

EFT@LHC

Tilman Plehn

Higgs couplings

Higgs EFT

Top EFT

DM EFT

Higgs models

The Rise of Effective Lagrangians at the LHC

Tilman Plehn

Universität Heidelberg

Weizmann, June 2016

Theory in data-driven era

Same old theory motivation

- WIMP dark matter still best choice [Hooperon@Fermi]
- hierarchy problem (probably) a problem
- but: data in driving seat [750 GeV]

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- Lagrangian language established by Higgs discovery
 - 1- full new physics model [built to solve problems]
 - 2- simplified models [capturing experimental features, theoretically poor]
 - 3- effective field theory [symmetries and particles fixed, non-renormalizable operators]
- ⇒ matter of convenience and taste

	bottom-up EFT	simplified models	full models
agnostic	(x)		
data-driven		(x)	(x)
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agnostic	(x)	dishonest	pre-LHC
data-driven	boring	(x)	(x)
theory-driven	pointless	(x)	pre-LHC

Higgs couplings

Agnostic: why super-simple SM-Higgs sector? [SFitter]

- or: all couplings proportional to masses?
- assume: narrow CP-even scalar
Standard Model operators
- total production/decay rates only
- Lagrangian

$$\begin{aligned} \mathcal{L} = & \mathcal{L}_{\text{SM}} + \Delta_W g m_W H W^\mu W_\mu + \Delta_Z \frac{g}{2c_w} m_Z H Z^\mu Z_\mu - \sum_{\tau,b,t} \Delta_f \frac{m_f}{v} H (\bar{f}_R f_L + \text{h.c.}) \\ & + \Delta_g F_G \frac{H}{v} G_{\mu\nu} G^{\mu\nu} + \Delta_\gamma F_A \frac{H}{v} A_{\mu\nu} A^{\mu\nu} + \text{invisible} + \text{unobservable} \end{aligned}$$

- electroweak renormalizability through some UV completion
- QCD renormalizability not an issue

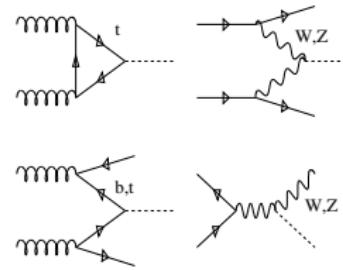
$$\begin{aligned} gg &\rightarrow H \\ qq &\rightarrow qqH \\ gg &\rightarrow ttH \\ qq' &\rightarrow VH \end{aligned}$$



$$g_{HXX} = g_{HXX}^{\text{SM}} (1 + \Delta_X)$$



$$\begin{aligned} H &\rightarrow ZZ \\ H &\rightarrow WW \\ H &\rightarrow b\bar{b} \\ H &\rightarrow \tau^+ \tau^- \\ H &\rightarrow \gamma\gamma \end{aligned}$$

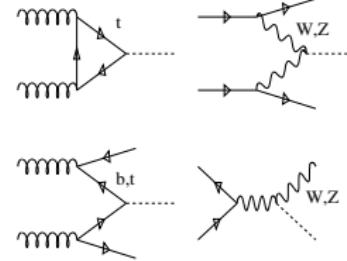


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Total width

- coupling extraction impossible without width assumption
- observed partial widths:

$$N = \sigma BR \propto \frac{g_p^2}{\sqrt{\Gamma_{\text{tot}}}} \frac{g_d^2}{\sqrt{\Gamma_{\text{tot}}}} \sim \frac{g^4}{g^2 \sum \frac{\Gamma_i(g^2)}{g^2} + \Gamma_{\text{unobs}}} \xrightarrow{g^2 \rightarrow 0} 0$$

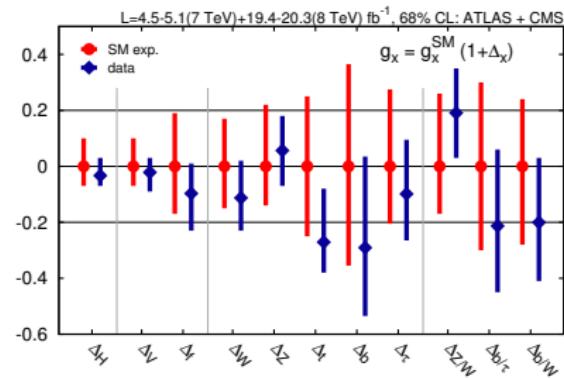
gives constraint from $\sum \Gamma_i(g^2) < \Gamma_{\text{tot}} \rightarrow \Gamma_H|_{\text{min}}$

