

MSSM HIGGS BOSON PRODUCTION VIA BOTTOM PARTONS

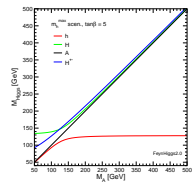
Tilman Plehn
CERN

- Charged Higgs: bottom induced process
- Total rate
- Top and Higgs distributions

MSSM HIGGS BOSONS AT THE LHC

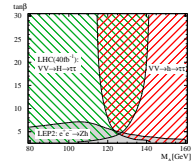
MSSM Higgs Sector

- Softly broken supersymmetric anomaly-free theory
- two doublets, coupling to up and down type fermions
 - five physical states h^0, H^0, A^0, H^\pm
 - mixing of scalars to mass eigenstates (mixing angle α)
 - more predictive than Standard Model (upper h^0 mass limit)
- conveniently expressed as function of m_A and $\tan \beta \equiv v_2/v_1$
- Yukawa couplings to H, A, H^\pm : $m_b \tan \beta, m_t/\tan \beta$ (large m_A)
- typically one light, many heavy scalars [Heinemeyer, Weiglein]



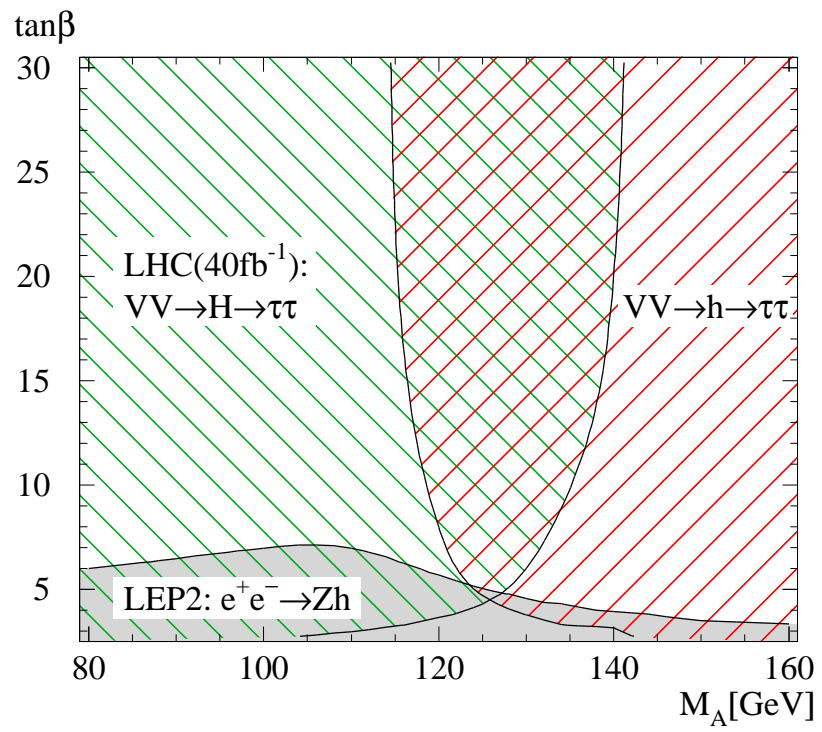
Find first Higgs boson

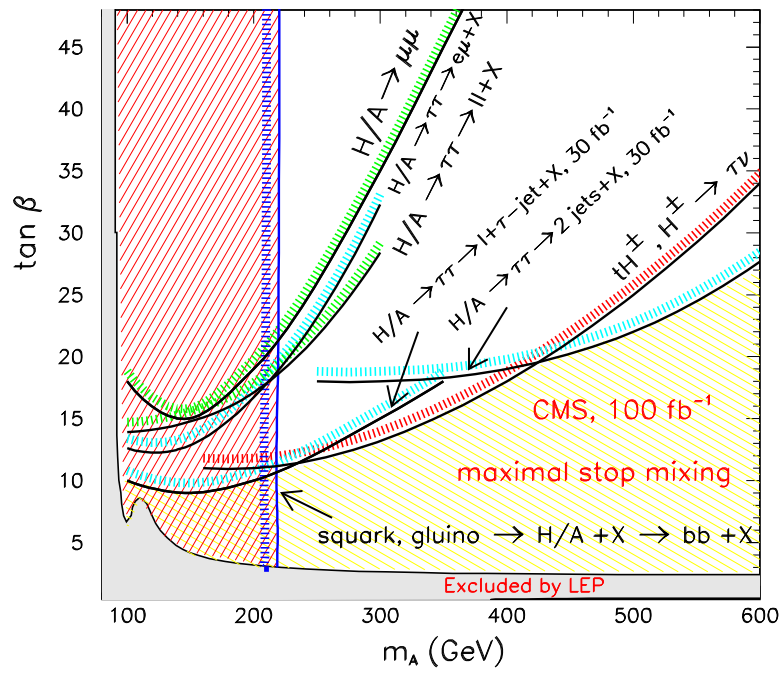
- complete coverage by WBF $h \rightarrow \tau\tau$ [TP, Rainwater, Zeppenfeld; Schumacher]
- problem: mass degeneracy [Boos, Djouadi, Mühlleitner, Nikitenko]
 - $\Delta m_h/m_h \sim \sigma/\sqrt{N}$ ($\sigma \sim 1.5$ GeV for $\mu\mu, \gamma\gamma$ and $\sigma \sim 15$ GeV for $\tau\tau$)



