

Talking about Four Generations

Tilman Plehn

Heidelberg University

Single Top Workshop, DESY, 2009

Outline

Chiral 4th Generation

Electroweak precision data

Flavor constraints

Perturbativity

Direct Searches

Higgs physics

Supersymmetric fourth generation

Chiral 4th Generation

Some questions [Sher; Holdom; Hou;... ;Kribs, TP, Spannowsky, Tait]

- simply phenomenological: why three generations? [review: Framton, Hung, Sher]
- anomaly cancellation?
light neutrinos and LEP?
Majorana neutrinos in neutrinoless double beta decay?
electroweak precision data?
flavor constraints?
- ⇒ none of the constraints convincing ['Why there should not be a fourth generation'; Feyeraabend]
 - strongly interacting theory?
electroweak baryogenesis?
dark matter?
- ⇒ **at least as interesting as everything else**

Chiral 4th Generation

Some questions [Sher; Holdom; Hou;... ;Kribs, TP, Spannowsky, Tait]

- simply phenomenological: why three generations? [review: Framton, Hung, Sher]
- anomaly cancellation?
light neutrinos and LEP?
Majorana neutrinos in neutrinoless double beta decay?
electroweak precision data?
flavor constraints?
- ⇒ none of the constraints convincing ['Why there should not be a fourth generation'; Feyereabend]
 - strongly interacting theory?
electroweak baryogenesis?
dark matter?
- ⇒ **at least as interesting as everything else**

The model [old story]

- complete additional generation $[Q_4, U_4, D_4, L_4, e_4, \nu_4]$
- masses from Yukawas
- representations as Standard Model: no FCNC
- charged currents: (4×4) fermion–mixing matrices [single-top (D0) $V_{bt} \gtrsim 0.68$]
- neutrino mass: $\mathcal{L} \sim y_4 \tilde{H} \bar{L}_4 \nu_{4R} + M \bar{\nu}_{4R}^c \nu_{4R} / 2$

Electroweak precision data

Electroweak precision data [LEPEWWG]

- Particle Data Group:

An extra generation of ordinary fermions is excluded at the 6σ level on the basis of the S parameter alone... [Erler & Langacker]

Electroweak precision data

Electroweak precision data [LEPEWWG]

- Particle Data Group:

An extra generation of ordinary fermions is excluded at the 6σ level on the basis of the S parameter alone... [Erler & Langacker]

This result assumes that...any new families are degenerate [Erler & Langacker]

Electroweak precision data

Electroweak precision data [LEPEWWG]

- Particle Data Group:

An extra generation of ordinary fermions is excluded at the 6σ level on the basis of the S parameter alone... [Erler & Langacker]

This result assumes that...any new families are degenerate [Erler & Langacker]

Just as our 3rd generation??? [Holdom; Kribs, TP, Spannowsky, Tait]

Electroweak precision data

Electroweak precision data [LEPEWWG]

- Particle Data Group:

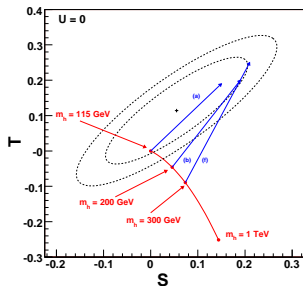
An extra generation of ordinary fermions is excluded at the 6σ level on the basis of the S parameter alone... [Erlar & Langacker]

This result assumes that...any new families are degenerate [Erlar & Langacker]

Just as our 3rd generation??? [Holdom; Kribs, TP, Spannowsky, Tait]

- okay, got it, some people prefer a boring Z'
but for the purpose of this talk let's be honest and open minded
- for our purpose: leading S and T [$\Delta U \sim 0$ as in SM]
- neutrino with Dirac mass [$\Delta S < 0$ for Majorana neutrinos: Kniehl, Kohrs]
- fermion doublet: $\Delta S = N_f / (6\pi) (1 - 2Y_\ell \log m_U^2 / m_D^2)$ [$Y_\ell = -1/2$; $Y_q = 1/6$]
- old trick: compensate $\Delta S \sim \Delta T > 0$ [Hill]
small m_H : $\Delta T \sim \Delta S \sim 0.2$
large m_H : $\Delta T \sim \Delta S + 0.2 \sim 0.3$

