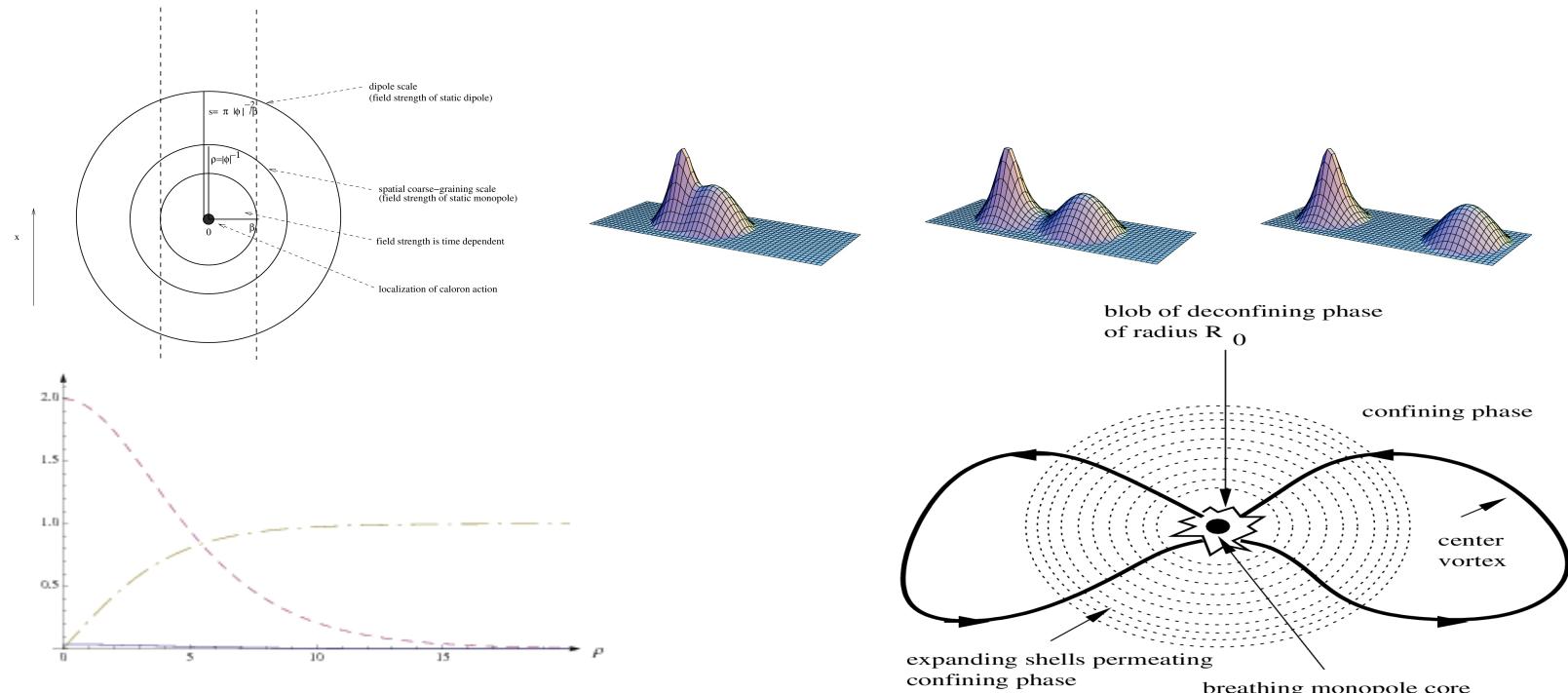




The isolated electron: De Broglie's „hidden“ Thermodynamics, SU(2) Quantum Yang-Mills theory, and a strongly perturbed BPS Monopole

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outline

- Louis de Broglie's arguments on existence of „thermodynamics of the isolated electron“

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 - deconfining, preconfining, and confining phases
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 - value of coupling e , caloron action
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- electron as self-intersecting, spatial center-vortex loop: size estimates

Louis de Broglie: reminder of derivation of wavelength λ , I

- consider particle at rest as a ‘clock’ with frequency ν_0 which extends through a large portion of space, that is, amplitude Ψ of oscillation is the same everywhere:

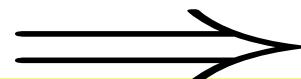
$$\Psi = a_0 \sin(2\pi\nu_0 t_0)$$

- by the rules of QM: rest energy equals quantum of action times ν_0 :

$$m_0 c^2 = h\nu_0$$

- consider Lorentz-trafo in \mathcal{Z} direction:

$$t_0 = \frac{t - \frac{\beta}{c} z}{\sqrt{1 - \beta^2}} \quad (\beta \equiv \frac{v}{c})$$



spatially “flat” oscillation becomes propagating wave with

$$\nu = \frac{\nu_0}{\sqrt{1 - \beta^2}}, \quad v_p = \frac{c}{\beta}, \quad \lambda = \frac{v_p}{\nu} = \frac{c\sqrt{1 - \beta^2}}{\nu_0 \beta}$$

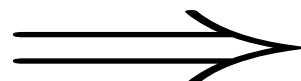
Louis de Broglie: reminder of derivation of wavelength λ , II

- on the other hand:

$$\frac{p}{h} = \frac{m_0 v}{h \sqrt{1 - \beta^2}}$$

or upon substitution of $m_0 = \frac{h\nu_0}{c^2}$

$$\frac{p}{h} = \frac{\nu_0 \beta}{c \sqrt{1 - \beta^2}} = \lambda^{-1}$$



existence of 'clock' frequency ν_0 essential in derivation of
de Broglie wavelength λ !

