

Quantum Gravity at LHC Scales

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Outline

LHC basics

Collider searches

Large dimensions

Effective theory

String theory

Fixed point

Warped dimensions

Fixed point

Standard–Model effective theory

Hadron colliders — the big guys

Large extra dimensions

Effective KK theory

String theory completion

Fixed-point completion

Warped extra dimensions

Fixed-point completion

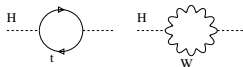
Standard–Model effective theory

Data vs renormalizable Standard Model

- dark matter? [only solid evidence for new physics, weak–scale?]
 - $(g - 2)_\mu$? [loop effects around weak scale?]
 - flavor physics? [new operators above 10^4 GeV?]
 - neutrino masses? [see-saw at 10^{11} GeV?]
 - gauge–coupling unification? [something happening above 10^{16} GeV?]
 - gravity? [mostly negligible below 10^{19} GeV]
- ⇒ obviously effective theory, cutoff negotiable

Problem with fundamental Higgs

- mass driven to cutoff: $\delta m_H^2/m_H^2 \propto g^2(2m_W^2 + m_Z^2 + m_H^2 - 4m_t^2) \Lambda^2$
 - easy solution: tune counter term
 - whole idea of gauge theories betrayed, evil
 - or new physics at TeV scale:
 - supersymmetry
 - extra dimensions
 - little Higgs, Higgsless, composite Higgs...
 - typically cancellation by new states or discussing away high scale
 - beautiful concepts, challenged at TeV scale
- ⇒ whatever is there - LHC's job to sort it out



Collider searches

Real-particle production

- produce searched-for particle
 - observe decay [before hadronization]
 - reconstruct decay products [or missing energy]
 - measure mass, spin, branching ratios
- ⇒ high-energy colliders

Virtual-particle effects

- produce and measure something known [like $pp \rightarrow \ell^+ \ell^-$]
 - compare to Standard Model predictions
 - trust in quantum theory and error estimates
 - study deviations
- ⇒ high-precision colliders

Rare effect or rare decays [B physics, EDMs]

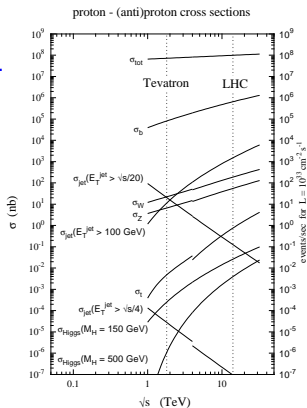
- produce something known [like B_S, \dots]
 - find effect forbidden in Standard Model
- ⇒ well-chosen experiment, not multi-purpose

Collider searches

Everything you always wanted to know about LHC...

- signal: everything new, exciting and rare
background: yesterday's signal
- Standard Model: theory of background
QCD: evil background theory trying to kill us
- trigger: no leptons/photons — not on tape
- $N_{\text{events}} = \sigma \cdot \mathcal{L} \cdot \epsilon$

⇒ discovery statistical $N_S / \sqrt{N_B} > 5$



Collider searches

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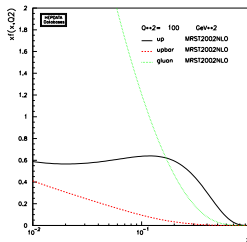
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- ⇒ **discovery statistical $N_S/\sqrt{N_B} > 5$**

Computing stuff for hadron colliders

- protons to first approximation valence quarks

$$\sigma_{AB} = \sum_{a,b} \int_0^1 dx_a dx_b f_{a/A}(x_a) f_{b/B}(x_b) \hat{\sigma}_{ab}$$

- (1) parton density $f_{a/A}(x_a)$:
probability to find a with momentum fraction x_a in A
 - (2) partonic cross section $\hat{\sigma}_{ab}$:
perturbative in QCD ["hard process"]
 - integration over partonic energy scale
- ⇒ **energetic valence quarks ahead of many gluons**



Large extra dimensions

Remember the hierarchy problem

- fundamental scalars bad with high scale present
 - weakness of gravity means large Planck scale $G_N = 1/(16\pi M_{\text{Planck}})^2$
- ⇒ solution: another reason why we see huge non-fundamental M_{Planck}

