

Higgs Physics

Tilman Plehn

Lagrangian

Tagging jets

Higgs couplings

D6 Lagrangians

Weak scale

High scale

Higgs Physics after the Discovery

Tilman Plehn [and lots of Karlsruhe people]

Universität Heidelberg

Karlsruhe, 5/2013

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Immediate questions

1. What is the 'Higgs' Lagrangian?

- psychologically: looked for Higgs, so found a Higgs
- CP-even spin-0 scalar expected
 - spin-1 vector unlikely
 - spin-2 graviton unexpected

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2. What are the coupling values?

- requires fixed operator basis
- Standard Model structure?
- anomalous couplings?

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2. What are the coupling values?

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- anomalous couplings?

3. What can we expect in the future?

- WBF analyses still sub-leading
- VH and $t\bar{t}H$ missing
- self coupling not accessible?

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Lagrangian

Equivalent questions

- what are the Higgs quantum numbers?
- what is the structure of the Higgs Lagrangian?
- can the Higgs give mass to heavy states?

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Lagrangian

Equivalent questions

- what are the Higgs quantum numbers?
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Heavy flavor inspiration

- for any observed Higgs coupling there exists a renormalizable operator
- except Higgs production in gluon fusion
- except Higgs decay to photons
- except g_{WWH} might mean $HW^{\mu\nu}W_{\mu\nu}$
- Higgs Lagrangian all but trivial

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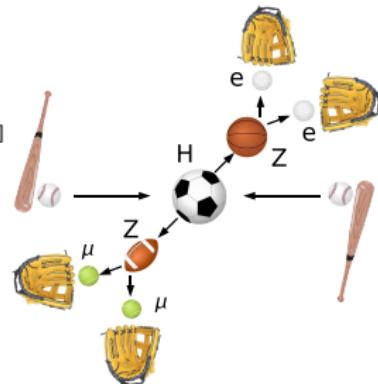
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- ⇒ **analyze Higgs kinematics** [in as many channels as possible]



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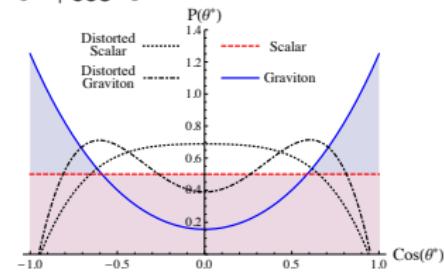
High scale

Model independent angles

- first step: Higgs polar angle for spin-0 vs spin-2

$$\frac{d\Gamma_0}{d \cos \theta^*} \sim P_0(\theta^*) = 1$$

$$P_2(\theta^*) \sim 1 + 6 \cos^2 \theta^* + \cos^4 \theta^*$$



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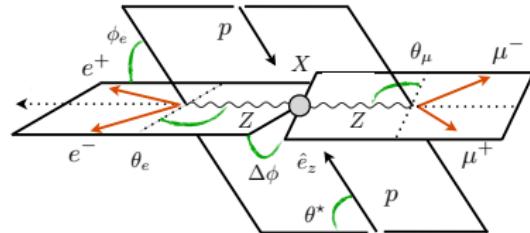
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Model independent angles

- $H \rightarrow ZZ$ decays [Choi et al; Melnikov et al; Lykken et al; v d Bij et al; Englert, Spannowsky, Takeuchi] classic Cabibbo–Maksymowicz–Dell’Aquila–Nelson angles [all over LHCb]

$$\begin{aligned}\cos \theta_e &= \hat{p}_{e-} \cdot \hat{p}_{Z_\mu} \Big|_{Z_e} & \cos \theta_\mu &= \hat{p}_{\mu-} \cdot \hat{p}_{Z_e} \Big|_{Z_\mu} & \cos \theta^* &= \hat{p}_{Z_e} \cdot \hat{p}_{\text{beam}} \Big|_X \\ \cos \phi_e &= (\hat{p}_{\text{beam}} \times \hat{p}_{Z_\mu}) \cdot (\hat{p}_{Z_\mu} \times \hat{p}_{e-}) \Big|_{Z_e} \\ \cos \Delta\phi &= (\hat{p}_{e-} \times \hat{p}_{e+}) \cdot (\hat{p}_{\mu-} \times \hat{p}_{\mu+}) \Big|_X\end{aligned}$$



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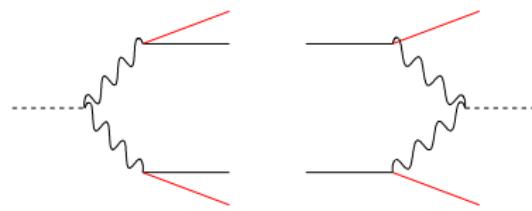
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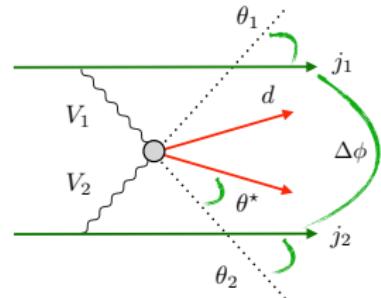
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$$\begin{aligned} \cos \theta_1 &= \hat{p}_{j_1} \cdot \hat{p}_{V_2} \Big|_{V_1 \text{ Breit}} & \cos \theta_2 &= \hat{p}_{j_2} \cdot \hat{p}_{V_1} \Big|_{V_2 \text{ Breit}} & \cos \theta^* &= \hat{p}_{V_1} \cdot \hat{p}_d \Big|_X \\ \cos \phi_1 &= (\hat{p}_{V_2} \times \hat{p}_d) \cdot (\hat{p}_{V_2} \times \hat{p}_{j_1}) \Big|_{V_1 \text{ Breit}} \\ \cos \Delta\phi &= (\hat{p}_{q_1} \times \hat{p}_{j_1}) \cdot (\hat{p}_{q_2} \times \hat{p}_{j_2}) \Big|_X . \end{aligned}$$



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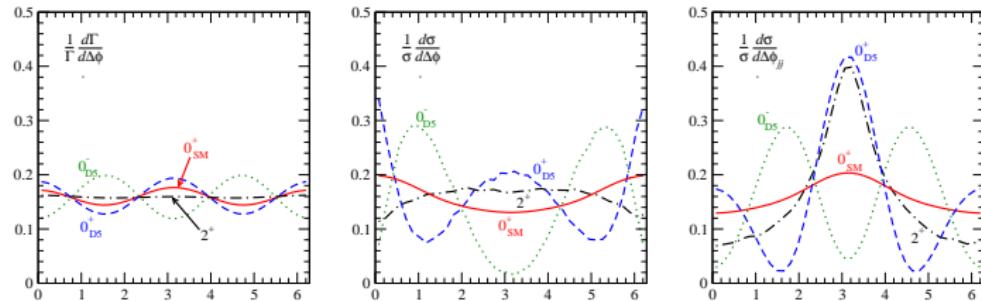
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⇒ different approaches with similar physics