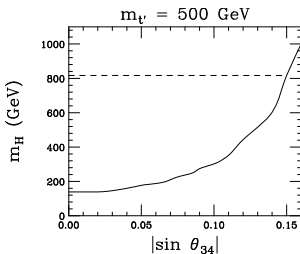
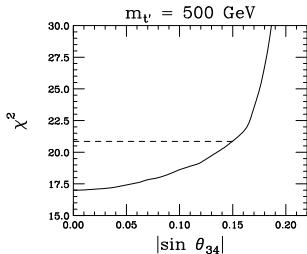
m_{U4}	m_{D4}	m_h	ΔS_{tot}	ΔT_{tot}
310	260	115	0.15	0.19
310	260	200	0.19	0.20
330	260	300	0.21	0.22
400	350	115	0.15	0.19
400	340	200	0.19	0.20
400	325	300	0.21	0.25



Flavor constraints

More precision data [Alex?]

- proper ew precision fit: $|\sin \theta_{34}| < 0.11$ for $m_H \lesssim 280$ GeV [Chanowitz]



- just had flavored coffee with Frau Prof Hiller
 $B_s \rightarrow \ell^+ \ell^-$ allowed in m_{U_4} - $\sin \theta_{34}$ plane for small θ_{34}
- all allowed region

$$m_{\ell_4}, m_{\nu_4} \gtrsim 100 \text{ GeV} \quad m_{\ell_4} - m_{\nu_4} \simeq 30 - 60 \text{ GeV}$$

$$m_{U_4}, m_{d_4} \gtrsim 260 \text{ GeV} \quad m_{U_4} - m_{d_4} \simeq \left(1 + \frac{1}{5} \ln \frac{m_H}{115 \text{ GeV}}\right) \times 50 \text{ GeV}$$

$$|V_{ud_4}|, |V_{cd_4}|, |V_{u_4 d}| \lesssim 0.04 \dots$$

$$|U_{e4}|, |U_{\mu 4}| \lesssim 0.01 \dots$$

- **small but finite θ_{34} reasonable** [CKMFitter? Heiko?]

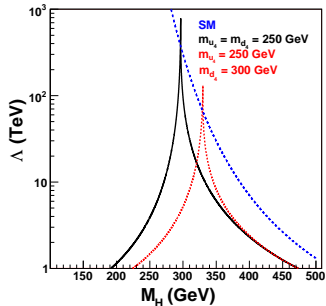
Perturbativity

Perturbative validity [review: Sher]

- Higgs mass and potential [dominant RGE effect in SM]:

$$m_H^2 = \lambda v^2 \quad 16\pi^2 \frac{d\lambda}{d \log \mu} \sim 12\lambda^2 + 4 \sum_f N_c^2 (\lambda y_f^2 - y_f^4) + \dots$$

- stability bound (m_H^{\min}) vs Landau pole or triviality bound (m_H^{\max})
- valid to scales comparable to Little Higgs
- different for 4MSSM [later]



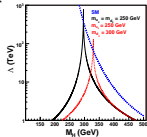
Perturbativity

Perturbative validity [review: Sher]

- Higgs mass and potential [dominant RGE effect in SM]:

$$m_H^2 = \lambda v^2 \quad 16\pi^2 \frac{d\lambda}{d \log \mu} \sim 12\lambda^2 + 4 \sum_f N_c^2 (\lambda y_f^2 - y_f^4) + \dots$$

- stability bound (m_H^{\min}) vs Landau pole or triviality bound (m_H^{\max})
- valid to scales comparable to Little Higgs
- different for 4MSSM [later]



Turn this into virtue [Hill...; Holdom]

