

Yang-Mills thermodynamics

Ralf Hofmann
Universität Karlsruhe (TH)

“Symmetry in Nonlinear Mathematical Physics”
Institute of Mathematics, National Academy of Science, Kyiv, Ukraine

June 28, 2007



Universität Karlsruhe (TH)

Forschungsuniversität • gegründet 1825

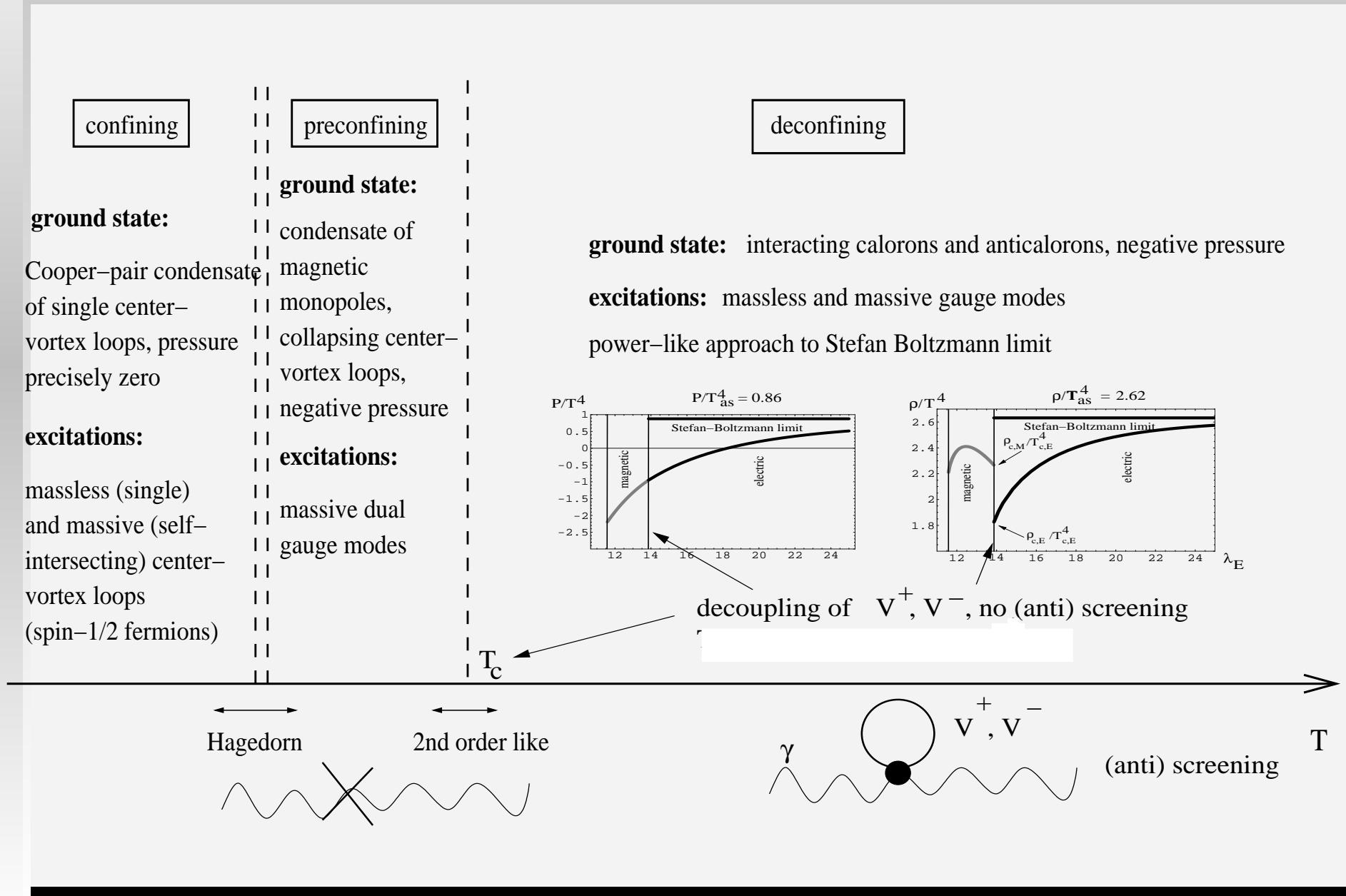
plan

- ▶ brief motivation and preview on phase diagram
[hep-th: 0411214, 0504064, 0609033, 0609172, 0702027]
- ▶ deconfining ground-state physics:
coarse-grained, interacting calorons
- ▶ coarse-grained excitations:
Legendre-trafos and loop expansion
- ▶ preconfinement:
cond. magn. monopoles, dual Meissner effect
- ▶ low temperatures:
Hagedorn, flip of statistics, Borel summation
- ▶ summary, conclusions, mention of applications

Why nonpert. YM TD?

- ▶ infrared instability of PT even for $T \gg \Lambda$ in magnetic sector
[Linde 1980]
- ▶ highly nonpert. ground-state physics even for $T \gg \Lambda$:
 - $\theta_{\mu\mu} \propto T$
[Miller 1998]
 - spatial string tension: $\sigma \propto T^2$
[Philipsen 1998, Korthals-Altes 1998, ...]
- ▶ no lattice control at low temperature:
 - correlation length larger than linear lattice size
 - analytical grasp \Rightarrow equilibrium violated

preview: phase diagram SU(2)



deconfining ground state

- ▶ coarse-grained (anti)calorons of $|Q| = 1$
⇒ adjoint scalar field ϕ^a , $|\phi|$ spatially homogeneous
- ▶ strategy:
 - thermodynamics $\Rightarrow \phi^a$ periodic in eucl. time
in any admissible gauge \Rightarrow
phase $\hat{\phi}^a$ determined by *classical* configs.
 - stable configs.: $|Q| = 1$ HS (anti)calorons (BPS)
of trivial holonomy (only these enter!)

