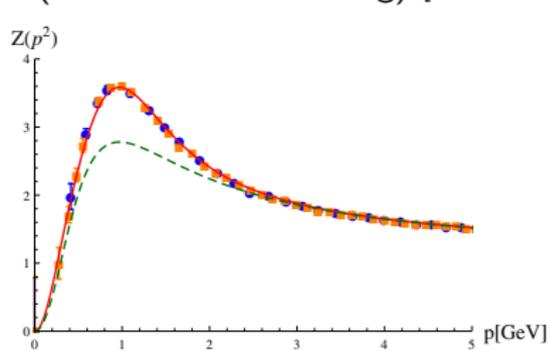
However, one-loop truncation discards some terms.

The diagrammatic equation shows the resummation of higher-order diagrams for the gluon propagator. The left side is a gluon line with a black dot vertex. The right side is a sum of terms: a gluon line with a black dot vertex, minus half of a gluon line with a gluon loop, minus half of a gluon line with a ghost loop, plus a gluon line with a ghost loop and a gluon line with a ghost loop and a gluon line.

→ Puts constraints on UV behavior of vertices [von Smekal, Hauck, Alkofer '97].  
Way out: Include in models (for now).

# Propagators and ghost-gluon vertex with three-gluon vertex model

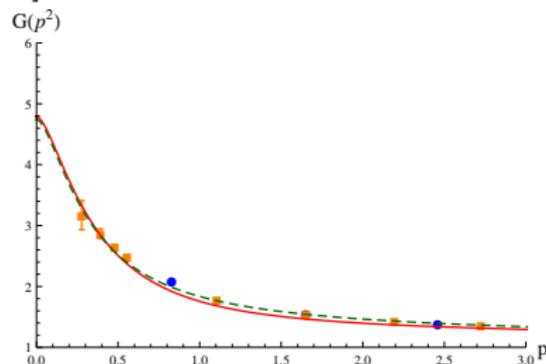
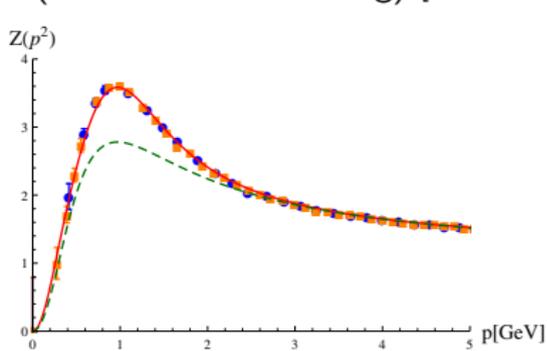
One-loop truncation of gluon propagator with an optimized effective model (contains zero crossing) [MQH, von Smekal '13]:



Good quantitative agreement for ghost and gluon dressings.

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$$\mathcal{L} = -\frac{1}{2} \text{Tr}(F_{\mu\nu} F^{\mu\nu}) + \sum_j \bar{\psi}_j [i \gamma^\mu D_\mu - m_j] \psi_j$$

$$\text{WOBEL} \quad F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu + ig[A_\mu, A_\nu]$$

$$\text{UND} \quad D_\mu = \partial_\mu + igA_\mu$$

# Resummed behavior

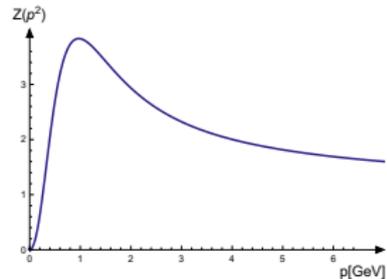
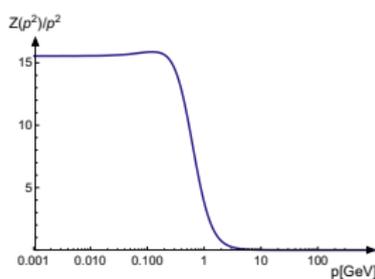
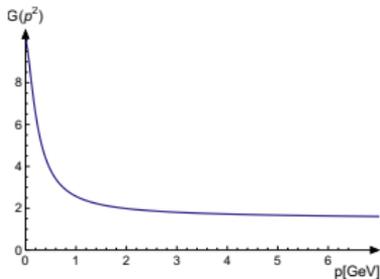
Minimal requirements to obtain one-loop resummed behavior:

- Squint diagram
- Correct anomalous dimensions of three-point functions
- Correct renormalization (constants)

# Resummed behavior

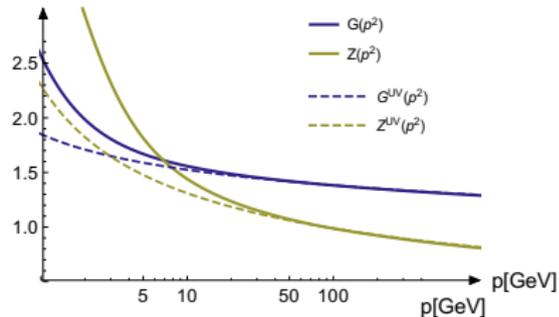
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[propagator+ghost-gluon eqs. full, 3-gluon vertex model, bare 4-gluon vertex]

- Resummed behavior is recovered [MQH '17].



# Lower dimensional Yang-Mills theories as testing ground

## Advantages:

- UV finite: no renormalization, no anomalous running
- Spurious divergences easier to handle
- UV behavior 'easier':  $\propto \frac{g^2}{p^n}$  instead of resummed logarithm

→ Many complications from  $d = 4$  absent.

→ Disentanglement of UV easier.

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Historically interesting because cheaper on the lattice → easier to reach the IR.

# Vertices in two dimensions

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Aspects of two-dimensional Yang-Mills theory:

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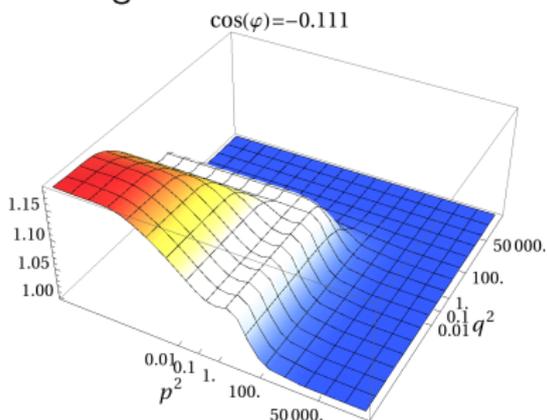
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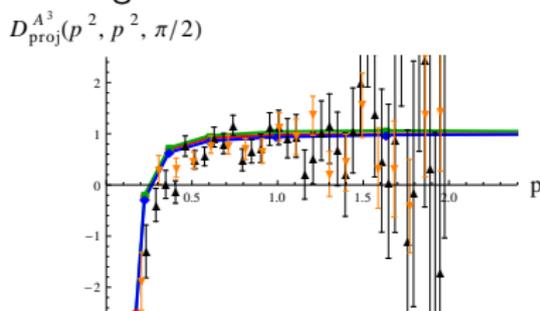
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Ghost-gluon vertex:



