

SUPERSYMMETRY – AN INTRODUCTION

Tilman Plehn

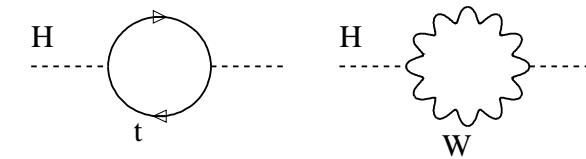
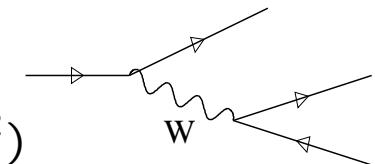
Max Planck Institute for Physics (upstairs)

- Fine tuning in the Standard Model
- Supersymmetry as one solution
- The MSSM
- Supersymmetry at the Tevatron and the LHC
- Supersymmetry at the Linear Collider
- and lots of little side remarks...

WHAT IS WRONG WITH THE STANDARD MODEL: 1

A brief history of the mess we are in...

- Fermi 1934: theory of weak interactions $[n \rightarrow p e^- \bar{\nu}_e \text{ and } \mu \rightarrow e^- \bar{\nu}_e \nu_\mu]$
 - divergent four fermion reaction amplitude: $\mathcal{A} \propto G_F E^2$
 - unitarity violation [transition probability $\propto |\mathcal{A}|^2$]
 - ‘effective theory’ valid for $E < 600 \text{ GeV}$
 - Yukawa 1935: massive virtual particle exchange
 - Fermi’s theory for $E \ll M$
 - four fermions unitary for large energies: $\mathcal{A} \propto g^2 E^2 / (E^2 - M^2)$
 - unitarity violation in $WW \rightarrow WW$ $[\mathcal{A} \propto G_F E^2]$
 - Higgs & Kibble 1967: spontaneous symmetry breaking
 - unitary gauge theory with massive W, Z [discovered 1983]
 - fermion masses linked to Yukawa couplings [top quark discovered 1995]
 - fundamental Higgs scalar below TeV scale [mass unknown]
 - Higgs couplings fixed $g_{HXX} \propto g m_X$
 - **desaster strikes:** scalar mass perturbatively unstable
 - quadratic divergences $\delta m_H^2 \propto g^2 \Lambda^2$ [Λ loop cutoff]
 - all-orders Higgs mass driven to cutoff $m_H \rightarrow \Lambda$
- ⇒ new approach needed to solve hierarchy problem



WHAT IS WRONG WITH THE STANDARD MODEL: 2

Starting from data...

- ...which seem to indicate a light Higgs
- problem of light Higgs: mass driven to cutoff of theory [remember pain to get that up]
$$\delta m_H^2 \propto g^2 (2m_W^2 + m_Z^2 + m_H^2 - 4m_t^2) \Lambda^2$$
Veltman's condition ($\dots = 0$) would be fun
problem preferably solved to arbitrary loop order
- ⇒ solution: counter term for exact cancellation ⇒ **artificial, unmotivated, ugly**
- ⇒ or new physics at TeV scale: **supersymmetry**
extra dimensions
little Higgs (pseudo–Goldstone Higgs)
Higgsless/composite Higgs
YourFavoriteNewPhysics...
- ⇒ typically either cancellation with new particles or discarding away high scale
- ⇒ all really beautiful concepts and symmetries
- ⇒ in general problematic to realize at TeV scale [data seriously in the way]

Idea of supersymmetry: cancellation of divergences through statistics factor (-1)
[scalars vs. SM fermions; fermions vs. SM gauge bosons; fermions vs. SM scalars]

TEV SCALE SUPERSYMMETRY: 1

Great idea: solve hierarchy problem by only doubling particle number?!

- stops (scalar) cancel top loop [couplings also protected]
 - gauginos (neutral or charged) cancel W, Z loop
 - higgsinos cancel Higgs loop [mix with neutralinos]
 - since we are at it, let's postulate gluino
 - top not special, so postulate sleptons and squarks
- ⇒ hierarchy problem gone
- ⇒ rich collider and non-collider phenomenology guaranteed

Slight problem with Higgs sector

- adjoint Higgs field not allowed in \mathcal{L}
→ how to give mass to t and b?
→ two doublets
- ⇒ SUSY Higgs sector interesting by itself
[no SUSY knowledge needed]

		spin	d.o.f.	
fermion → sfermion	f_L, f_R \tilde{f}_L, \tilde{f}_R	1/2 0	1+1 1+1	
gluon → gluino	G_μ \tilde{g}	1 1/2	n-2 2	Majorana
gauge bosons Higgs bosons → neutralinos	γ, Z h^0, H^0, A^0 $\tilde{\chi}_i^0$	1 0 1/2	2+3 3 4 · 2	Majorana
gauge bosons Higgs bosons → charginos	W^\pm H^\pm $\tilde{\chi}_i^\pm$	1 0 1/2	2 · 3 2 2 · 4	Dirac
graviton → gravitino	G \tilde{G}	2 3/2	5 5	hard to catch

SUPERSYMMETRIC HIGGS SECTOR: 1

Required by supersymmetry: two Higgs doublet model

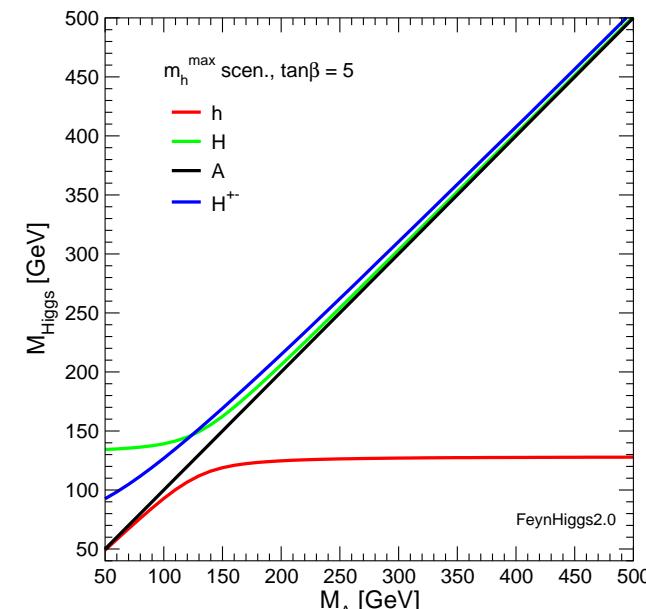
- one (complex) Higgs doublet: 4 degrees of freedom
→ three for longitudinal W, Z , one for scalar Higgs
- two Higgs doublets: 8 degrees of freedom
→ three for longitudinal W, Z , five for Higgs particles
→ scalars h^0, H^0 , pseudoscalar A^0 , charged H^\pm
- free parameters
 - (1) still only one free mass scale: m_A
 - (2) two vacuum expectation values: $\tan \beta = v_t/v_b$

Plateau structure

- only plateau Higgs coupling to W, Z
- heavy Higgs with $\tan \beta$ -enhanced coupling to b, τ
- light Higgs mass limited from above

$$2m_h^2 \simeq m_{AZ}^2 - \sqrt{m_{AZ}^4 - 4m_A^2 m_Z^2 c_{2\beta}^2 - 4 \frac{3G_F}{\sqrt{2}\pi^2} \frac{m_t^4}{s_\beta^2} \log \frac{m_t^2}{m_b^2} (m_A^2 s_\beta^2 + m_Z^2 c_\beta^2)} \lesssim 2 \cdot 135 \text{ GeV}$$

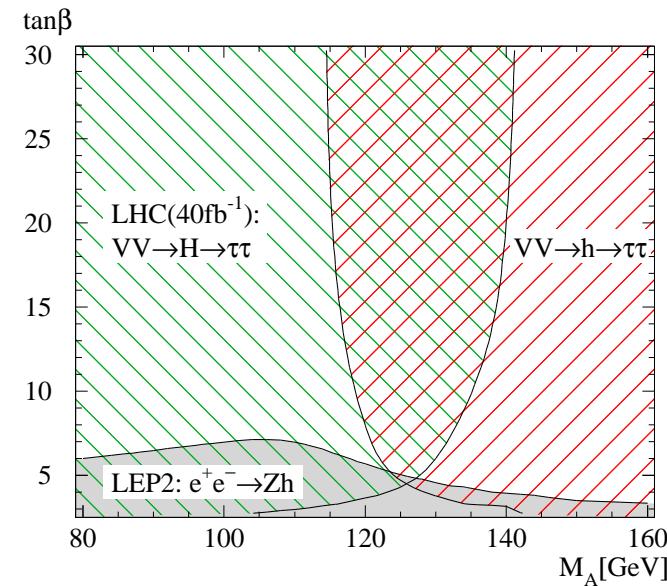
⇒ Little hierarchy problem $m_{\tilde{t}} \gg m_t$ because $m_h > 114 \text{ GeV a la LEP2}$



SUPERSYMMETRIC HIGGS SECTOR: 2

Challenge: find one Higgs at the LHC [very difficult at Tevatron]

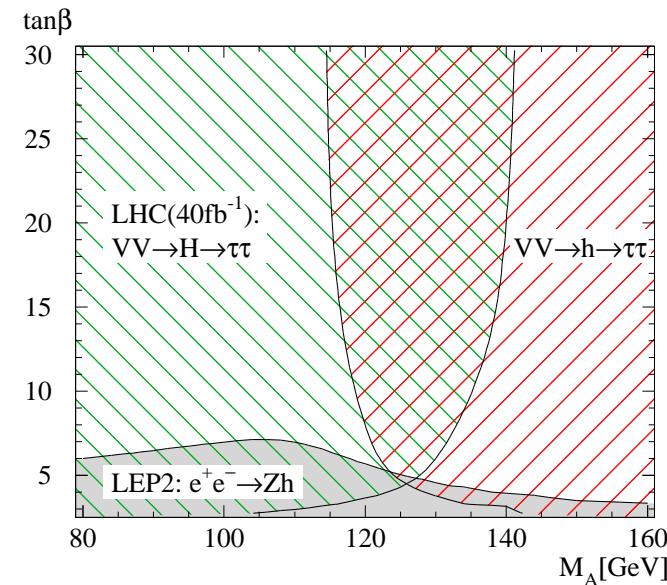
- ‘decoupling regime’ $m_A \gtrsim 160$ GeV
→ h^0 looks like SM Higgs [of corresponding mass]
 - opposite case $m_A \lesssim 120$ GeV
→ $qq \rightarrow qqH^0 \rightarrow qq\tau\tau$ as in SM
 - in between: maybe even two bums
- ⇒ No-lose theorem: $qq \rightarrow qq\{h^0, H^0\} \rightarrow qq\tau\tau$