- compute $\hat{\phi}^a \in (\text{Kernel of } \mathcal{D})$ by respecting isotropy and $S_{\text{HS}} = \frac{8\pi^2}{g^2} \neq f(T, \Lambda)$ in *inf.-vol.* average over magnetic-magnetic correlation mediated by *single* (anti)caloron
- fixes \mathcal{D} uniquely \Rightarrow winding number
- impose BPS $\Rightarrow \hat{\phi}^a$
- average saturates rapidly
 \Rightarrow scale Λ and analyticity in ϕ^a
 \Rightarrow RHS of BPS eq. for ϕ^a
 \Rightarrow ϕ 's potential and saturation scale $|\phi|$
 \Rightarrow inertness of $|\phi|$

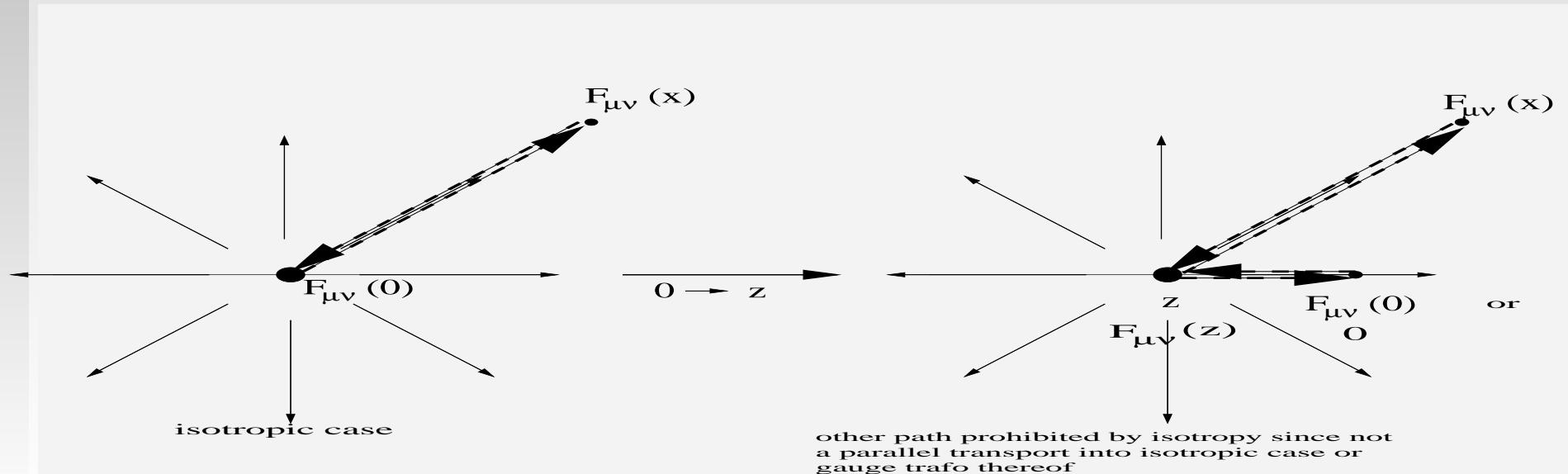
technically:

(integration over $S_3^{R=\infty}$)

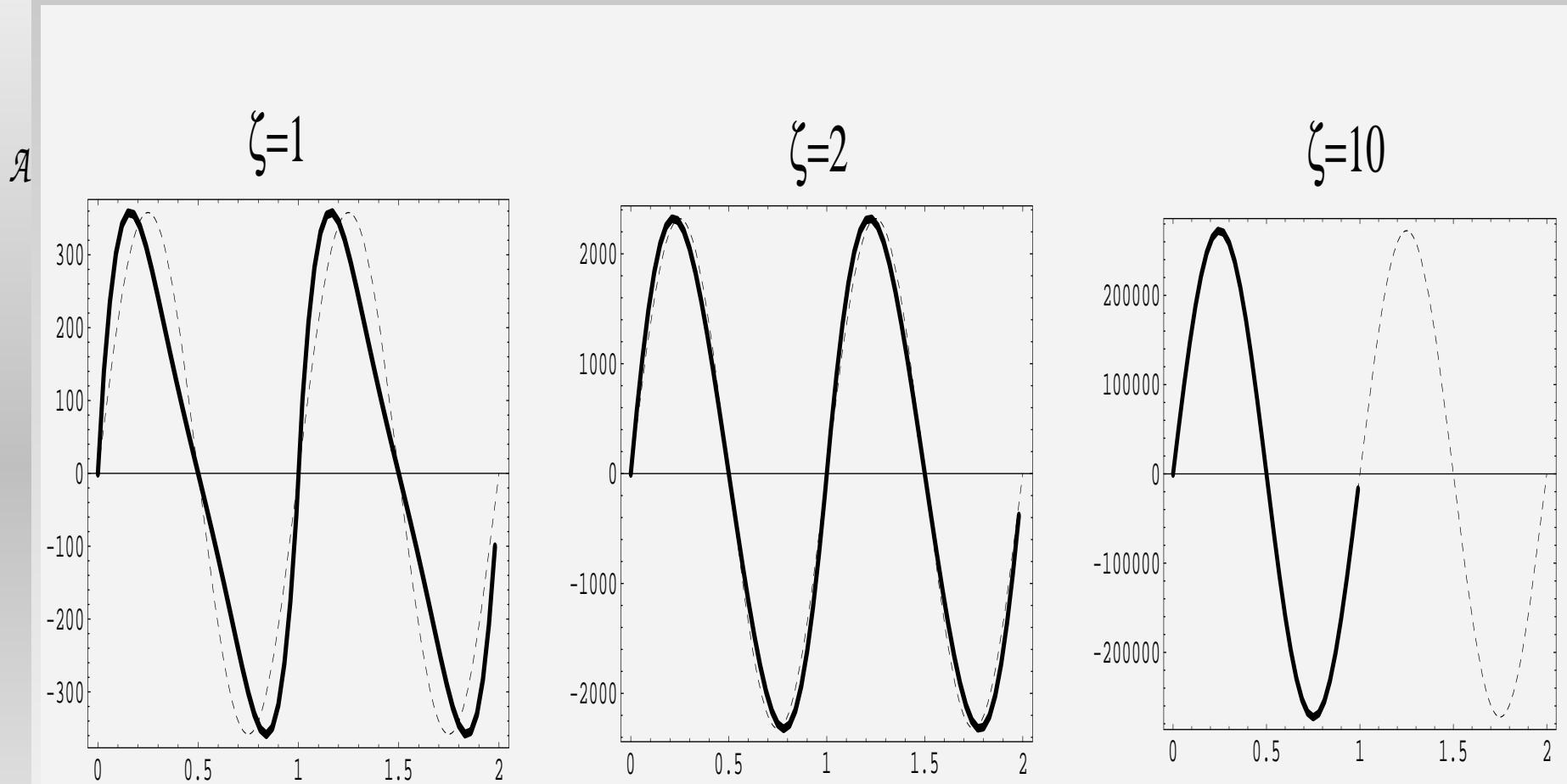
$$\hat{\phi}^a(\tau) \in \sum_{\text{HS (anti)caloron}} \text{tr} \int d^3x \int d\rho \frac{\lambda^a}{2} \times$$