Tell it is 2HDM (MSSM?) \Rightarrow look for heavy Higgs bosons

- $H^0, A^0 \rightarrow \tau\tau, \mu\mu$ inclusive $gg \rightarrow H$ and $gg \rightarrow b\bar{b}H$
- $H^\pm \rightarrow \nu\tau, tb$ in $pp \rightarrow tH^-, W^+H^-, H^+H^-$
 - (n.b. SUSY loops) [Hollik et al, Kniehl et al]
- appearance in SUSY cascades [Datta, Djouadi, Guchait, Moortgat]
- no other conclusive way but to find these particles





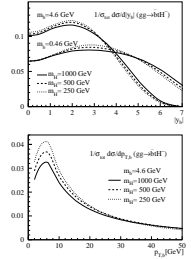
(HEAVY) CHARGED HIGGS

Most promising channel

- associated production $pp \rightarrow tH^- + X$ for large $\tan \beta$
- decay $H^\pm \rightarrow \nu\tau$ most promising [Assamagan, Coadou]

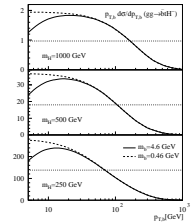
Exclusive production $gg \rightarrow \bar{b}tH^-$

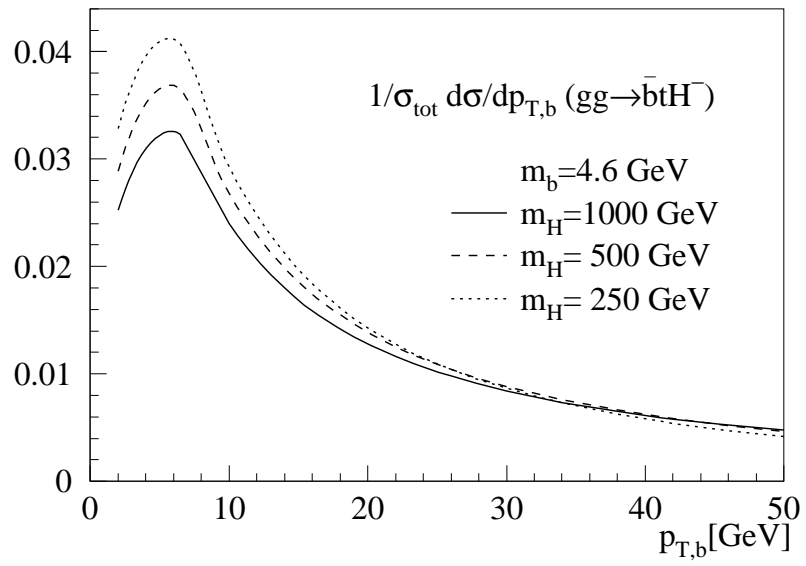
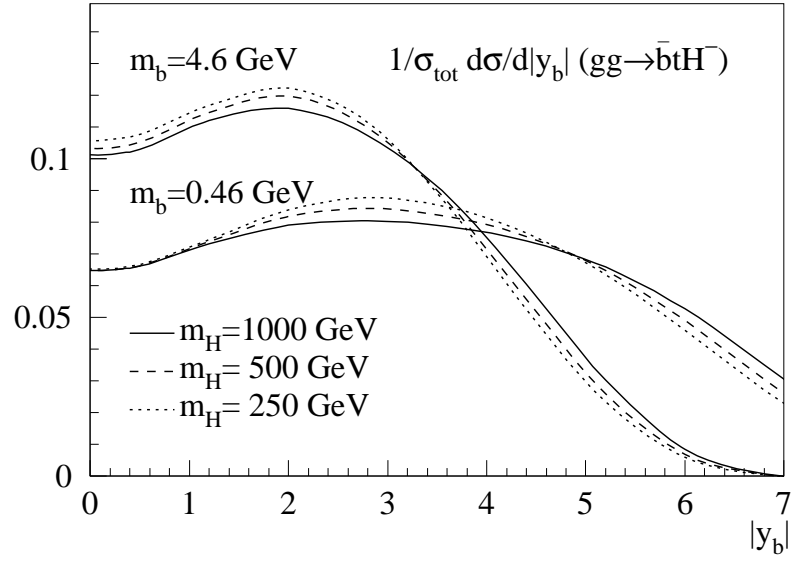
- collinear bottom jets from gluon splitting, regularized by m_b
- experiment: forward jets, $p_{T,b}$ peaked at m_b (factor 1/6 for each tagged b)
- use bottom-inclusive cross section
- check asymptotic cross section behavior $d\sigma/dp_{T,b} \propto 1/p_{T,b}$
- inclusive total rate $\sigma \propto \log(p_{T,b}^{\max}/p_{T,b}^{\min}) = \log(p_{T,b}^{\max}/m_b)$
- how large logarithms? resum?