[MQH, Maas, von Smekal '12]

Three-gluon vertex:



[Maas '07; MQH, Maas, von Smekal '12]

# Three dimensions

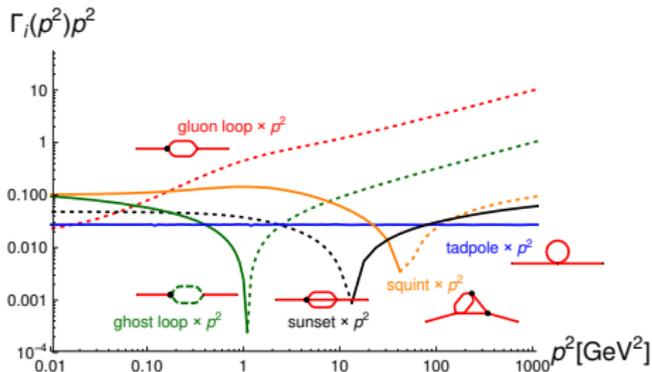
- Four-point functions numerically cheaper.
- Perturbation theory works.

## Continuum results:

- Coupled propagator DSEs: [Maas, Wambach, Grüter, Alkofer '04]
- (R)GZ: [Dudal, Gracey, Sorella, Vandersickel, Verschelde '08]
- DSEs of PT-BFM: [Aguilar, Binosi, Papavassiliou '10]
- YM + mass term: [Tissier, Wschebor '10, '11]

# Gluon propagator: Single diagrams

$$\begin{array}{c}
 \text{---} \bullet \text{---}^{-1} = \text{---} \text{---}^{-1} - \frac{1}{2} \text{---} \text{---} \text{---}^{-1} - \frac{1}{2} \text{---} \text{---} \text{---}^{-1} + \text{---} \text{---} \text{---}^{-1} - \frac{1}{6} \text{---} \text{---} \text{---} \text{---}^{-1} - \frac{1}{2} \text{---} \text{---} \text{---} \text{---}^{-1}
 \end{array}$$



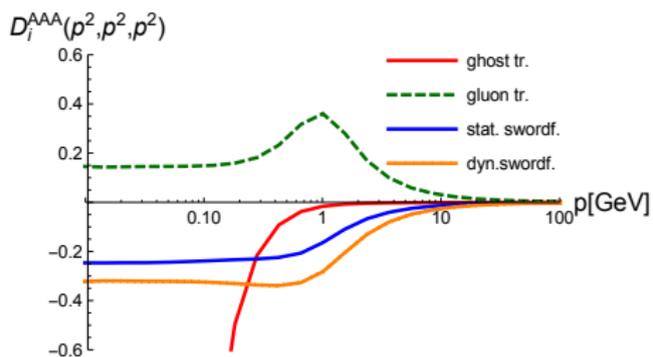
[MQH '16]

→ Clear **hierarchies** identified.

- UV: as expected perturbatively
- non-perturbative: squint important, sunset small  
( $d=4$ :  
[Mader, Alkofer '13; Meyers, Swanson '14])

# Cancellations in gluonic vertices

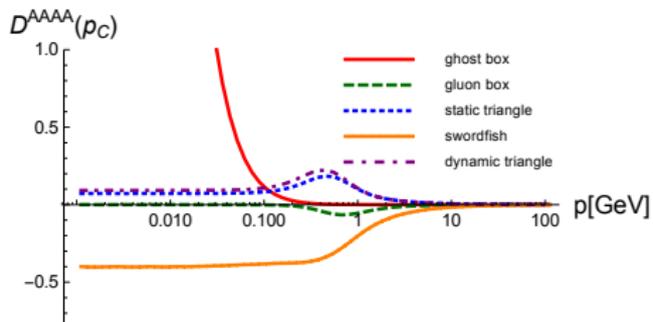
## Three-gluon vertex:



- Individual contributions large.
- **Sum is small!**

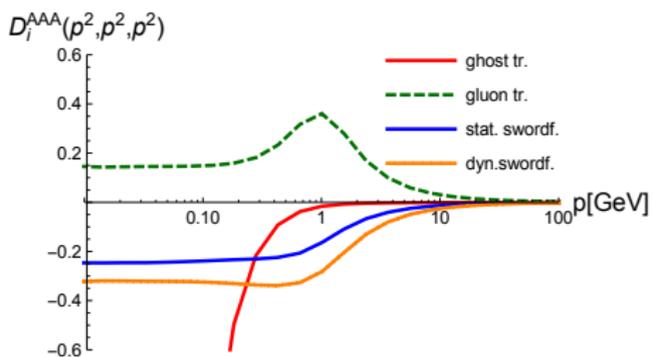
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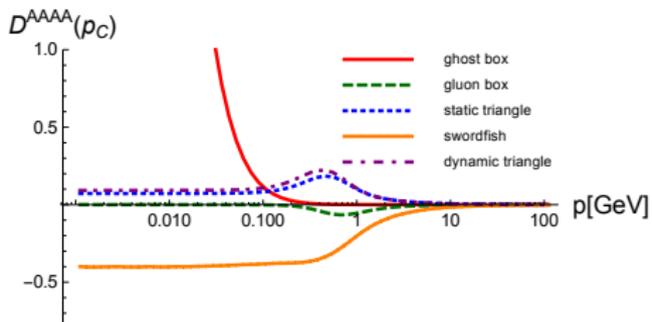


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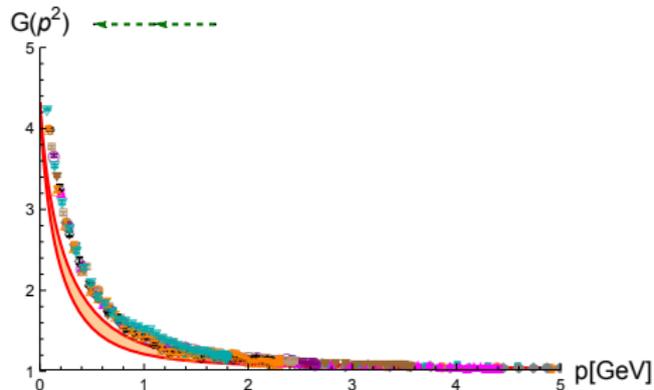
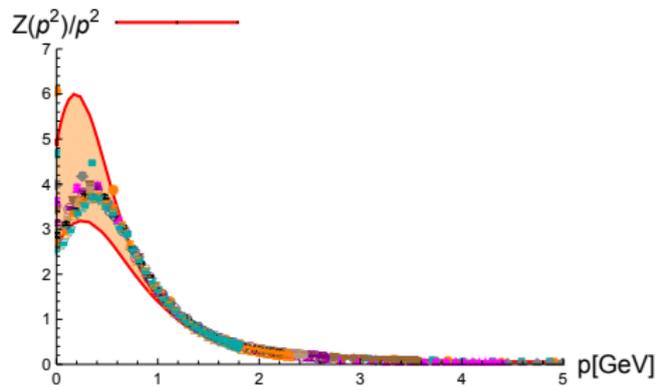
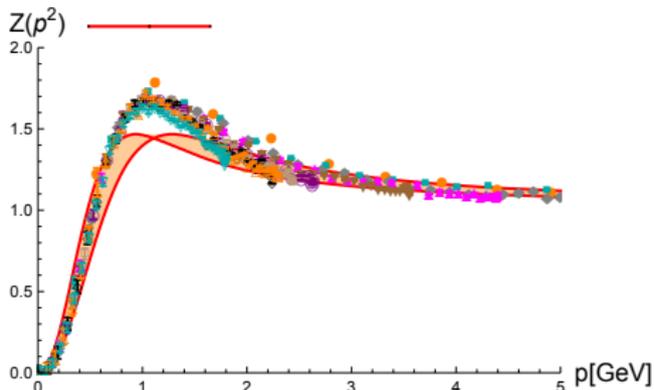
## Four-gluon vertex:



