

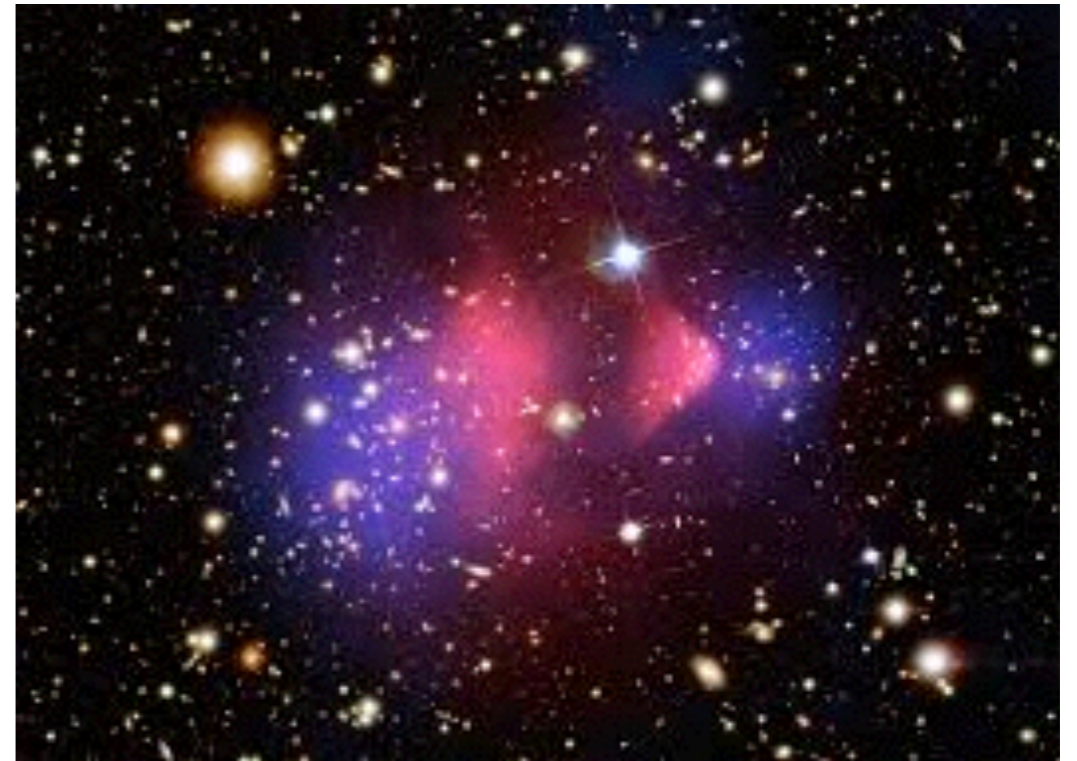
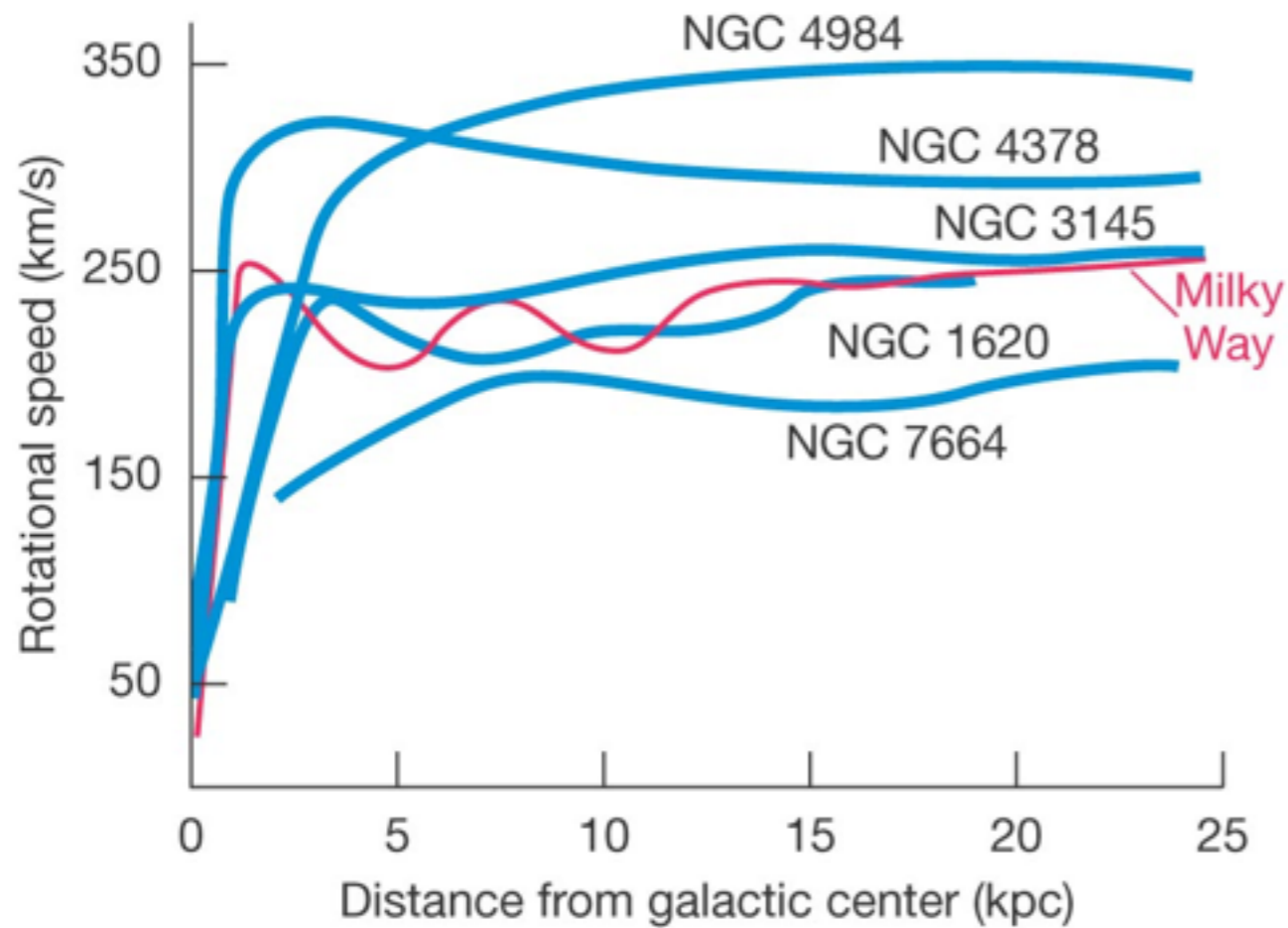
COLLIDING PIZZA SEMINAR

DARK MATTER @ LHC

Nishita Desai

**WHAT DO WE
KNOW SO FAR?**

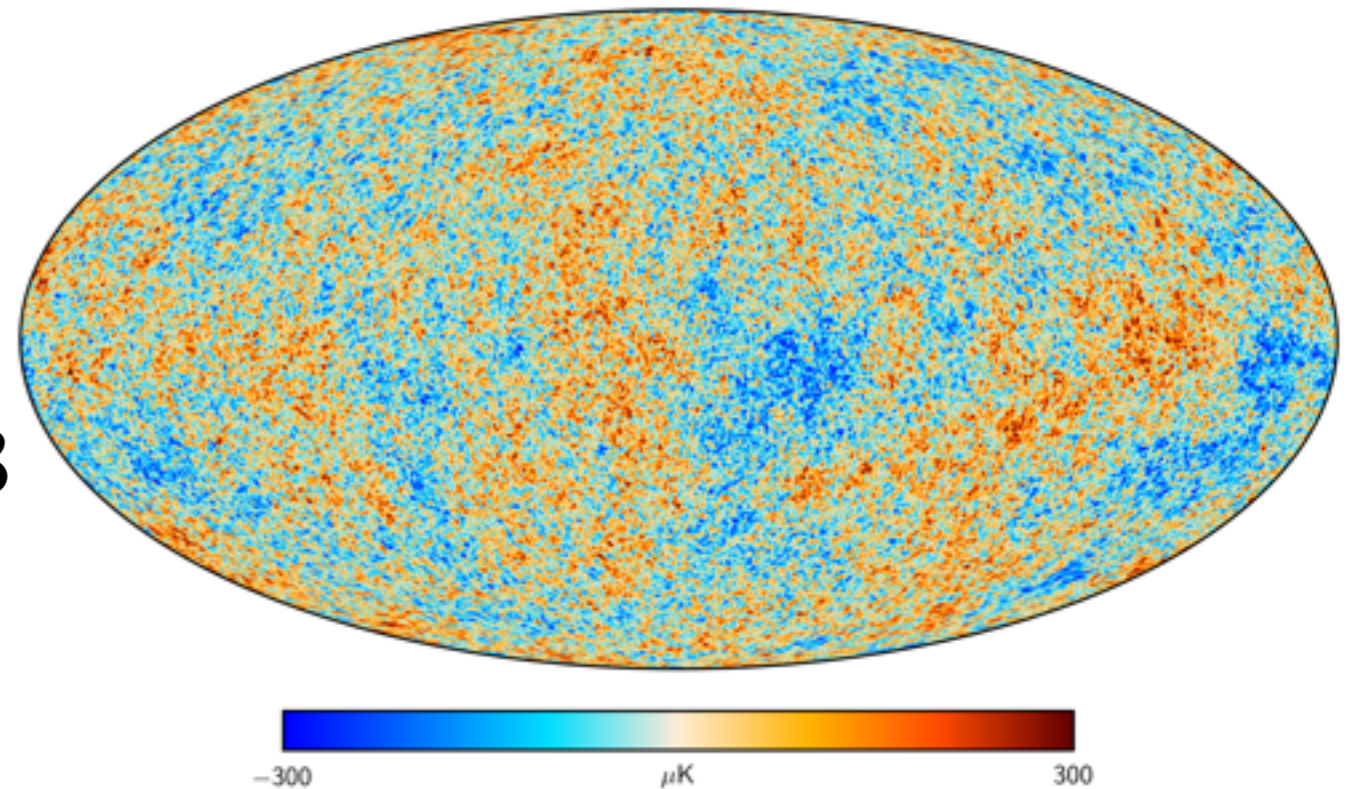
HINTS FROM ASTROPHYSICS



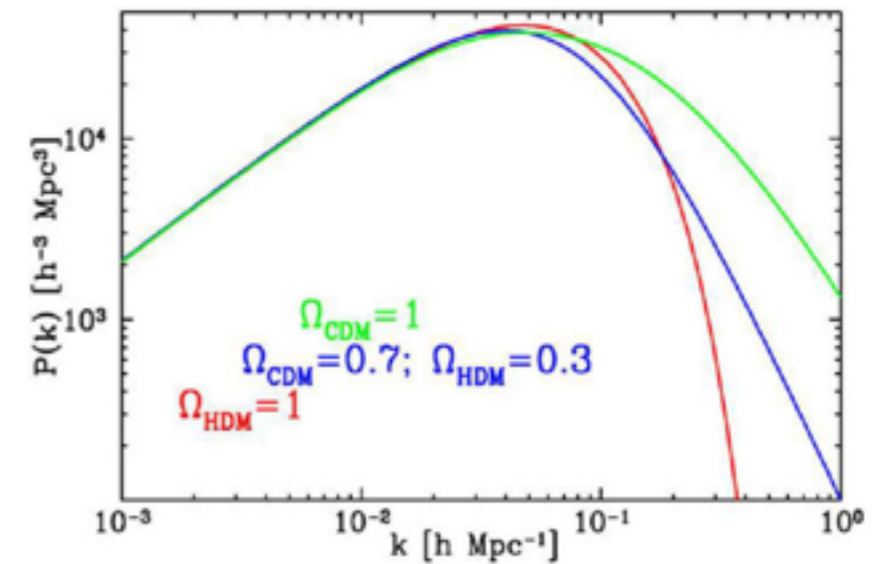
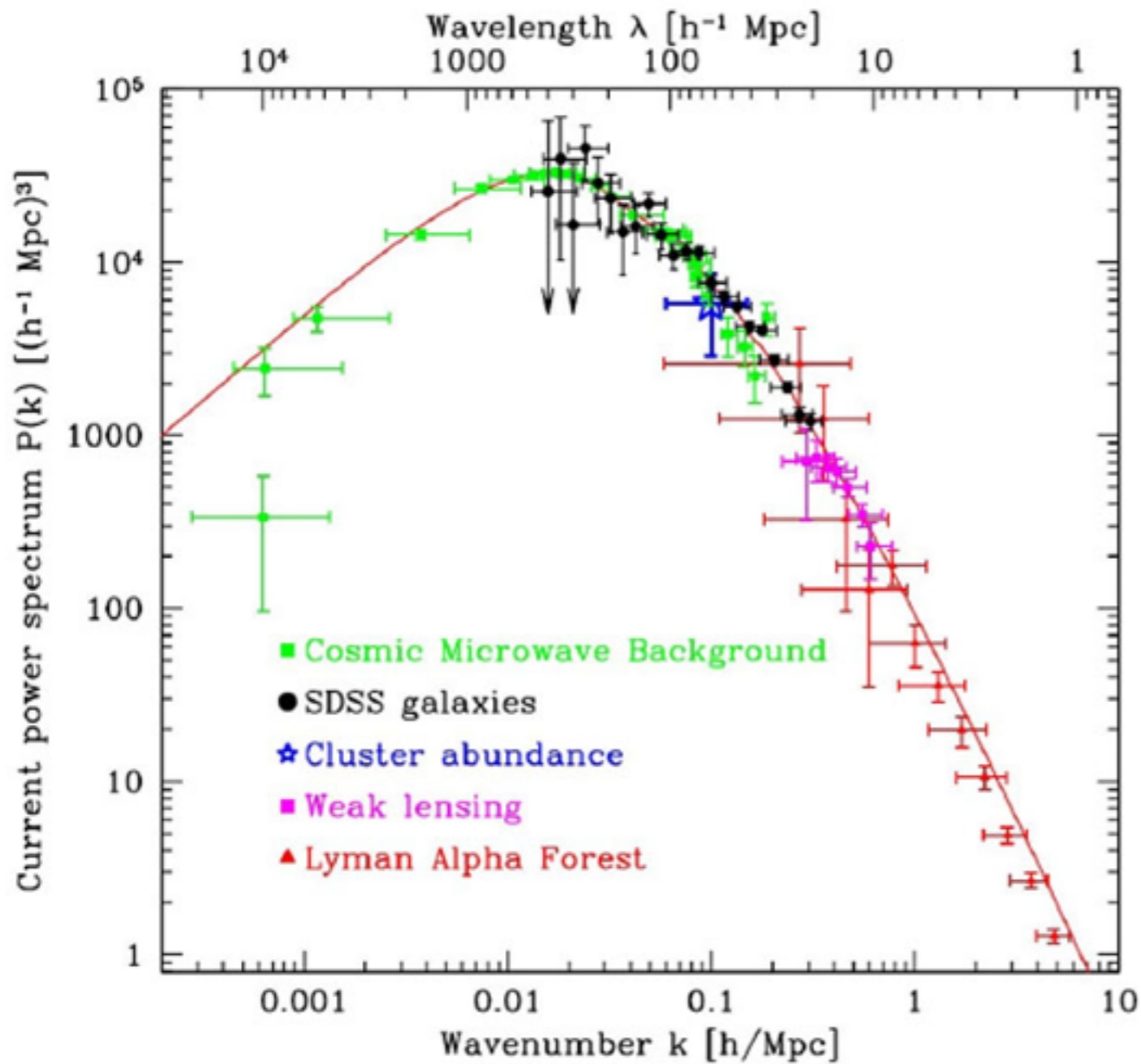
(b) Eric Chaisson, Steve McMillan

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$$\begin{aligned}\Omega_m h^2 &= 0.1415 \pm 0.0019 \\ \Omega_b h^2 &= 0.02226 \pm 0.00023 \\ \Omega_c h^2 &= 0.1186 \pm 0.0020\end{aligned}$$



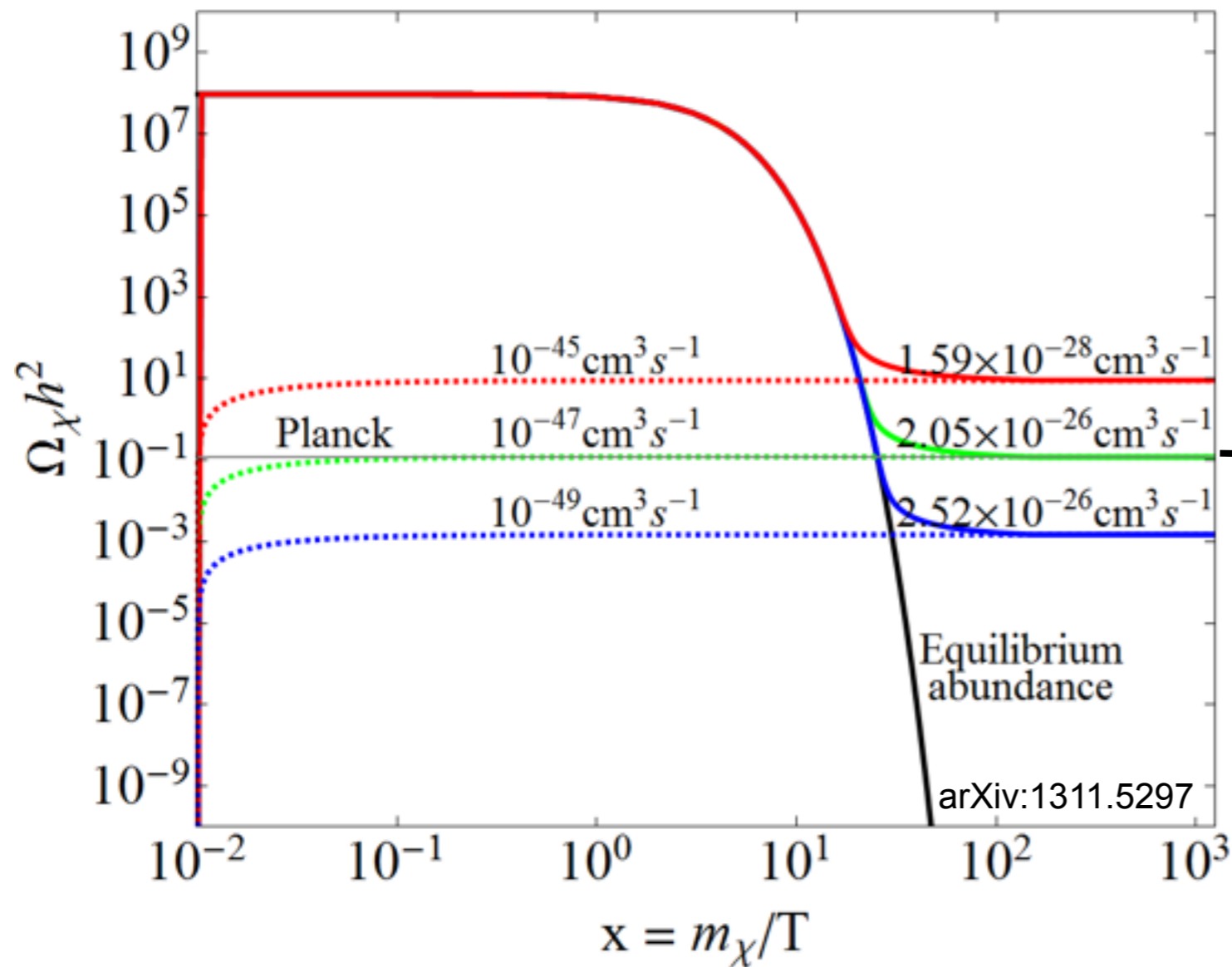
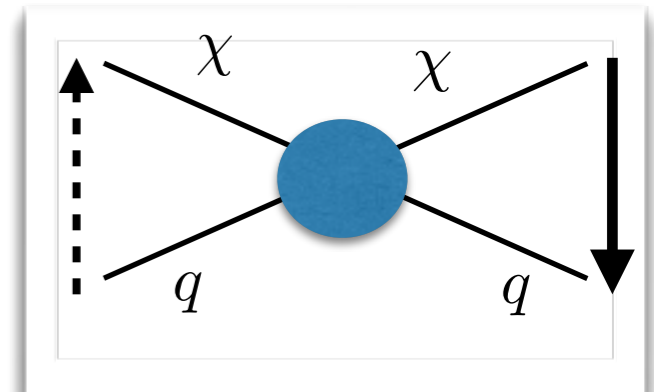
(DARK) MATTER POWER SPECTRUM



The assumption of
CDM is essential
to explain our
Universe

HOW TO CALCULATE THE RELIC DENSITY?

