

# Issues with RAI

- usually no unique choice of physical scale

→ competing effects/multiscale physics  
not captured well

→ if there are dimensional classical quantities,  
then different identifications are possible

e.g.  $k \sim 1/r$   
or

$$k \sim \frac{1}{r} \left( \frac{G_N M}{r} \right)^n$$

only need  
 $k \nearrow \leftrightarrow r \searrow$

$k \searrow \leftrightarrow r \nearrow$

- where to implement RGI?

- action?

- equations of motion?

- solution to eqs. of motion?

↑  
more  
difficult

- in gravity: — backreaction can become important  
interplay of classical singular metric  
and quantum fluctuations that are  
supposed to regularise

many attempts  
in the literature  
to fix these  
problems

- coordinate invariance  
scale identification incompatible  
with diffeomorphism invariance

→ we discuss some ideas

# ① constraints from Bianchi identity

it's a consistency condition  
↑

recall Bianchi identity:  $D^\mu G_{\mu\nu} = 0$

↑  
Einstein tensor  
 $R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu}$

geometric  
identity, has  
nothing to do  
with dynamics

→ perform RGI at the level of  $\begin{cases} \text{action} & d=1 \\ \text{EOM} & d=0 \end{cases}$

$$G_{\mu\nu} = 8\pi G_N(k) T_{\mu\nu} - \Lambda(k) g_{\mu\nu} + \alpha \Delta t_{\mu\nu}$$

$$\Delta t_{\mu\nu} = G_N(k) \left[ D_\mu D_\nu - g_{\mu\nu} D^2 \right] G_N(k)^{-1}$$

assume  $D^\mu T_{\mu\nu} = 0$  and impose Bianchi identity:

a)  $d=1$ : compute  $D^\mu G_{\mu\nu} = 0$  and insert EoM (solve for  $T_{\mu\nu}$ )

$$G_N'(k) R = 2 (G_N'(k) \Lambda(k) - \Lambda'(k) G_N(k))$$

→ specific RH trajectory fixes  $k = k(R)$

$$\text{near FP: } G_N(k) \sim g_* k^2, \quad \Lambda(k) \sim \lambda_* k^2 \\ \Rightarrow k^2 \propto R / 4 \lambda_*$$

b)  $d=0$ : compute  $D^\mu G_{\mu\nu} = 0$  and insert EoM (solve for  $R_{\mu\nu}$ )

$$[8\pi G_N'(k) T_{\mu\nu} - \Lambda'(k) g_{\mu\nu}] D^\nu k(x) = 0$$

perfect fluid:  $\rho = w \mathcal{S}$ ,  $T_\mu{}^\nu = \text{diag}(-\mathcal{S}, p, p, p)$

$$\text{near FP, } k^4 \propto \frac{8\pi g_* \mathcal{S}}{\lambda_*}$$

## ② Self-consistency and iterative RGI

Platzman 1903.10411

- observation: we use  $g_{\mu\nu}^{cl}$  to identify  $k(x)$  and obtain  $g_{\mu\nu}^?$   
→ should use  $g_{\mu\nu}^?$  to identify scale self-consistently

↳ iterative RGI:

$$g_{\mu\nu}^{cl} \xrightarrow{\text{RGI}} g_{\mu\nu}^{q,1} \xrightarrow{\text{RGI}} g_{\mu\nu}^{q,2} \rightarrow \dots \rightarrow g_{\mu\nu}^?$$