SUPERSYMMETRIC HIGGS SECTOR: 2

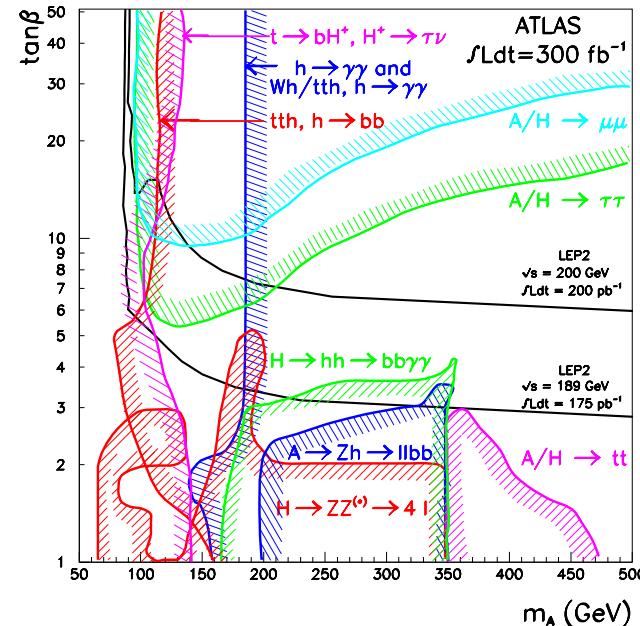
Challenge: find one Higgs at the LHC [very difficult at Tevatron]

- ‘decoupling regime’ $m_A \gtrsim 160$ GeV
→ h^0 looks like SM Higgs [of corresponding mass]
 - opposite case $m_A \lesssim 120$ GeV
→ $qq \rightarrow qqH^0 \rightarrow qq\tau\tau$ as in SM
 - in between: maybe even two bums
- ⇒ **No-lose theorem:** $qq \rightarrow qq\{h^0, H^0\} \rightarrow qq\tau\tau$



Nightmare: confirm SUSY Higgs sector

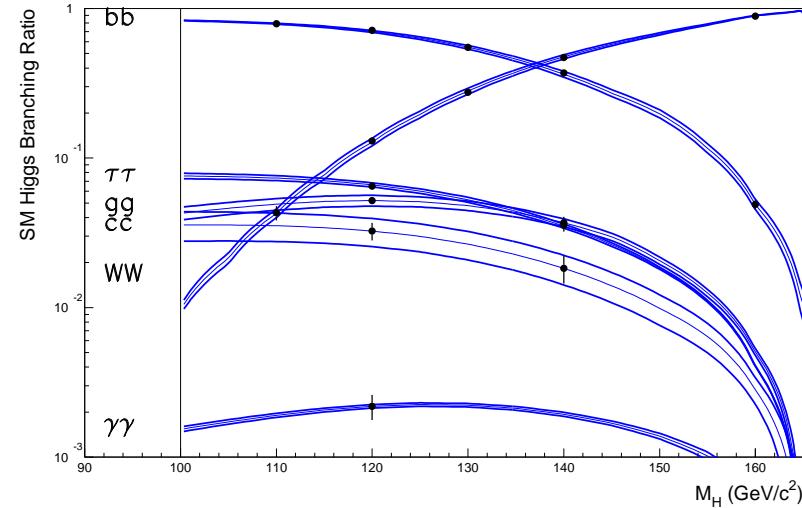
- decoupling regime $m_A \gtrsim 160$ GeV
→ Yukawa coupling for H^0, A^0, H^\pm : $m_b \tan \beta$
→ production e.g. $b\bar{b} \rightarrow H^0$ or $gb \rightarrow tH^-$
→ decays e.g. $H^0 \rightarrow \tau\tau, \mu\mu$ or $H^- \rightarrow \tau\bar{\nu}$
- intermediate $m_A \sim 120$ GeV:
→ lots of Higgs bosons h^0, H^0, A^0 observable
→ problem distinguishing them: $H^0 \rightarrow \mu\mu$
→ top quark decays $t \rightarrow bH^+$? [Tevatron sample small]



SUPERSYMMETRIC HIGGS SECTOR: 3

Higgs searches at the ILC

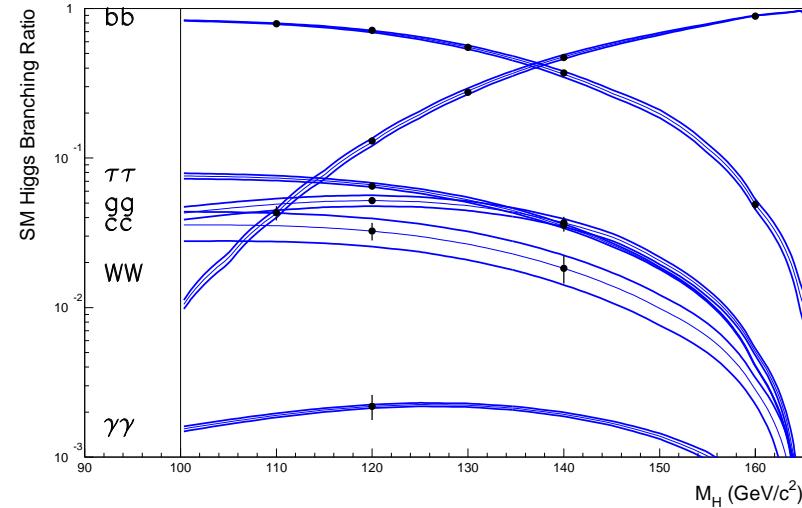
- production rate the challenge (energy), backgrounds under control
- Higgs-strahlung $e^+e^- \rightarrow Zh^0$ for SM-like Higgs
- pair production H^+H^- , A^0h^0 , ...
single production $WW \rightarrow h^0$, $\gamma\gamma \rightarrow h^0, H^0$, ...
- precision the key



SUPERSYMMETRIC HIGGS SECTOR: 3

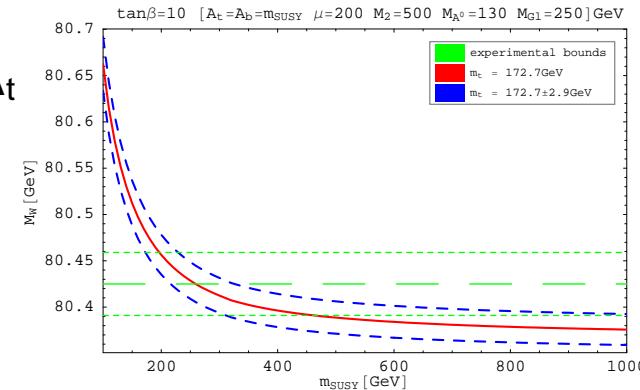
Higgs searches at the ILC

- production rate the challenge (energy), backgrounds under control
- Higgs-strahlung $e^+e^- \rightarrow Zh^0$ for SM-like Higgs
- pair production H^+H^- , A^0h^0 , ...
- single production $WW \rightarrow h^0$, $\gamma\gamma \rightarrow h^0, H^0$, ...
- precision the key



Electroweak analyses at the ILC [Wolfgang Hollik's talk]

- similar to LEP–electroweak fit
 - input SUSY-Higgs, electroweak sector with SUSY loops
 - logarithmic dependence on SUSY masses
- ⇒ measure W mass and some SUSY masses, derive A_t
measure SUSY masses and tell people m_W is wrong

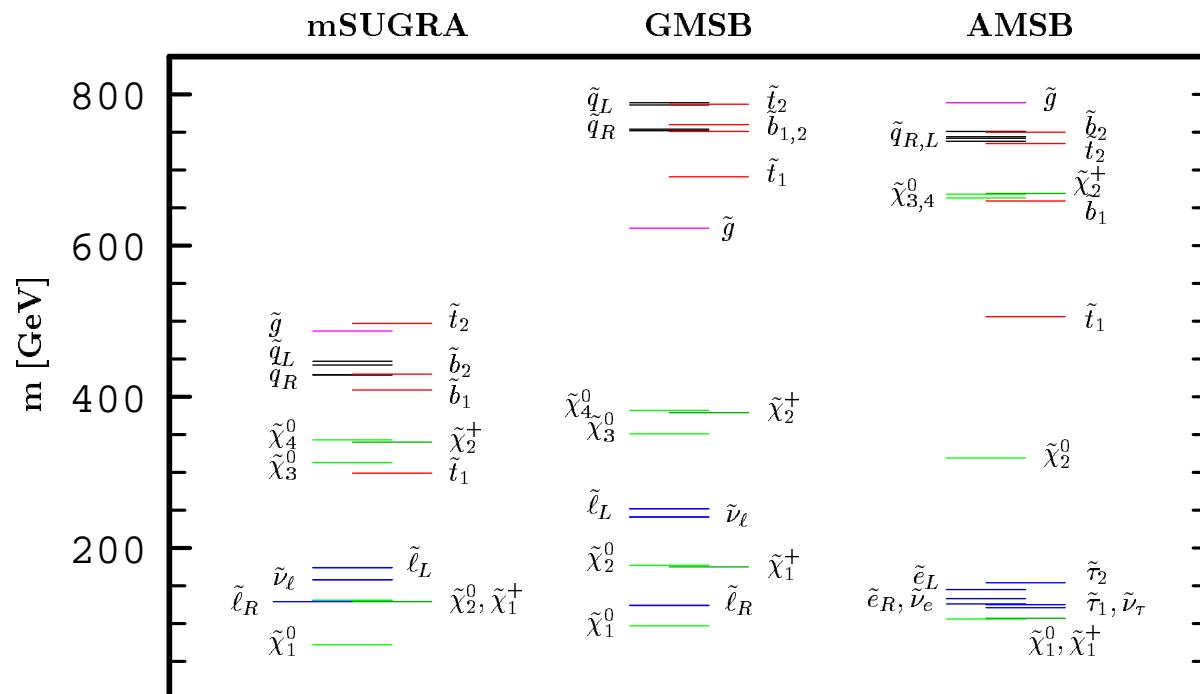


SUSY particle masses

- mechanism of mass generation unknown [soft breaking leaves quadratic divergences alone]
- link to flavor physics and baryogenesis/leptogenesis unknown
- assume SUSY breaking in hidden sector
+ mediation to visible sector [gravitationally coupled, branes in extra dimensions]
- **maximally blind mediation at high scale: mSUGRA**
scalars: m_0 , fermions: $m_{1/2}$, tri-scalar term: A_0
plus $\text{sign}(\mu)$ and $\tan \beta$ from Higgs sector [Higgs masses let free: NUHM]

Alternatives to mSUGRA

- gauge mediation
 - anomaly mediation
 - gaugino mediation
- ⇒ none of them convincing
- ⇒ **measure spectrum**