Thermodynamics of isolated electron: Louis de Broglie

- from demand that entropy

$$S = \int \frac{dQ}{T}$$

is a fundamental invariant under the Poincare group and from

$$dQ = dQ_0 \sqrt{1 - \beta^2}$$

(considering an adiabatic variation of rest mass due to input of heat)

\Rightarrow

$$T = T_0 \sqrt{1 - \beta^2}$$

- increase of rest energy $Q_0 = m_0 c^2$ under boost can be decomposed into a **decrease** of internal heat to

and work A invested to boost the particle

$$\frac{Q_0}{\sqrt{1 - \beta^2}} = Q + A$$

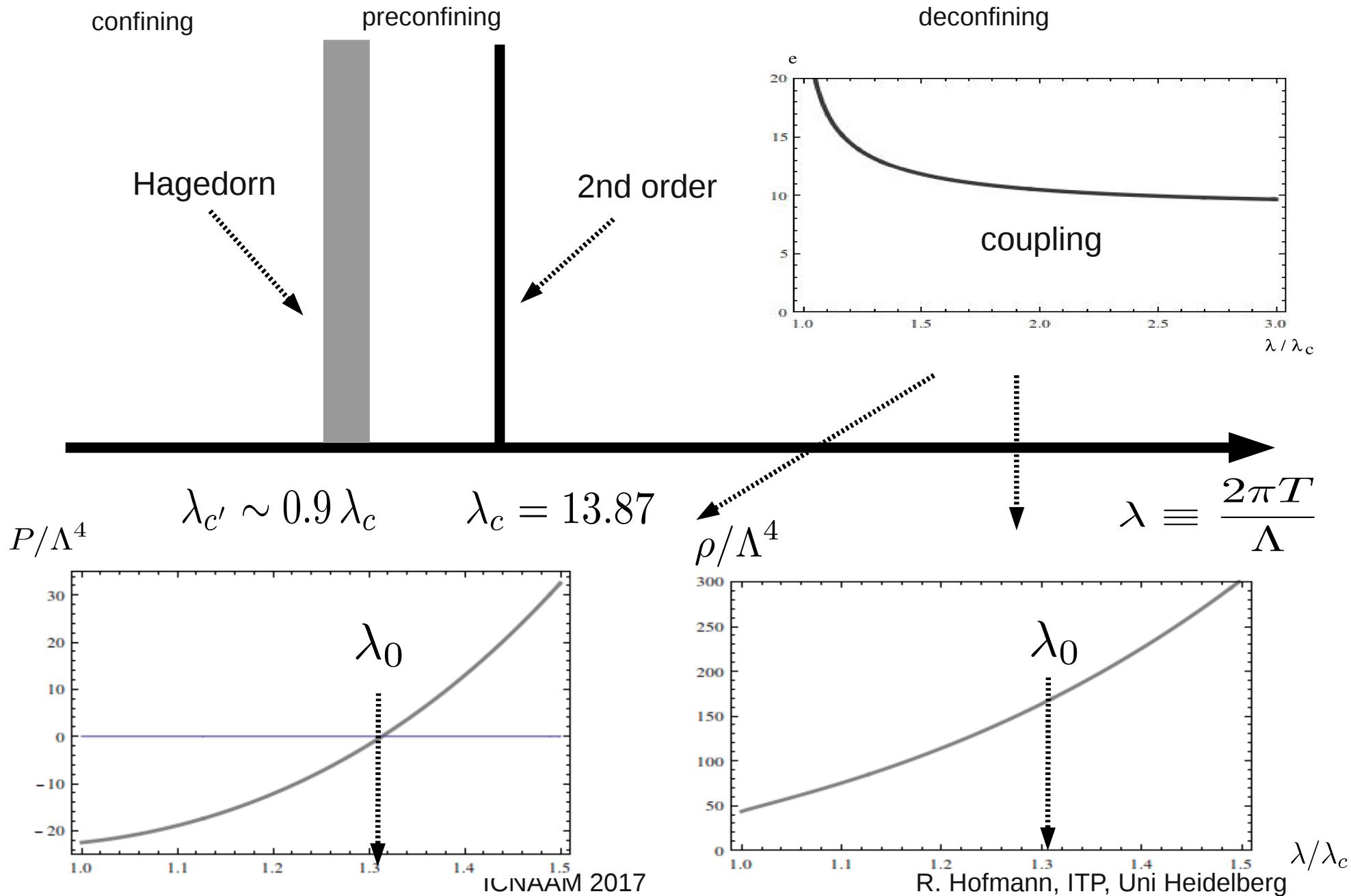
$$Q = \sqrt{1 - \beta^2} Q_0$$

$$A = \frac{m_0 v^2}{\sqrt{1 - \beta^2}}$$

Particle with internal heat
can not be point-like.