Large extra dimensions (ADD) [Antoniadis, Arkani-Hamed, Dimopoulos, Dvali]

- Einstein–Hilbert action for low fundamental Planck scale

$$\begin{aligned} S &= -\frac{1}{2} \int d^4x \sqrt{|g|} M_D^2 R \rightarrow -\frac{1}{2} \int d^{4+n}x \sqrt{|g|} M_D^{2+n} R \\ &= -\frac{1}{2} (2\pi r)^n \int d^4x \sqrt{|g|} M_D^{2+n} R \\ &\equiv -\frac{1}{2} \int d^4x \sqrt{|g|} M_{\text{Planck}}^2 R \end{aligned}$$

- ⇒ express M_{Planck} in terms of fundamental M_D

$$M_{\text{Planck}} = M_D (2\pi r M_D)^{n/2}$$

Numbers to make it work

- free parameter $rM_D \gg 1$
 - constraints from gravity tests above $\mathcal{O}(\text{mm})$
- ⇒ signatures of strong gravity in extra dimension?

$M_D = 1 \text{ TeV}$	
n	r
1	10^{12} m
2	10^{-3} m
3	10^{-8} m
...	...
6	10^{-11} m

Effective KK theory

Toy model: only gravitons in extra dimensions

- expand the metric [graviton field h]

$$ds^2 = g_{MN}^{(4+n)} dx^M dx^N = \left(\eta_{MN} + \frac{1}{M_D^{n/2+1}} h_{MN} \right) dx^M dx^N$$

- matter in Einstein's equation [no cosmological constant]

$$R_{AB} - \frac{1}{2+n} g_{AB} R = -\frac{1}{M_D^{2+n}} \begin{pmatrix} T_{\mu\nu}(x) \delta^{(n)}(y) & 0 \\ 0 & 0 \end{pmatrix}$$

- Fourier transformation of extra dimensions [KK excitations for periodic boundary conditions]

$$h_{AB}(x; y) = \sum_{m_1=-\infty}^{\infty} \cdots \sum_{m_j=-\infty}^{\infty} \frac{h_{AB}^{(m)}(x)}{\sqrt{(2\pi r)^n}} e^{i \frac{m_j y_j}{r}}$$

- universal interactions of KK gravitons [$h_{AB} \rightarrow G_{\mu\nu}$, QCD massless]

$$(\square + m_k^2) G_{\mu\nu}^{(k)} = \frac{1}{M_{\text{Planck}}} \left[-T_{\mu\nu} + \left(\frac{\partial_\mu \partial_\nu}{\hat{m}^2} + \eta_{\mu\nu} \right) \frac{T_\lambda^\lambda}{3} \right] = \frac{-T_{\mu\nu}}{M_{\text{Planck}}}$$

- tiny KK mass splitting [$M_D = 1 \text{ TeV}$]

$$\delta m \sim \frac{1}{r} = 2\pi M_D \left(\frac{M_D}{M_{\text{Planck}}} \right)^{2/n} = \begin{cases} 0.003 \text{ eV} & (n=2) \\ 0.1 \text{ MeV} & (n=4) \\ 0.05 \text{ GeV} & (n=6) \end{cases}$$

⇒ continuum of weakly interacting gravitons at the LHC

Large extra dimensions

Effective KK theory [Giudice, Rattazzi, Wells; Han, Lykken, Zhang;...]

- real radiation of continuous KK tower [$dm/d|k| = 1/r$; $(d\sigma) \propto 1/M_{\text{Planck}}^2$]

$$(d\sigma) \rightarrow \int dm (d\sigma) S_{n-1} m^{n-1} r^n = \int dm (d\sigma) \frac{S_{n-1} m^{n-1}}{(2\pi M_D)^n} \left(\frac{M_{\text{Planck}}}{M_D} \right)^2$$