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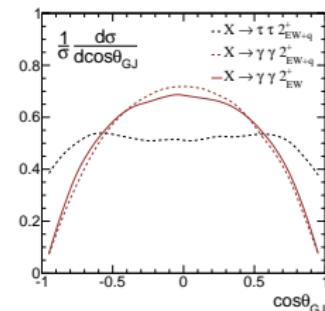
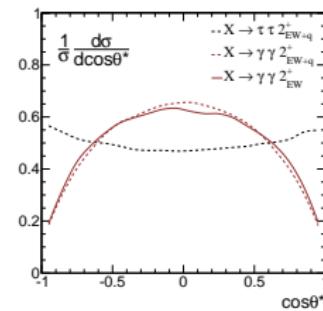
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Spin-2 test? [Englert, Mawatari, Netto, TP]

- unitarization affecting all energy variables
- try Gottfried-Jackson angle $[\hat{p}_{X,lab} \text{ vs } \hat{p}_{d,X}]$; Frank, Rauch, Zeppenfeld; Schumi]



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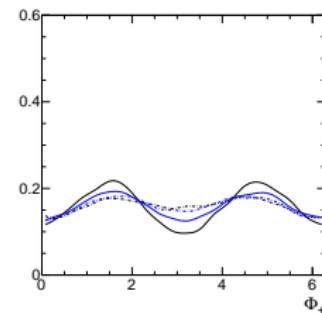
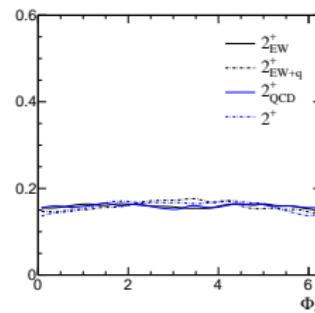
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- unitarization affecting all energy variables
- try Gottfried-Jackson angle $[\hat{p}_{X,lab} \text{ vs } \hat{p}_{d,X}; \text{Frank, Rauch, Zeppenfeld; Schumi}]$
- alternatively $\phi_1 + \phi_2$ [Hagiwara, Li, Mawatari]



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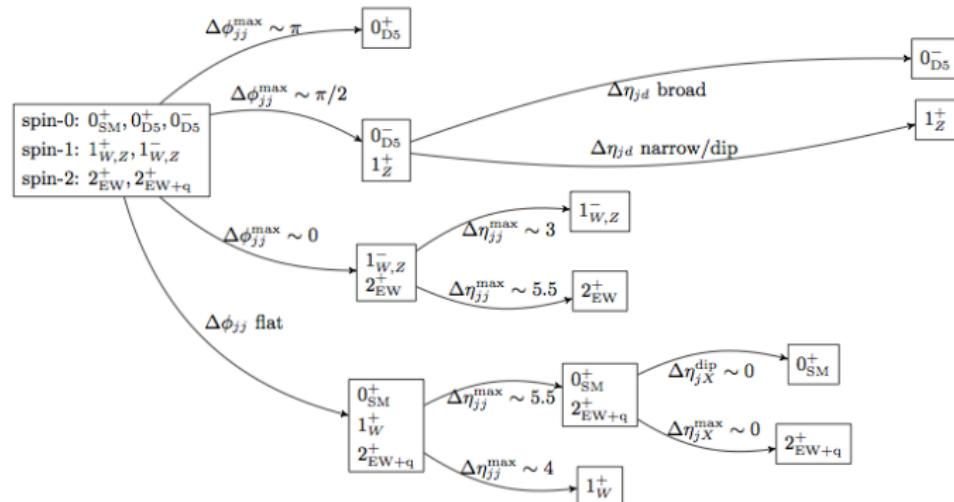
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- diagrammatic analysis for WBF $[\Delta\eta_{jj} \text{ crucial}]$



⇒ angular observables in most channels

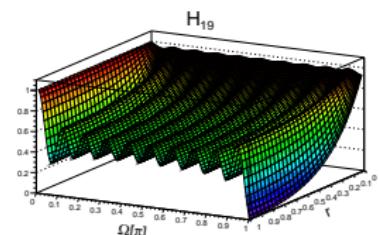
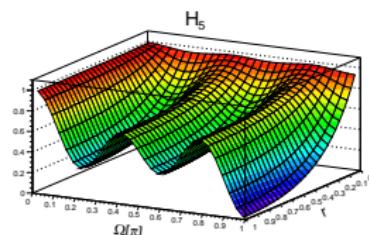
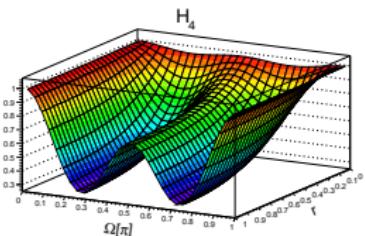
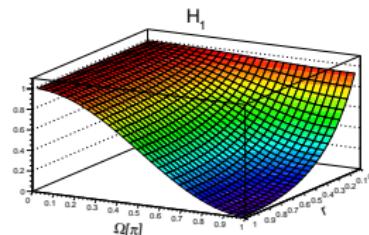
Fox-Wolfram moments

Series in spherical harmonics [Field, Kanev, Tayebnejad; BaBar; Bernaciak, Buschmann, Butter, TP]

- originally alternative to event shapes

$$H_\ell^T = \frac{4\pi}{2\ell+1} \sum_{m=-\ell}^{\ell} \left| \sum_{i=1}^N Y_\ell^m(\Omega_i) \frac{p_{T,i}}{p_{T,\text{tot}}} \right|^2 = \sum_{i,j=1}^N \frac{p_{T,i} p_{T,j}}{p_{T,\text{tot}}^2} P_\ell(\cos \Omega_{ij}) ,$$

- defined on separated jets for a start



	$H_\ell < 0.3$	$0.3 < H_\ell < 0.7$	$0.7 < H_\ell < 1$
even ℓ	forbidden	democratic	ordered, collinear, back-to-back
odd ℓ	back-to-back	democratic	collinear, ordered

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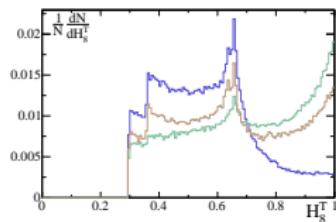
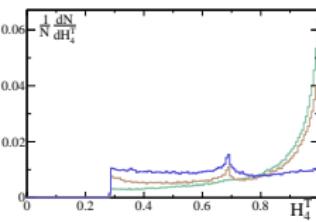
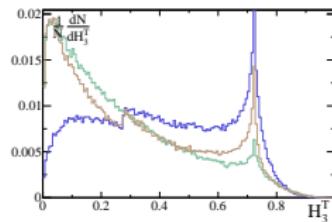
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- defined on separated jets for a start
- applied to tagging jets in WBF [$m_{jj} > 600 \text{ GeV}$]



