

# UV-Completing the Standard Model at the LHC

Tilman Plehn

University of Edinburgh

University of Manchester, 3/2008

# Outline

TeV scale in the LHC era

Masses from cascades

Underlying parameters

Spins from cascades

Spins from jets

UV-complete extra dimensions

# Standard-Model effective theory

## Remember the Standard Model

- gauge theory  $SU(3) \times SU(2) \times U(1)$
  - massless  $SU(3)$  and  $U(1)$  gauge bosons  
massive  $SU(2)$  gauge bosons [spontaneous symmetry breaking]
  - massive Dirac fermions [via Yukawas]
  - perturbatively renormalizable Lagrangian [no  $1/M$  terms]
  - one missing piece: Higgs [fundamental? minimal? mass?]
- ⇒ truly fundamental theory

## How complete experimentally?

- dark matter? [solid evidence! — for weak-scale new physics?]
  - $(g - 2)_\mu$ ? [possible evidence for weak-scale new physics, review by Dominik]
  - quark mixing — flavor physics? [new operators above  $10^4$  GeV?]
  - neutrino masses and mixing? [see-saw at  $10^{11}$  GeV?]
  - matter-antimatter asymmetry? [universe mostly matter]
  - gauge-coupling unification? [almost perfect, but proton stable]
  - gravity? [mostly negligible but perturbatively non-renormalizable]
- ⇒ cut-off scale unavoidable: SM effective theory

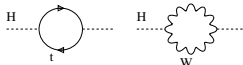
# Standard-Model effective theory

## Remember the Standard Model

- gauge theory  $SU(3) \times SU(2) \times U(1)$
- massless  $SU(3)$  and  $U(1)$  gauge bosons  
massive  $SU(2)$  gauge bosons [spontaneous symmetry breaking]
- massive Dirac fermions [via Yukawas]
- perturbatively renormalizable Lagrangian [no  $1/M$  terms]
- one missing piece: Higgs [fundamental? minimal? mass?]

⇒ truly fundamental theory

## How consistent theoretically?



- problem of light Higgs: mass driven to cutoff of effective Standard Model  
$$\delta m_H^2 \propto g^2(2m_W^2 + m_Z^2 + m_t^2 - 4m_t^2) \Lambda^2$$
- cancelled by finely tuned counter term? [ugly, against spirit of symmetries]
- why fundamental Higgs at all, if such poor high-energy behavior?
- tied in with new physics at TeV scale:
  - supersymmetry [my favorite]
  - extra dimensions [cool idea]
  - little Higgs [old idea, now working]
  - no fundamental Higgs... [not pretty]

⇒ TeV scale: beautiful ideas — complicated realistic models

# Effective Standard Model in the LHC era

## UV completions from the LHC [Uli Baur's rule: 'there is always new physics at higher scales']

- find light Higgs?
- find new physics stabilizing Higgs mass?
- see dark-matter candidate?

## Particle theory and new physics

- model-independent analyses not helpful to understand TeV scale
- early running data probably with little information [just a pheno game]
- testable TeV-scale models [e.g. Higgs sector vs. underlying theory?]
- link to other observations [DM+Tevatron: Hooper, TP, Valinotto]
- reconstruction of Lagrangian

# Effective Standard Model in the LHC era

## UV completions from the LHC [Uli Baur's rule: 'there is always new physics at higher scales']

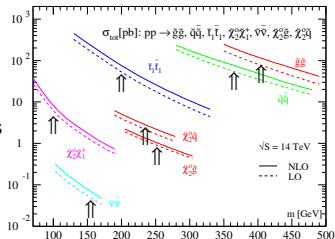
- find light Higgs?
- find new physics stabilizing Higgs mass?
- see dark-matter candidate?