Structures in the SUSY spectrum

- gauginos–higgsinos mixing: $m_{\tilde{\chi}_2^0} \sim m_{\tilde{\chi}_1^+}$ or $m_{\tilde{\chi}_1^0} \sim m_{\tilde{\chi}_1^+}$ in **MSSM**

$$\begin{pmatrix} m_{\tilde{B}} & 0 & -m_Z s_w c_\beta & m_Z s_w s_\beta \\ 0 & m_{\tilde{W}} & m_Z c_w c_\beta & -m_Z c_w s_\beta \\ -m_Z s_w c_\beta & m_Z c_w c_\beta & 0 & -\mu \\ m_Z s_w s_\beta & -m_Z c_w s_\beta & -\mu & 0 \end{pmatrix} \begin{pmatrix} m_{\tilde{W}} & \sqrt{2} m_w s_\beta \\ \sqrt{2} m_w c_\beta & -\mu \end{pmatrix}$$

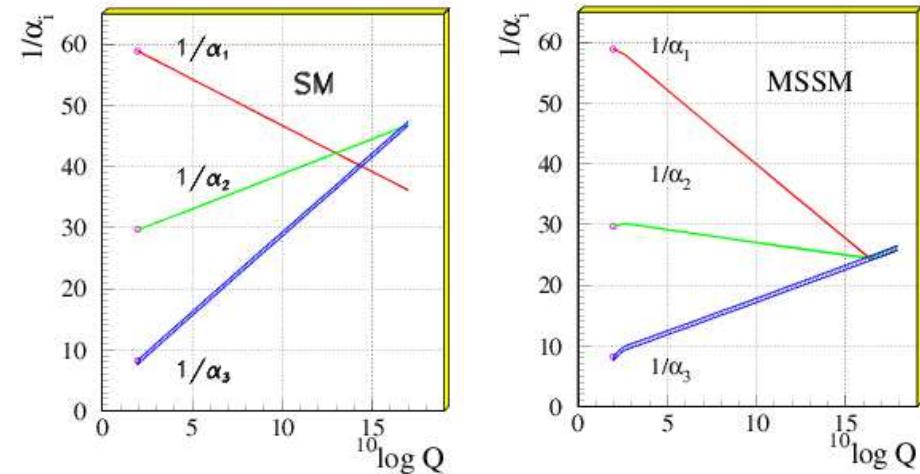
- stop and sbottom mixing in **MSSM**

$$\begin{pmatrix} m_Q^2 + m_t^2 + \left(\frac{1}{2} - \frac{2}{3}s_w^2\right) m_Z^2 c_{2\beta} & -m_t (A_t + \mu \cot \beta) \\ -m_t (A_t + \mu \cot \beta) & m_U^2 + m_t^2 + \frac{2}{3}s_w^2 m_Z^2 c_{2\beta} \end{pmatrix}$$

- heavy gluinos and squarks through **unification**: $m_{\tilde{B}, \tilde{W}, \tilde{g}} / m_{1/2} \sim 0.4, 0.8, 2.6$
 $m_{\tilde{\ell}, \tilde{q}} / m_{1/2} \sim 0.7, 2.5$ [$m_0 \ll m_{1/2}$]
[mass and coupling unification independent]
- lightest SUSY partner $\tilde{\chi}_1^0, \tilde{\nu} \Rightarrow$ after dark matter data $\tilde{\chi}_1^0 \sim \tilde{B}, \tilde{W}$ [gravitinos?]

Unification of couplings

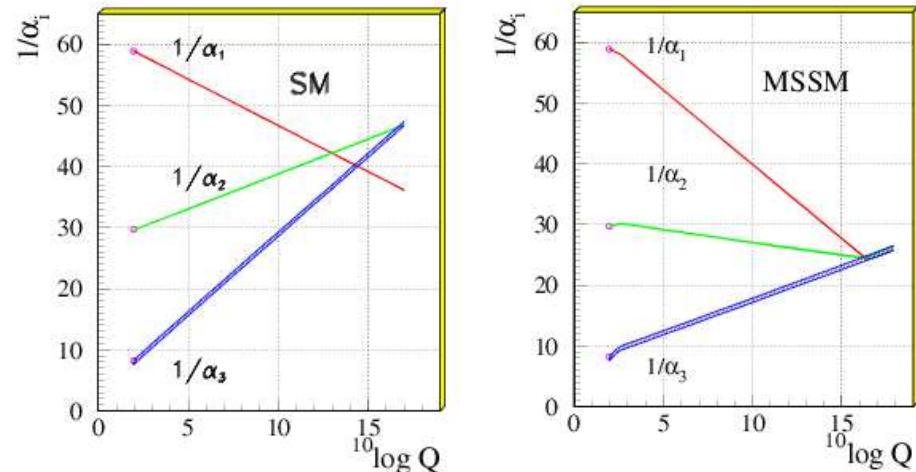
- gauge couplings run: $g_{U(1)}, g_{SU(2)}, g_{SU(3)}$
- remember observed running of α_s
- running described as $\mu dg/d\mu = \beta_g$
- beta functions from particle content
- at one loop $\beta(n_H, n_s, n_f, \dots)$
 $[\beta_{\text{SQCD}} = 11/3N - 2/3N - 2/3n_f - 1/3n_s]$



Unification of couplings

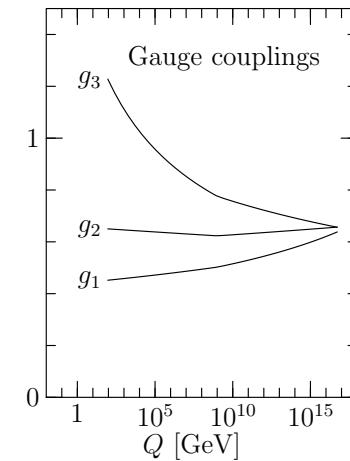
- gauge couplings run: $g_{U(1)}, g_{SU(2)}, g_{SU(3)}$
- remember observed running of α_s
- running described as $\mu dg/d\mu = \beta_g$
- beta functions from particle content
- at one loop $\beta(n_H, n_s, n_f, \dots)$

$$[\beta_{\text{SQCD}} = 11/3N - 2/3N - 2/3n_f - 1/3n_s]$$



Unification of mass parameters

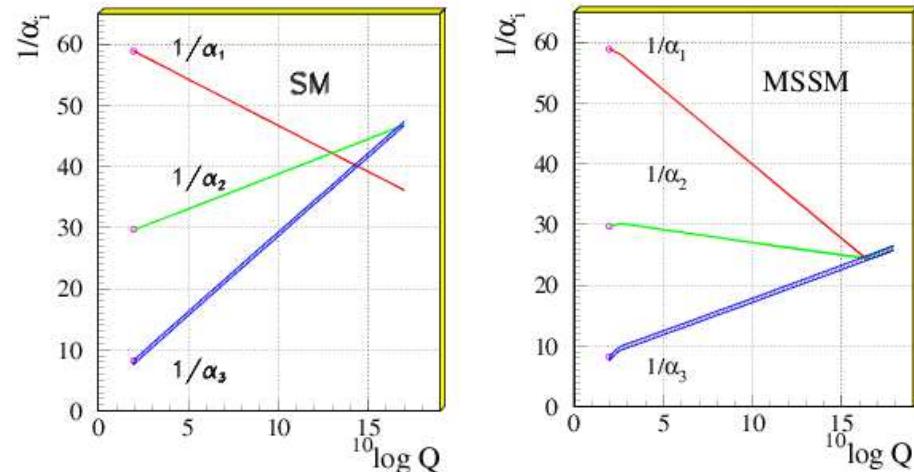
- parameters which gets renormalized run because of $\mu^{2\epsilon} \int dq^{4-2\epsilon}$
 - Standard Model masses run: $m_b, m_t, m_{\tilde{g}}, \dots$
 - SUSY-breaking parameters run: $m_{\tilde{g}}, m_{\tilde{B}}, m_{\tilde{W}}, \dots$
- ⇒ insight into SUSY breaking at some scale?



Unification of couplings

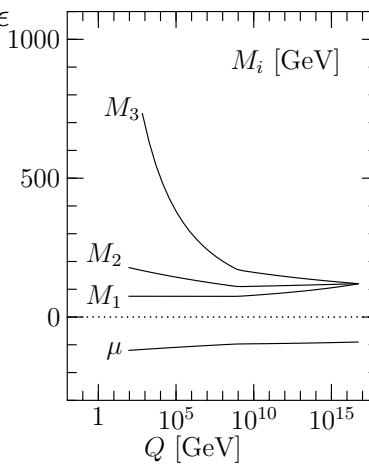
- gauge couplings run: $g_{U(1)}, g_{SU(2)}, g_{SU(3)}$
- remember observed running of α_s
- running described as $\mu dg/d\mu = \beta_g$
- beta functions from particle content
- at one loop $\beta(n_H, n_s, n_f, \dots)$

$$[\beta_{\text{SQCD}} = 11/3N - 2/3N - 2/3n_f - 1/3n_s]$$



Unification of mass parameters

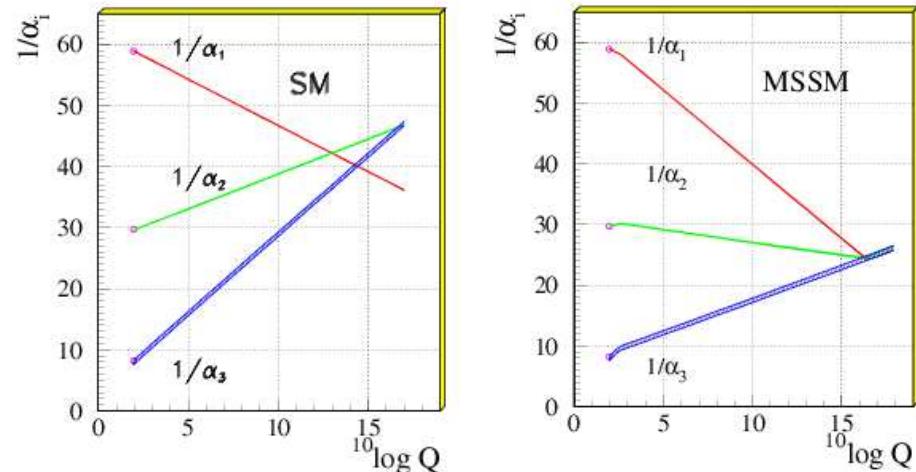
- parameters which gets renormalized run because of $\mu^{2\epsilon} \int dq^{4-2\epsilon}$
 - Standard Model masses run: $m_b, m_t, m_{\tilde{g}}, \dots$
 - SUSY-breaking parameters run: $m_{\tilde{g}}, m_{\tilde{B}}, m_{\tilde{W}}, \dots$
- ⇒ insight into SUSY breaking at some scale?