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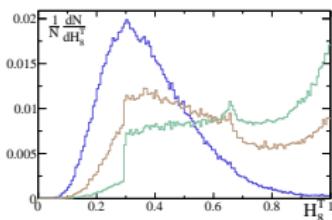
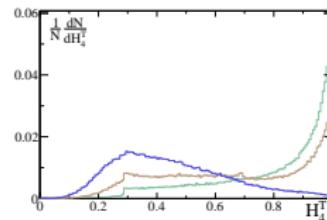
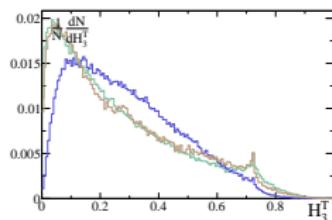
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- applied to all jets in WBF



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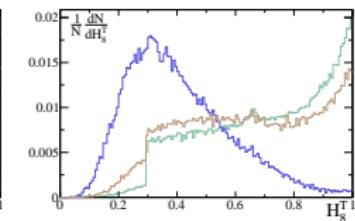
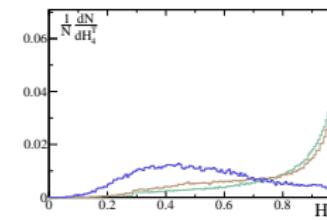
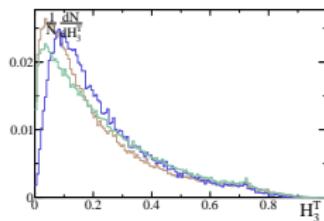
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- applied to tagging jets in WBF [$m_{jj} > 600$ GeV]
- applied to all jets in WBF
- applied to all jets after WBF cuts



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- applied to all jets in WBF
- applied to all jets after WBF cuts
- useful information left
tuned resolution via variable ℓ [not too correlated]
- adjust weight factor?
adjust objects entering FWMs?

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- ⇒ might be useful eventually

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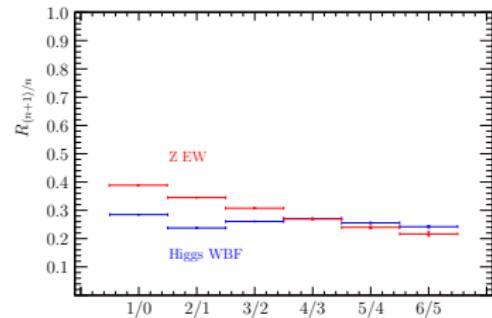
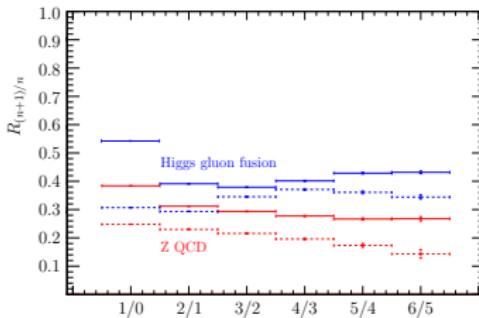
Weak scale

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Jet counting

Jets with Higgs [Englert, Gerwick, TP, Schichtel, Schumann]

- example: WBF $H \rightarrow \tau\tau$
- staircase scaling before WBF cuts [QCD and e-w processes]
- e-w Zjj production with too many structures



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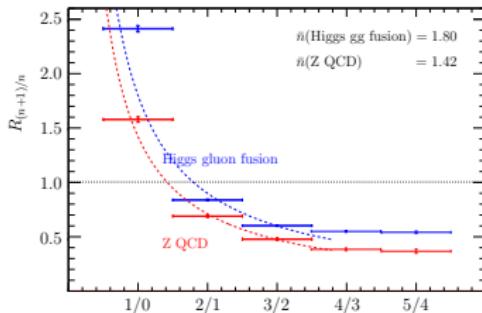
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Understanding a jet veto from QCD [simulated with SHERPA, also ask Stefan Gieseke]

- count add'l jets to reduce backgrounds

$$p_T^{\text{veto}} > 20 \text{ GeV} \quad \min y_{1,2} < y^{\text{veto}} < \max y_{1,2}$$

- Poisson for QCD processes ['radiation' pattern]



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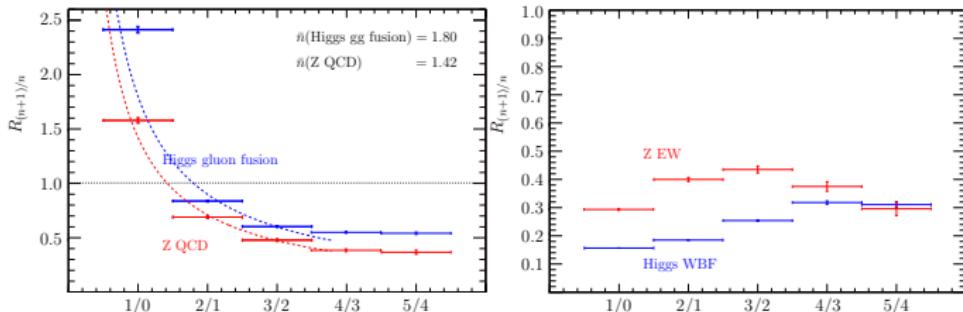
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- count add'l jets to reduce backgrounds
- $p_T^{\text{veto}} > 20 \text{ GeV}$ $\min y_{1,2} < y^{\text{veto}} < \max y_{1,2}$
- Poisson for QCD processes ['radiation' pattern]
- (fairly) staircase for e-w processes [cuts keeping signal]
- **distribution of number of jets understood**



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Higgs rates

The model

- assume: we see a scalar [ZZ and WBF correlations]
it is a narrow resonance
SM-like D4 structures
benchmarks useless
- production & decay combinations
- signal strength vs couplings?