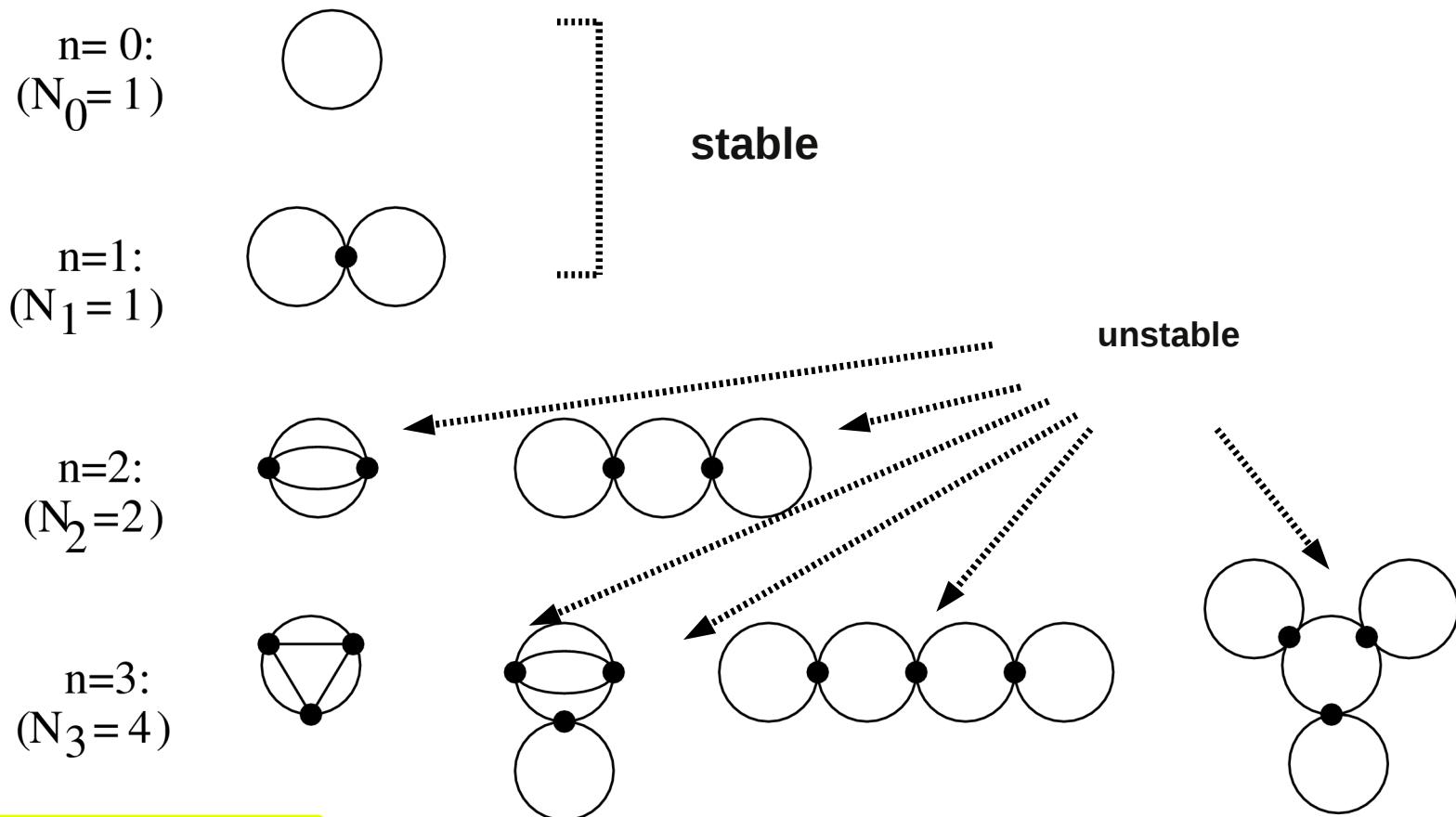
How does $SU(2)$
Yang-Mills thermodynamics
come in?

