High-Scale SUSY and the Higgs from a Stringy Perspective

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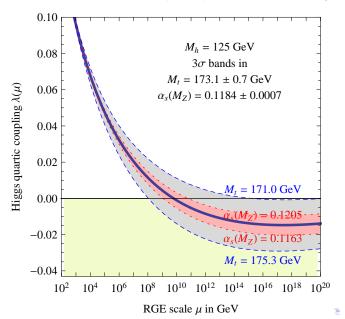
cf. 1204.2551 and 1304.2767 with A. Knochel and T. Weigand

<u>Outline</u>

- In the SM, the 'vacuum stabilty scale'  $\mu_\lambda$  has emerged as new, important piece of data.
- The vanishing of  $\lambda$  could be related to the breaking approximately shift-symmetric SUSY at that scale.
- This situation arises naturally in the stringy context.
- Alternatively, a SUSY (NMSSM-type) UV completion might appear only far above the scale  $\mu_{\lambda}$ .

NNLO, from Degrassi,.., Espinosa et al., 1205.6497

see also Bezrukov, Kniehl et al. '12; for top-mass precision: Moch et al. '15/'16



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Our perspective:

- The critical scale  $\mu_{\lambda}$  could be anywhere between 10<sup>9</sup> and 10<sup>17</sup> GeV.
- Below  $\mu_{\lambda}$ : just SM.
- The weak scale is fine-tuned; The UV completion is stringy; SUSY is motivated by the stability of known string models.
- $\lambda = 0$  is the result of a shift-symmetry in the Higgs sector, together with SUSY -breaking at that scale.

The subject has a long history...

• It has always been well-known that, for low  $m_h$ ,  $\lambda$  runs to zero at some scale  $< M_P$  (vacuum stability bound)

> Lindner, Sher, Zaglauer '89 Froggatt, Nielsen '96 Gogoladze, Okada, Shafi '07

Shaposhnikov, Wetterich 09' Giudice, Isidori, Strumia, Riotto, ... Masina '12

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- Many attempts were made to turn this into an  $m_h$  prediction
- Let us mention two ideas in more detail...

Higgs mass prediction from  $\lambda = 0$  at  $M_P$ 

(Shaposhnikov, Wetterich, 0912.0208)

- One might hope that  $\lambda = 0$  emerges in at  $M_P$  as a result of quantum gravity
- In 2009, with  $m_t \simeq 171$  GeV, this gave a prediction of  $m_h = 126$  GeV
- With today's data, this works less well, but is still an option
- The underlying theory assumes a non-perturbative UV fixed point of gravity (asymptotic safety)

Weinberg '79; Reuter '98; Reuter et al. '98...'11

• But it is far from clear why  $\lambda = 0$  should come out and who tunes the Higgs mass small....

Higgs mass prediction from  $\lambda = 0$  at 'unification scale'

(Gogoladze, Okada, Shafi, 0705.3035 and 0708.2503)

- 5d Gauge-Higgs unification  $\rightarrow$  flat Higgs potential
- Based on non-SUSY SM gauge unification (with non-canonical U(1)), one finds a unification scale of 10<sup>16</sup> GeV
- A prediction of  $m_h = 125 \pm 4$  GeV was made
- Obviously, there is strong model dependence in the non-SUSY GUT sector, .... other 'predictions' are possible ....

for another related suggestion see Redi/Strumia '12

String-phenomenologist's perspective

- No strong preference for a particular SUSY breaking scale
- Natural guess: The special scale  $\mu(\lambda = 0)$  is the SUSY-breaking scale
- Crucial formula:

$$\lambda(m_s) = \frac{g^2(m_s) + g^{\prime 2}(m_s)}{8} \cos^2(2\beta)$$

• Reminder:

$$M_{H}^{2} = \begin{pmatrix} |\mu|^{2} + m_{H_{d}}^{2} & b \\ b & |\mu|^{2} + m_{H_{u}}^{2} \end{pmatrix} = \begin{pmatrix} m_{1}^{2} & m_{3}^{2} \\ m_{3}^{2} & m_{2}^{2} \end{pmatrix}$$

$$\sin(2\beta) = \frac{2m_3^2}{m_1^2 + m_2^2}$$

Need this to be 1!

• Our goal:

Identify a special structure/symmetry leading to  $\tan\beta=1$  (i.e. to  $\lambda=0$  )

Indeed, such a structure is known in heterotic orbifolds:

Shift symmetric Kahler potential:

$$K_H \sim |H_u + \overline{H}_d|^2$$

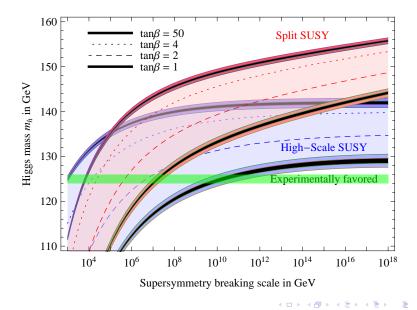
Lopes-Cardoso, Lüst, Mohaupt '94 Antoniadis, Gava, Narain, Taylor '94 Brignole, Ibanez, Munoz, Scheich, '95...'97

Note: The actual shift symmetry transformation is

$$H_u \to H_u + \alpha$$
,  $H_d \to H_d - \overline{\alpha}$ .