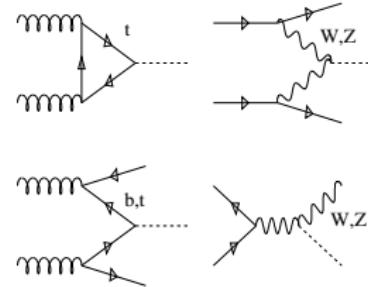
$gg \rightarrow H$
 $qq \rightarrow qqH$
 $gg \rightarrow ttH$
 $q\bar{q}' \rightarrow WH$
 plus a little problem



$H \rightarrow ZZ$
 $H \rightarrow WW$
 $H \rightarrow b\bar{b}$
 $H \rightarrow \tau_\ell^+ \tau_\ell^-$
 $H \rightarrow \gamma\gamma$
 $H \rightarrow Z\gamma$
...



signal \times trigger
 backgrounds
 Gauss/Poisson statistics
 systematics
 theory errors



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Why 126 GeV is just perfect [Zeppenfeld et al; Dührssen et al; SFitter 2009/2012]

- because it wants a weakly interacting top partner

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Why 126 GeV is just perfect [Zeppenfeld et al; Dührssen et al; SFitter 2009/2012]

- parameters: Higgs couplings to $W, Z, t, b, \tau, g, \gamma$ [SM-like D4 operators]

$$g_{HXX} = g_{HXX}^{\text{SM}} (1 + \Delta x) \quad g_{HWW} > 0$$

- measurements:
 - $GF : H \rightarrow ZZ, WW, \gamma\gamma$
 - $WBF : H \rightarrow ZZ, WW, \gamma\gamma, \tau\tau$
 - $VH : H \rightarrow b\bar{b}$
 - $t\bar{t}H : H \rightarrow \gamma\gamma, b\bar{b}$

⇒ perfect application for SFitter

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Error analysis

Sources of uncertainty

- statistical error: Poisson
- systematic error: Gaussian, if measured
- theory error: not Gaussian
- simple argument
- LHC rate 10% off: no problem
- LHC rate 30% off: no problem
- LHC rate 300% off: Standard Model wrong
- theory likelihood flat centrally and zero far away
- profile likelihood construction: RFit [CKMFitter]

$$-2 \log \mathcal{L} = \chi^2 = \vec{\chi}_d^T C^{-1} \vec{\chi}_d$$

$$\chi_{d,i} = \begin{cases} 0 & |d_i - \bar{d}_i| < \sigma_i^{(\text{theo})} \\ \frac{|d_i - \bar{d}_i| - \sigma_i^{(\text{theo})}}{\sigma_i^{(\text{exp})}} & |d_i - \bar{d}_i| > \sigma_i^{(\text{theo})} \end{cases}$$

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- theory error: not Gaussian
- profile likelihood construction: RFit [CKMFitter]

$$-2 \log \mathcal{L} = \chi^2 = \vec{\chi}_d^T C^{-1} \vec{\chi}_d$$

$$\chi_{d,i} = \begin{cases} 0 & |d_i - \bar{d}_i| < \sigma_i^{(\text{theo})} \\ \frac{|d_i - \bar{d}_i| - \sigma_i^{(\text{theo})}}{\sigma_i^{(\text{exp})}} & |d_i - \bar{d}_i| > \sigma_i^{(\text{theo})} \end{cases}$$

Efficient combination of errors [different from Michael's ATLAS analysis]

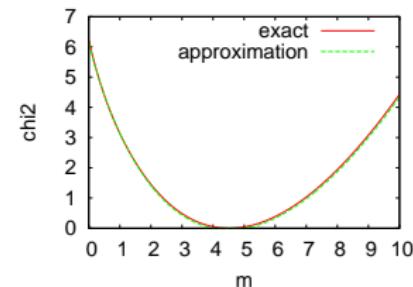
- Gaussian \otimes Gaussian: half width added in quadrature
- Gaussian/Poisson \otimes flat: RFit scheme
- Gaussian \otimes Poisson: ??

- approximate formula

$$\frac{1}{\log \mathcal{L}_{\text{comb}}} = \frac{1}{\log \mathcal{L}_{\text{Gauss}}} + \frac{1}{\log \mathcal{L}_{\text{Poisson}}}$$

- modified Minuit gradient fit last step

\Rightarrow error bars from toy measurements



Lagrangian

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D6 Lagrangians

Weak scale

High scale

Error analysis

Sources of uncertainty

- statistical error: Poisson
- systematic error: Gaussian, if measured
- theory error: not Gaussian
- profile likelihood construction: RFit [CKMFitter]

$$-2 \log \mathcal{L} = \chi^2 = \vec{\chi}_d^T C^{-1} \vec{\chi}_d$$

$$\chi_{d,i} = \begin{cases} 0 & |d_i - \bar{d}_i| < \sigma_i^{(\text{theo})} \\ \frac{|d_i - \bar{d}_i| - \sigma_i^{(\text{theo})}}{\sigma_i^{(\text{exp})}} & |d_i - \bar{d}_i| > \sigma_i^{(\text{theo})} \end{cases}$$

$$|d_i - \bar{d}_i| < \sigma_i^{(\text{theo})}$$

$$|d_i - \bar{d}_i| > \sigma_i^{(\text{theo})}$$

Systematic uncertainties

luminosity measurement	5 %
detector efficiency	2 %
lepton reconstruction efficiency	2 %
photon reconstruction efficiency	2 %
WBF tag-jets / jet-veto efficiency	5 %
<i>b</i> -tagging efficiency	3 %
τ -tagging efficiency (hadronic decay)	3 %
lepton isolation efficiency ($H \rightarrow 4\ell$)	3 %

	$\Delta B^{(\text{syst})}$
$H \rightarrow ZZ$	1 %
$H \rightarrow WW$	5 %
$H \rightarrow \gamma\gamma$	0.1 %
$H \rightarrow \tau\tau$	5 %
$H \rightarrow b\bar{b}$	10 %

Higgs couplings

Higgs sector at LHC [Zeppenfeld et al; Dührssen et al; SFitter 2009/2012; Contino et al]

- light Higgs around 126 GeV: over 10 channels ($\sigma \times BR$)
- measurements: $GF : H \rightarrow ZZ, WW, \gamma\gamma$ [first analyses]
 $WBF : H \rightarrow ZZ, WW, \gamma\gamma, \tau\tau$ [just starting]
 $VH : H \rightarrow b\bar{b}$ [BDRS-like analyses only]
 $t\bar{t}H : H \rightarrow \gamma\gamma, WW, b\bar{b}\dots$ [useful but later]
- parameters: couplings $W, Z, t, b, \tau, g, \gamma$ [plus eventually Higgs mass]

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 $VH : H \rightarrow b\bar{b}$ [BDRS-like analyses only]
 $t\bar{t}H : H \rightarrow \gamma\gamma, WW, b\bar{b}$... [useful but later]
- parameters: couplings $W, Z, t, b, \tau, g, \gamma$ [plus eventually Higgs mass]

Total width

- myths about scaling

$$N = \sigma BR \propto \frac{g_p^2}{\sqrt{\Gamma_{\text{tot}}}} \frac{g_d^2}{\sqrt{\Gamma_{\text{tot}}}} \sim \frac{g^4}{g^2 \frac{\sum \Gamma_i(g^2)}{g^2} + \Gamma_{\text{unobs}}} \xrightarrow{g^2 \rightarrow 0} 0$$

gives constraint from $\sum \Gamma_i(g^2) < \Gamma_{\text{tot}} \rightarrow \Gamma_H|_{\text{min}}$

- $WW \rightarrow WW$ unitarity: $g_{WWH} \lesssim g_{WWH}^{\text{SM}} \rightarrow \Gamma_H|_{\text{max}}$
- **SFitter assumption** $\Gamma_{\text{tot}} = \sum_{\text{obs}} \Gamma_j$ [plus generation universality]