- $WW \rightarrow WW$ unitarity: $g_{WWH} \lesssim g_{WWH}^{\text{SM}} \rightarrow \Gamma_H|_{\text{max}}$ [HiggsSignals]
- our assumption $\Gamma_{\text{tot}} = \sum_{\text{obs}} \Gamma_j$ [plus generation universality]

Higgs couplings after Run I

Run I legacy [Corbett, Eboli, Goncalves, Gonzalez-Fraile, Lopez-Val, TP, Rauch]

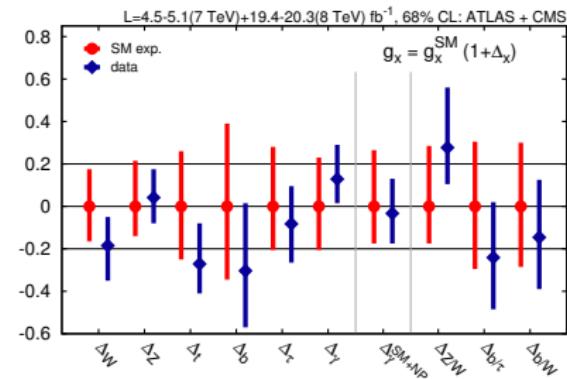
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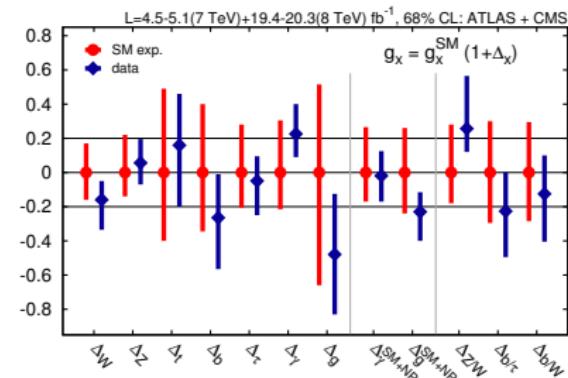
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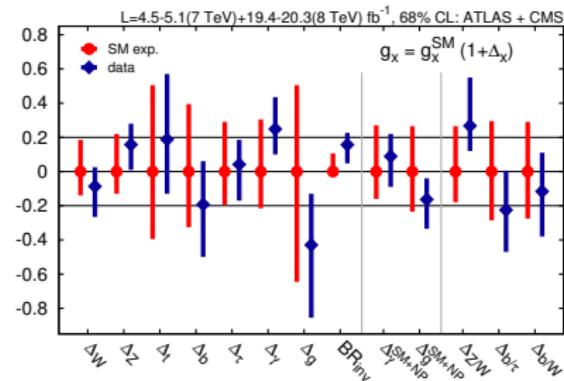
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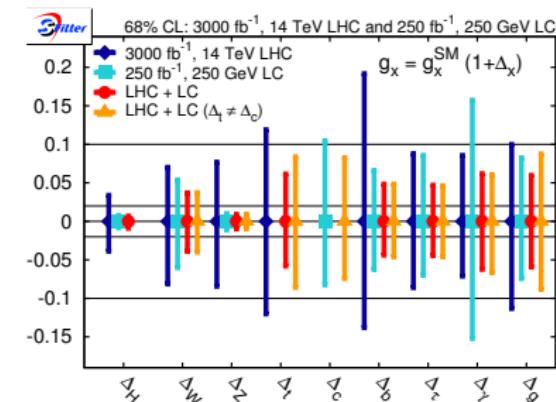
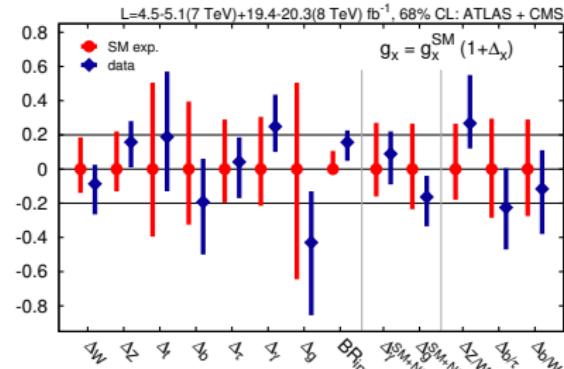
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Future [SFitter; Cranmer, Kreiss, Lopez-Val, TP]