## Particle theory and new physics

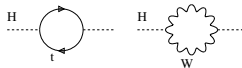
- model-independent analyses not helpful to understand TeV scale
- early running data probably with little information [just a pheno game]
- testable TeV-scale models [e.g. Higgs sector vs. underlying theory?]
- link to other observations [DM+Tevatron: Hooper, TP, Valinotto]
- reconstruction of Lagrangian

## Special about LHC

- beyond inclusive searches [that was Tevatron]  
millions of new strongly interacting particles
- ⇒ (1) aim at underlying theory  
(2) try to survive QCD [ask your guys]



# TeV-scale supersymmetry



## Supersymmetry

- give each Standard-Model particle a partner [with different spin, including strong interactions]
  - symmetry obviously broken by masses [soft breaking, mechanism unknown]
  - assume dark matter, stable lightest partner [R parity]
- ⇒ **measure BSM spectrum with missing energy at LHC**

## LHC searches: MSSM

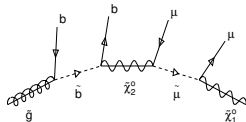
- SUSY-Higgs alone interesting...  
...but not conclusive  
...and another talk [ask Apostolos]
- ⇒ **list of SUSY partners**

		spin	d.o.f.	
fermion	$f_L, f_R$	1/2	1+1	
→ sfermion	$\tilde{f}_L, \tilde{f}_R$	0	1+1	
gluon	$G_\mu$	1	n-2	
→ gluino	$\tilde{g}$	1/2	2	Majorana
gauge bosons	$\gamma, Z$	1	2+3	
Higgs bosons	$h^0, H^0, A^0$	0	3	
→ neutralinos	$\tilde{\chi}_i^0$	1/2	4 · 2	LSP
gauge bosons	$W^\pm$	1	2 · 3	
Higgs bosons	$H^\pm$	0	2	
→ charginos	$\tilde{\chi}_i^\pm$	1/2	2 · 4	

# Masses from cascades

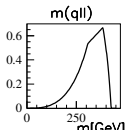
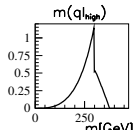
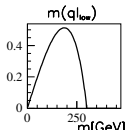
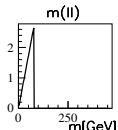
## Cascade decays [Atlas-TDR, Cambridge]

- if new particles strongly interacting and LSP weakly interacting
- like Tevatron: jets + missing energy
- easiest: cascade kinematics [ $10^7 \dots 10^8$  events, rates tough because of QCD]
- thresholds & edges [RAL school exercise]



$$0 < m_{\mu\mu}^2 < \frac{m_{\tilde{\chi}_2^0}^2 - m_{\tilde{\ell}}^2}{m_{\tilde{\ell}}} \frac{m_{\tilde{\ell}}^2 - m_{\tilde{\chi}_1^0}^2}{m_{\tilde{\ell}}}$$

⇒ **new-physics mass spectrum from cascade kinematics**

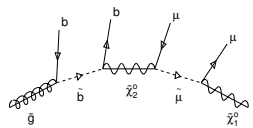




# Masses from cascades

## Cascade decays [Atlas-TDR, Cambridge]

- if new particles strongly interacting and LSP weakly interacting
- like Tevatron: jets + missing energy
- easiest: cascade kinematics [10<sup>7</sup> . . . 10<sup>8</sup> events, rates tough because of QCD]
- thresholds & edges [RAL school exercise]



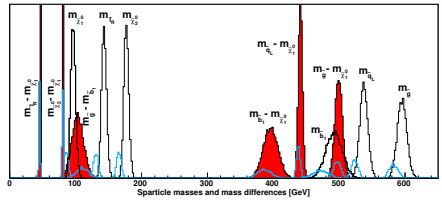
$$0 < m_{\mu\mu}^2 < \frac{m_{\tilde{\chi}_2^0}^2 - m_{\tilde{\ell}}^2}{m_{\tilde{\ell}}} \frac{m_{\tilde{\ell}}^2 - m_{\tilde{\chi}_1^0}^2}{m_{\tilde{\ell}}}$$

⇒ new-physics mass spectrum from cascade kinematics

## Glino decay [Gjelsten, Miller, Osland, Raklev]

- no problem: additional jets [Rainwater, TP, Skands; Alwall, Wacker,...]
- no problem: off-shell effects [Catpiss collaboration; Kauer & Rainwater]
- all decay jets b quarks [otherwise dead by QCD]
- limited by jet energy scale
- alternative methods?

⇒ but why physical masses?