- coupling to weak Goldstone modes strong [$m_{q4} \gtrsim 550$ GeV and $\Lambda \gtrsim 1$ TeV]
- condensates $\langle \bar{u}_4 u_4 \rangle = \langle \bar{d}_4 d_4 \rangle \neq 0$ [while $\langle \bar{t} t \rangle = \langle \bar{b} b \rangle = 0$]
- bottom-up: invoke series of $U(1)$ symmetries for

$$(\bar{q}_{4L} d_{4R}) \cdot (\bar{q}_{3L} t_R) \rightarrow (\bar{d}_{4L} d_{4R}) (\bar{t}_L t_R) \rightarrow m_t (\bar{t}_L t_R)$$

$$(\bar{q}_{4L} u_{4R}) \cdot (\bar{q}_{3L} b_R) \rightarrow (\bar{u}_{4L} u_{4R}) (\bar{b}_L b_R) \rightarrow m_b (\bar{b}_L b_R)$$

- improved compared to technicolored boson exchange:
 - no contribution to Zbb vertex ala $(\bar{T}_L T_R) (\bar{t}_R t_L) \sim (\bar{q}_L \gamma_\mu T_L) (\bar{T}_L \gamma^\mu b_L)$
 - similarly, no large u_4 - d_4 mass splitting contributing to T
- computable as weakly interacting RS model [Burdman & De Rold]

⇒ simple theory without Higgs

Direct searches

Direct searches [Holdom,...]

- assume $m_{u_4} \gtrsim m_{d_4} + 50 \text{ GeV}$ [precision constraints]
 - tree level via CKM: $d_4 \rightarrow tW$
pair production $d_4 \bar{d}_4 \rightarrow b \bar{b} W W W W$
 - loop-induced $d_4 \rightarrow bZ$ [CDF $m_U \gtrsim 270 \text{ GeV}$]
 - tree level via CKM: $u_4 \rightarrow bW$ [CDF $m_U > 260 \text{ GeV}$]
pair production $u_4 \bar{u}_4 \rightarrow b \bar{b} W W$
 - small mixing: $u_4 \rightarrow d_4 W$
pair production $u_4 \bar{u}_4 \rightarrow b \bar{b} W W W W W W W W$
 - single production better? [Fabio's talk]
the more W s the better...
 - bread-and-butter leptons plus missing energy
- ⇒ (Majorana) neutrinos more interesting?

Higgs physics

Dimension-5 Higgs couplings [e.g. SFitter-Higgs; got a hacked HDecay]

– loop effects of new particles [Arik, Arik, Cetin, Conca, Mailov, Sultansoy; Kribs, TP, Spannowsky, Tait]

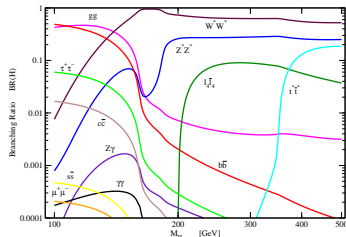
$$\Gamma_{H \rightarrow \gamma\gamma} = \frac{G_\mu \alpha^2 m_H^3}{128 \sqrt{2} \pi^3} \left| \sum_f N_c Q_f^2 A_f(\tau_f) + A_W(\tau_W) \right|^2$$

$$\Gamma_{H \rightarrow gg} = \frac{G_\mu \alpha_s^2 m_H^3}{36 \sqrt{2} \pi^3} \left| \frac{3}{4} \sum_f A_f(\tau_f) \right|^2 \quad \text{with} \quad \tau_i = \frac{m_H^2}{4m_i^2}$$

$$A_f(\tau) = \frac{2}{\tau^2} [\tau + (\tau - 1)f(\tau)]$$

$$A_W(\tau) = -\frac{1}{\tau^2} [2\tau^2 + 3\tau + 3(2\tau - 1)f(\tau)] \quad \text{with} \quad f(\tau \rightarrow 0) \rightarrow \tau$$