- LHC extrapolations unclear
- systematic/theory uncertainties large
- e^+e^- linear collider much better
unobserved decays avoided
width measured from σ_{ZH}
 $H \rightarrow c\bar{c}$ accessible
invisible decays hugely improved
QCD theory error bars avoided



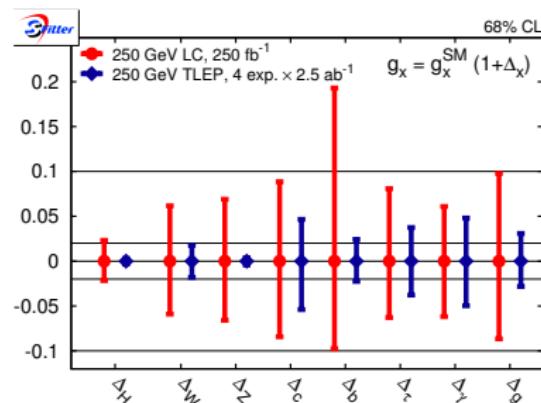
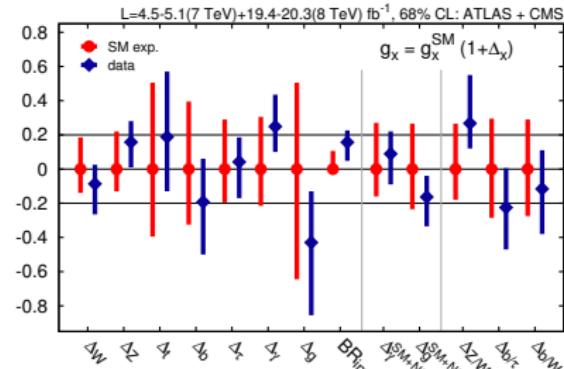
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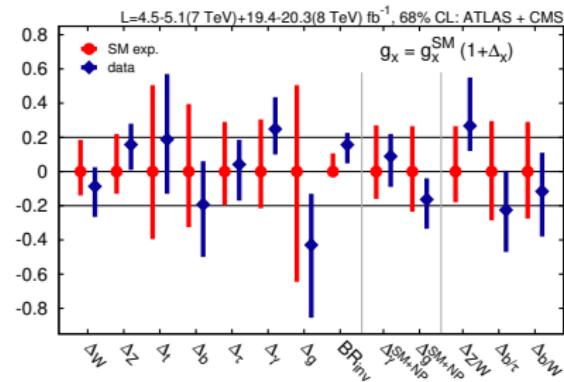
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- ⇒ Higgs factory case obvious



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Three major problems with approach

- 1– theory: no electroweak renormalizability
- 2– experiment: no kinematic distributions
- 3– phenomenology: no link to other sectors

D6 Higgs operators

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Higgs sector effective field theory [HISZ, polish, Eboli, Goncales-Garcia,...]

- set of Higgs operators [renormalizable, #1 solved]

$$\mathcal{O}_{GG} = \phi^\dagger \phi G_{\mu\nu}^a G^{a\mu\nu} \quad \mathcal{O}_{WW} = \phi^\dagger \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \phi \quad \mathcal{O}_{BB} = \dots$$

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- relevant part after equation of motion, etc

$$\mathcal{L}^{HVV} = - \frac{\alpha_S v}{8\pi} \frac{f_g}{\Lambda^2} \mathcal{O}_{GG} + \frac{f_{BB}}{\Lambda^2} \mathcal{O}_{BB} + \frac{f_{WW}}{\Lambda^2} \mathcal{O}_{WW} + \frac{f_B}{\Lambda^2} \mathcal{O}_B + \frac{f_W}{\Lambda^2} \mathcal{O}_W + \frac{f_{\phi,2}}{\Lambda^2} \mathcal{O}_{\phi,2}$$