$$F_{\mu\nu}((\tau, 0)) \{(\tau, 0), (\tau, \vec{x})\} \times$$

$$F_{\mu\nu}((\tau, \vec{x})) \{(\tau, \vec{x}), (\tau, 0)\} .$$



saturation:



\implies

$$-\mathcal{D} = \partial_\tau^2 + \left(\frac{2\pi}{\beta}\right)^2$$

$$-\partial_\tau\phi = \pm i \Lambda^3 \lambda_3 \phi^{-1}, \text{ (fixed global gauge)}$$

$$\text{where } \phi^{-1} \equiv \frac{\phi}{|\phi|^2}$$

$$\Rightarrow V(\phi) = \text{tr } \Lambda^6 \phi^{-2} \text{ by squaring RHS}$$

$$\Rightarrow |\phi| = \sqrt{\frac{\Lambda^3}{2\pi T}}$$

$$\Rightarrow \text{unique, coarse-grained action for } |Q| = 1 \text{ HS (anticalorons)}$$

$$\Rightarrow \phi's \text{ inertness}$$

What about $Q = 0$?

- ▶ perturbative renormalizability:
[’t Hooft, Veltman 1971-73]
⇒ coarse-graining yields
same form as fundamental action
- ▶ gauge invariance glues $Q = 0$ to $|Q| = 1 \Rightarrow$

$$S = \text{tr} \int_0^\beta d\tau \int d^3x \left(\frac{1}{2} G_{\mu\nu} G_{\mu\nu} + D_\mu \phi D_\mu \phi + \Lambda^6 \phi^{-2} \right)$$

- ▶ subject to offshellness constraints in
unitary-Coulomb gauge
(coarse-graining down to resolution $|\phi|$)

full ground state

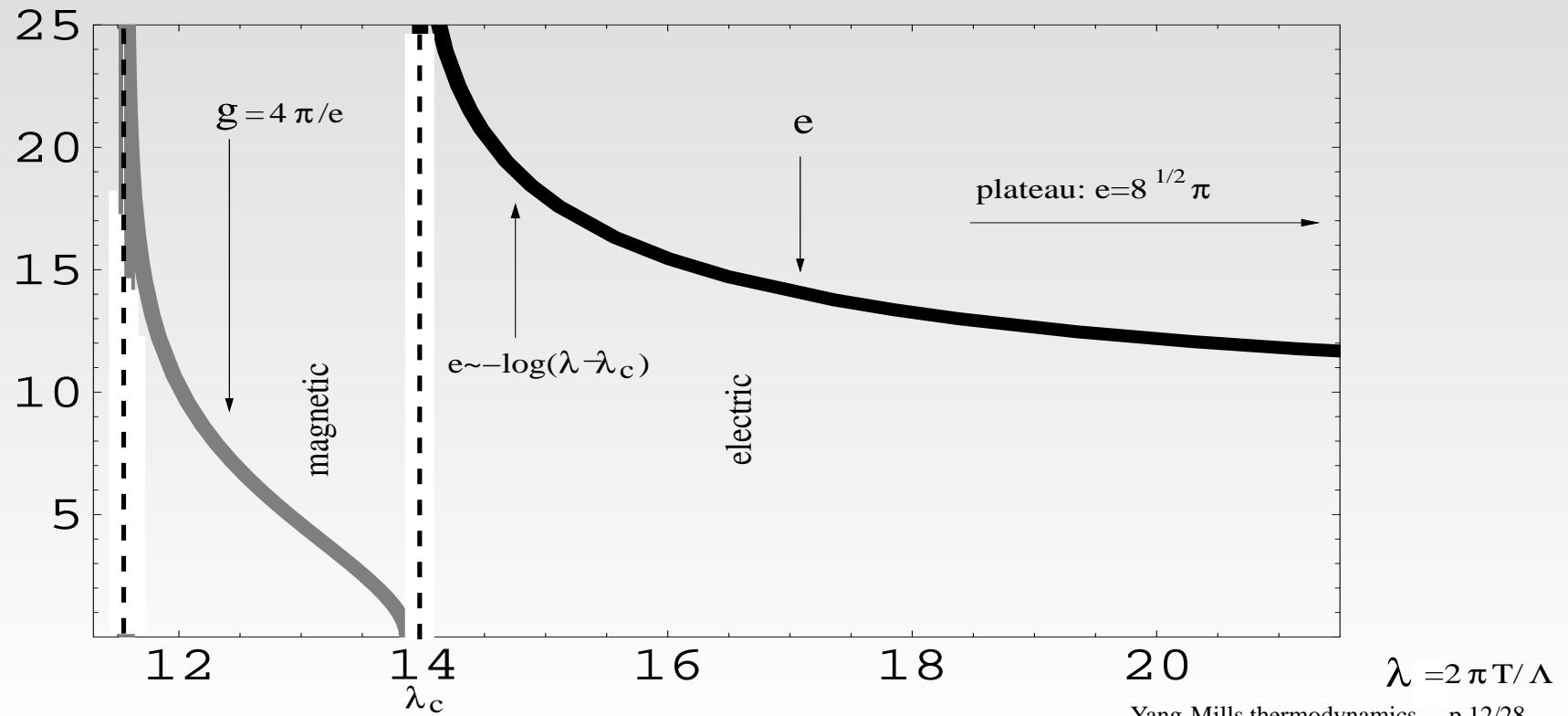
- ▶ from $D_\mu G_{\mu\nu} = 2ie[\phi, D_\nu \phi]$:
 - pure gauge $a_\mu^{bg} = \frac{\pi}{e} T \delta_{\mu 4} \lambda_3$
 \Rightarrow ground-state energy-density and pressure
$$\rho^{g.s} = 4\pi \Lambda^3 T = -P^{g.s} \neq 0$$
- ▶ rotation to unitary gauge $a_\mu^{bg} = 0$:
 - gauge transformation singular but admissible
(does not affect periodicity of fluct. δa_μ)
 - but: $\text{Pol}[a^{bg}] = -1 \xrightarrow{GT} \text{Pol}[a^{bg}] = +1$
 $\Rightarrow Z_2^{\text{el}}$ degeneracy
 \Rightarrow deconfinement