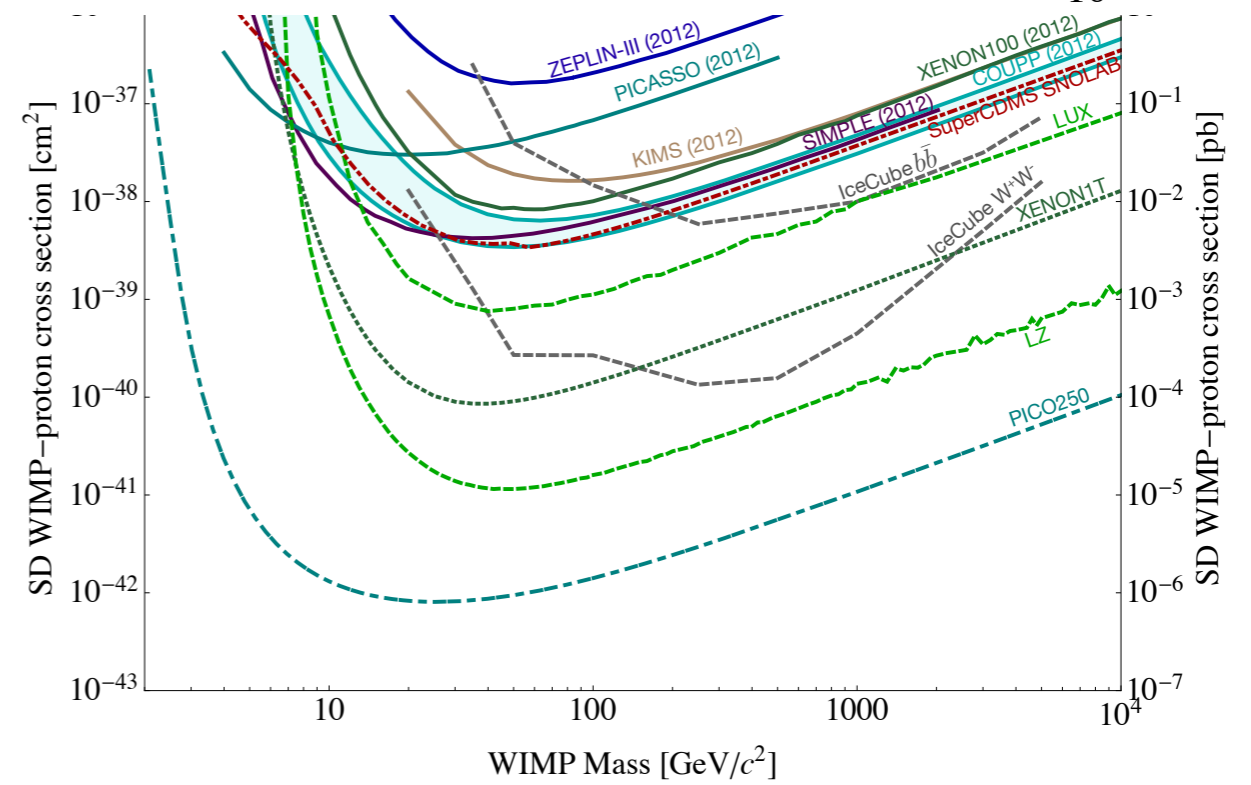
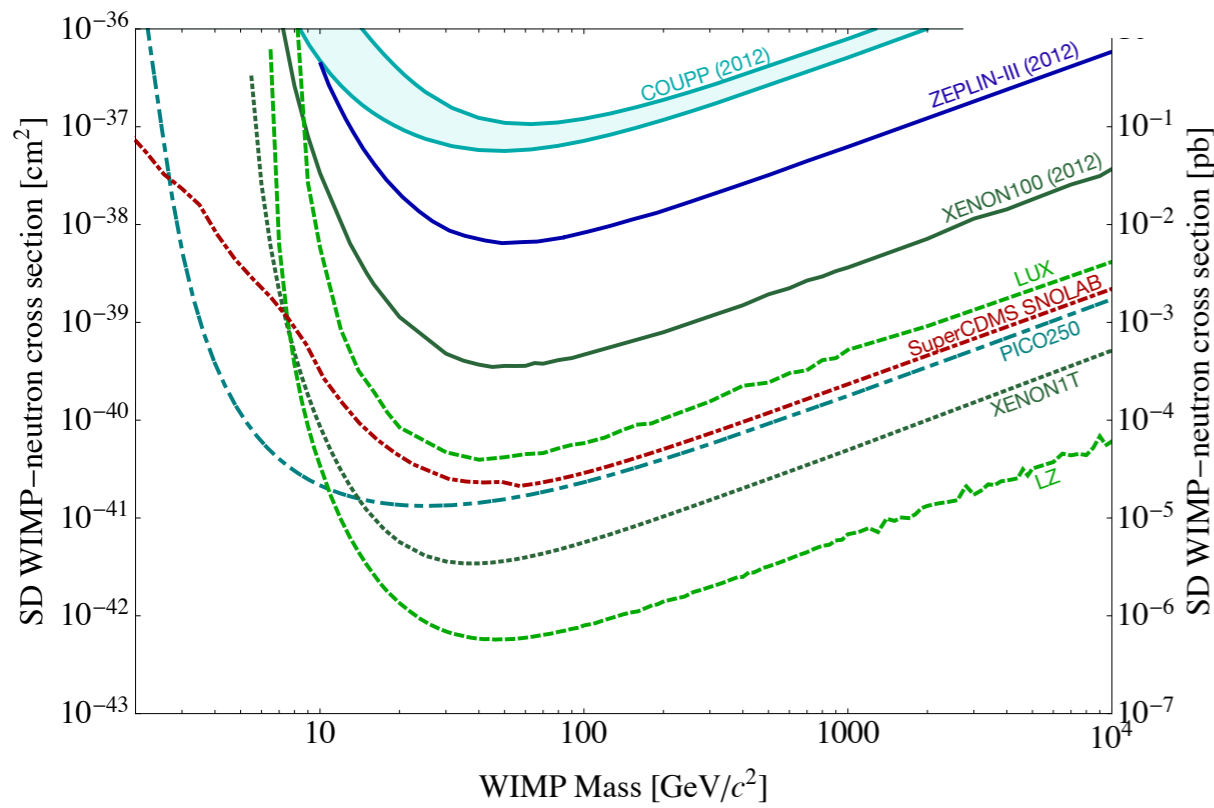
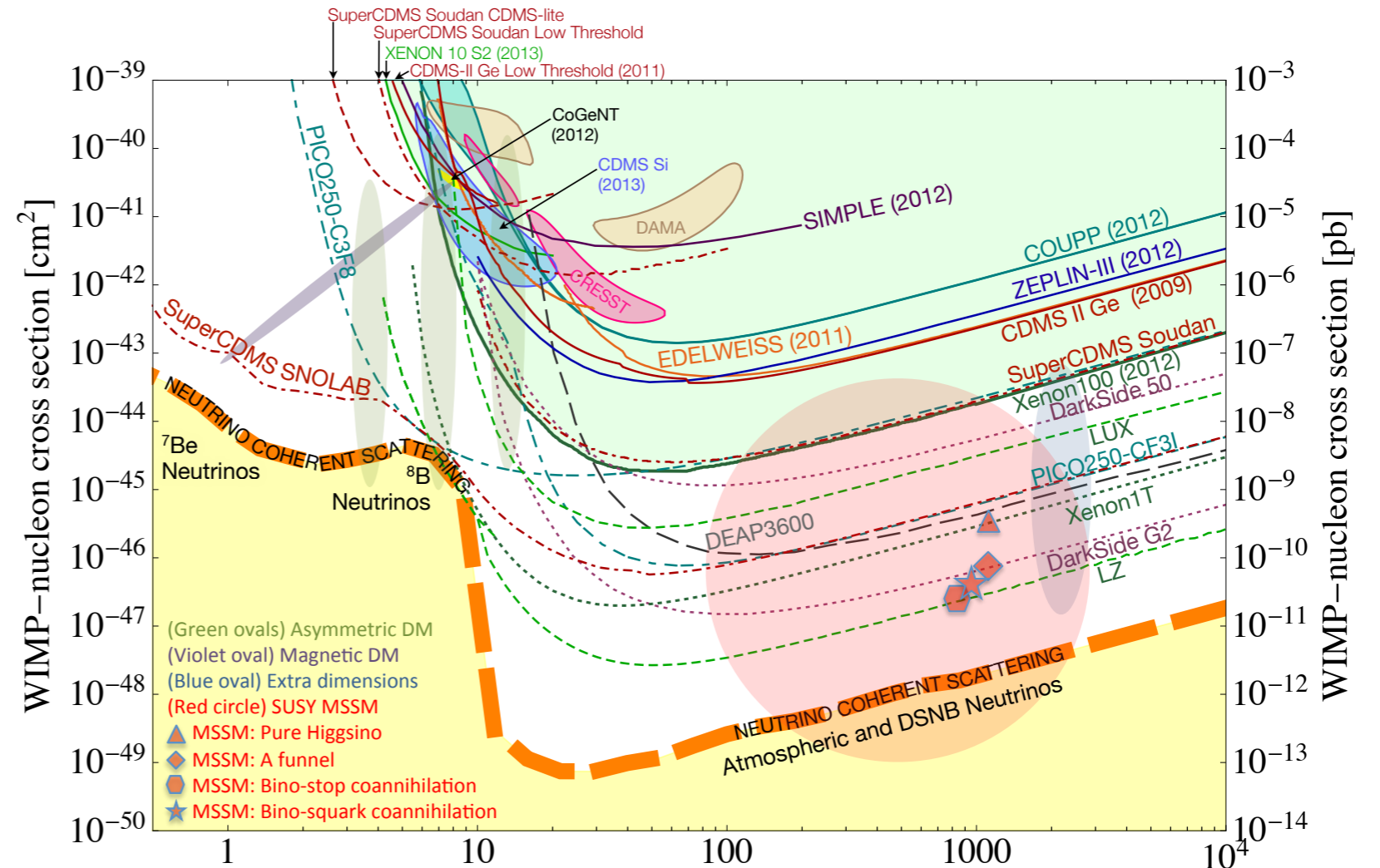
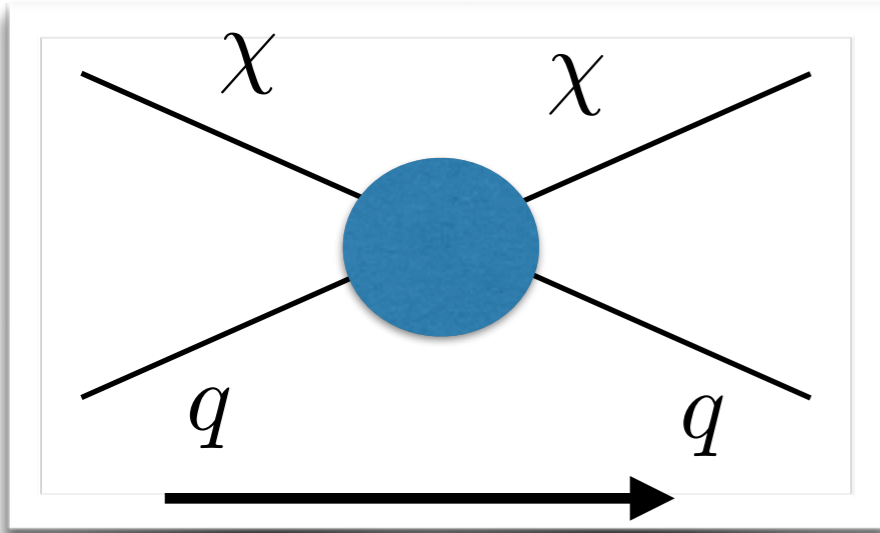
$$\frac{dn}{dt} + 3nH = -\frac{1}{2} \langle \sigma v \rangle (n^2 - (n^{eq})^2)$$



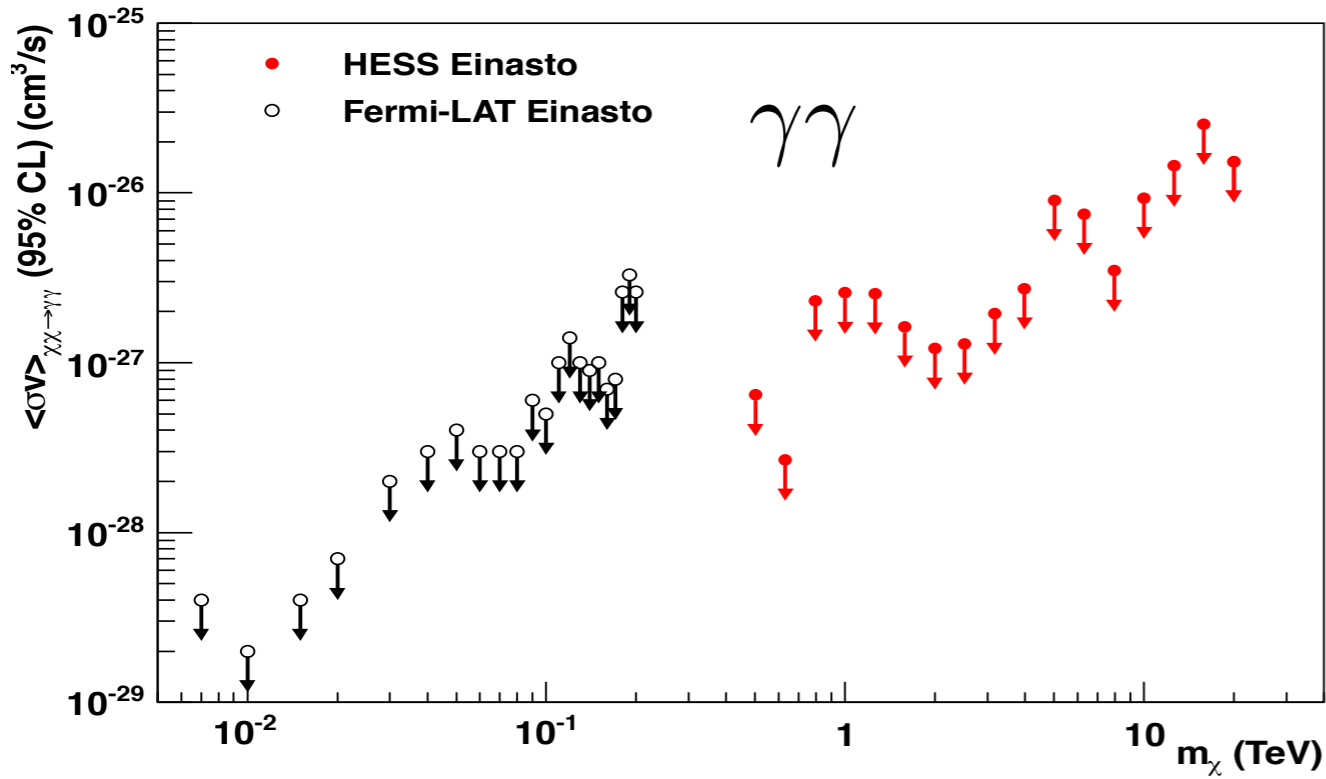
$$m_\chi \sim 100 \text{ GeV}$$

$$g \sim g_{EW}$$

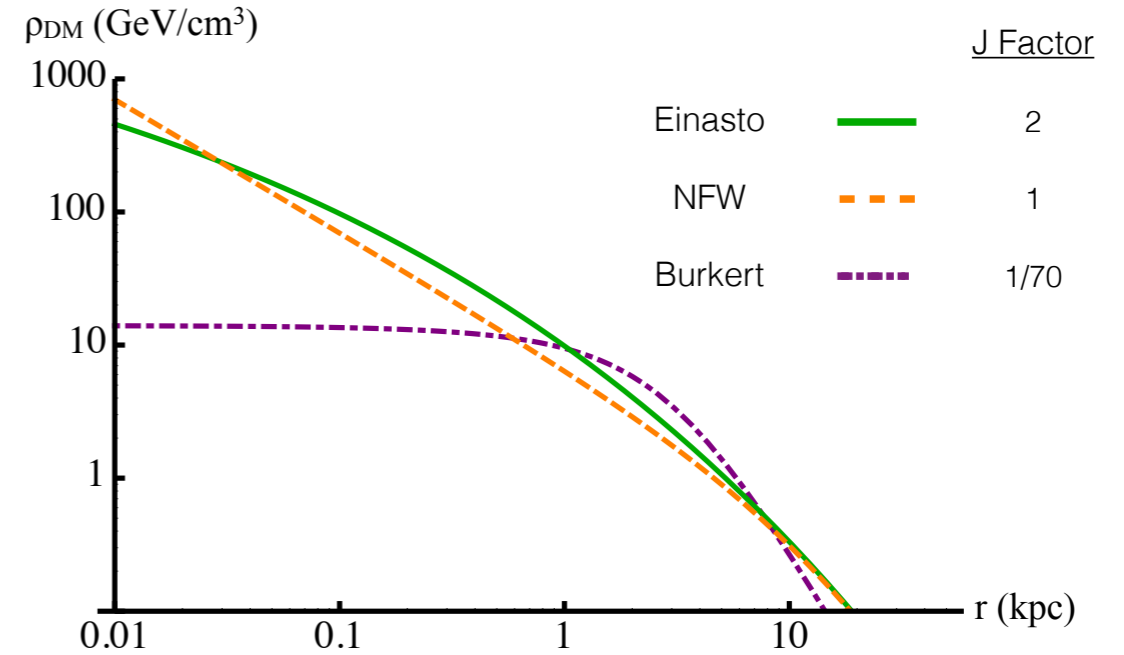
LOOKING FOR DM WINDS: DIRECT DETECTION



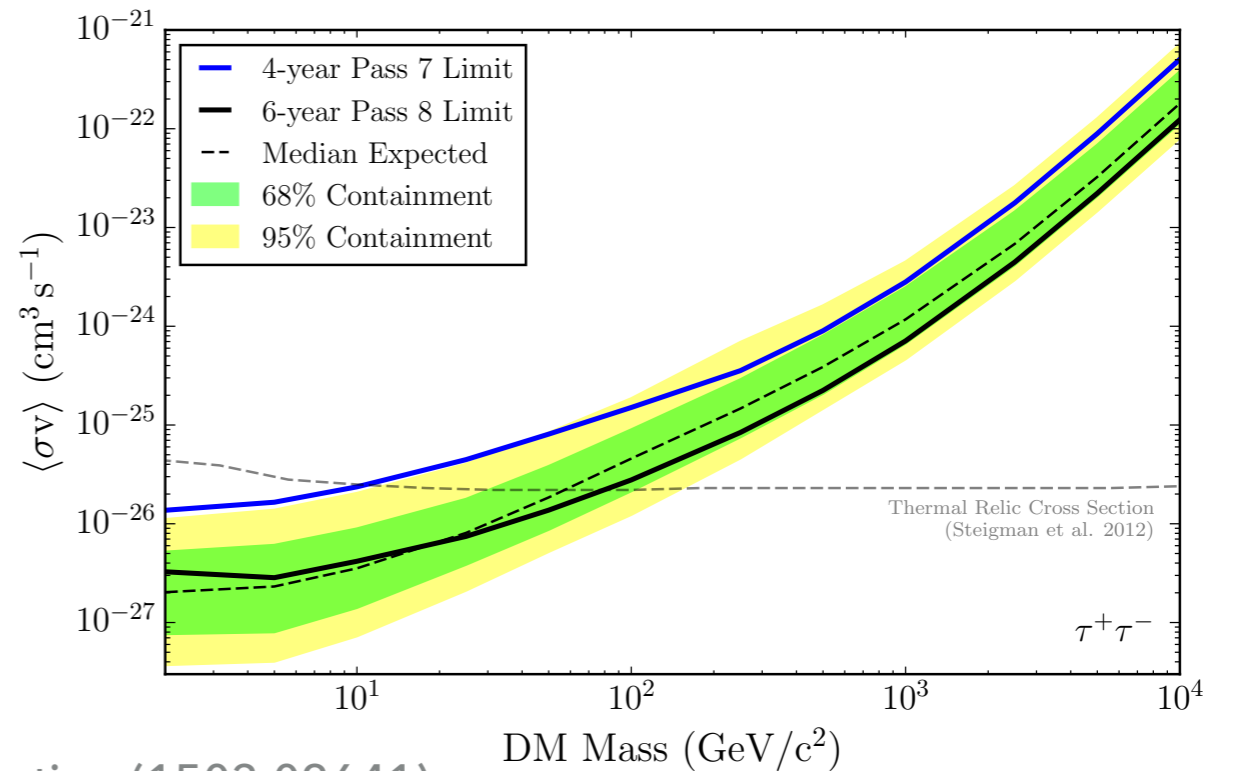
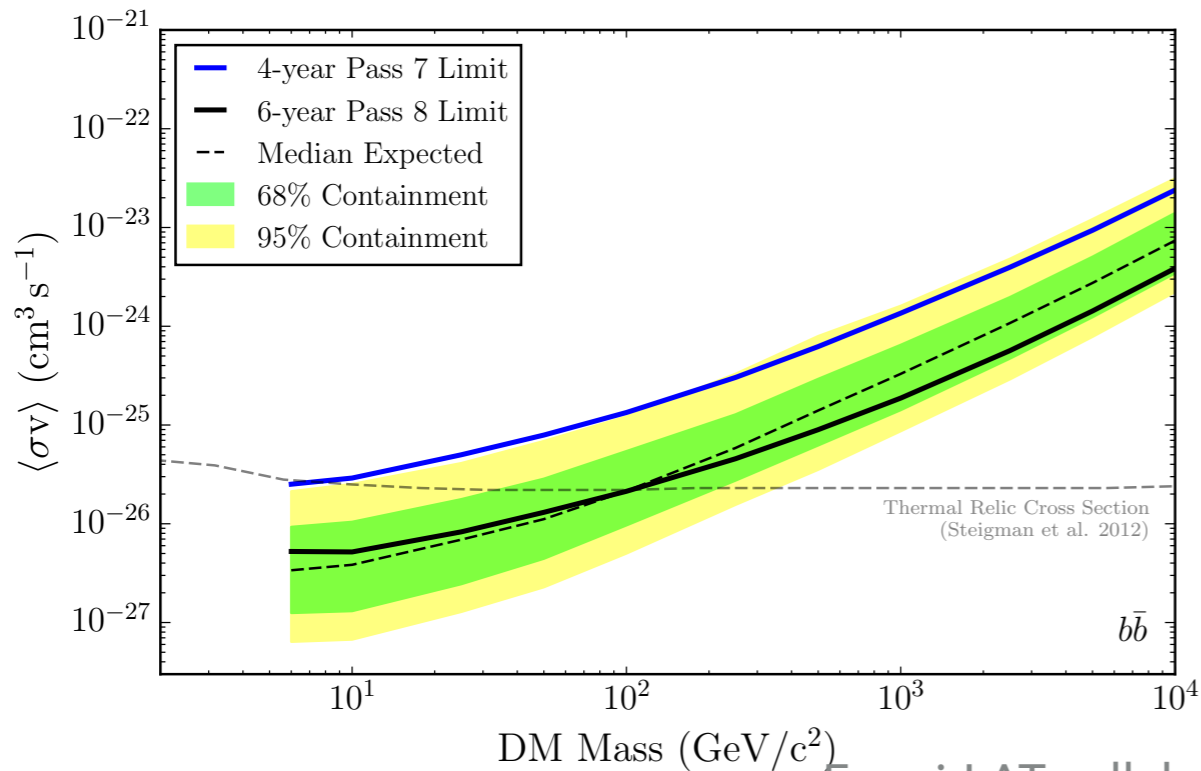
LOOKING FOR LINES IN THE SKY: INDIRECT DETECTION



$$\rho_{\text{NFW}}(r) = \frac{\rho_\odot}{(r/R)(1+r/R)^2}, \quad \rho_{\text{Ein}}(r) = \rho_\odot \exp\left[-\frac{2}{\alpha}\left(\left(\frac{r}{R}\right)^\alpha - 1\right)\right]$$



Dark Matter Halo Profiles



WHY USE COLLIDERS?

WE HAVE ONE, MIGHT AS WELL USE IT!