Phase diagram of SU(2) Quantum Yang-Mills thermodynamics



Phase diagram of SU(2) Quantum Yang-Mills thermodynamics: Confining phase

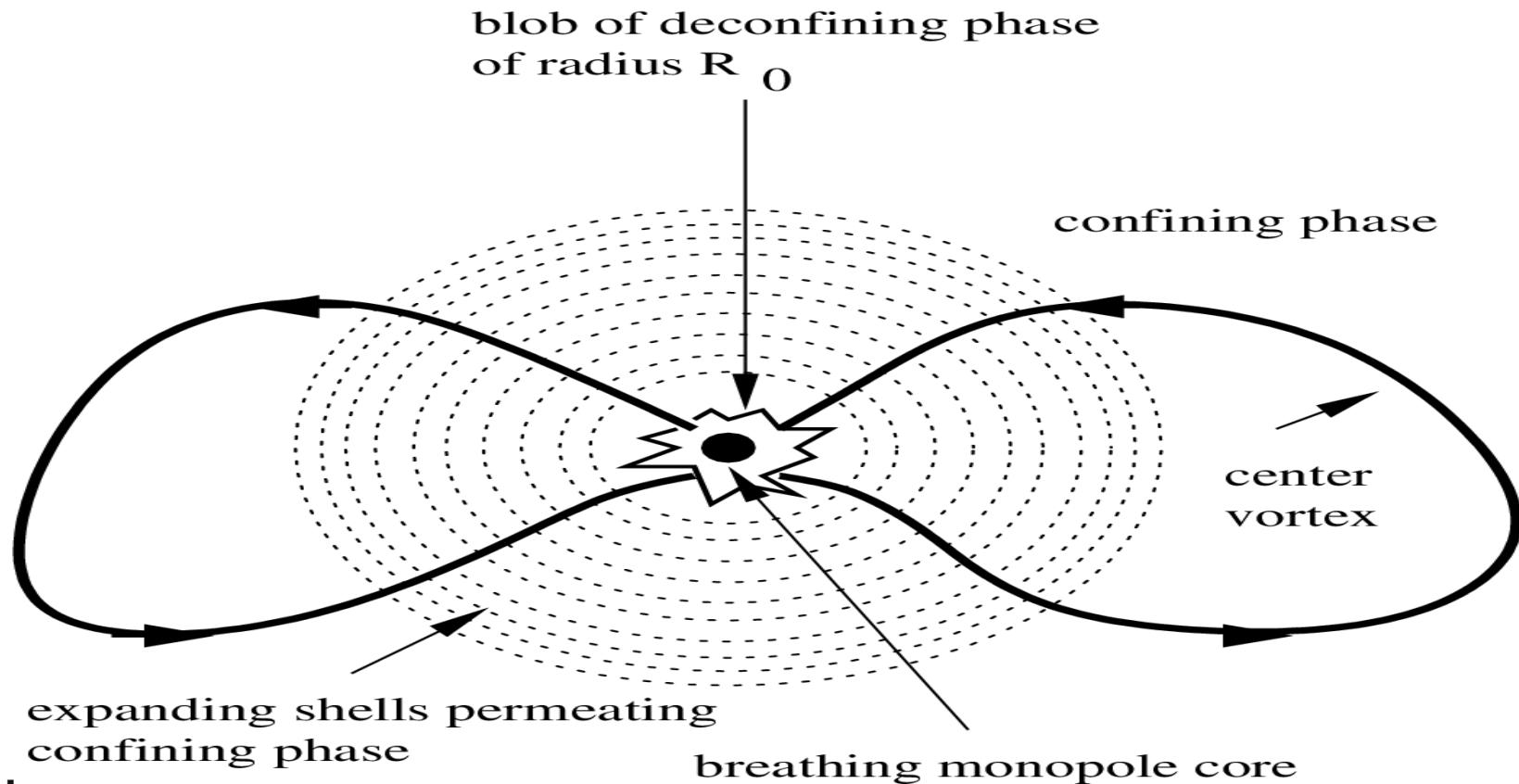
excitations: n-fold self-intersecting center-vortex loops



multiplicity $> n!$

[Bender & Wu, 1969]