## Higher contributions:

- Higher vertices close to 'tree-level'?  
→ Small.
- If pattern changes (higher vertices large): cancellations required.

# Results: Propagators

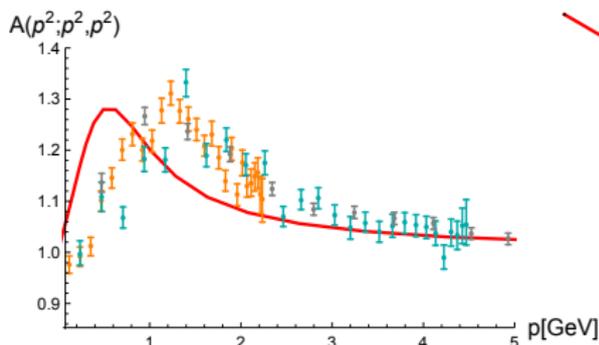


Bands from uncertainty in setting the physical scale.

[MQH '16; lattice: Maas '14]

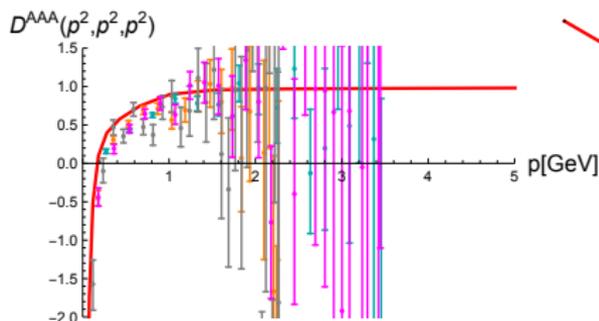
# Comparison of three-point functions with lattice results

## Dressings:



[MQH '16; lattice: Maas, unpublished]

- Deviation from tree-level 'small'
- Position of maximum shifted (as observed with other continuum methods for  $SU(2)$ , e.g., [Pelaez, Matthieu, Wschebor '13; Cyrol, Fister, Mitter, Pawłowski, Strodthoff '15; Corell '18]).

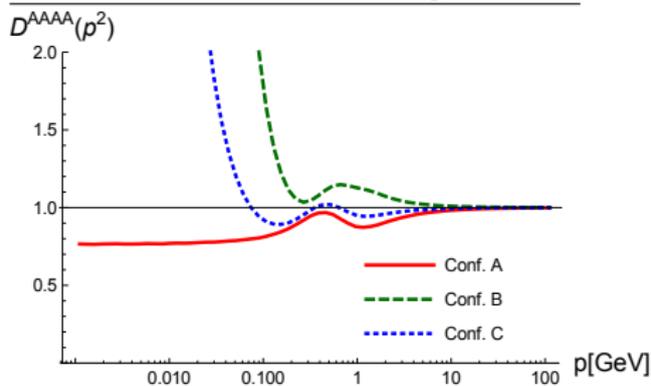


[MQH '16; lattice: Cucchieri, Maas, Mendes '08]

- Close to tree-level above 1 GeV
- Good agreement with lattice data.
- Linear IR divergence [Pelaez, Tissier, Wschebor '13; Aguilar et al. '13]

# Four-gluon vertex

Different momentum configurations:

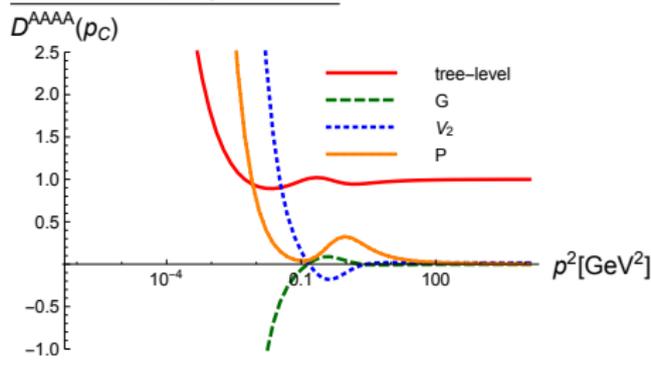


[MQH '16]

Four-gluon vertex:

- Close to tree-level down to 1 GeV

Different projections:



# Solution from the 3PI effective action

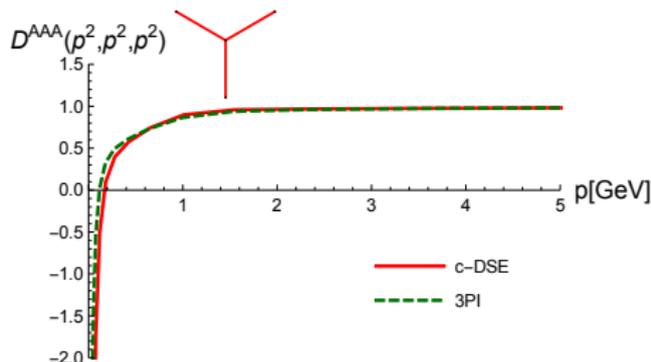
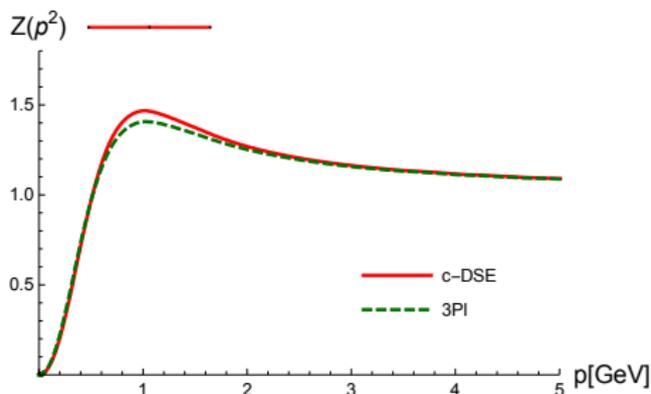
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Equations of motion from 3PI effective action (at three-loop level)