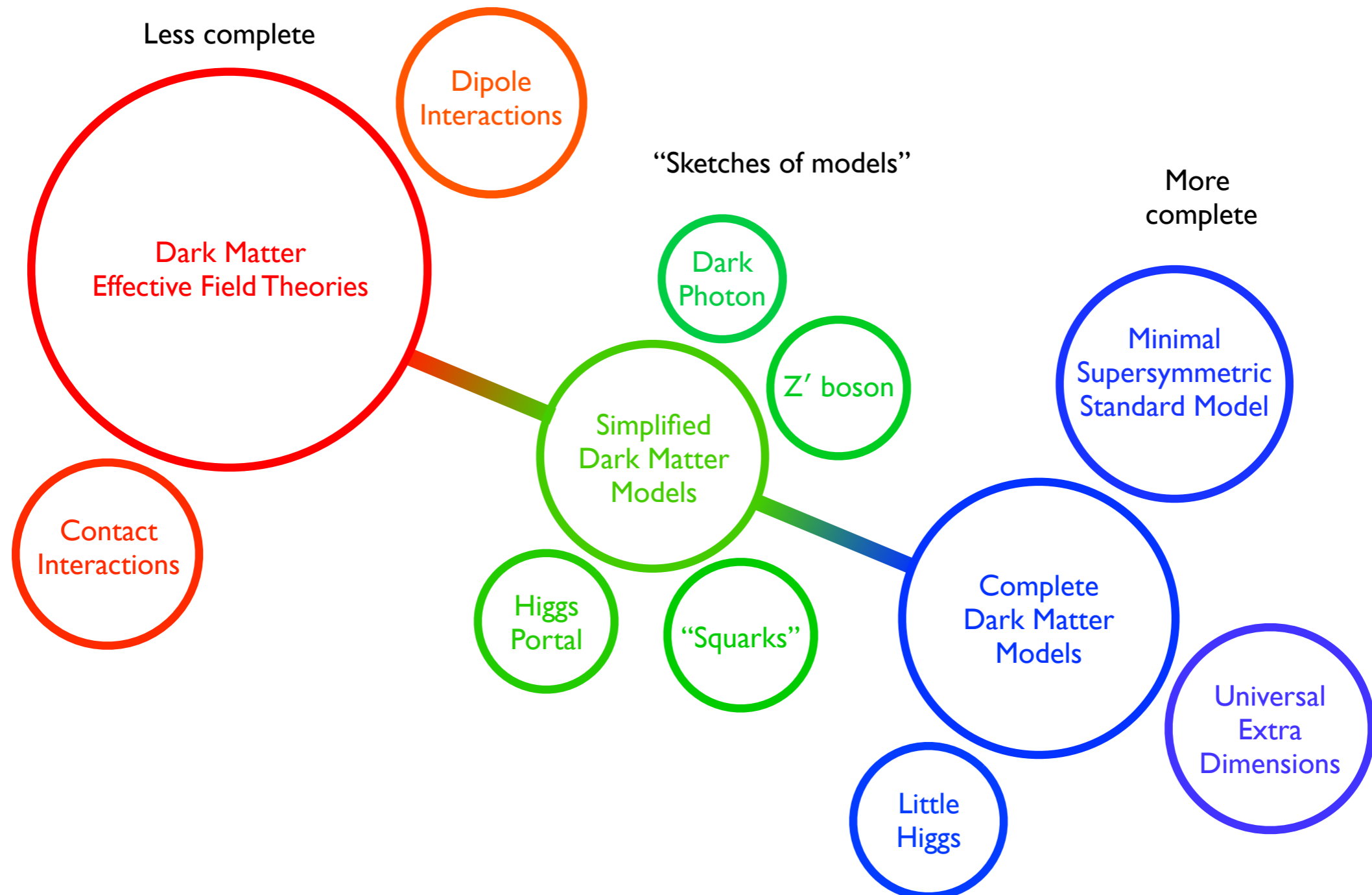
BUT ALSO . . .

IT TURNS OUT, COLLIDERS CAN DO THINGS OTHER
EXPERIMENTS CAN'T — BETTER SPIN DEPENDENT
SENSITIVITY + CONSTRAIN LOW RECOIL REGION +
LOOK FOR ACCOMPANYING PARTICLES

TO MODEL DM WE NEED TO KNOW:

Does it couple directly to some SM particle?

If there is a mediator, how does the mediator couple to SM? to DM?



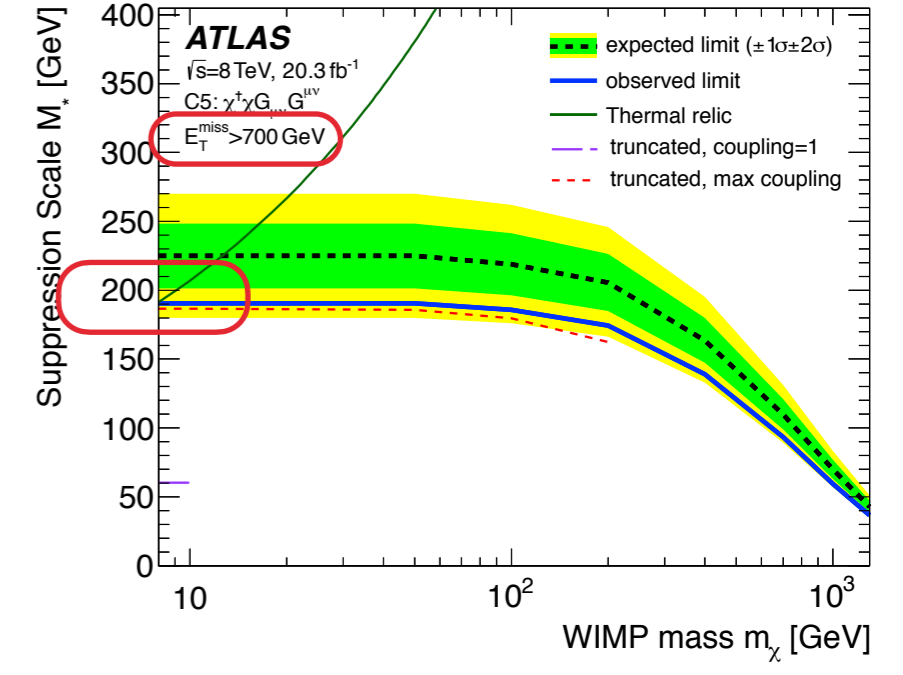
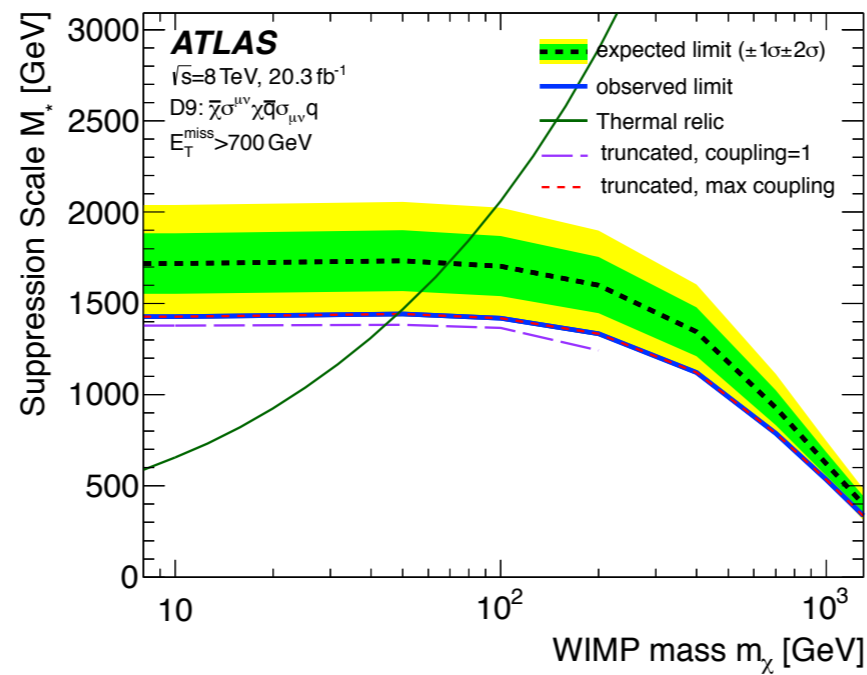
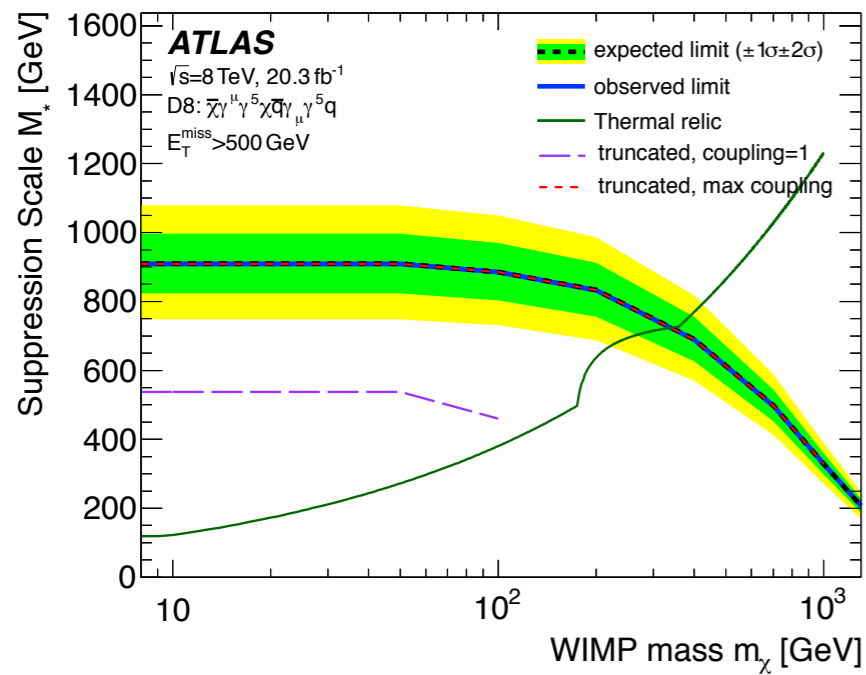
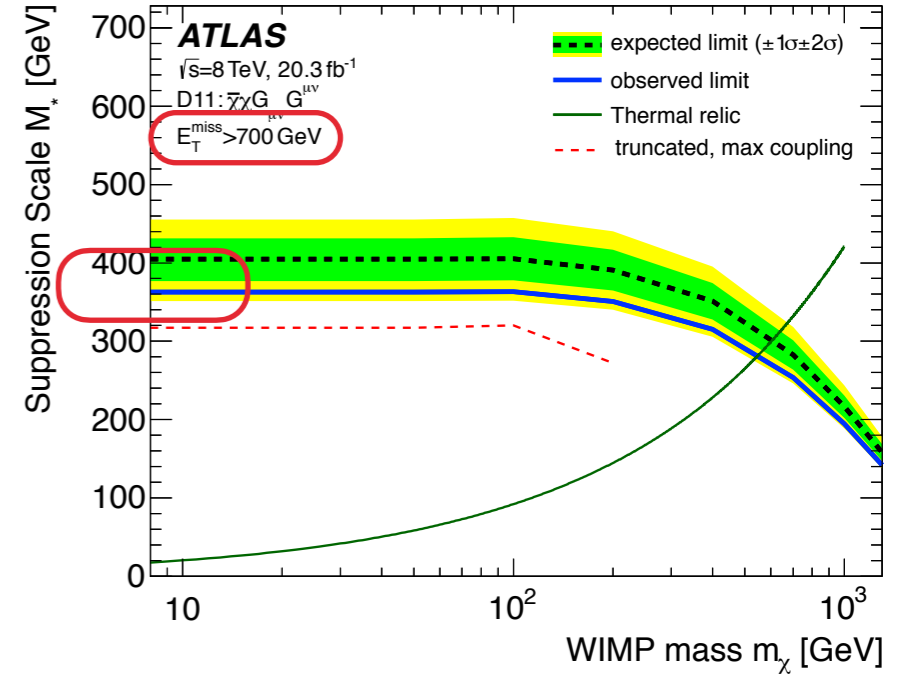
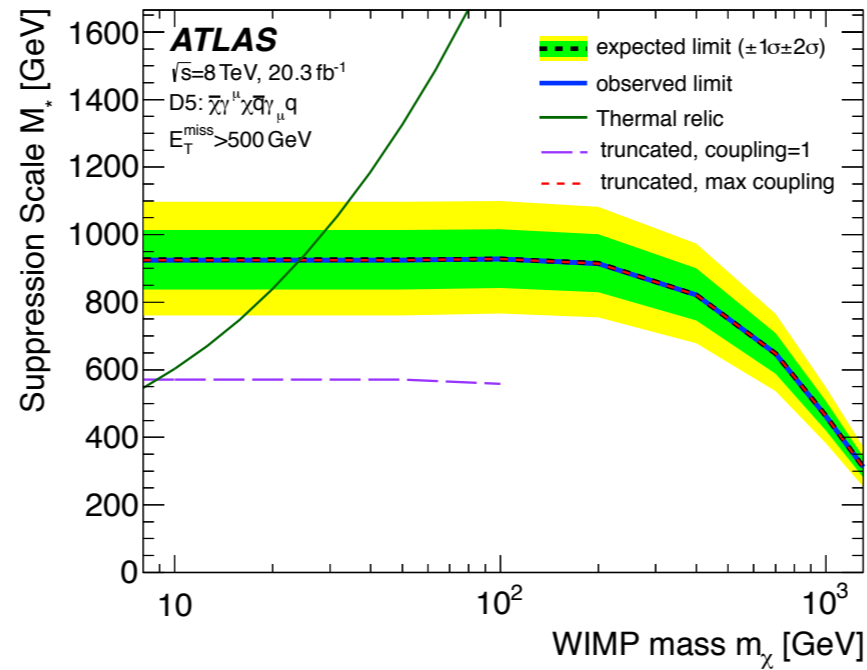
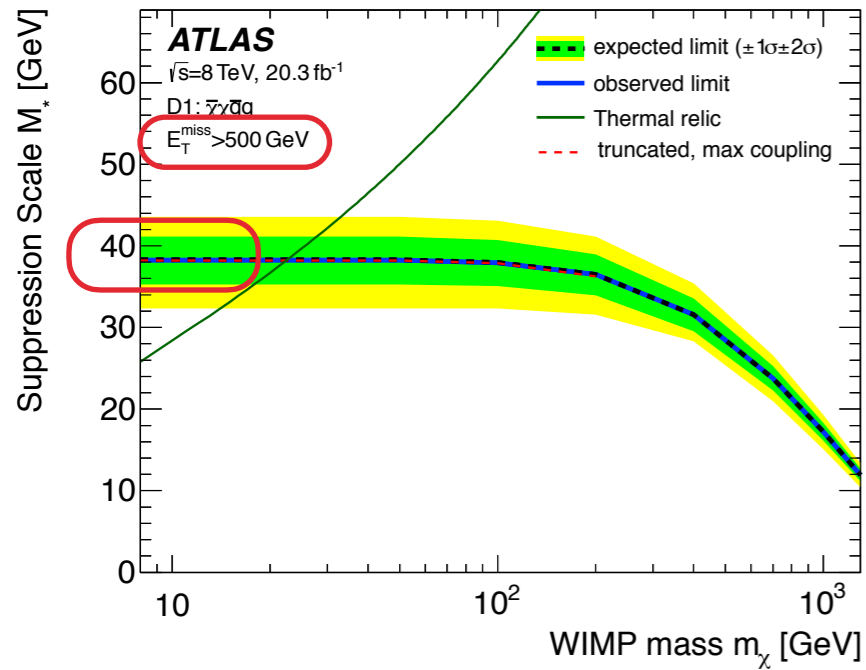
COLLIDER SEARCHES: A COMPLEMENTARY VIEW

Goodman et al. (2010)

Name	Operator	Coefficient
D1	$\bar{\chi}\chi\bar{q}q$	m_q/M_*^3
D2	$\bar{\chi}\gamma^5\chi\bar{q}q$	im_q/M_*^3
D3	$\bar{\chi}\chi\bar{q}\gamma^5q$	im_q/M_*^3
D4	$\bar{\chi}\gamma^5\chi\bar{q}\gamma^5q$	m_q/M_*^3
D5	$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q$	$1/M_*^2$
D6	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q$	$1/M_*^2$
D7	$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu\gamma^5q$	$1/M_*^2$
D8	$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$	$1/M_*^2$
D9	$\bar{\chi}\sigma^{\mu\nu}\chi\bar{q}\sigma_{\mu\nu}q$	$1/M_*^2$
D10	$\bar{\chi}\sigma_{\mu\nu}\gamma^5\chi\bar{q}\sigma_{\alpha\beta}q$	i/M_*^2
D11	$\bar{\chi}\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/4M_*^3$
D12	$\bar{\chi}\gamma^5\chi G_{\mu\nu}G^{\mu\nu}$	$i\alpha_s/4M_*^3$
D13	$\bar{\chi}\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/4M_*^3$
D14	$\bar{\chi}\gamma^5\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$\alpha_s/4M_*^3$

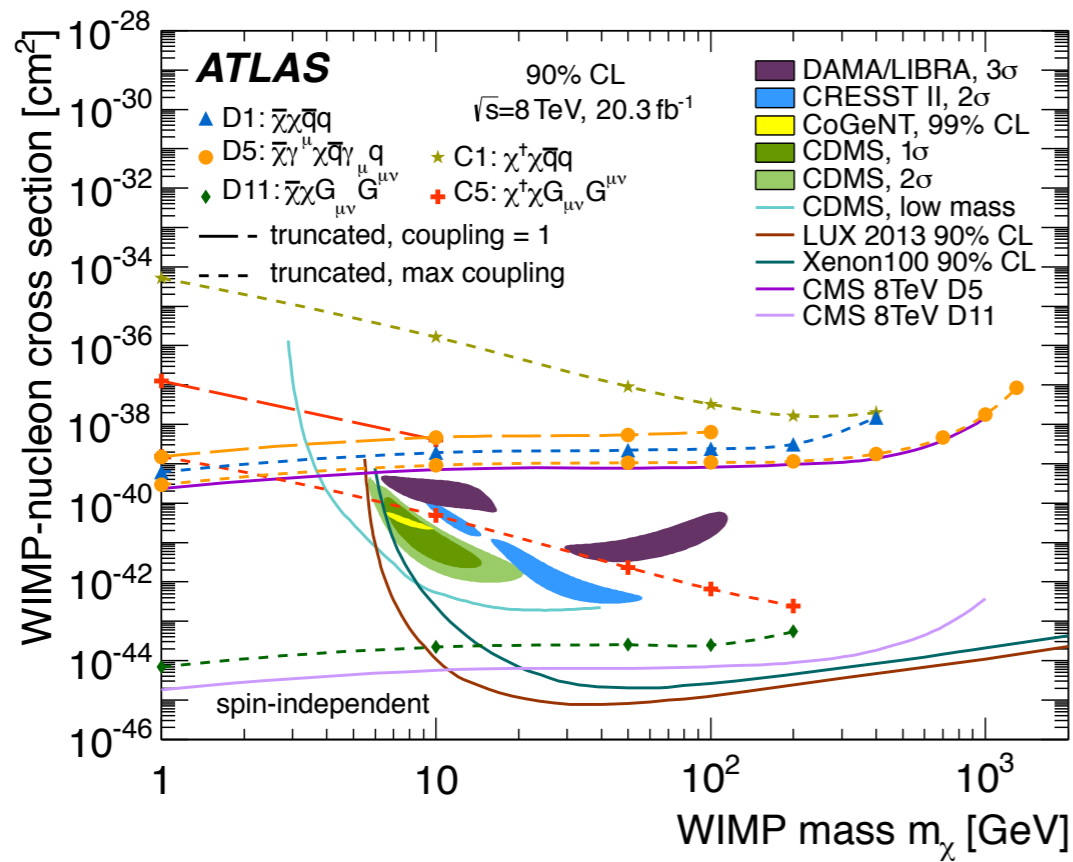
Name	Operator	Coefficient
C1	$\chi^\dagger\chi\bar{q}q$	m_q/M_*^2
C2	$\chi^\dagger\chi\bar{q}\gamma^5q$	im_q/M_*^2
C3	$\chi^\dagger\partial_\mu\chi\bar{q}\gamma^\mu q$	$1/M_*^2$
C4	$\chi^\dagger\partial_\mu\chi\bar{q}\gamma^\mu\gamma^5q$	$1/M_*^2$
C5	$\chi^\dagger\chi G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/4M_*^2$
C6	$\chi^\dagger\chi G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/4M_*^2$
R1	$\chi^2\bar{q}q$	$m_q/2M_*^2$
R2	$\chi^2\bar{q}\gamma^5q$	$im_q/2M_*^2$
R3	$\chi^2 G_{\mu\nu}G^{\mu\nu}$	$\alpha_s/8M_*^2$
R4	$\chi^2 G_{\mu\nu}\tilde{G}^{\mu\nu}$	$i\alpha_s/8M_*^2$

ATLAS LIMITS ON EFT OPERATORS

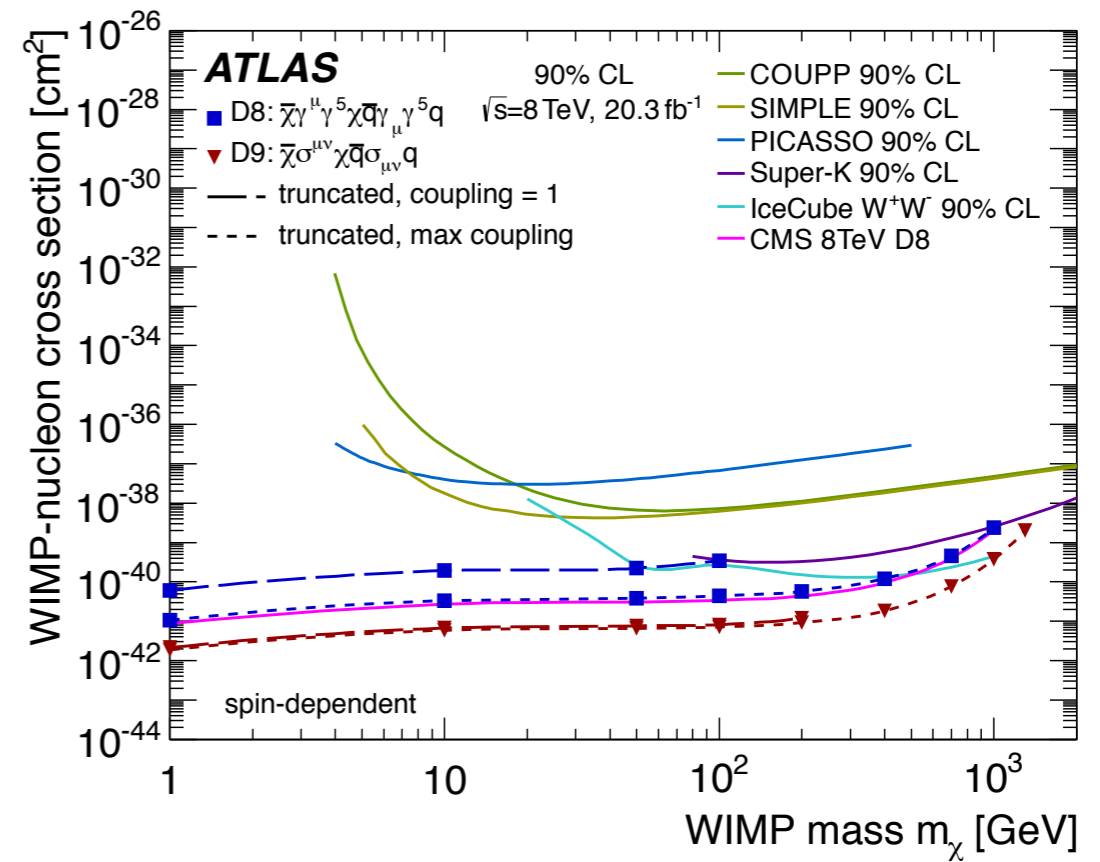


(c)