excitations and loop expansion

- adjoint Higgs mechanism:

2 out of 3 directions massive with $m = e \sqrt{\frac{\Lambda_E^3}{2\pi T}}$

- T evolution of eff. coupl. e :
requiring that P, ρ, \dots from partition function



► counting of d.o.f.:

fundamentally:

$$\begin{aligned} & 3 \text{ species (gluons)} \times 2 \text{ pols.} + \\ & 1 \text{ species (monop)} \times 2 \text{ charges} = 8 \end{aligned}$$

after coarse-graining:

$$\begin{aligned} & 2 \text{ species (gluons)} \times 3 \text{ pols.} + \\ & 1 \text{ species (gluon)} \times 2 \text{ pols.} = 8 \\ \Rightarrow & 8 \text{ (fund)} = 8 \text{ (coarse-grained).} \end{aligned}$$

same way for SU(3):

$$\Rightarrow 22 \text{ (fund)} = 22 \text{ (coarse-grained)}$$

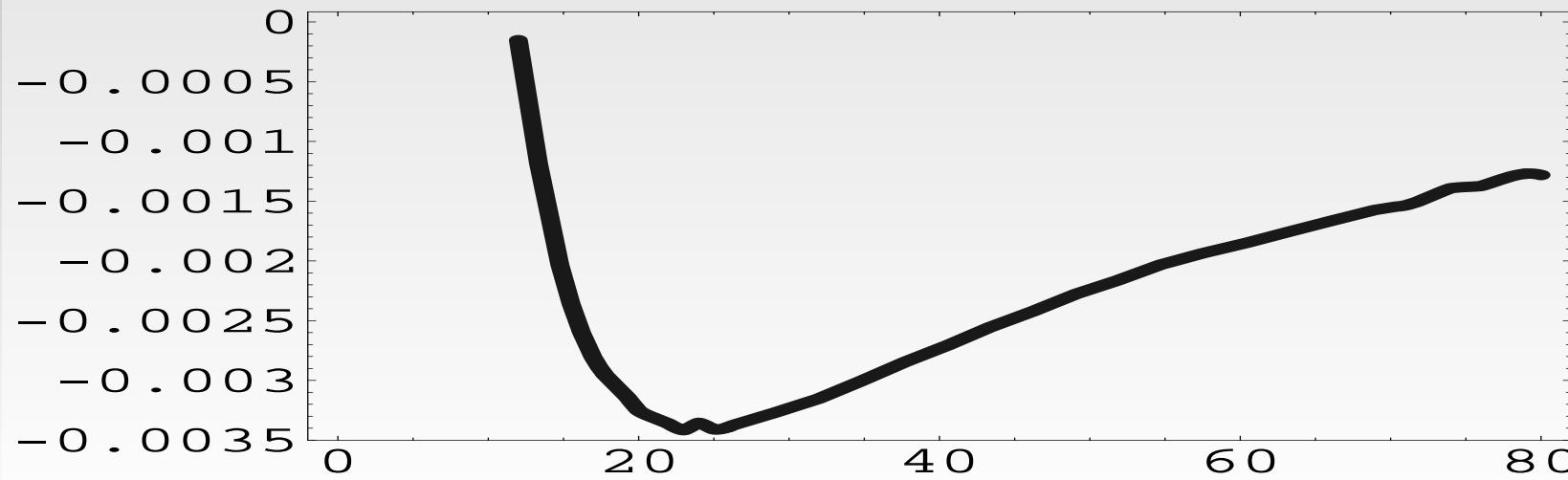
► loop expansion:

2-loop:

[Rohrer,Herbst,RH 2004; Schwarz,RH,Giacosa 2006]