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SFitter ansatz [Dührssen, Klute, Lafaye, TP, Rauch, Zerwas]

- SM operators $g_{HXX} = g_{HXX}^{\text{SM}} (1 + \Delta_X)$
D5 couplings $g_{ggH}, g_{\gamma\gamma H}$ free?
electroweak correction currently negligible
- experimental/theory errors on signal and backgrounds
ATLAS and CMS both included
- exclusive likelihood map
each coupling from profile likelihoods
best-fit point with Minuit
complete error analysis

Lagrangian

Tagging jets

Higgs couplings

D6 Lagrangians

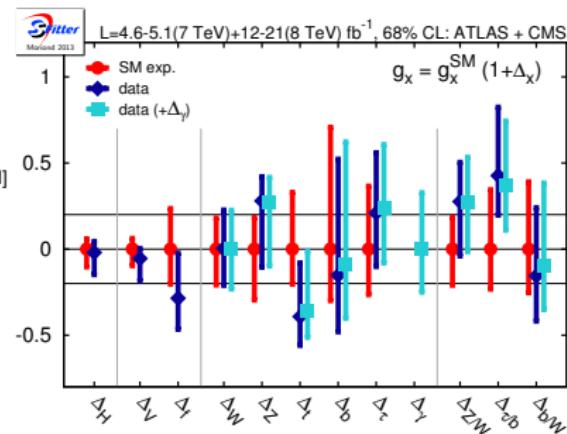
Weak scale

High scale

SFitter results

LHC including Moriond/Aspen data [SFitter: Klute, Lafaye, TP, Rauch, Zerwas]

- focus SM-like [secondary solutions possible]
- six couplings and ratios from data
 - g_b from width
 - g_g vs g_t not yet possible
 - [similar: Ellis et al, Djouadi et al, Strumia et al, Grojean et al]
- poor man's analyses: $\Delta_H, \Delta_V, \Delta_f$
- Tevatron $H \rightarrow b\bar{b}$ with little impact



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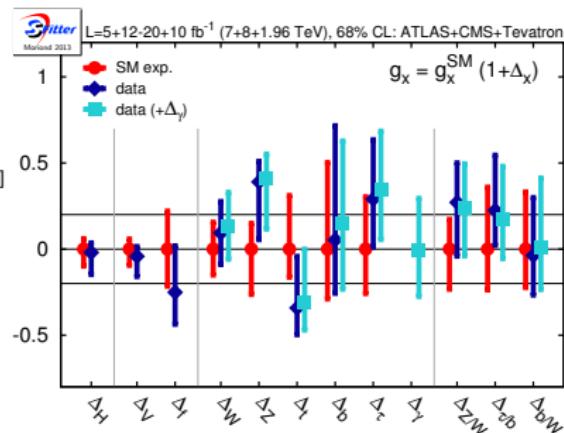
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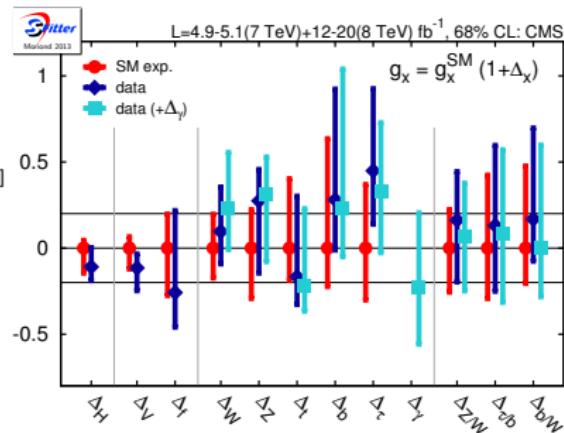
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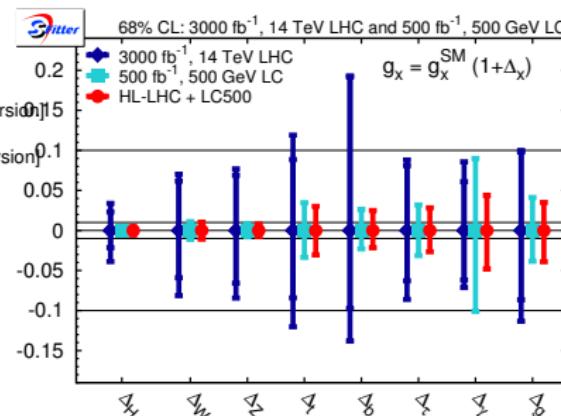
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Future dinosaurs

- LHC extrapolations unclear [SFitter version 0.15]
- theory extrapolations tricky [SFitter version 0.15]
- ILC case obvious [500 GeV for now]
- interplay in loop-induced couplings



Lagrangian

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Theorists ideas

Light Higgs as a Goldstone boson [Giudice, Grojean, Pomarol, Rattazzi]

- strongly interacting models predicting heavy broad resonance(s)
- light state if protected by Goldstone's theorem [Georgi & Kaplan]
- interesting if $v \ll f < 4\pi f$ [little Higgs $v \sim g^2 f / (2\pi)$]
- postulate new $f \gtrsim v$ and $m_\rho \rightarrow 4\pi f$ [$c_j \sim 1$] [assume custodial symmetry]
- adding D6 weak operators with relative strengths

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- adding D6 weak operators with relative strengths

$$\begin{aligned}
 \mathcal{L}_{\text{SILH}} = & \frac{c_H}{2f^2} \partial^\mu (H^\dagger H) \partial_\mu (H^\dagger H) + \frac{c_T}{2f^2} (H^\dagger \overleftrightarrow{D}^\mu H) (H^\dagger \overleftrightarrow{D}_\mu H) \\
 & - \frac{c_6 \lambda}{f^2} (H^\dagger H)^3 + \left(\frac{c_y y_f}{f^2} H^\dagger H \bar{f}_L H f_R + \text{h.c.} \right) \\
 & + \frac{i c_W g}{2m_\rho^2} (H^\dagger \sigma^i \overleftrightarrow{D}^\mu H) (D^\nu W_{\mu\nu})^i + \frac{i c_B g'}{2m_\rho^2} (H^\dagger \overleftrightarrow{D}^\mu H) (\partial^\nu B_{\mu\nu}) \\
 & + \frac{i c_{HW} g}{16\pi^2 f^2} (D^\mu H)^\dagger \sigma^i (D^\nu H) W_{\mu\nu}^i + \frac{i c_{HB} g'}{16\pi^2 f^2} (D^\mu H)^\dagger (D^\nu H) B_{\mu\nu} \\
 & + \frac{c_\gamma g'^2}{16\pi^2 f^2} \frac{g^2}{g_\rho^2} H^\dagger H B_{\mu\nu} B^{\mu\nu} + \frac{c_g g_S^2}{16\pi^2 f^2} \frac{y_t^2}{g_\rho^2} H^\dagger H G_{\mu\nu}^a G^{a\mu\nu}.
 \end{aligned}$$