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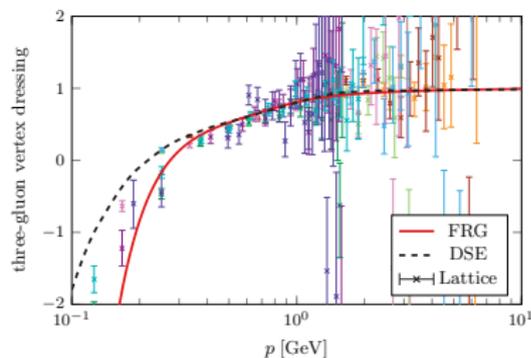
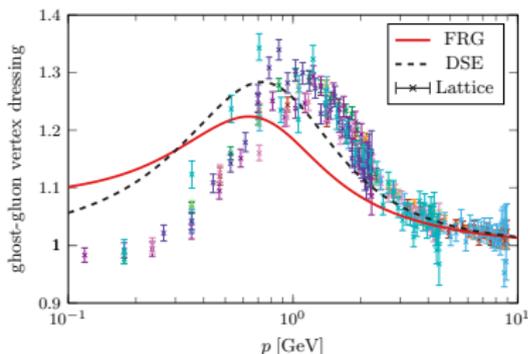
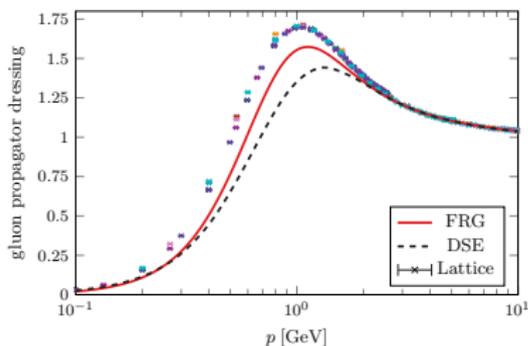


→ Very similar results.

[MQH '16]

# Solutions from the FRG

FRG calculations by Corell, Cyrol, Mitter, Pawłowski, Strodthoff,  
arXiv:1803.10092



NB: Scaling (FRG) and decoupling (DSEs)

- FRG has 'additional' diagrams (tadpoles).
- Equivalence of truncations not trivial.

[Cucchieri, Maas, Mendes '08; MQH '16; Corell et al. '18; Maas, unpublished]

# Conclusions from three dimensions

- **Hierarchy** of correlation functions and diagrams
- **Cancellations** lead to small deviations from the perturbative behavior above 2 GeV.
- Some degree of **stability** (but no complete list of checks done) when
  - varying *system* of equations.
  - varying *equations* of system.
- Discrepancies with lattice results:
  - Nonperturbative gauge fixing?
  - Missing diagrams for vertices?
  - Incomplete tensor bases for some vertices?

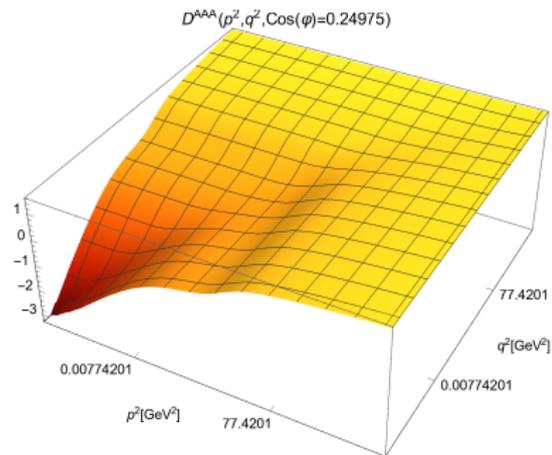
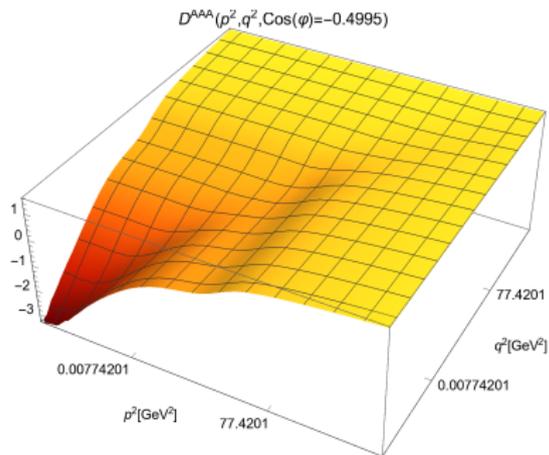
# Extending truncations

Various ways to extend truncations:

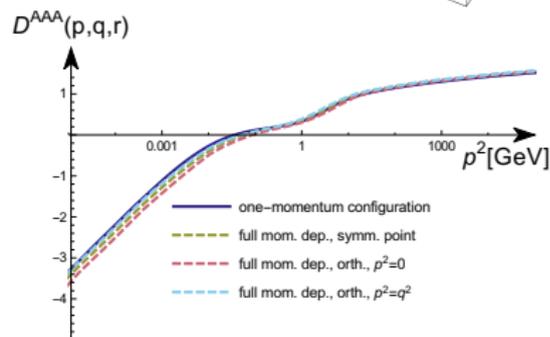
- Vertex tensors beyond tree-level
- Neglected diagrams
- Neglected correlation functions

Extensions also test the previous truncations!

# Three-gluon vertex: Kinematic dependence



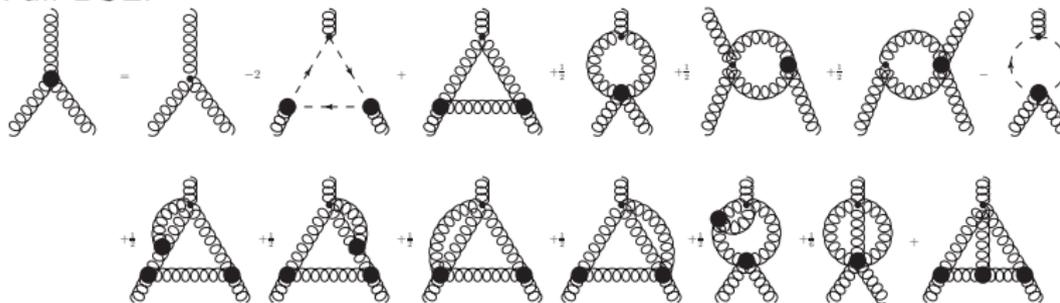
- Kinematic dependence weak.
- In the following: **One-momentum approximation**



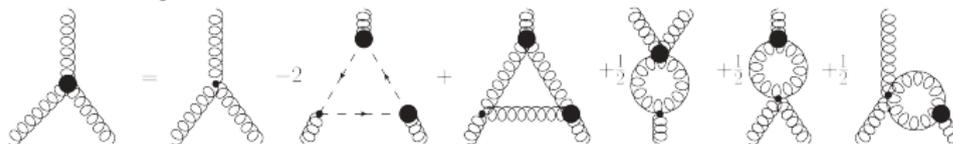


# Three-gluon vertex DSE

Full DSE:

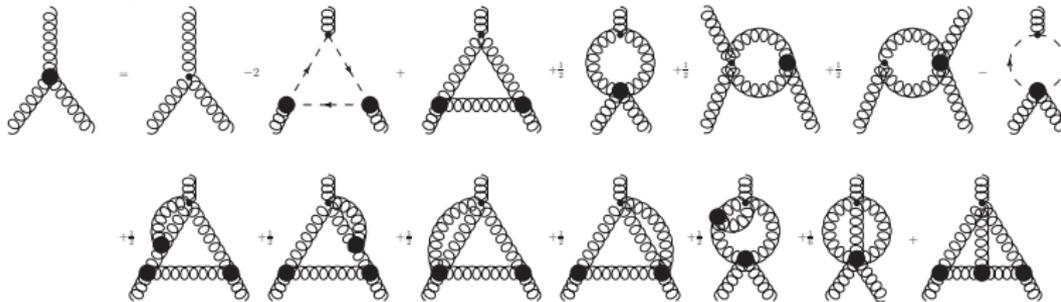


Perturbative one-loop truncation [Blum, MQH, Mitter von Smeikal '14; Eichmann, Alkofer, Vujanovic '14]:

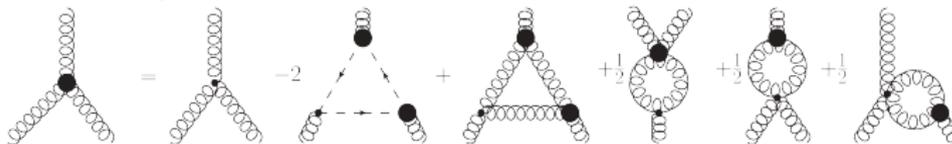


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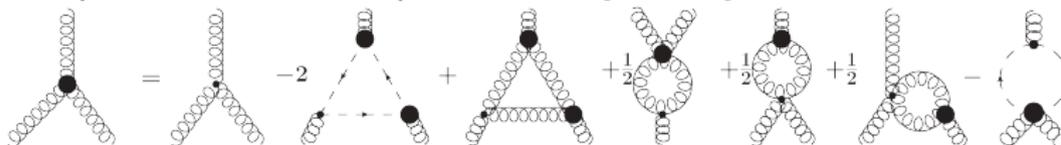
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Non-perturbative one-loop truncation [MQH '17]:



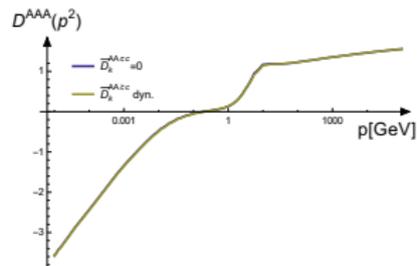
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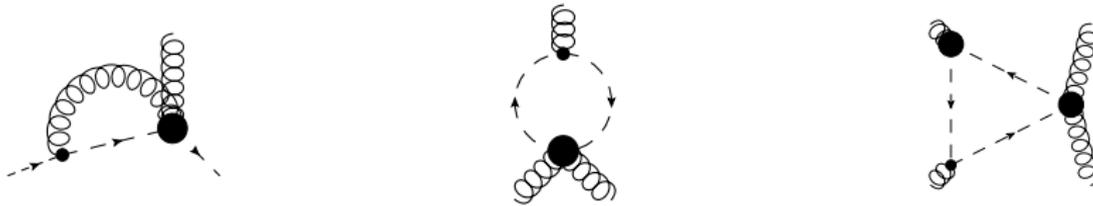
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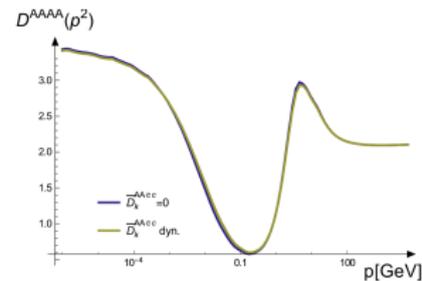
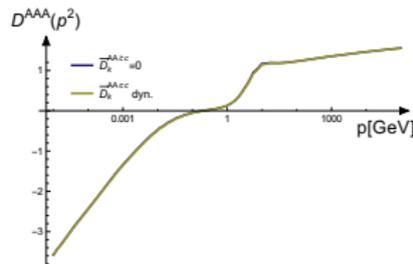
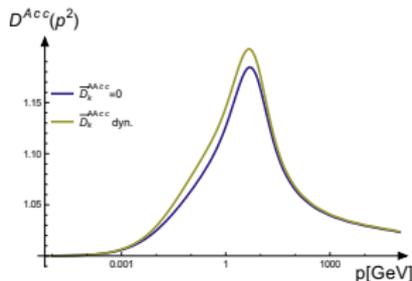
Coupled system of ghost-gluon, three-gluon and four-gluon vertices **with and without** two-ghost-two-gluon vertex [MQH '17]:



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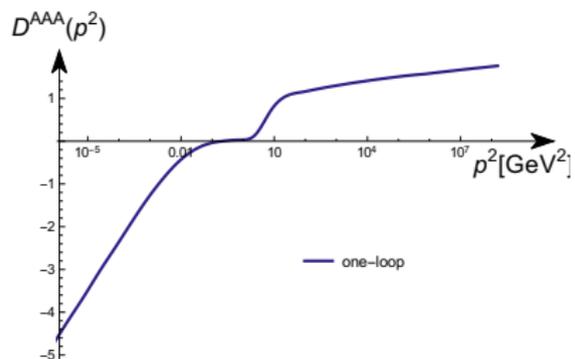


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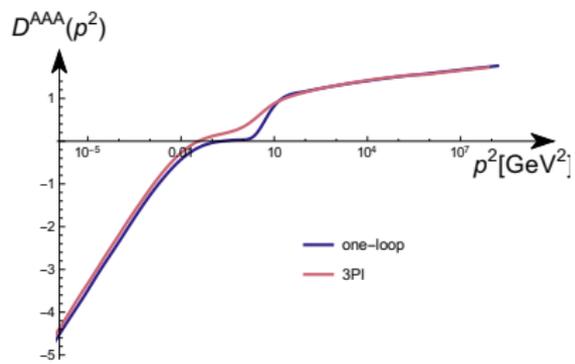


- **Small** influence on ghost-gluon vertex ( $< 1.7\%$ )
- **Negligible** influence on three- and four-gluon vertices.