Unification of couplings

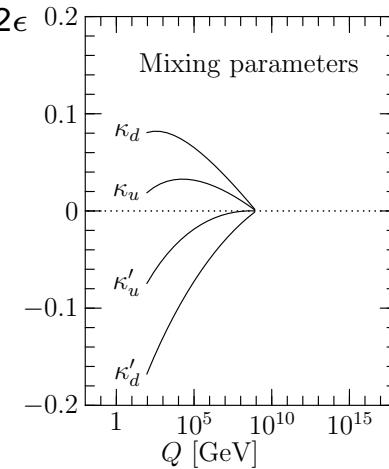
- gauge couplings run: $g_{U(1)}, g_{SU(2)}, g_{SU(3)}$
- remember observed running of α_s
- running described as $\mu dg/d\mu = \beta_g$
- beta functions from particle content
- at one loop $\beta(n_H, n_s, n_f, \dots)$

$$[\beta_{\text{SQCD}} = 11/3N - 2/3N - 2/3n_f - 1/3n_s]$$



Unification of mass parameters

- parameters which gets renormalized run because of $\mu^{2\epsilon} \int dq^{4-2\epsilon}$
 - Standard Model masses run: $m_b, m_t, m_{\tilde{g}}, \dots$
 - SUSY-breaking parameters run: $m_{\tilde{g}}, m_{\tilde{B}}, m_{\tilde{W}}, \dots$
- ⇒ insight into SUSY breaking at some scale?



Supersymmetric parameter conventions

- comparison of specialized codes crucial [remember: e.g. Comphep–Pythia–Isajet]
- ⇒ fix SUSY conventions once for all
 - soft breaking parameters [e.g. $\pm A_t$]
 - scale dependence of couplings, masses [e.g. $m(q = \text{TeV}, v, m_t)$?]
 - definitions of mass matrixes, mixing angles [e.g. $\tilde{t}_{L,R}$ up or down?]

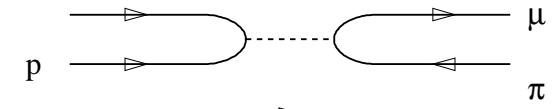
SUSY Les Houches Accord [P. Skands et al.]

- spectrum generators: SoftSusy, SPheno, FeynHiggs,...
- multi-purpose Monte Carlos: Pythia, Herwig, Sherpa
- matrix element generators: Whizard, Smadgraph
- NLO cross sections: Prospino2
- NLO decay rates: Sdecay
- SUSY parameter extraction: Fittino, Sfitter
- dark matter: Micromegas
- ⇒ **fixed parameter convention and read-write format** [list to be extended]

Problem with proton decay

- SUSY allows vertices like $qq\tilde{q}$, $q\ell\tilde{\ell}$, $\ell\ell\tilde{\ell}$
 - problem with dimension-4 proton decay
 - problem with leptoquark-type constraints: atomic parity violation
 - problem with $(e - \mu)$ conversion
 - problem with K and B decays and mixing
- ⇒ **define ad-hoc multiplicative quantum number $R = (-1)^{2S-L+3B}$**

[+1 for SM, -1 for SUSY partners, +1 for Higgses]



Impact on collider phenomenology

- every vertex involving 2 or 4 SUSY partners
 - no single SUSY-partner production
 - lightest SUSY partner stable (LSP)
 - stable LSP at the end of each cascade
- ⇒ **SUSY searches always for leptons+jets+missing energy** [no pure QCD backgrounds]

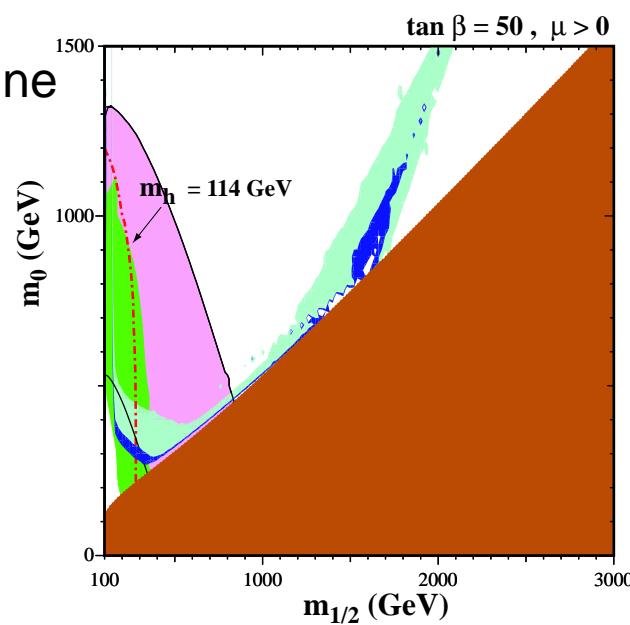
DARK MATTER IN SUPERSYMMETRY

Dark matter [Georg Raffelt's talk]

- total matter density same as critical density
- observed (baryonic) matter much less
- light neutrinos too relativistic to match equation of state
- rotation of galaxies too fast (only explained by additional matter)
- cannot have charge or strong interaction, because of missing evidence

SUSY candidates with conserved R parity

- Weakly Interacting Massive Particles at weak scale fine
 - sneutrinos excluded by searches
 - lightest neutralino prime candidate
 - thermally produced in equilibrium, then frozen out
 - annihilation $\tilde{\chi}\tilde{\chi} \rightarrow ff$ necessary
s-channel annihilation through A^0 at large $\tan\beta$?
co-annihilation $\tilde{\chi}\tilde{\tau} \rightarrow Z\tau$ helpful?
- ⇒ few observables in a large parameter space, though



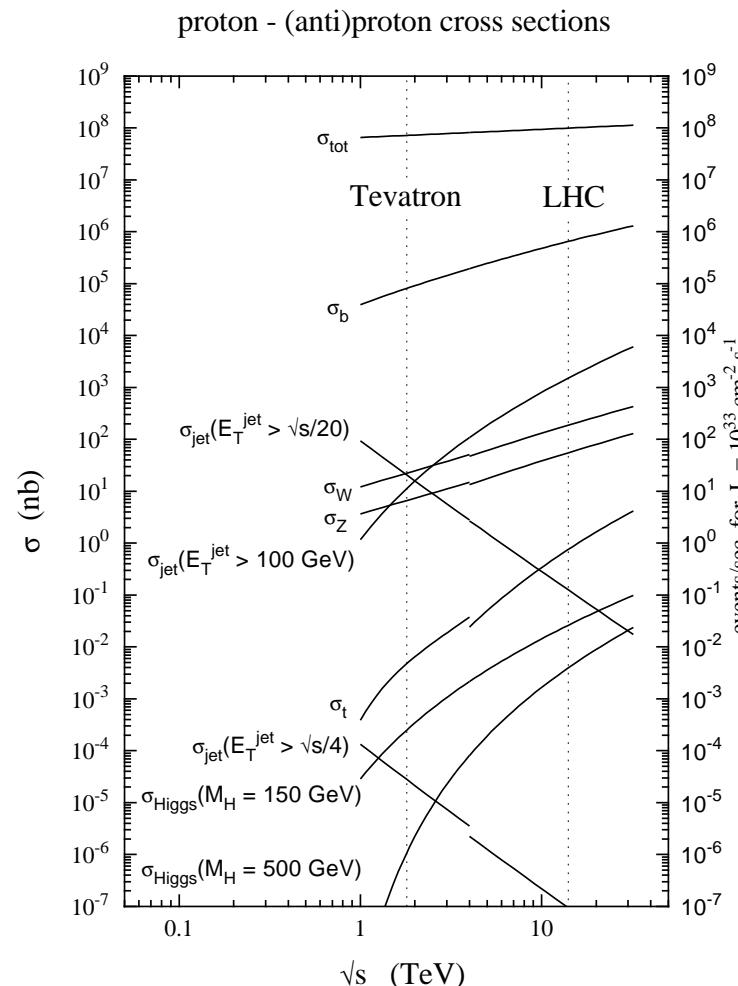
HADRON COLLIDERS: TEVATRON & LHC

Conversion of energy into mass [$E = mc^2$ with $c=1$]

- search for new particles [easier if real particle produced]
→ highest possible energies required
- electron colliders: collide two highly relativistic e
LEP: 46 GeV each to produce Z
LEP2: 103 GeV each for e.g. ZH
ILC/CLIC: 1...4 TeV in the future
- hadron colliders: collide (anti)protons
Tevatron: $p\bar{p}$ collision with 2 TeV [valence quarks]
LHC: pp collisions with 14 TeV [gluons]
- **LHC mass reach ~ 3 TeV** [trade luminosity for energy]

New physics with hadron colliders

- what is a jet and what is inside? [bottom tag, τ tag]
- trigger: ‘no leptons — no data’ [e.g. $pp \rightarrow t\bar{t} \rightarrow$ jets]
- backgrounds rates $pp \rightarrow jj$ or $pp \rightarrow WZ +$ jets
- **statistical significance: $S/\sqrt{B} > 5$ is discovery**



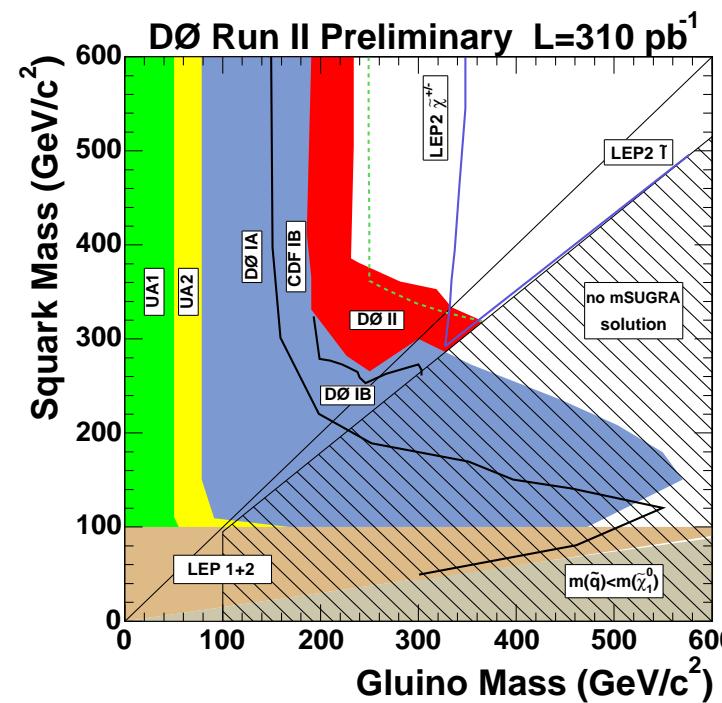
SUPERSYMMETRY AT THE TEVATRON: 1

Promising SUSY channels at the Tevatron

- until now: best limit $m(\tilde{\chi}^+) > 103 \text{ GeV}$ from LEP
[captures essentially all leptons+LSP channels]
- strongly interacting squarks and gluinos
 $p\bar{p} \rightarrow \tilde{q}\tilde{q}^*, \tilde{q}\tilde{g}, \tilde{g}\tilde{g}$ [best if $m(\tilde{q}) \sim m(\tilde{g})$]
- decays to jets and LSP and
 $\tilde{g} \rightarrow \tilde{q}\bar{q}, \tilde{q}_L \rightarrow q\tilde{\chi}_2^0, \tilde{q}_R \rightarrow q\tilde{\chi}_1^0$
 additional jets and leptons possible
- ⇒ **search as inclusive as possible**