- higher-dimensional operator from virtual graviton exchange [s-channel in DY]

$$\mathcal{A} = \frac{1}{M_{\text{Planck}}^2} \frac{1}{s - m_{\text{KK}}^2} \rightarrow \frac{S_{\delta-1}}{2} \frac{\Lambda^{n-2}}{M_D^{n+2}}$$

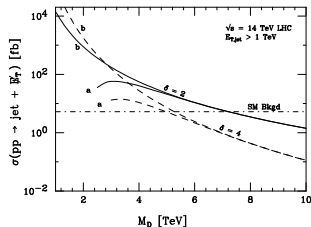
- $1/M_D^2$ interactions for KK tower

⇒ like any effective theory valid for $E < M_D$

Real emission $pp \rightarrow G_{\text{KK}} + \text{jets}$ [Giudice, Rattazzi, Wells; Vacavant, Hichliffe...]

- recoil against hard jet [$E_j \sim M_D$]
background: radiation of $Z \rightarrow \nu\bar{\nu}$
- $\mathcal{M} = 0$ for $E_{\text{parton}} > \Lambda_{\text{cutoff}}$
 $\mathcal{M} = 0$ automatically for $m_{\text{KK}} > E_{\text{parton}}$
- effective cutoff: steep gluon density
- little UV sensitivity for $\Lambda_{\text{cutoff}} \rightarrow \infty$

⇒ explicit cutoff not crucial



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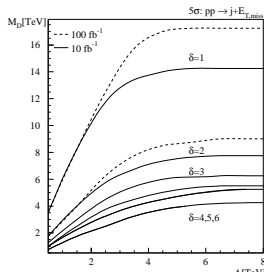
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Virtual gravitons at LHC

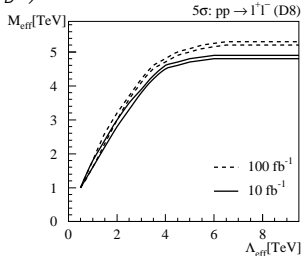
Effective theory of virtual gravitons [Giudice & Strumia; Giudice, Strumia, TP; Kachelries & Plümacher,...]

- virtual graviton in s channel $pp \rightarrow \mu^+ \mu^-$ [$q\bar{q}$ and gg initial states]
- reconstructed $m_{\mu\mu}$ for photon, Z, graviton
- loop-induced D6 operator [axial vector-axial vector]

$$\mathcal{O} = \frac{1}{16} \frac{1}{M_D^2} \frac{\pi^{n-2}}{\Gamma^2(n/2)} \left(\frac{\Lambda_{\text{cutoff}}}{M_D} \right)^{2n+2}$$

- tree-induced D8 operator [leading constant in $\sqrt{s}/\Lambda_{\text{cutoff}}$]

$$S = \frac{S_{n-1}}{M_D^{2+n}} \int dm \frac{m^{n-1}}{s+m^2} = \begin{cases} \frac{4\pi}{M_{\text{eff}}^4} & \text{(effective scale)} \\ \frac{S_{n-1}}{M_D^4} \frac{1}{2} \left(\frac{\Lambda_{\text{cutoff}}}{M_D} \right)^{n-2} & \text{(NDA)} \\ \frac{S_{n-1}}{M_D^4} \frac{1}{n-2} \left(\frac{\Lambda_{\text{cutoff}}}{M_D} \right)^{n-2} & \text{(cutoff } \Theta) \end{cases}$$



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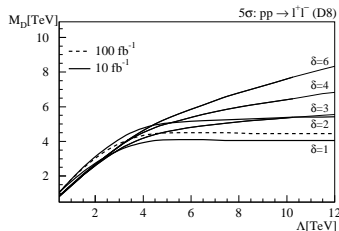
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- two-dimensional integration over (m, s)
cutoff requiring $m \lesssim M_D$
additional cutoff for $\sqrt{s} \lesssim M_D$
 - rates scaling $M_D^{\text{max}} \sim \Lambda_{\text{cutoff}}^{(n-2)/(n+2)}$
- ⇒ **breakdown of effective theory**



String theory completion

String theory as UV completion [e.g. Cullen, Perelstein, Peskin; Antoniadis, Benakli, Laugier...]