- for Schwarzschild as starting point, this leads to the  
Dymnikova BH

$$G_N(r) = G_N \left[ 1 - e^{-r^3/r_s e^2} \right]$$

$$r_s = 2G_N M$$
$$e^2 = \frac{3\epsilon}{8\pi g_x} \ell_p^2$$

### ③ Invariant RGI

Held 2105.11458

- observation: RGI at the level of EoM (solution) is not coordinate-invariant

→  $g_{\mu\nu}$  transforms as a tensor

- idea: take curvature invariant and RG-improve there

- for Schwarzschild:  $R_{\mu\nu} R^{\mu\nu} = 48 \frac{G_N^2 M^2}{r^6}$



compute in terms

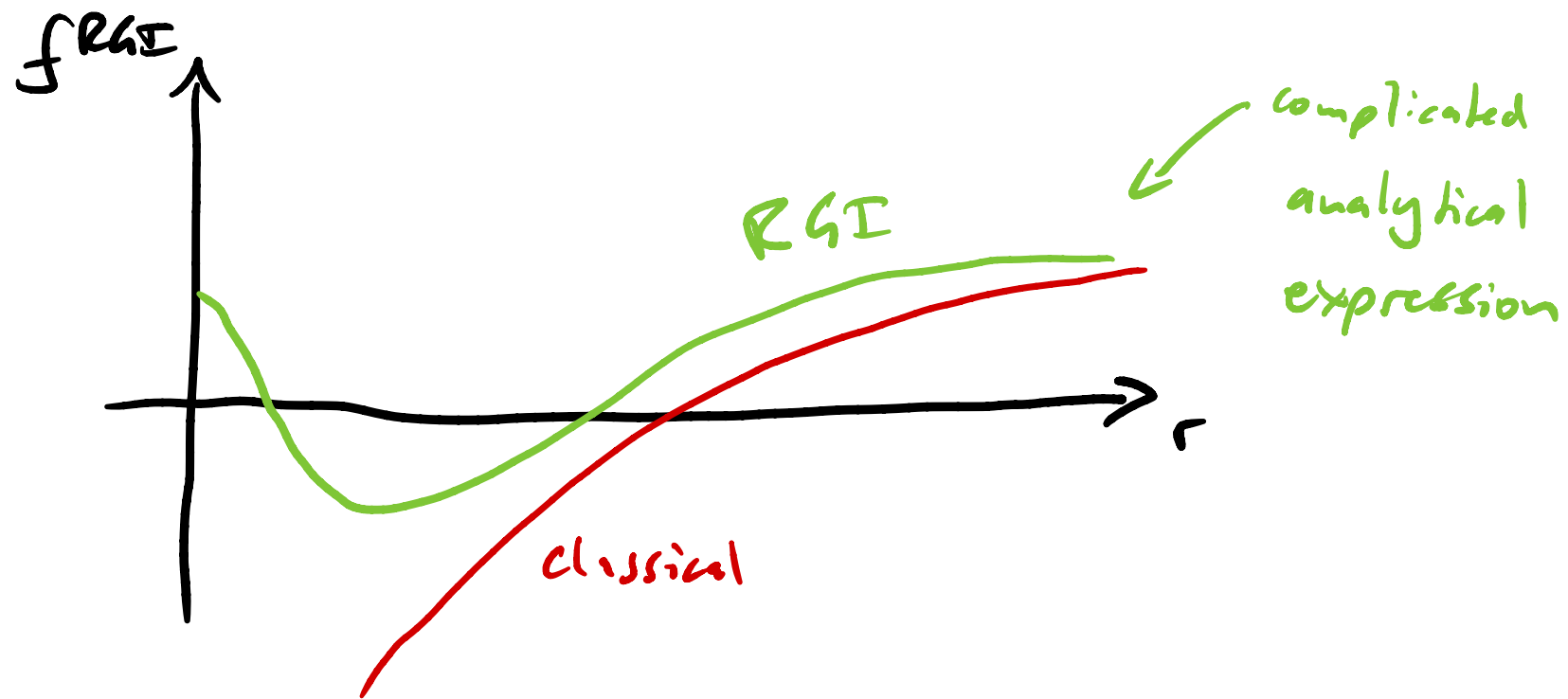
of

$$f^{RGI} = 1 - \frac{2 G_N(k(r)) M}{r}$$



RGI:  $G_N \rightarrow G_N(k(r))$

⇒ differential equation for  $k(r)$



obvious problem: what if there are several independent curvature invariants?