Necessary model assumptions

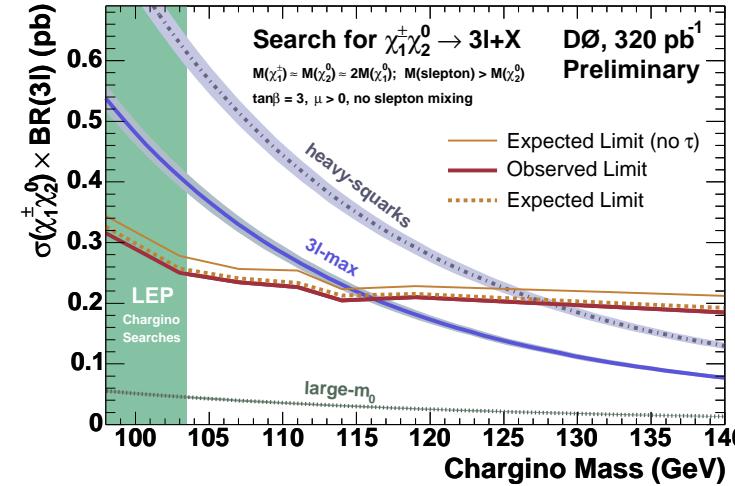
- assume 100% branching into inclusive jets+LSP
- for detector efficiencies $m(\tilde{\chi}_2^0), \dots$
- gluino decay into squark or vice versa?
- ⇒ **mSUGRA assumed, but effect moderate**



SUPERSYMMETRY AT THE TEVATRON: 2

More specialized: trilepton channels

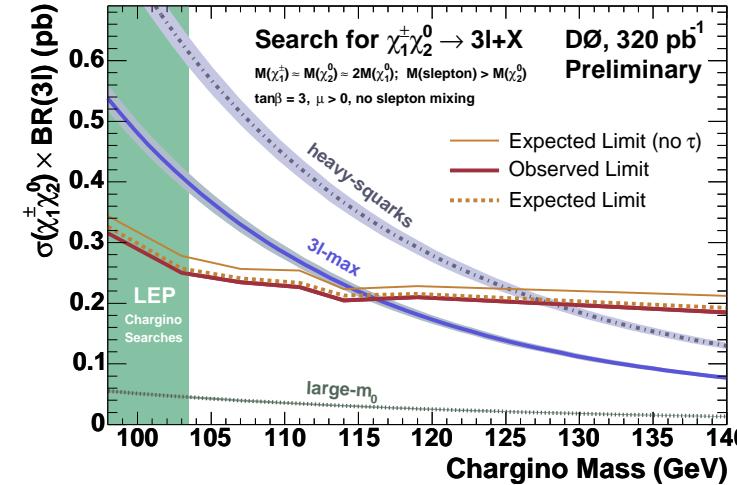
- generally assumed that charginos are lighter than aluinos
- largest cross section $p\bar{p} \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_2^0$
decays $\tilde{\chi}_1^+ \rightarrow \ell^+ \nu \tilde{\chi}_1^0$ and $\tilde{\chi}_2^0 \rightarrow \ell^+ \ell^- \tilde{\chi}_1^0$
- gaugino rate determined by t channel squark
- trilepton signature plagued by W, Z background
- ⇒ about to pass LEP chargino limits



SUPERSYMMETRY AT THE TEVATRON: 2

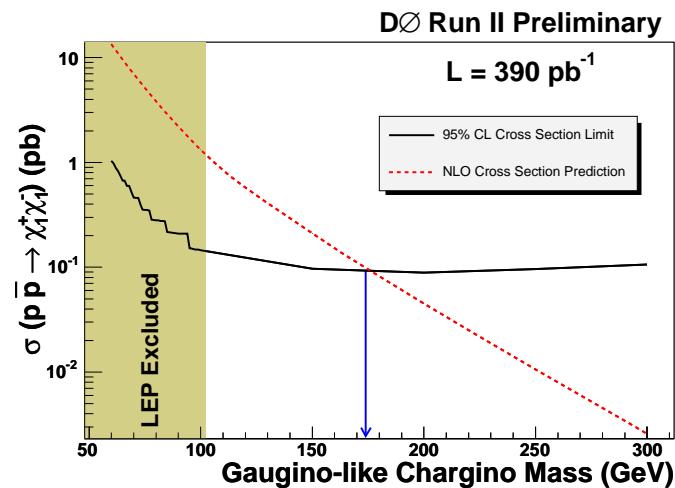
More specialized: trilepton channels

- generally assumed that charginos are lighter than aluinos
- largest cross section $p\bar{p} \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_2^0$
decays $\tilde{\chi}_1^+ \rightarrow \ell^+ \nu \tilde{\chi}_1^0$ and $\tilde{\chi}_2^0 \rightarrow \ell^+ \ell^- \tilde{\chi}_1^0$
- gaugino rate determined by t channel squark
- trilepton signature plagued by W, Z background
- ⇒ about to pass LEP chargino limits



Even more specialized: long-lived gauginos

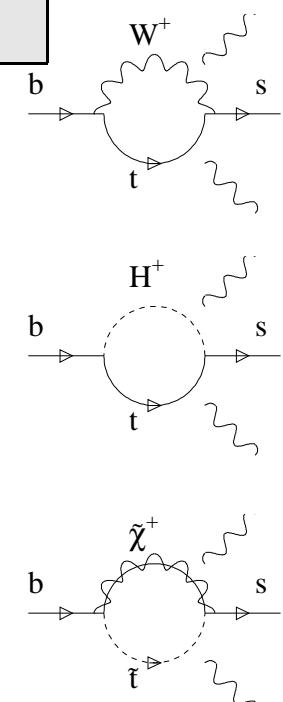
- stable on the detector time scale of ns
- like massive muons in tracker–muon chamber
- gauginos motivated by GMSB or AMSB



TEVATRON: B PHYSICS AND SUPERSYMMETRY

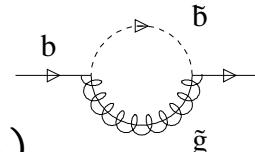
Tevatron channel: $B_s \rightarrow \mu\mu$

- s-channel exchanges dominant: H, Z, γ
suppressed in Standard Model $[BR_{SM} \sim (2.4 \pm 0.5) \times 10^{-9}]$
- more Higgs bosons in 2HDM
 $\tan \beta$ enhancement of s channel Higgses $[BR_{2HDM} \propto \tan^6 \beta / m_A^4]$
additional Higgs loop
- charginos in MSSM
 $\tan \beta$ enhancement for Higgsinos
gluino loop for non-minimal flavor physics...



Bottom Yukawa in the MSSM

- gluino-sbottom loops universal: $y_b \rightarrow y_b / (1 + \Delta_b)$
- large, leading in $\tan \beta$ & resummable $\Delta_b \sim \alpha_s \tan \beta m_{\tilde{g}} \mu / \max^2(m_{\tilde{b}, \tilde{g}})$
 \Rightarrow decoupling in MSSM, but not in MSSM+μ
[similar terms for chargino/neutralino exchange]
- easy to implement in MC, numerically great for $\tan \beta > 10$
- \Rightarrow enhancement good for SUSY signals, but pain in analyses

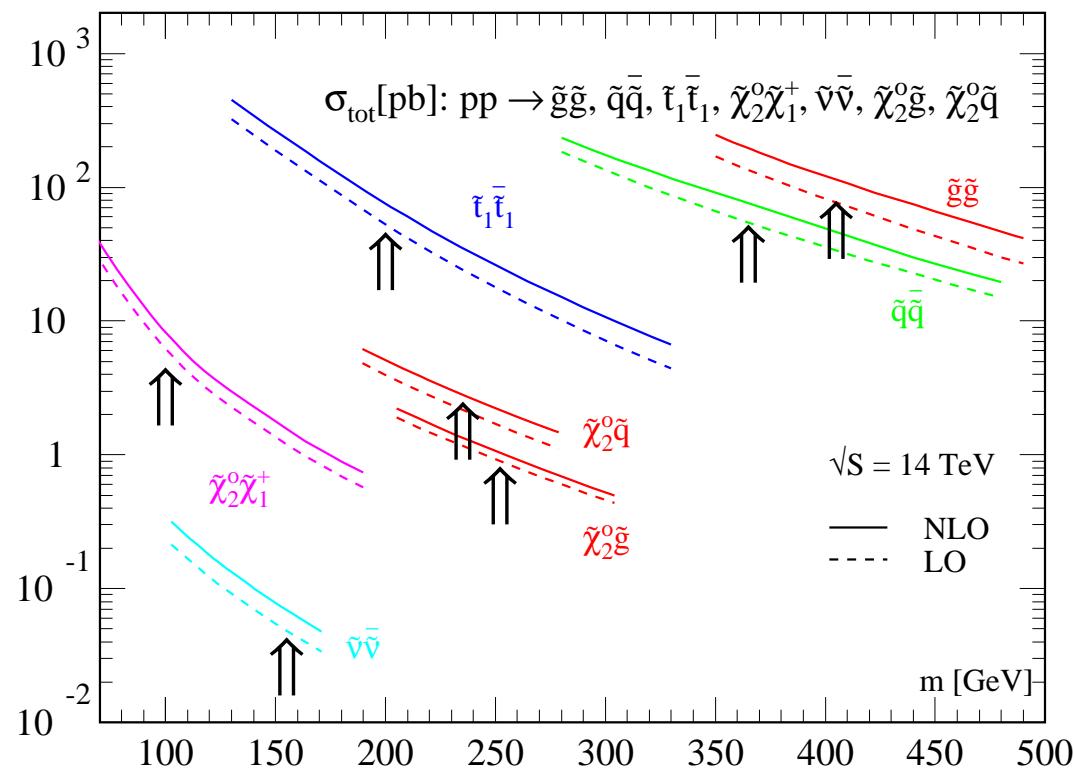


Supersymmetry at the LHC

- (1) **possible discovery** — signals for new physics, exclusion of parameter space
- (2) **measurements** — masses, cross sections, decays
- (3) **parameter studies** — MSSM Lagrangean, SUSY breaking