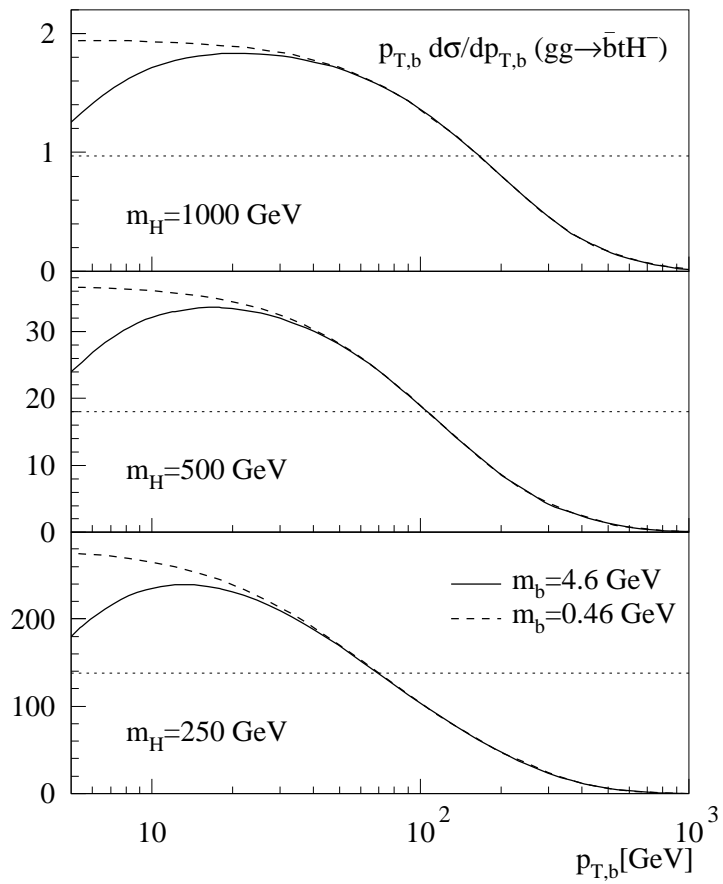


Inclusive process $bg \rightarrow tH^-$

- resum large logarithms $\log(p_{T,b}/m_b)$ in exclusive process $gg \rightarrow \bar{b}tH^-$
- equivalent to bottom parton density and inclusive process $bg \rightarrow tH^-$
- $\mu_{F,b}$ ‘transverse momentum size’ of bottom parton
($\mu_{F,b} \equiv p_{T,b}^{\max}$; usually hard scale $\mu_{F,b} = M$)
- numerical improvement or overestimate?
- (1) check bottom-inclusive total rate
- (2) check bottom-inclusive t, H distributions





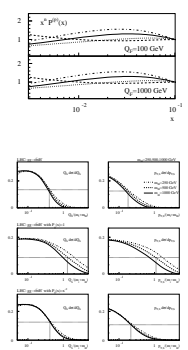


TOTAL RATE: BOTTOM FACTORIZATION SCALE

Perturbative bottom factorization scale from exclusive process [Boos, TP]

- two steps: first bottom virtuality Q_b^{\max}
- general exclusive process: $gg \rightarrow \bar{b}X_M$
- asymptotic behavior $|\overline{\mathcal{M}}|^2 = S^2 \sigma_0 / Q_b^2$; $\mathcal{L} = \mathcal{L}_0 / x^2$

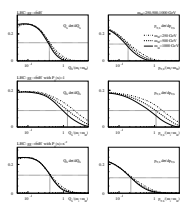
$$\sigma = \frac{2\sigma_0 \mathcal{L}_0}{16\pi} \int_0^{S-M^2} \frac{dQ_b}{Q_b} F(Q_b)$$



- $F(Q_b)$ known correction to asymptotic behavior $d\sigma/dQ_b \sim 1/Q_b$
- define Q_b^{\max} at turning point $d^2F(Q_b)/d(\log Q_b)^2 = 0$
- $Q_b^{\max} \sim M/2$ (hard scale argument $Q_b^{\max} \propto M$, not more than that!)

Second step: transverse momentum $p_{T,b}^{\max}$

- check explicitly: $Q_b \sim Q_b^{\max}$ also yields $p_{T,b} \sim p_{T,b}^{\max}$
- translate Q_b into $p_{T,b}$ point by point
- $p_{T,b}^{\max}/Q_b^{\max} \sim Q_b^{\max}/M$ yields $p_{T,b}^{\max} \sim Q_b^{\max}/2 \sim M/4$
(numerical study of $gg \rightarrow \bar{b}tH^-$: $\mu_{F,b} \sim M/5$)

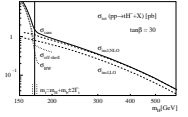


So what did we learn from exclusive process?

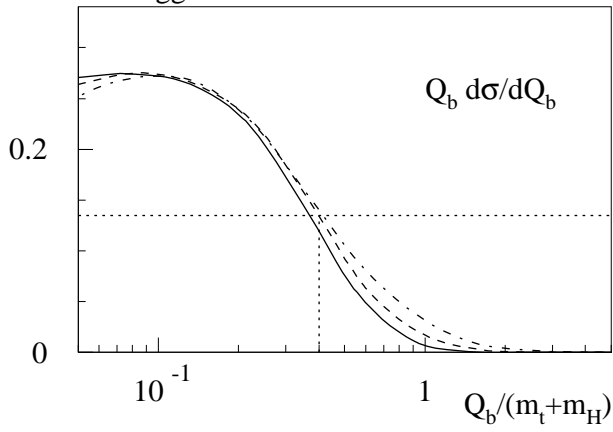
- $\log(p_{T,b}/m_b)$ after integrating over bottom jet
- but 'large' logs at maximum $\log(M/(5m_b))$ [TP; Maltoni, Willenbrock]
- gg and bg processes: $\mu_{F,b} \sim M/5$ from partonic phase space