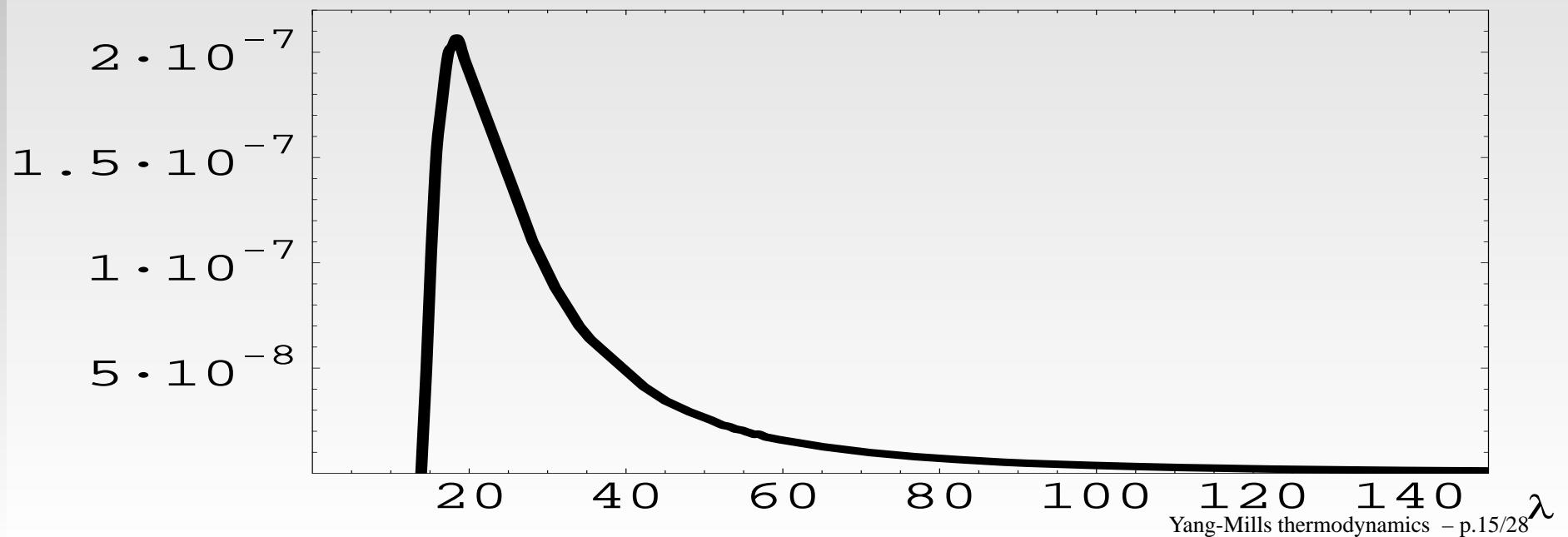
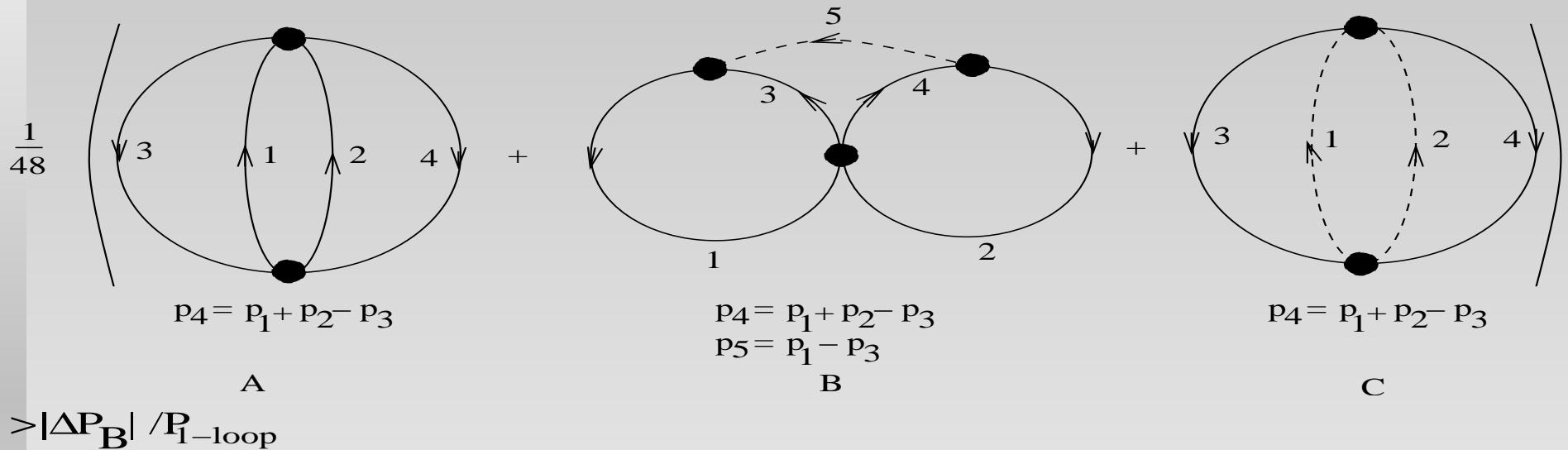
$$\Delta P = \frac{1}{4} \left(\text{Diagram A} \right) + \frac{1}{8} \left(\text{Diagram B} + \text{Diagram C} \right)$$

$(\Delta P_{ttv}^{\text{HJM}} + \Delta P_{ttc}^{\text{HJM}})/P_{\text{1-loop}}$



irreducible 3-loop:

[Kaviani,RH 2007]



arguments on loop expansion in general:

[RH 2006]

- resummation of 1PI diagrams
 \Rightarrow **no pinch singularities**
- irreducible diagrams **terminate** at finite loop order

(Euler characteristics for spherical polygon,
constraints on loop momenta in effective theory

\Rightarrow

number of constraints **exceeds**
number of independent radial loop variables
at **sufficiently large number of loops**)

preconfining phase

- ▶ condensation of monopoles:
 - phase of complex scalar = magnetic flux through $S_2^{R=\infty}$ of M-A pair at rest ($e \rightarrow \infty$)
 - modulus as in dec. phase
 - no change of form of action for free dual gauge modes by coarse-graining
⇒ **unique effective action**
 - Polyakov loop always unique
 - pressure exact at one loop
 - evol. of magnetic coupling g by requiring der. of pressure from fund. partition function

► counting of d.o.f.:

fundamentally:

$$1 \text{ species ('photon')} \times 2 \text{ pols.} + \\ 1 \text{ species (center-vortex loop)} = 3$$

after coarse-graining:

$$1 \text{ species (massive 'photon')} \times 3 \text{ pols.} = 3 \\ \Rightarrow 3 \text{ (fund)} = 3 \text{ (coarse-grained).}$$

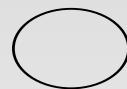
same way for SU(3):

$$\Rightarrow 6 \text{ (fund)} = 6 \text{ (coarse-grained)}$$

low temperature

- ▶ condensation of center-vortex loops (CVL's)
 - discrete values of phase of complex scalar field
 - = center flux through $S_1^{R=\infty}$ of spin-0 vortex pair at rest ($g \rightarrow \infty$)
 - spectrum: single and selfintersecting CVL's