- 'In particular, the only known framework that allows for a self-consistent description of quantum gravity is string theory'

- Regge excitations of Standard Model [Burikham, Figy, Han]

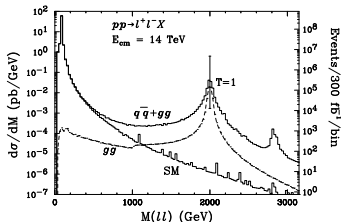
- compute different helicity contributions, like

$$\begin{aligned} A(e_L^+ e_R^- \rightarrow \gamma_L \gamma_R) &= -2e^2 \sqrt{\frac{u}{t}} \left(\frac{u}{s} + \frac{t}{s} - 1 \right) = 2e^2 \sqrt{\frac{u}{t}} \\ &= -2e^2 \sqrt{\frac{u}{t}} \left(\frac{u}{s} S(s, t) + \frac{t}{s} S(s, u) - S(t, u) \right) \end{aligned}$$

- with Veneziano form factor

$$\begin{aligned} S(s, t) &= \frac{\Gamma(1 - \alpha' s) \Gamma(1 - \alpha' t)}{\Gamma(1 - \alpha' (s + t))} = \frac{\Gamma(1 - s/M_S^2) \Gamma(1 - t/M_S^2)}{\Gamma(1 - (s + t)/M_S^2)} \\ &= 1 - \frac{\pi^2}{6} \frac{st}{M_S^4} + \mathcal{O}(M_S^{-6}) \end{aligned}$$

- dominant over KK because $M_{\text{eff}} \gg M_S$
closed-string gravitons suppressed
- effective operator below string scale
- string resonances: $\sqrt{n} M_S$



Fixed-point completion

Modified graviton propagator [Reuter; Fischer & Litim]

- effective action: $\Gamma_k = 1/(16\pi G_k) \int d^{4+n}x \sqrt{g} [-R(g) + \dots]$
- gravity weak enough at high energies?
- IR — no running; M_D regime — strong effects; UV — fixed point
- iterative approach: start with anomalous dimension of graviton

$$P(s, m) = \begin{cases} \frac{1}{s + m^2} & \sqrt{s}, m < k_{\text{trans}} \\ \frac{M_D^{n+2}}{(s + m^2)^{n/2+2}} & \sqrt{s}, m > k_{\text{trans}} \end{cases}$$

- IR and UV contributions to D8 operator [leading in \sqrt{s}/m , matched to $1/s^2$]

$$S^{(\text{FP})} = \frac{S_{n-1}}{M_D^4} \left(\frac{k_{\text{trans}}}{M_D} \right)^{n-2} \frac{n-1}{n-2} = (1 + (n-2)) S^{(\Theta)}$$

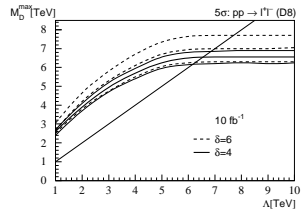
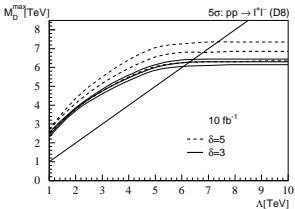
- needed and not needed:
 - (1) transition scale $k_{\text{trans}} \sim M_D$ [anomalous dimension modeled]
 - (2) no artificial Λ_{cutoff} for m integration
 - (3) matching/cutoff for s integration

⇒ UV fixed point indeed solution to our problems

Fixed-point completion

test with artificial Λ_{cutoff} setting $\mathcal{M}_{\text{KK}} = 0$

- perfect decoupling, as expected [similar to real emission]
- mild effects for $k_{\text{trans}} = M_D \pm 10\%$ [more details to be studied]
- reach largely independent of n



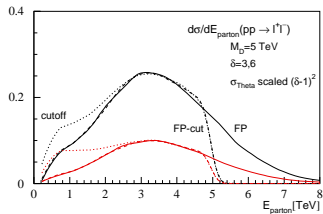
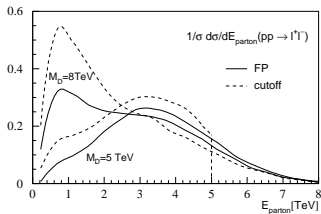
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Shape of graviton kernel