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$$\begin{aligned} \mathcal{L}_{SILH} \sim & \frac{c_H}{f^2} \partial^\mu (H^\dagger H) \partial_\mu (H^\dagger H) + \frac{c_T}{f^2} (H^\dagger \overleftrightarrow{D}^\mu H) (H^\dagger \overleftrightarrow{D}_\mu H) \\ & - \frac{c_6}{(3f)^2} (H^\dagger H)^3 + \left(\frac{c_y y_f}{f^2} H^\dagger H \bar{f}_L H f_R + \text{h.c.} \right) \\ & + \frac{i c_W}{(16f)^2} (H^\dagger \sigma^i \overleftrightarrow{D}^\mu H) (D^\nu W_{\mu\nu})^i + \frac{i c_B}{(16f)^2} (H^\dagger \overleftrightarrow{D}^\mu H) (\partial^\nu B_{\mu\nu}) \\ & + \frac{i c_{HW}}{(16f)^2} (D^\mu H)^\dagger \sigma^i (D^\nu H) W_{\mu\nu}^i + \frac{i c_{HB}}{(16f^2)} (D^\mu H)^\dagger (D^\nu H) B_{\mu\nu} \\ & + \frac{c_\gamma}{(256f)^2} H^\dagger H B_{\mu\nu} B^{\mu\nu} + \frac{c_g}{(256f)^2} H^\dagger H G_{\mu\nu}^a G^{a\mu\nu}. \end{aligned}$$

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- collider phenomenology of mostly $(H^\dagger H)$ terms [Mühlleitner et al]
- ⇒ remember what your operators are!

Lagrangian

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Anomalous couplings

Anomalous Higgs couplings [Hagiwara et al; Corbett, Eboli, Gonzales-Fraile, Gonzales-Garcia]

- assume Higgs is largely Standard Model
- additional higher-dimensional couplings

$$\begin{aligned} \mathcal{L}_{\text{eff}} = & -\frac{\alpha_s V}{8\pi} \frac{f_g}{\Lambda^2} (\Phi^\dagger \Phi) G_{\mu\nu} G^{\mu\nu} + \frac{f_{WW}}{\Lambda^2} \Phi^\dagger W_{\mu\nu} W^{\mu\nu} \Phi \\ & + \frac{f_W}{\Lambda^2} (D_\mu \Phi)^\dagger W^{\mu\nu} (D_\nu \Phi) + \frac{f_B}{\Lambda^2} (D_\mu \Phi)^\dagger B^{\mu\nu} (D_\nu \Phi) + \frac{f_{WWW}}{\Lambda^2} \text{Tr}(W_{\mu\nu} W^{\nu\rho} W_\rho^\mu) \\ & + \frac{f_b}{\Lambda^2} (\Phi^\dagger \Phi) (\bar{Q}_3 \Phi d_{R,3}) + \frac{f_\tau}{\Lambda^2} (\Phi^\dagger \Phi) (\bar{L}_3 \Phi e_{R,3}) \end{aligned}$$

- plus e-w precision data and triple gauge couplings

Lagrangian

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High scale

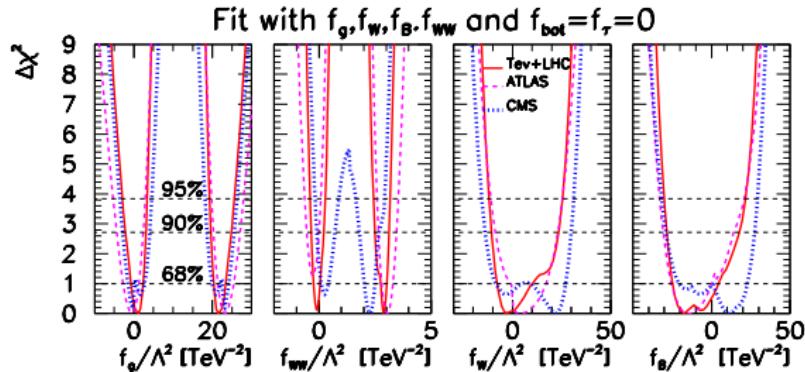
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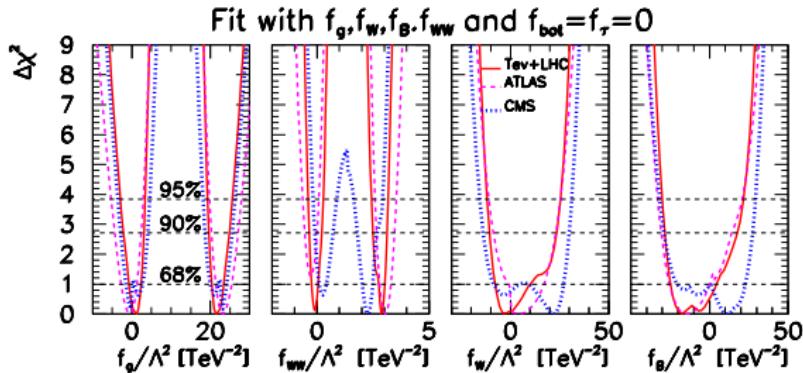
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Top Yukawa

Direct measurement $t\bar{t}H, H \rightarrow b\bar{b}$ [Atlas-Bonn: Jochen Cammin]

- crucial to understand Higgs sector [details later]
- trigger: $t \rightarrow bW^+ \rightarrow b\ell^+\nu$
reconstruction and rate: $\bar{t} \rightarrow \bar{b}W^- \rightarrow \bar{b}jj$
- continuum background $t\bar{t}b\bar{b}, t\bar{t}jj$ [weighted by b-tag]

Lagrangian

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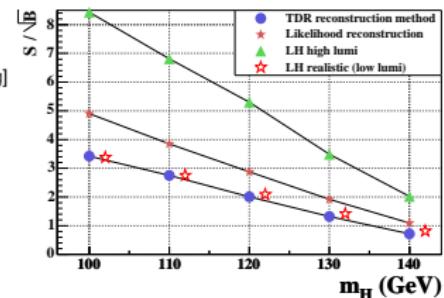
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- not a chance:
 - 1– combinatorics: m_H in $pp \rightarrow 4b_{tag} 2j \ell\nu$
 - 2– kinematics: peak-on-peak
 - 3– systematics: $S/B \sim 1/9$