# Underlying parameters

## From kinematics to weak-scale parameters [Fittino; SFitter: Lafaye, TP, Rauch, Zerwas]

- parameters: weak-scale Lagrangian
- measurements: kinematic endpoints, branching ratios, rates [Prospino2]  
 $B$  decays,  $(g - 2)_\mu$ , dark matter, e-w precision data...
- errors: general correlation, statistics & systematics & theory [flat theory errors!]
- problem in grid/fit: no local/global maximum!  
**problem in physics: secondary maxima!**

## Probability maps of new physics [Baltz,...; Roszkowski,...; Allanach,...; SFitter]

- fully exclusive likelihood map  $p(d|m)$  over  $m$  [hard part]
- LHC problem: remove directions
- Bayesian:  $p(m|d) \sim p(d|m) p(m)$  with theorists' bias  $p(m)$  [cosmology, BSM]  
frequentist: best-fitting point  $\max_m p(d|m)$  [flavor]
- LHC era: (1) compute high-dimensional map  $p(d|m)$   
(2) find and rank local maxima in  $p(d|m)$   
(3) Bayesian-frequentist dance to reduce dimensions

# Underlying parameters

## From kinematics to weak-scale parameters [Fittino; SFitter: Lafaye, TP, Rauch, Zerwas]

- parameters: weak-scale Lagrangian
- measurements: kinematic endpoints, branching ratios, rates [Prospino2]  
*B* decays,  $(g - 2)_\mu$ , dark matter, e-w precision data...
- errors: general correlation, statistics & systematics & theory [flat theory errors!]
- problem in grid/fit: no local/global maximum!  
**problem in physics: secondary maxima!**

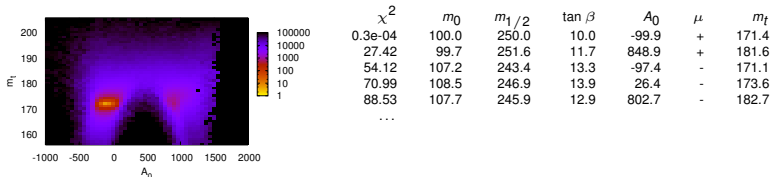
## Toy model: MSUGRA map from LHC [LHC endpoints with free $y_f$ ]

- weighted Markov chains: several times faster [similar to: Ferrenberg & Swendsen]

$$P_{\text{bin}}(p \neq 0) = \frac{N}{\sum_{i=1}^N 1/p}$$

- SFitter output #1: fully exclusive likelihood map
- SFitter output #2: ranked list of local maxima

⇒ **strong correlations even in MSUGRA**

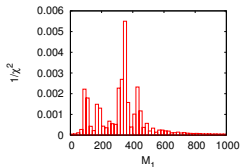
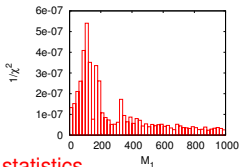


# Underlying parameters

## MSSM map at LHC

- shifting from 6D to 19D parameter space [killing grids, Minuit, laptop-style fits...]
- SFitter outputs still the same, but best points degenerate
- e.g. 3 neutralinos observed

Bayesian pdf noisy  
profile likelihood no pdf



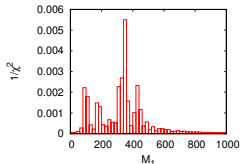
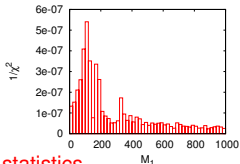
⇒ no golden approach to BSM statistics

# Underlying parameters

## MSSM map at LHC

- shifting from 6D to 19D parameter space [killing grids, Minuit, laptop-style fits...]
- SFitter outputs still the same, but best points degenerate
- e.g. 3 neutralinos observed

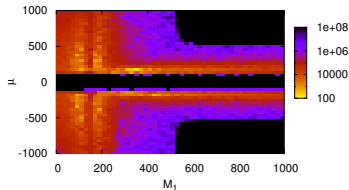
Bayesian pdf noisy  
profile likelihood no pdf