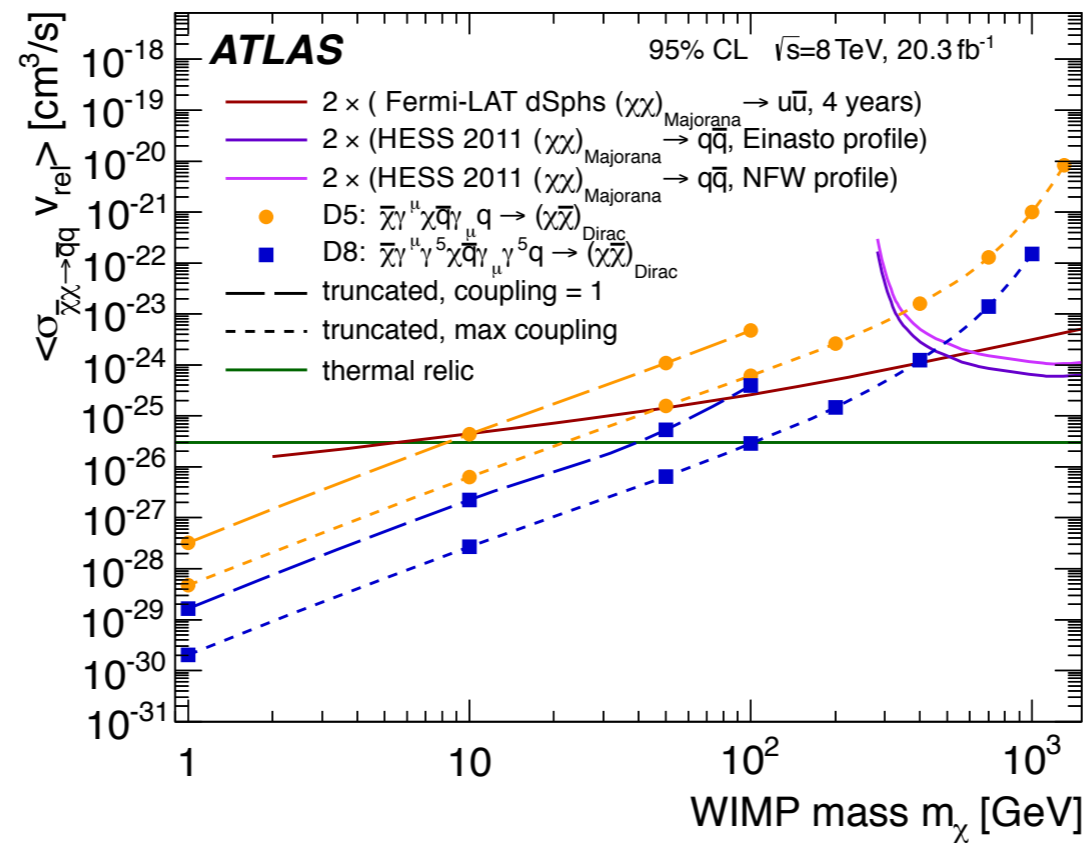
(d)



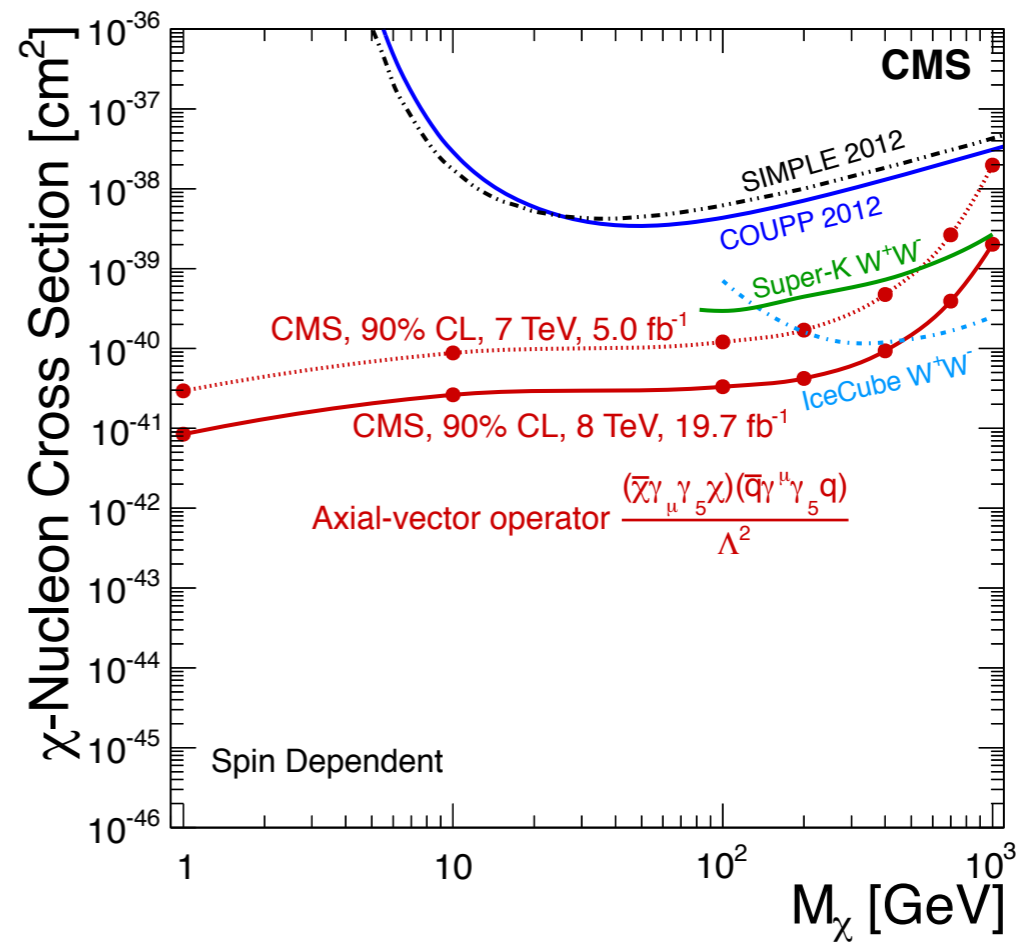
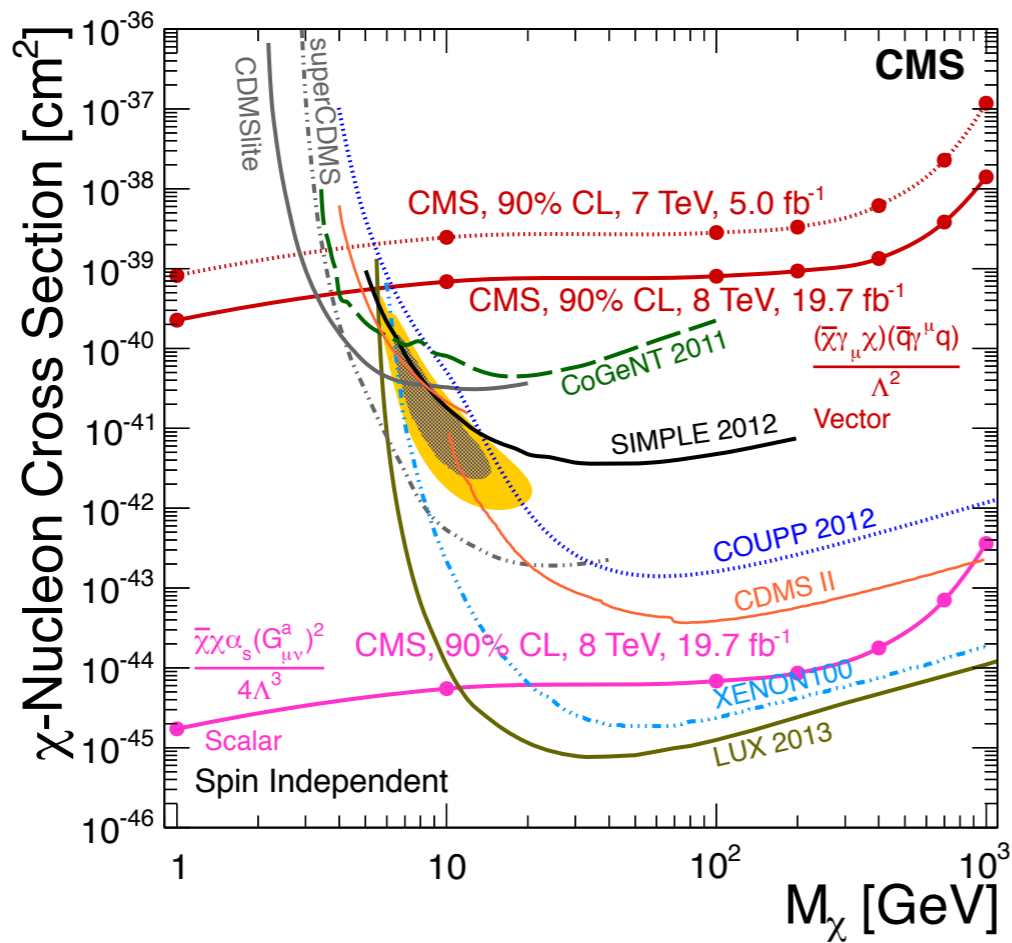
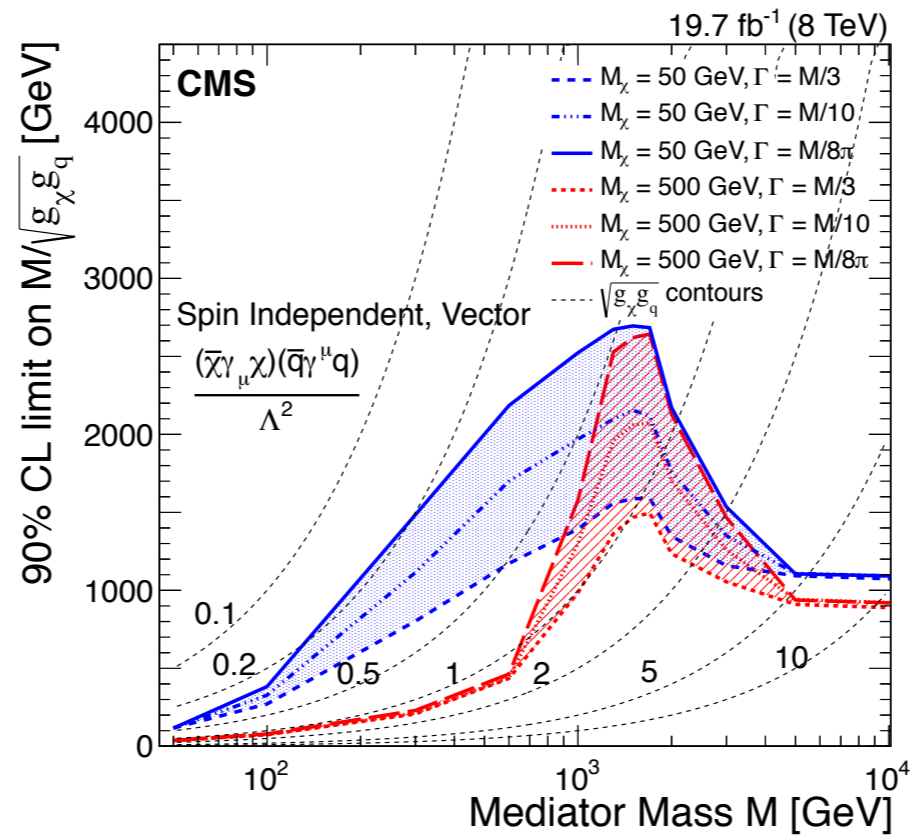
(a)



(b)



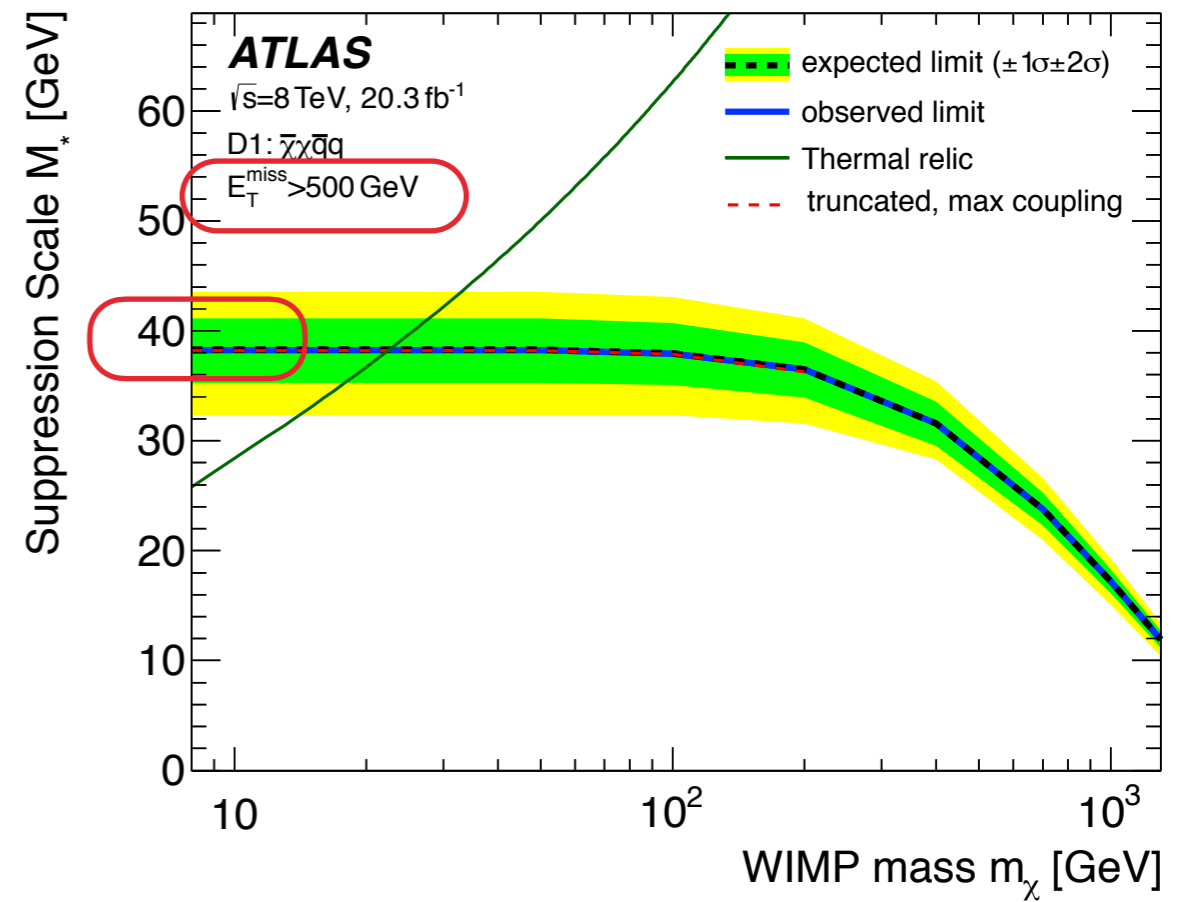
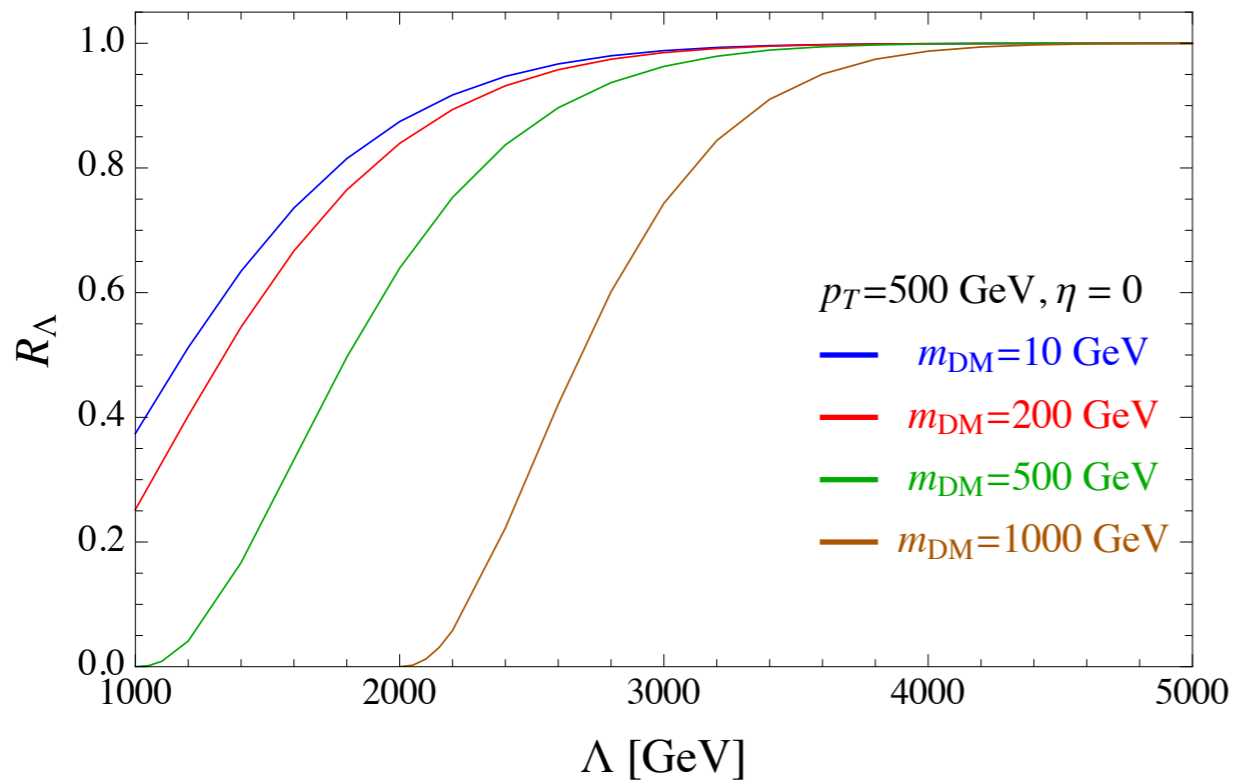
CMS LIMITS ON EFT OPERATORS



Original idea of “EFT”:

$$\bar{q}q \frac{g}{p^2 - M^2} \bar{\psi}\psi \xrightarrow{M \gg p} \frac{g}{M^2} \bar{q}q \bar{\psi}\psi$$

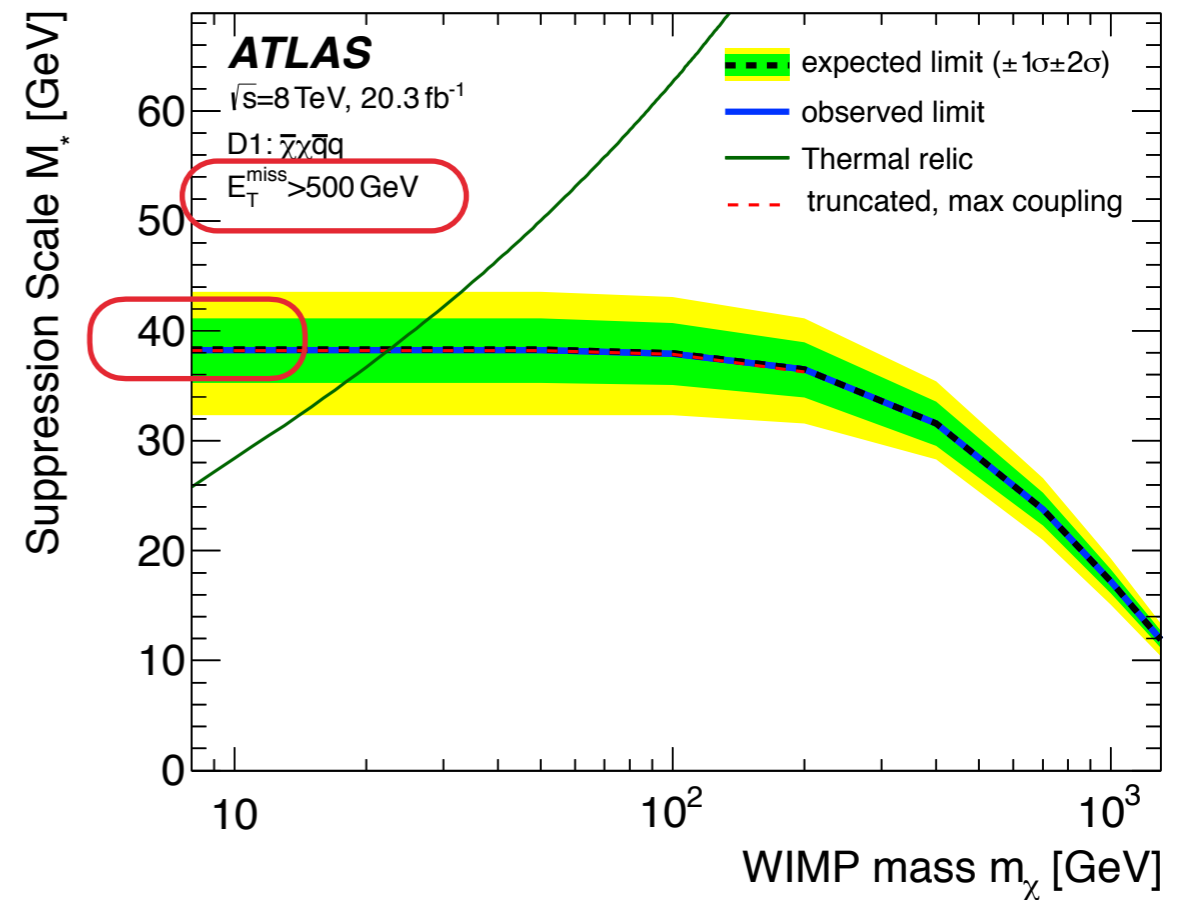
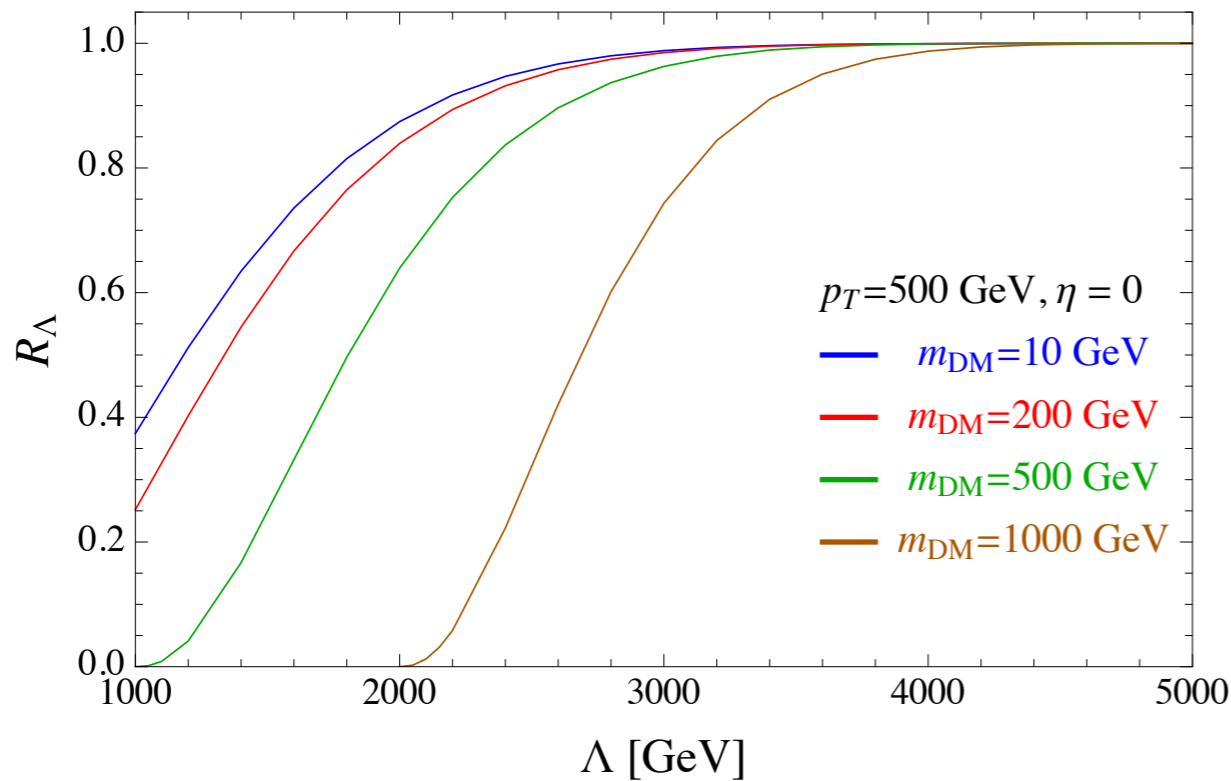
So how does one live with:



Original idea of “EFT”:

$$\bar{q}q \frac{g}{p^2 - M^2} \bar{\psi}\psi \xrightarrow{M \gg p} \frac{g}{M^2} \bar{q}q \bar{\psi}\psi$$

So how does one live with:

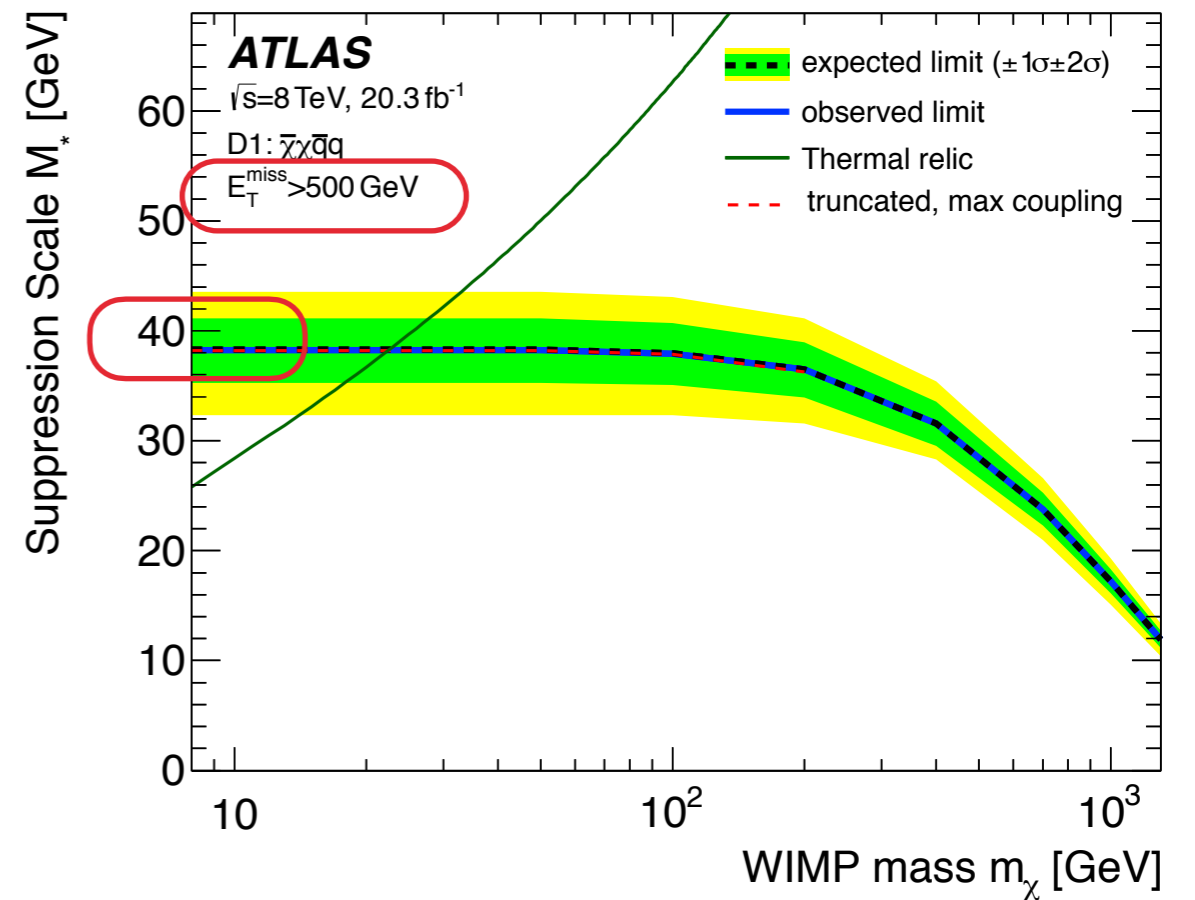
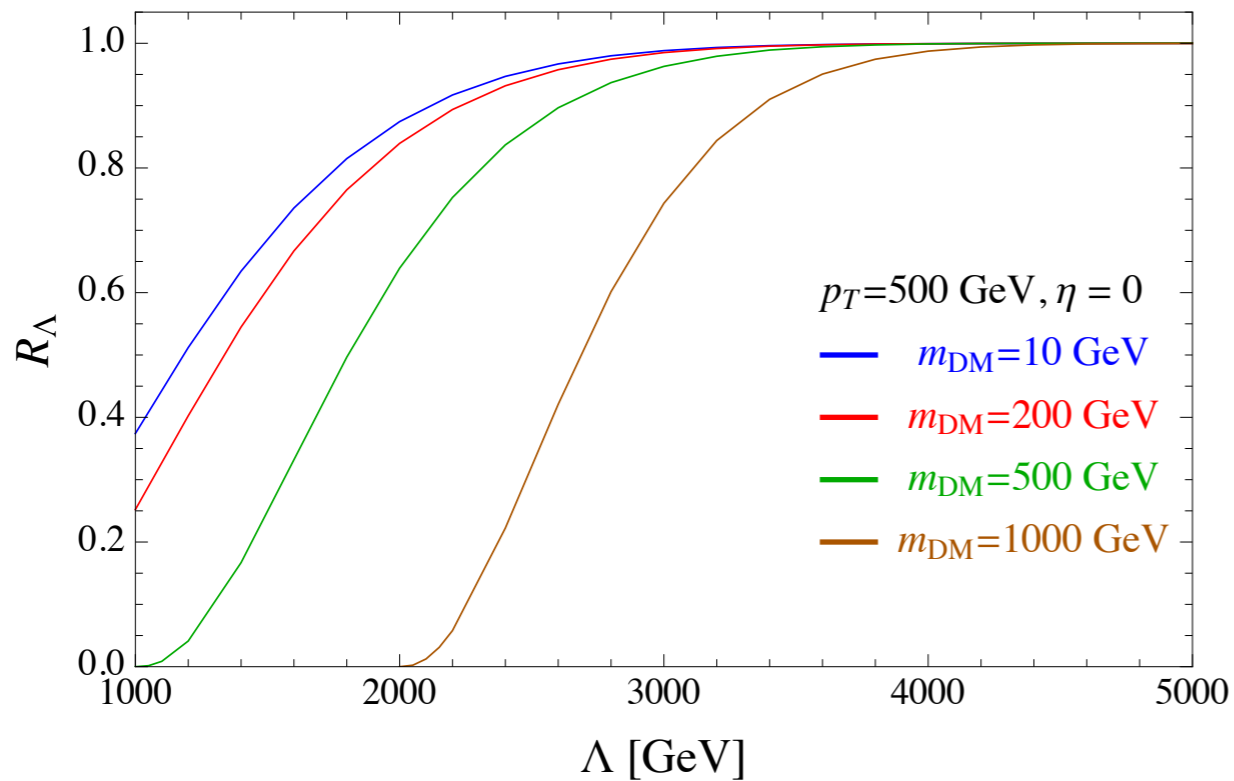


Option 1: Use only those events with $p \ll \Lambda$

Original idea of “EFT”:

$$\bar{q}q \frac{g}{p^2 - M^2} \bar{\psi}\psi \xrightarrow{M \gg p} \frac{g}{M^2} \bar{q}q \bar{\psi}\psi$$

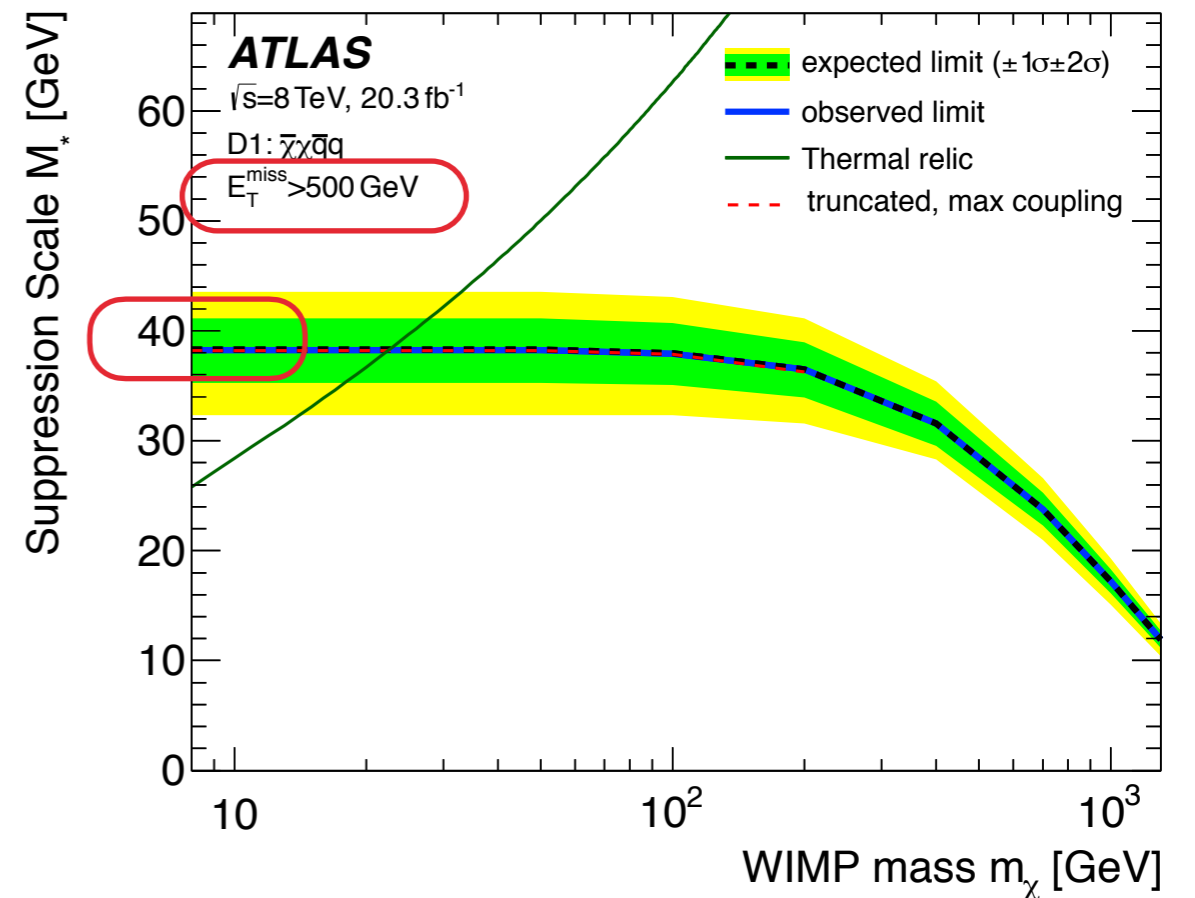
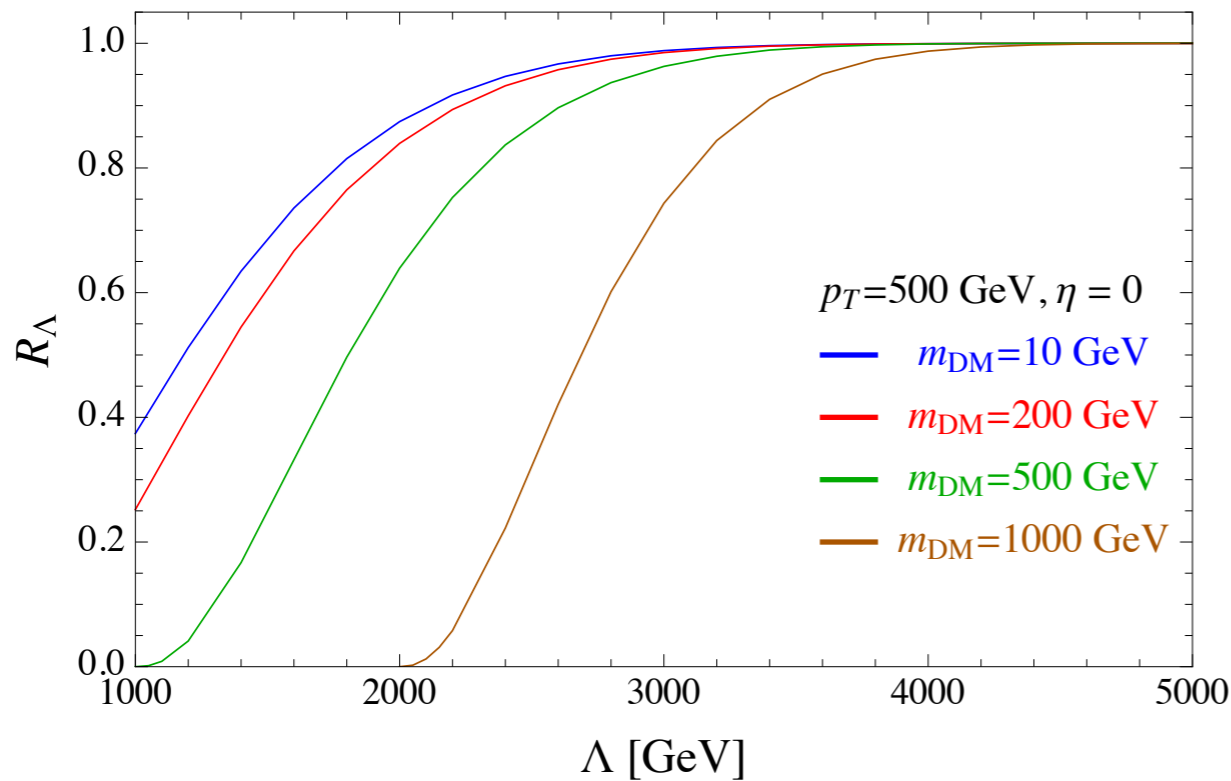
So how does one live with:



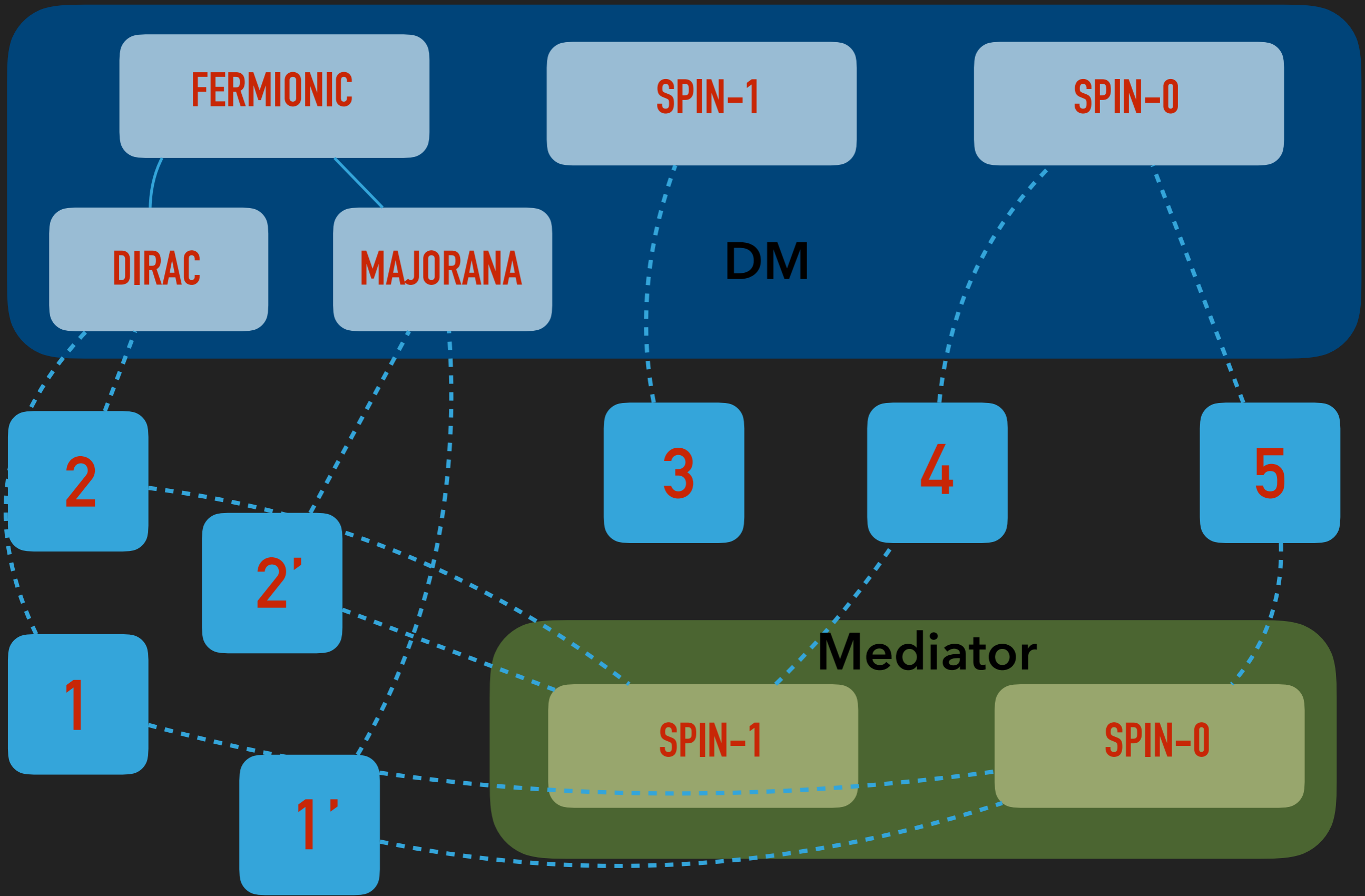
Original idea of “EFT”:

$$\bar{q}q \frac{g}{p^2 - M^2} \bar{\psi}\psi \xrightarrow{M \gg p} \frac{g}{M^2} \bar{q}q \bar{\psi}\psi$$

So how does one live with:



Option 2: Treat the EFT as simply a basis of operators and not as a result of integrating out massive particles.



HOW TO WRITE A SIMPLIFIED MODEL?

1

$$\mathcal{L}_S = \frac{1}{2} \partial_\mu S \partial^\mu S - m_S^2 S^2 + \sum g_{s\chi\bar{\chi}} \bar{\chi} \chi S + \sum g_{sq\bar{q}} \bar{q} q S + \bar{\chi} (i\partial_\mu \gamma^\mu - m_\chi) \chi$$

$$\mathcal{L}_P = \frac{1}{2} \partial_\mu P \partial^\mu P - m_P^2 P^2 + \sum g_{s\chi\bar{\chi}} \bar{\chi} \gamma^5 \chi P + \sum g_{sq\bar{q}} \bar{q} \gamma^5 q P + \bar{\chi} (i\partial_\mu \gamma^\mu - m_\chi) \chi$$

$$\begin{aligned} \mathcal{L}_T = & \frac{1}{2} D_\mu T D^\mu T - m_T^2 T^2 + \sum g_{T\chi\bar{\chi}} (\bar{\chi} q T^* + \text{c.c.}) \\ & + \bar{\chi} (i\partial_\mu \gamma^\mu - m_\chi) \chi \end{aligned}$$

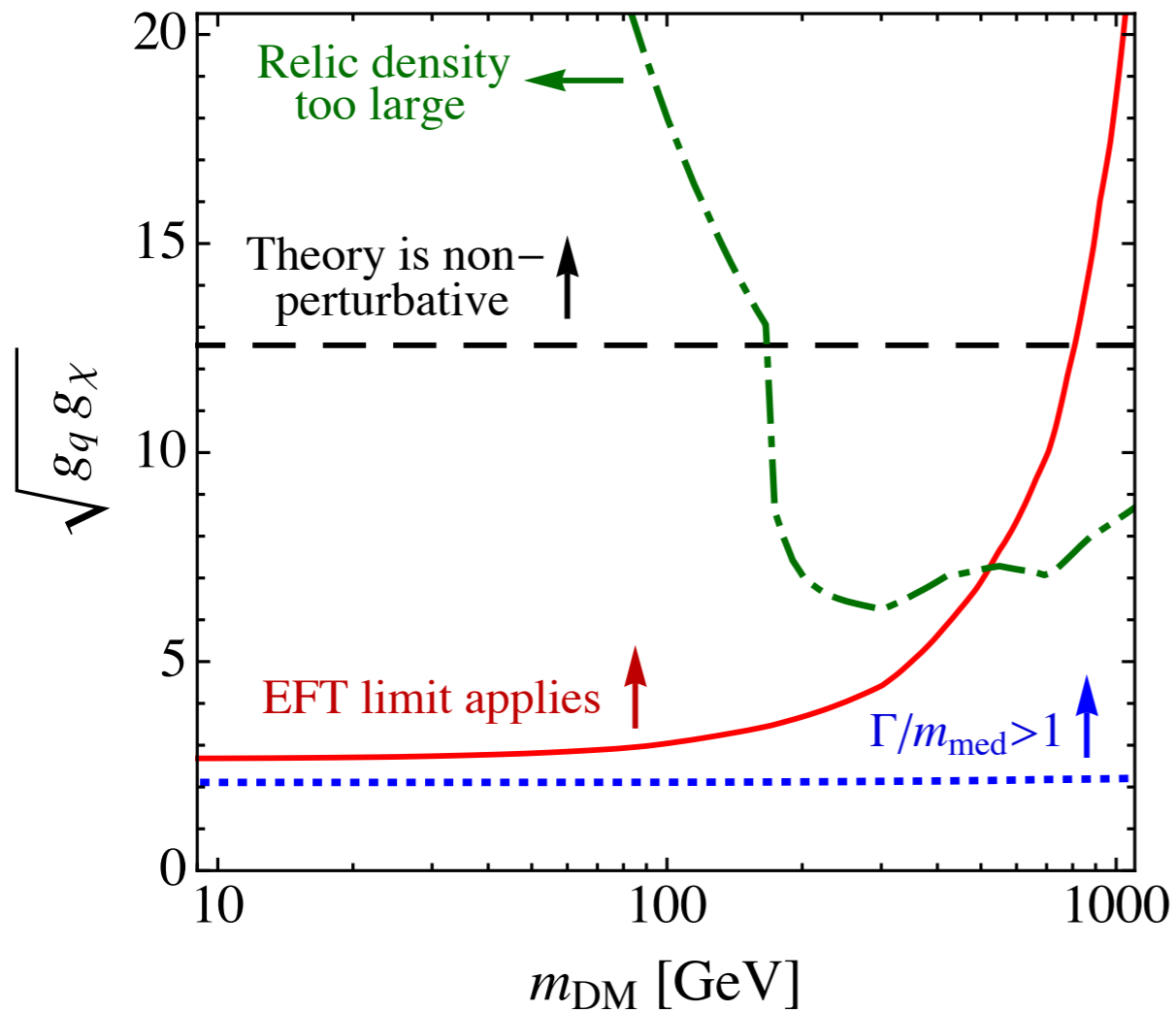
$$\begin{aligned} \mathcal{L}_{Z'} = & \sum g_{Z'\chi\bar{\chi}} \bar{\chi} \gamma^\mu \chi Z'^\mu + \sum g_{Z'q\bar{q}} \bar{q} \gamma^\mu q Z'^\mu \\ & + \bar{\chi} (i\partial_\mu \gamma^\mu - m_\chi) \chi + \text{gaugeterms} \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{A'} = & \sum g_{A'\chi\bar{\chi}} \bar{\chi} \gamma^\mu \gamma^5 \chi A'^\mu + \sum g_{A'q\bar{q}} \bar{q} \gamma^\mu \gamma^5 q A'^\mu \\ & + \bar{\chi} (i\partial_\mu \gamma^\mu - m_\chi) \chi + \text{gaugeterms} \end{aligned}$$

