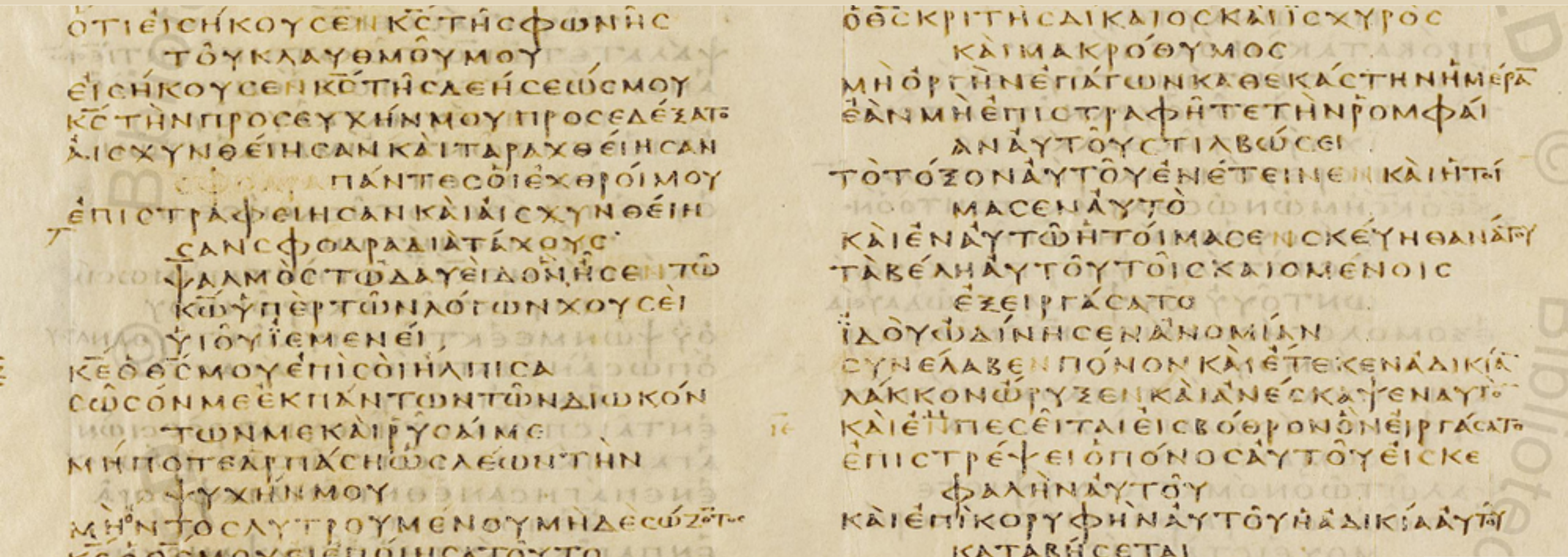


Searching for long lived particles

Why & how?



Vladimir V Gligorov

Thanks to S. Knapen, M. Papucci, D. Robinson, B. Dey for slides

Pizza seminar, Heidelberg

03-07-2018

Possibilities & Capabilities

Why long lived particle searches?

Long lifetimes arise from a **hierarchy of scales** or a **small coupling***

Three mechanisms:

- Off-shell decay
- Small splitting (phase space)
- Small coupling

Lessons from the SM:

- **generic** if there is more than one scale
- Often 3 body decays
- Weak theory prior on lifetime

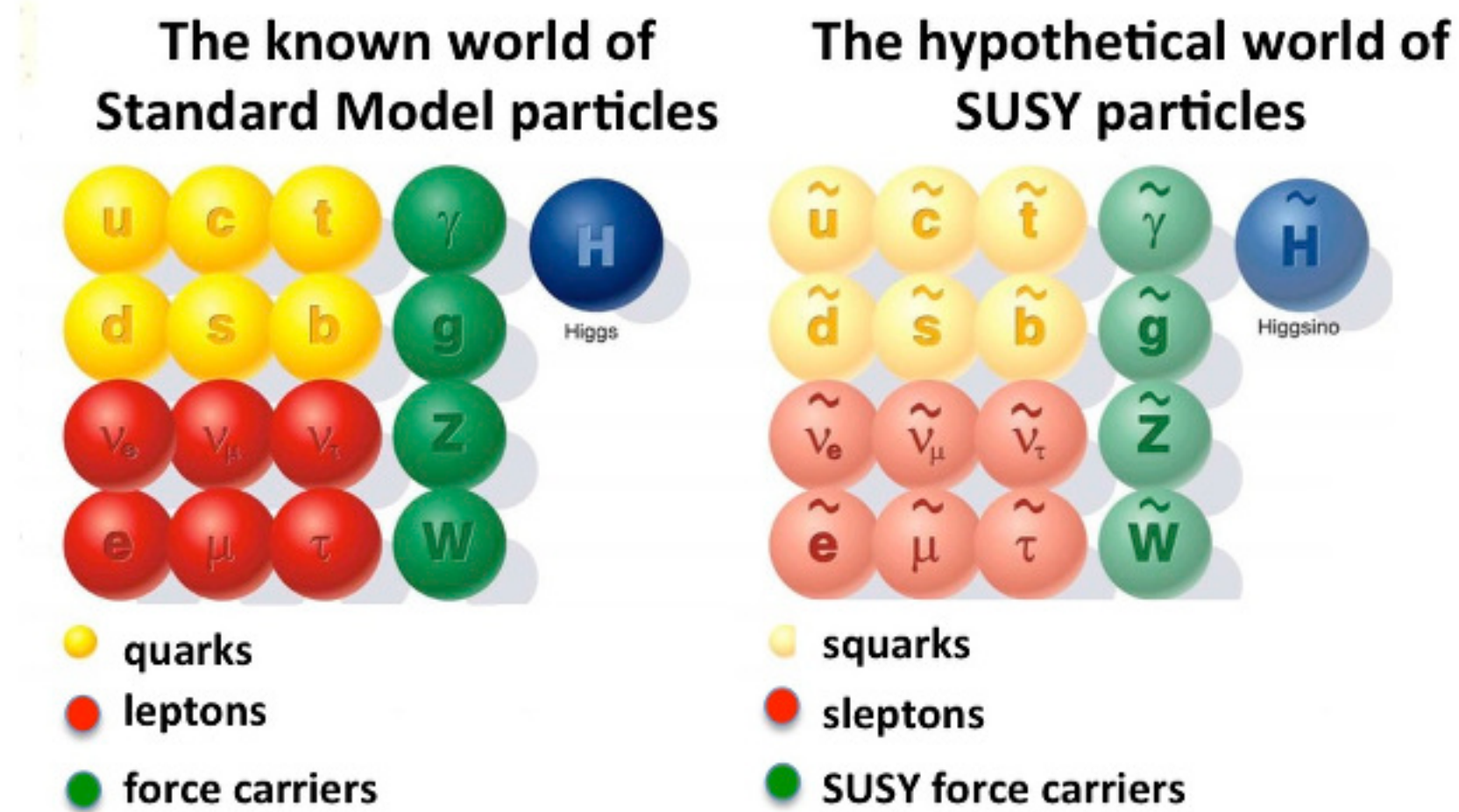
(e.g. proton decay!)

The diagram illustrates the decay rate formula $\Gamma \sim y^2 \left(\frac{m}{M}\right)^n m$. Annotations include:

- small coupling**: An arrow points from the coupling y^2 to this text.
- hierarchy of scales**: An arrow points from the scale ratio $\left(\frac{m}{M}\right)$ to this text.
- Set by symmetry structure, typically $n \geq 4$** : An arrow points from the exponent n to this text.

* could either be a hierarchy or loop suppression

Long-lived particles are generic



Other

R-parity violation
Gauge mediation
(mini-)split SUSY
stealth SUSY

Asymmetric Dark Matter
Freeze-in
composite Dark Matter
...

Baryogenesis
Neutrino masses
Neutral Naturalness
Hidden Valleys

A very wide range of BSM models introduce long-lived particles

LLP mass vs lifetime vs production

broken sym
weak mixing/ marginal operator
technically natural

$$\Gamma \sim \varepsilon^2 \left(\frac{m}{M} \right)^n \text{PS}$$

$m \ll M$, typically $n \geq 4$
loop factors

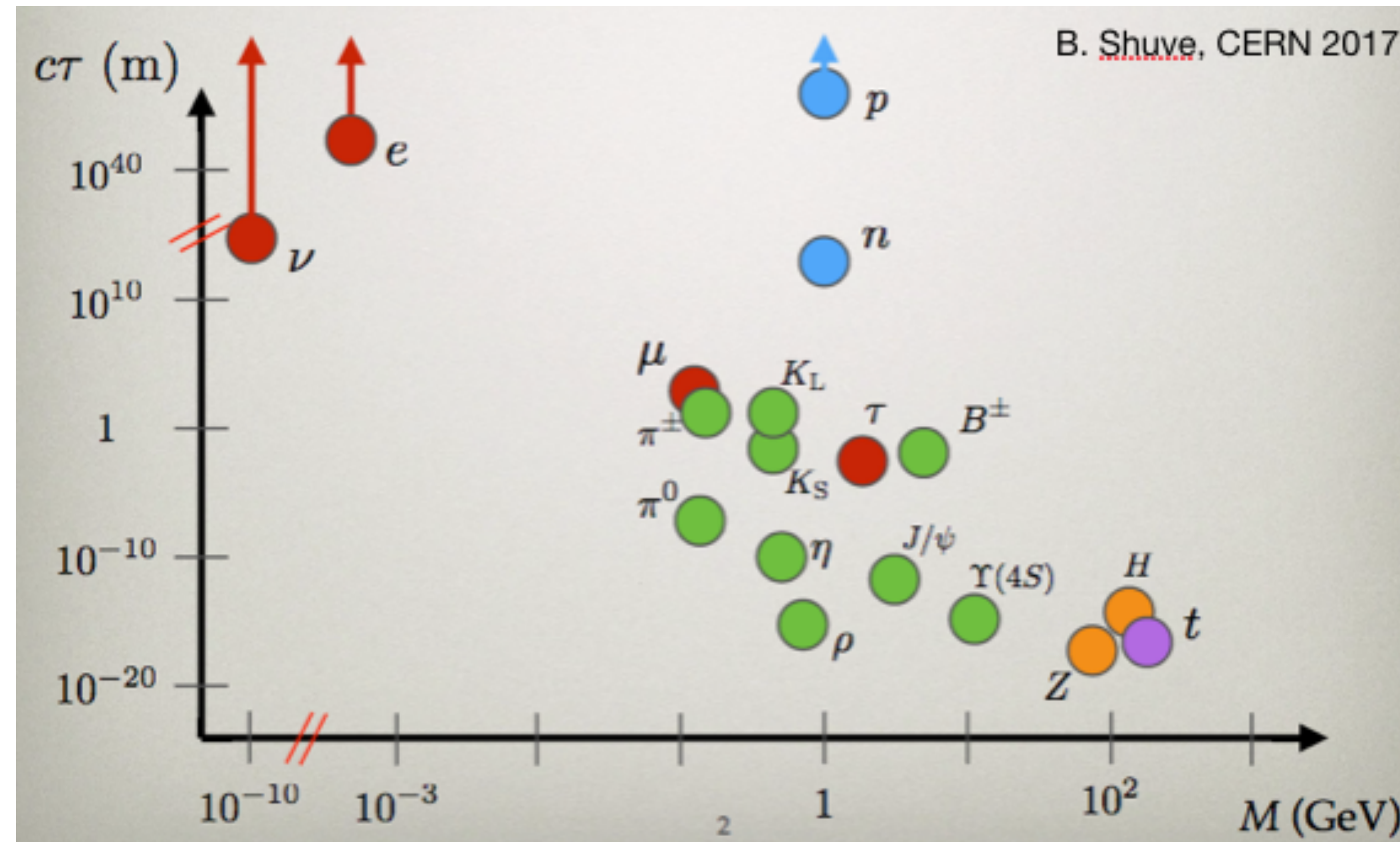
squeezed spectra
approx sym
multibody decays

The diagram illustrates the decay width Γ as a function of mass m and scale M . The formula is $\Gamma \sim \varepsilon^2 \left(\frac{m}{M} \right)^n \text{PS}$. Annotations explain the terms: ε^2 is linked to 'broken sym', 'weak mixing/ marginal operator', and 'technically natural'; $\left(\frac{m}{M} \right)^n$ is linked to ' $m \ll M$, typically $n \geq 4$ ' and 'loop factors'; and 'PS' is linked to 'squeezed spectra', 'approx sym', and 'multibody decays'.

The bigger the mass, the smaller in general the coupling you have to impose to get a narrow width (long lifetime)

The details linking production and decay in this heavily depend on the specific LLP and the portal used to access it.

LLP mass vs lifetime vs production



The bigger the mass, the smaller in general the coupling you have to impose to get a narrow width (long lifetime)

The details linking production and decay in this heavily depend on the specific LLP and the portal used to access it.

So how do we search for them?

No theory guidance on lifetime → large detectors

Many possible decay modes → hermeticity, particle ID

Small coupling and production rate → zero background

Small coupling and production rate → huge integrated lumi

Very hard for any single detector to meet all these criteria!

Collider vs. fixed target mode

Fixed target

Collider

Advantages

Disadvantages

Collider vs. fixed target mode

Fixed target

Collider

Advantages

Production rate

Collimated

production & decay

Disadvantages

Collider vs. fixed target mode

Fixed target

Collider

Advantages

**Production rate
Collimated
production & decay**

Disadvantages

**No access to very
heavy LLPs
Big shielding
required for bkg**

Collider vs. fixed target mode

Fixed target

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Access to higher
mass LLPs via e.g.
Higgs portal

Disadvantages

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Collider vs. fixed target mode

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Disadvantages

No access to very
heavy LLPs
Big shielding
required for bkg

Uncollimated
production
Hard to instrument
Hard to shield

Collider vs. fixed target mode

To put the production argument in some context,
consider the SPS vs. HL-LHC, each over 5 years

Charm Hadrons @ SPS : $O(10^{18})$

Charm Hadrons @ HL-LHC : $O(10^{16})$

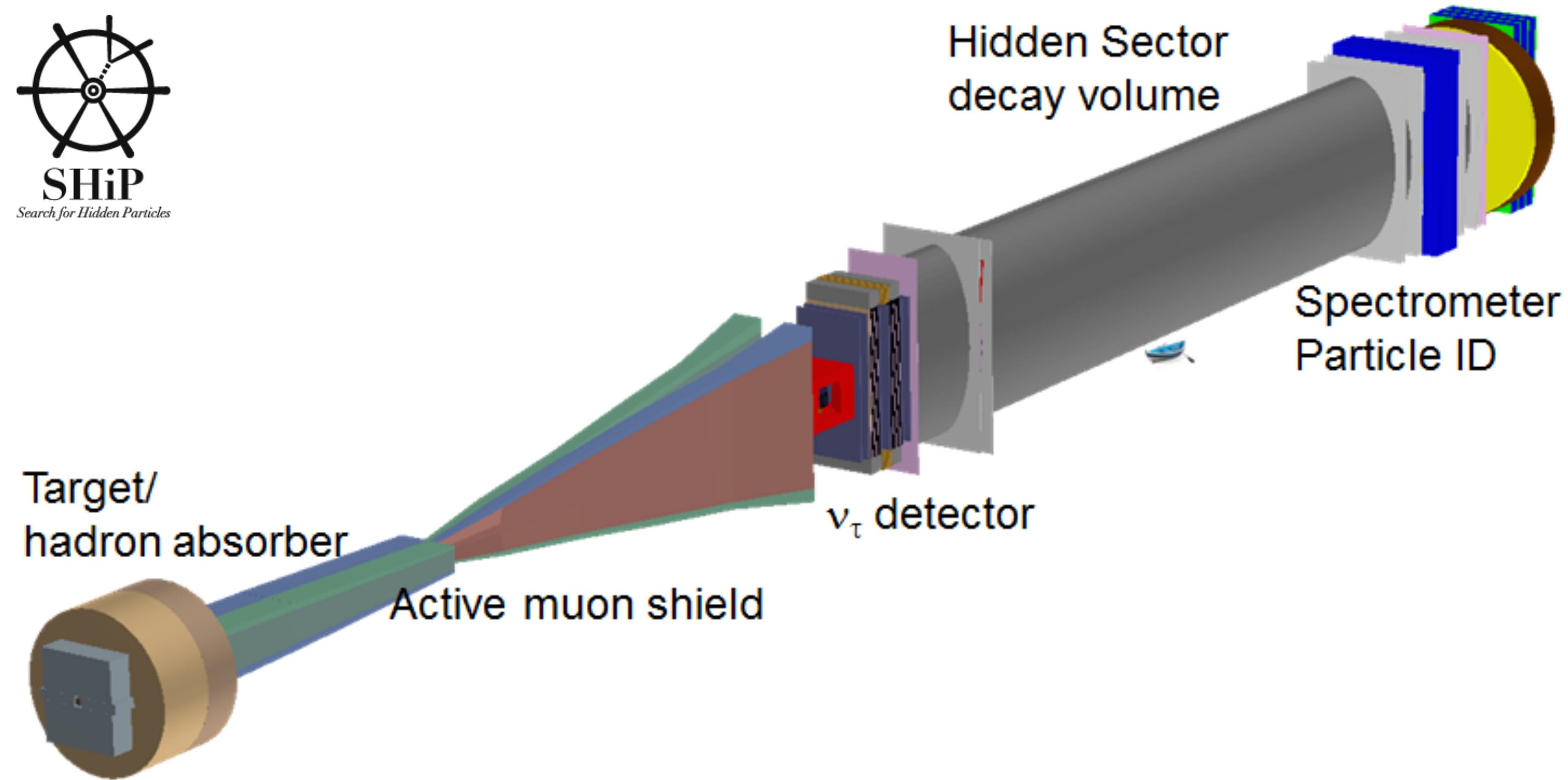
Beauty Hadrons @ SPS : $O(10^{14})$

Beauty Hadrons @ HL-LHC : $O(10^{15})$

This is why SHIP is so great at LLPs produced in
charm decays, while HL-LHC can compete for
beauty and dominates for anything heavier

Distance versus solid angle coverage

Fixed target : collimated production

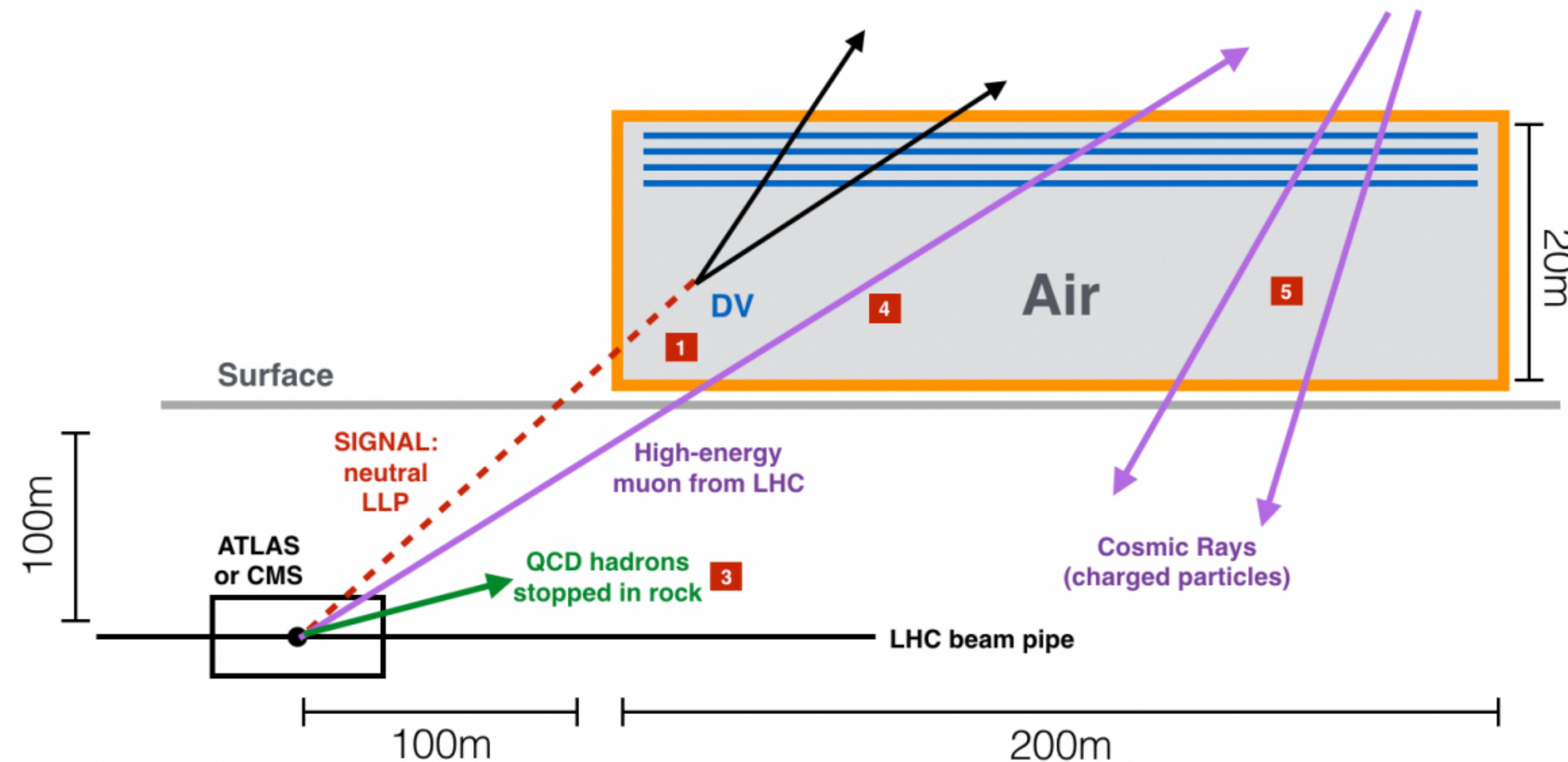


Collimated production and decay mean that solid angle coverage is largely independent of optimal decay volume. The geometry is dominated more by the required size of shield.

Distance versus solid angle coverage

Collider mode : solid angle is critical!

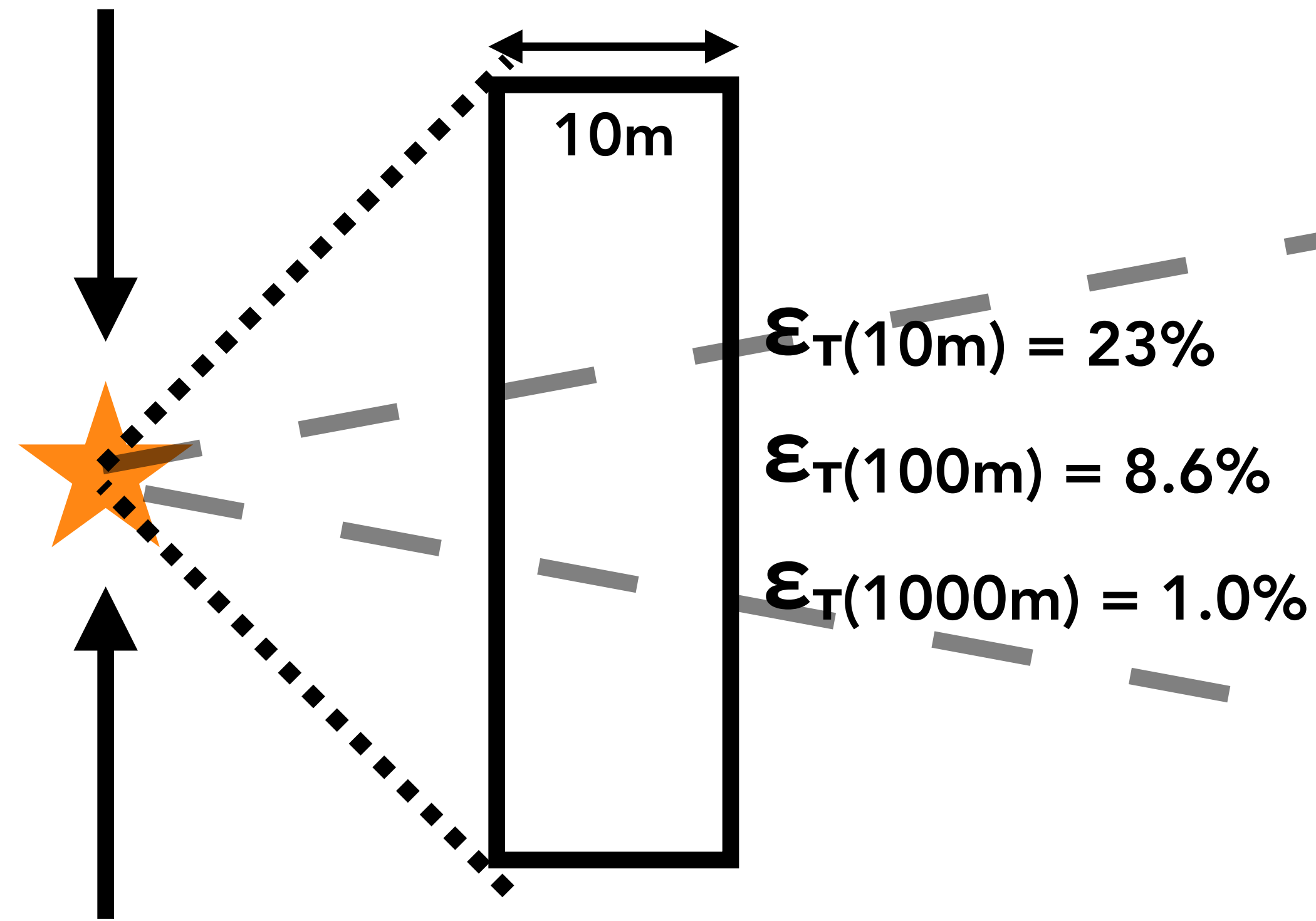
~~MATHUSLA~~



Uncollimated production means that unless you go forward like FASER, the size of your detector goes quadratically with the distance from collision point. Hence MATHUSLA's 200x200 m²...