→ not obvious that consistent

RG I procedure exists

# Conceptual issues of RGI

Anber, Donoghue 1111.2875

- all of the above ideas hint towards that RGI has conceptual problems in gravity

- Q: Can RGI work at all in QG?

in the sense of: there is a regime where RGI  
does provide the correct qualitative  
answers, similar to the velocity potential

for this, we go back to perturbative QG and consider

different  $2 \rightarrow 2$  scattering processes involving gravity

$$\Lambda = 0$$

quick reminder :  $2 \rightarrow 2$  scattering

→ 4 momenta  $p_1, p_2, p_3, p_4$

Convention : all ingoing  $p_1 + p_2 + p_3 + p_4 = 0$

→ external lines are on-shell  $\Rightarrow p_i^2 = m_i^2 = 0$

for gravitons

→ Mandelstam variables

$$s = (p_1 + p_2)^2$$

$$t = (p_1 + p_3)^2$$

$$u = (p_1 + p_4)^2$$

with 
$$s + t + u = \sum_{i=1}^4 m_i^2$$

→ the amplitude is a function of  $(s, t, u)$ , or equivalently,  
in the centre-of-mass frame,  $(E^2, \theta)$

$\uparrow$  (energy transfer)<sup>2</sup>       $\uparrow$  scattering angle in COM frame

⇒ if RGI makes sense, then there should be a  
replacement  $G_N \rightarrow G_N(E^2)$  that, inserted into the  
tree-level result, gives the one-loop result

→ this should be universal, i.e., valid for  
all processes

to maximise chances for success, we consider the differential cross section  $\frac{d\sigma}{d\Omega}$  (measure for the effective area of a scattering target)

1) graviton-graviton scattering

✓ helicities (don't worry about it)

$$\mathcal{A}_{\text{tree}}^{++--} \propto G_N \frac{s^3}{tu}$$



match RGI to one-loop  $\frac{d\sigma}{d\Omega}$

$$G_N(E^2) \approx G_N [1 + c_g G_N E^2]$$

$$c_g > 0$$

grows with  $E$

2) grav. scatt. of massless scalars

$$\mathcal{A}_{\text{tree}} \propto G_N \left[ \frac{s^4}{u} + \frac{s^4}{t} + \frac{t^4}{s} \right]$$



$$G_N(E^2) \approx G_N [1 - c_s G_N E^2]$$

$$c_s > 0$$



decreases with  $E$

(similar issue if other helicity channels are considered)

→ suggests that there is **no** reasonable def. of  $R_h$  running for  $G_N$  ??? what did we do then?

→ is  $R_h$  nonsense, or is gravity special?

to answer this, we have to go back to our starting point for renormalisation:

→ original observation: physics changes if a physical energy scale changes

→ what we compute with  $FR_h$ : change when an IR cutoff changes

"prunning" vs. "krunning"

~> physical energy scale is e.g.  $\rightarrow$  momentum  
 $\longleftrightarrow$  derivative in the action

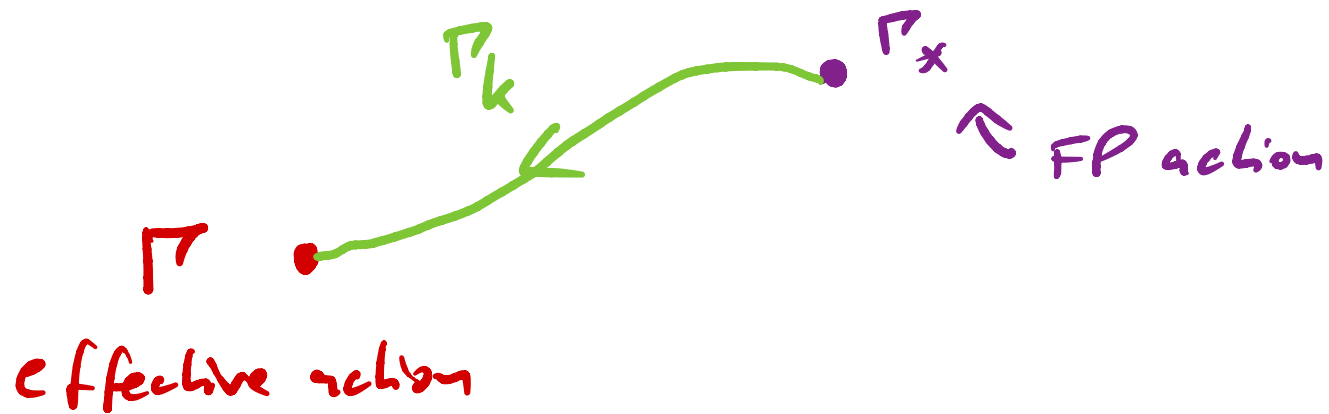
~>  $k$  is a fiducial momentum scale

in some cases,  $k$ -running mimics  $p$ -running, but  
there are exceptions

$\rightarrow$  when talking about RG, one has to specify  
which is meant!