⇒ Total cross section with bottom partons understood

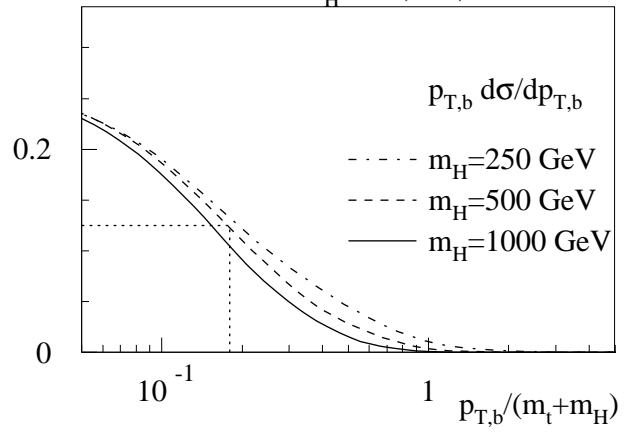
[Dittmaier, Spira, Krämer]



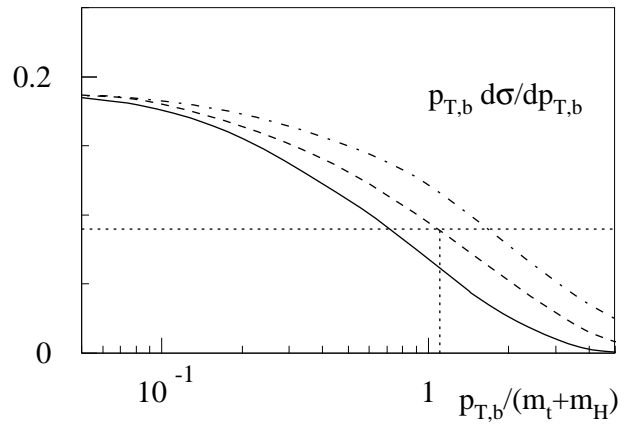
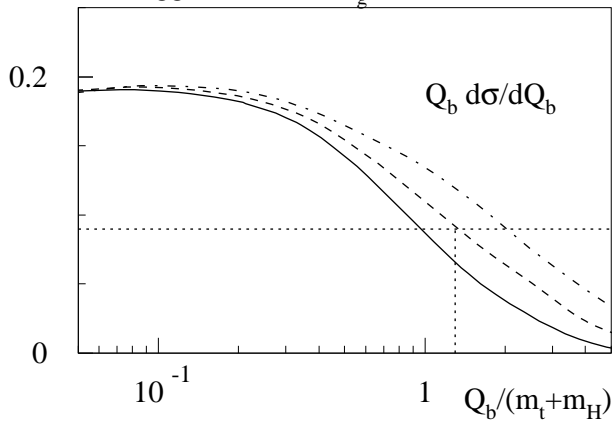
LHC: $gg \rightarrow \bar{b}tH^-$



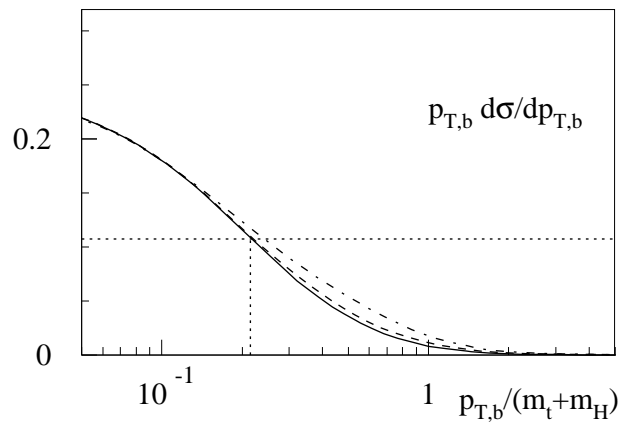
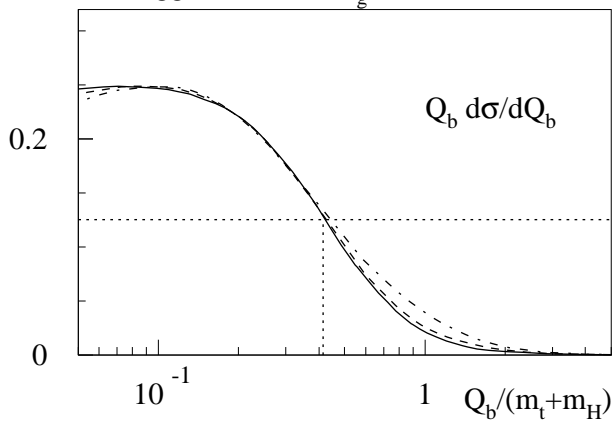
$m_H = 250, 500, 1000$ GeV



LHC: $gg \rightarrow \bar{b}tH^-$ with $P_g(x)=1$



LHC: $gg \rightarrow \bar{b}tH^-$ with $P_g(x)=x^{-2}$



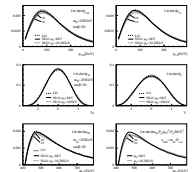
DISTRIBUTIONS FOR INCLUSIVE PROCESS

On to the distributions (preliminary)

- bottom parton description appropriate for total rate
- Higgs and top distributions?
- bottom partons established for exclusive cross sections?