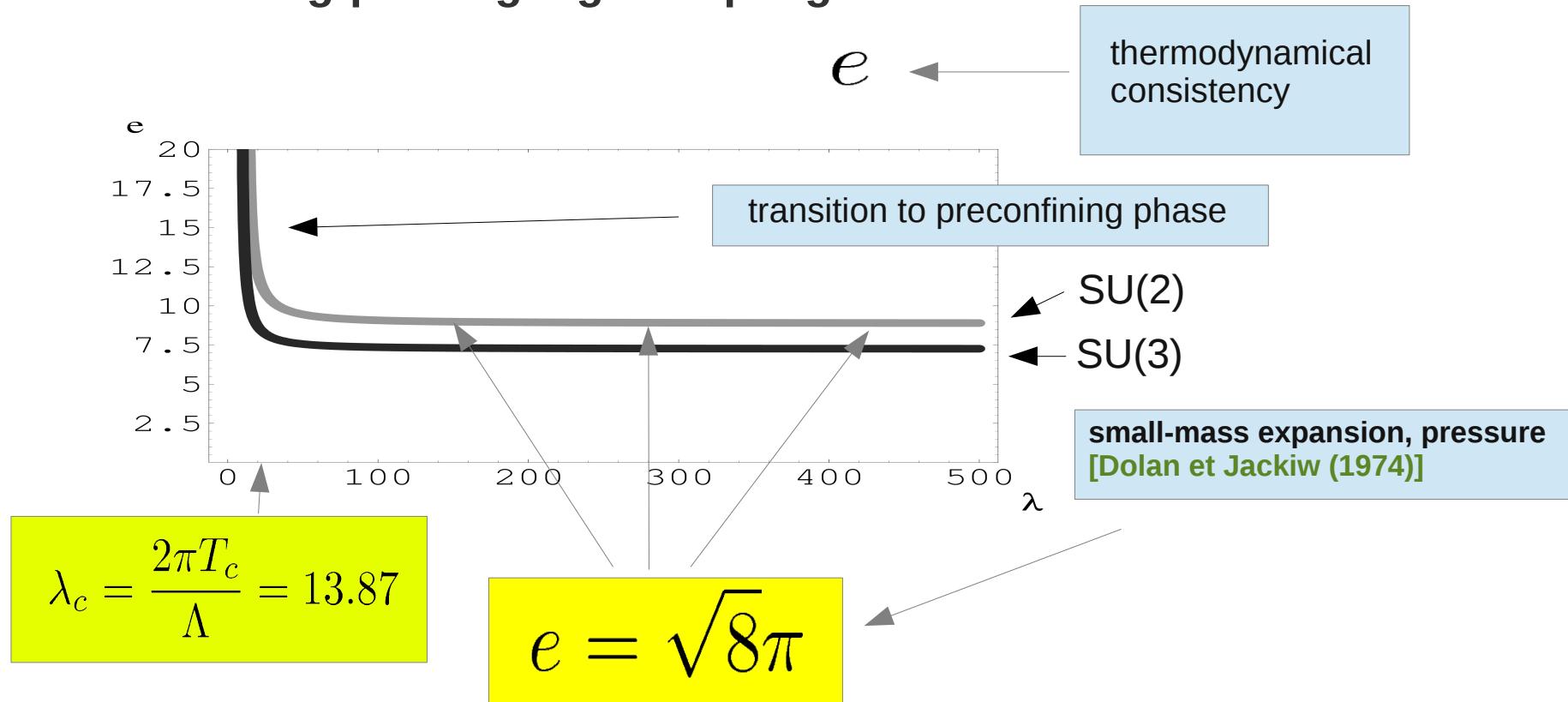
Proposal for electron's anatomy



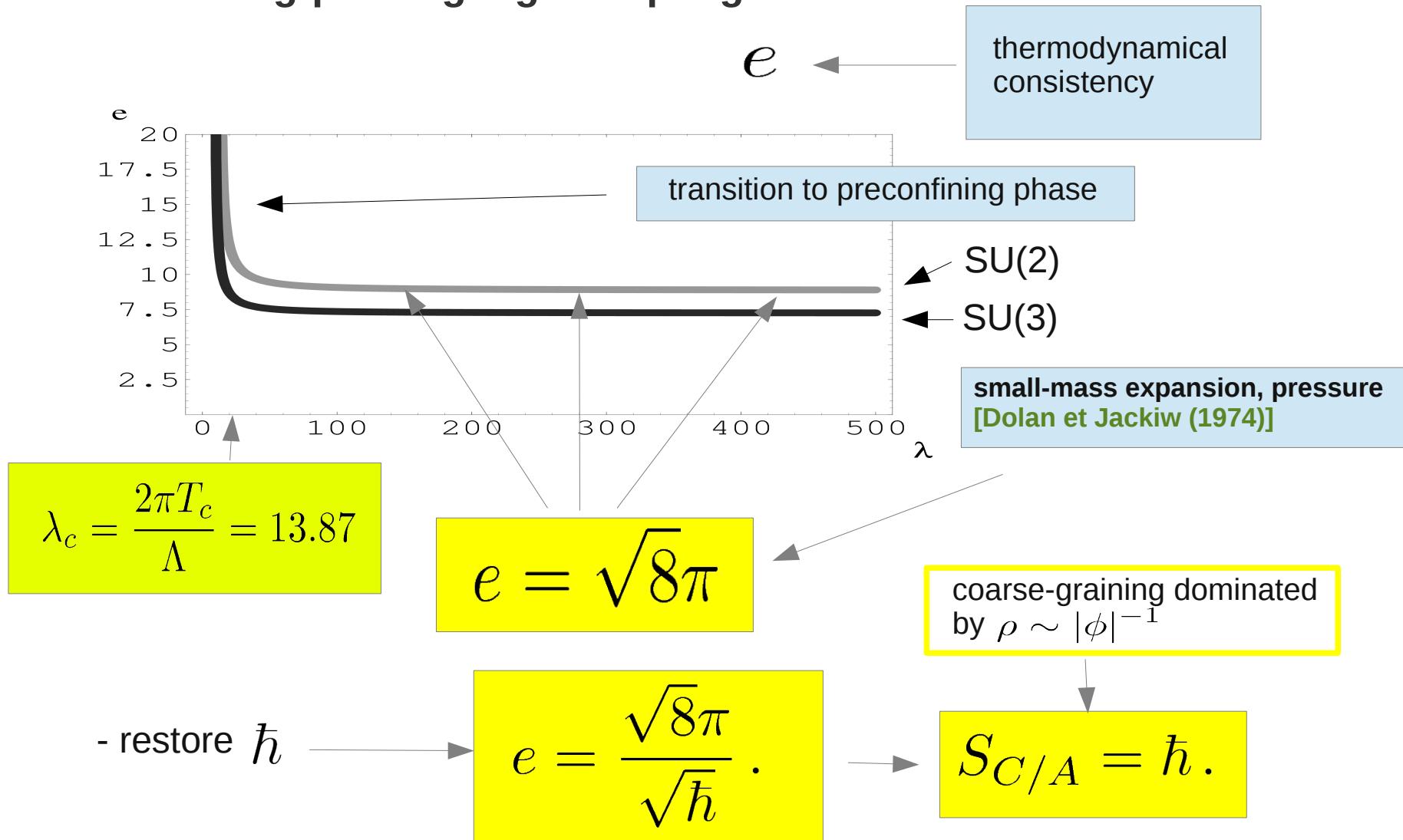
based on:

- (i) electric-magnetically dual interpretation of U(1) charges in SU(2)
- (ii) classical dynamics of strongly perturbed BPS monopole
[Fodor & Racz, 2004]
- (iii) structure of deconfining thermal ground state, small-holonomy (anti)calorons,
dissociation of large-holonomy (anti)calorons **[Herbst & RH, 2004; RH, TQYM, 2016]**