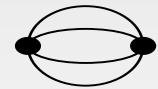
$n=0:$
 $(N_0=1)$



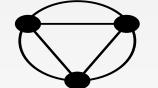
$n=1:$
 $(N_1=1)$



$n=2:$
 $(N_2=2)$



$n=3:$
 $(N_3=4)$

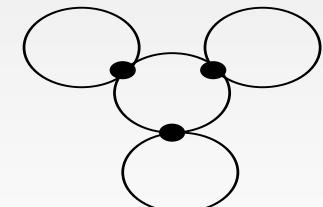
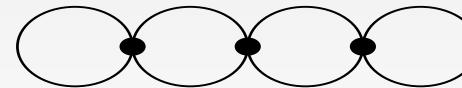
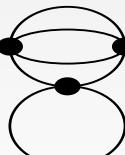
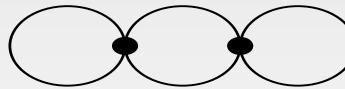


multiplicity of mass-n soliton:
number of connected bubble diagrams
at order n in $\lambda\phi^4$

large-order counting:
anharmonic oscillator
[Bender, Wu 1969–76]

→ **over-exponentially in energy
growing density of states**

→ **Hagedorn transition to
preconfining phase**



► counting of d.o.f.:

fundamentally:

1 species ('very massive photon') \times 3 pols.=3

after coarse-graining:

1 species (massless CVL)+

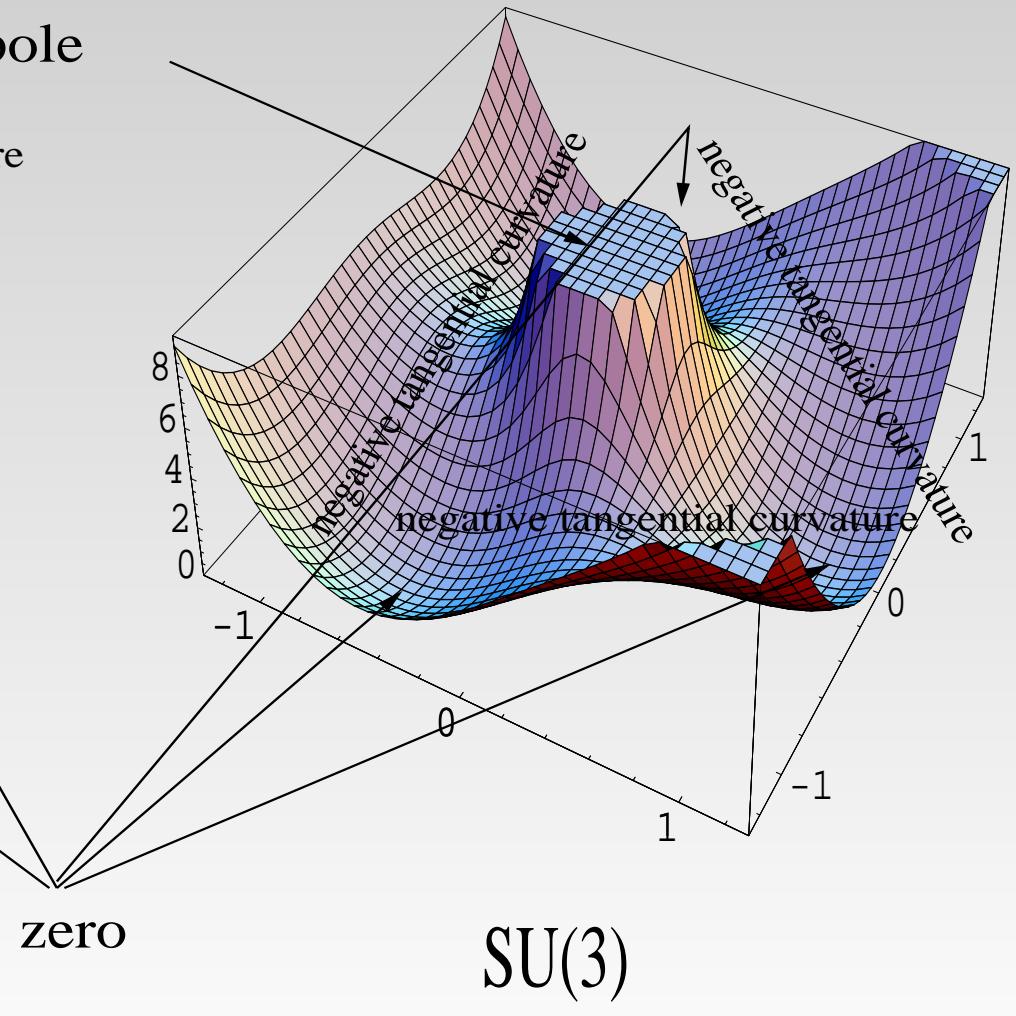
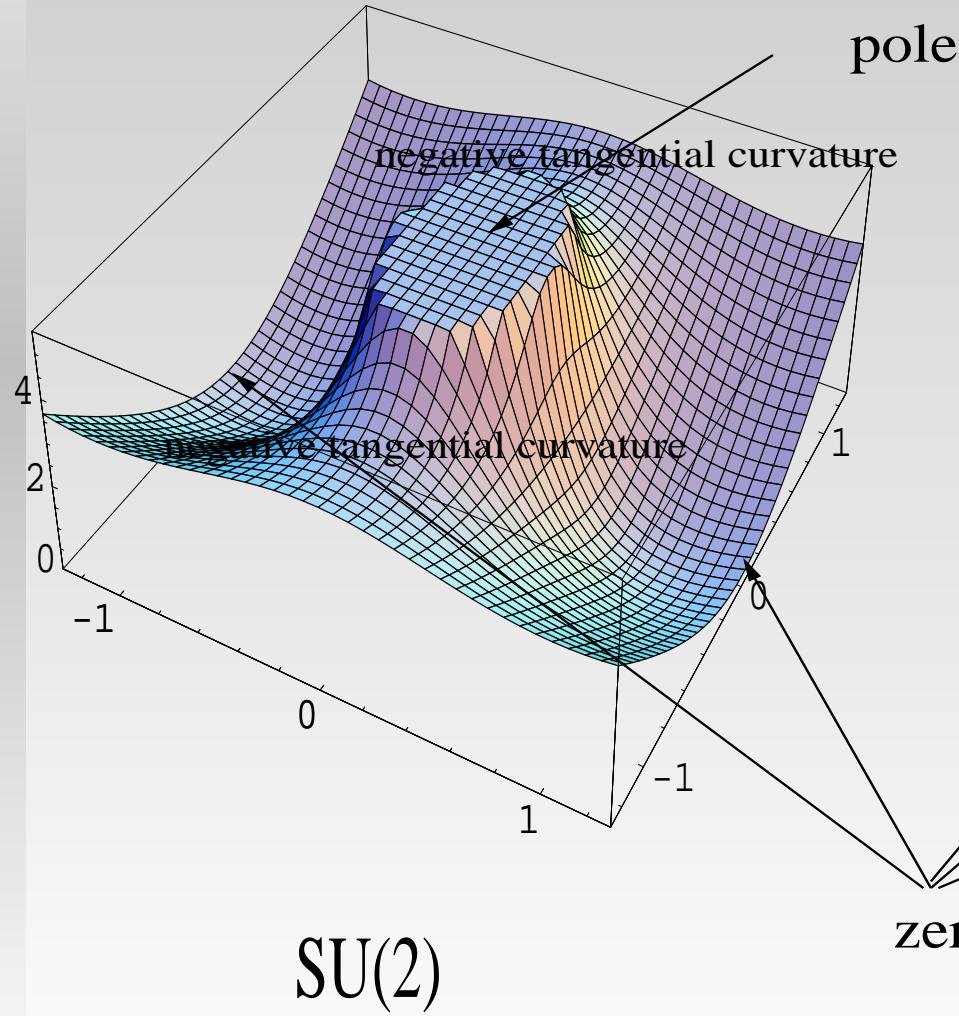
1 species (massive CVL) \times 2 charges=3