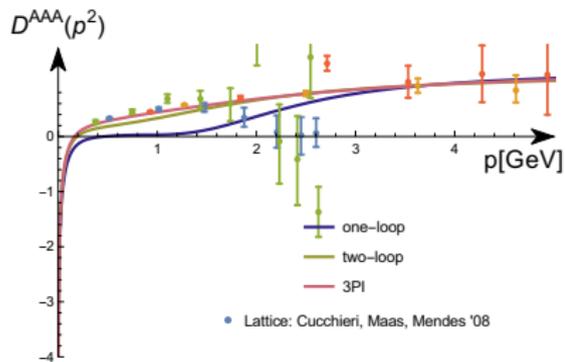
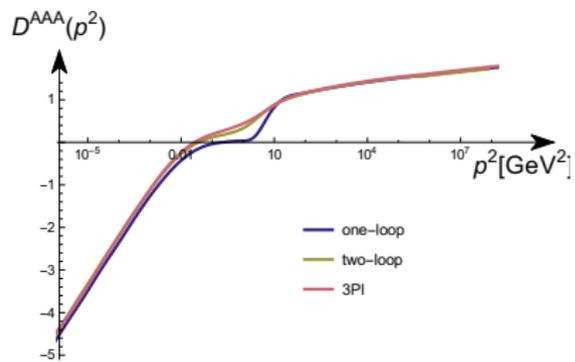
# Three-gluon vertex results



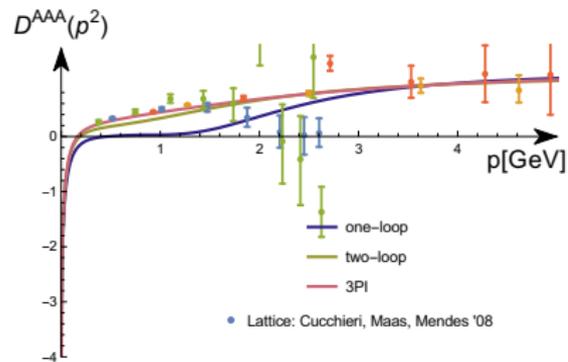
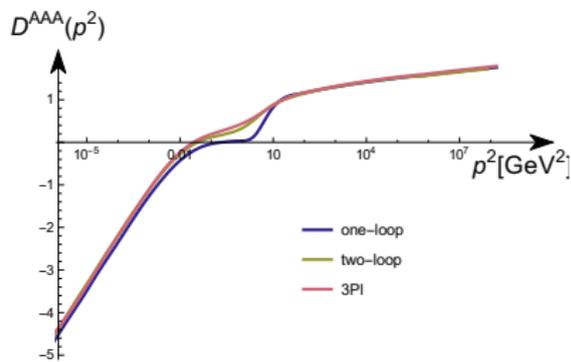
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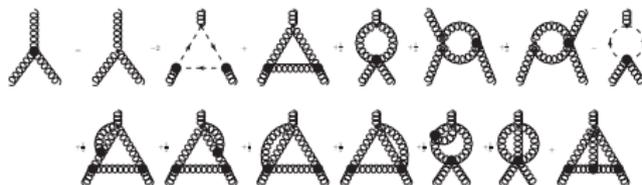
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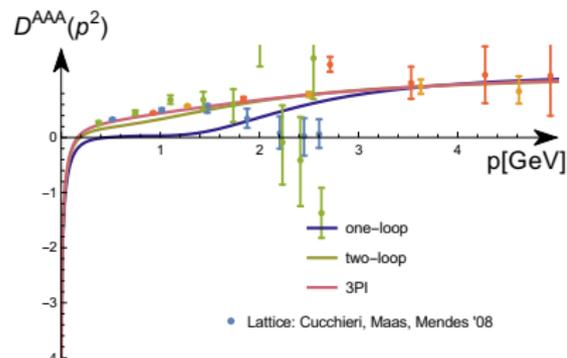
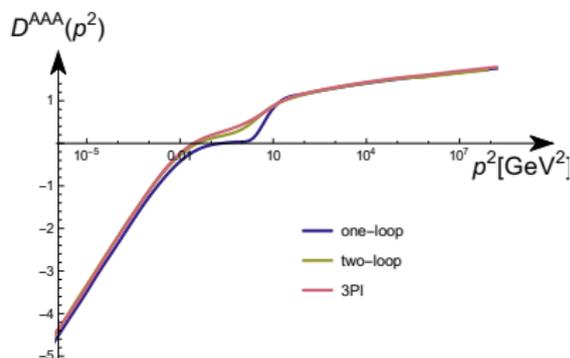
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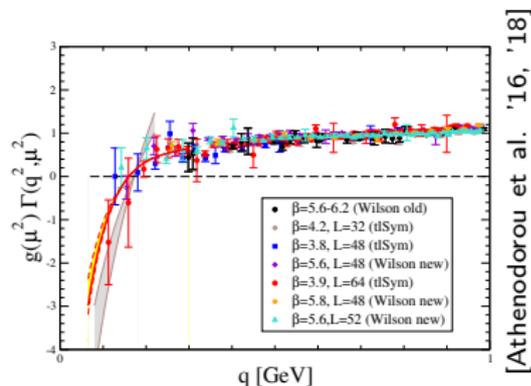
- **Two-loop truncation:** All diagrams except the one with a five-point function.



# Three-gluon vertex results



- Difference between two-loop DSE and 3PI smaller than lattice error.
- Resolves ambiguity in zero crossing due to RG improvement [Blum et al. '14; Eichmann et al. '14; Williams et al. '16]
- Zero crossing in agreement with other approaches, e.g., [Pelaez et al. '13; Aguilar et al. '13; Athenodorou et al. '16; Duarte et al. '16; Sternbeck et al. '17]



[Athenodorou et al. '16, '18]

# Four-point functions: Color space

15 possibilities:

$\delta\delta$  : 3 combinations

$ff$  : 3 combinations

$dd$  : 3 combinations

$df$  : 6 combinations

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 [Pascual, Tarrach '80].

$SU(3)$ :  $\{\sigma_1, \dots, \sigma_8\}$  chosen with these symmetries:

	$\sigma_1$	$\sigma_2$	$\sigma_3$	$\sigma_4$	$\sigma_5$	$\sigma_6$	$\sigma_7$	$\sigma_8$
$a \leftrightarrow b$	+	+	+	-	-	-	-	+
$c \leftrightarrow d$	+	+	+	-	-	+	-	-

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## Four-ghost vertex



$$\Gamma^{\bar{c}\bar{c}cc,abcd}(p, q, r, s) = g^4 \sum_{k=1}^8 \sigma^{k,abcd} E_k^{\bar{c}\bar{c}cc}(p, q, r, s).$$