- (1) increase $g_{ggH} \rightarrow 3 \times g_{ggH}$
 decrease $g_{\gamma\gamma H} \rightarrow 1/3 \times g_{\gamma\gamma H}$
 light-Higgs BRs suppressed by $H \rightarrow \text{jets}$



Higgs physics

Dimension-5 Higgs couplings [e.g. SFitter-Higgs; got a hacked HDecay]

– loop effects of new particles [Arik, Arik, Cetin, Conca, Mailov, Sultansoy; Kribs, TP, Spannowsky, Tait]

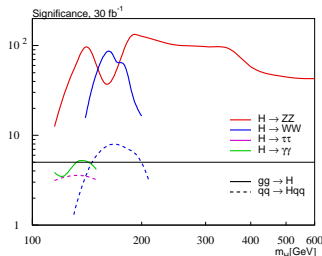
$$\Gamma_{H \rightarrow \gamma\gamma} = \frac{G_\mu \alpha^2 m_H^3}{128 \sqrt{2} \pi^3} \left| \sum_f N_c Q_f^2 A_f(\tau_f) + A_W(\tau_W) \right|^2$$

$$\Gamma_{H \rightarrow gg} = \frac{G_\mu \alpha_s^2 m_H^3}{36 \sqrt{2} \pi^3} \left| \frac{3}{4} \sum_f A_f(\tau_f) \right|^2 \quad \text{with} \quad \tau_i = \frac{m_H^2}{4m_i^2}$$

$$A_f(\tau) = \frac{2}{\tau^2} [\tau + (\tau - 1)f(\tau)]$$

$$A_W(\tau) = -\frac{1}{\tau^2} [2\tau^2 + 3\tau + 3(2\tau - 1)f(\tau)] \quad \text{with} \quad f(\tau \rightarrow 0) \rightarrow \tau$$

- increase $g_{ggH} \rightarrow 3 \times g_{ggH}$
decrease $g_{\gamma\gamma H} \rightarrow 1/3 \times g_{\gamma\gamma H}$
light-Higgs BRs suppressed by $H \rightarrow$ jets
- factor 9 enhancement of $gg \rightarrow H$ [Tevatron!?]
 $\sigma_{gg} \text{BR}_{\gamma\gamma} \rightarrow \sigma_{gg} \text{BR}_{\gamma\gamma}$
 $\sigma_{gg} \text{BR}_{ZZ} \rightarrow (5 \dots 8) \sigma_{gg} \text{BR}_{ZZ}$



Higgs physics

Dimension-5 Higgs couplings [e.g. SFitter-Higgs; got a hacked HDdecay]

– loop effects of new particles [Arik, Arik, Cetin, Conca, Mailov, Sultansoy; Kribs, TP, Spannowsky, Tait]

$$\Gamma_{H \rightarrow \gamma\gamma} = \frac{G_\mu \alpha^2 m_H^3}{128 \sqrt{2} \pi^3} \left| \sum_f N_c Q_f^2 A_f(\tau_f) + A_W(\tau_W) \right|^2$$

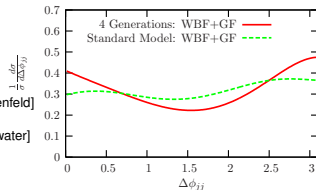
$$\Gamma_{H \rightarrow gg} = \frac{G_\mu \alpha_s^2 m_H^3}{36 \sqrt{2} \pi^3} \left| \frac{3}{4} \sum_f A_f(\tau_f) \right|^2 \quad \text{with} \quad \tau_i = \frac{m_H^2}{4m_i^2}$$

$$A_f(\tau) = \frac{2}{\tau^2} [\tau + (\tau - 1)f(\tau)]$$

$$A_W(\tau) = -\frac{1}{\tau^2} [2\tau^2 + 3\tau + 3(2\tau - 1)f(\tau)] \quad \text{with} \quad f(\tau \rightarrow 0) \rightarrow \tau$$