2

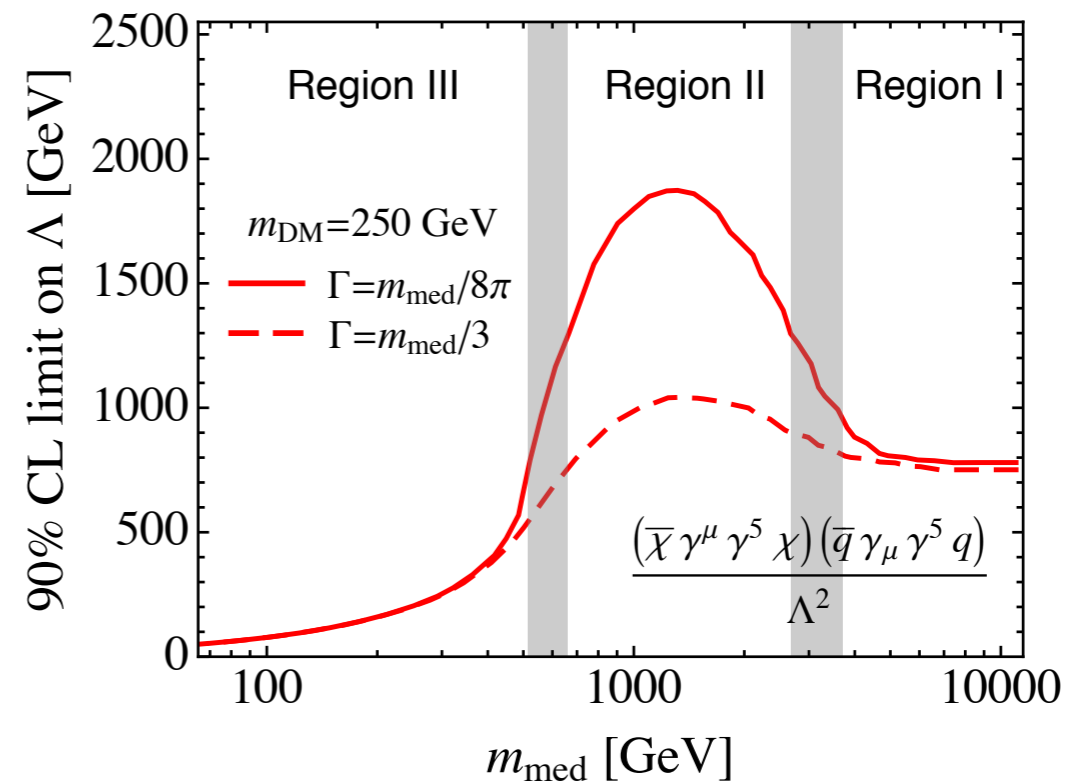
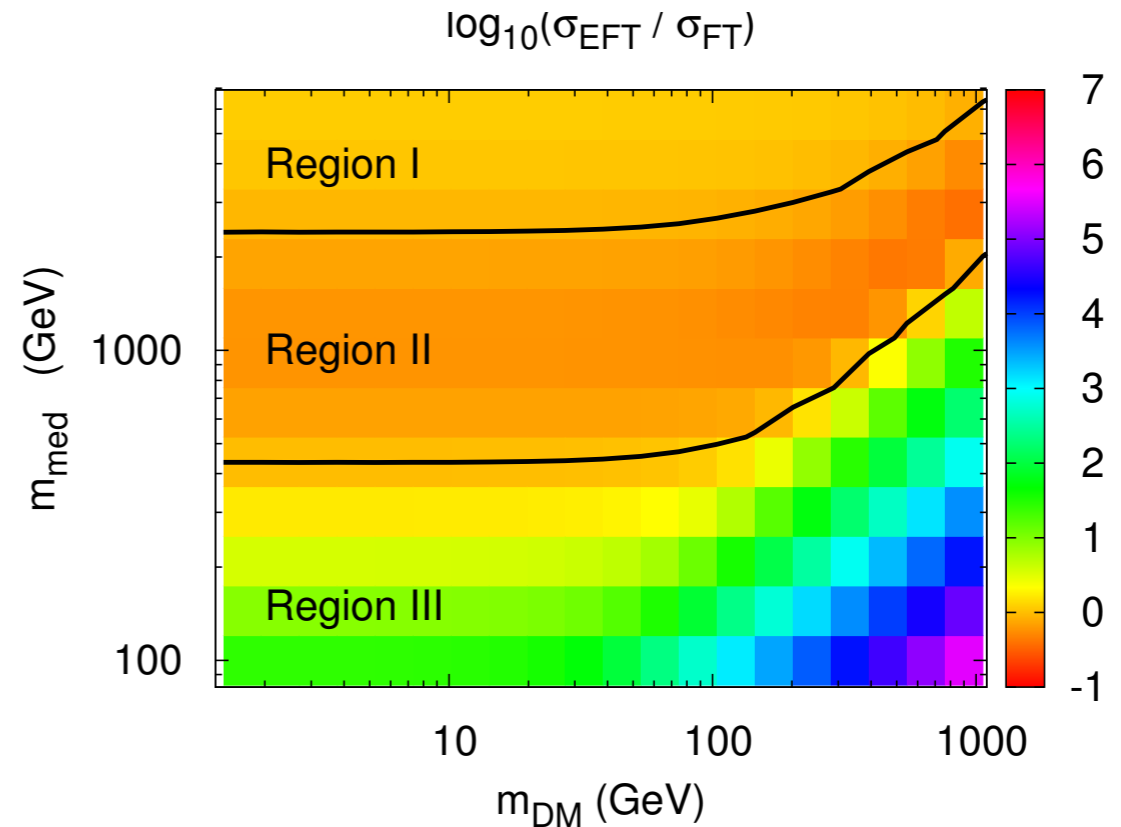
EXAMINING THE EFT-SIMP LINK MORE CLOSELY

arXiv:1308.6799 Buchmueller et al.



$$\frac{\Gamma}{m_{med}} = \frac{N_C g_f^2}{12\pi} \left(1 - \frac{4m_f^2}{m_{med}^2} \right)^{3/2}$$

$$\Lambda \equiv \frac{m_{med}}{\sqrt{g_q g_\chi}}$$



LOOKING FOR THE MEDIATOR

ATLAS Exotics Searches* - 95% CL Exclusion

Status: March 2016

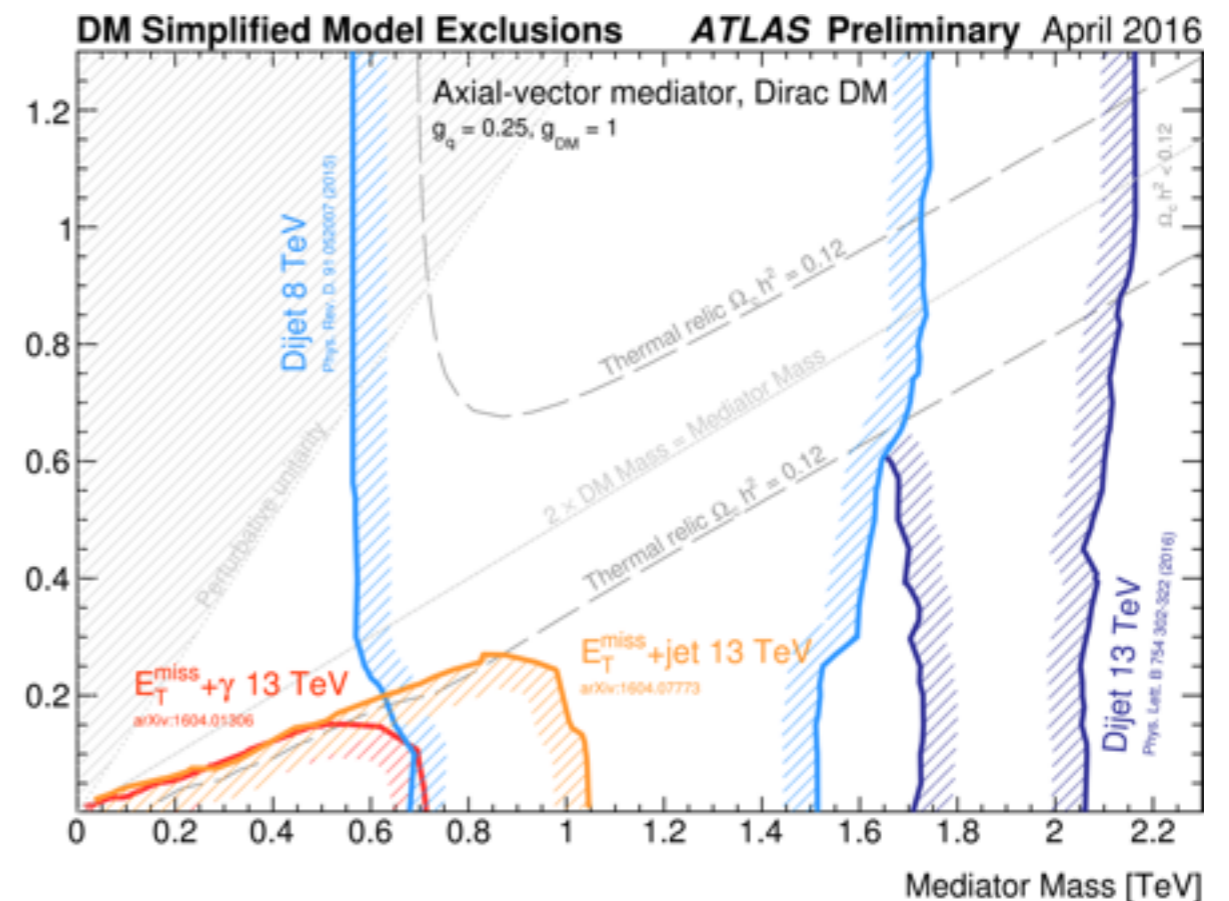
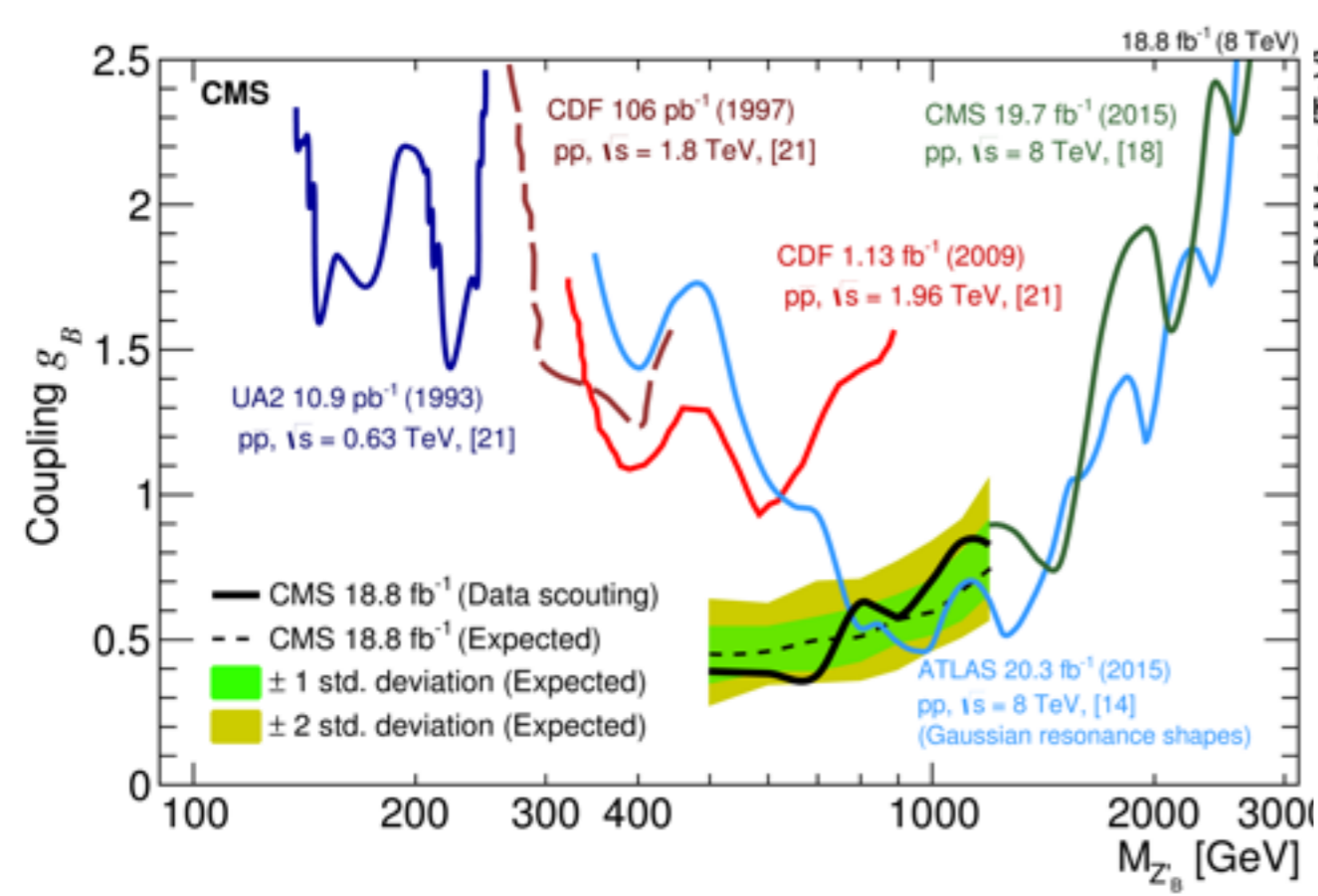
$\int \mathcal{L} dt = (3.2 - 20.3) \text{ fb}^{-1}$

ATLAS Preliminary

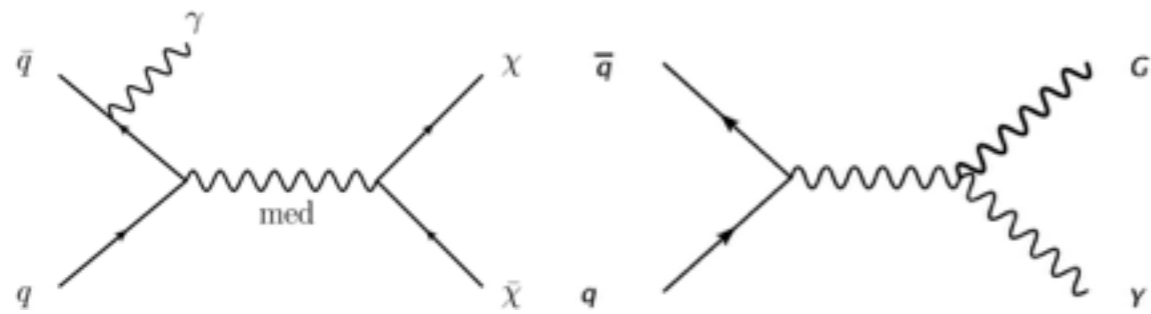
$\sqrt{s} = 8, 13 \text{ TeV}$

Gauge bosons	Search				Exclusion		Reference	
	Final State	Model	Signature	Mass [TeV]	Mass [TeV]	Mass [TeV]		
Gauge bosons	SSM $Z' \rightarrow \ell\ell$	$2 e, \mu$	-	-	3.2	Z' mass	3.4 TeV	ATLAS-CONF-2015-070 1502.07177 Preliminary ATLAS-CONF-2015-063 ATLAS-CONF-2015-068 ATLAS-CONF-2015-073 ATLAS-CONF-2015-074 ATLAS-CONF-2015-074 1410.4103 1408.0886
	SSM $Z' \rightarrow \tau\tau$	2τ	-	-	19.5	Z' mass	2.02 TeV	
	Leptophobic $Z' \rightarrow bb$	-	$2 b$	-	3.2	Z' mass	1.5 TeV	
	SSM $W' \rightarrow \ell\nu$	$1 e, \mu$	-	Yes	3.2	W' mass	4.07 TeV	
	HVT $W' \rightarrow WZ \rightarrow qq\nu\nu$ model A	$0 e, \mu$	$1 J$	Yes	3.2	W' mass	1.6 TeV	
	HVT $W' \rightarrow WZ \rightarrow qqqq$ model A	-	$2 J$	-	3.2	W' mass	1.38-1.6 TeV	
	HVT $W' \rightarrow WH \rightarrow \ell\nu bb$ model B	$1 e, \mu$	$1-2 b, 1-0 j$	Yes	3.2	W' mass	1.62 TeV	
	HVT $Z' \rightarrow ZH \rightarrow \nu\nu bb$ model B	$0 e, \mu$	$1-2 b, 1-0 j$	Yes	3.2	Z' mass	1.76 TeV	
	LRSM $W'_R \rightarrow tb$	$1 e, \mu$	$2 b, 0-1 j$	Yes	20.3	W' mass	1.92 TeV	
	LRSM $W'_R \rightarrow tb$	$0 e, \mu$	$\geq 1 b, 1 J$	-	20.3	W' mass	1.76 TeV	

$g_V = 1$
 $g_V = 1$
 $g_V = 3$
 $g_V = 3$

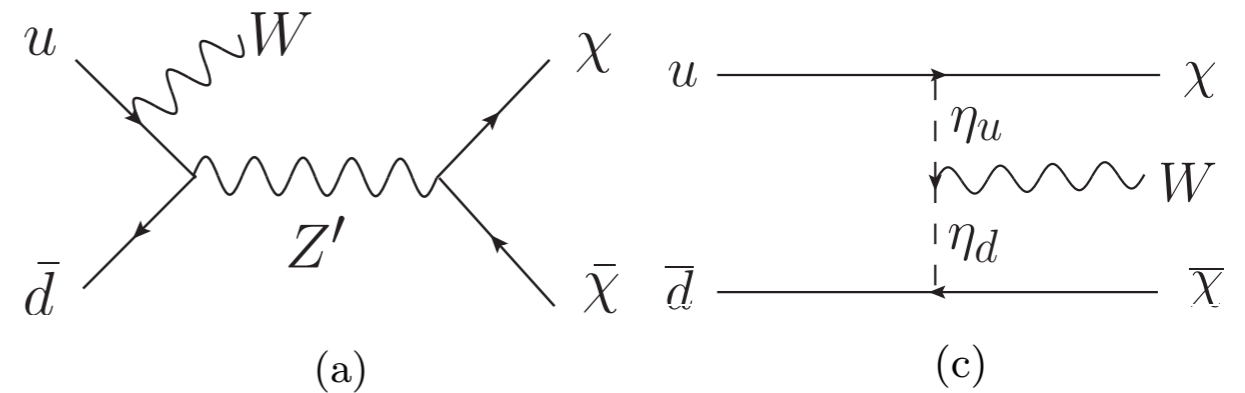


MONO-PHOTON



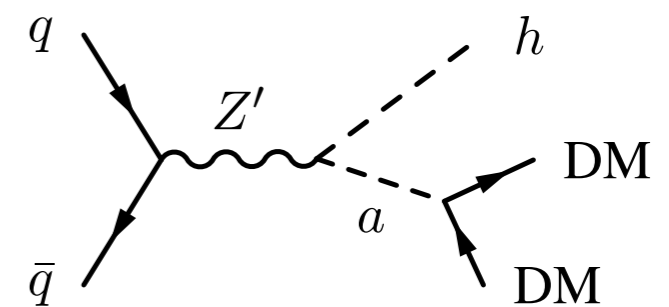
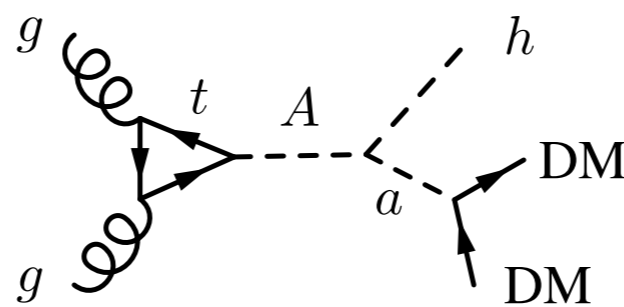
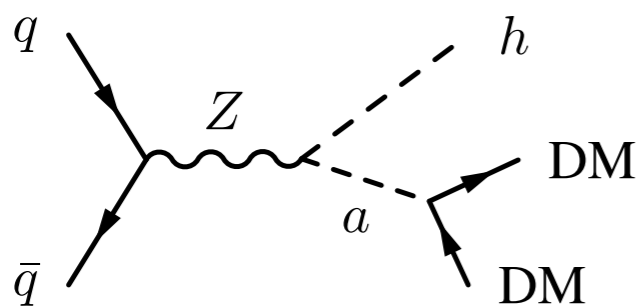
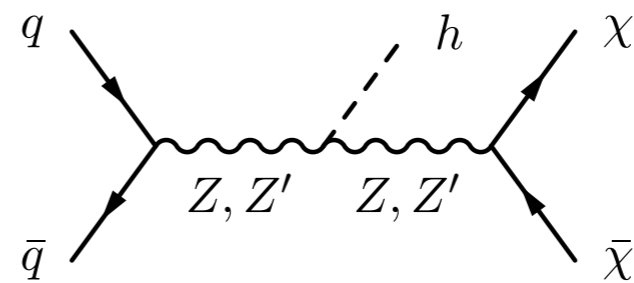
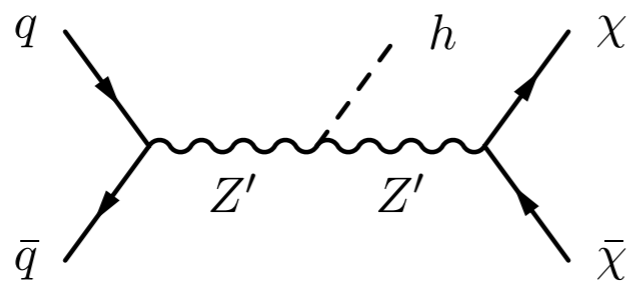
MONO-W

arXiv:1512.00476



MONO-H

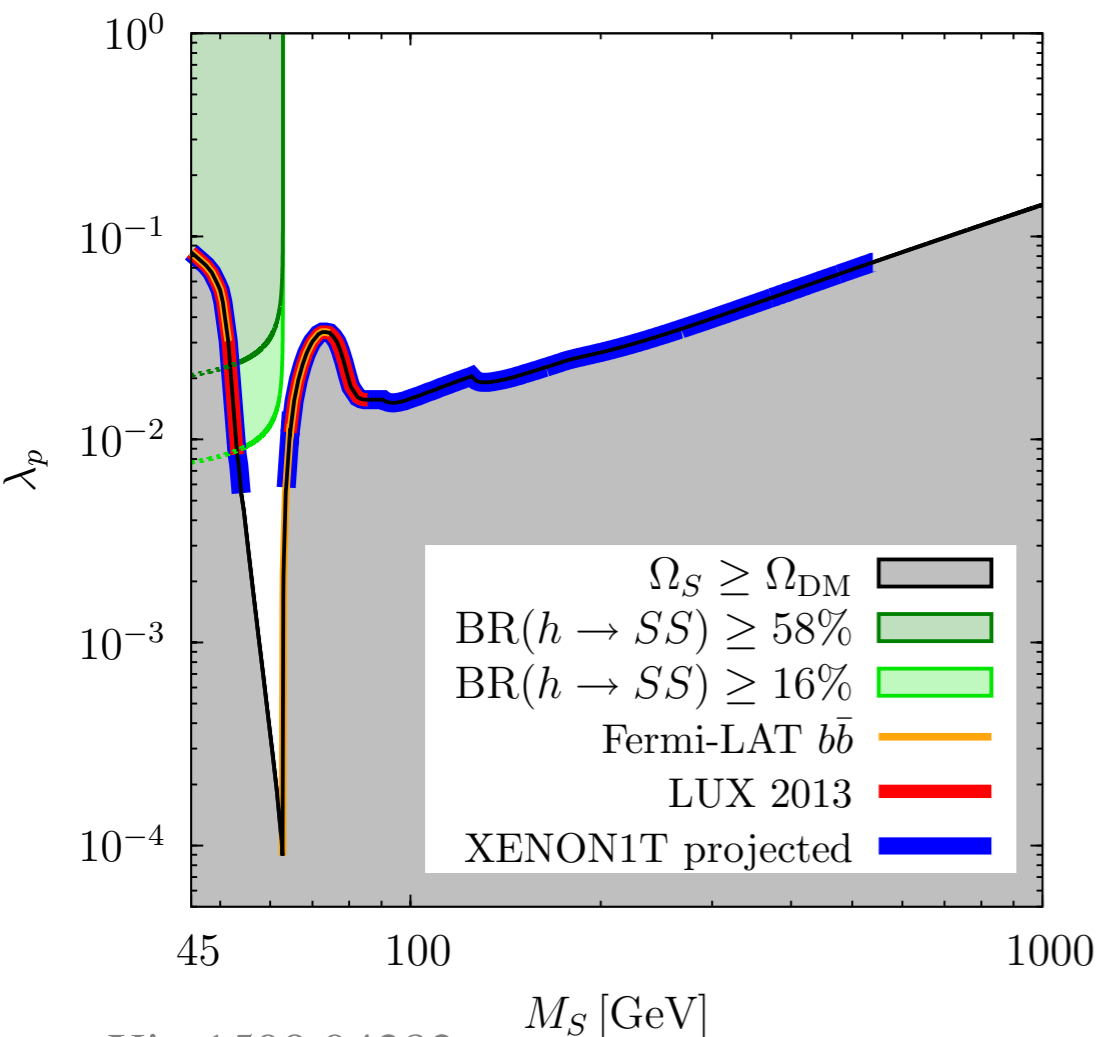
Carpenter et al. (2013)