SUSY signals included

- QCD coupling $g\tilde{q}\tilde{q}$, $q\tilde{q}\tilde{q}$, $g\tilde{g}\tilde{g}$
- jets and \cancel{E}_T : $pp \rightarrow \tilde{q}\tilde{q}^*, \tilde{g}\tilde{g}, \tilde{q}\tilde{g}$
- funny tops: $pp \rightarrow \tilde{t}_1\tilde{t}_1^*$
- like sign dileptons: $pp \rightarrow \tilde{g}\tilde{g}$
 $[\tilde{g} \rightarrow \tilde{u}\bar{u} \rightarrow \tilde{\chi}_1^+ d\bar{u} \text{ or } \tilde{g} \rightarrow \tilde{u}^* u \rightarrow \tilde{\chi}_1^- \bar{d}u]$
- tri-leptons: $pp \rightarrow \tilde{\chi}_2^0\tilde{\chi}_1^-$
 $[\tilde{\chi}_2^0 \rightarrow \tilde{\ell}\bar{\ell} \rightarrow \tilde{\chi}_1^0\ell\bar{\ell}; \tilde{\chi}_1^- \rightarrow \tilde{\chi}_1^0\ell\bar{\nu}]$

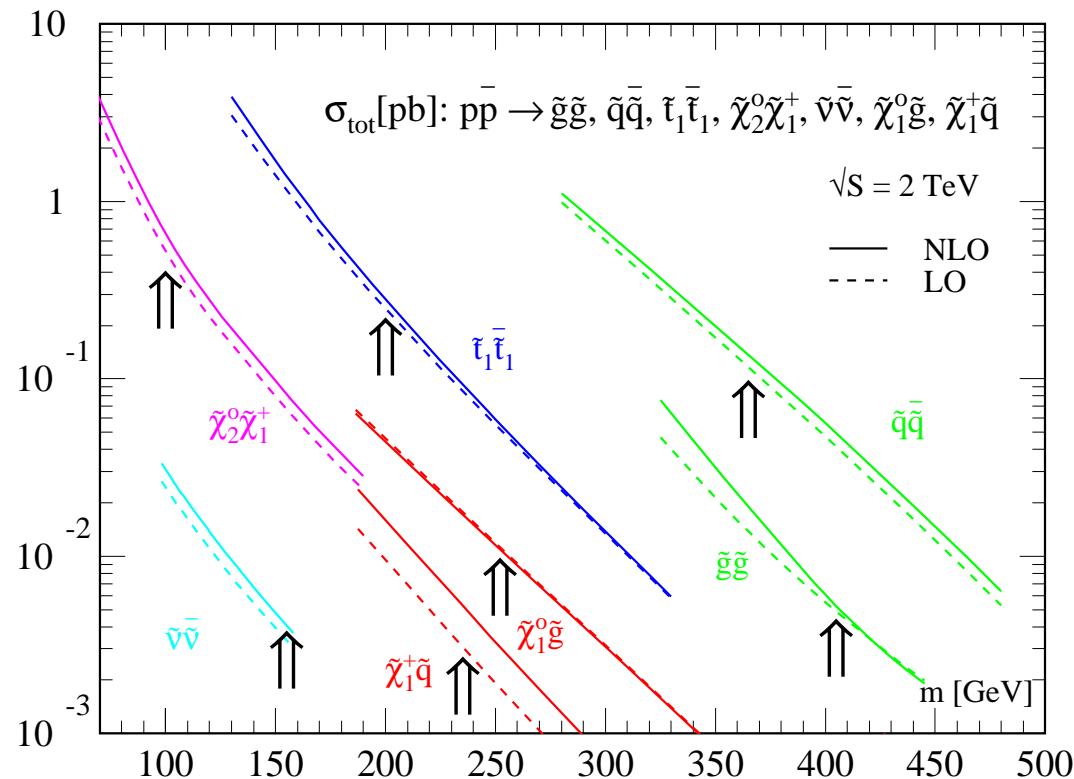


Supersymmetry at the LHC

- (1) **possible discovery** — signals for new physics, exclusion of parameter space
- (2) **measurements** — masses, cross sections, decays
- (3) **parameter studies** — MSSM Lagrangean, SUSY breaking

SUSY signals included

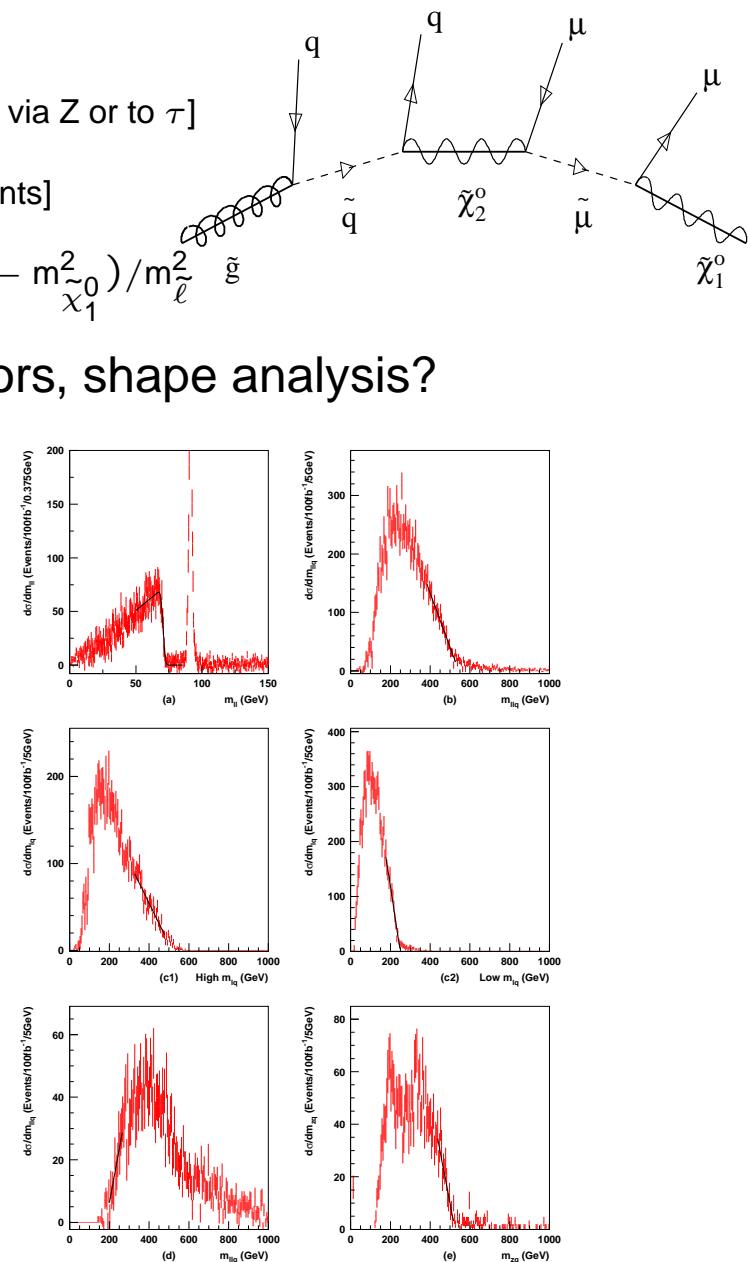
- QCD coupling $g\tilde{q}\tilde{q}$, $q\tilde{q}\tilde{q}$, $g\tilde{g}\tilde{g}$
- jets and \not{E}_T : $p\bar{p} \rightarrow \tilde{q}\tilde{q}^*, \tilde{g}\tilde{g}, \tilde{q}\tilde{g}$
- funny tops: $p\bar{p} \rightarrow \tilde{t}_1\tilde{t}_1^*$
- like sign dileptons: $p\bar{p} \rightarrow \tilde{g}\tilde{g}$
 $[\tilde{g} \rightarrow \tilde{u}\bar{u} \rightarrow \tilde{\chi}_1^+ d\bar{u} \text{ or } \tilde{g} \rightarrow \tilde{u}^* u \rightarrow \tilde{\chi}_1^- \bar{d}u]$
- tri-leptons: $p\bar{p} \rightarrow \tilde{\chi}_2^0\tilde{\chi}_1^-$
 $[\tilde{\chi}_2^0 \rightarrow \tilde{\ell}\bar{\ell} \rightarrow \tilde{\chi}_1^0\ell\bar{\ell}; \tilde{\chi}_1^- \rightarrow \tilde{\chi}_1^0\ell\bar{\nu}]$



SUPERSYMMETRY AT LHC: 2

Spectra from cascade decays

- decay $\tilde{g} \rightarrow \tilde{q}\bar{q} \rightarrow \tilde{\chi}_2^0 q\bar{q} \rightarrow \mu^+ \mu^- q\bar{q} \tilde{\chi}_1^0$ [better not via Z or to τ]
 - cross sections some 100 pb [more than 3×10^5 events]
 - thresholds & edges classical $m_{\ell\ell}^2 < (m_{\tilde{\chi}_2^0}^2 - m_\ell^2)(m_\ell^2 - m_{\tilde{\chi}_1^0}^2)/m_{\tilde{g}}^2$
 - detector resolution, calibration, systematic errors, shape analysis?
 - cross sections as additional input?
- ⇒ \tilde{q}_L cascade reconstruction great for SPS1a