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- plus Yukawa structure $f_{\tau,b,t}$

- 9 operators for Run I data

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- linked to Higgs couplings

$$\begin{aligned} g_g &= \frac{f_{GG} v}{\Lambda^2} \equiv -\frac{\alpha_s}{8\pi} \frac{f_g v}{\Lambda^2} & g_\gamma &= -\frac{g^2 v s_w^2}{2\Lambda^2} \frac{f_{BB} + f_{WW}}{2} \\ g_Z^{(1)} &= \frac{g^2 v}{2\Lambda^2} \frac{c_w^2 f_W + s_w^2 f_B}{2c_w^2} & g_W^{(1)} &= \frac{g^2 v}{2\Lambda^2} \frac{f_W}{2} \\ g_Z^{(2)} &= -\frac{g^2 v}{2\Lambda^2} \frac{s_w^4 f_{BB} + c_w^4 f_{WW}}{2c_w^2} & g_W^{(2)} &= -\frac{g^2 v}{2\Lambda^2} f_{WW} \\ g_Z^{(3)} &= M_Z^2 (\sqrt{2} G_F)^{1/2} \left(1 - \frac{v^2}{2\Lambda^2} f_{\phi,2} \right) & g_W^{(3)} &= M_W^2 (\sqrt{2} G_F)^{1/2} \left(1 - \frac{v^2}{2\Lambda^2} f_{\phi,2} \right) \\ g_f &= -\frac{m_f}{v} \left(1 - \frac{v^2}{2\Lambda^2} f_{\phi,2} \right) + \frac{v^2}{\sqrt{2}\Lambda^2} f_f \end{aligned}$$

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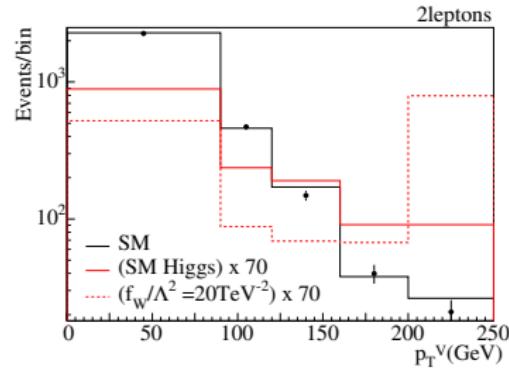
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- kinematics: $p_{T,V}, \Delta\phi_{jj}$ [#2 solved]