HIGGS PORTAL DM

$$\mathcal{L}_{\text{scalar},H} \supset -\lambda_\chi \chi^4 - \lambda_p \chi^2 |H|^2,$$

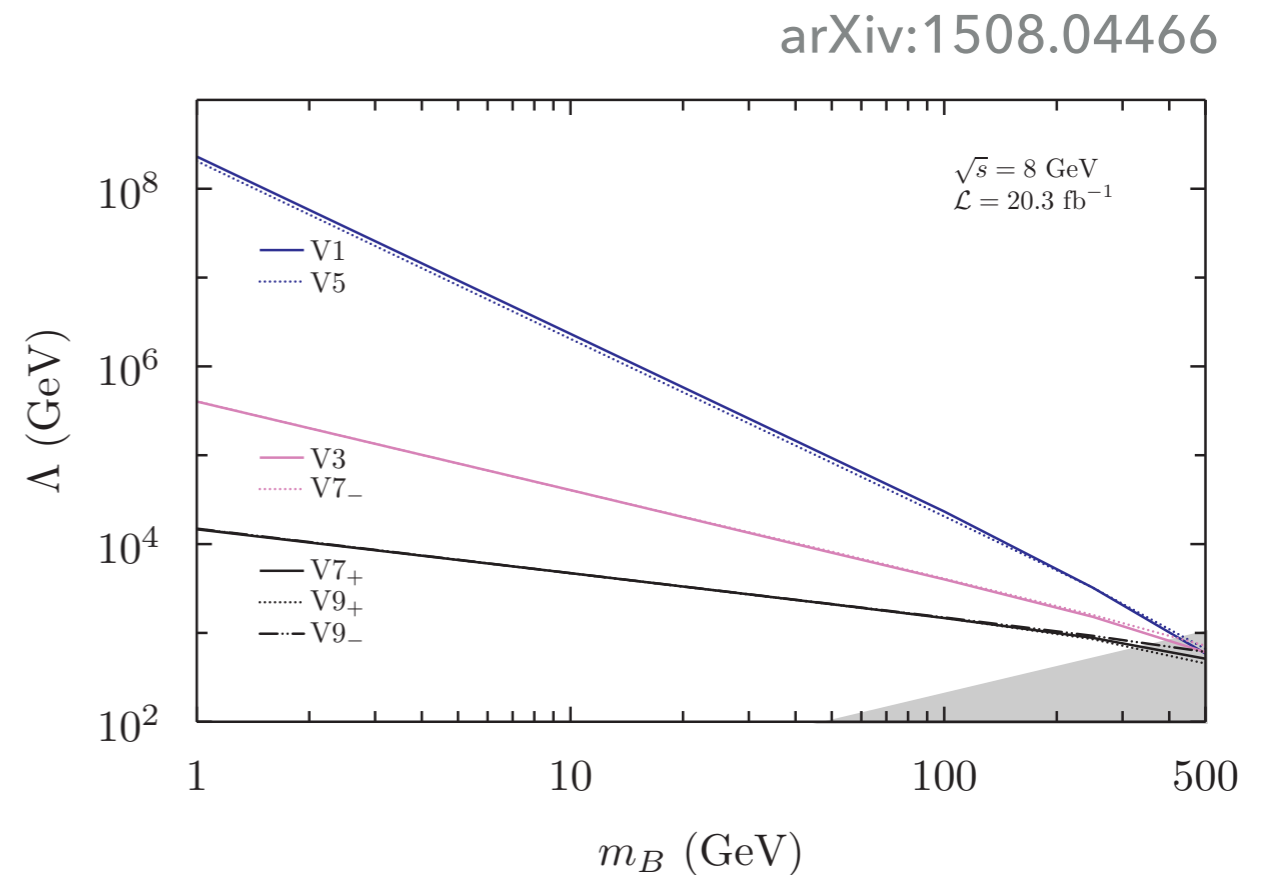
$$\Gamma(h \rightarrow \chi\chi) = \frac{\lambda_p^2 v^2}{2\pi m_h} \left(1 - \frac{4m_\chi^2}{m_h^2}\right)^{1/2}$$



arXiv:1509.04282

VECTOR DM

Operator	Structure	Dim D
V1	$(1/\Lambda) B_\mu^\dagger B^\mu \bar{q}q$	5
V2	$(1/\Lambda) i B_\mu^\dagger B^\mu \bar{q}\gamma^5 q$	5



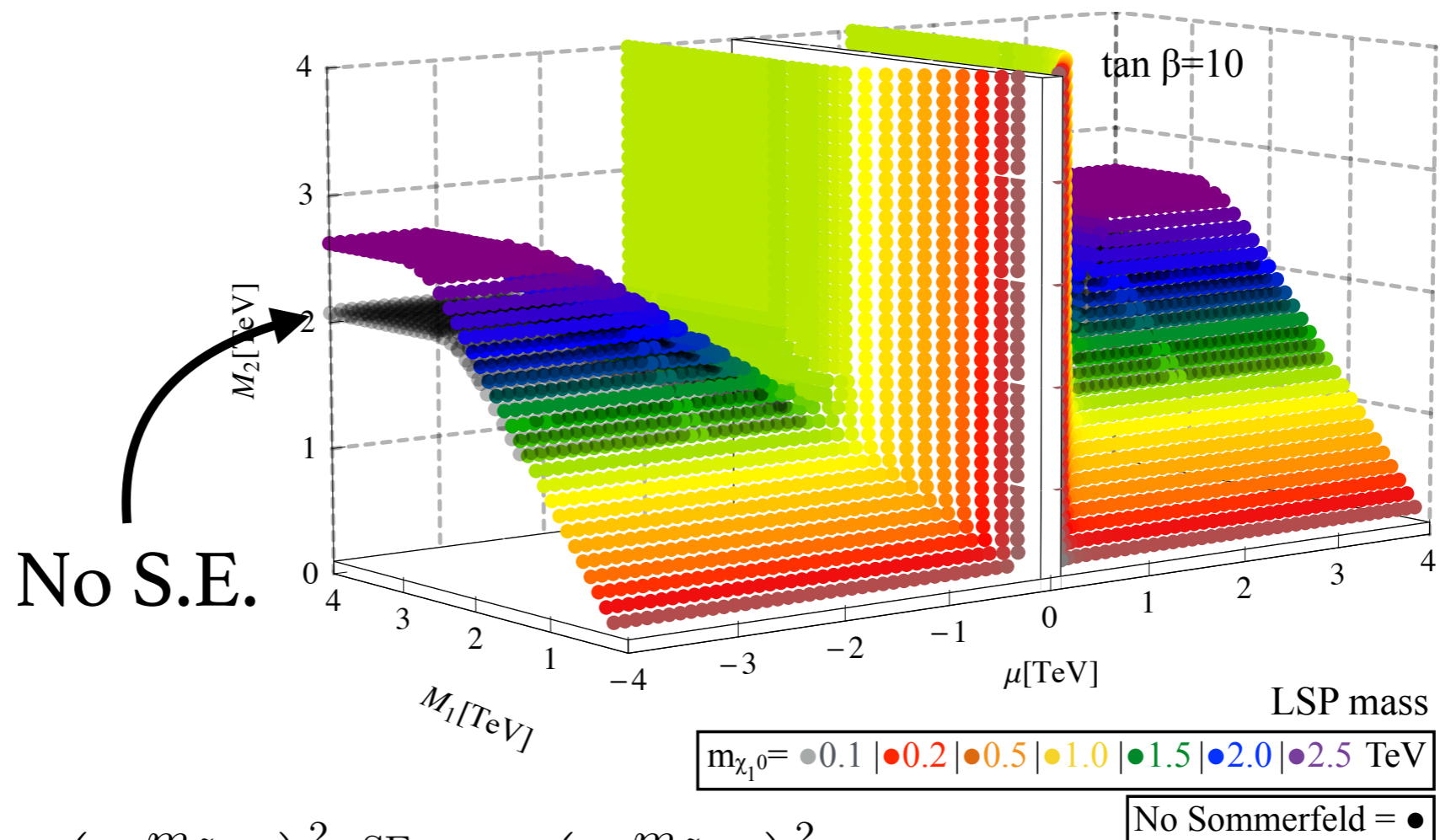
EXAMPLE OF A UV-COMPLETE MODEL: SUSY

- What kinds of interactions?
- Co-annihilation
- Sommerfeld enhancement
- Direct detection constraints
- Indirect detection constraints

$$g_{Z\tilde{\chi}_1^0\tilde{\chi}_1^0} = \frac{g}{2 \cos \theta_w} (|N_{13}|^2 - |N_{14}|^2)$$

$$g_{h\tilde{\chi}_1^0\tilde{\chi}_1^0} = (gN_{11} - g'N_{12}) (\sin \alpha N_{13} + \cos \alpha N_{14})$$

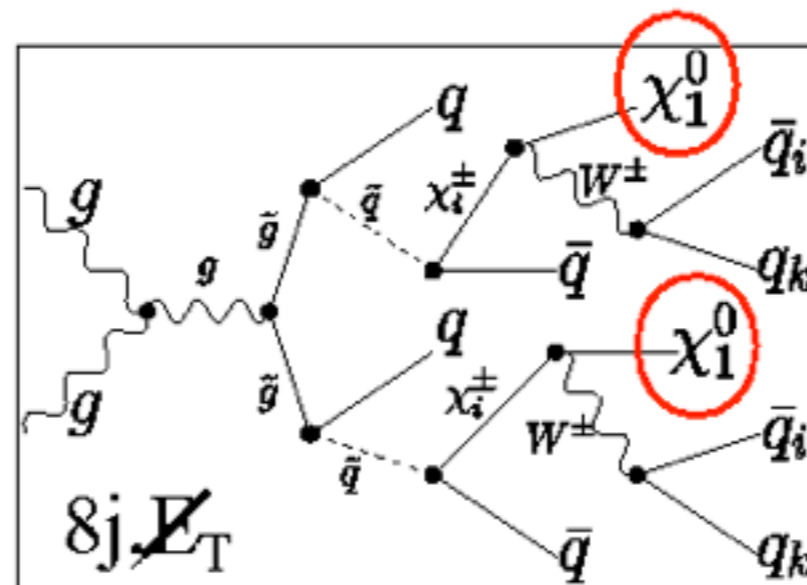
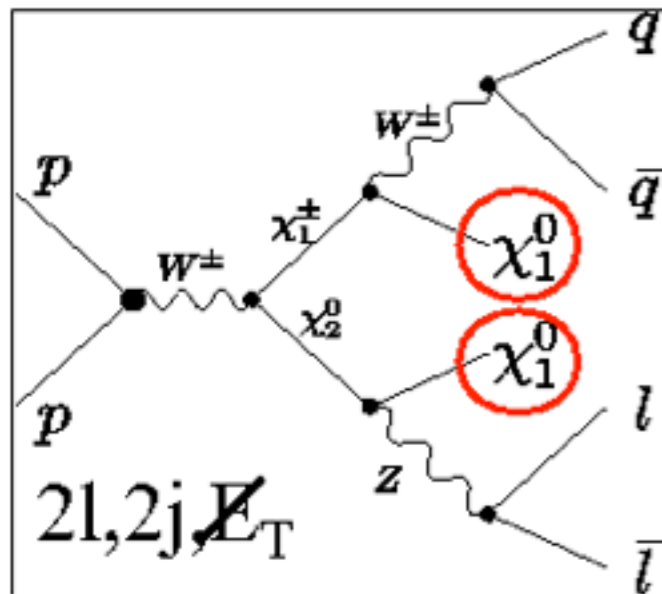
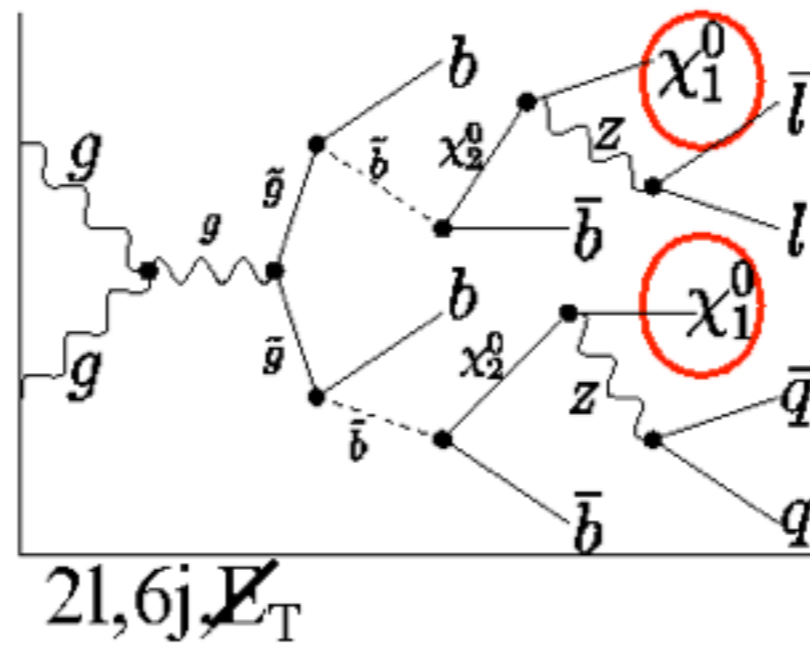
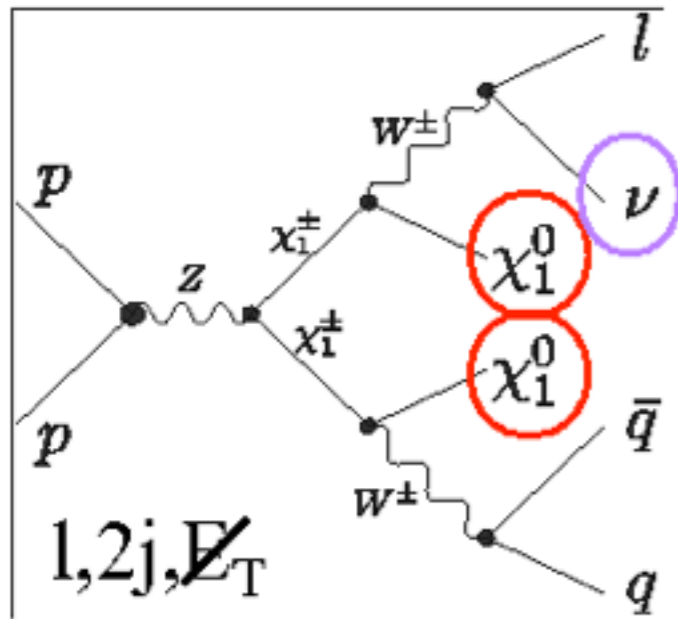
$$g_{W\tilde{\chi}_1^0\tilde{\chi}_1^+} = \frac{g \sin \theta_w}{\sqrt{2} \cos \theta_w} (N_{14}V_{12}^* - \sqrt{2}N_{12}V_{11}^*) ,$$



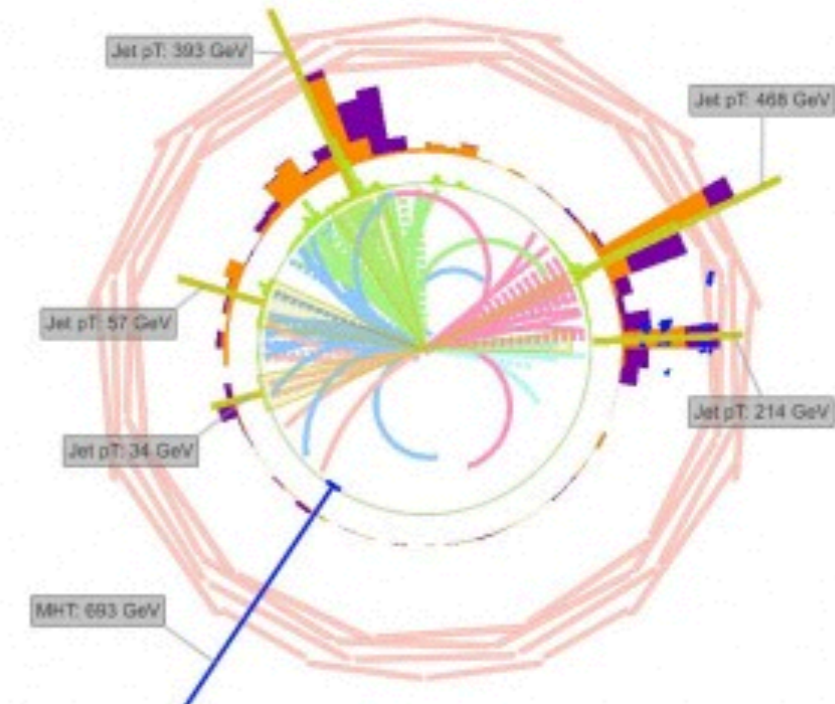
$$\Omega_{\tilde{W}} h^2 \simeq 0.12 \left(\frac{m_{\tilde{\chi}}}{2.1 \text{ TeV}} \right)^2 \xrightarrow{\text{SE}} 0.12 \left(\frac{m_{\tilde{\chi}}}{2.6 \text{ TeV}} \right)^2 .$$

$$\Omega_{\tilde{H}} h^2 \simeq 0.12 \left(\frac{m_{\tilde{\chi}}}{1.13 \text{ TeV}} \right)^2 \xrightarrow{\text{SE}} 0.12 \left(\frac{m_{\tilde{\chi}}}{1.14 \text{ TeV}} \right)^2 .$$