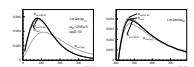
(1) Inclusive kinematics

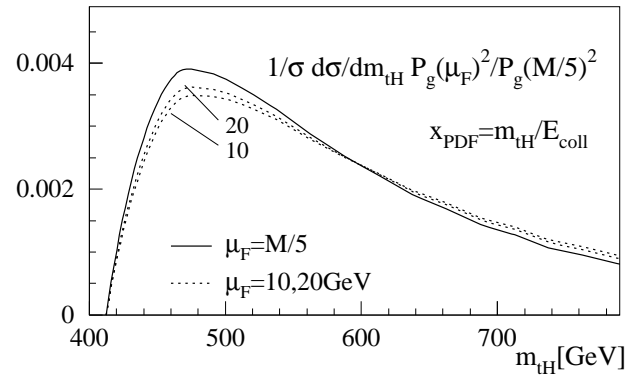
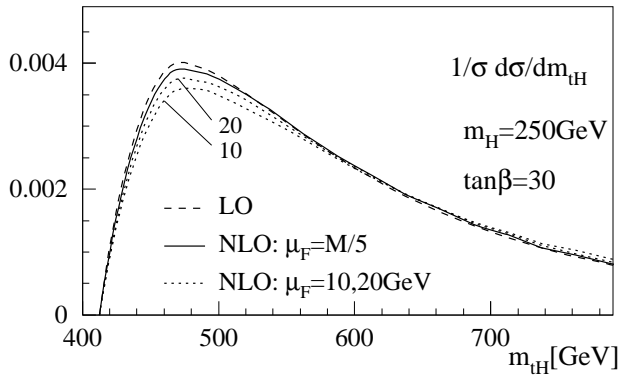
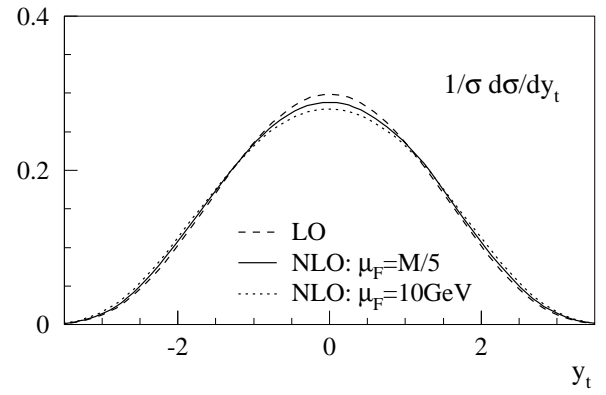
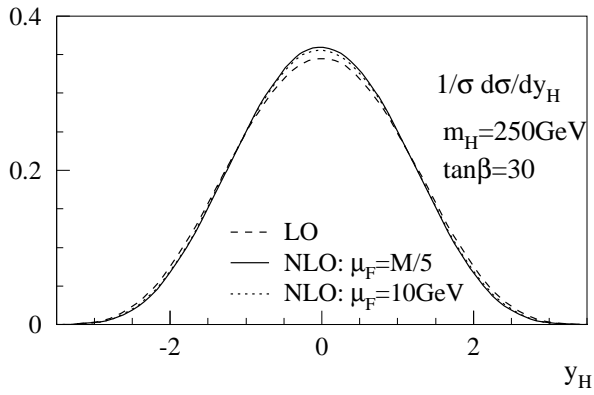
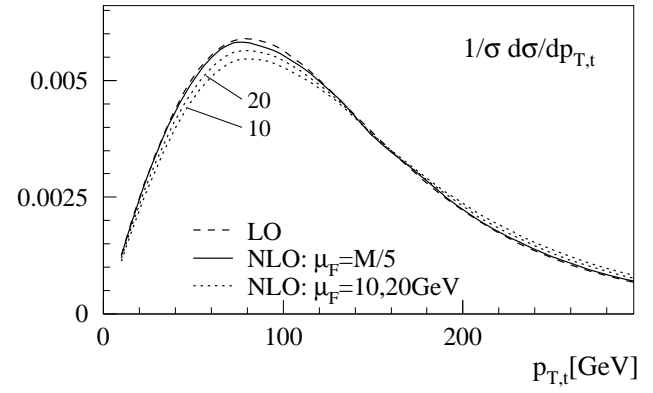
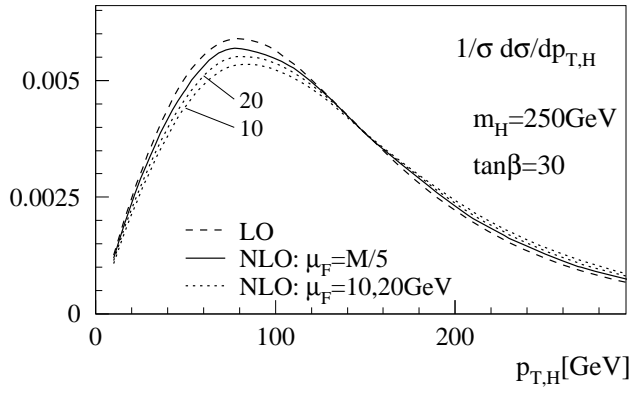
- bottom partons assuming small $p_{T,b} \ll p_{z,b}$
- compare to exclusive ($2 \rightarrow 3$) process which is part of NLO rate
- run bottom factorization scale $\mu_F \rightarrow m_b$
switch off incoming bottoms, left with $gg \rightarrow \bar{b}tH^-$
- slightly harder distributions (due to x dependence of bottom PDF)