\Rightarrow 3 (fund)=3 (coarse-grained).

same way for SU(3):

\Rightarrow 6 (fund)=6 (coarse-grained)

- potential unique up to inessential,
 $U(1)$ invariant rescaling



$SU(3)$

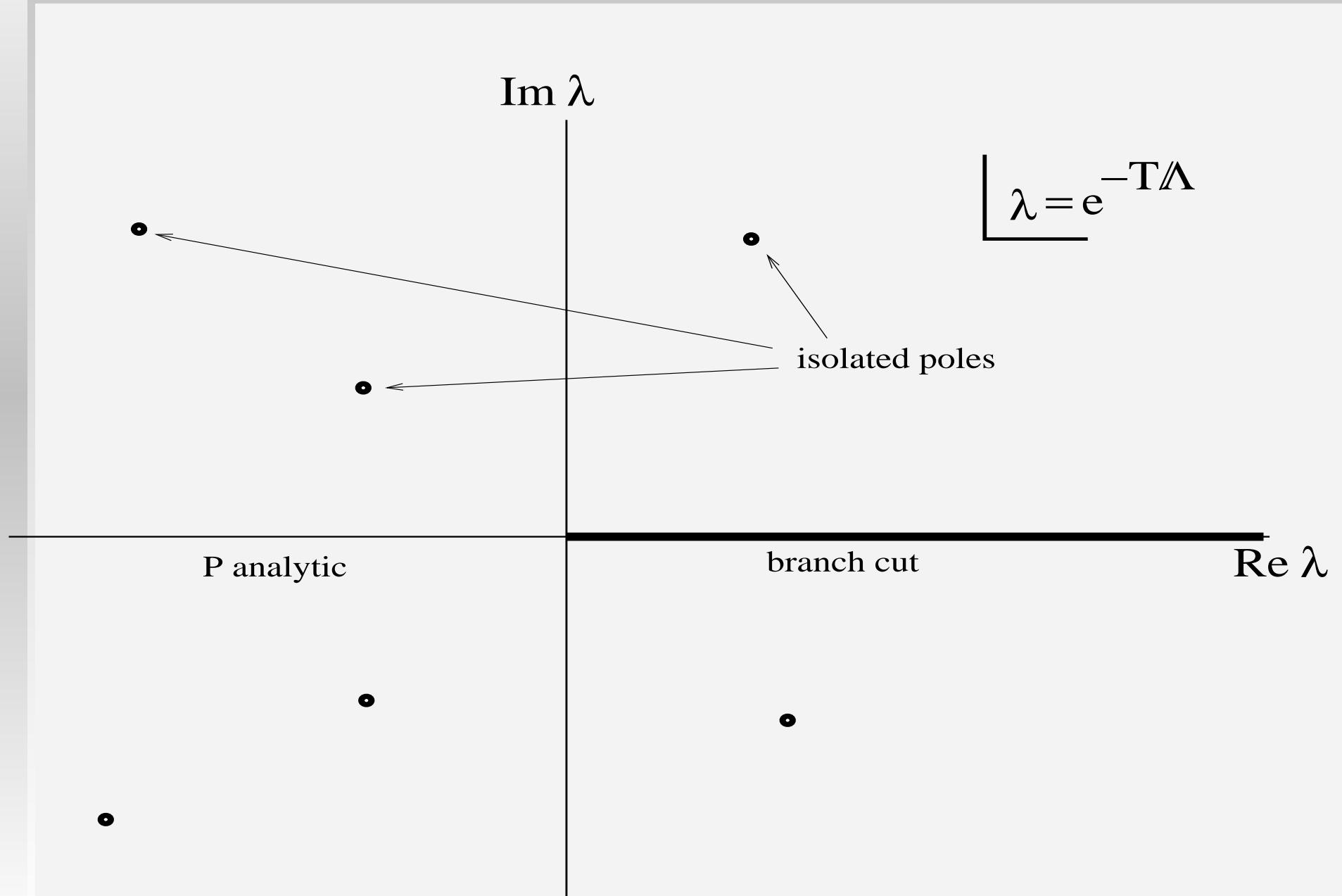
- **asymptotic-series** representation of pressure:

$$P_{\text{as}} = \frac{\Lambda^4}{2\pi^2} \hat{\beta}^{-4} \times \\ \left(\frac{7\pi^4}{180} + \sqrt{2\pi} \hat{\beta}^{\frac{3}{2}} \sum_{l=0}^L a_l \sum_{n \geq 1} (32\lambda)^n n! n^{\frac{3}{2}+l} \right) ,$$

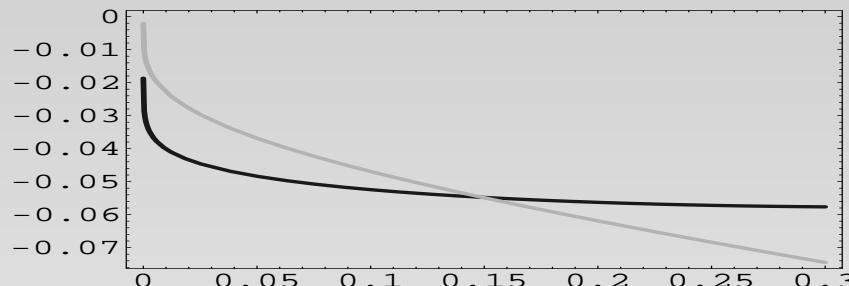
where $\hat{\beta} \equiv \Lambda/T$ and $\lambda \equiv \exp[-\hat{\beta}]$.