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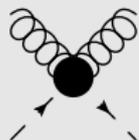
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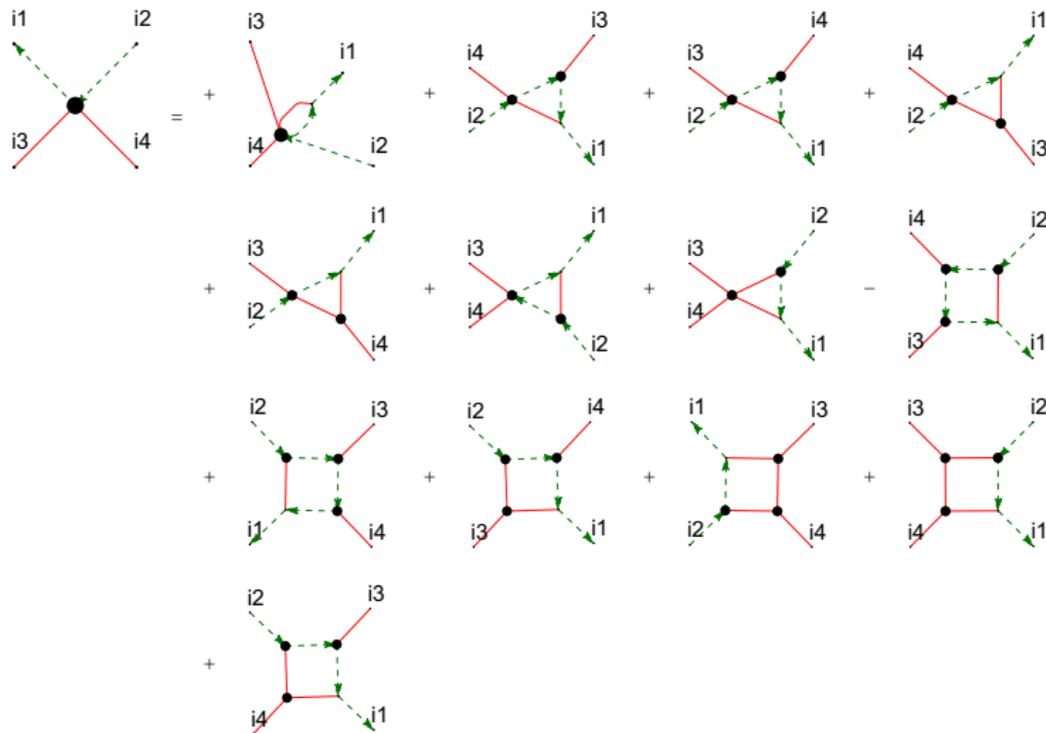
$$\Gamma_{\mu\nu}^{AA\bar{c}c,abcd}(p, q; r, s) = g^4 \sum_{k=1}^{40} \rho_{\mu\nu}^{k,abcd} D_{k(i,j)}^{AA\bar{c}c}(p, q; r, s)$$

with

$$\rho_{\mu\nu}^{k,abcd} = \sigma_i^{abcd} \tau_{\mu\nu}^j, \quad k = k(i, j) = 5(i - 1) + j$$

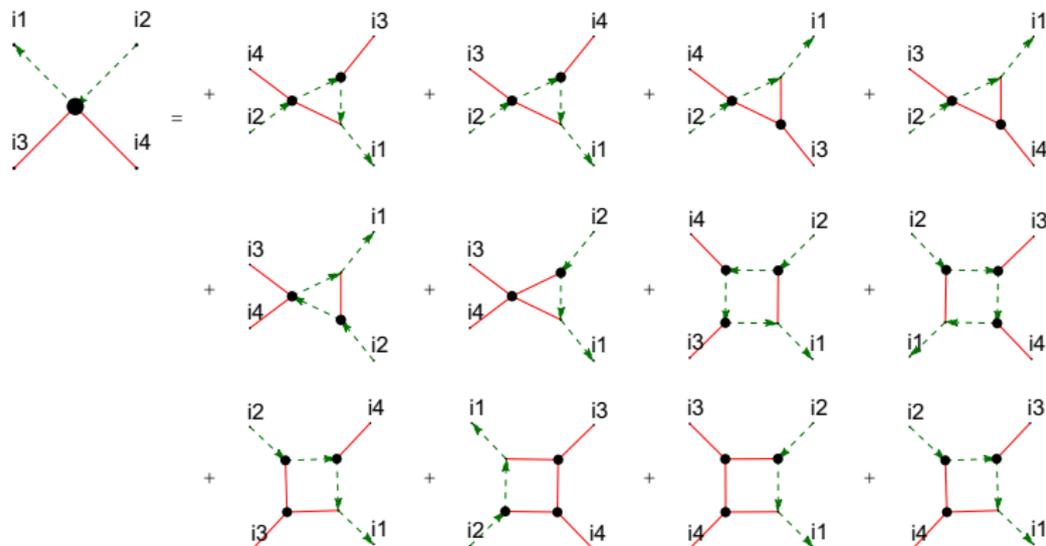
# The two-ghost-two-gluon vertex DSE

2 DSEs, choose the one with the ghost leg attached to the bare vertex  
 → Truncation discards only one diagram.



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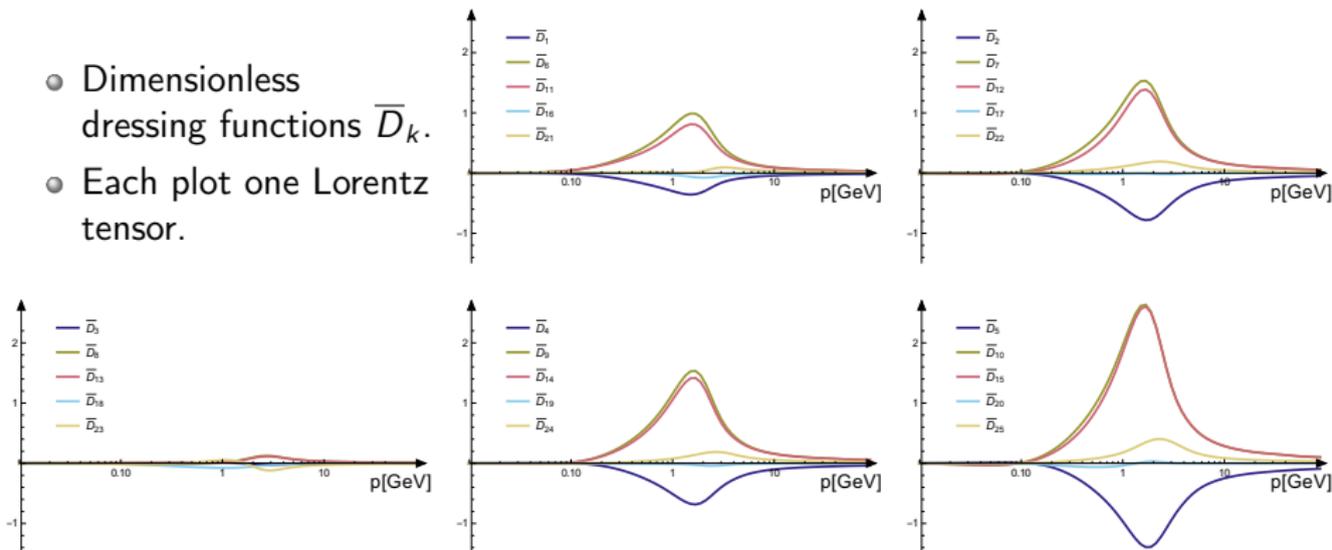
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# Results for the two-ghost-two-gluon vertex

Kinematic approximation: one-momentum configuration

- Dimensionless dressing functions  $\bar{D}_k$ .
- Each plot one Lorentz tensor.