⇒ no golden approach to BSM statistics

## MSSM map in LHC era [SFitter+friends]

- LHC:  $\text{sign}(\mu)$  unknown and  $\tan\beta$  extraction believe-based

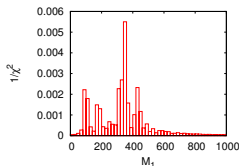
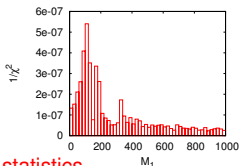


# Underlying parameters

## MSSM map at LHC

- shifting from 6D to 19D parameter space [killing grids, Minuit, laptop-style fits...]
- SFitter outputs still the same, but best points degenerate
- e.g. 3 neutralinos observed

Bayesian pdf noisy  
profile likelihood no pdf



⇒ no golden approach to BSM statistics

## MSSM map in LHC era [Sfitter+friends]

- LHC:  $\text{sign}(\mu)$  unknown and  $\tan \beta$  extraction believe-based
- $\tan \beta$  and  $\text{sign}(\mu)$  from  $(g-2)_\mu$  or  $B_s \rightarrow \mu\mu$  [Les Houches 2007: Alexander et al]
- including  $(g-2)$  very promising

	LHC	LHC $\otimes (g-2)$	SPS1a
$\tan \beta$	$10.0 \pm 4.5$	$10.3 \pm 2.0$	10.0
$M_1$	$102.1 \pm 7.8$	$102.7 \pm 5.9$	103.1
$M_2$	$193.3 \pm 7.8$	$193.2 \pm 5.8$	192.9
$M_3$	$577.2 \pm 14.5$	$578.2 \pm 12.1$	577.9
$M_{\tilde{\mu}_L}$	$193.2 \pm 8.8$	$194.0 \pm 6.8$	194.4
$M_{\tilde{\mu}_R}$	$135.0 \pm 8.3$	$135.6 \pm 6.3$	135.8
$M_{\tilde{q}_L}$	$524.6 \pm 14.5$	$525.5 \pm 10.6$	526.6
$M_{\tilde{q}_R}$	$507.3 \pm 17.5$	$507.6 \pm 15.8$	508.1
$\mu$	$350.5 \pm 14.5$	$352.5 \pm 10.8$	353.7

# Underlying parameters

## MSSM map at LHC

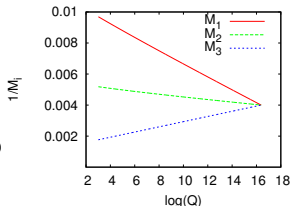
- shifting from 6D to 19D parameter space [killing grids, Minuit, laptop-style fits...]
  - SFitter outputs still the same, but best points degenerate
- ⇒ **no golden approach to BSM statistics**

## MSSM map in LHC era [SFitter+friends]

- LHC:  $\text{sign}(\mu)$  unknown and  $\tan\beta$  extraction believe-based
- $\tan\beta$  and  $\text{sign}(\mu)$  from  $(g-2)_\mu$  or  $B_s \rightarrow \mu\mu$  [Les Houches 2007: Alexander et al]
- including  $(g-2)$  very promising

## Fundamental theory [SFitter + Kneur]

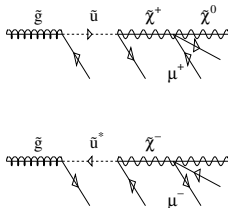
- SUSY breaking?
  - unification, GUT?
  - scale-invariant sum rules? [Cohen, Schmalz]
  - new & crucial: renormalization group bottom-up
- ⇒ **LHC sensitive to UV models**



# Spins from cascades

## Supersymmetry at Tevatron/LHC: Majorana gluinos [Barger,...; Barnett,...; Baer,...]

- LHC: first jet ( $q$  or  $\bar{q}$ ) fixes lepton charge
  - same-sign dileptons in 1/2 of events
  - similar:  $t$ -channel gluino in  $pp \rightarrow \tilde{q}\tilde{q}$
- $\Rightarrow$  gluino = like-sign dileptons in SUSY sample