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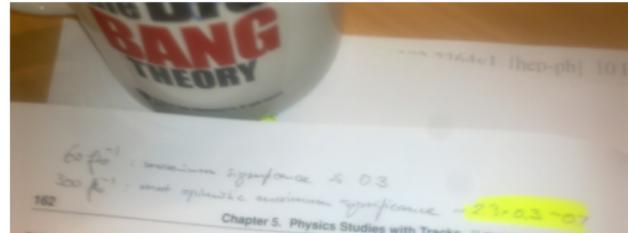
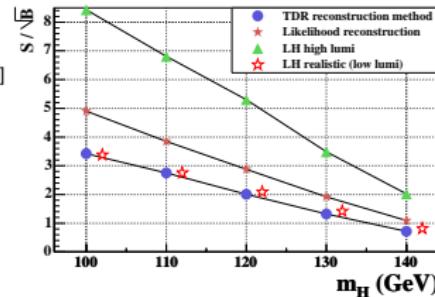


Table 5.30: Significance before and after taking into account the uncertainty dB in the total number of background events due to systematics.

mass	S/B	S/\sqrt{B}	$S/\sqrt{B} + dB^2$
$t\bar{t}H$ (115)	0.082	2.2	0.20
$t\bar{t}H$ (120)	0.043	1.8	0.15
$t\bar{t}H$ (130)	0.030	1.3	0.11
$\Delta_{btag} > 0.75$ (c_{btag})			
$t\bar{t}H$ (115)	0.108	2.0	0.44
$t\bar{t}H$ (120)	0.082	1.6	0.34
$t\bar{t}H$ (130)	0.060	1.1	0.24
$\Delta_{btag} > 0.55$ (c_{btag})			
electron	S/B	S/\sqrt{B}	$S/\sqrt{B} + dB^2$
$t\bar{t}H$ (118)	0.028	0.7	0.27
hadron	S/B	S/\sqrt{B}	$S/\sqrt{B} + dB^2$
$t\bar{t}H$ (115)	0.069	1.4	0.42
$t\bar{t}H$ (120)	0.045	0.9	0.27
$t\bar{t}H$ (130)	0.029	0.6	0.18
dilepton	S/B	S/\sqrt{B}	$S/\sqrt{B} + dB^2$
4-7 jets, 3-4 b-tagged (c_{btag})			
$t\bar{t}H$ (115)	0.018	1.8	0.19
$t\bar{t}H$ (120)	0.015	1.4	0.08
$t\bar{t}H$ (130)	0.009	0.9	0.05
4-6 jets, 4-6 b-tagged (c_{btag})			
$t\bar{t}H$ (115)	0.069	1.4	0.42
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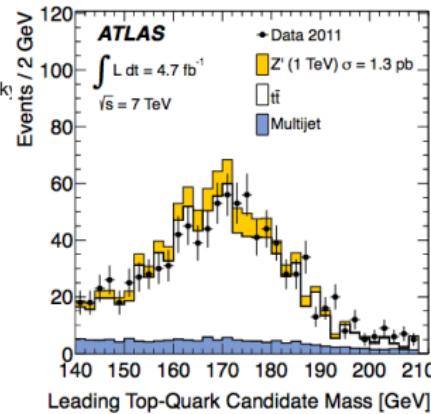
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 - 3– systematics: $S/B \sim 1/9$

New(ish) approaches

- semi-leptonic fat jets analysis [TP, Salam, Spannowski]
require tagged top and Higgs
trigger on lepton
only continuum $t\bar{t}b\bar{b}$ left [with sidebands]
top tagger working [Atlas-Heidelberg]



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 - 3– systematics: $S/B \sim 1/9$

New(ish) approaches

- semi-leptonic fat jets analysis [TP, Salam, Spannowsky, Takeuchi]
require tagged top and Higgs
trigger on lepton
only continuum $t\bar{t}b\bar{b}$ left [with sidebands]
top tagger working [Atlas-Heidelberg]
- purely leptonic matrix element method [Artoisenet, de Aquino, Maltoni, Mattelaer]
 $4b-2\ell$ -MET final state
combinatorics from matrix element

Top Yukawa

Direct measurement $t\bar{t}H, H \rightarrow b\bar{b}$ [Atlas-Bonn: Jochen Cammin]

- crucial to understand Higgs sector [details later]
- trigger: $t \rightarrow bW^+ \rightarrow b\ell^+\nu$
reconstruction and rate: $\bar{t} \rightarrow \bar{b}W^- \rightarrow \bar{b}jj$
- continuum background $t\bar{t}b\bar{b}, t\bar{t}jj$ [weighted by b-tag]
- not a chance:
 - 1– combinatorics: m_H in $pp \rightarrow 4b_{tag} \ 2j \ \ell\nu$
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- ⇒ good ideas welcome

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Weak scale

High scale

Weak scale models

Higgs portal [Englert, TP, Rauch, Zerwas, Zerwas]

- only few renormalizable links to a new physics world
general standard-hidden ansatz $H_1 = \cos \chi H_s + \sin \chi H_h$
- visible and hidden decays [plus $H_2 \rightarrow H_1 H_1$ cascade decays]

$$\Gamma_1^{\text{tot}} = \cos^2 \chi \Gamma_{\text{tot},1}^{\text{SM}} + \sin^2 \chi \Gamma_1^{\text{hid}}$$

- constraints on event rate

$$\frac{\sigma[H_1 \rightarrow XX^*]}{\sigma[H_1 \rightarrow XX^*]^{\text{SM}}} = \frac{\cos^2 \chi}{1 + \tan^2 \chi \frac{\Gamma_1^{\text{hid}}}{\Gamma_{\text{tot},1}^{\text{SM}}}}$$

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⇒ invisible Higgs needed for final answer [Eboli & Zeppenfeld]

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Form factor Higgs [Kaplan & Georgi; Contino, Espinosa, Giudice, Grojean, Mühlleitner, Pomarol, Rattazzi]

- simple trick: $\xi \equiv v/f \gtrsim 0.3$ while $m_\rho = g_\rho f \gg f$ [also not calculable]
- 1– all couplings scaled $g \rightarrow g\sqrt{1-\xi}$
 - one-parameter fit in SFitter
 - from 8 TeV data $\Delta_H = 0 \pm 0.15$
- 2– gauge couplings $g \rightarrow g\sqrt{1-\xi}$
Yukawas $g \rightarrow g(1-2\xi)/\sqrt{1-\xi}$
 - sign change of Yukawas, $g_{\gamma\gamma H}$ correlated

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Non-decoupling D6 operators

- SM: chiral fermions $g_{Hgg} \sim \alpha_s / (12\pi v)$
- new particle with charge Q and SU(3) Casimir $C(R)$ [Reece]

$$R_\gamma = \frac{g_{H\gamma\gamma}}{g_{H\gamma\gamma}^{\text{SM}}} = \left[1 + 0.28\xi \left(1 \mp \sqrt{R_g} \right) \right]^2, \quad \xi = \frac{3Q^2}{C_2(R)}$$

⇒ end of a fourth chiral generation [Lenz et al]