$\nwarrow$  subtlety that leads to a lot  
of confusion

recall what the main idea of the FRG was:



↑  
this is where the physics is

→ all quantum fluctuations included

= "path integral performed"

→ once we have  $\Gamma$ , we can forget that we  
used the FRG to compute it

$\Rightarrow$  running  $\sim$  momentum dependence in  $\Gamma$

$\rightarrow$  in gravity,  $\Gamma \approx \int d^4x \sqrt{g} \left[ \frac{2\Lambda - R}{16\pi G_N} + R \underline{f_R(\Box)} R + R_{\mu\nu} \underline{f_{Ric}(\Box)} R^{\mu\nu} + \dots \right]$

*no running!!!* (pointing to  $\frac{2\Lambda - R}{16\pi G_N}$ )

*running here* (pointing to  $\underline{f_{Ric}(\Box)}$ )

$\Rightarrow G_N$  (and  $\Lambda$ ) do not run with physical momenta!

Suppose they would, then e.g.

Q: what about  
 $\sqrt{R} \frac{1}{G_N(\Box)} \sqrt{R}$ ?

$$\int d^4x \sqrt{g} \frac{1}{G_N(\Box)} R$$

$\Rightarrow$  all potential non-derivative  $\Box$ -dependence are surface terms

$\Rightarrow$  krunning = k-dependence of  $\Gamma_k$

$$\hookrightarrow \Gamma_k = \int d^d x \sqrt{g} \left[ \frac{2\lambda_k - R}{16\pi G_{N,k}} + R f_{R,k}(\Box) R + R_{\mu\nu} f_{R_{\mu\nu},k}(\Box) R^{\mu\nu} + \dots \right]$$

krunning everywhere!

Subject to constraints:

$$k \rightarrow 0: G_{N,k} \rightarrow G_N, \lambda_k \rightarrow \lambda, f_{R,k} \rightarrow f_R, \dots$$

$$k \rightarrow \infty: g_k \rightarrow g_* , \lambda_k \rightarrow \lambda_* , f_{R,k} \rightarrow f_{R*}, \dots$$

dimensionless in units of  $k$

learning: •  $G_N$  and  $\Lambda$  have running

but no pruning

$\Rightarrow$  RGI questionable  $\nearrow$  this explains the problem with the amplitude

idea for RGI of  $f_{R,k}$ : •  $f_R, f_{Ric}, \dots$  have both running and pruning

$$f_{R,k}(0) \xleftrightarrow["k \rightarrow p"]{} f_R(p^2)$$

$\Rightarrow$  RGI potentially more promising

but still issue with Wick rotation:  $p^2 \in \mathbb{R}$  but  $k \in \mathbb{R}_+$

$\leadsto$  for the amplitude, relevant one-loop contributions come from  $R^2, R_{\mu\nu} R^{\mu\nu}$ ; these have a different tensor

structure than the Einstein-Hilbert term

$\Rightarrow$  cannot be both captured by a single coupling

Comments to put this into perspective:

- computing proving is **hard**

$\rightarrow$  many computations rely on expansions in powers of  $g$   
or  $RhI$  where it makes sense

- pheno: people do what they can, analyzing complicated actions with form factors.  $\beta$

often unfeasible

↙ i.e. for  $G_n \& 1$

→ simple R4I often only way  
to make any statement at all

→ nevertheless: "AS-inspired" instead  
of "predictions of AS"

- large part of the literature focusses on  
establishing the FP in the first place

→ crucially important for this

- answering some of the most pressing open  
questions **need** to consider pruning,

e.g. is the theory unitary & causal?

e.g. check bounds on scattering  
amplitudes

→ also strictly needs Lorentzian computations  
of which we have some now

some of these developments came out of constructive criticism

from outside the community

Donoghue 1911.02967

⇒ having an open mind is important for progress,  
scientific or otherwise