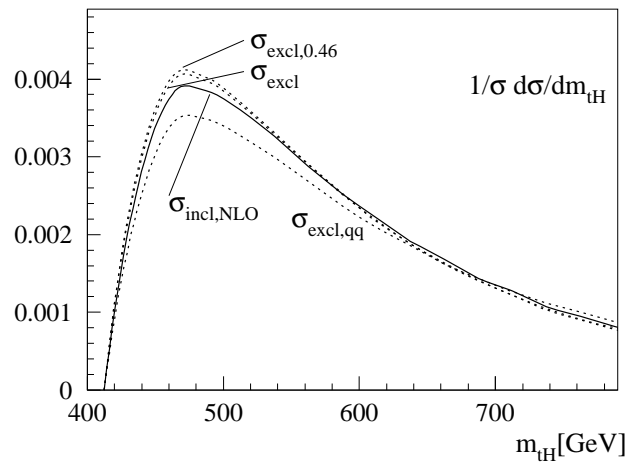
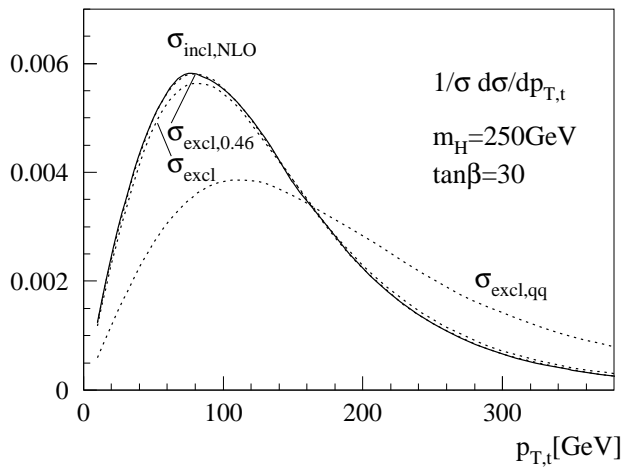


(2) Zero bottom mass

- agreement exclusive vs. inclusive cross section established
- check with bottom mass dependent $pp \rightarrow \bar{b}tH^-$
- perfect agreement with exclusive process for small m_b
very good agreement with physical bottom mass case
- Bottom parton picture altogether appropriate







SUSY-QCD CORRECTIONS

SUSY-QCD Loop Contributions [TP; Berger, Han, Jiang, TP]

- infrared finite but ultraviolet divergent SUSY loop contributions
- (1) universal corrections $y_b/(1 + \Delta_b)$
[Carena, Garcia, Nierste, Wagner; Guasch, Häflinger, Spira]
- (2) remaining explicit SUSY loop diagrams

	m_0	$m_{1/2}$	$\tan \beta$	μ	m_H			$(\Delta_b)_{\text{resum}}$	non- Δ_b
1a	100	250	10	420	477			-10.2%	3.0%
1b	200	400	30	511	535			-23.5%	-0.1%
2	1450	300	10	425	1503			-0.9%	-1.0%
3	90	400	10	633	719			-9.5%	3.0%
4	400	300	50	389	357			-31.0%	-0.4%
5	150	300	5	637	697			-8.0%	10.0%
	m_0	$m_{1/2}$	$\tan \beta$	μ	m_H	M_1	$M_{2,3}$		
6	150	300	10	402	476	480	300	-9.5%	3.0%
	Λ	M_{mes}	N_{mes}	$\tan \beta$	μ	m_H			
7	40×10^3	80×10^3	3	15	316	476			-8.1% 0.5%
8	100×10^3	200×10^3	1	15	421	538			-7.1% 0.5%

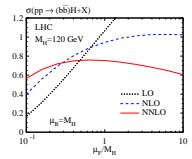
→ Δm_b corrections dominant for $\tan \beta \gtrsim 10$ (dependent on sign of μ)

→ explicit loop corrections negligible $\lesssim 10\%$ for generic mSUGRA

(HEAVY) NEUTRAL HIGGS

Bottom induced production of neutral Higgses

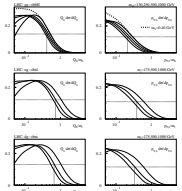
- rate enhanced by $\tan^2 \beta$
- $gg \rightarrow b\bar{b}H$ exclusive versus $bg \rightarrow bH$ inclusive
 $bg \rightarrow bh$ exclusive versus $b\bar{b} \rightarrow H$ inclusive
- appropriate factorization scale $\mu_{F,b} \sim M/5 = m_h/5$
- check: $b\bar{b} \rightarrow H$ NNLO scale dependence [Harlander & Kilgore]
 $\mu_{R,b}$ variation for fixed $\mu_{F,b} \sim m_h/4$ well under control
 $\mu_{F,b}$ variation for fixed $\mu_{R,b} \sim m_h$ almost fixed point
- check: exclusive vs. inclusive total rate [Dittmaier, Spira, Krämer]



	M_H	$\sigma(q\bar{q}, gg \rightarrow b\bar{b}H + X)$ [fb]		$\sigma(b\bar{b} \rightarrow H + X)$ [fb]	
		LO	NLO	LO	NNLO
Tevatron	120	$3.9^{+3.5}_{-1.7}$	$8.0^{+3.1}_{-2.4}$	$8.6^{+4.7}_{-5.0}$	$10.5^{+0.3}_{-1.1}$
	200	$0.22^{+0.19}_{-0.09}$	$0.56^{+0.23}_{-0.18}$	$0.69^{+0.20}_{-0.26}$	$0.79^{+0.02}_{-0.03}$
LHC	120	$(5.3^{+2.7}_{-1.7}) \times 10^2$	$(7.3^{+2.0}_{-1.6}) \times 10^2$	$(4.8^{+4.3}_{-3.2}) \times 10^2$	$(7.2^{+0.4}_{-1.6}) \times 10^2$
	400	$4.3^{+2.4}_{-1.4}$	$8.1^{+2.2}_{-1.9}$	$7.4^{+2.4}_{-2.5}$	$9.8^{+0.2}_{-0.4}$