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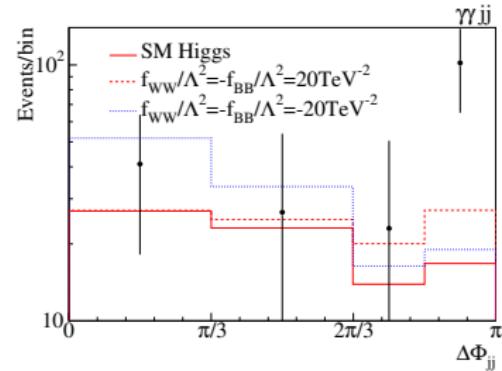
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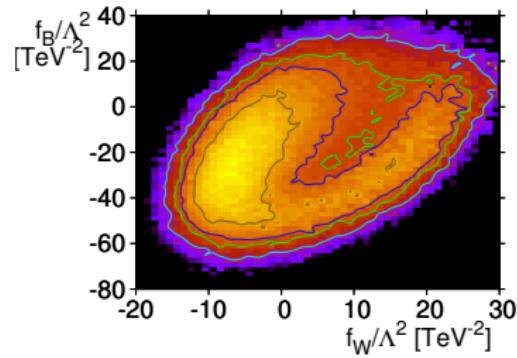
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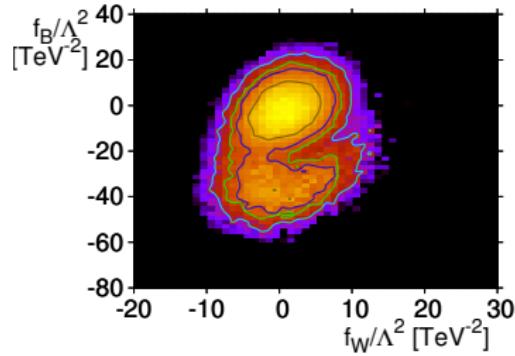
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$$\begin{aligned} \mathcal{O}_{GG} &= \phi^\dagger \phi G_{\mu\nu}^a G^{a\mu\nu} & \mathcal{O}_{WW} &= \phi^\dagger \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \phi & \mathcal{O}_{BB} &= \dots \\ \mathcal{O}_{BW} &= \phi^\dagger \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \phi & \mathcal{O}_W &= (D_\mu \phi)^\dagger \hat{W}^{\mu\nu} (D_\nu \phi) & \mathcal{O}_B &= \dots \\ \mathcal{O}_{\phi,1} &= (D_\mu \phi)^\dagger \phi \phi^\dagger (D^\mu \phi) & \mathcal{O}_{\phi,2} &= \frac{1}{2} \partial^\mu (\phi^\dagger \phi) \partial_\mu (\phi^\dagger \phi) \\ \mathcal{O}_{\phi,3} &= \frac{1}{3} (\phi^\dagger \phi)^3 & \mathcal{O}_{\phi,4} &= (D_\mu \phi)^\dagger (D^\mu \phi) (\phi^\dagger \phi) \end{aligned}$$

Run 1 legacy

- kinematics: $p_T, v, \Delta\phi_{jj}$ [#2 solved]
- with impact...
- ...in last bin



D6 Higgs operators

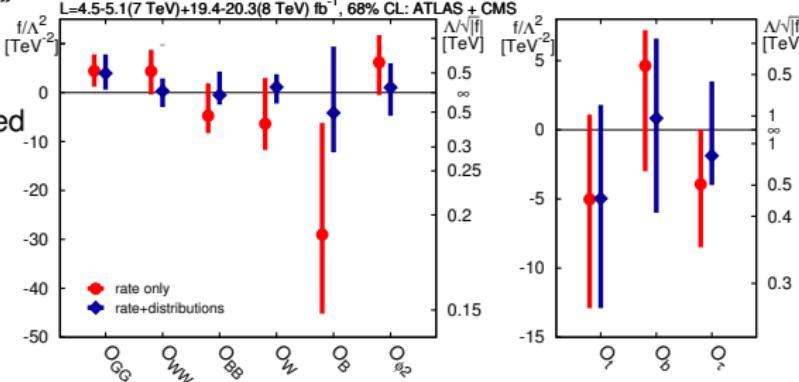
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Run 1 legacy

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 - with impact...
 - ...in last bin
 - Run I sensitivity limited
- ⇒ consistent EFT???



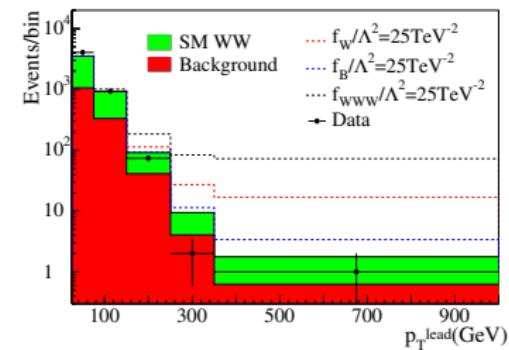
D6 Higgs-gauge operators

Triple gauge couplings

- one more Higgs-gauge operator [#3 solved]

$$\mathcal{O}_W = (D_\mu \phi)^\dagger \hat{W}^{\mu\nu} (D_\nu \phi) \quad \mathcal{O}_B = (D_\mu \phi)^\dagger \hat{B}^{\mu\nu} (D_\nu \phi) \quad \mathcal{O}_{WWW} = \text{Tr} \left(\hat{W}_{\mu\nu} \hat{W}^{\nu\rho} \hat{W}_\rho^\mu \right)$$