This guarantees a light doublet, even after SUSY breaking.

NNLO, from Degrassi,.., Espinosa et al., 1205.6497 Predicted range for the Higgs mass



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In more detail:

$$K_H = f(S,\overline{S})|H_u + \overline{H}_d|^2$$

Assuming  $F_S \neq 0$  and  $m_{3/2} \neq 0$  this gives

$$m_1^2 = m_2^2 = m_3^2 = \left| m_{3/2} - \overline{F}^S f_{\overline{S}} \right|^2 + m_{3/2}^2 - F^S \overline{F}^S (\ln f)_{S\overline{S}}$$

• One of the concrete modern realizations is through F-theory GUTs, with a D7-brane 'bulk Higgs' à la

Donagi, Wijnholt, '11

• In this setting one expects  $F^{S} = 0$  and hence

$$m_i^2 = 2m_{3/2}^2$$
.

• To understand the physical origin of the shift symmetry, the simplest context is that of orbifold GUTs

K. Choi et al. '03 AH, March-Russell, Ziegler '08 Brümmer et al. '09...'10 Lee, Raby, Ratz, Ross, ... '11

• Indeed, one has the 5d breaking pattern

 $SU(6) \rightarrow SU(5) \times U(1); \quad 35 = 24 + 5 + \overline{5} + 1.$ 

• The Higgs arises from the scalar and vector of the 5d multiplet

 $\mathsf{Higgs}_{5,\overline{5}} = \Sigma + iA_5$ 

(the shift symmetry comes from 5d gauge trfs.,  $A_5 \rightarrow A_5 + c$ ).

cf. Gogoladze, Okada, Shafi '07

- This also happens in full-fledged 10d heterotic string constructions, but it is much more generic: heterotic WLs ↔ type IIA / D6-WLs ↔ type IIB / D7-WLs or positions
- These and other origins of the Higgs-shift-symmetry and of  $\tan\beta=1$  have also been explored in

Ibanez, Marchesano, Regalado, Valenzuela '12 Ibanez, Valenzuela '13

In particular, they observe that to get tan β = 1,
a Z<sub>2</sub> exchange symmetry acting on H<sub>u</sub>, H<sub>d</sub> is sufficient;
the rest is done by the usual tuning...

$$M_H^2 = \left(\begin{array}{cc} m_1^2 & m_3^2 \\ m_3^2 & m_2^2 \end{array}\right)$$

## <u>Comments</u>

- Clearly, we eventually need more phenomenological implications of 'stringy high-scale SUSY' (e.g. in cosmology)
- For example, axion(s), cosmological moduli and a possible 'dark radiation sector' can be potentially related to the high SUSY-breaking scale

Chatzistavrakidis, Erfani, Nilles, Zavala '1206... Higaki, Hamada, Takahashi '1206... Cicoli, Conlon, Quevedo,... Angus,... '12...'13

• The situation concerning non-SUSY F-theory unification in this context is interesting but complicated....

Ibanez et at. '12 AH, Unwin, '14

Returning to our shift-symmetry proposal we now ask about

Corrections? Precision?

- The superpotential (e.g. top Yukawa) breaks the shift symmetry
- The crucial point is compactification

Shift symmetry is exact (gauge symmetry!) in 10d. The shift corresponds to switching on a WL. This is not a symmetry in 4d (4d-zero modes 'feel' the WL). 4d-loops destroy the shift symmetry of Kähler potential.

• Optimistic approach to estimating the 'goodness' of our symmetry:

Symmetry-violating running between  $m_c$  and  $m_s$  $\Rightarrow$  Correction  $\delta \sim \ln(m_c/m_s)$ 

## More explicitly:

$$M_{H}^{2} = (|\mu|^{2} + m_{H}^{2}) \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} + \begin{pmatrix} \delta |\mu|^{2} + \delta m_{H_{d}}^{2} & \delta b \\ \delta b & \delta |\mu|^{2} + \delta m_{H_{u}}^{2} \end{pmatrix}$$

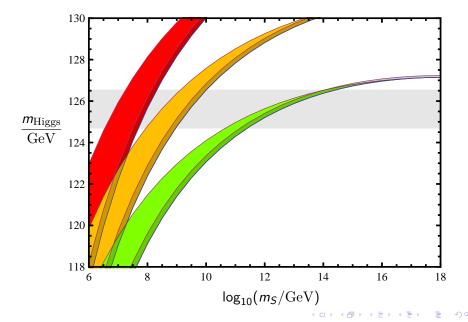
= symmetric + loop violation

• Leading effects: y<sub>t</sub> and gauge

$$\delta M_H^2 = f(\epsilon_y, \epsilon_g, m_{\text{soft}})$$
 ;  $\epsilon_y = \int_{\ln m_s}^{\ln m_c} dt \, \frac{6|y_t|^2}{16\pi^2}$ 