Distance versus lifetime coverage

10 m from IP



50 m from IP

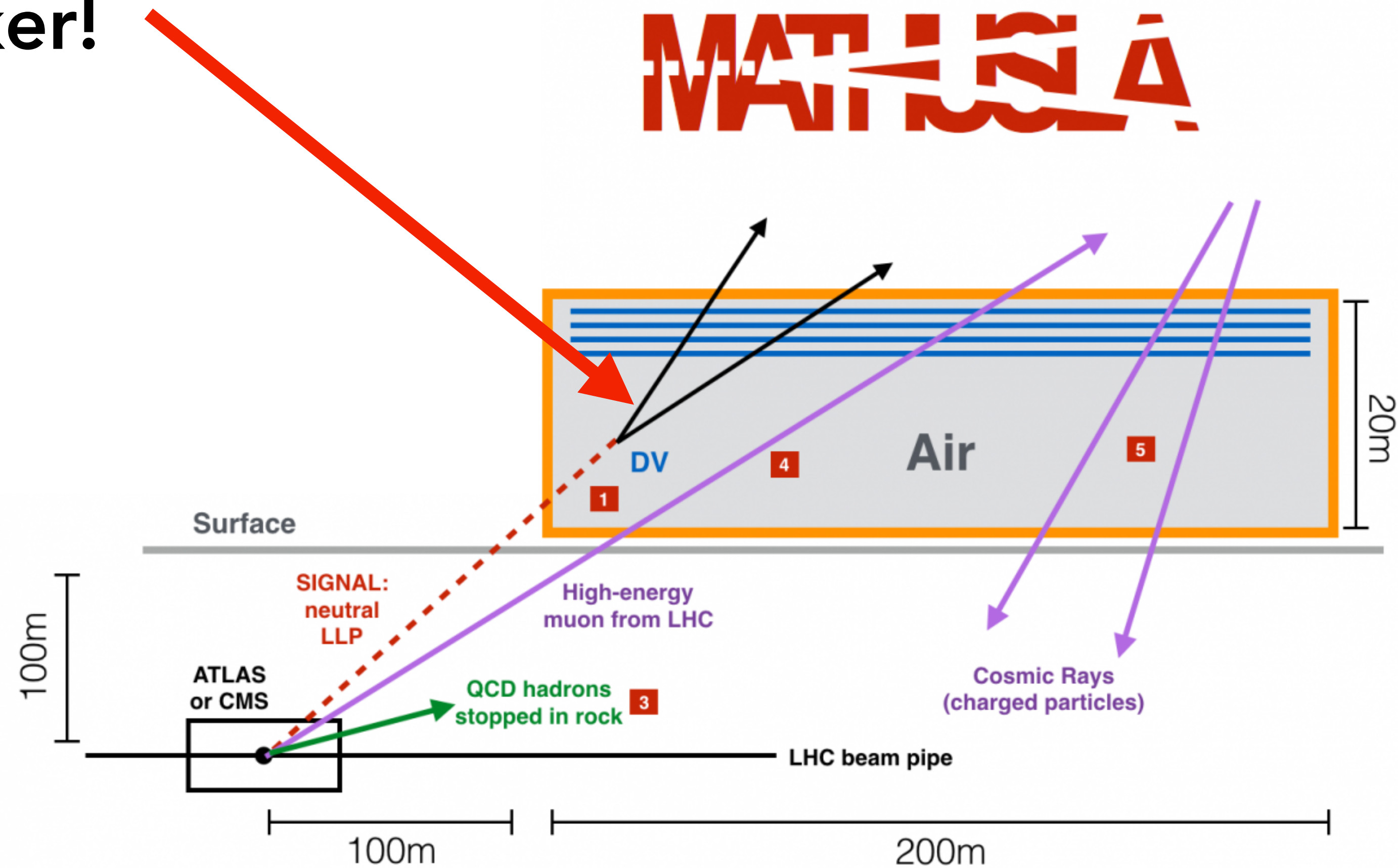
A diagram showing a detector represented by a vertical rectangle. A dashed line representing a particle trajectory passes through an orange star (the Interaction Point, IP) and enters the detector. Dashed lines radiate from the IP, representing the angular coverage of the detector. The diagram illustrates the geometry for calculating lifetime coverage at a distance of 50 m from the IP.

Distance from IP	ϵ_T
10 m	0.4%
100 m	5.8%
1000 m	1.0%

Being far away isn't even really helpful for probing longer lifetimes, since for very long lifetimes the exponential looks almost flat anyway. What really matters is your volume/lumi. Of course if you see a signal, you'll struggle to measure its lifetime without a deep detector or precise timing..

Side effects of that kind of size

Huge distance to first measured
point inside tracker!

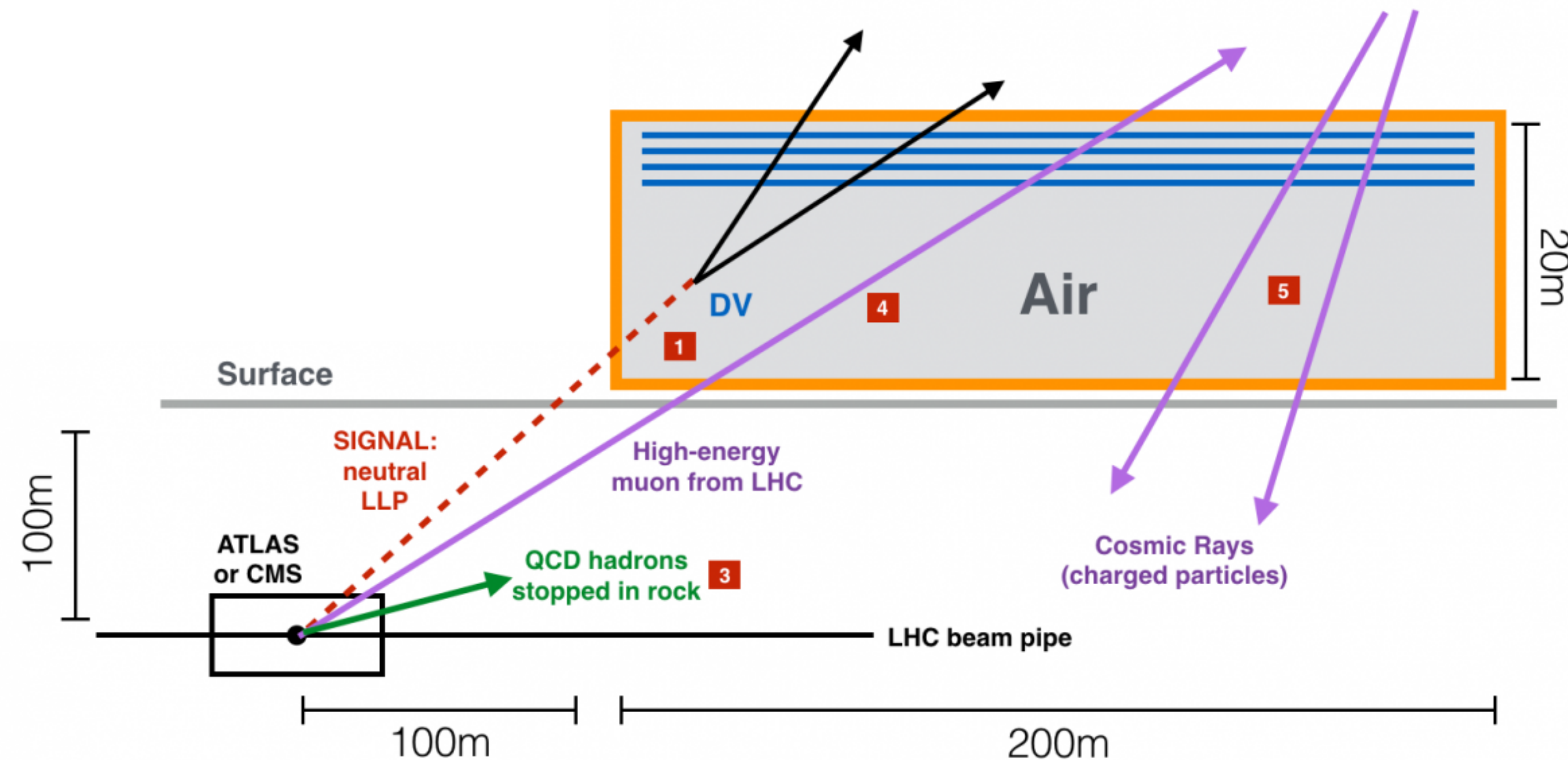


This also has an interesting impact on your vertex resolution, which is shared by the fixed-target layout. Prepare to have distances of closest approach $0(\text{cm})$ for your signal products...

A kingdom for a magnet

Collider mode : good luck...

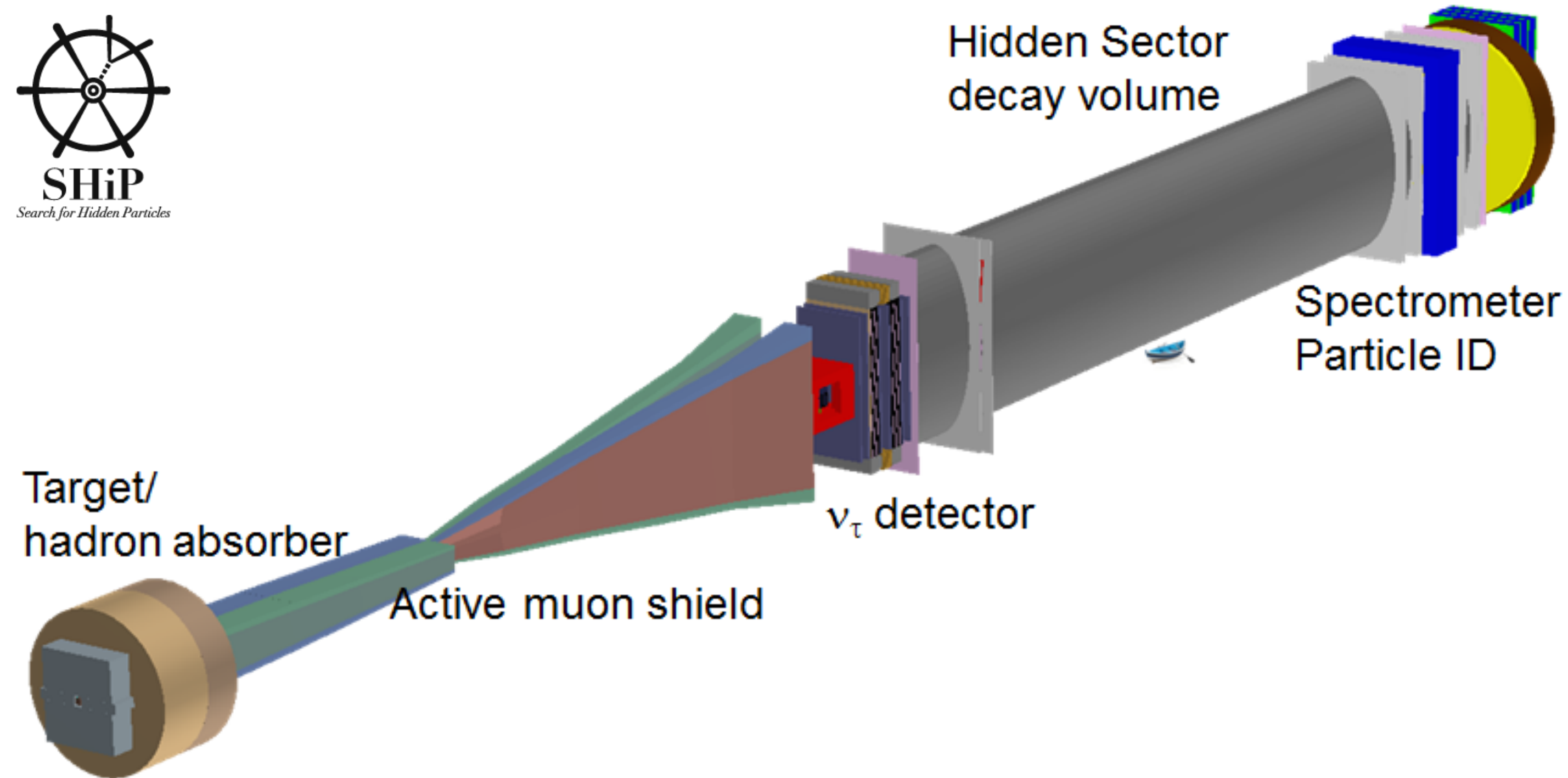
~~MATHUSLA~~



The other problem with uncollimated production is that unless you want to do something crazy with permanent magnets, you are not really going to be able to install one to cover the volume

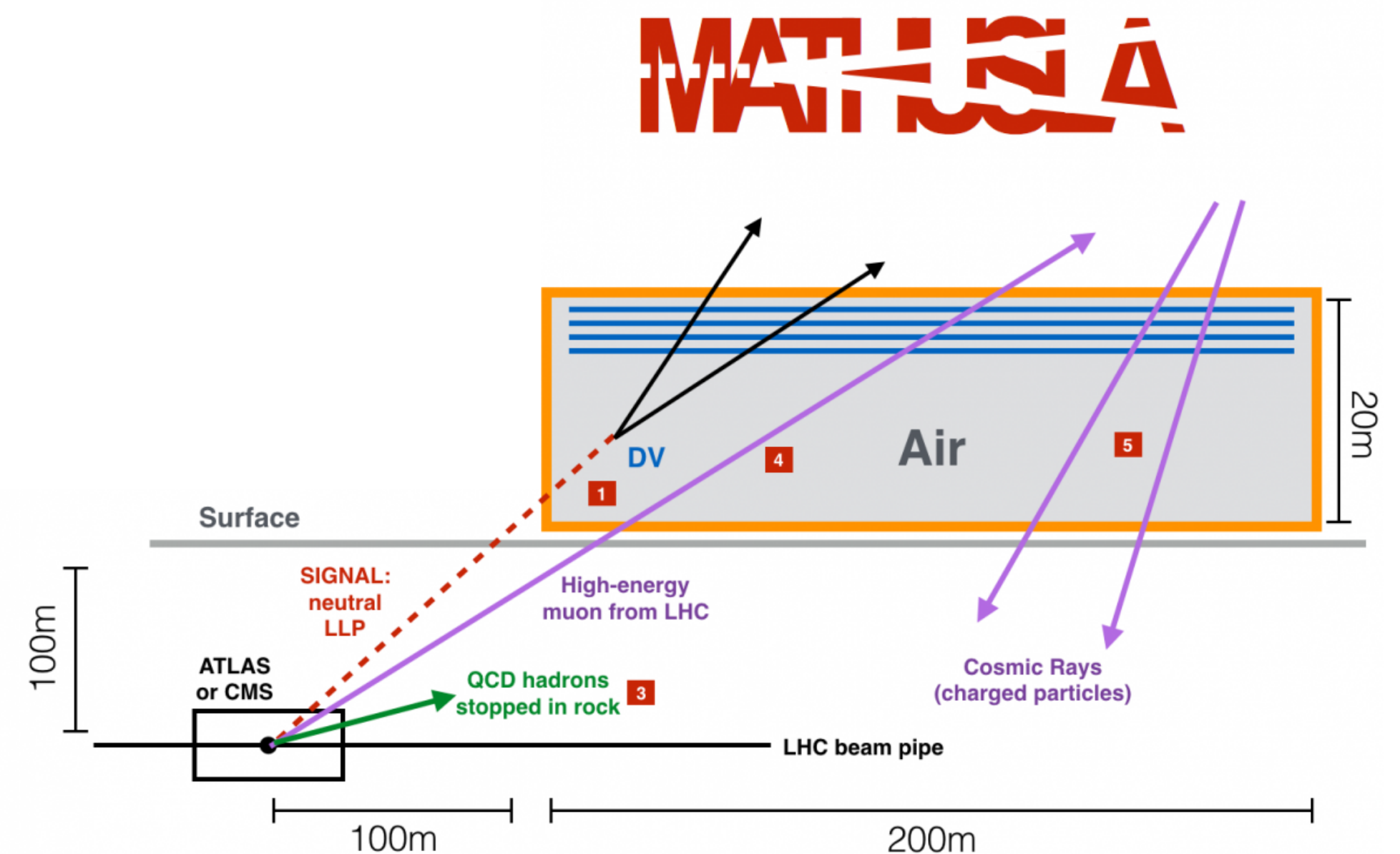
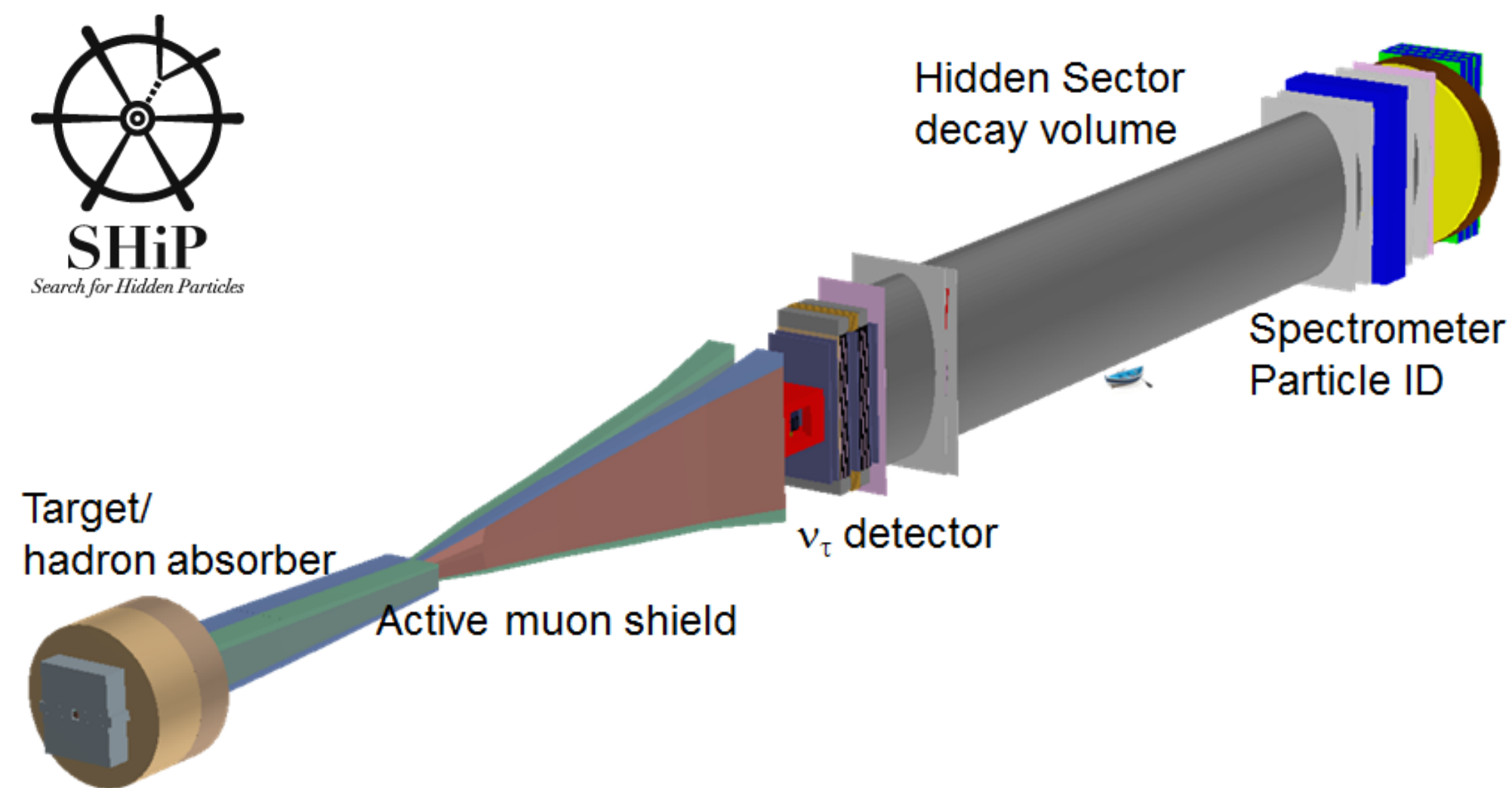
A kingdom for a magnet

Fixed target : easy!



In fixed target mode on the other hand, even if your distance to the first measured point is large, all decay products go in a small geometrical cone, so quite possible to add a magnet

The quest for zero background



Considerations : size of shield, active layer for in-shield secondary production, vacuum decay vessel or calorimeter style detector (?), magnet or timing/calorimetry for reconstruction?