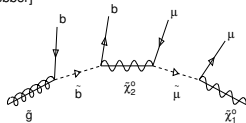
# Spins from cascades

## Supersymmetry at Tevatron/LHC: Majorana gluinos [Barger,...; Barnett,...; Baer,...]

- LHC: first jet ( $q$  or  $\bar{q}$ ) fixes lepton charge
  - same-sign dileptons in 1/2 of events
  - similar:  $t$ -channel gluino in  $pp \rightarrow \tilde{q}\tilde{q}$
- ⇒ gluino = like-sign dileptons in SUSY sample

## Loop hole: gluino is Majorana if fermion [Alves, Eboli, TP]

- start with mass-measurement cascade
- now: physics between the endpoints
- model-independent analysis difficult [Barr, Lester, Smillie, Webber]



# Spins from cascades

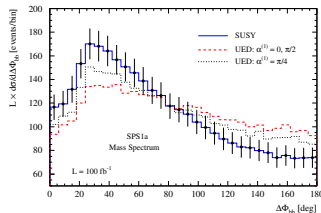
## Supersymmetry at Tevatron/LHC: Majorana gluinos [Barger,...; Barnett,...; Baer,...]

- LHC: first jet ( $q$  or  $\bar{q}$ ) fixes lepton charge
  - same-sign dileptons in 1/2 of events
  - similar:  $t$ -channel gluino in  $pp \rightarrow \tilde{q}\tilde{q}$
- ⇒ **gluino = like-sign dileptons in SUSY sample**

## Loop hole: gluino is Majorana if fermion [Alves, Eboli, TP]

- start with mass-measurement cascade
- now: physics between the endpoints
- model-independent analysis difficult [Barr, Lester, Smillie, Webber]

- UV-completion hypotheses tests instead
  - 'gluino' a boson: universal extraD
  - compare SUSY vs. KK  $g, b, Z, l, \gamma$
  - simple distributions  $\Delta\phi_{bb}$  [3-body decays: Csaki,...]
- ⇒ **gluino = fermion with like-sign dileptons**



# Spins from cascades

## Supersymmetry at Tevatron/LHC: Majorana gluinos [Barger,...; Barnett,...; Baer,...]

- LHC: first jet ( $q$  or  $\bar{q}$ ) fixes lepton charge
  - same-sign dileptons in 1/2 of events
  - similar:  $t$ -channel gluino in  $pp \rightarrow \tilde{q}\tilde{q}$
- ⇒ **gluino = like-sign dileptons in SUSY sample**

## Loop hole: gluino is Majorana if fermion [Alves, Eboli, TP]

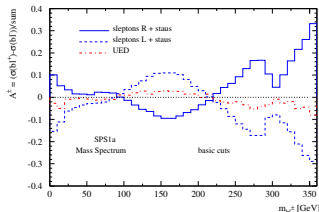
- start with mass-measurement cascade
- now: physics between the endpoints
- model-independent analysis difficult [Barr, Lester, Smillie, Webber]

- UV-completion hypotheses tests instead
- 'gluino' a boson: universal extraD
- compare SUSY vs. KK  $g, b, Z, l, \gamma$
- simple distributions  $\Delta\phi_{bb}$  [3-body decays: Csaki,...]

⇒ **gluino = fermion with like-sign dileptons**

- sensitive to model's details

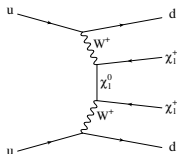
⇒ **LHC requiring testable TeV-scale models**



## Spins from jets

### Beyond cascades: spin of LSP [Alwall, TP, Rainwater]

- hypotheses: like–sign charginos (Majorana neutralino)  
like–sign scalars (stable scalars)  
like–sign vector bosons (little–Higgs inspired)
- chargino decay/kinematics not used
- jet correlations the key



# Weak boson fusion and unitarity

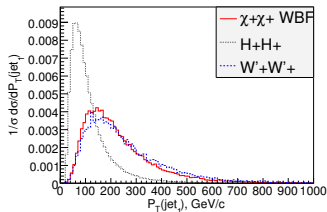
## Like-sign scalars instead of fermions