- (1) increase $g_{ggH} \rightarrow 3 \times g_{ggH}$
decrease $g_{\gamma\gamma H} \rightarrow 1/3 \times g_{\gamma\gamma H}$
light-Higgs BRs suppressed by $H \rightarrow \text{jets}$
- (2) factor 9 enhancement of $gg \rightarrow H$ [Tevatron!?]
 $\sigma_{gg} \text{BR}_{\gamma\gamma} \rightarrow \sigma_{gg} \text{BR}_{\gamma\gamma}$
 $\sigma_{gg} \text{BR}_{ZZ} \rightarrow (5 \cdots 8) \sigma_{gg} \text{BR}_{ZZ}$
- (3) misleading WBF correlations [TP, Rainwater, Zeppenfeld]
- (4) Higgs pair production the winner [Baur, TP, Rainwater]

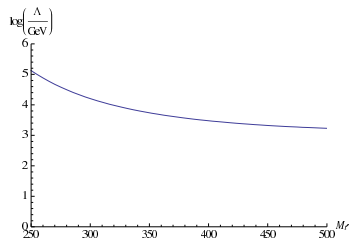
⇒ **if nothing else — what a great straw man!**



Supersymmetric fourth generation

Motivated by baryogenesis and little hierarchy [Fok, Kribs]

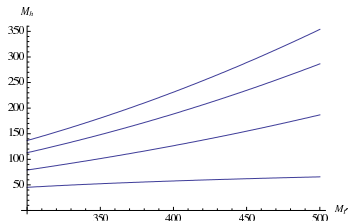
- cutoff dominantly from y_f Landau pole
preferable: $\tan \beta = 1, X_{b,t,d_4,u_4} \rightarrow 0$



Supersymmetric fourth generation

Motivated by baryogenesis and little hierarchy [Fok, Kribs]

- cutoff dominantly from y_f Landau pole
preferable: $\tan \beta = 1, X_{b,t,d_4,u_4} \rightarrow 0$
- Higgs mass all generated by heavy loops [shown pole masses]
 $H \rightarrow WW$ discovery channel for 4MSSM
light charged Higgs possible [m_A small] [Litsey & Sher]

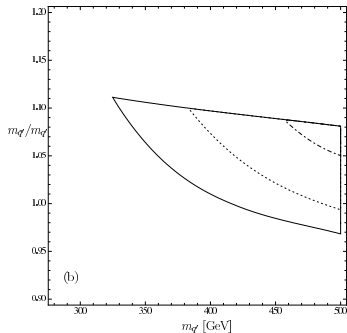


Supersymmetric fourth generation

Motivated by baryogenesis and little hierarchy [Fok, Kribs]

- cutoff dominantly from y_f Landau pole
preferable: $\tan \beta = 1$, $X_{b,t,d_4,u_4} \rightarrow 0$
- Higgs mass all generated by heavy loops [shown pole masses]
 $H \rightarrow WW$ discovery channel for 4MSSM
light charged Higgs possible [m_A small] [Litsey & Sher]
- electroweak baryogenesis [strongly first order]
net zero-temperature effects zero in SUSY limit
finite-temperature effects only helpful from bosons
finite-temperature W, Z too small for $m_H \lesssim m_W$ [Hebecker; Laine & Rummukainen]
sufficient from \tilde{u}_4 and \tilde{d}_4 [$2 \times 2 \times 3$ d.o.f.]
mass range $1 \lesssim m_{\tilde{q}_4} / m_{q_4} \lesssim 1.1$ preferred

⇒ for once an actual motivation?



A fourth generation at the LHC

4th Generation

Precision data

Flavor

Perturbativity

Direct Searches

Higgs physics

4MSSM

- it's fun
- it's not ruled out
- it has many interesting faces

Talking about Four Generations

Tilman Plehn

4th Generation

Precision data

Flavor

Perturbativity

Direct Searches

Higgs physics

4MSSM