Fixed target case
study : SHIP

Detector design

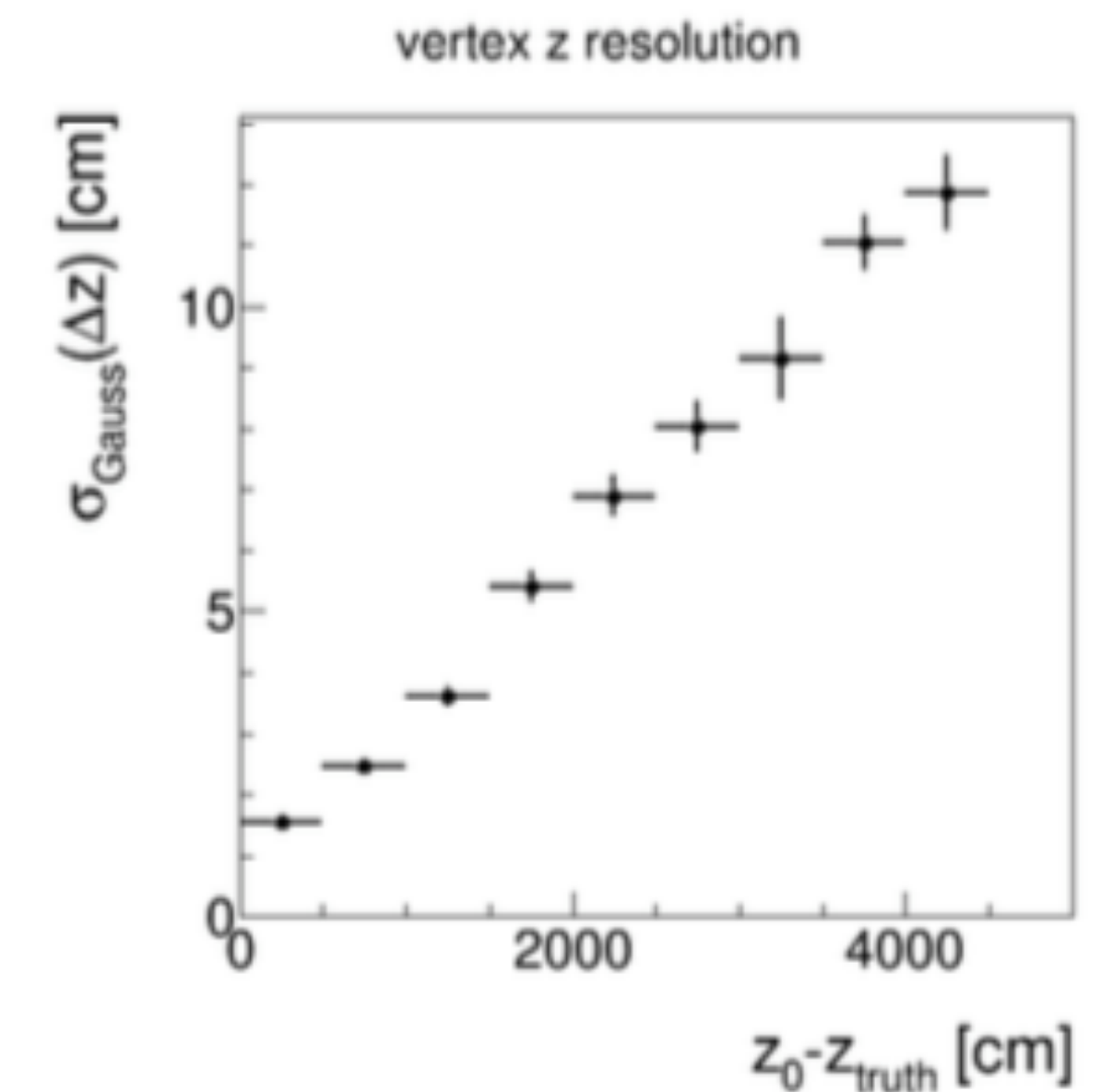
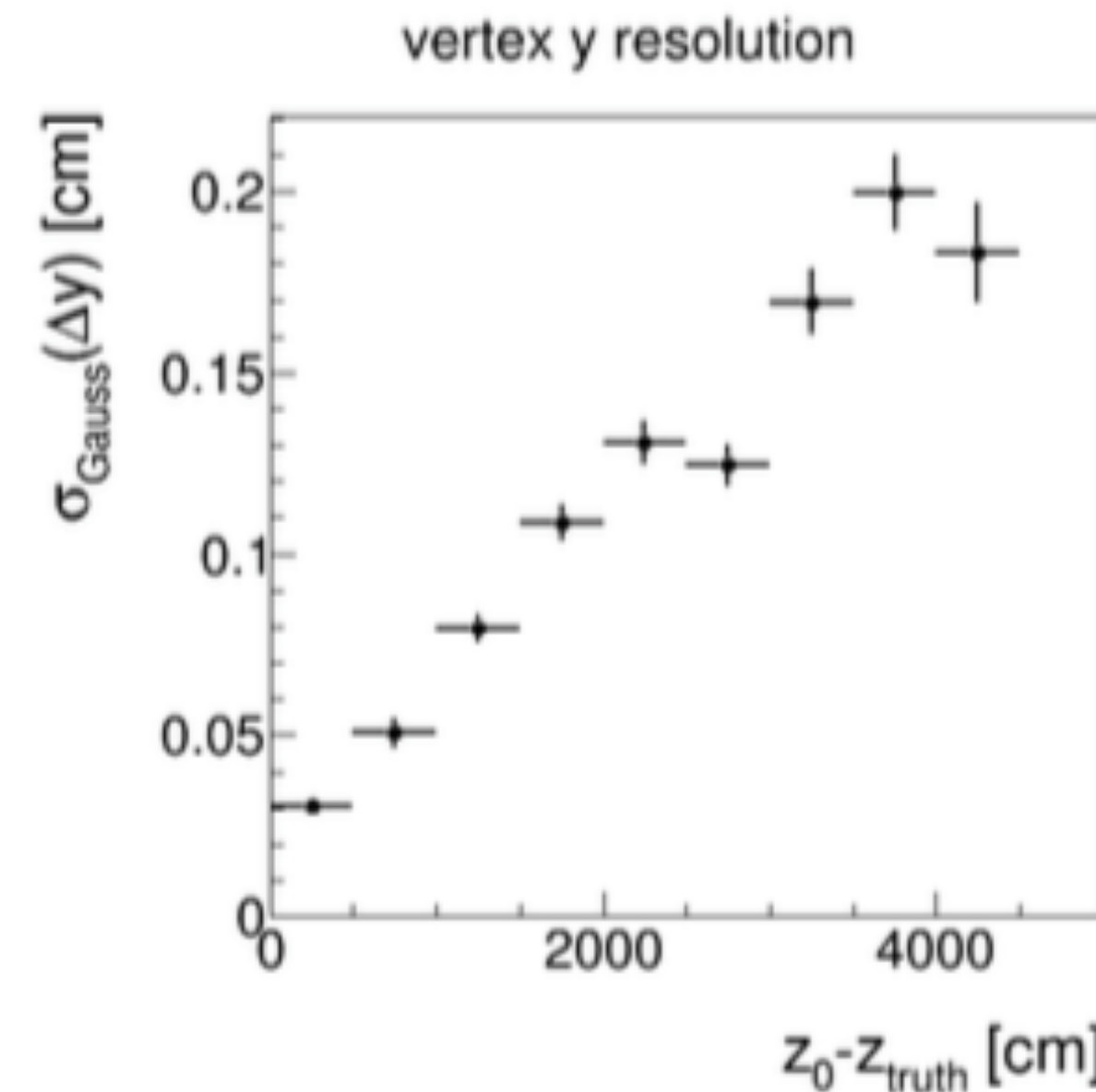
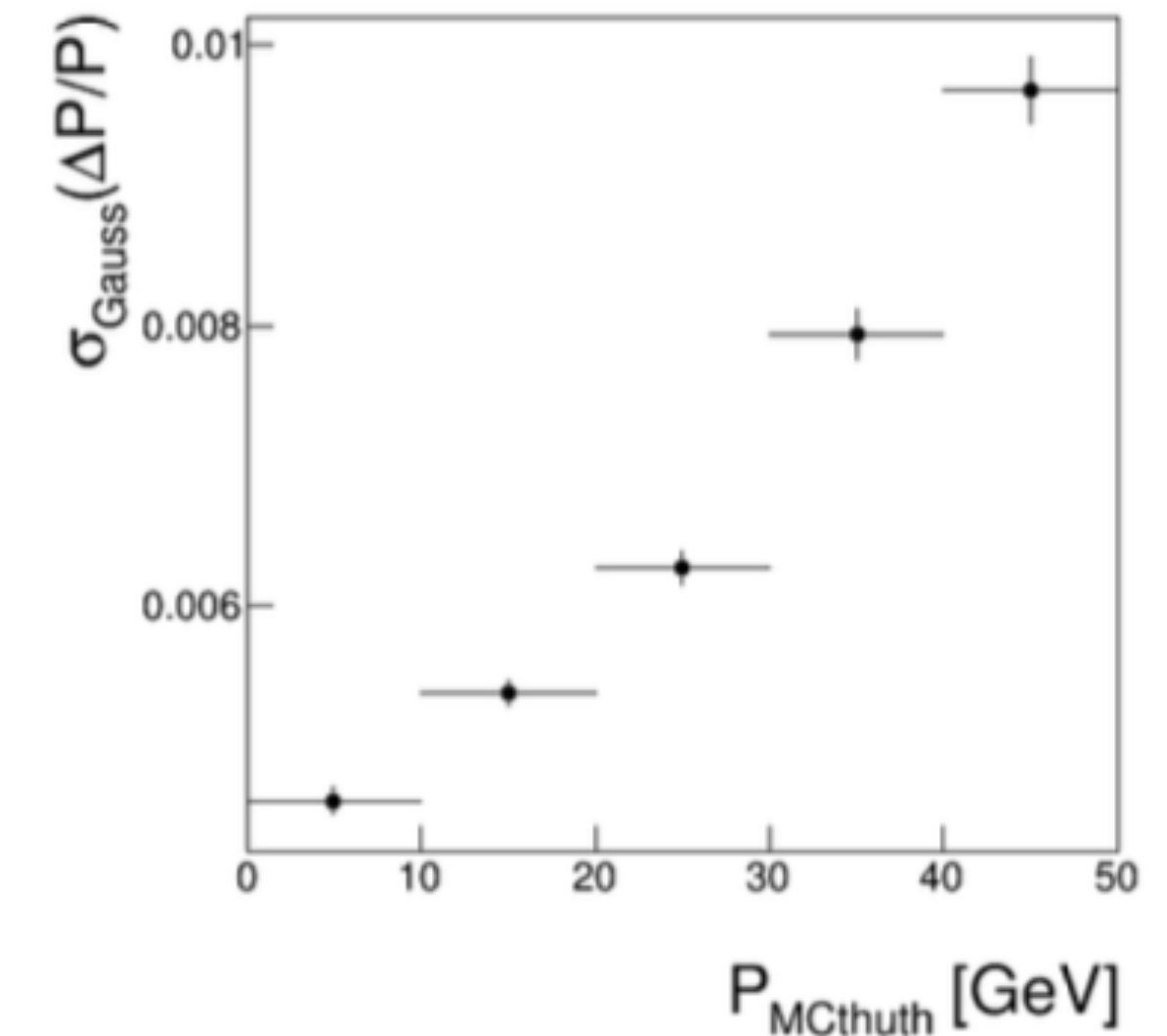
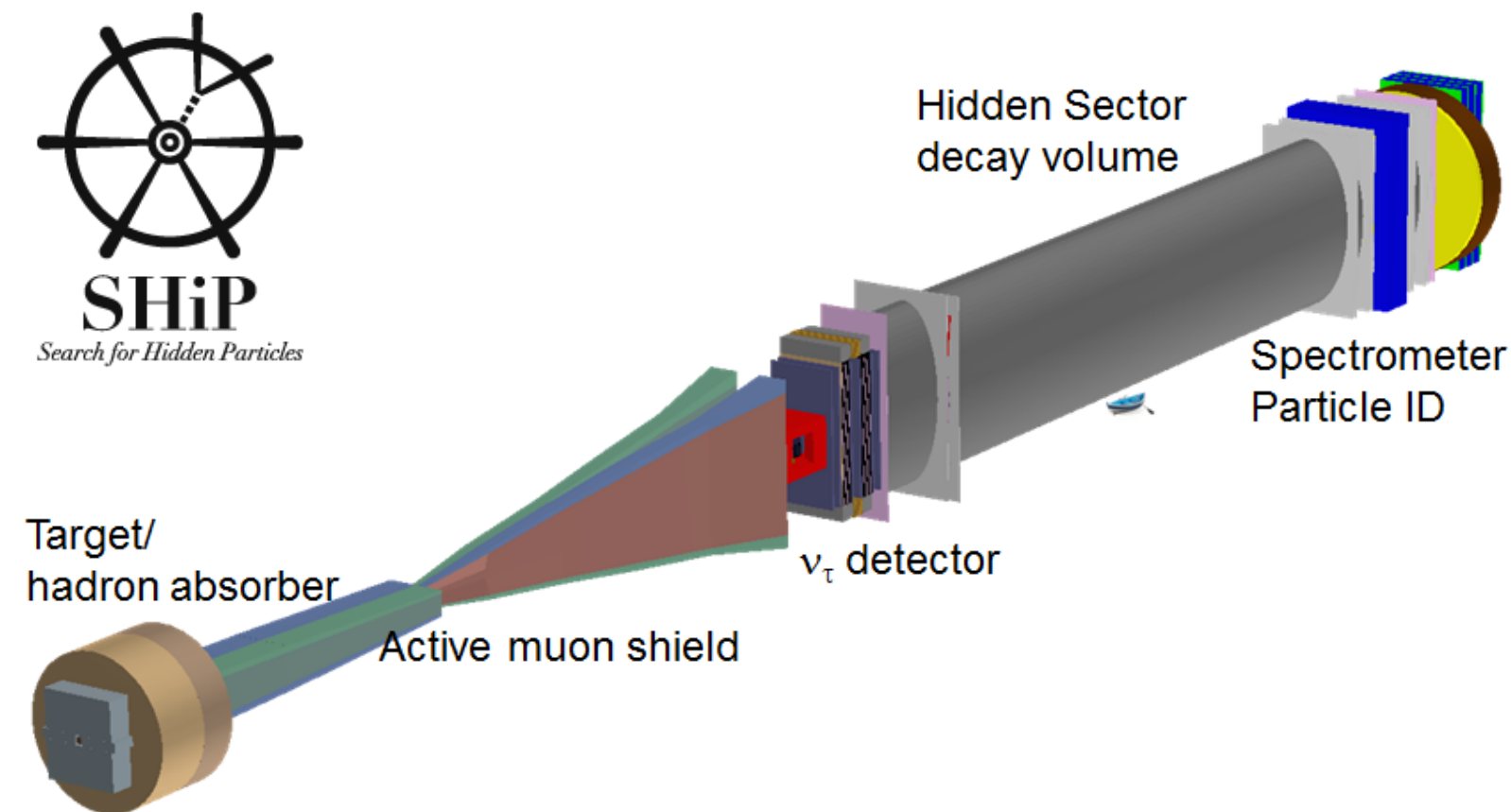
Key points :

Active shield and vacuum decay volume to minimize backgrounds

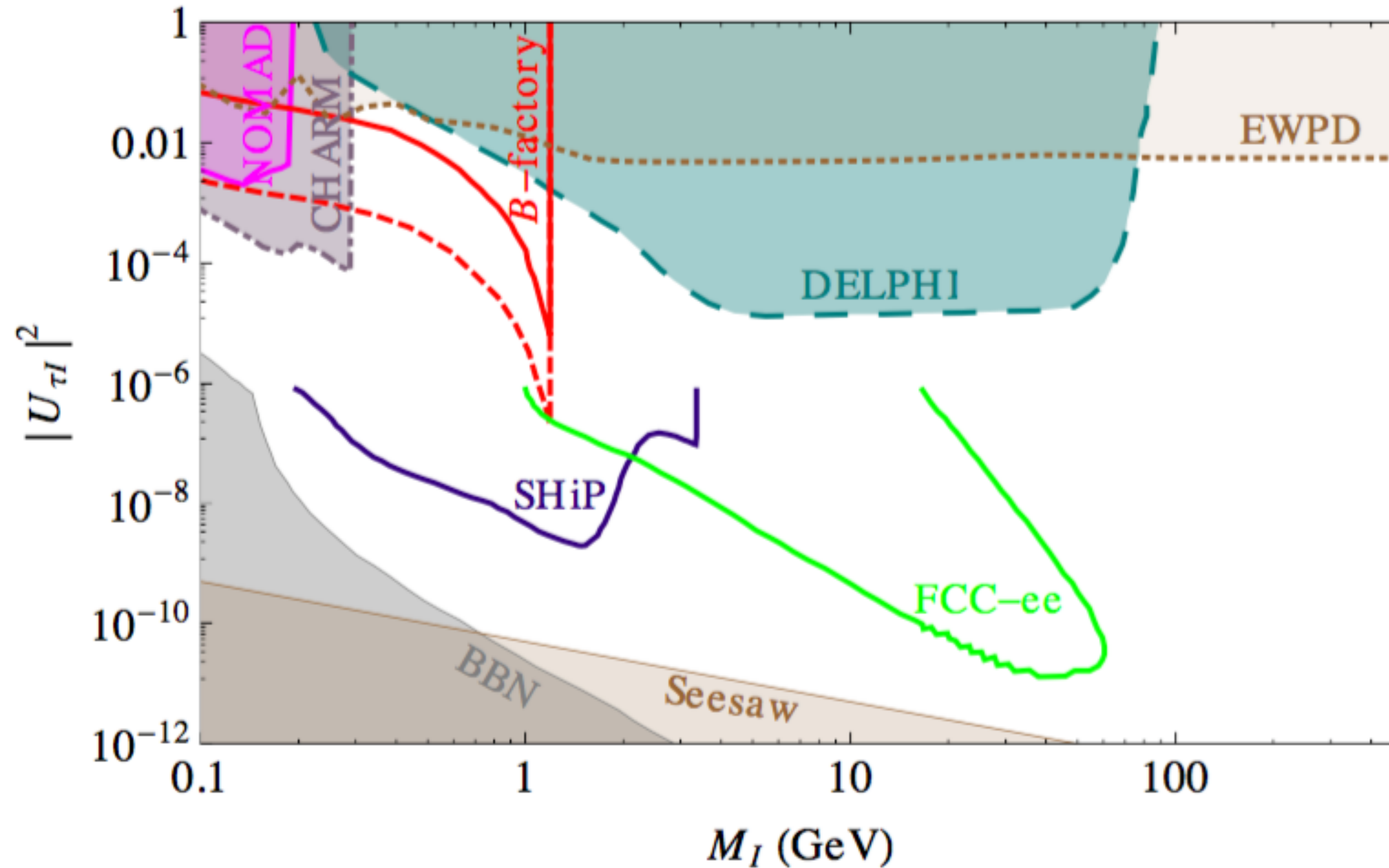
Sub percent momentum resolution, particle ID, mm vertex resolution in the transverse plane

Timing coincidence (a la NA62) used to suppress backgrounds

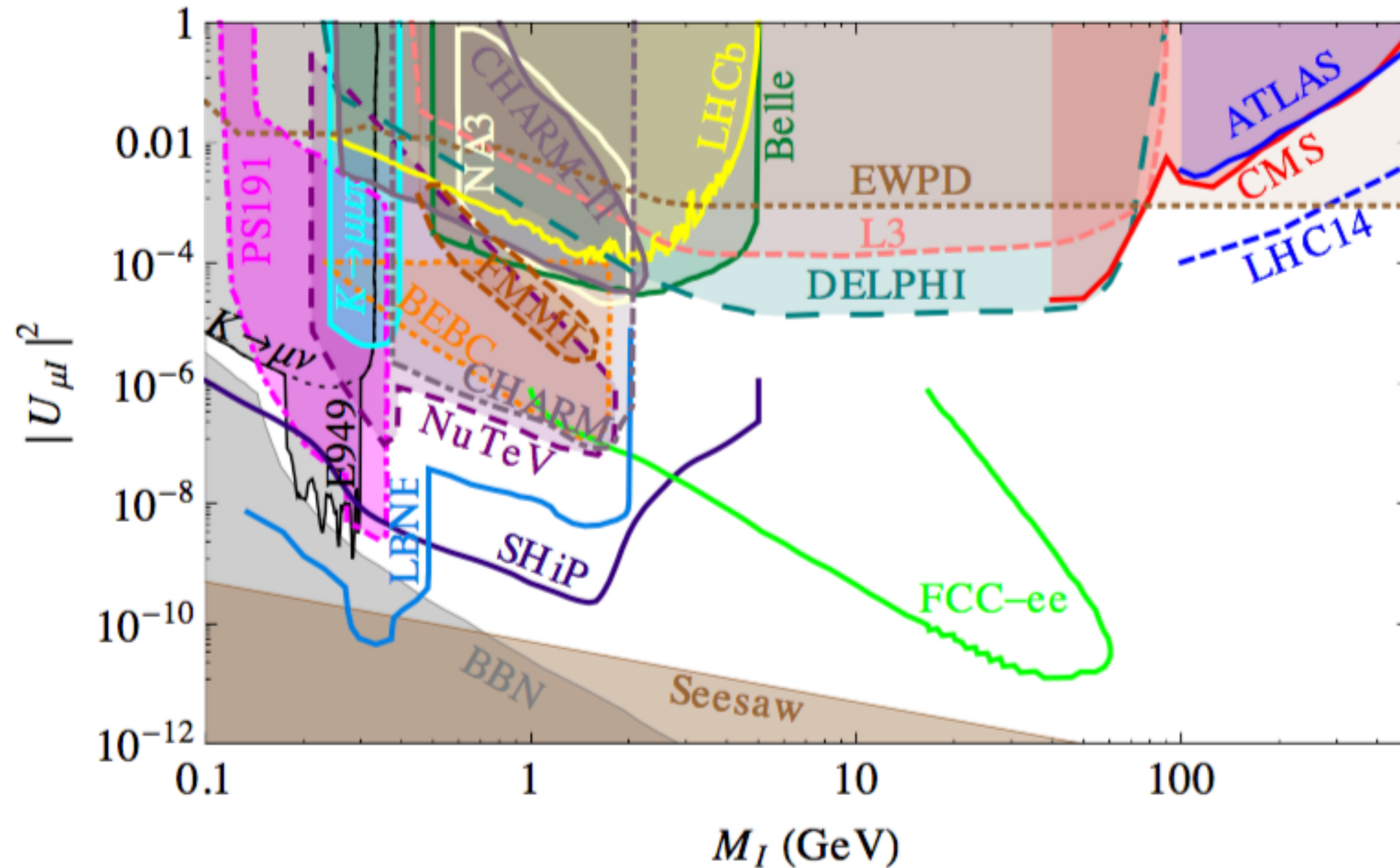
Exploits boost of produced heavy flavour to improve acceptance for LLPs, particularly shorter lived ones



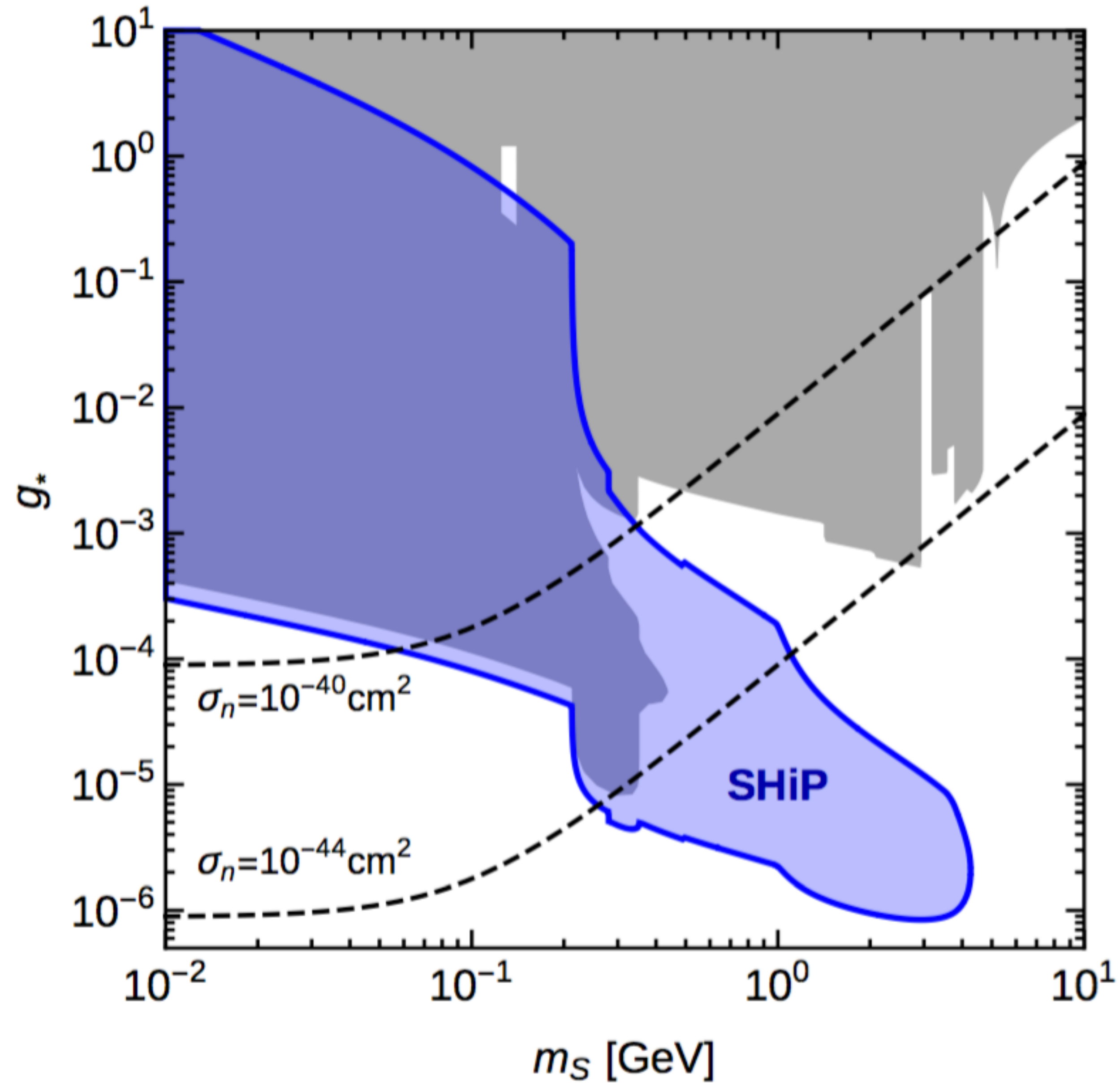
Reach estimates for HNLs



Reach estimates for HNLs



Reach estimates for $b \rightarrow sX$



Collider case
study : MATHUSLA

Detector design

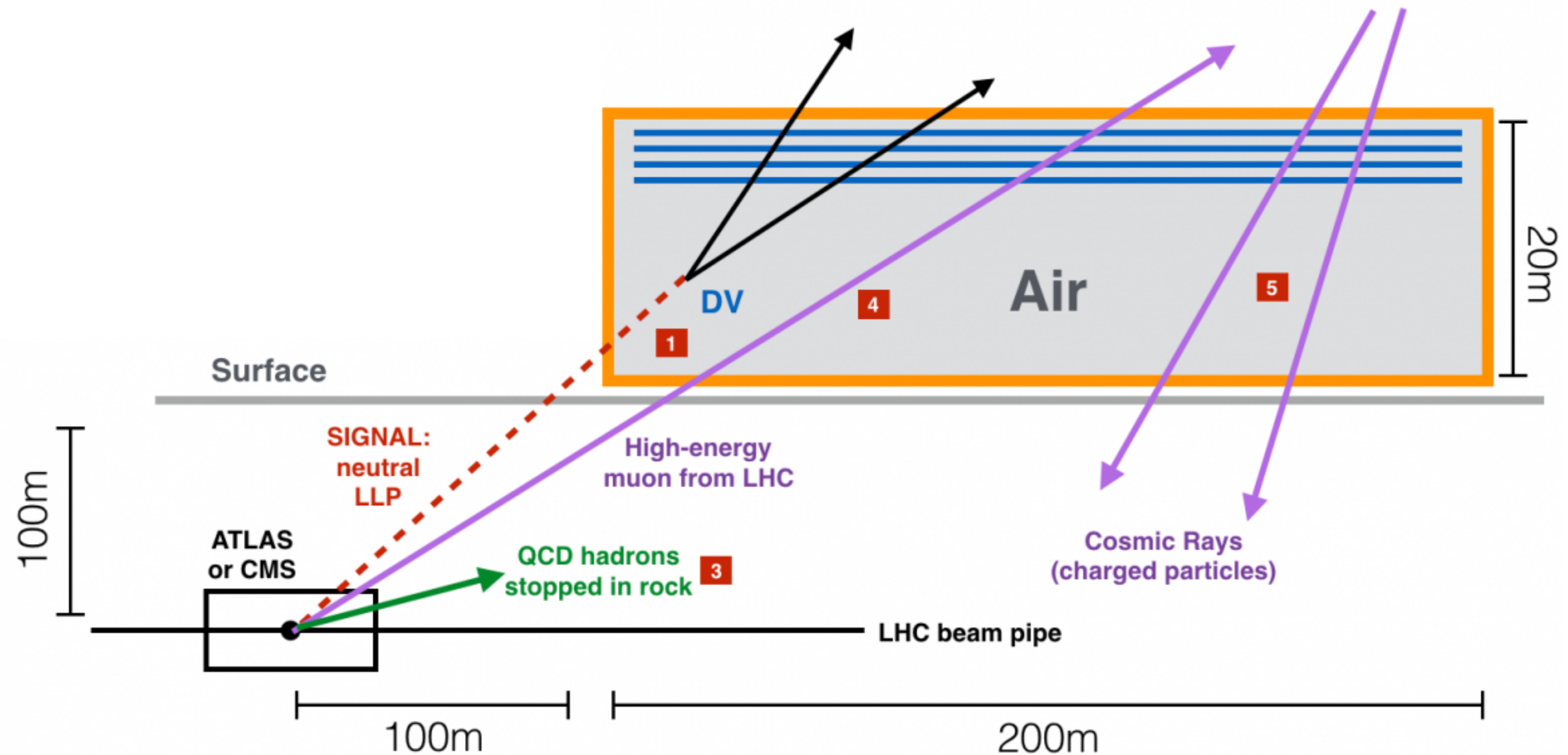
MATIS A

Key points :