Side remark: single top production $qg \rightarrow \bar{b}tq'$ [Willenbrock et al]

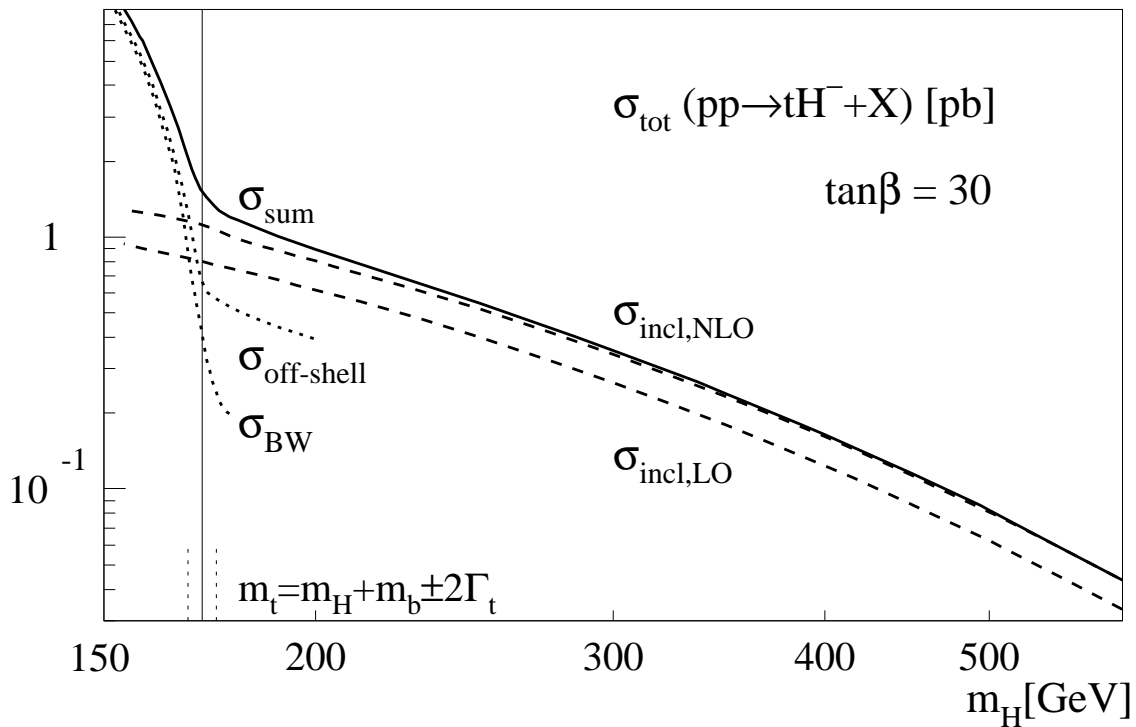
- less steep quark densities, $x_1 \neq x_2$
- production above threshold
- $Q_b^{\max} \sim m_t$
- generally $p_{T,b}^{\max} \sim Q_b^{\max}/2$
- $\mu_{F,b} \sim m_t/2$ covered by quoted theoretical uncertainty