- shift to larger $m_{\ell\ell}$ [shown $n = 3$]
- small $m_{\ell\ell}$: factor $\mathcal{S} \propto (n - 1)$
- large $m_{\ell\ell}$: factor $\mathcal{S}^2 \propto (n - 1)^2$
- UV contribution predicted and not negligible



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predicted LHC rates

- $\mathcal{S}^{(\text{FP})}$: UV regime in addition to $\mathcal{S}^{(\Theta)} \equiv \mathcal{S}^{\text{IR}}$

σ [fb]	$n = 3$			$n = 6$		
	2 TeV	5 TeV	8 TeV	2 TeV	5 TeV	8 TeV
M_D						
$\mathcal{S}^{(\text{NDA})}$	43.6	0.18	0.0053	263	1.11	0.031
$\mathcal{S}^{(\Theta)}$	173	0.72	0.0204	66	0.28	0.008
$\mathcal{S}^{(\text{FP})}$	408	1.24	0.0317	398	1.21	0.031

⇒ proof of concept quantitatively very promising

Fixed point

Alternative: start with form factor [Hewett & Rizzo]

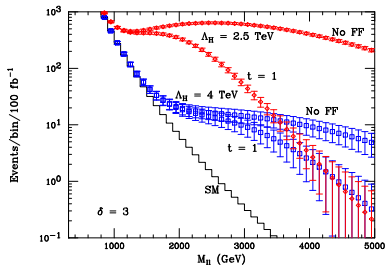
- dress s-channel coupling with form factor

$$\frac{1}{M_D^{2+n}} \rightarrow \frac{1}{M_D^{2+n}} F(s) = \frac{1}{M_D^{2+n}} \left(1 + \left(\frac{\sqrt{s}}{tM_D} \right)^{2+n} \right)^{-1}$$

- avoids matching of s integration (unitarity)

$$\frac{1}{M_D^{2+n}} F(s) \sim \frac{1}{M_D^{2+n}} \left(\frac{tM_D}{\sqrt{s}} \right)^{2+n} \sim \frac{1}{s^{1+n/2}}$$

- improvement of s integration/matching
- still cutoff in KK-mass integration



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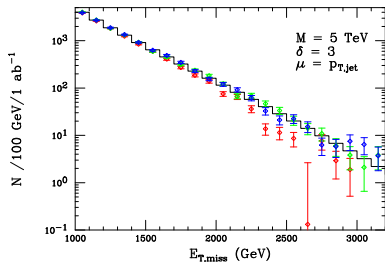
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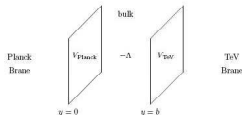
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- improvement of s integration/matching
- still cutoff in KK-mass integration
- also applicable to graviton emission [less impressive effect]

⇒ both, KK and energy integrations understood for fixed point



Warped extra dimensions



Another Solution to hierarchy problem [Randall & Sundrum]

- one extra dimension, not flat [TeV brane at $y = b$]

$$ds^2 = e^{-2k|y|} \eta_{\mu\nu} dx^\mu dx^\nu - dy^2 \Leftrightarrow g_{AB} = \begin{pmatrix} e^{-2k|y|} \eta_{\mu\nu} & 0 \\ 0 & \eta_{jk} \end{pmatrix}$$

- integration measure in our usual Lagrangian $d^4\tilde{x} e^{-4kb}$, $\tilde{g}_{\mu\nu} = \eta_{\mu\nu}$

$$S = \int dy \delta(y) d^4\tilde{x} e^{-4kb} \mathcal{L} = \int d^4\tilde{x} e^{-4kb} \left[|D_\mu H|^2 - \lambda(|H|^2 - v^2)^2 + \dots \right]$$