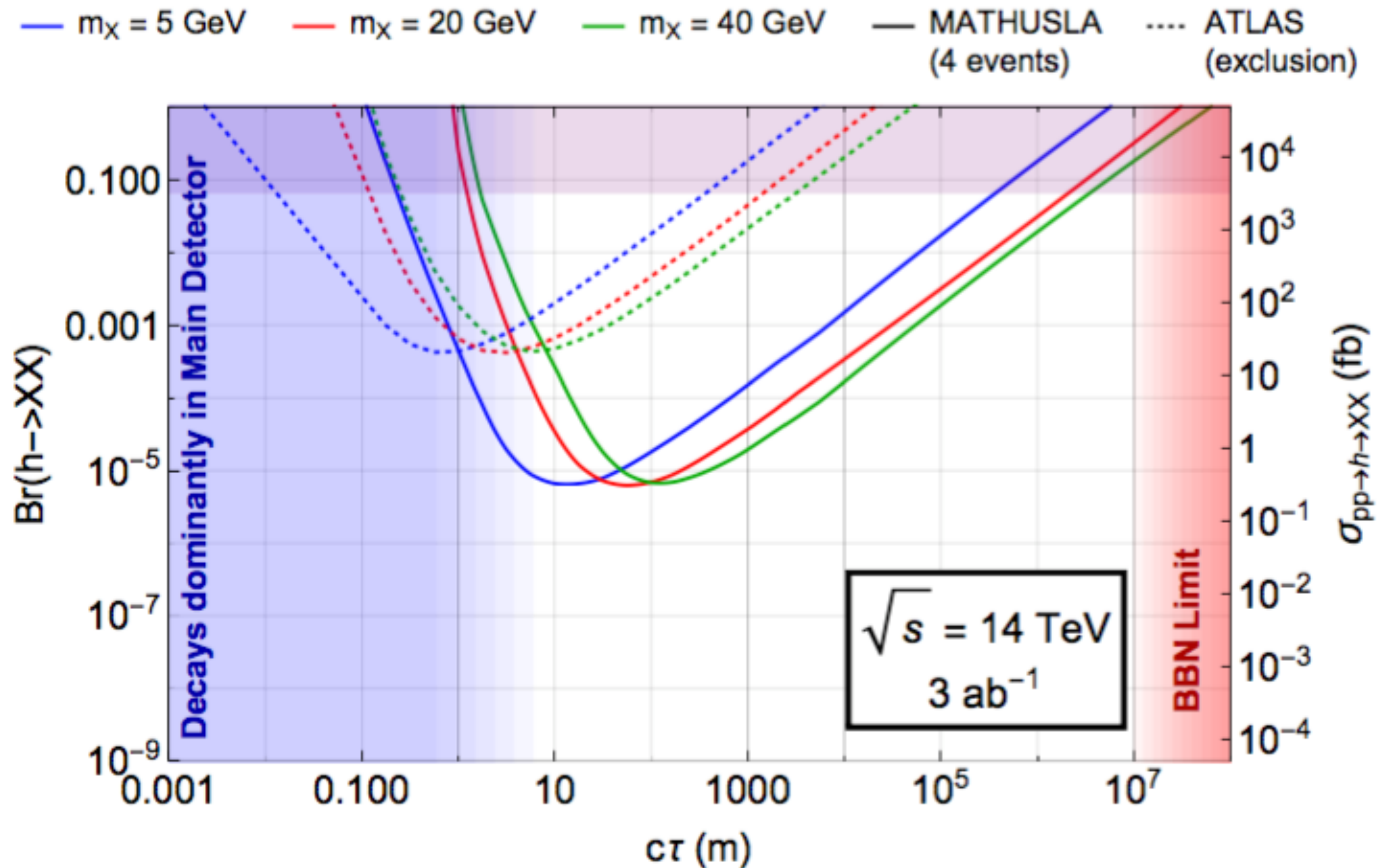
Access full HL-LHC luminosity

“Natural” shielding from LHC backgrounds, active vetoes on sides for cosmes and similar

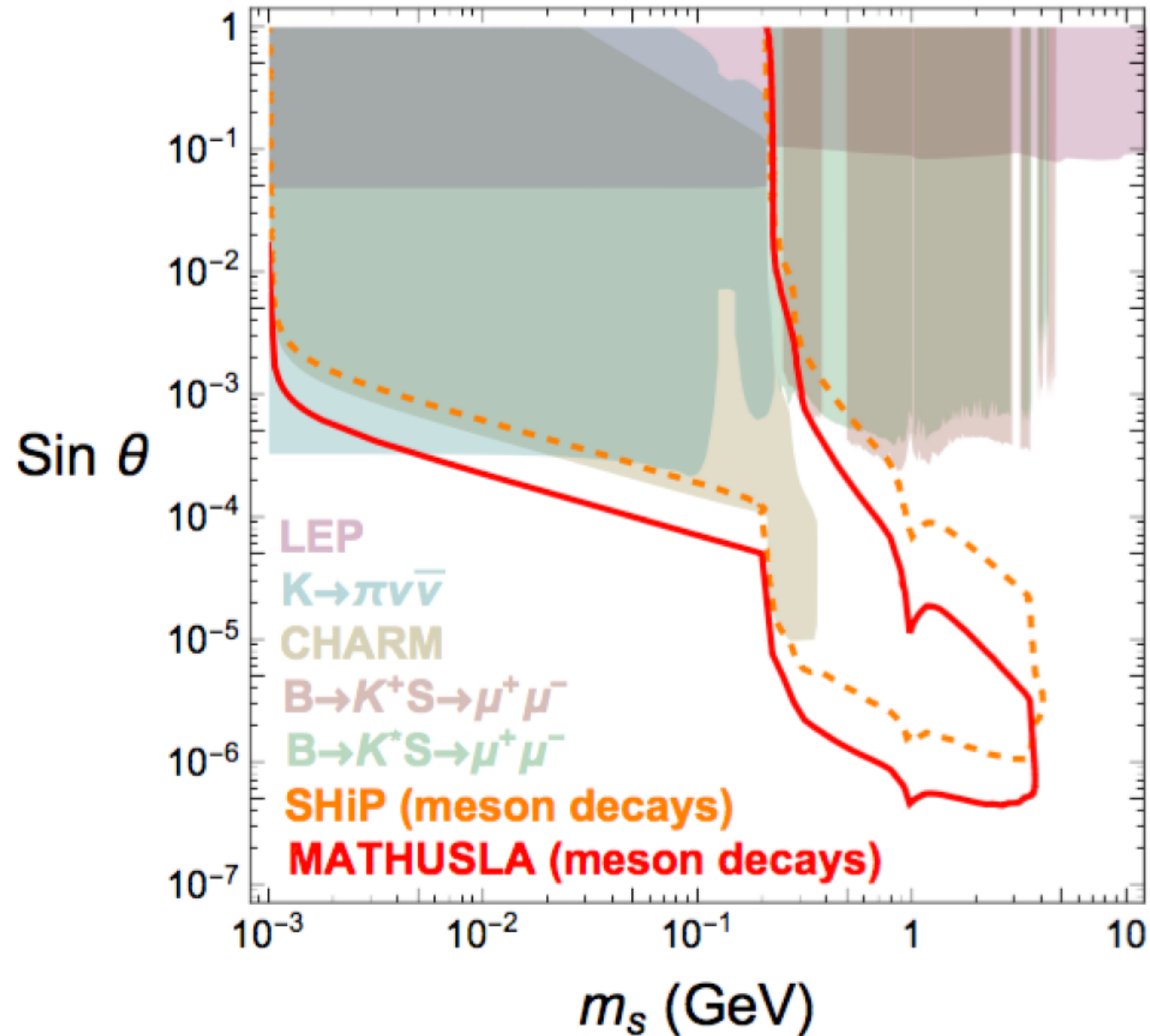
Enormous size : several tracking layers of $200 \times 200 \text{ m}^2$ each



Reach estimates for Higgs portal

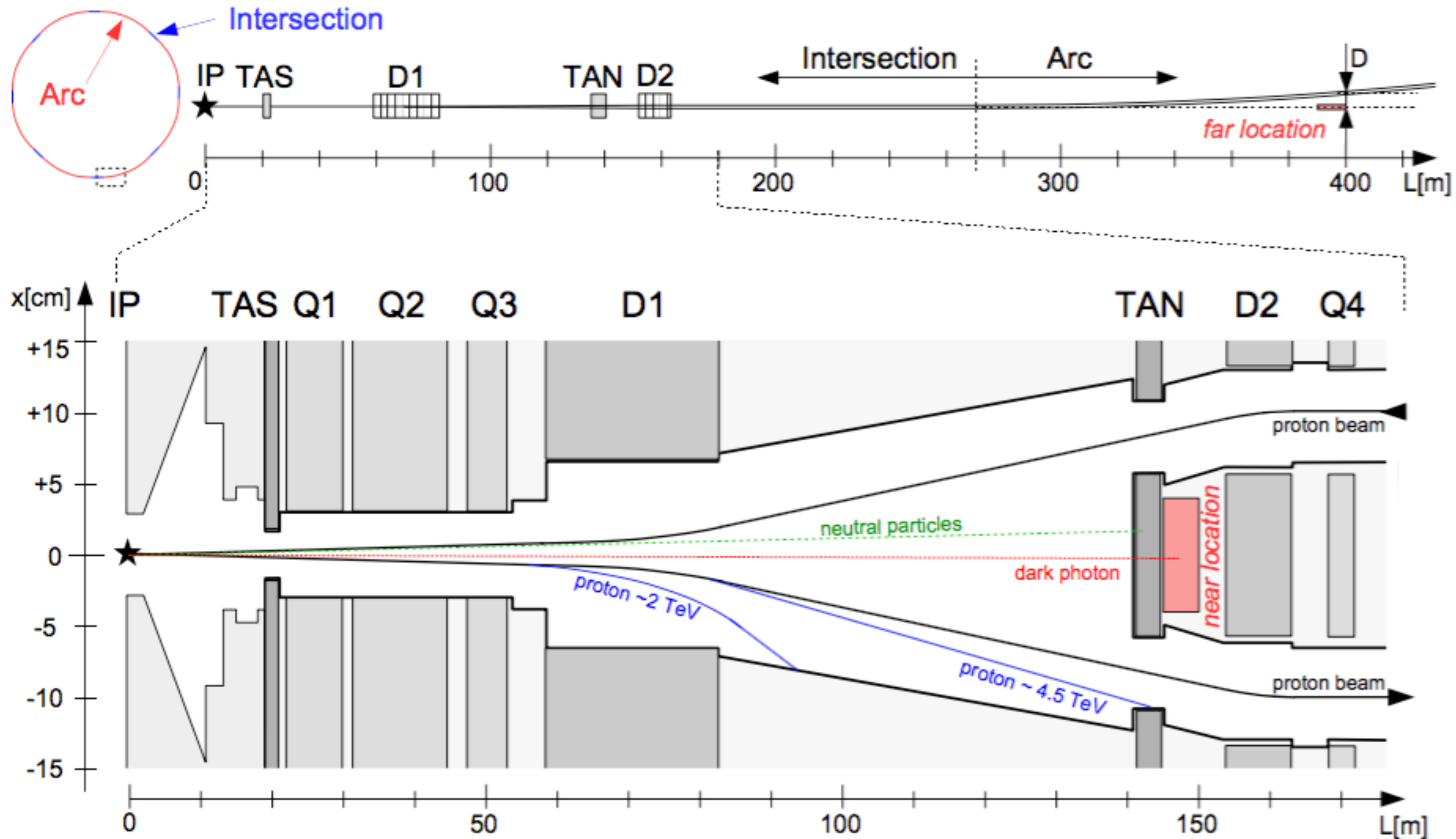


Reach estimates for $b \rightarrow sX$



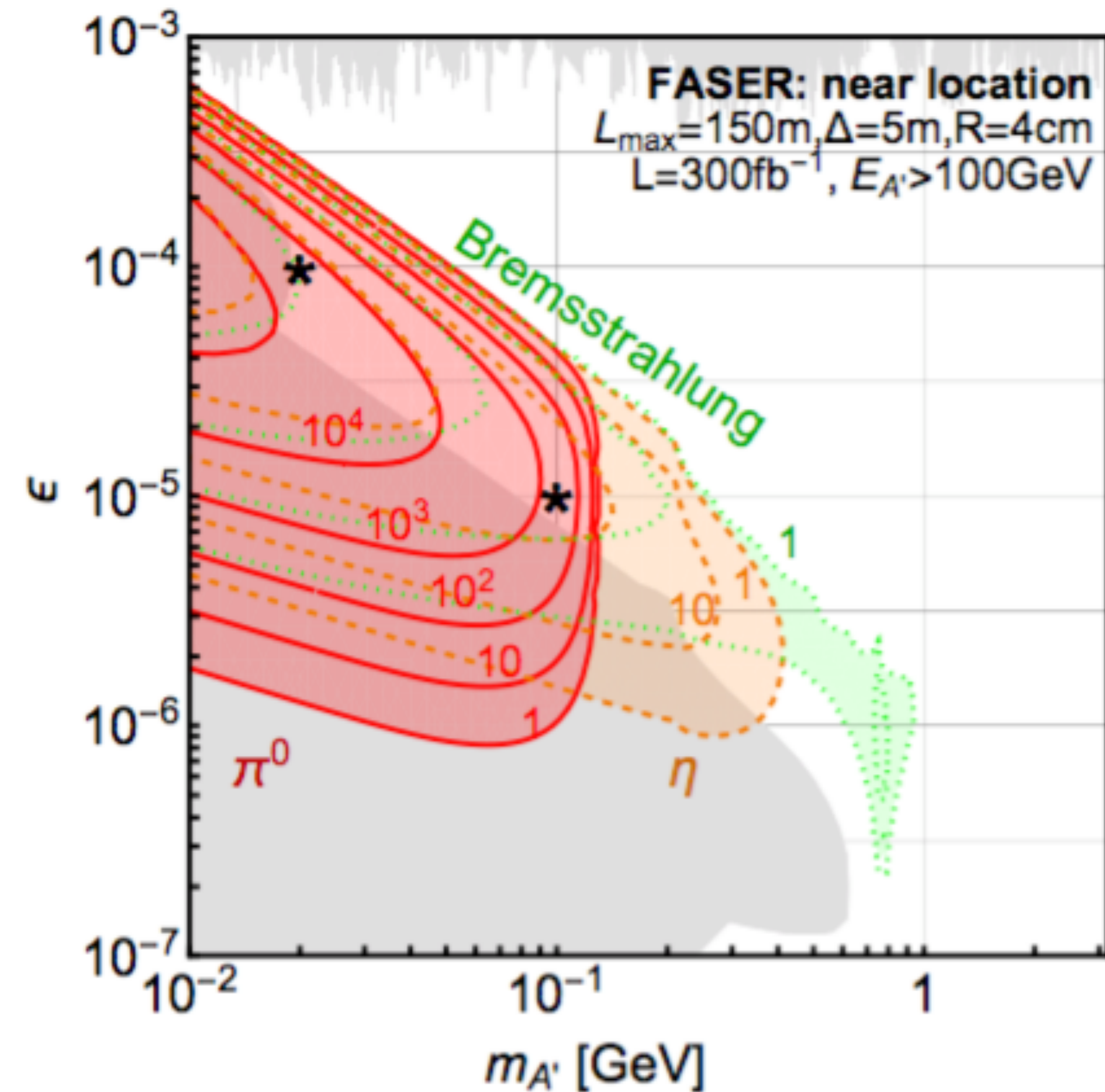
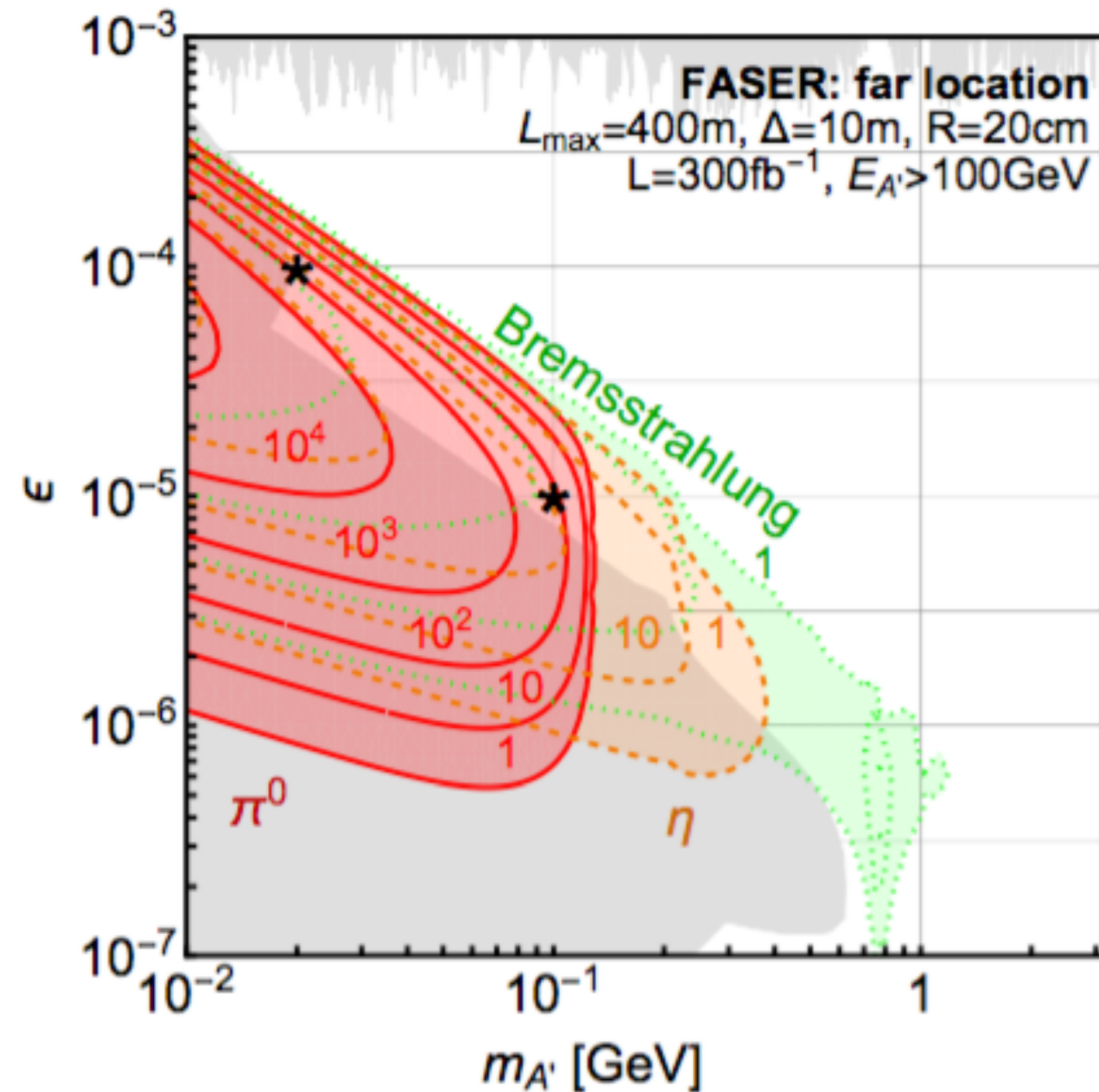
Collider case
study : FASER