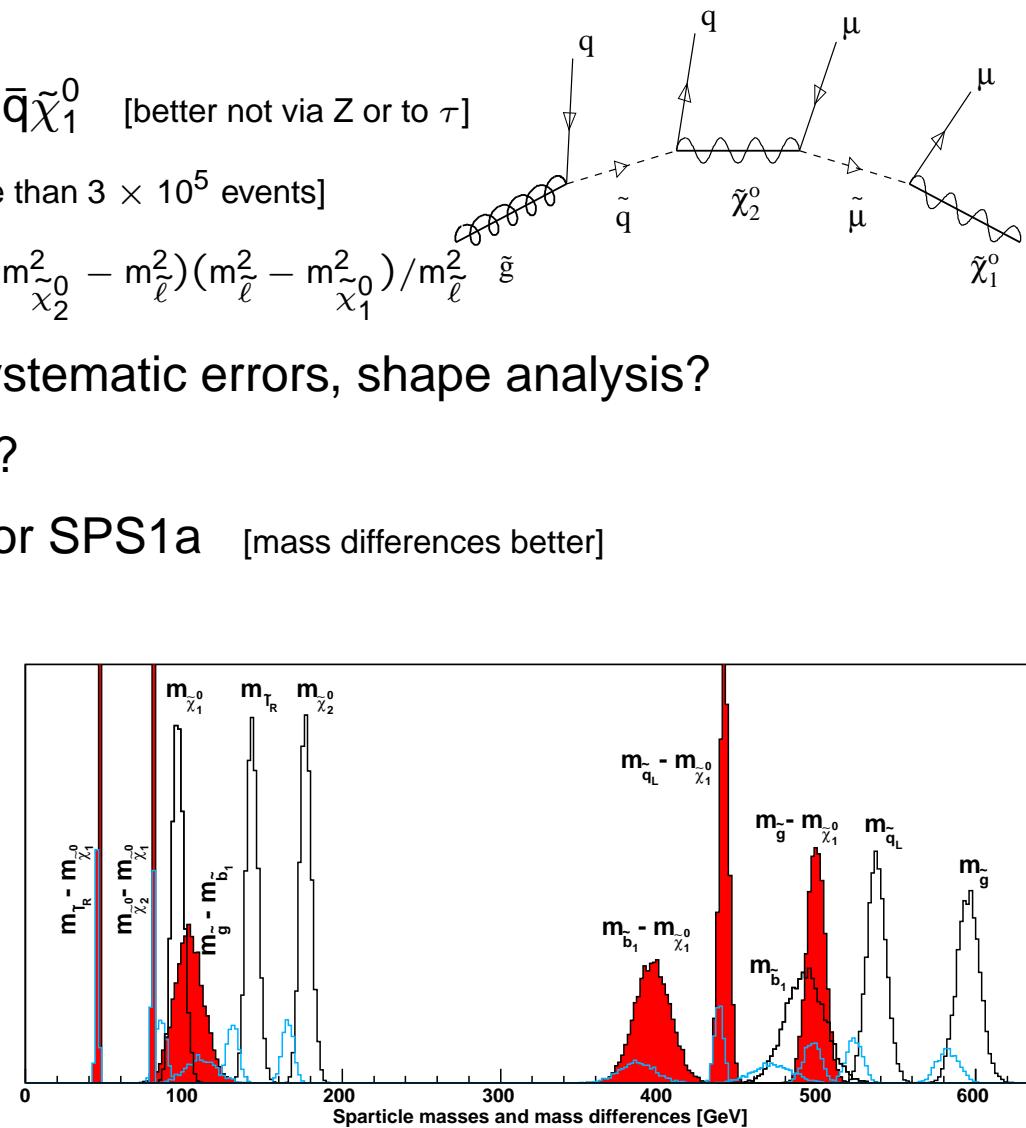


Spectra from cascade decays

- decay $\tilde{g} \rightarrow \tilde{q}\bar{q} \rightarrow \tilde{\chi}_2^0 q\bar{q} \rightarrow \mu^+ \mu^- q\bar{q} \tilde{\chi}_1^0$ [better not via Z or to τ]
 - cross sections some 100 pb [more than 3×10^5 events]
 - thresholds & edges classical $m_{\ell\ell}^2 < (m_{\tilde{\chi}_2^0}^2 - m_\ell^2)(m_\ell^2 - m_{\tilde{\chi}_1^0}^2)/m_{\tilde{g}}^2$
 - detector resolution, calibration, systematic errors, shape analysis?
 - cross sections as additional input?
- ⇒ \tilde{q}_L cascade reconstruction great for SPS1a [mass differences better]

Gluino mass

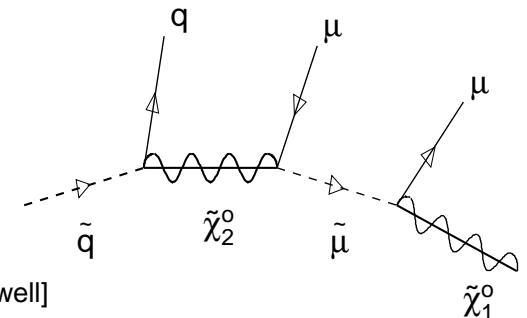
- now four jets instead of two
 - \tilde{b}_L instead, all jets b-tagged
- ⇒ gluino mass to $\sim 1\%$
statistical error dominant



How to make sure it is SUSY

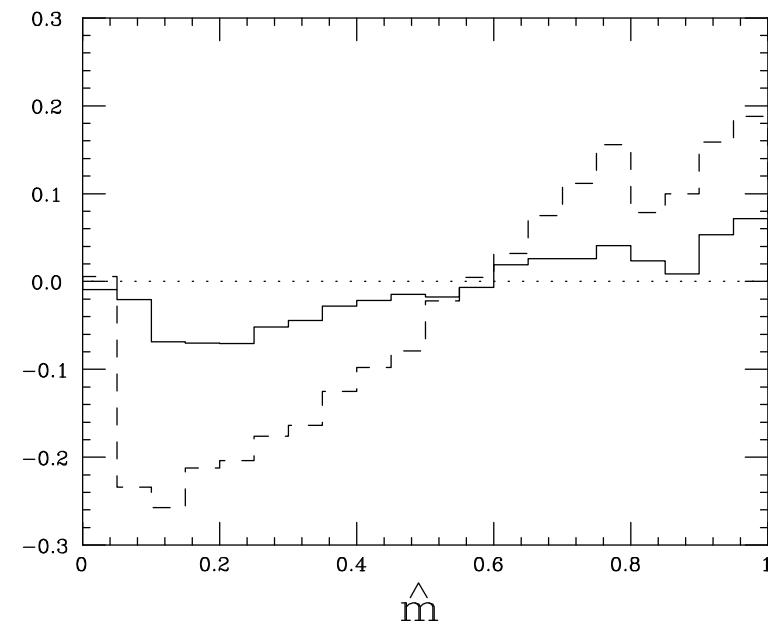
- assume neutralino is found in cascades
- compare with a model where gluino is a boson,...
- straw man: universal extra dimensions

[mass spectra degenerate —ignore this information; cross section factor 10 larger —ignore this as well]



Slepton cascade

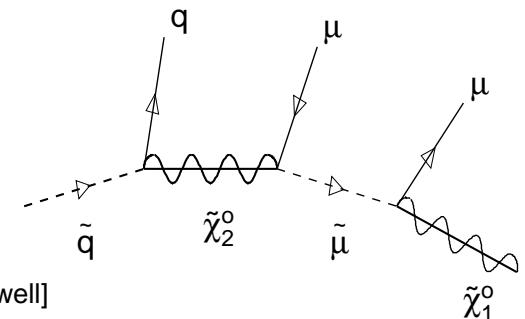
- decay chain $\tilde{\chi}_2^0 \rightarrow \ell\bar{\ell}^* \rightarrow \ell\bar{\ell}\tilde{\chi}_1^0$
- compare with first excited Z and ℓ
- typically largest $pp \rightarrow \tilde{q}\tilde{q}$ [$\tilde{q}/\tilde{q} \sim 2$]
- ⇒ $\hat{m} = m_{j\ell}/m_{j\ell}^{\max}$ most promising
 $\mathcal{A} = [\sigma(j\ell^+) - \sigma(j\ell^-)]/[\sigma(j\ell^+) + \sigma(j\ell^-)]$
- assume hierarchical SPS1a spectrum
[dashed SUSY, solid UED]
- ⇒ we can tell spin in cascade



How to make sure it is SUSY

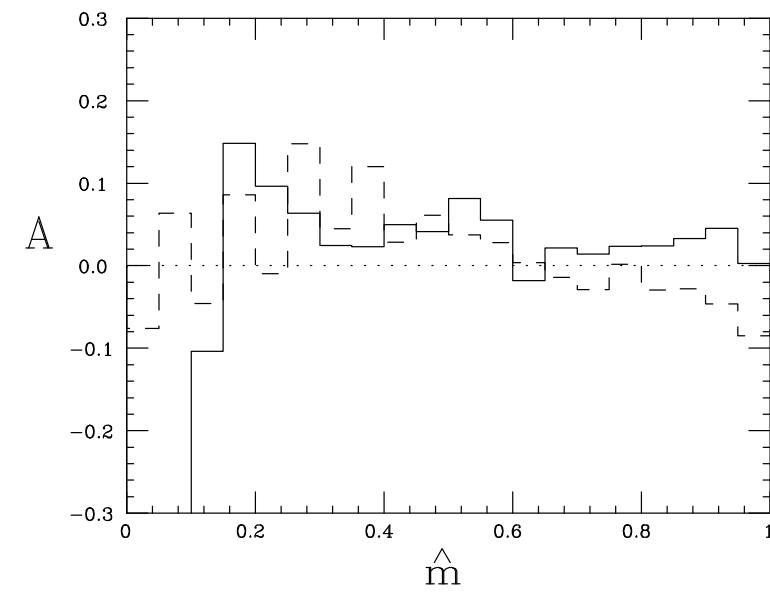
- assume neutralino is found in cascades
- compare with a model where gluino is a boson,...
- straw man: universal extra dimensions

[mass spectra degenerate —ignore this information; cross section factor 10 larger —ignore this as well]



Slepton cascade

- decay chain $\tilde{\chi}_2^0 \rightarrow \ell\bar{\ell}^* \rightarrow \ell\bar{\ell}\tilde{\chi}_1^0$
- compare with first excited Z and ℓ
- typically largest $pp \rightarrow \tilde{q}\tilde{q}$ [$\tilde{q}/\tilde{q} \sim 2$]
- ⇒ $\hat{m} = m_{j\ell}/m_{j\ell}^{\max}$ most promising
 $A = [\sigma(j\ell^+) - \sigma(j\ell^-)]/[\sigma(j\ell^+) + \sigma(j\ell^-)]$
- assume non-hierarchical UED spectrum
[dashed SUSY, solid UED]
- ⇒ we can tell spin in cascade