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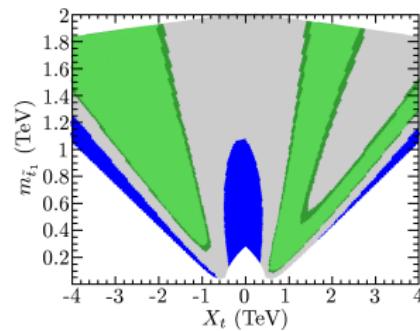
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Supersymmetry

- MSSM Higgs mass the best-predicted LHC observable [Hahn et al + Stal]
- production rates mix of form factor and D6 [e.g. Hollik, TP, Rauch, Rzezhak]
- stop mass/mixing crucial [$m_A = 1$ TeV, $\tan \beta = 20$]



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- SUSY particles in eff couplings [everyone]
stop mixing destructive [Reece]

$$\frac{g_{Hgg}}{g_{Hgg}^{\text{SM}}} = 1 + \frac{1}{4} \left(\frac{m_t^2}{m_{\tilde{t}_1}^2} + \frac{m_t^2}{m_{\tilde{t}_2}^2} - \frac{m_t^2 X_t^2}{m_{\tilde{t}_1}^2 m_{\tilde{t}_2}^2} \right)$$

– move towards NMSSM always an option...
 ⇒ no final verdict on the MSSM (ever?)

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Moving towards the Standard Model?

- should we worried by SM $\pm 20\%$?
- what is expected in BSM models [Gupta, Rzehak, Wells]

	ΔhVV	$\Delta h\bar{t}t$	Δhbb
mixed-in singlet	6%	6%	6%
composite Higgs	8%	tens of %	tens of %
MSSM	< 1%	3%	depends...

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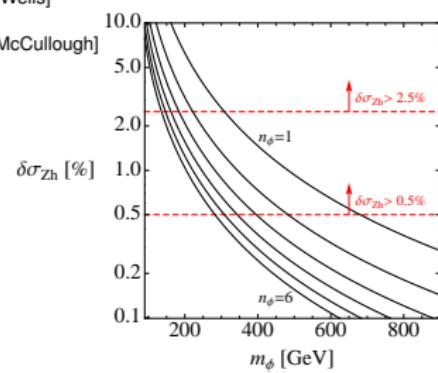
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- m_ϕ and λ_j given by hierarchy problem
- contributing to σ_{ZH} through self energy



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 - contributing to σ_{ZH} through self energy
- ⇒ trivial: want best possible precision

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What if it is essentially the Standard Model

- many theories decouple in Higgs sector [custodial symmetry, suppressed D6]
- any handle on high-scale physics needed

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Renormalization group

- Higgs mass related to self coupling: $m_H = v\sqrt{2\lambda}$
top mass related to Yukawa: $y_t = \sqrt{2}m_t/v$

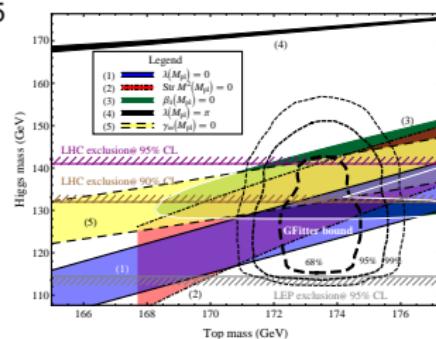
$$\frac{d\lambda}{d \log Q^2} = \frac{1}{16\pi^2} \left[12\lambda^2 + 6\lambda y_t^2 - 3y_t^4 - \frac{3}{2}\lambda \left(3g_2^2 + g_1^2 \right) + \frac{3}{16} \left(2g_2^4 + (g_2^2 + g_1^2)^2 \right) \right]$$

- IR fixed point for λ/y_t^2 fixing $m_H^2/m_t^2 = 1/2$ [with gravity: Shaposhnikov, Wetterich]

$$m_H = 126.3 + \frac{m_t - 171.2}{2.1} \times 4.1 - \frac{\alpha_s - 0.1176}{0.002} \times 1.5$$

- Planck-scale conditions [Holthausen, Lim, Lindner]

⇒ **Higgs and top crucial in combination**



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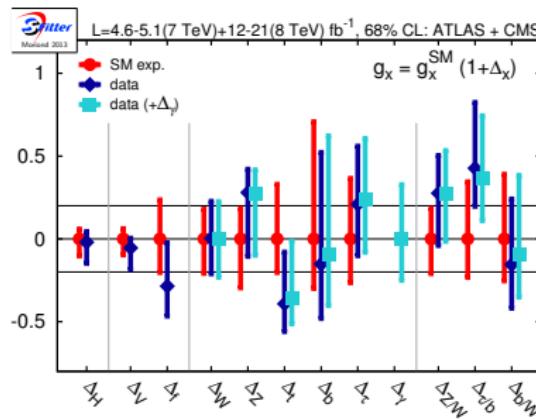
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Outlook

LHC Higgs program

- experimentalists: do not listen to theorists
- phenomenologists: enjoy fun problems
- theorists: give us some fun interpretations



Higgs Physics

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Markov chains

Probability maps [statistics unexpectedly hard...]

- honest LHC parameters: weak-scale Lagrangean [Higgs, MSSM, dark matter,...]
- likelihood map: data given a model $p(d|m) \sim |\mathcal{M}|^2(m)$
- Bayes' theorem: $p(m|d) = p(d|m) p(m)/p(d)$ [$p(d)$ normalization, $p(m)$ prejudice]

Markov chains

- problem in grid: huge phase space, find local best points?
problem in fit: domain walls, find global best points?
- construct ‘representative’ poll
- classical: representative set of spin states
compute average energy on this reduced sample
- BSM or Higgs: map $p(d|m)$ of parameter points
evaluate whatever you want
- Metropolis-Hastings
starting probability $p(d|m)$ vs suggested probability $p(d|m')$
 - 1– accept new point if $p(d|m') > p(d|m)$
 - 2– or accept with $p(d|m')/p(d|m) < 1$

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Markov chains

Weighted Markov chains [Lafaye, TP, Rauch, Zerwas; Ferrenberg, Swendsen]

- special situation
measure of ‘representative’: probability itself
- example with 2 bins, probability 9:1
10 entries needed for good Markov chain
2 entries needed if weight kept
- binning with weight would double count
bin with inverse averaging

$$P_{\text{bin}}(p \neq 0) = \frac{\text{bincount}}{\sum_{i=1}^{\text{bincount}} p^{-1}}$$

- good choice for $\mathcal{O}(6)$ dimensions

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Cooling Markov chains [Lafaye, TP, Rauch, Zerwas]

- zoom in on peak structures [inspired by simulated annealing]
- modified condition
Markov chain in partitions, numbered by j

$$p(d|m') > p(d|m) r^{10/j} \quad r \in [0, 1] \quad \text{random number}$$

- check for parameter coverage with many Markov chains
- \Rightarrow exclusive likelihood map first result

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