Detector design



Very forward, exploits tail of the boost distribution

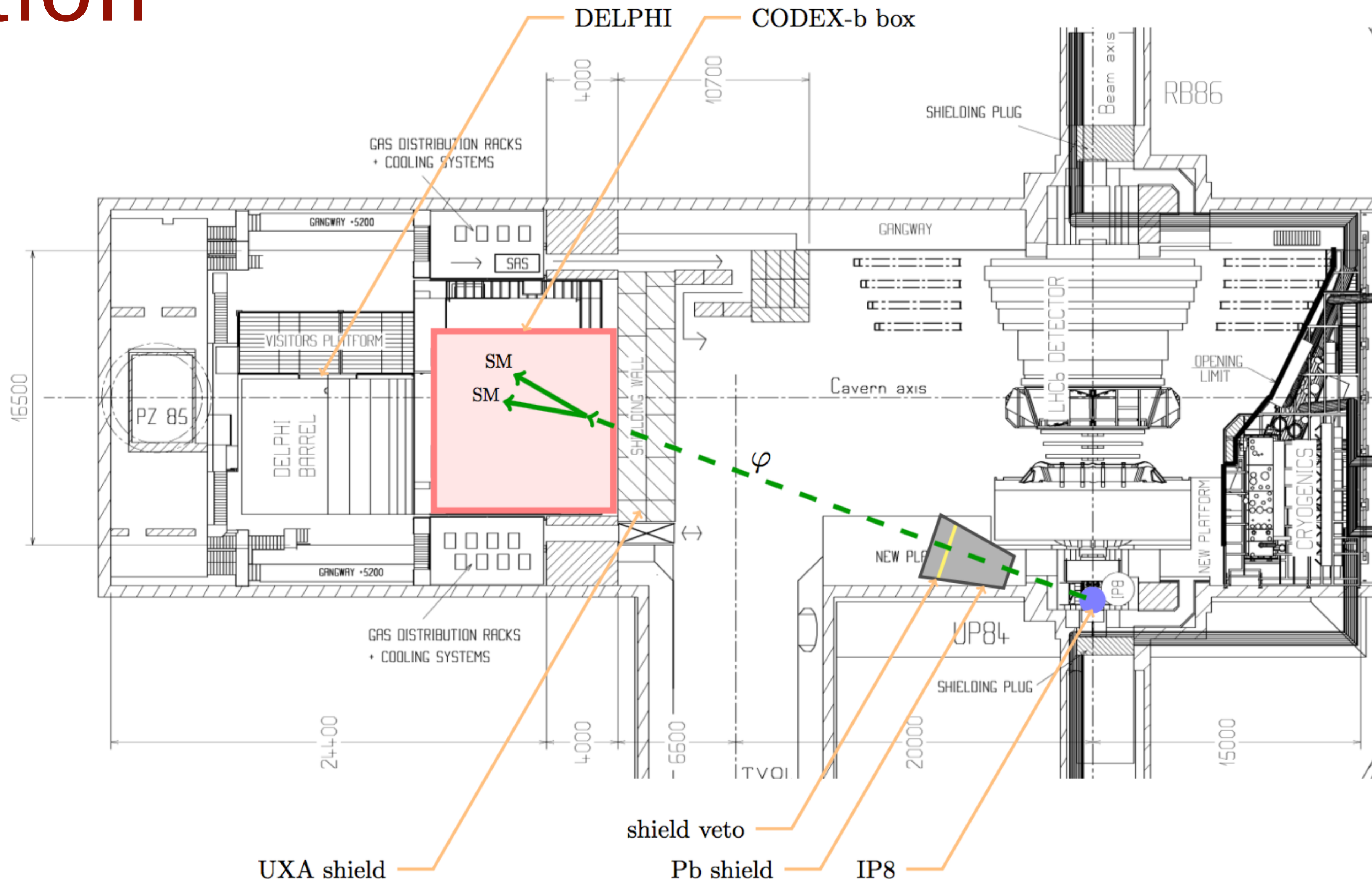
Reach estimates for dark photons



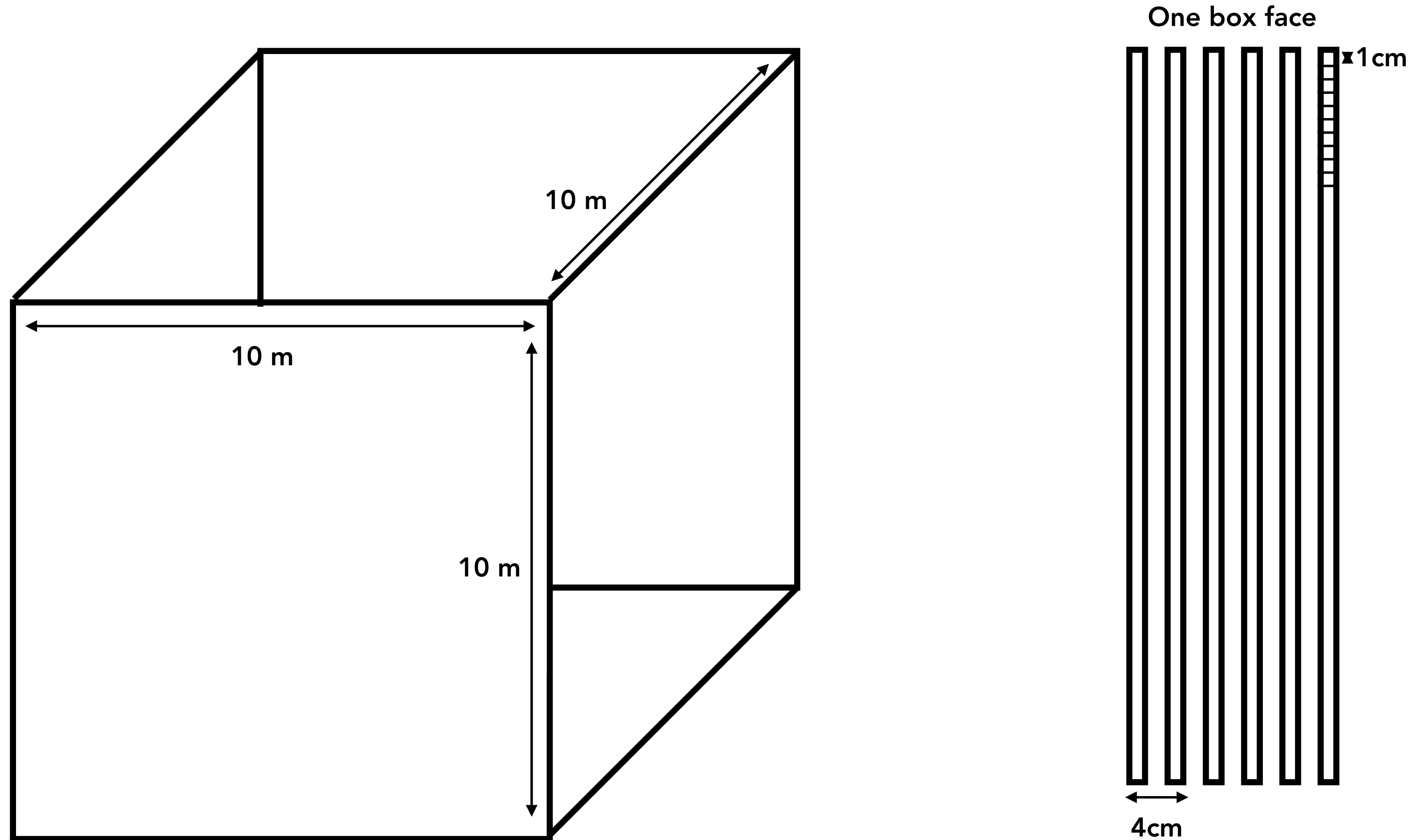
Production of proton brems (!) highlights unique forward regime

Collider case
study : CODEX-b

Location

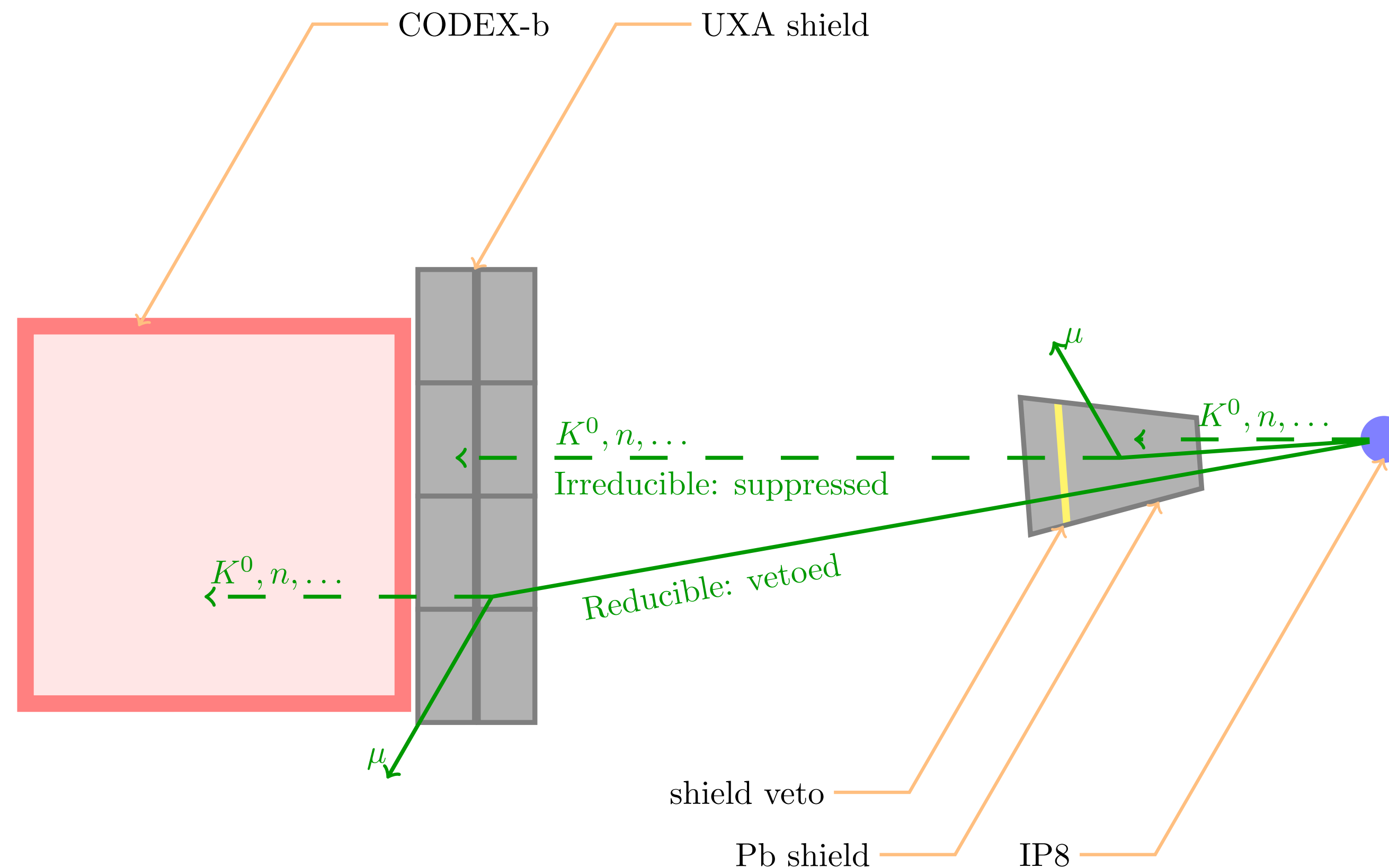


Minimal proof-of-concept geometry



10x10x10 metre box, with 6 RPC layers on each box face. Assume 1 cm granularity for the RPCs, and possibility of timing information (explored later in talk). Add 5 other triplets of RPC layers equally spaced in box to minimize the distance to the first measured point for the decay vertex determination.

Minimal shield & veto design



Simple design : use first part of the shield to attenuate muon & neutral hadron backgrounds which could enter the detector volume and scatter or decay within it, faking a signal. Then use a thin veto layer to eliminate secondary production of backgrounds within the shield itself.

Basic GEANT background estimate

BG species	Particle yields		Baseline Cuts
	irreducible by shield veto	reducible by shield veto	
$n + \bar{n}$	7	$5 \cdot 10^4$	$E_{\text{kin}} > 1 \text{ GeV}$
K_L^0	0.2	870	$E_{\text{kin}} > 0.5 \text{ GeV}$
$\pi^\pm + K^\pm$	0.5	$3 \cdot 10^4$	$E_{\text{kin}} > 0.5 \text{ GeV}$
$\nu + \bar{\nu}$	0.5	$2 \cdot 10^6$	$E > 0.5 \text{ GeV}$

Simulate initial background flux with Pythia 8, propagate through shield, air, and detector using GEANT4. A few things to note :

- Nominally largest background is neutrons entering the box
- Muon-air interactions can be vetoed using front detector faces
- Neutrino backgrounds are entirely negligible.

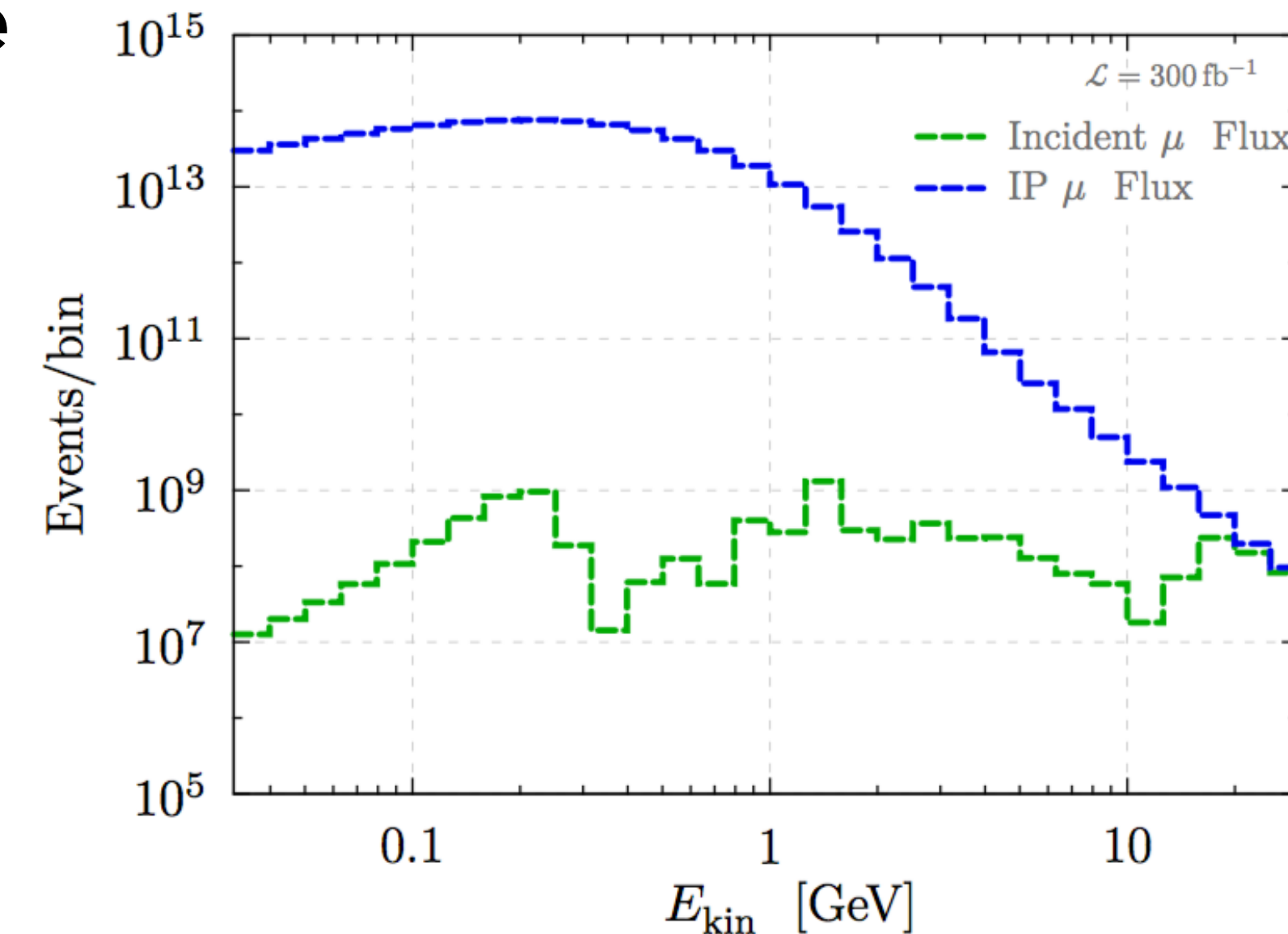
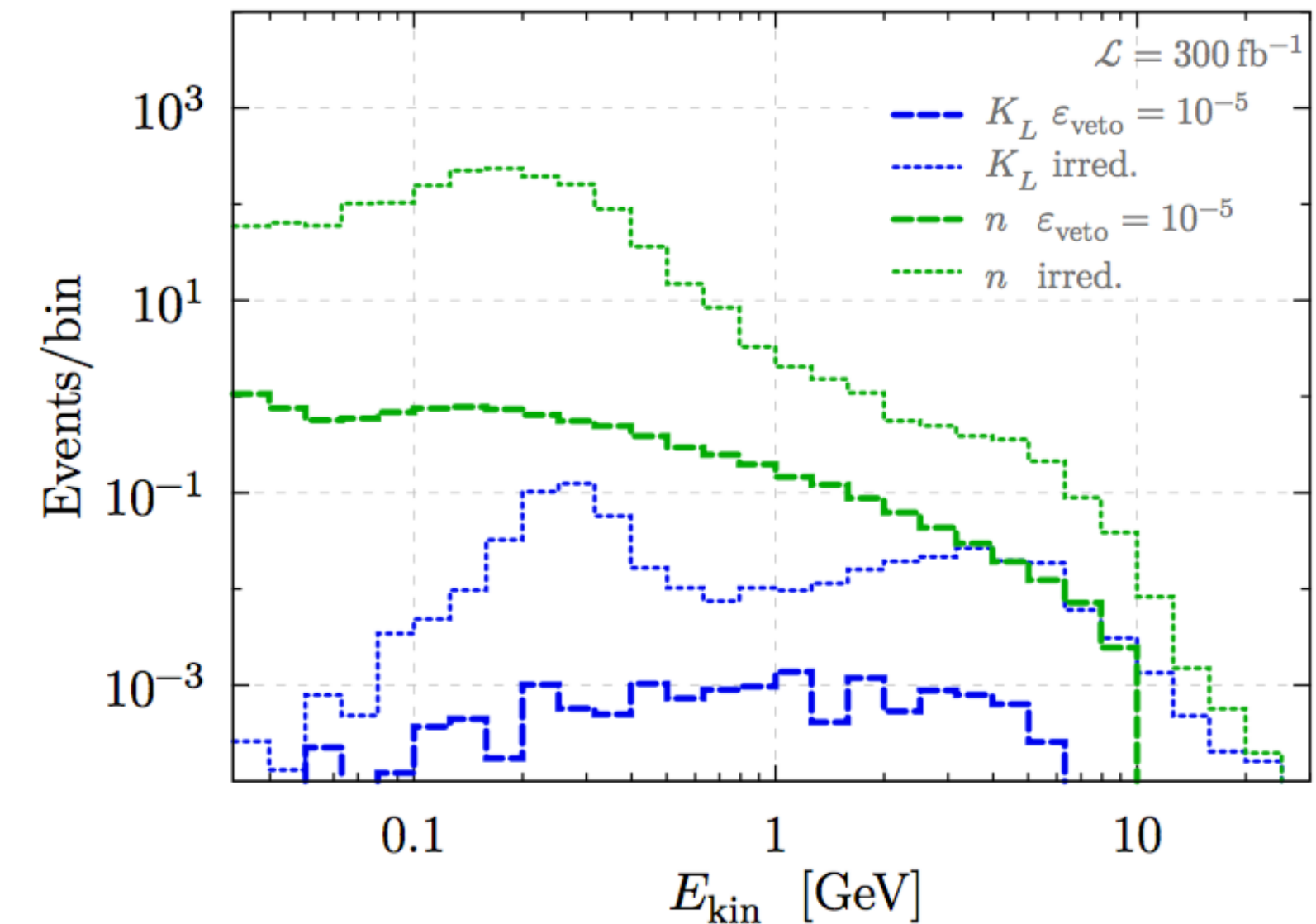
No attempt yet to use any properties of reconstructed backgrounds to reject them, but timing + spatial information should help there.

Energy spectrum of backgrounds

BG species	Particle yields		Baseline Cuts
	irreducible by shield veto	reducible by shield veto	
$n + \bar{n}$	7	$5 \cdot 10^4$	$E_{\text{kin}} > 1 \text{ GeV}$
K_L^0	0.2	870	$E_{\text{kin}} > 0.5 \text{ GeV}$
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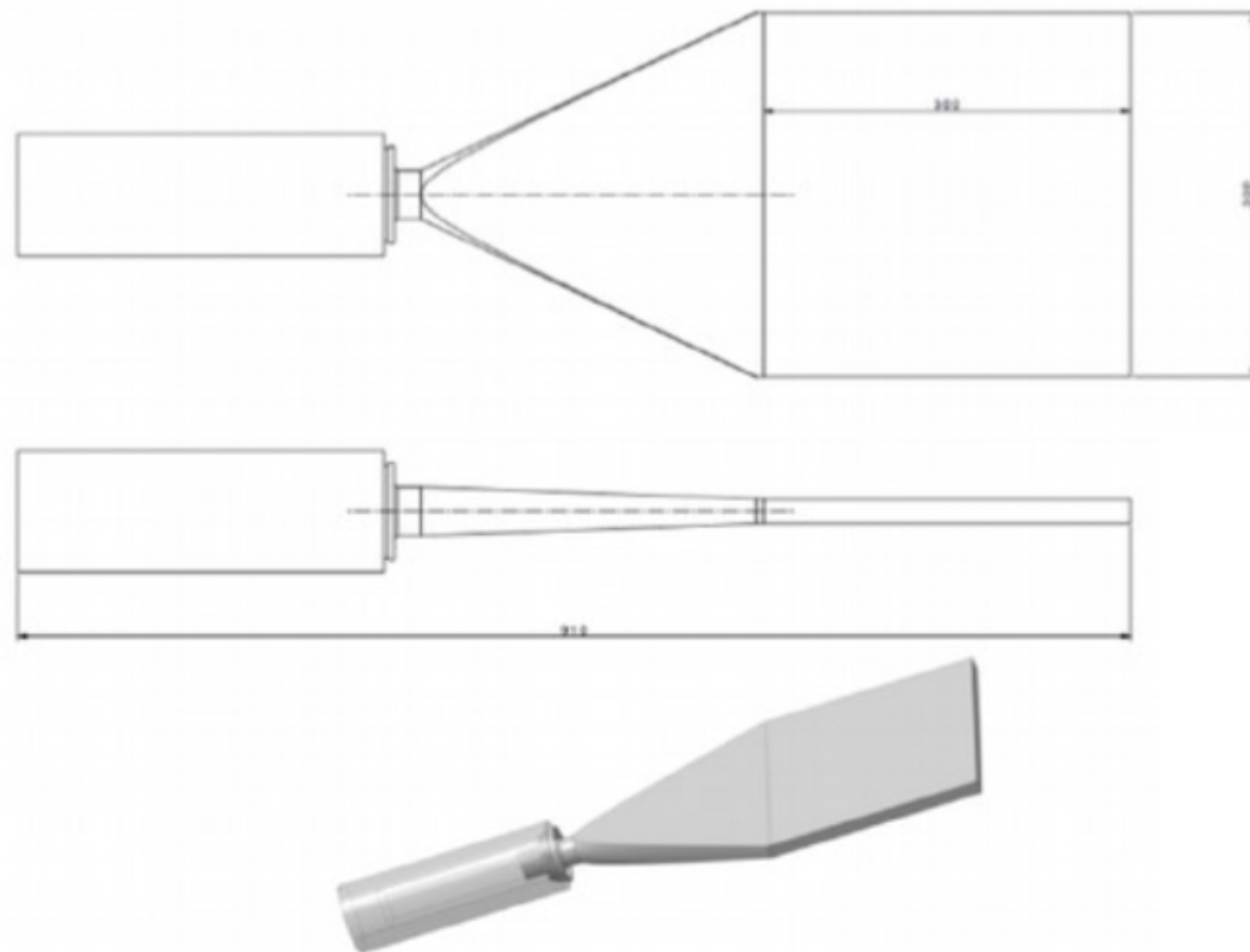
These are the numbers of unvetoable particles entering the box, the estimated number of scatters in box is <1 for all particle species!

Also notice the energy spectrum of these particles : most of them, especially the neutrons, are very soft!