SUPERSYMMETRY AT THE ILC: 1

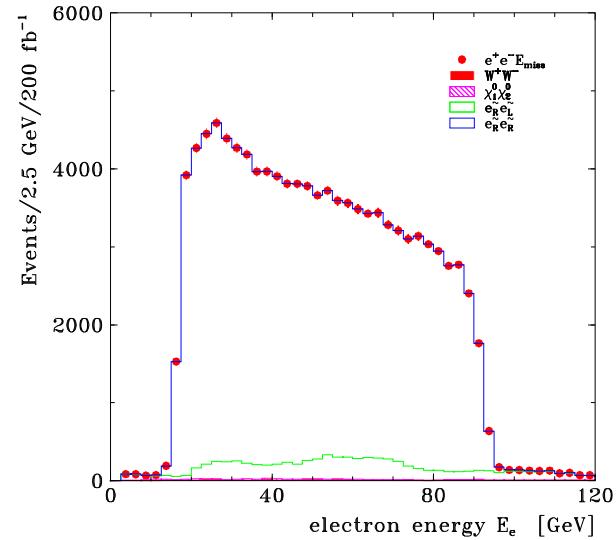
All information in the kinematics

- very high luminosity: 1 ab^{-1}
- excellent mass resolution to reject backgrounds
- mass determination from energy spectrum

$$m_{\tilde{\ell}} = \sqrt{s} \sqrt{E_+ E_-} / (E_+ + E_-)$$

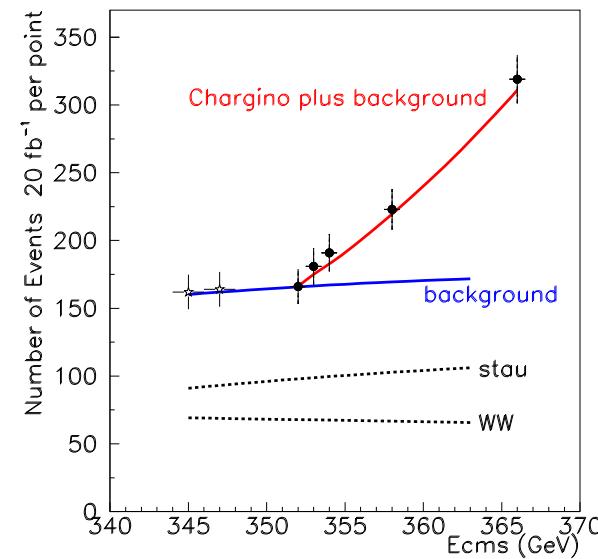
$$m_{\tilde{\chi}_1^0} = m_{\tilde{\ell}} \sqrt{1 - 2(E_+ + E_-) / \sqrt{s}}$$

- higher-order/off-shell corrections necessary
- ⇒ **mass determination to better than 0.2%**



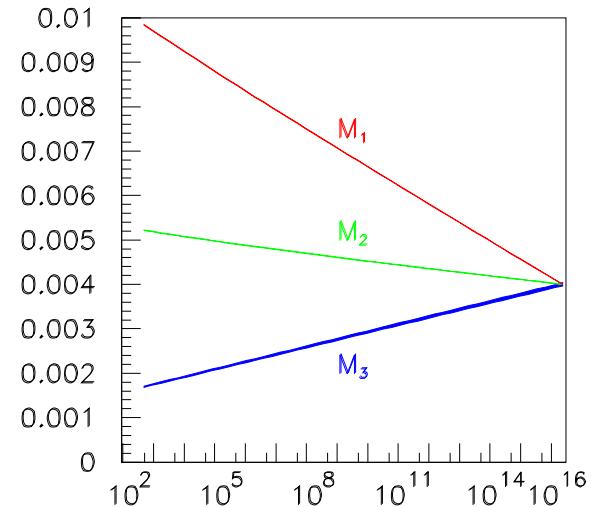
Similarly good: threshold scan

- required luminosity around 100 fb^{-1} per particle
 - theoretical threshold description: what mass?
 - necessary for example for $\tilde{\chi}_1^0 \nu$ in decay
- ⇒ **mass determination to better than 0.5%**



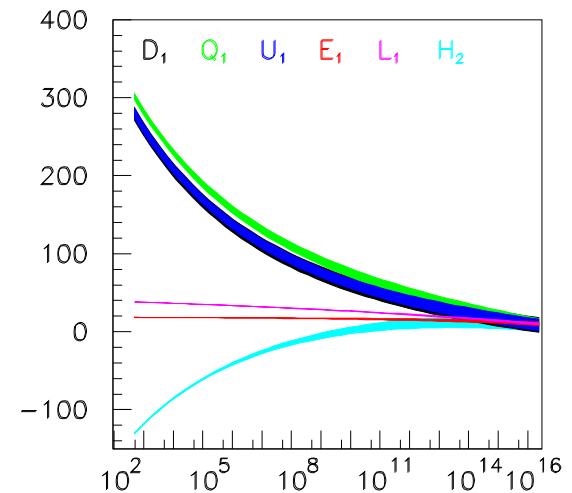
Renormalization group bottom-up

- high-scale parameters without assumptions?
- all parameters in beta functions necessary
neutralino/chargino sector from ILC (NLO)
gluino mass from LHC
- gaugino mass unification ‘SUGRA light’



Slepton masses more involved

- again test of unified masses
- quark flavor crucial input
- slepton mass running flat
squark mass determination weak
u-d-s-c flavor tagging in squark decays



SUPERSYMMETRIC PARAMETERS

SUSY parameters from observables

- parameters: weak-scale MSSM Lagrangean
- measurements: masses
branching fractions
cross sections
- errors: general correlation, statistics & systematics & theory
- problem in grid: huge phase space, local minimum?
problem in fit: domain walls, starting values, global minimum

Combination of methods

- (1) grid for closed subset
(2) fit of remaining parameters
(3) complete fit
 - more modern alternatives:
simulated annealing
Markov Chains
- ⇒ LHC+ILC with no assumptions

	LHC	ILC	LHC+ILC	SPS1a
$\tan\beta$	10.22 ± 9.1	10.26 ± 0.3	10.06 ± 0.2	10
M_1	102.45 ± 5.3	102.32 ± 0.1	102.23 ± 0.1	102.2
M_3	578.67 ± 15	fix 500	588.05 ± 11	589.4
$M_{\tilde{\tau}_L}$	fix 500	197.68 ± 1.2	199.25 ± 1.1	197.8
$M_{\tilde{\tau}_R}$	129.03 ± 6.9	135.66 ± 0.3	133.35 ± 0.6	135.5
$M_{\tilde{\mu}_L}$	198.7 ± 5.1	198.7 ± 0.5	198.7 ± 0.5	198.7
$M_{\tilde{q}_3_L}$	498.3 ± 110	497.6 ± 4.4	521.9 ± 39	501.3
$M_{\tilde{t}_R}$	fix 500	420 ± 2.1	411.73 ± 12	420.2
$M_{\tilde{b}_R}$	522.26 ± 113	fix 500	504.35 ± 61	525.6
A_τ	fix 0	-202.4 ± 89.5	352.1 ± 171	-253.5
A_t	-507.8 ± 91	-501.95 ± 2.7	-505.24 ± 3.3	-504.9
A_b	-784.7 ± 35603	fix 0	-977 ± 12467	-799.4

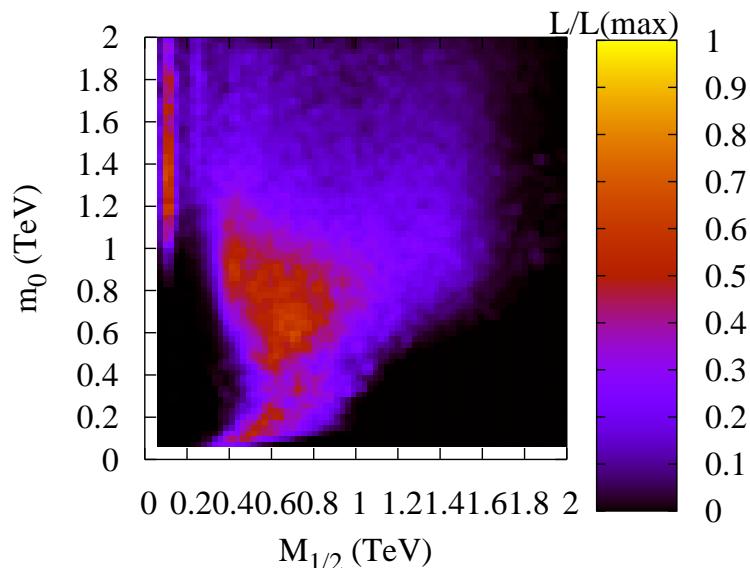
SUPERSYMMETRIC PARAMETERS

SUSY parameters from observables

- parameters: weak-scale MSSM Lagrangean
- measurements: masses or edges
branching fractions
cross sections
- errors: general correlation, statistics & systematics & theory
- problem in grid: huge phase space, local minimum?
problem in fit: domain walls, starting values, global minimum?

First go at problem

- ask a friend who knows how SUSY is broken
- ⇒ mSUGRA
- fit m_0 , $m_{1/2}$, A_0 , $\tan \beta$, $\text{sign}(\mu)$
 - no problem, include indirect constraints
- ⇒ who would bet a month's salary on mSUGRA?
- admittedly: **mSUGRA a very useful testing ground for methods**



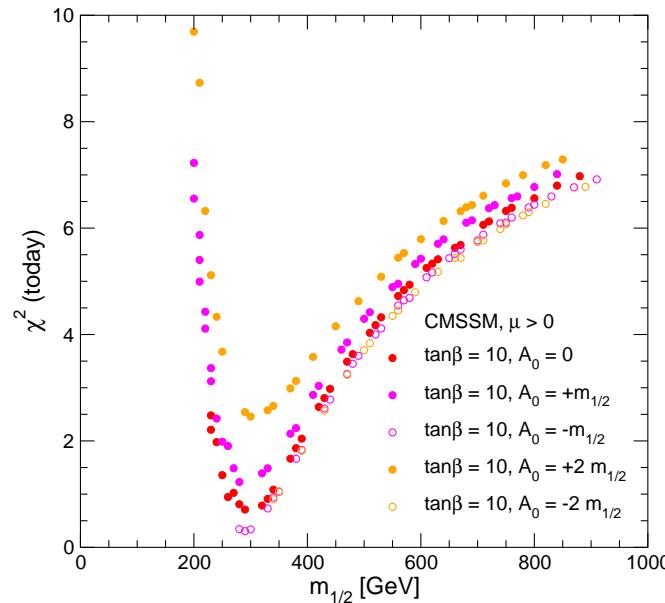
SUPERSYMMETRIC PARAMETERS

SUSY parameters from observables

- parameters: weak-scale MSSM Lagrangean
- measurements: masses or edges
branching fractions
cross sections
- errors: general correlation, statistics & systematics & theory
- problem in grid: huge phase space, local minimum?
problem in fit: domain walls, starting values, global minimum?

First go at problem

- ask a friend who knows how SUSY is broken
- ⇒ mSUGRA
- fit $m_0, m_{1/2}, A_0, \tan \beta, \text{sign}(\mu)$
 - no problem, include indirect constraints
- ⇒ who would bet a month's salary on mSUGRA?
- admittedly: **mSUGRA a very useful testing ground for methods**



OUTLOOK

Supersymmetry a good solution to hierarchy problem

- dark matter motivation getting stronger
 - no hint of SUSY-flavor physics
 - Tevatron searches promising put tough
 - LHC searches much more powerful
 - ILC analyses perfect tool, luminosity the issue
- ⇒ Super-cool times ahead