- charged Higgs in 2HDM [TP, Rainwater, Zeppenfeld; Buszello, Marquard, v.d.Bij]
- $W$  radiated off quarks [Goldstone coupling to Higgs]

$$P_T(x, p_T) \sim \frac{1 + (1-x)^2}{2x} \frac{1}{p_T^2}$$

⇒ scalars with softer  $p_{T,j}$

$$P_L(x, p_T) \sim \frac{(1-x)^2}{x} \frac{m_W^2}{p_T^4}$$



# Weak boson fusion and unitarity

## Like-sign scalars instead of fermions

- charged Higgs in 2HDM [TP, Rainwater, Zeppenfeld; Buszello, Marquard, v.d.Bij]
- $W$  radiated off quarks [Goldstone coupling to Higgs]

$$P_T(x, p_T) \sim \frac{1 + (1-x)^2}{2x} \frac{1}{p_T^2}$$

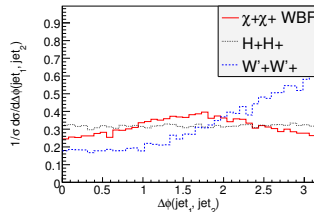
$$P_L(x, p_T) \sim \frac{(1-x)^2}{x} \frac{m_W^2}{p_T^4}$$

⇒ scalars with softer  $p_{T,j}$

## Like-sign vectors instead of fermions

- little-Higgs inspired [ $T$ -type parity]
- start with copy of SM, heavy  $W'$ ,  $Z'$ ,  $H'$ ,  $f'$  [ $H'$  necessary for unitarity, but irrelevant at LHC]
- Lorentz structure reflected in angle between jets

⇒ vectors with peaked  $\Delta\phi_{jj}$



# Weak boson fusion and unitarity

## Like-sign scalars instead of fermions

- charged Higgs in 2HDM [TP, Rainwater, Zeppenfeld; Buszello, Marquard, v.d.Bij]
- $W$  radiated off quarks [Goldstone coupling to Higgs]

$$P_T(x, p_T) \sim \frac{1 + (1-x)^2}{2x} \frac{1}{p_T^2}$$

$$P_L(x, p_T) \sim \frac{(1-x)^2}{x} \frac{m_W^2}{p_T^4}$$

⇒ scalars with softer  $p_{T,j}$

## Like-sign vectors instead of fermions

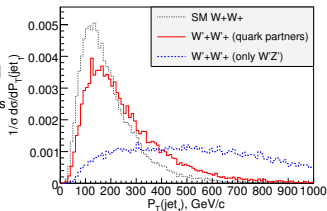
- little-Higgs inspired [ $T$ -type parity]
- start with copy of SM, heavy  $W'$ ,  $Z'$ ,  $H'$ ,  $f'$  [ $H'$  necessary for unitarity, but irrelevant at LHC]
- Lorentz structure reflected in angle between jets

⇒ vectors with peaked  $\Delta\phi_{jj}$

## Heavy fermions in little-Higgs models

- part of unitary UV completion [Englert, Zeppenfeld]
- huge effects on distributions [strongly interacting  $W$ s]

⇒ LHC requiring testable TeV-scale models



# UV-complete extra dimensions

## Elegant solution to hierarchy problem [Arkani-Hamed, Dimopoulos, Dvali]

- highest scale: Planck scale  $G_N \sim 1/M_{\text{Planck}}^2$  [ $M_{\text{Planck}} \sim 10^{19}$  GeV]
- Einstein–Hilbert action in  $4 + n$  dimensions [on torus — periodic boundaries]

$$\int d^4 x \sqrt{|g|} M_{\text{Planck}}^2 R \rightarrow \int d^{4+n} x \sqrt{|g|} M_D^{2+n} R = (2\pi r)^n \int d^4 x \sqrt{|g|} M_D^{2+n} R$$

$$M_{\text{Planck}} = M_D (2\pi r M_D)^{n/2} \gg M_D \sim 1 \text{ TeV}$$

- to get numbers right:  $r = 10^{12}, 10^{-3}, \dots 10^{-11}$  m for  $n = 1, 2, \dots 6$

⇒ **fundamental Planck scale at TeV**

## Kaluza–Klein gravitons [Giudice, Ratazzi, Wells]

- periodic boundaries: Fourier–transform in extra dimensions [QCD massless]

$$(\square + m_k^2) G_{\mu\nu}^{(k)} = -\frac{T_{\mu\nu}}{M_{\text{Planck}}} \quad \delta m \sim \frac{1}{r} = 2\pi M_D \left( \frac{M_D}{M_{\text{Planck}}} \right)^{2/n} \lesssim 0.05 \text{ GeV}$$