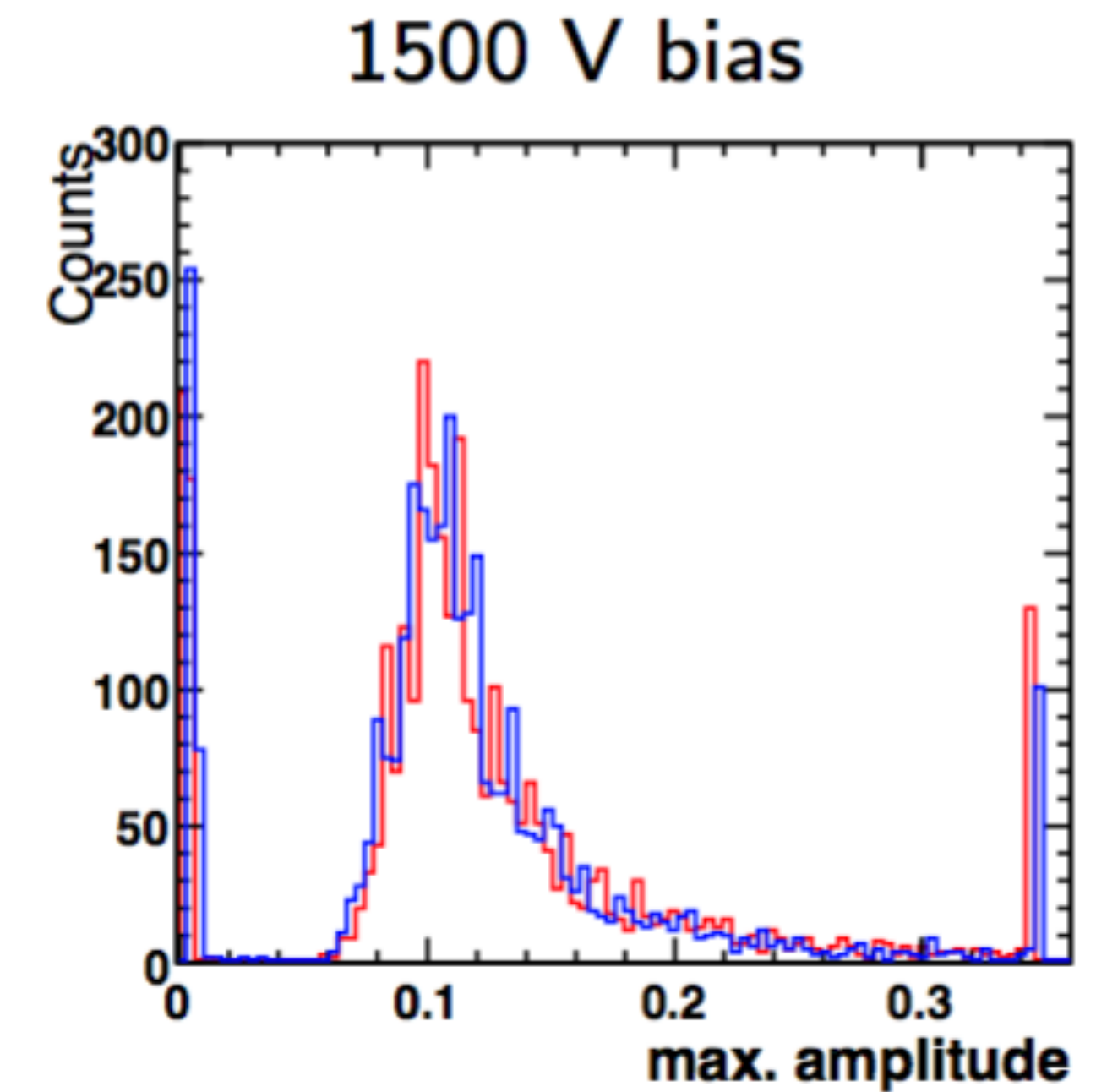
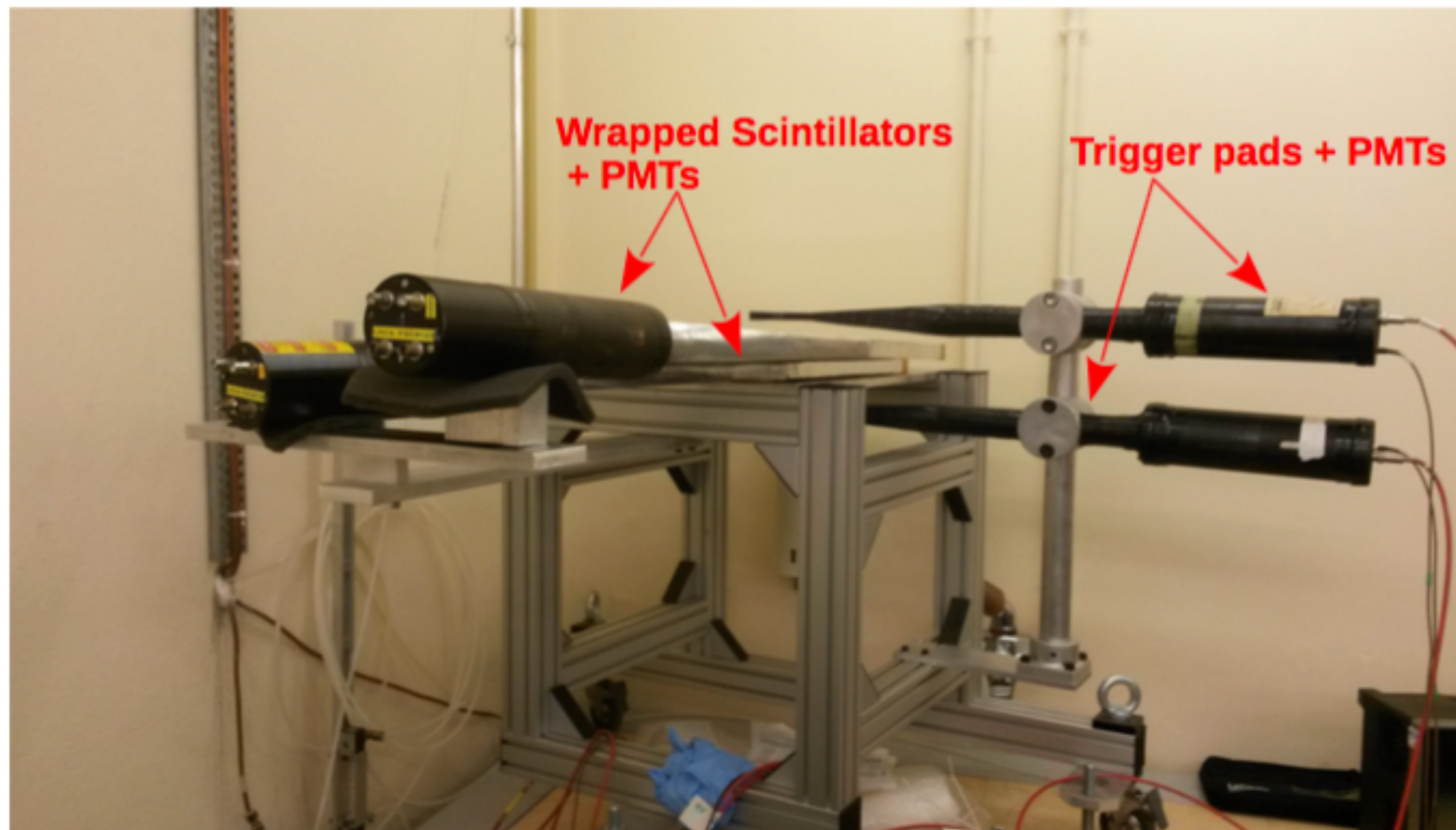
Backgrounds from data

- Two $30 \times 30 \times 2$ cm wrapped plastic scintillators + PMT + mechanical stand.



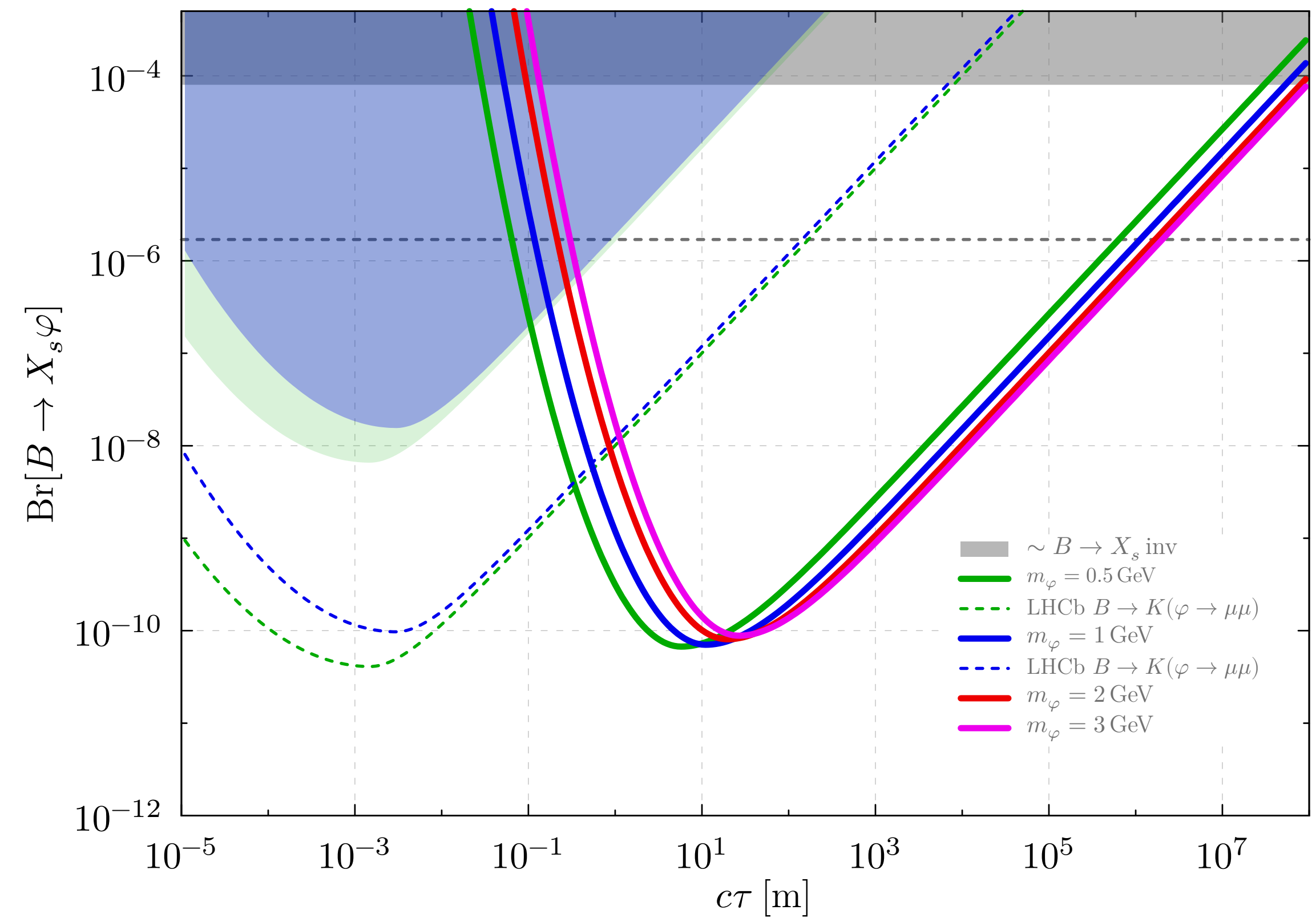
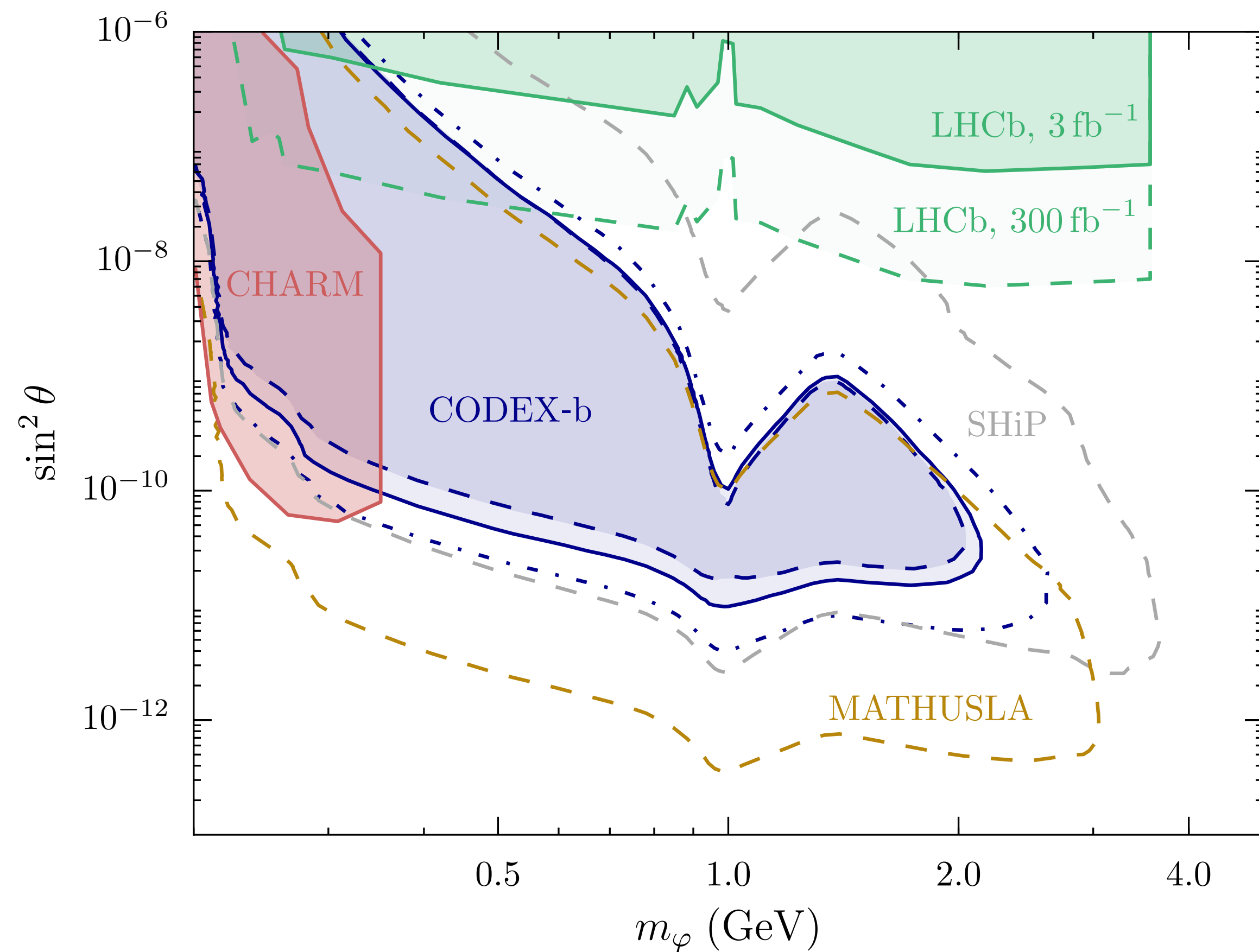
Backgrounds from data

- Setup tested with cosmics. $\mathcal{O}(3000)$ in a couple of hours.

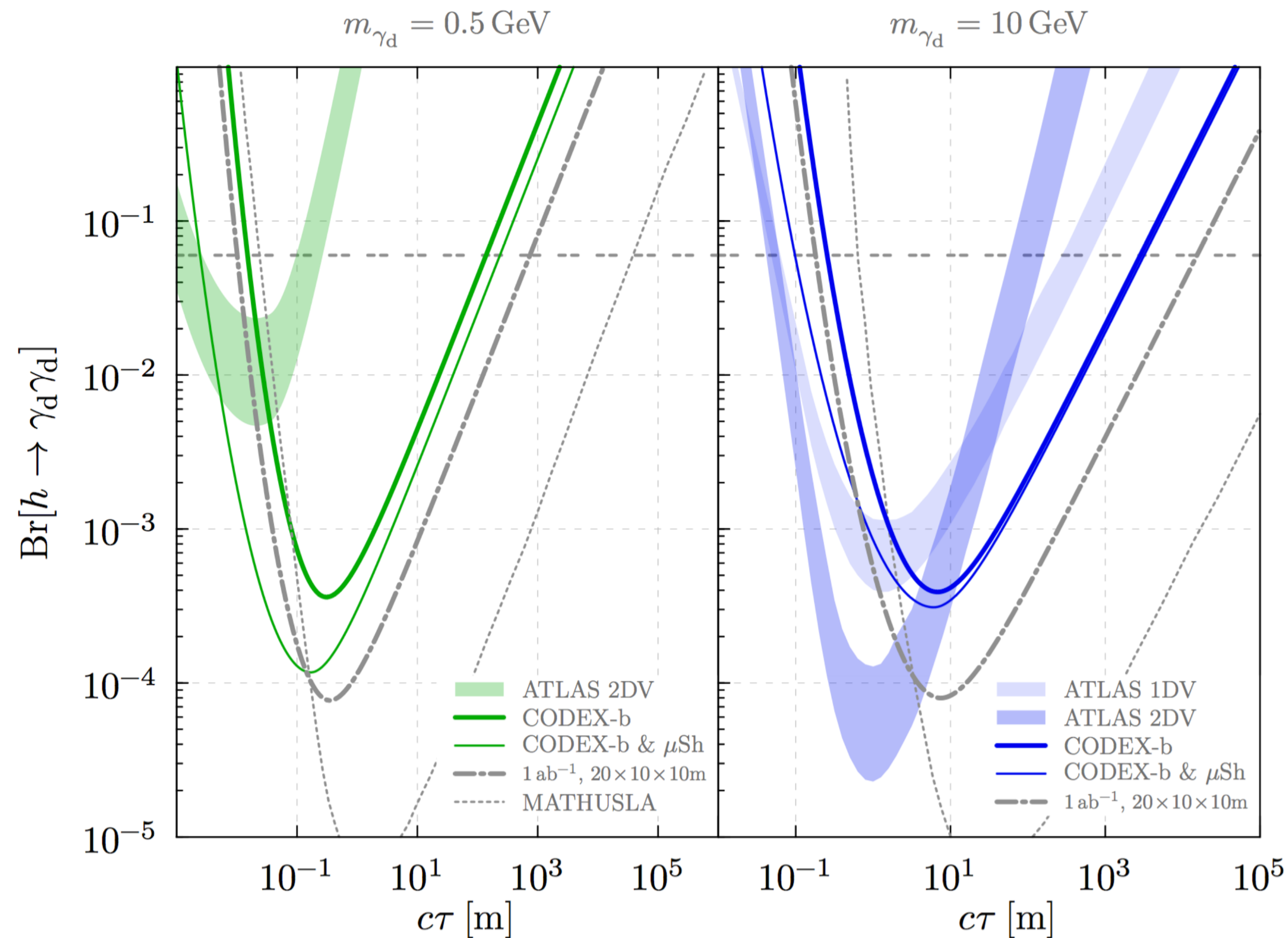


- Hope to get get enough events within around a week in the cavern.

Example model 1 — $b \rightarrow sX$

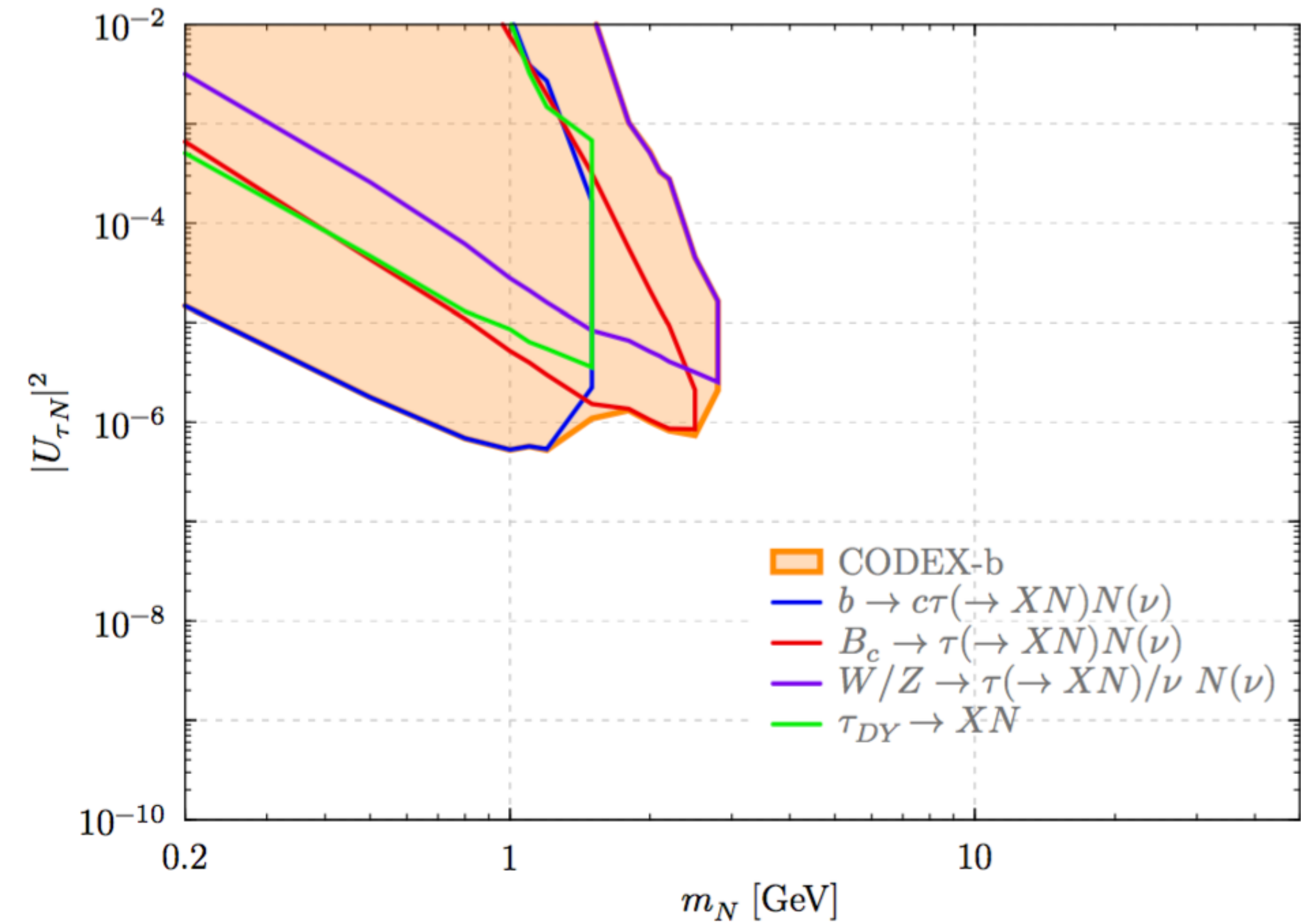
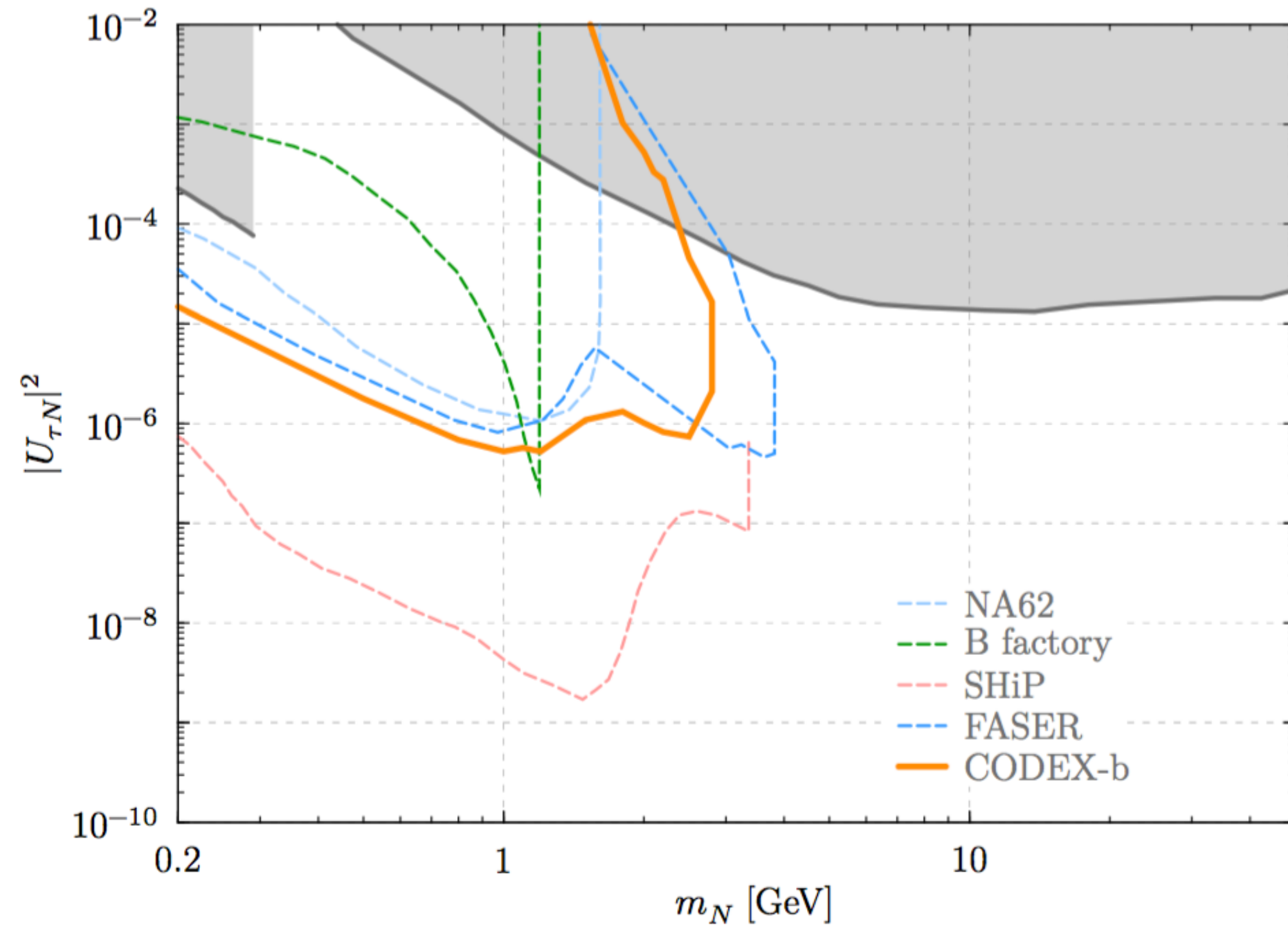


Example model 2 — $H \rightarrow \varphi\varphi$



Extends LHCb coverage far beyond ATLAS at low masses, competitive & complementary at higher ones. MATHUSLA has greater reach but backgrounds are uncorrelated.