• Enforce det  $M_H^2 = 0$  after corrections  $\Rightarrow \epsilon_y, \epsilon_g, m_{\text{soft}}$  are related

 $\cos 2\beta = \epsilon_y \times \{ \text{calculable } \mathcal{O}(1) \text{ factor} \}$ 



Another type of corrections:

with

$$\delta\lambda_{TH}(m_{S}) = \frac{3y_{t}^{4}}{16\pi^{2}} \Big[ \frac{X_{t}^{2}}{m_{S}^{2}} \Big( 1 - \frac{X_{t}^{2}}{12m_{S}^{2}} \Big) + 2\log(\frac{m_{\tilde{t}}}{m_{S}}) \Big]$$

$$X_t = A_t - \mu \cot eta pprox A_t - \mu$$

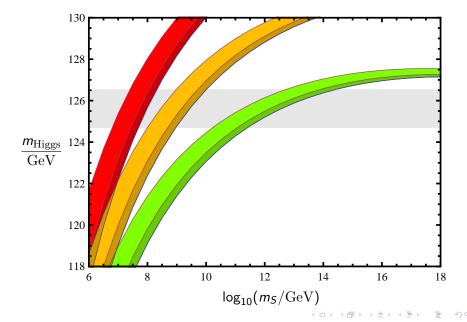
• For  $X_t^2 = 0 \dots 6m_S^2$ , they are in the range

$$\delta\lambda_{TH}(m_S) = 0\dots 3 imes rac{3y_t^4}{16\pi^2}$$

 These are qualitatively different from SUSY thresholds and should hence presumably not be absorbed in an 'effective SUSY breaking scale'

Drees, priv. comm.

A-term corrections for 
$$X_t^2 = m_S^2$$
 and  $X_t^2 = 6m_S^2$ 



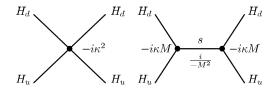
From unstable high-scale

to metastable low-scale theories:

- So far, we argued that SUSY should appear at least at the scale μ<sub>λ</sub>.
- However, this can avoided with very little extra effort:
- Let string theory produce a high-scale NMSSM, with a large supersymmetric mass *M* for the singlet *S*,

$$W = \kappa S H_u H_d + \frac{1}{2} M S^2 \,.$$

• Clearly, integrating out S will not induce a quartic coupling due to a supersymmetric cancellation...



 However, adding additionally a negative soft mass-squared upsets this cancellation and gives a negative quartic effect:

Giudice/Strumia '11

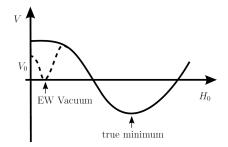
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$$V_{\Lambda=M} \quad \supset \quad \kappa^2 rac{m_s 2}{M^2 + m_s^2} \left| H_u H_d \right|^2.$$

- We propose to make this effect large enough to produce a non-negligible λ < 0 at the scale m<sub>s</sub>.
- We also still have a shift-symmetric Kahler potential and hence tan β = 1 at LO.

- Our theory is now weakly unstable at the SUSY breaking scale.
- This is cured in the UV $\rightarrow$ IR RGE-running:



- 'Our' minimum is generated only radiatively, as  $\lambda$  runs from negative to positive values in a loop-calculation based on an unstable vacuum.
- Thus, we have a simple UV completion of the meta-stable SM.

• This setting is reminicsent of situations with tachyonic high-scale soft masses

see e.g. Dermisek/Kim '06 Ellis/Lebedev/Olive/Srednicki '08

• It is interesting to work out the cosmology of this setting in more detail...

Abel/Chu/Jaeckel/Khoze '06 Lebedev/Westphal '12

see also recent work by Enqvist, Lebedev, Karciauskas, Rusak, Zatta, Gross, ... Espinosa et al.

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## Conclusions / Summary

- It is conceivable that we have to expect new physics not at a TeV, but only at the 'vacuum stability scale' μ<sub>λ</sub>.
- Well-motivated guess: SUSY broken with tan  $\beta = 1$  at  $\mu_{\lambda}$
- Possible structural reason: shift symmetry in Higgs sector (Predictivity, i.e. m<sub>h</sub>, m<sub>t</sub>, α<sub>s</sub> ⇒ m<sub>s</sub> remains strong, even if shift symmetry is only approximate)

- But: SUSY breaking above  $\mu_{\lambda}$  with  $\lambda < 0$  is also possible
- This is a very natural UV-completion for the minimalist 'metastability scenario' without new physics near μ<sub>λ</sub>.