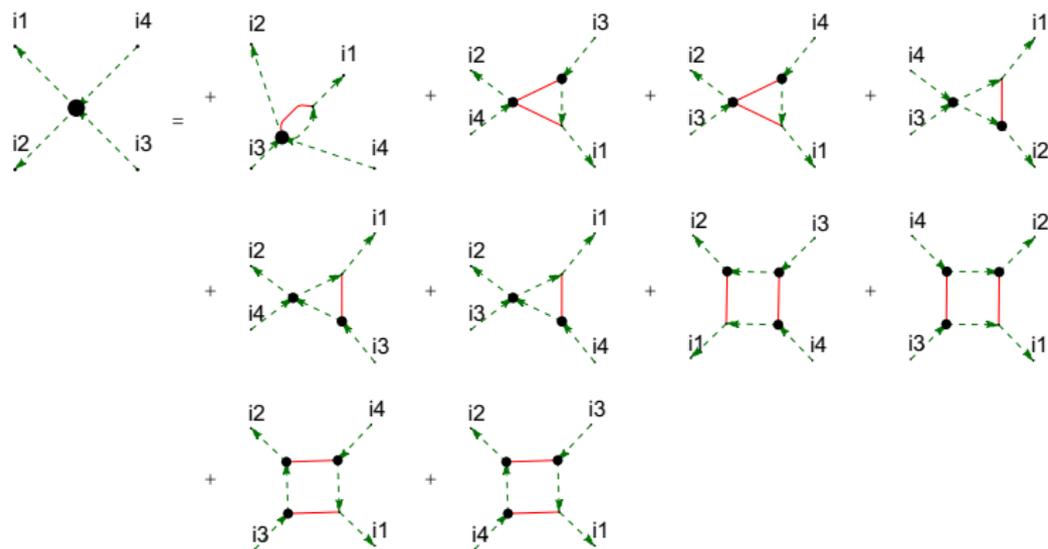
→ Two classes of dressings: 13 very small, 12 not small

→ No nonzero solution for  $\{\sigma_6, \sigma_7, \sigma_8\}$  found.

[MQH '17]

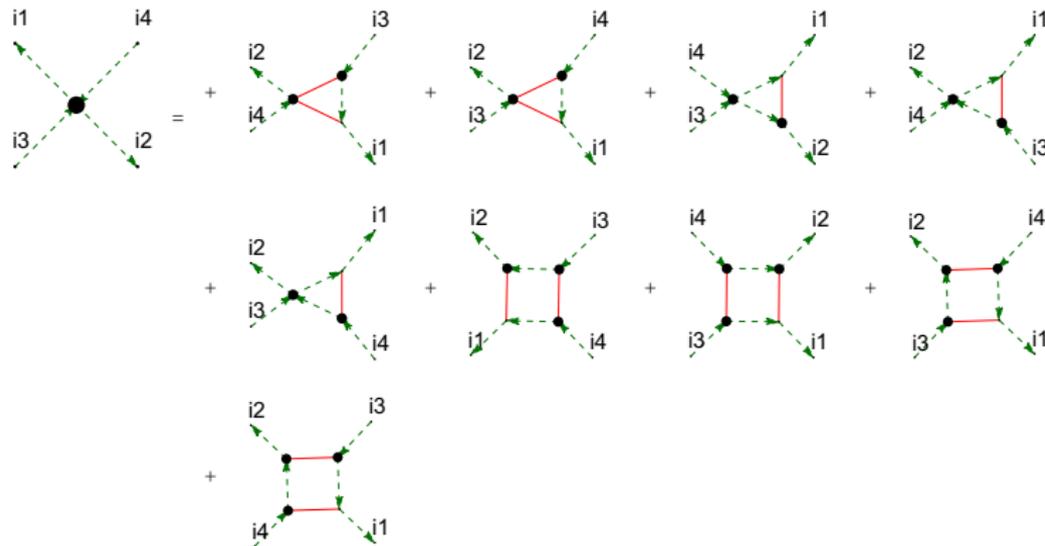
# The four-ghost vertex DSE

→ Truncation discards only one diagram.



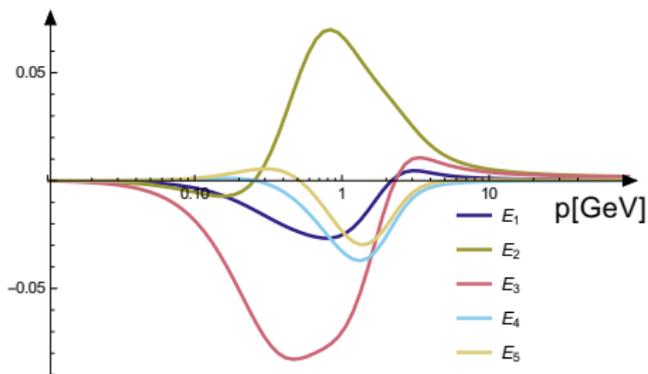
# The four-ghost vertex DSE

→ Truncation discards only one diagram.



# Results for the four-ghost vertex

Kinematic approximation: one-momentum configuration



→ All dressings very small.

[MQH '17]

$E_6, E_7, E_8 (\{\sigma_6, \sigma_7, \sigma_8\})$

Decouple into a homogeneous, linear equation. → Trivial solution always exists.  
Nontrivial one? → None found.

(Same applies to two-ghost-two-gluon vertex.)

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- Finite temperature
- Bound states
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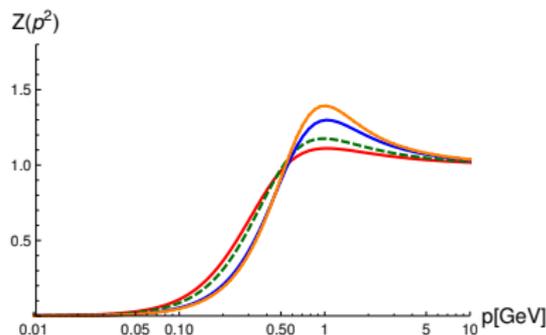
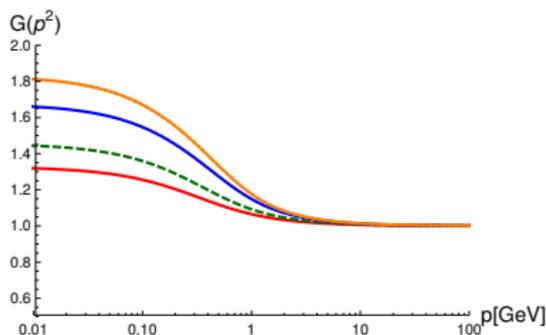
Thank you for your attention!

# Family of solutions in three dimensions

Cf. FRG results: Bare mass parameter from modified STIs [Cyrol, Fister, Mitter, Pawłowski, Strodthoff '15].

DSEs: Enforce family of solutions by fixing the gluon propagator at  $p^2 = 0$ .

Simple toy system with bare vertices [MQH, 1606.02068]:



⇒ Possibility of family of solutions.

NB: Effect overestimated here since vertices are fixed.