Example model 3 — HNL



- Only weak limits set by current data!
- Reach dominated by B , B_c decays; W , Z in short lifetime regime

Tracker efficiency estimate

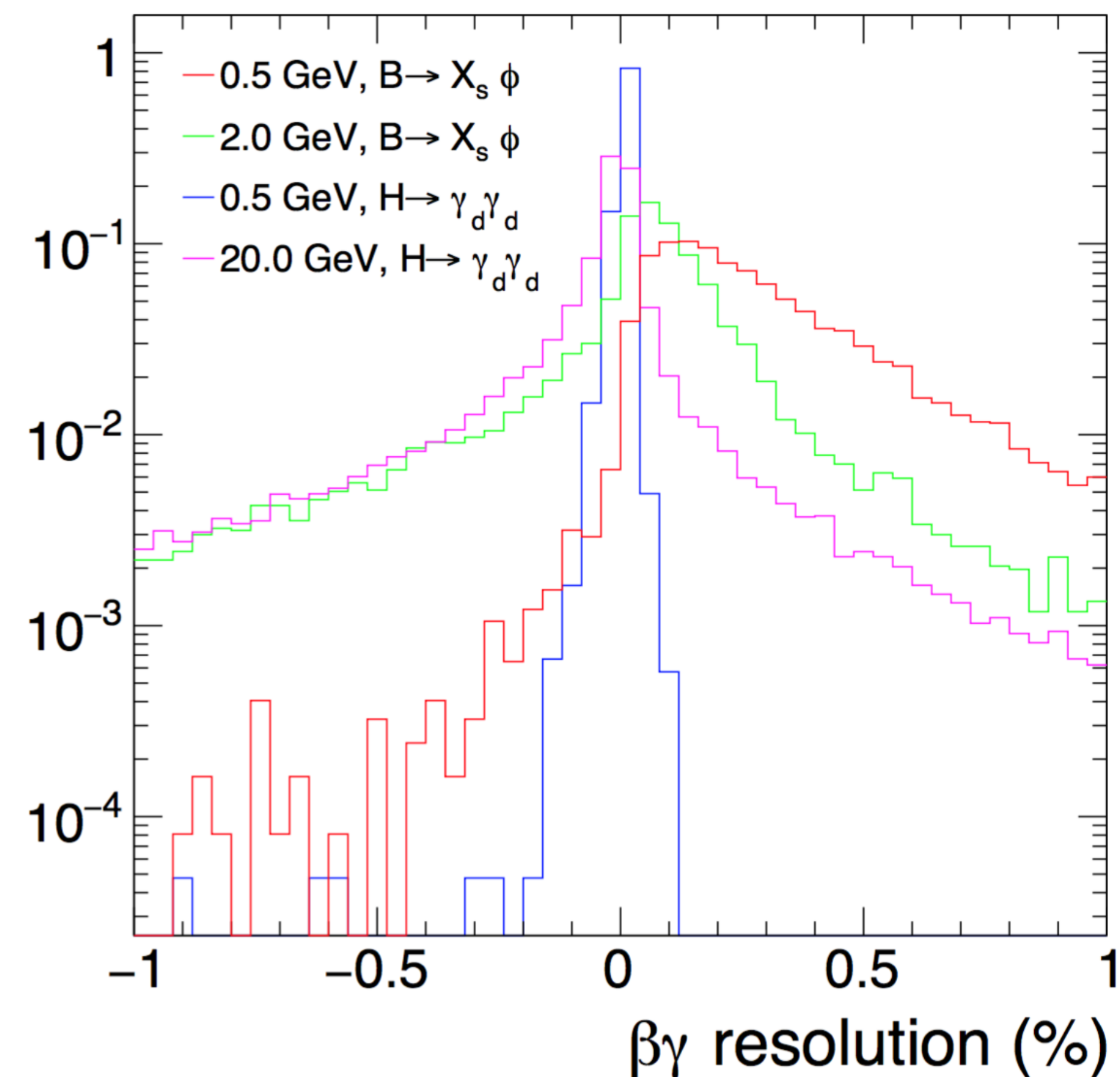
$c\tau$ (m)	$m_\varphi [B \rightarrow X_s \varphi]$			$m_{\gamma_d} [h \rightarrow \gamma_d \gamma_d]$				
	0.5	1.0	2.0	0.5	1.2	5.0	10.0	20.0
0.05	—	—	—	0.39	0.48	0.50	—	—
0.1	—	—	—	0.48	0.63	0.73	0.14	—
1.0	0.71	0.74	0.83	0.59	0.75	0.82	0.84	0.86
5.0	0.55	0.64	0.75	0.60	0.76	0.83	0.86	0.88
10.0	0.49	0.58	0.74	0.59	0.75	0.84	0.86	0.88
50.0	0.38	0.48	0.74	0.57	0.75	0.82	0.87	0.88
100.0	0.39	0.45	0.73	0.62	0.77	0.83	0.87	0.89
500.0	0.33	0.40	0.75	—	—	—	—	—

Dominated by partial overlap of decay products due to small opening angle, can be optimized using station spacing and granularity

Dominated by assumption that we don't track below 600 MeV of momentum, conservative since clearly we won't just fall off a cliff, but needs proper simulation

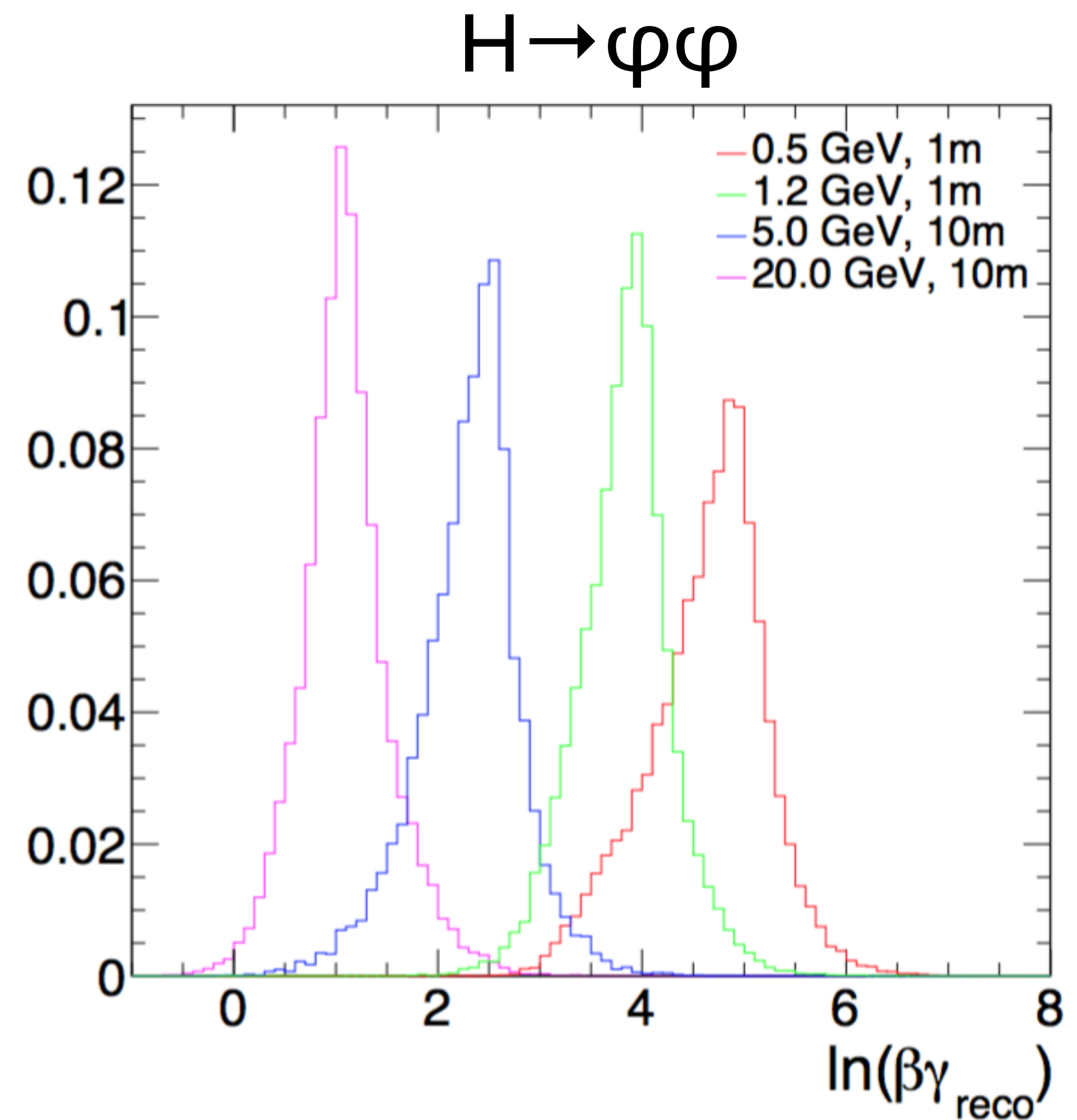
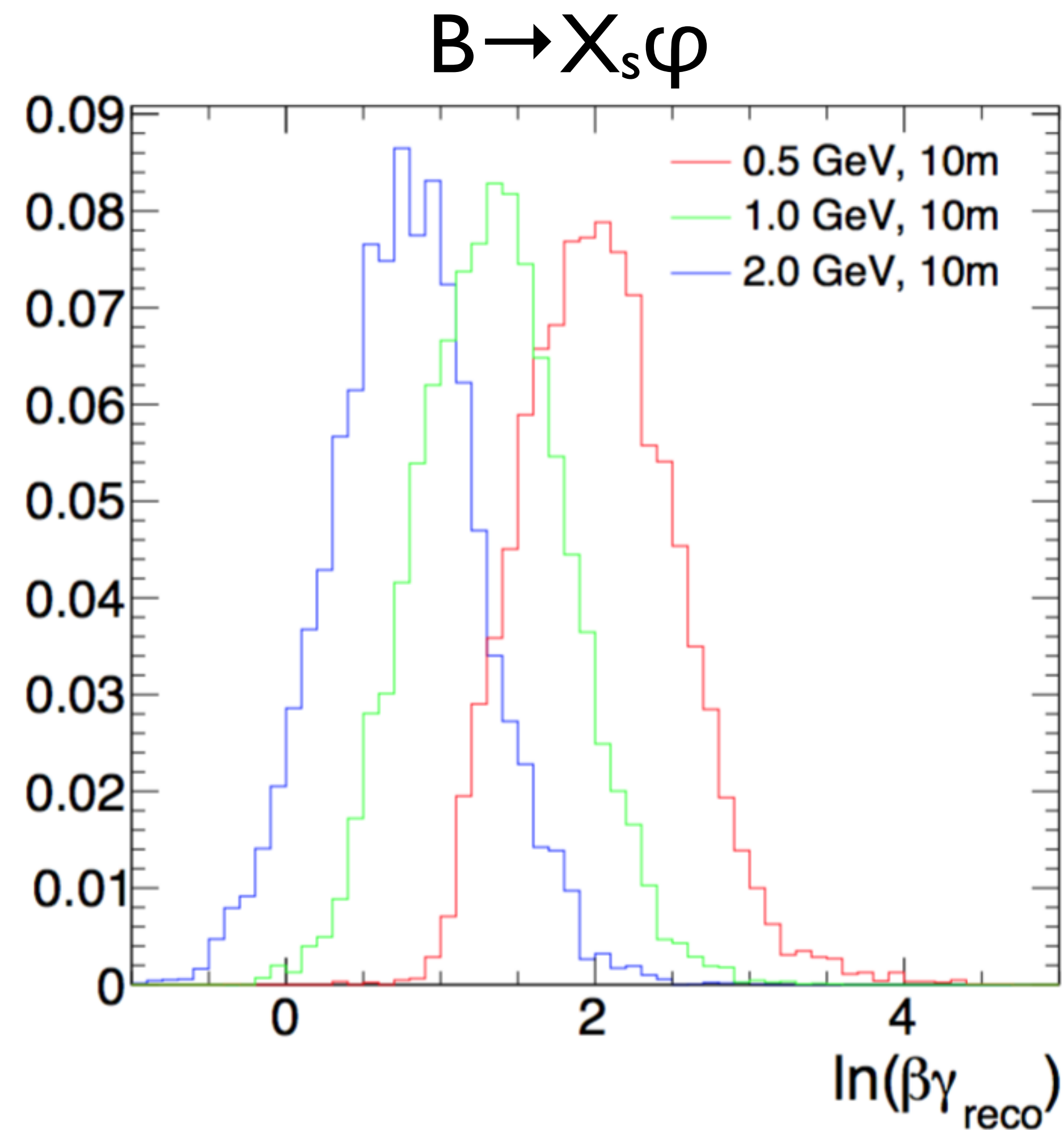
Bottom line : these are $O(1)$ numbers, not $O(\%)$, can be optimized further

Boost reconstruction



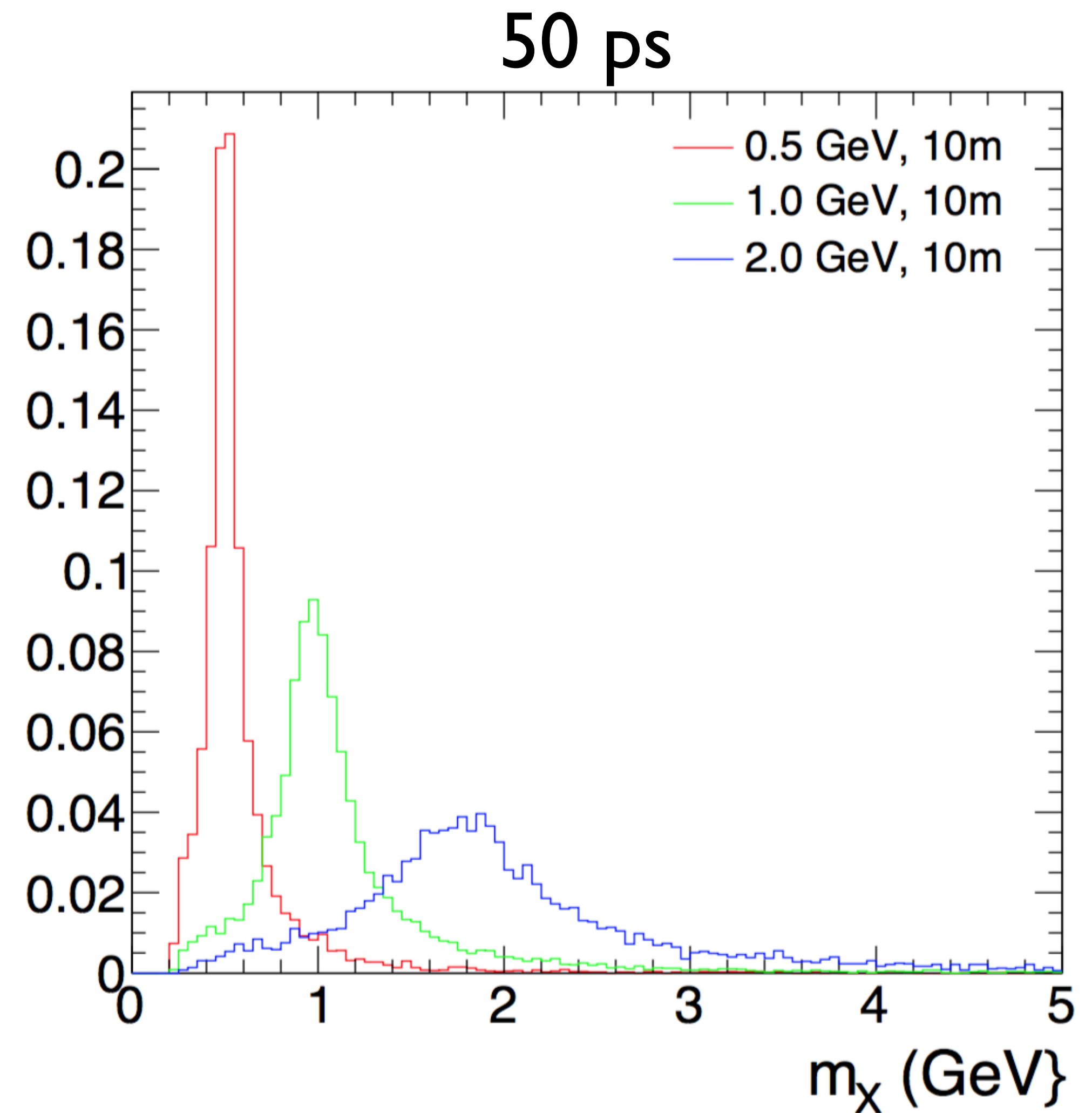
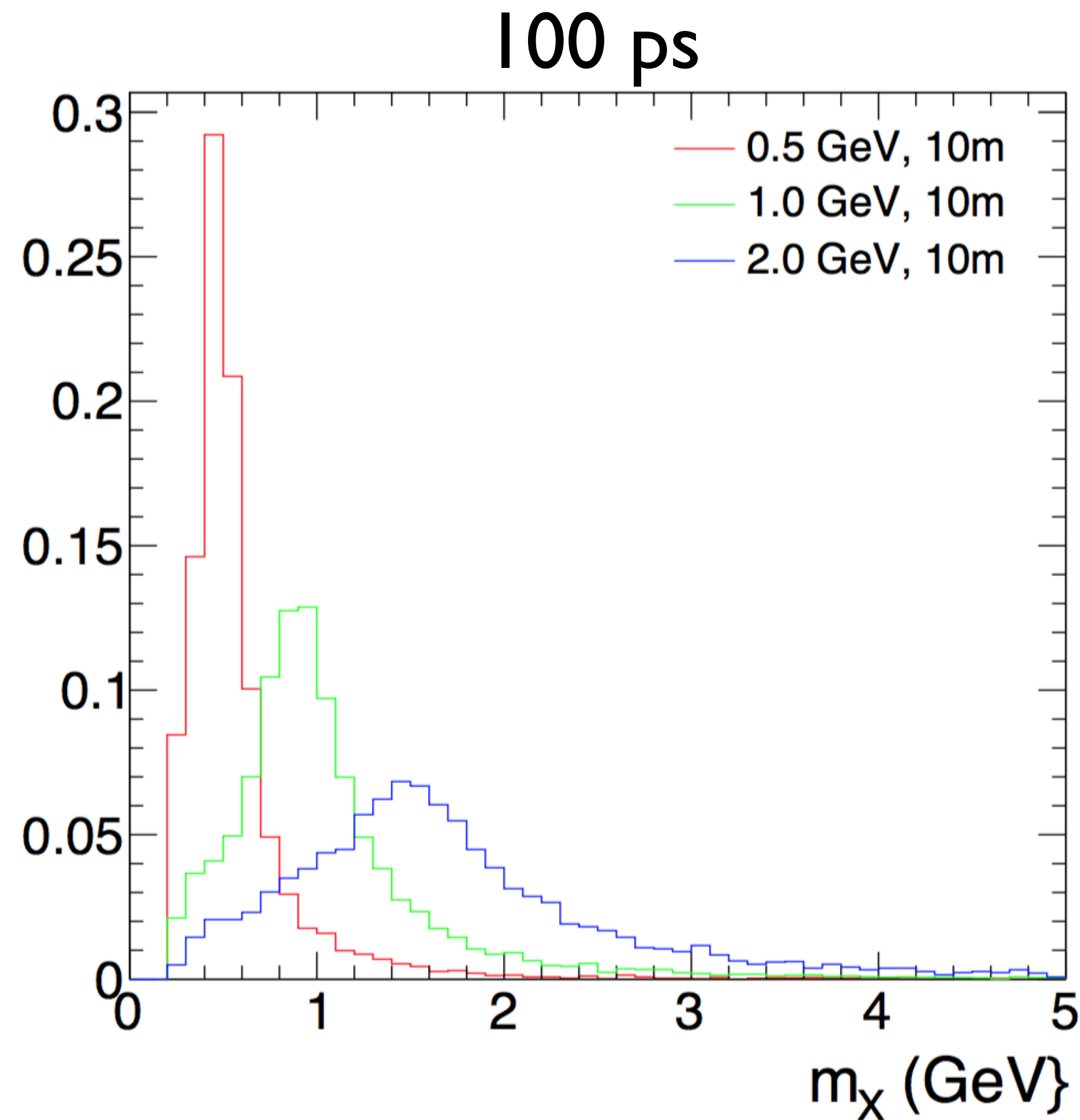
Reconstruct parent boost from the measured decay vertex (no timing!), assuming relativistic decay products. The resolution is $< 1\%$ (entirely dominated by distance to first measured point, not detector granularity) so the boost distribution is dominated by the generated spread of boosts, not resolution.

Boost reconstruction



Different initial states give different boost distributions; perhaps surprisingly we have some discriminating power between even the $B \rightarrow KX$ scenarios.

Mass reconstruction using time-of-flight



Now assume 100/50 ps time resolution (per hit) in the tracking stations. The $B \rightarrow KX$ signals are actually slow enough that we can reconstruct the X mass...

Conclusion

Outlook for LLP searches

No theory guidance on lifetime → large detectors

Many possible decay modes → hermeticity, particle ID

Small coupling and production rate → zero background

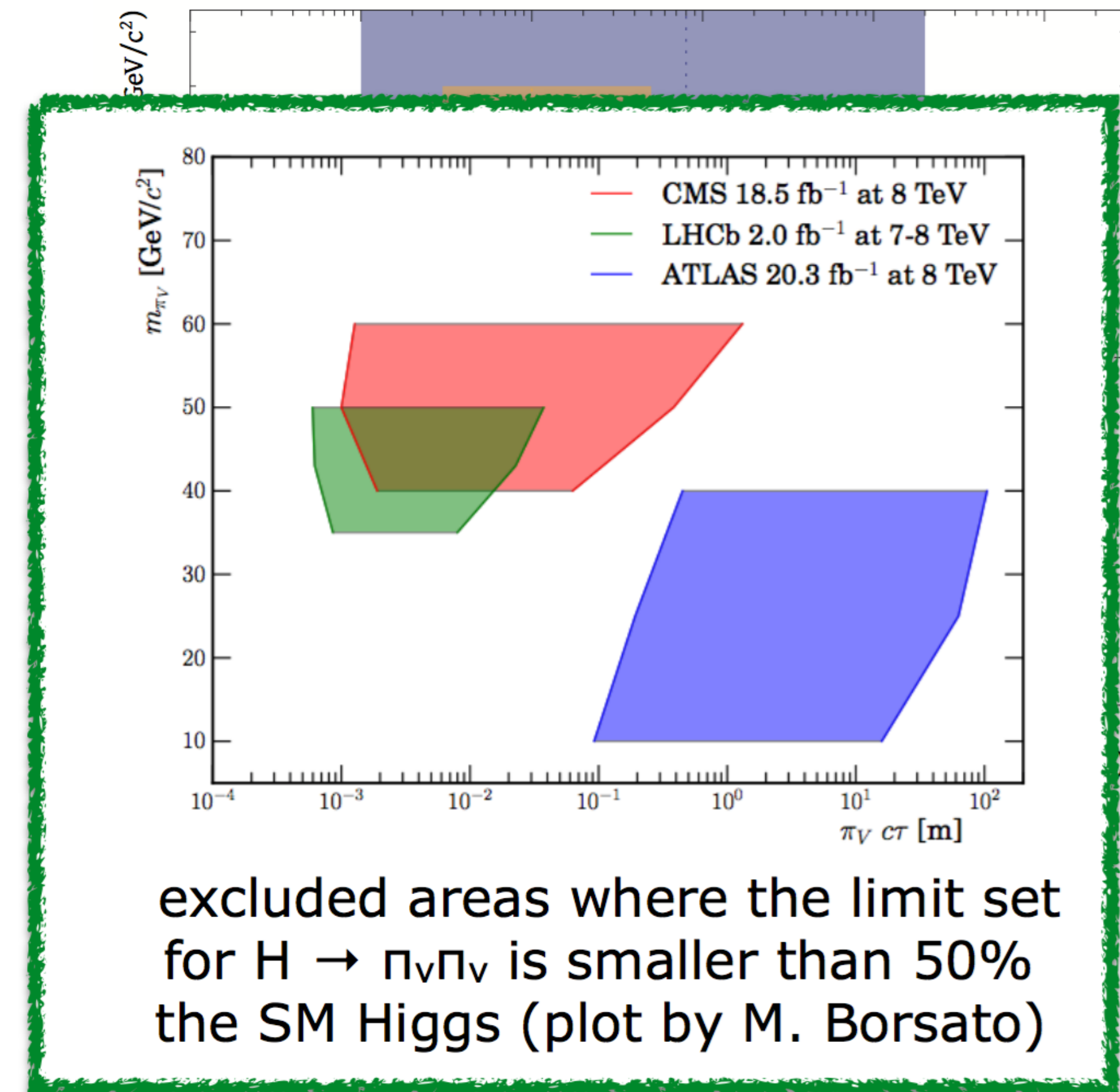
Small coupling and production rate → huge integrated lumi

Very hard for any single detector to meet all these criteria!
The proposed experiments overlap in reach but are complementary
in assumptions and backgrounds – critical if signal is seen.