(i) electric-magnetically dual interpretation: deconfining-phase gauge coupling and caloron action



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[Brodsky et al. (2011);
Kaviani & RH 2012;
RH 2012,2013]

(i) electric-magnetically dual interpretation: argument involving QED fine-structure constant

if SU(2) something to do with photons [RH (2005), Grandou & RH (2015), etc]
then **electric-magnetically dual** interpretation required:
in units $c = \epsilon_0 = \mu_0 = k_B = 1$ fine-structure constant

$$\alpha = \frac{Q^2}{4\pi\hbar},$$

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$$(e = \frac{\sqrt{8\pi}}{\sqrt{\hbar}} \cdot)$$

$$Q \propto \frac{1}{e}.$$

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But: magnetic coupling
in SU(2)

$$Q \propto \frac{1}{e}.$$

$$g = \frac{4\pi}{e}.$$

→ SU(2) to be interpreted in an **electric-magnetically dual way**.
(e.g., magnetic monopole \longleftrightarrow electric monopole, etc.)

(ii) classical dynamics of strongly perturbed BPS monopole

- hyperboloidal conformal transformation of the original field equations for the profiles

$$H(r, t) = h(r, t)/r + H_\infty$$

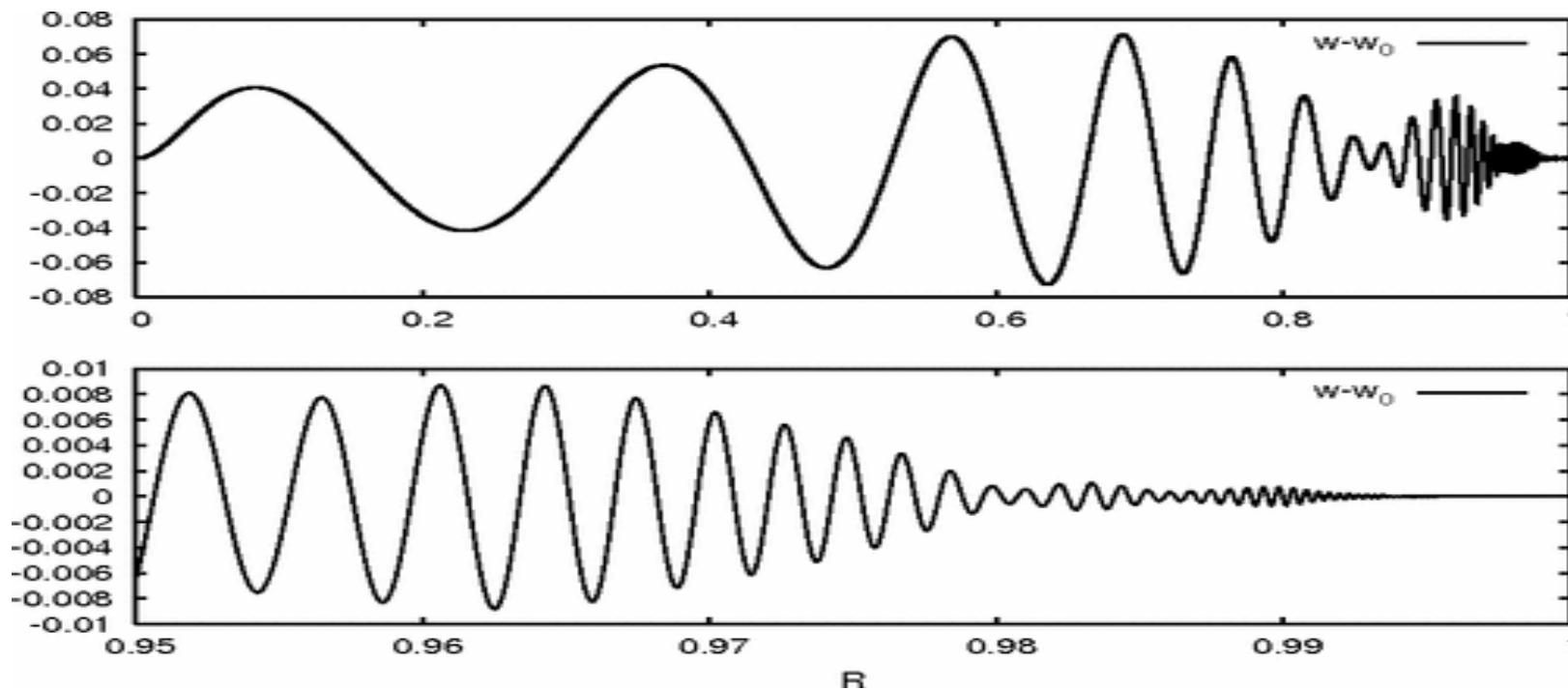
and

$$w(r, t)$$

of the adjoint Higgs and off-Cartan gauge fields fields, respectively

[Fodor & Racz, 2004]

- under a spherically symmetric initial pulse in energy density:
 - high-frequency oscillations in w forming expanding shells