- KK tower of single gravitons, each coupled as  $1/M_{\text{Planck}}$
- IR spectrum: constraints from supernova cooling
- UV spectrum: LHC effects at TeV scale [Giudice, Strumia, TP; cosmic rays]



# UV-complete extra dimensions

## Hope for collider searches [Giudice, Rattazzi, Wells; Han, Lykken, Zhang]

- real radiation of continuous KK tower

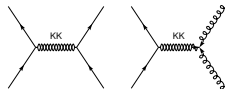
$$\sigma^{\text{tower}} \sim \sigma^{\text{graviton}} \int dm S_{n-1} m^{n-1} r^n = \sigma^{\text{graviton}} \int dm \frac{S_{n-1} m^{n-1}}{(2\pi M_D)^n} \left( \frac{M_{\text{Planck}}}{M_D} \right)^2$$

- higher-dimensional operator from virtual gravitons [UV dominated]

$$\mathcal{A}(s; m) = \frac{1}{M_{\text{Planck}}^2} T_{\mu\nu} T^{\mu\nu} \frac{1}{s - m^2} \rightarrow \frac{S_{n-1}}{2M_D^4} \left( \frac{\Lambda}{M_D} \right)^{n-2}$$

- KK tower coupling with  $1/M_D$  at LHC
- LHC rates (or reach) dependent on cut-off  $\Lambda$

⇒ **KK effective theory poor at LHC, UV completion needed**



# UV-complete extra dimensions

## Renormalization flow of gravity [Reuter; Litim; Wetterich;...]

- Einstein-Hilbert action with running  $G(\mu)$  and  $\Lambda(\mu)$
- dimensionless coupling  $g(\mu) = G(\mu) \mu^{2+n} = G_0 Z_G^{-1}(\mu) \mu^{2+n}$
- attractive finite UV fixed point [anomalous dimension:  $\eta = -\mu \partial_\mu \log Z_G \propto g$ ]

$$\mu \frac{\partial}{\partial \mu} g_*(\mu) = (2 + n + \eta(g_*)) g_*(\mu) = 0 \quad \text{for} \quad g_* \neq 0 \quad \eta(g_*) = -2 - n$$

- gravity asymptotically free in UV  $G(\mu) \sim g_* \mu^{-(2+n)}$

⇒ coupling weak enough for finite LHC predictions?

# UV-complete extra dimensions

## Renormalization flow of gravity [Reuter; Litim; Wetterich;...]

- Einstein-Hilbert action with running  $G(\mu)$  and  $\Lambda(\mu)$
- dimensionless coupling  $g(\mu) = G(\mu) \mu^{2+n} = G_0 Z_G^{-1}(\mu) \mu^{2+n}$
- attractive finite UV fixed point [anomalous dimension:  $\eta = -\mu \partial_\mu \log Z_G \propto g$ ]

$$\mu \frac{\partial}{\partial \mu} g_*(\mu) = (2 + n + \eta(g_*)) g_*(\mu) = 0 \quad \text{for} \quad g_* \neq 0 \quad \eta(g_*) = -2 - n$$

- gravity asymptotically free in UV  $G(\mu) \sim g_* \mu^{-(2+n)}$
- ⇒ coupling weak enough for finite LHC predictions?

## UV-safe gravity [Weinberg (1979)]

- gravity non-fundamental effective theory  $G \propto 1/M_D^{2+n}?$
- 't Hooft's perturbative renormalizability: finite number of counter terms
- Wilson's renormalizability: no unphysical UV divergences
- consistent theory beyond perturbation theory [no ghosts: Weinberg & Gomez]
- fixed point likely to  $\sqrt{|g|} R^8$  and including matter [no proof; not perturbative series]
- great idea for gravity — great for LHC