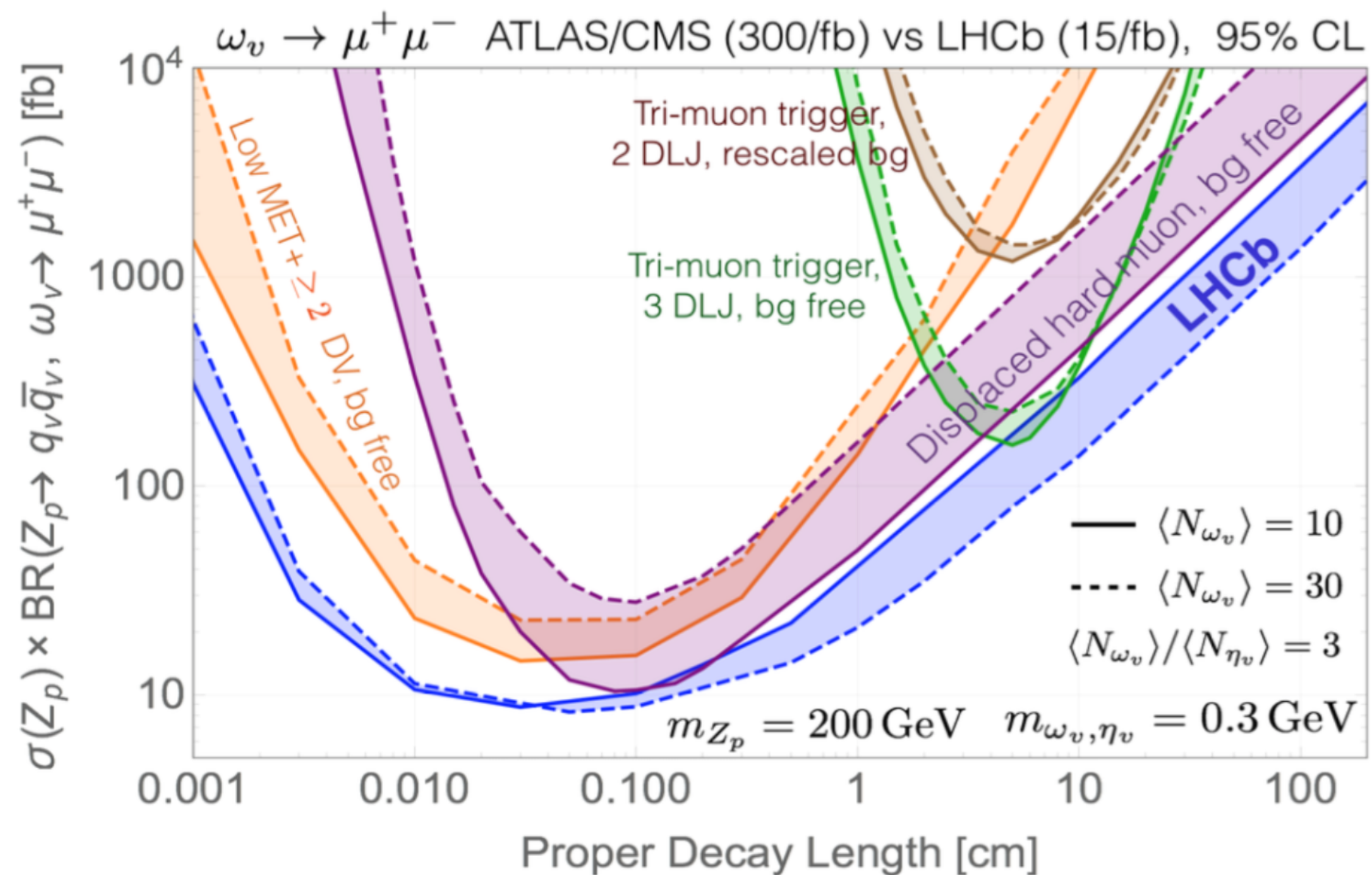
Backups

LHCb already complements ATLAS/CMS

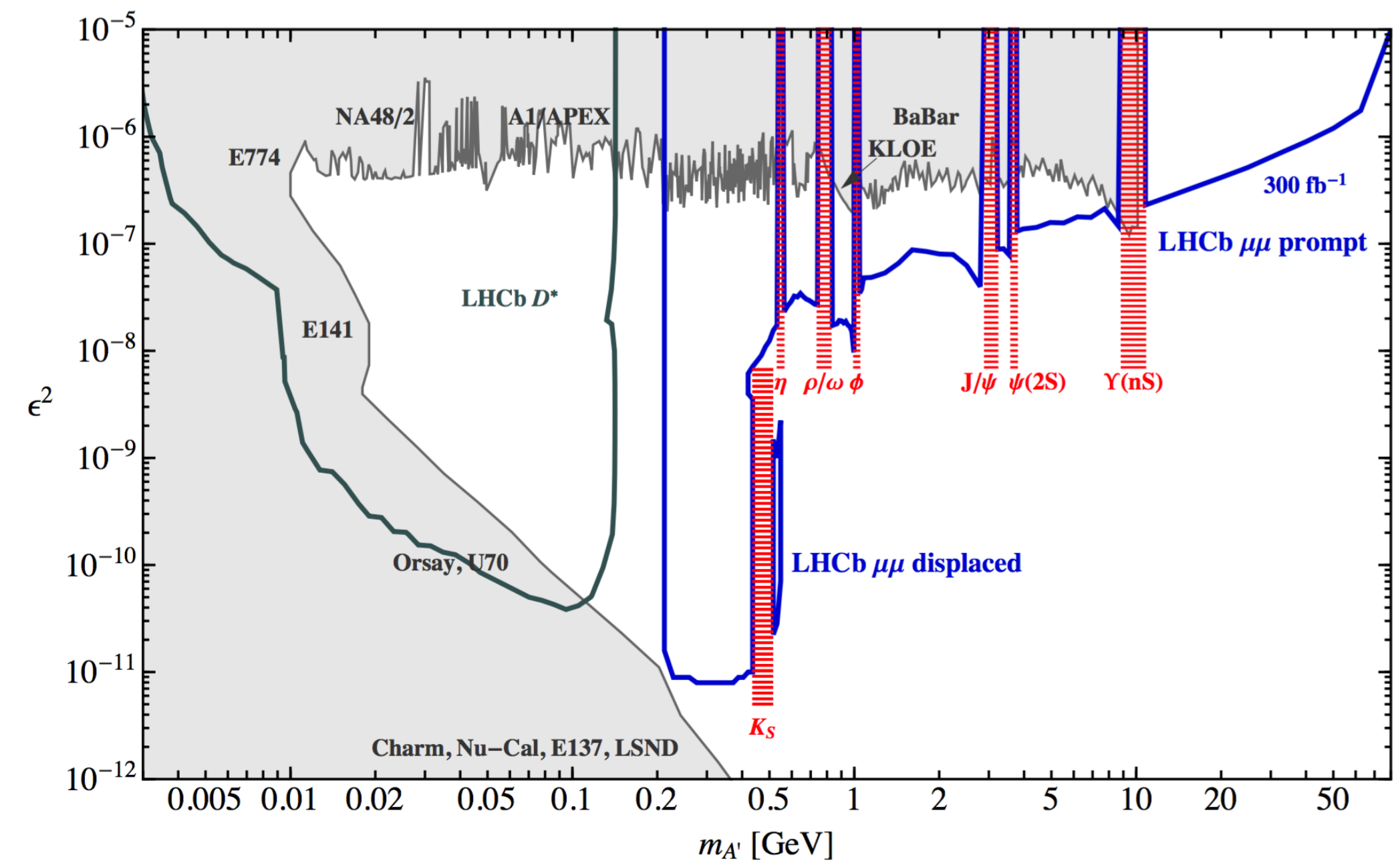
- ◆ Obvious disadvantage: LHCb collects less data than ATLAS/CMS and has worse acceptance for several searches
- ◆ But softer triggers (for instance, can trigger detached di-muons with $p_T \sim 1$ GeV/c), other advantages already mentioned
- ◆ In practice that means we can look into **complementary** phase space regions



So is something more needed?



Pierce et al. <https://arxiv.org/abs/1708.05389>



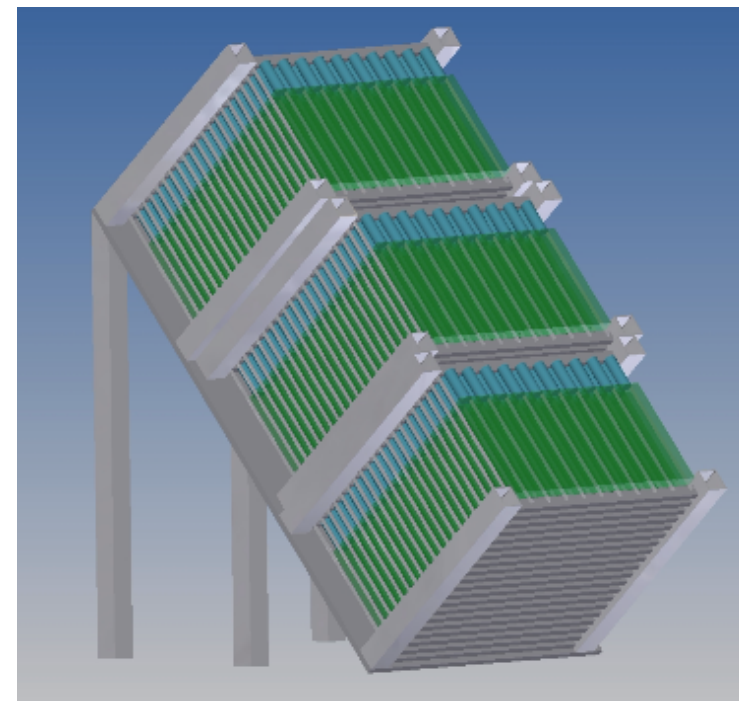
Itten et al.

<https://arxiv.org/abs/1509.06765>

<https://arxiv.org/abs/1603.08926>

LHCb reach worked out in certain scenarios, above showing two of them – you can see again that we can complement ATLAS/CMS for very light signals, up to a certain c τ region which is basically limited by the position of the TT where we need hits for a momentum measurement. Can we expand towards larger c τ values?

Other ideas targeting LLP's



MilliQan: 1607.04669

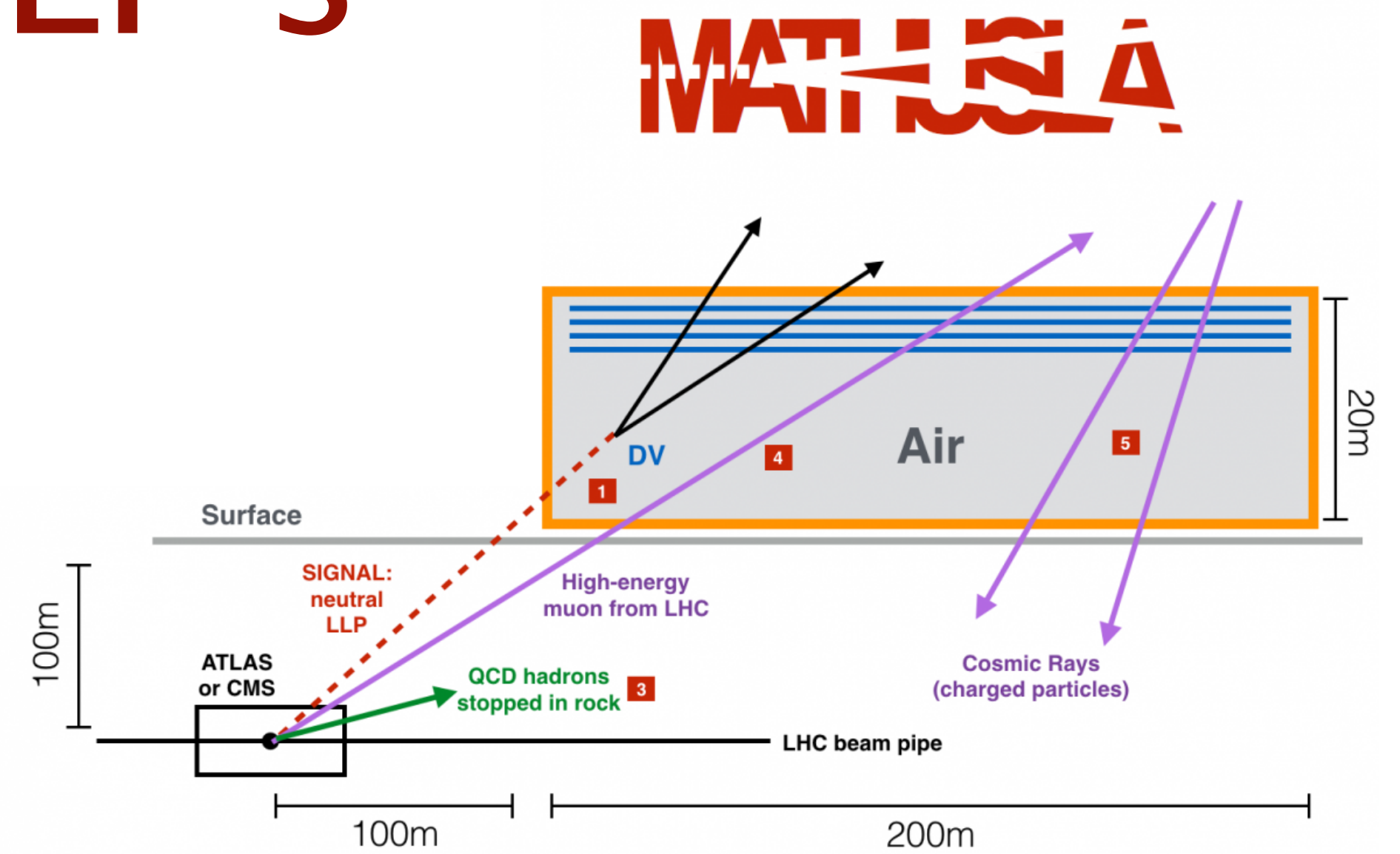
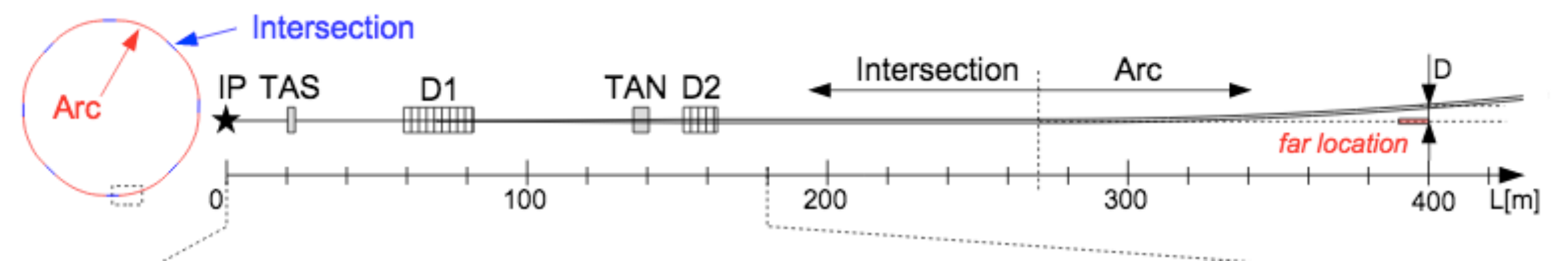
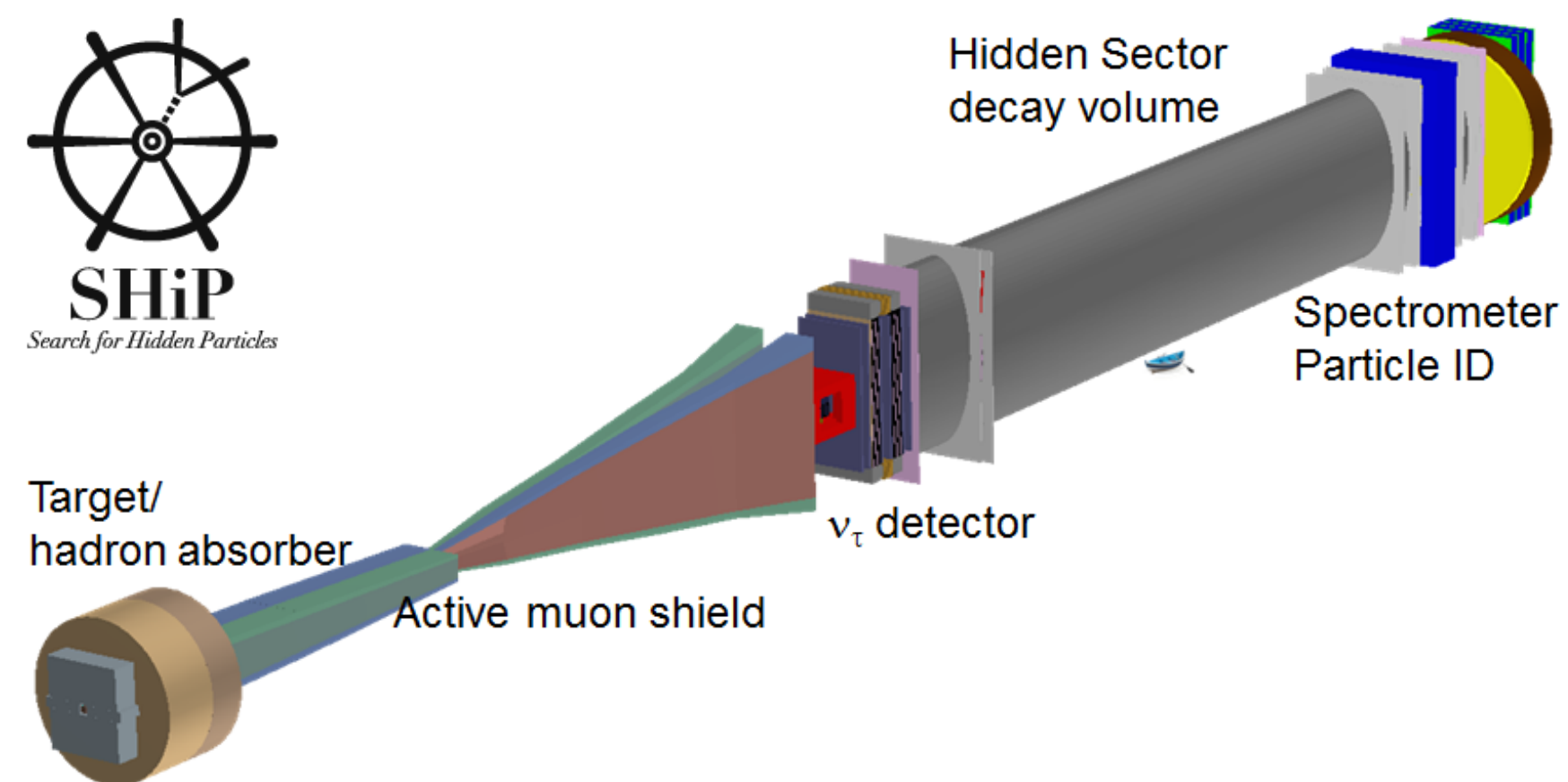


image by D. Curtin and R. Sundrum



FASER: 1708.09389

All proposals feature a substantial amount of shielding to suppress backgrounds. Useful geometric acceptance tends to require large detectors.

Integration with LHCb

It is highly desirable to treat CODEX-b as an additional subdetector of LHCb, and to integrate it into the DAQ & readout.

Allows events which look interesting in CODEX-b (whose rate is low by definition) to be saved in LHCb as well. If we see a signal we could then look at the event in LHCb and see if an interesting tag exists there.

You may think Phase II pileup would make this prohibitive, but that is not an immediate showstopper if both CODEX-b and LHCb give precise timing information.

A tricky bit is that CODEX-b “events” are offset by around ~80 ns wrt. the LHC collision which produced them, but should be manageable.

Data driven background calibration

Cosmics will be used for spatial & time detector alignment and their negligible contribution can be calibrated from this.

Other backgrounds can be measured by putting a small telescope in the LHCb cavern and measuring background rates with different shield thicknesses.

Could be done as an engineering run well ahead of full detector construction.

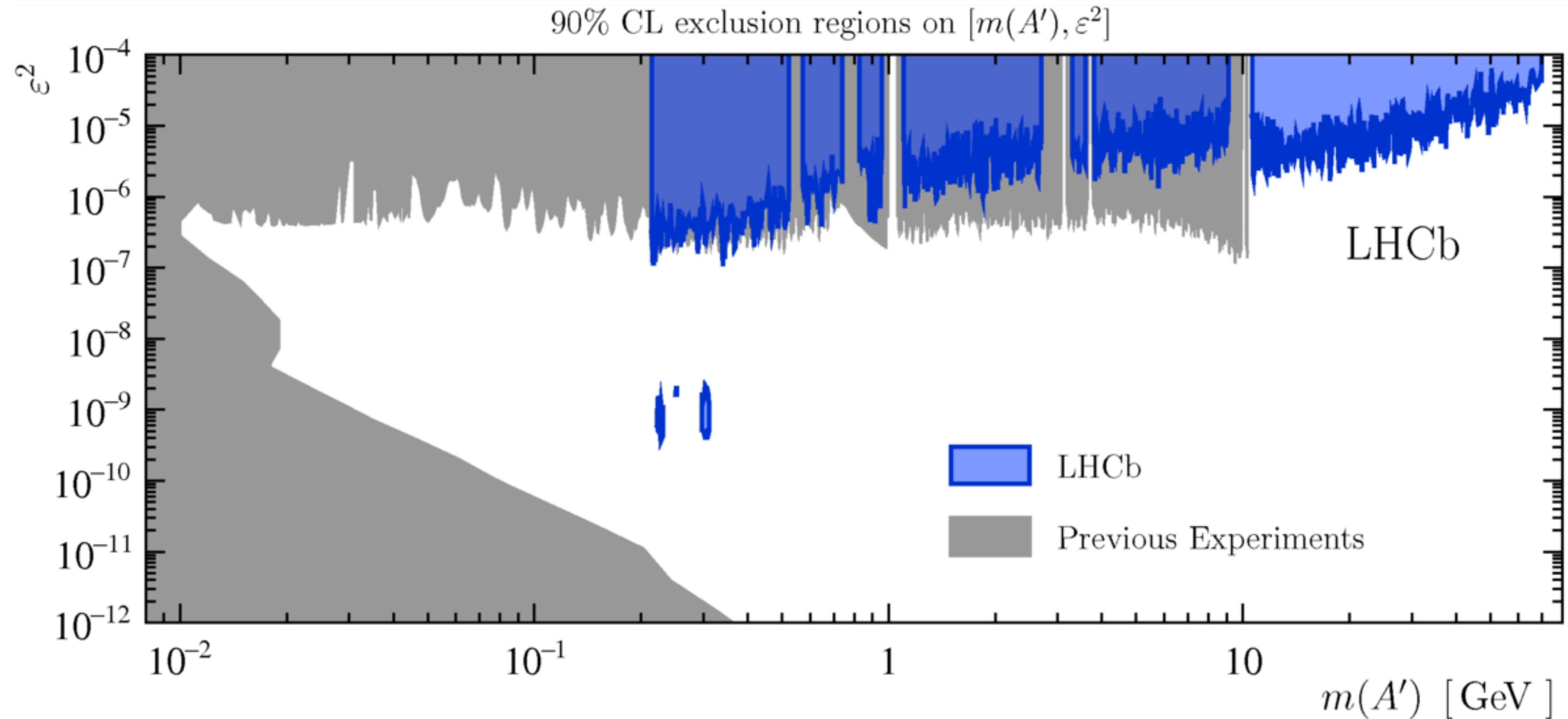
Complementarity with other searches

CODEX-b can cover a significant portion of parameter space for well-motivated, simple portals, and extend LHCb's reach for long lived particles well beyond ATLAS/CMS.

CODEX-b has to cover around 1/100th of MATHUSLA's tracking area (but of course does not have as large an absolute reach).

If you believe the physics case for LLP detection is worthwhile, allocating funds for a detector which is relatively simple to build, has complementary reach to more ambitious proposals, and has completely different backgrounds would seem prudent, particularly if someone sees a signal.

Dark photon example



The smaller the coupling, the smaller the production rate
Hence plots like this ([LHCb-PAPER-2017-038](#)) : no sensitivity to
directly produced long-lived dark photons above a certain mass.