OBSERVING LSP PRODUCTION AT THE LHC: JETS + MET

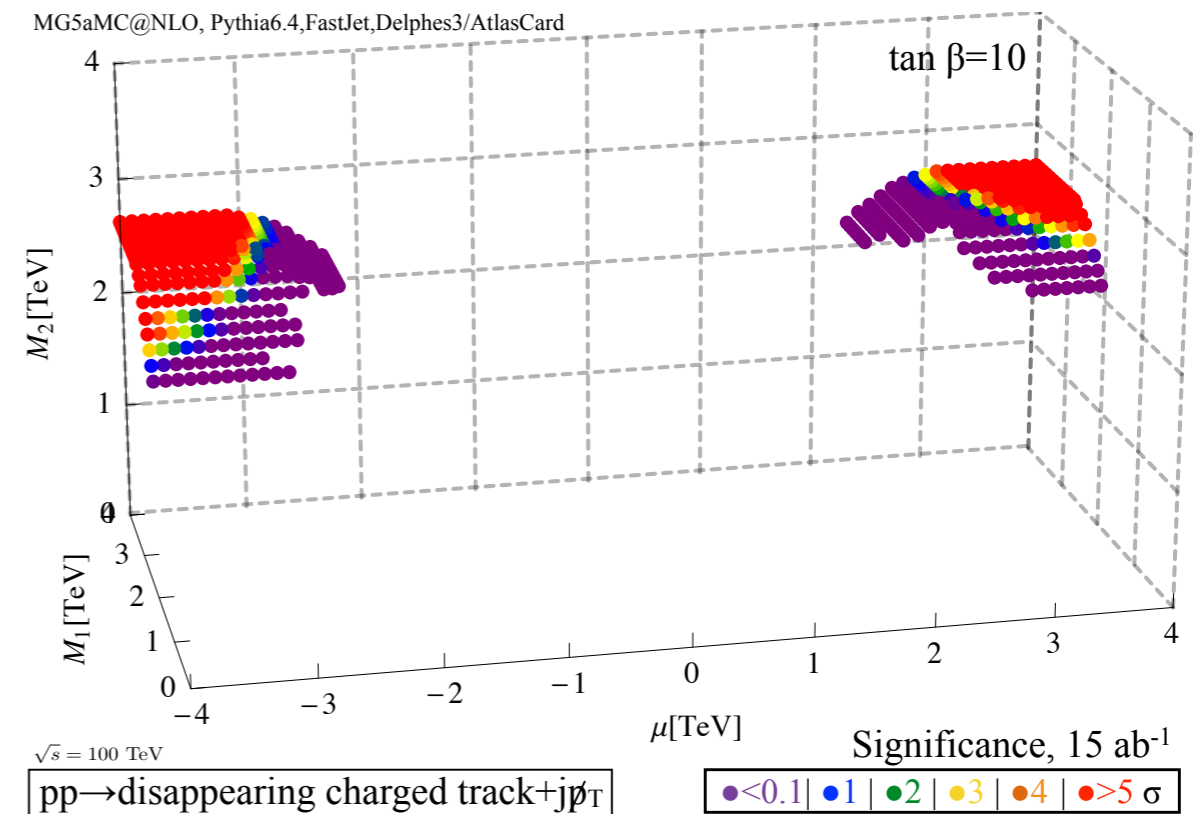
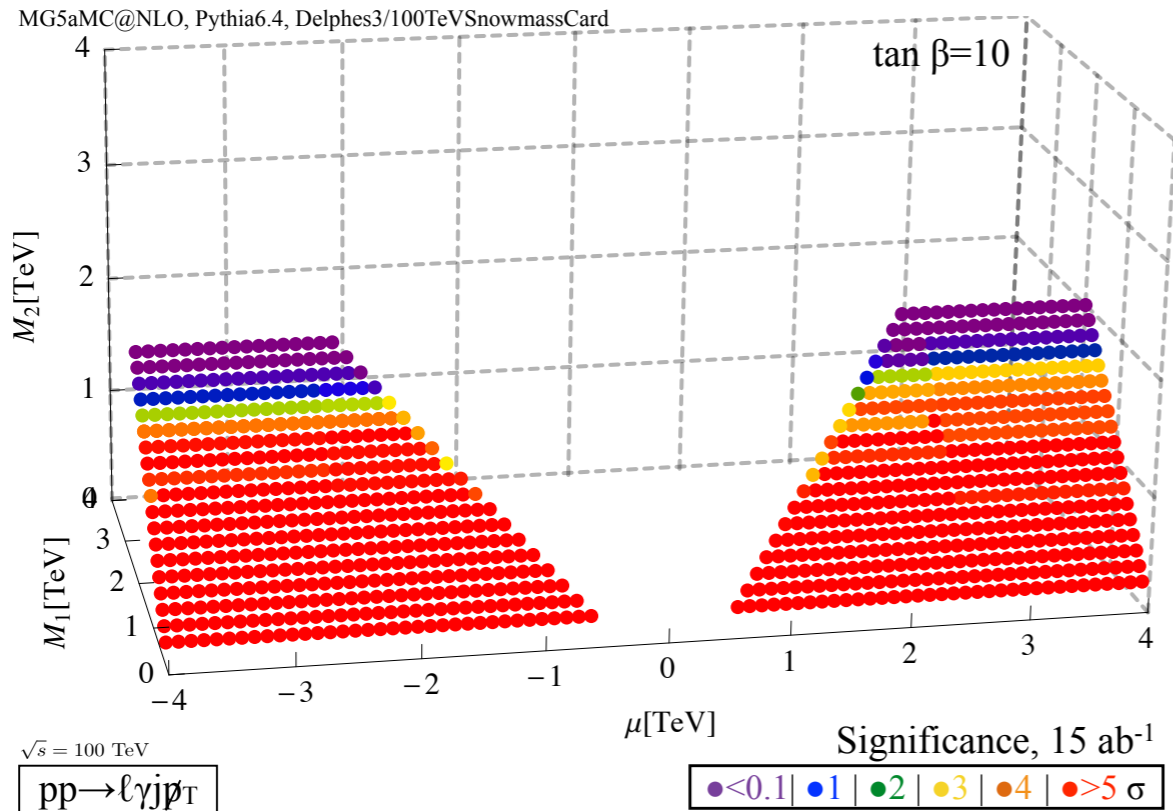


CMS
 CMS Experiment at LHC, CERN
 Data recorded: Tue Oct 26 07:13:54 2010 CEST
 Run/Event: 148953 / 70626194
 Lumi section: 49



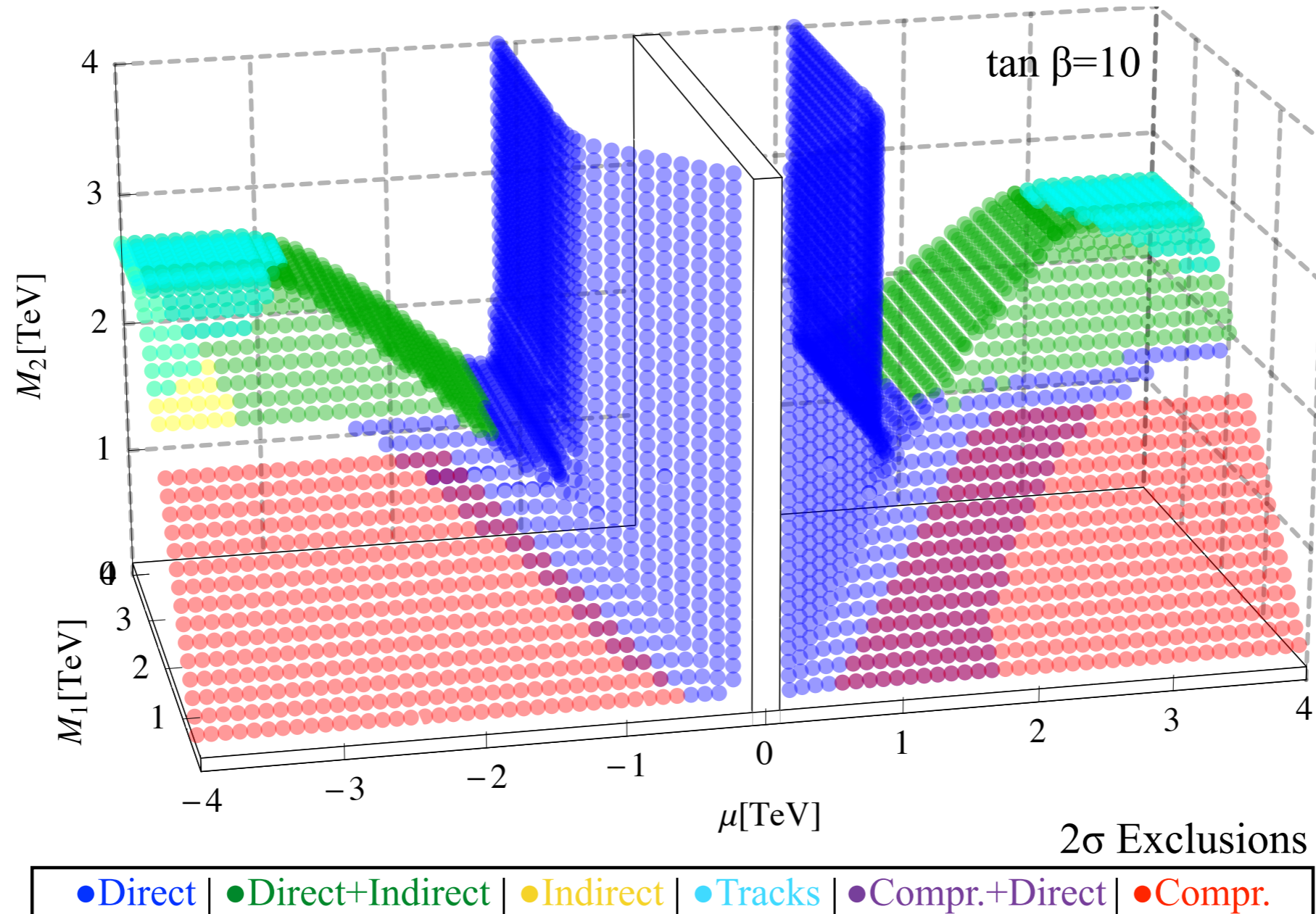
(POTENTIAL) COLLIDER SEARCHES

$$pp \rightarrow (\tilde{\chi}_2^0 \rightarrow \gamma \tilde{\chi}_1^0) (\tilde{\chi}_1^\pm \rightarrow \ell^\pm \nu_\ell \tilde{\chi}_1^0) j \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0 \ell^\pm \nu_\ell \gamma j ,$$

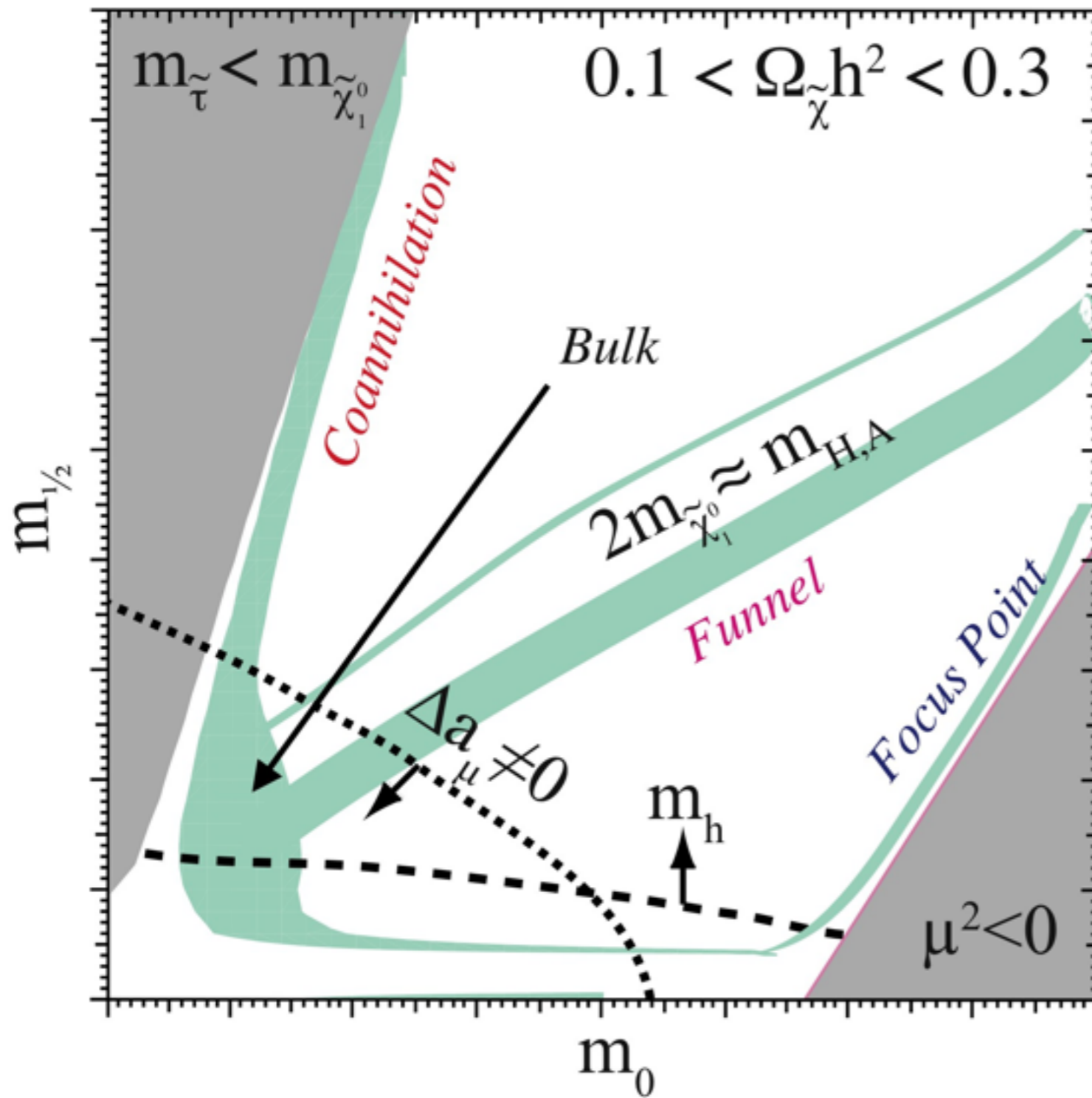


$$pp \rightarrow \chi_1^+ \chi_1^-, \chi_1^0 \chi_1^+$$

PUTTING IT ALL TOGETHER

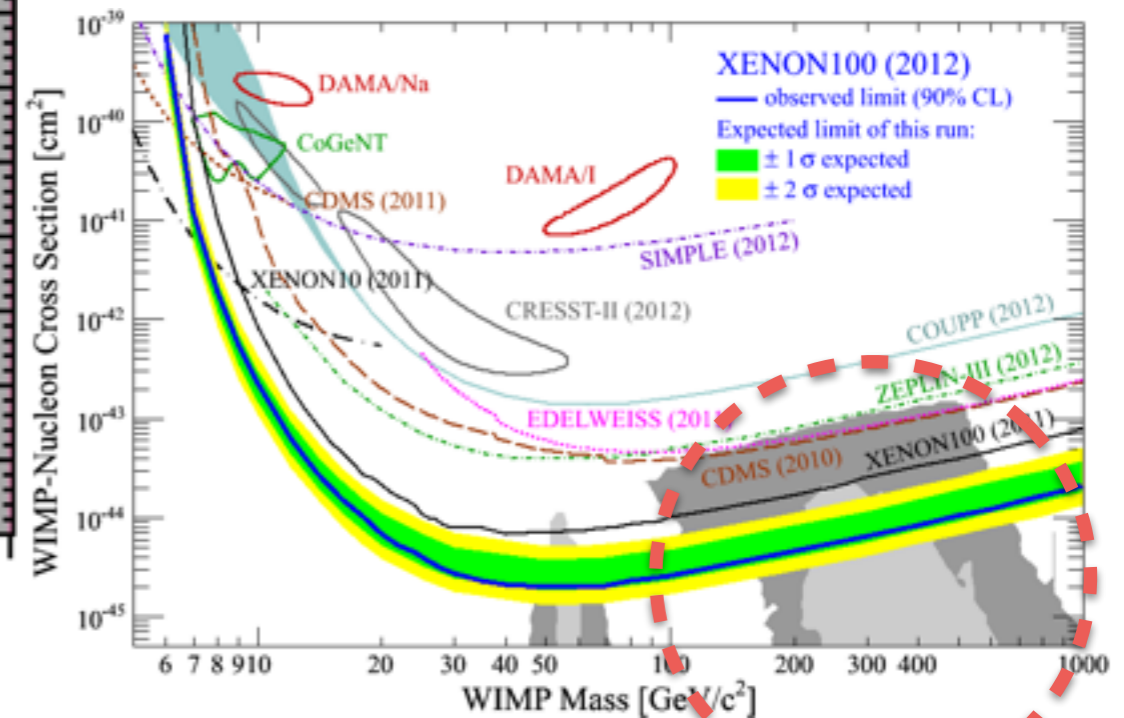


COANNIHILATION WITH SFERMIONS (STAUS)



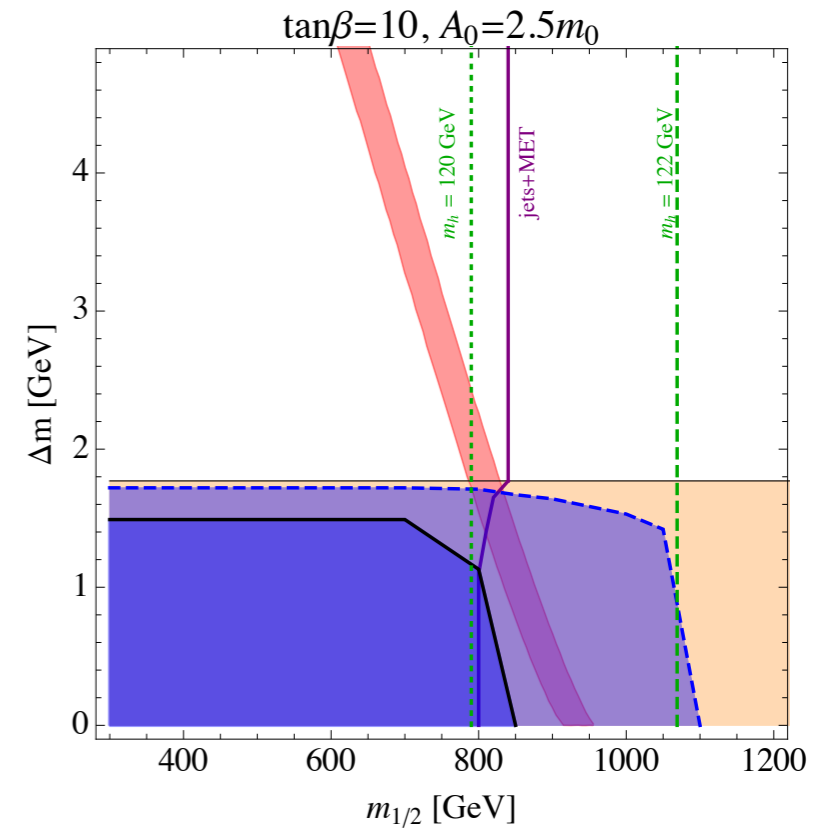
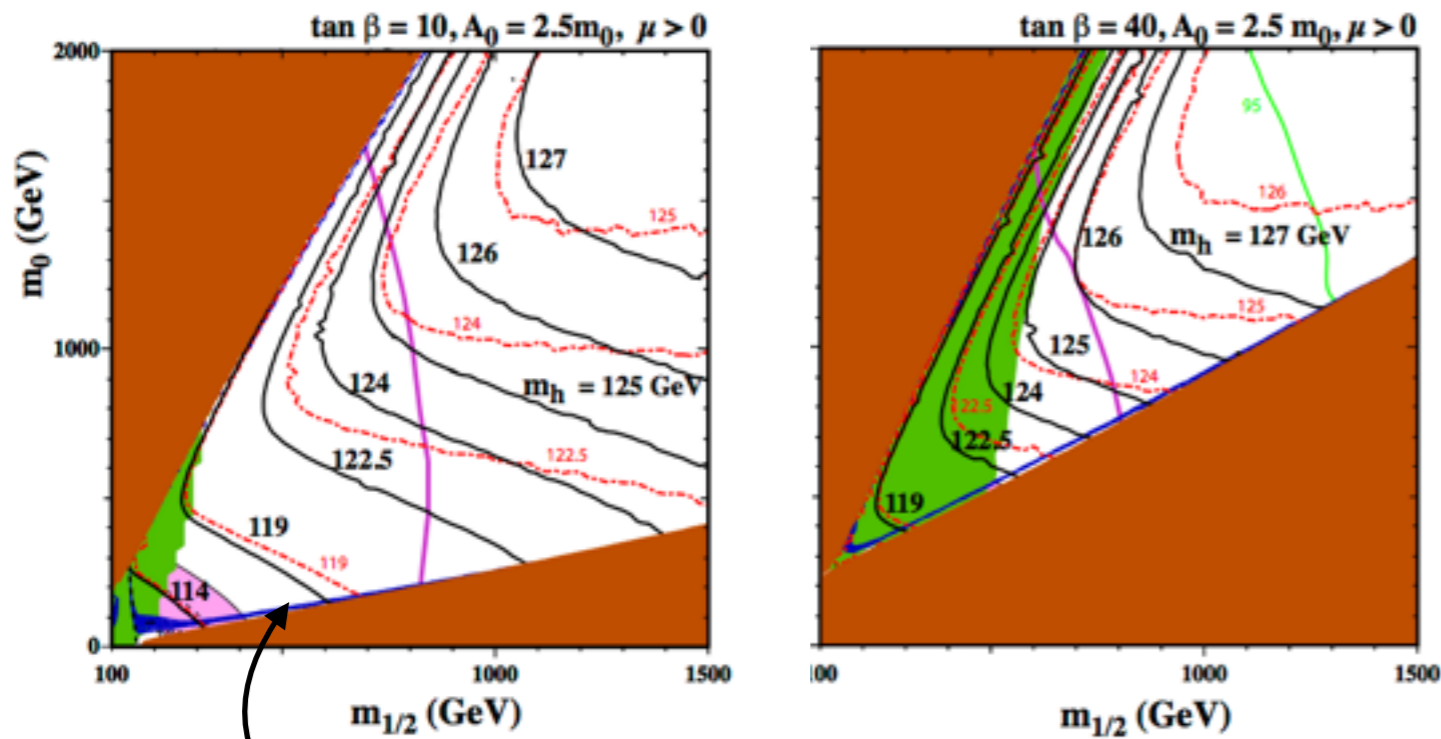
we use:

$$0.115 < \Omega_{\chi} h^2 < 0.125$$



http://www.nevis.columbia.edu/~mcooke/susy_page/susy.html

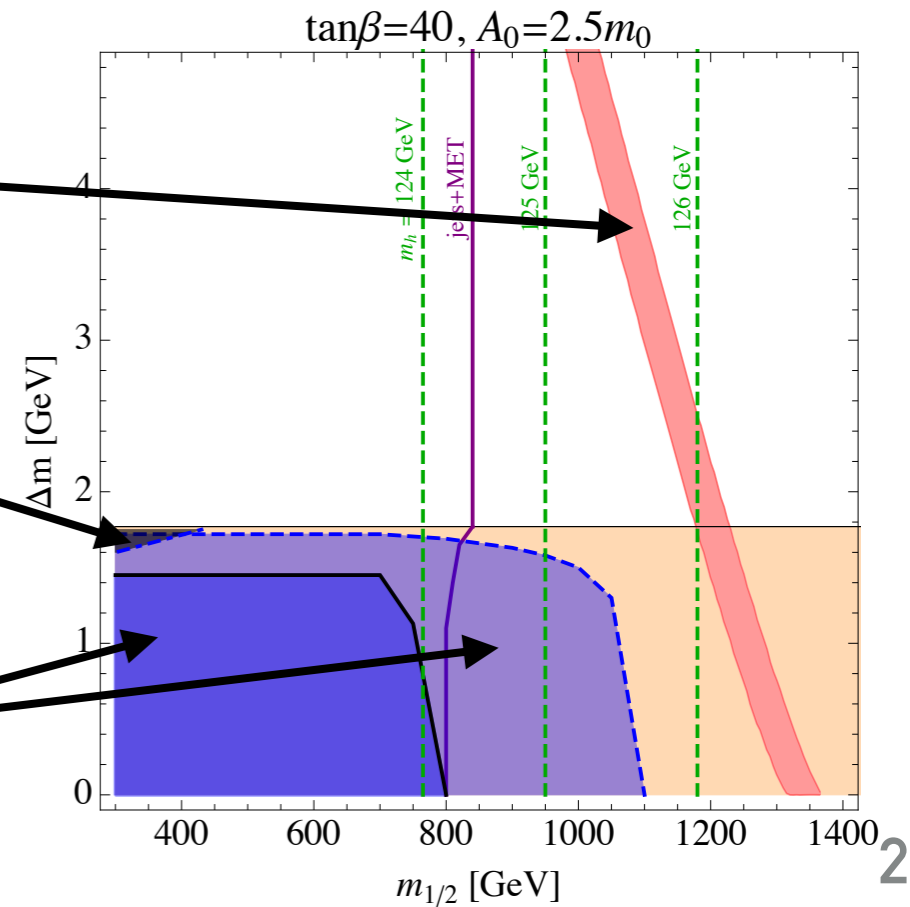
LIMITS FOR THE STAU-COANNIHILATION REGION



$m_\tau > \Delta m$
 = Long-lived or stable staus

Disappearing tracks

Charged tracks



- You can (and should) look for DM at the LHC
- EFT approach simple but not wholly unambiguous, tread with care.
- Simplified models are useful as a first strike, but also quite simplistic.
- Real UV complete theories may have more particles in the Dark Sector that one can probe.
- Much more to come in future ...