# UV-complete extra dimensions

## UV-completed graviton production [Litim & TP]

– form factor for  $G(\mu)$  [Hewett & Rizzo]  $\frac{1}{M_D^{2+n}} \rightarrow \frac{1}{M_D^{2+n}} \left[ 1 + \left( \frac{\sqrt{s}}{aM_D} \right)^{2+n} \right]^{-1}$

– alternative: changing anomalous dimension of graviton [QCD inspired]

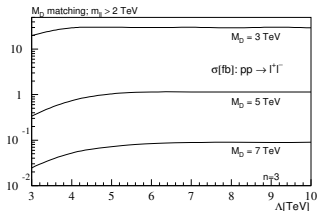
$$P(s, m) = \frac{1}{s + m^2} \rightarrow \frac{M_D^{n+2}}{(s + m^2)^{n/2+2}} \quad \text{around } m \sim M_D$$

– integration kernel after integration over  $m_{\text{KK}}$

$$\frac{1}{M_D^{2+n}} \int_0^\infty \frac{dm}{m^{n-1}} P(s, m) = \frac{1}{n-2} \frac{1}{M_D^4} \left( \frac{aM_D}{M_D} \right)^{n-2} \left[ 1 + \frac{n-2}{4} \right] \left[ 1 + \mathcal{O} \left( \frac{s}{M_D^2} \right) \right]$$

–  $\sqrt{s} > M_D$ : kernel only function of  $\sqrt{s}$  [matching at  $M_D$ ?, black-hole solutions?]

⇒ UV-safe gravity safe at LHC



# UV-complete extra dimensions

## UV-completed graviton production [Litim & TP]

– form factor for  $G(\mu)$  [Hewett & Rizzo]  $\frac{1}{M_D^{2+n}} \longrightarrow \frac{1}{M_D^{2+n}} \left[ 1 + \left( \frac{\sqrt{s}}{aM_D} \right)^{2+n} \right]^{-1}$

– alternative: changing anomalous dimension of graviton [QCD inspired]

$$P(s, m) = \frac{1}{s + m^2} \longrightarrow \frac{M_D^{n+2}}{(s + m^2)^{n/2+2}} \quad \text{around } m \sim M_D$$

– integration kernel after integration over  $m_{\text{KK}}$

$$\frac{1}{M_D^{2+n}} \int_0^\infty \frac{dm}{m^{n-1}} P(s, m) = \frac{1}{n-2} \frac{1}{M_D^4} \left( \frac{aM_D}{M_D} \right)^{n-2} \left[ 1 + \frac{n-2}{4} \right] \left[ 1 + \mathcal{O} \left( \frac{s}{M_D^2} \right) \right]$$

–  $\sqrt{s} > M_D$ : kernel only function of  $\sqrt{s}$  [matching at  $M_D$ ?, black-hole solutions?]

⇒ UV-safe gravity safe at LHC

## String theory as UV completion [e.g. Cullen, Perelstein, Peskin]

– Veneziano form factor

$$\frac{\Gamma(1 - \alpha' s) \Gamma(1 - \alpha' t)}{\Gamma(1 - \alpha' (s + t))} = \frac{\Gamma(1 - s/M_S^2) \Gamma(1 - t/M_S^2)}{\Gamma(1 - (s + t)/M_S^2)} = 1 - \frac{\pi^2}{6} \frac{st}{M_S^4} + \mathcal{O}(M_S^{-6})$$

– string resonances in UV:  $\sqrt{n} M_S$

⇒ quantum gravity testable at LHC

# New physics at the LHC

## Physics in the LHC era

- understand e-w symmetry breaking
- confirm new physics [dark matter]
- complete Standard Model



# New physics at the LHC

## Physics in the LHC era

- understand e-w symmetry breaking
- confirm new physics [dark matter]
- complete Standard Model

## LHC physics is fun physics!

- look for solid new-physics signals
- measure weak-scale Lagrangian
- determine fundamental physics
  
- construct testable new-physics hypotheses
- implement into realistic simulations
- avoid getting killed by QCD

⇒ **LHC + testable TeV-scale models: more than discovery machine!**



**UV-Completing  
the Standard  
Model at the LHC**

**Tilman Plehn**

TeV scale

Masses

Parameters

Spins & cascades

Spin & jets

**Extra dimensions**