- Borel transformation and analytic continuation
 \Rightarrow **analytic** dependence on Borel parameter t
 (polylogs) for $\lambda < 0$
- inverse Borel trafo:
 \Rightarrow analytic dependence for $\lambda < 0$ and
meromorphic in entire λ -plane except for $\lambda \geq 0$

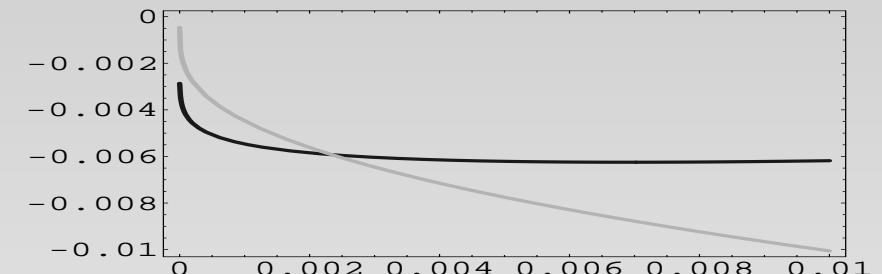
analyticity structure of physical pressure P :



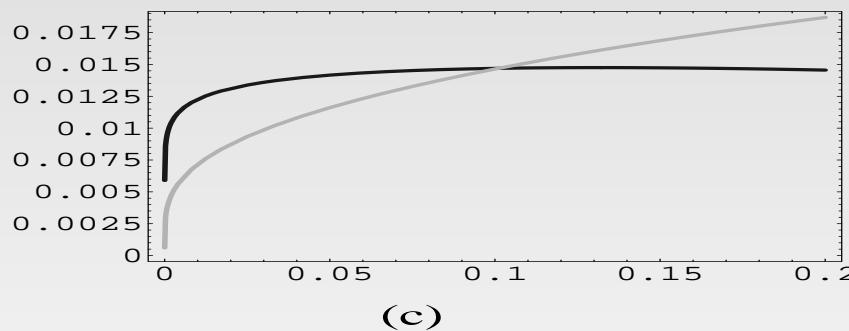
- $\text{Re } P$ continuous across cut:
- sign-ambiguous $\text{Im } P$ grows slower than $\text{Re } P$
- turbulences become relevant for sufficiently high T only



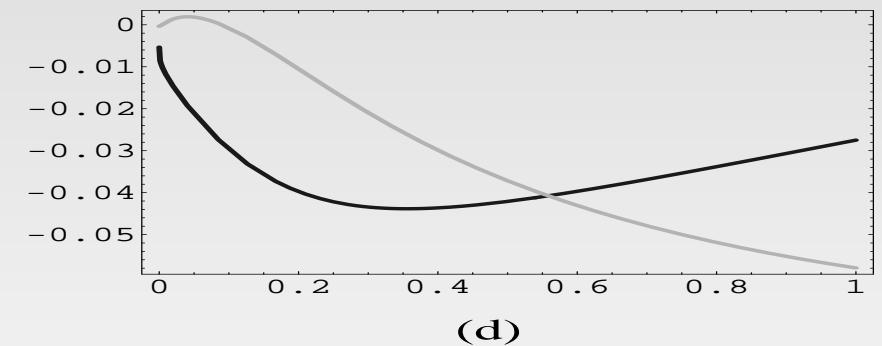
(a)



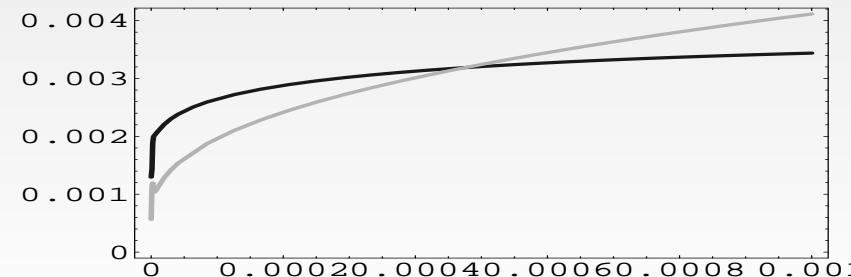
(b)



(c)



(d)



(e)

summary and conclusions

► deconf. phase:

- magnetic-magnetic correlations in (anti)calorons generate adj. Higgs field
- rapid saturation of average
- negative ground-state pressure by microscopic holonomy shifts (annihilating M-A pairs)
- thermal quasiparticles on tree level (adj. Higgs mech.)
- very small radiative corrections, termination of expansion in terms of irreducible loops

► **preconf. phase:**

- averaged magnetic flux of M-A pair through S_2^∞ generates phase of complex scalar
- dual gauge field Meissner massive
- loop expansion trivial

► **conf. phase:**

- averaged center flux of CVL pair through min. surface spanned by S_1^∞ generates discrete values of phase of complex scalar
- excitations are single or selfintersecting CVL's of factorially growing multiplicity
- asymptotic-series representation of pressure
- Borel summability for complex values of T
- analytic continuation: rapidly (slowly) rising modulus of real (imaginary) part
- interpretation: growing relevance of turbulences with increasing T

- ▶ physics applications:
 - CMB
 - late-time cosmology (axion + SU(2))
 - electroweak symmetry breaking

Thank you.