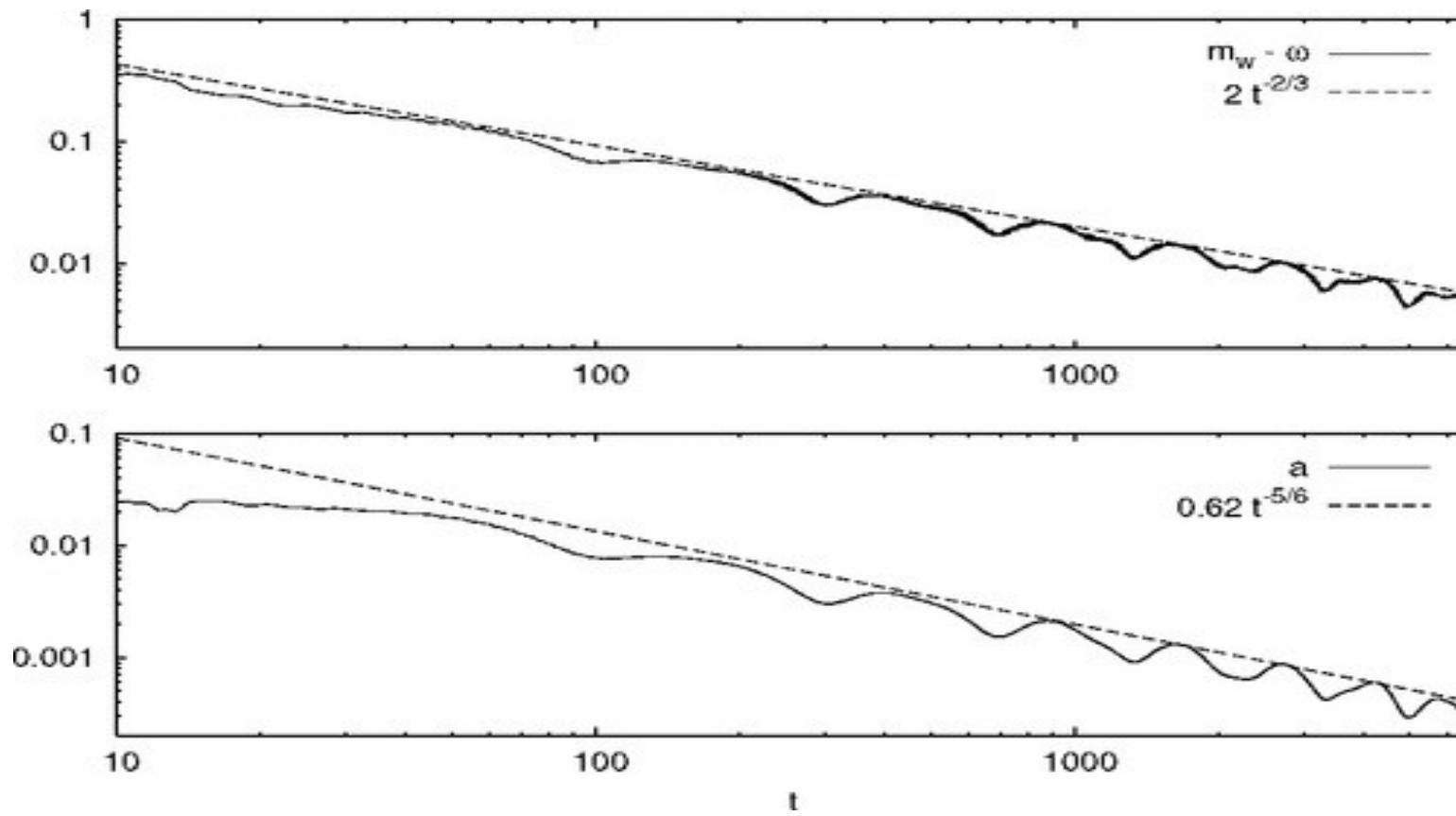
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- a localized **breathing state** appears in association with the energy density within the monopole core whose frequency ω_0 approaches the mass

$$m_w = eH_\infty$$

of the two off-Cartan modes in a power-like way in time

[Fodor & Racz, 2004]



(ii) classical dynamics of strongly perturbed BPS monopole

There is a generic circular frequency

$$\omega_0 = 2\pi\nu_0 = e H_\infty$$

associated with core breathing of the (initially perturbed) BPS monopole, that is, a **clock in de Broglie's sense**.