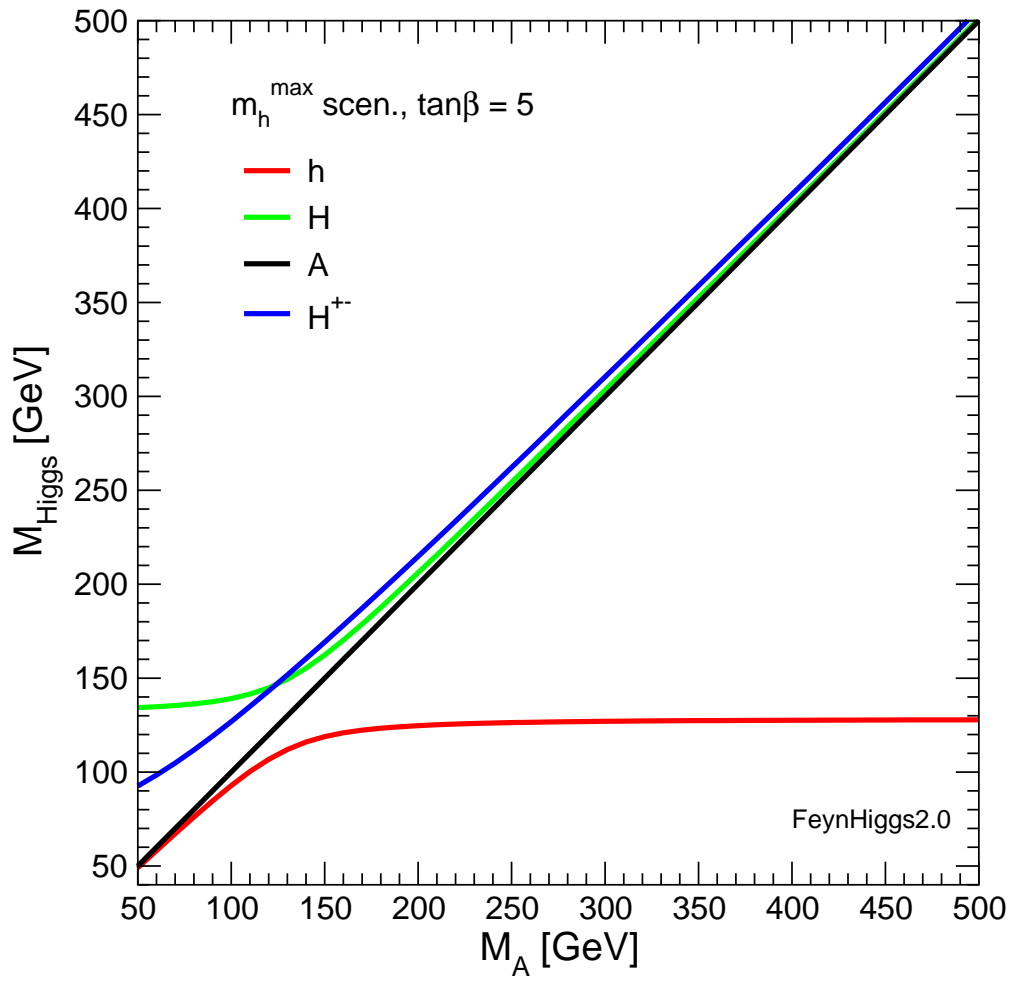


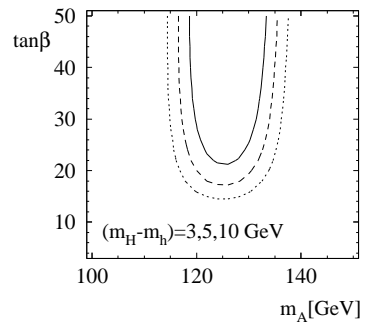
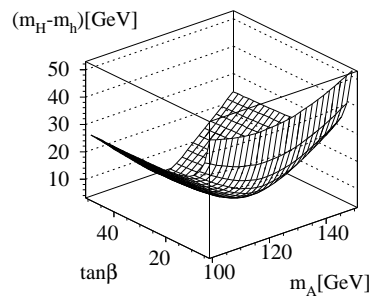
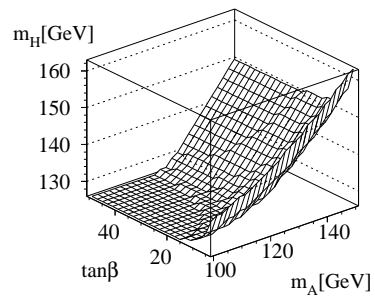
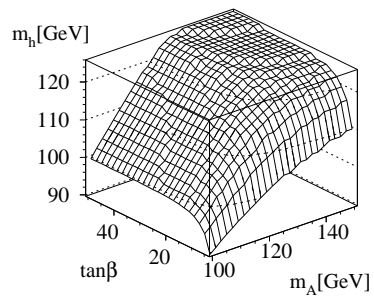
CONCLUSIONS

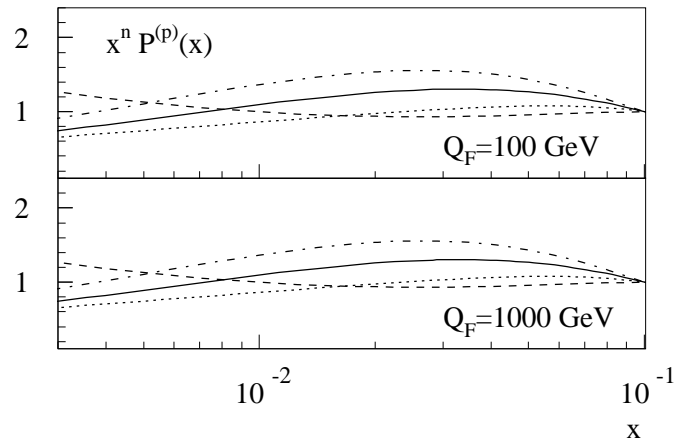
Bottom parton picture works fine

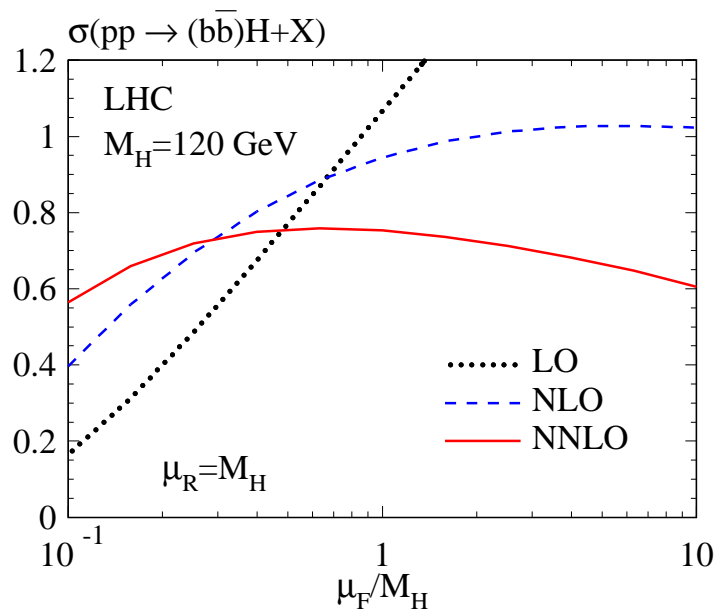
- total rate correct with appropriate factorization scale
- top and Higgs distributions correctly described
- we also understand why











Harlander, Kilgore