- write effective 4D theory on TeV brane scaling all fields

$$\tilde{H} = e^{-kb} H \quad \text{scalars}$$

$$\tilde{A}_\mu = e^{-kb} A_\mu \quad \text{or } \tilde{D}_\mu = e^{-kb} D_\mu$$

$$\tilde{\Psi} = e^{-3kb/2} \Psi \quad \text{fermions}$$

$$\tilde{m} = e^{-kb} m \quad \text{masses}$$

$$\tilde{v} = e^{-kb} v$$

$$\tilde{y} = y \quad \text{Yukawas}$$

- assume $kb \sim 35$ and large fundamental scales $M_D \sim M_{\text{Planck}} \sim k$

\Rightarrow all scales on TeV brane lowered

$$\tilde{v} \sim e^{-kb} M_{\text{Planck}} \lesssim 1 \text{ TeV}$$

Warped extra dimensions

Gravitons in one warped extra dimension

- re-write the metric including 4D graviton

$$ds^2 = \frac{1}{(1+kz)^2} \left(\eta_{\mu\nu} + h_{\mu\nu}(x, z) dx^\mu dx^\nu - dz^2 \right)$$

- solve Einstein's equations separating variables $\tilde{h}_{\mu\nu}(x, z) = \hat{h}_{\mu\nu}(x)\Phi(z)$

$$\partial_\mu \partial^\mu \hat{h}_{\mu\nu} = m^2 \hat{h}_{\mu\nu}$$

$$-\partial_z^2 \Phi + \frac{15}{4} \frac{k^2}{(kz+1)^2} \Phi = m^2 \Phi$$

- masses given by roots of Bessel functions $J_1(x_j) = 0$

$$m_j = x_j k e^{-kb} \sim \text{TeV} \quad x_j = 3.8, 7.0, 10.2, 16.5, \dots$$

- weak zero mode from wave-function overlap in z [approximately, neglect Bessel functions]

$$\frac{\Phi(z)|_{\text{TeV}}}{\Phi(z)|_{\text{Planck}}} \sim \frac{\sqrt{kz+1}|_{\text{Planck}}}{\sqrt{kz+1}|_{\text{TeV}}} \sim \frac{1}{\sqrt{e^{ky}}|_{\text{TeV}}} \sim \frac{1}{e^{kb/2}} \ll 1$$

- roughly universal couplings for higher modes

$$\mathcal{L} \sim \frac{1}{M_{\text{Planck}}} T^{\mu\nu} h_{\mu\nu}^{(0)} + \frac{1}{M_{\text{Planck}} e^{-kb}} T^{\mu\nu} \sum h_{\mu\nu}^{(m)}$$

⇒ Z' decays to leptons etc revisited...

Fixed point completion

Screened coupling again [Hewett & Rizzo]

- mass and energy integrations not problematic
- UV problem with widths of KK resonances

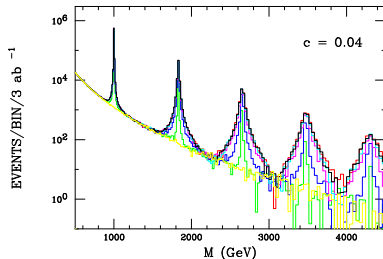
$$\frac{\Gamma_j}{m_j} \propto x_j^2$$

- form factor on mass pole implies constant widths

$$F^{-1} = 1 + \left(\frac{m_j}{tM_D e^{-kb}} \right)^3 \quad \text{implies} \quad \Gamma_j \propto k e^{-kb} t^3$$

- resonance structure also for high excitations at LHC

⇒ **RS models with fixed point well behaved in UV**



Outlook

KK Gravitons at LHC

- searches as part of puzzle of electroweak symmetry breaking
 - effective field theory:
 - real emission accidentally well defined
 - virtual-graviton predictions cutoff dependent
 - string theory:
 - solution not clear, but gravitons sub-leading anyhow
 - fixed-point picture:
 - gravity weak at large scales [asymptotic safety]
 - effect on KK mass: graviton anomalous dimension
 - effect on energy: coupling form factor
- ⇒ UV-complete ADD/RS models possible for LHC

**Quantum Gravity at
LHC Scales**

Tilman Plehn

LHC basics

Collider searches

Large dimensions

Effective theory

String theory

Fixed point

Warped dimensions

Fixed point