(iii) structure of deconfining thermal ground state

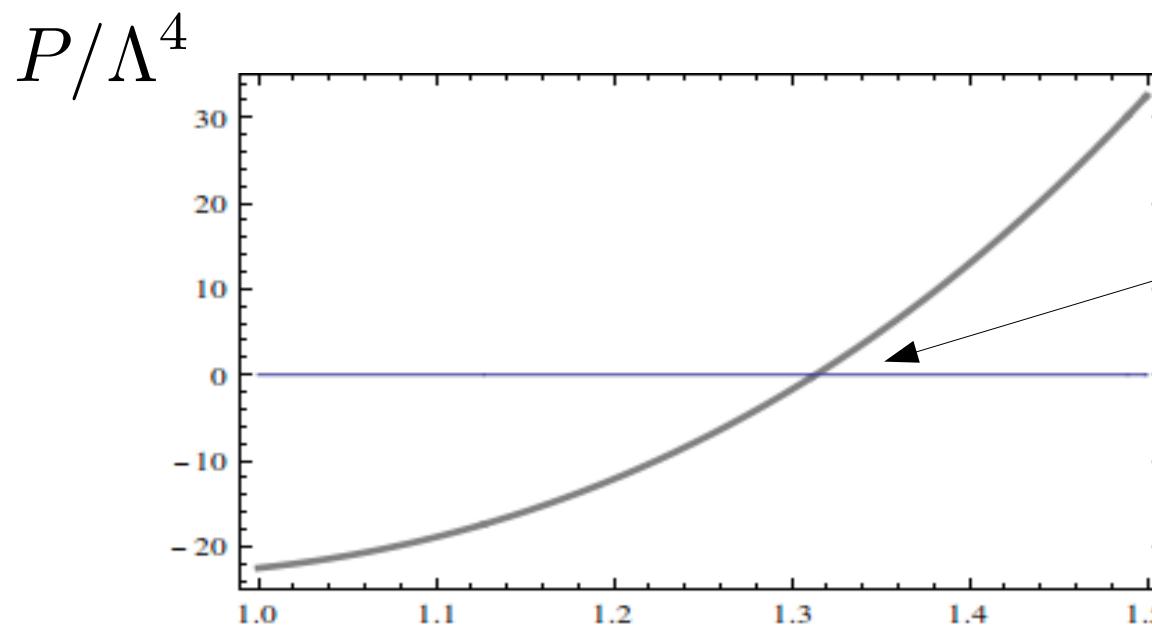
- Only (anti)calorons of scale $\rho \sim |\phi|^{-1}$ contribute to emergent a priori estimate of thermal ground state \Rightarrow coupling e relevant in describing caloron dissociation and ensuing monopole- antimonopole

[Herbst & RH 2004; RH 2016]

- If monopole liberated by dissociation of a large-holonomy caloron in deconfining phase at temperature T_0 then

$$H_\infty(T_0) = \pi T_0$$

[Diakonov et al 2004]



$$T_0 = \lambda_0 \frac{\Lambda}{2\pi}$$

$$\lambda_0/\lambda_c = 1.32$$

$$\lambda/\lambda_c$$

(iii) structure of deconfining thermal ground state

We have:

[Diakonov et al 2004]

$$e(1.32) = 12.96 \quad (\text{A})$$

$$\rho/\rho^{\text{gs}}(1.32) = 8.31 \quad (\text{B})$$

with:

$$\rho^{\text{gs}}(T_0) = 4\pi\Lambda^3 T_0 = 32 \frac{H_\infty^4(T_0)}{\lambda_0^3}$$

We assume static BPS monopole of mass

$$m_{\text{mon}} = \frac{8\pi^2}{e^2(T_0/T_c)} H_\infty(T_0) = \frac{8\pi^2}{12.96^2} H_\infty(T_0)$$

to be immersed into blob of deconfining phase of radius R_0 at T_0 .

Size estimate for self-intersection region:

Thus, from (A) and (B) we have:

$$\begin{aligned} m_0 &= e H_\infty(T_0) = 12.96 H_\infty(T_0) = m_{\text{mon}} + E_0 \\ &= \frac{8\pi^2}{12.96^2} H_\infty(T_0) + \frac{4}{3}\pi R_0^3 \rho(T_0) \\ &= H_\infty(T_0) \left(\frac{8\pi^2}{12.96^2} + 8.31 \times \frac{128\pi}{3} \left(\frac{R_0}{18.31} \right)^3 H_\infty^3(T_0) \right) \end{aligned}$$

$$\iff R_0 = 4.10 H_\infty^{-1}(T_0)$$

Comparing this to the Compton wave length $\lambda_C = m_0^{-1} = 2.426 \times 10^{-12}$ m,

$$\Rightarrow R_0 \sim 53.14 \lambda_C \sim 1.29 \times 10^{-10} \text{ m}$$

[RH, 2017]

(~ 4.12 core size of monopole)

hierarchy between Yang-Mills scale ($\sim T_c$) and electron mass m_0 and Coulomb self-energy:

from $H_\infty(T_0) \Rightarrow$

$$\Lambda = \frac{1}{118.6} m_0$$

\Rightarrow

$$T_c = 13.87 / (2\pi \times 118.6) m_0 = 0.019 m_0 = 9.49 \text{ eV}$$

For comparison: average electron temperature at ITER envisaged to be

$$\bar{T}_e = 8.8 \text{ eV !!}$$

Coulomb self-energy correction:

$$\Delta m_0 = \int_{R_0}^{\infty} dr \frac{r^2}{r^4} = R_0^{-1} = \frac{1}{4.10} H_\infty(T_0) \ll 12.96 H_\infty = m_0 \quad \text{or} \quad m_0 \sim 53.14 \Delta m_0 .$$

Summary and outlook

- brief review of Louis de Broglie's argument on electron's inherent thermodynamics
- phase structure of SU(2) YM TD, deconfining phase, confining phase
- electron as a one-fold self-intersecting center-vortex loop
- strongly perturbed BPS monopole
- structure of deconfining thermal ground state
- size estimates
- Coulomb self-energy and relation between Yang-Mills scale and crit. Temp.
- open questions:
 - spherically symmetric perturbation of BPS monopole generalizable to other perturbations
 - spatial extent of loops (depending on environmentally provided resolution?)
 - Schrödinger wave mechanics facilitated by small core sizes vs. R_0
 - What happens under acceleration? (radiation, radiation reaction,...)