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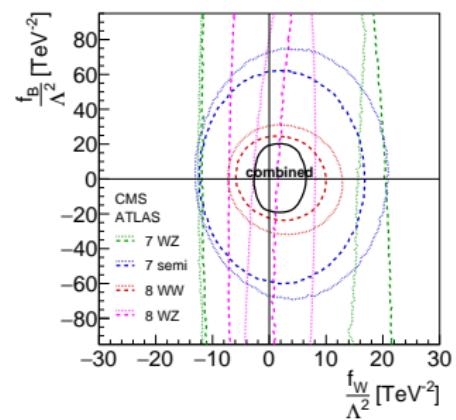
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- kinematics: $p_{T,\ell}$ in VV production
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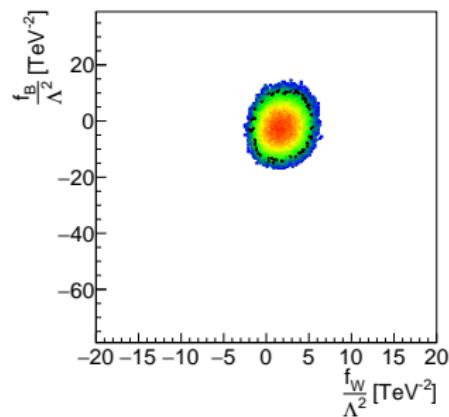
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- affecting correlations



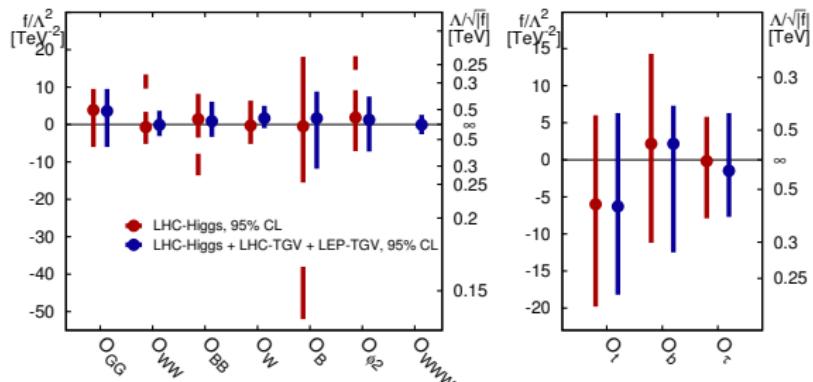
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- ⇒ complete Higgs-gauge analysis



D6 Higgs-gauge operators

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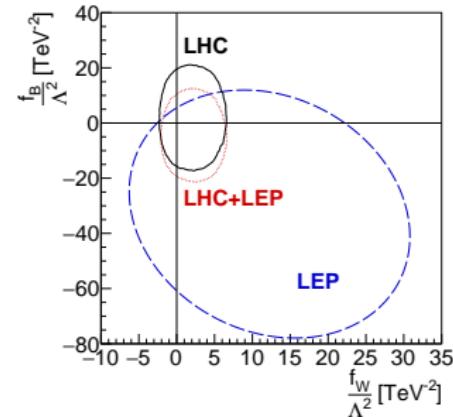
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- ⇒ complete Higgs-gauge analysis

LHC vs LEP

- triple gauge vertices g_1, κ, λ vs operators
 - semileptonic analyses missing for 8 TeV
- ⇒ Run I LHC beating LEP



Higgs couplings

Higgs EFT

Top EFT

DM EFT

Higgs models

Exercise: higher-dimensional operators

Higgs sector including dimension-6 operators

$$\mathcal{L}_{D6} = \sum_{i=1}^2 \frac{f_i}{\Lambda^2} \mathcal{O}_i \quad \text{with} \quad \mathcal{O}_{\phi,2} = \frac{1}{2} \partial_\mu (\phi^\dagger \phi) \partial^\mu (\phi^\dagger \phi), \quad \mathcal{O}_{\phi,3} = -\frac{1}{3} (\phi^\dagger \phi)^3$$

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first operator, wave function renormalization

$$\mathcal{O}_{\phi,2} = \frac{1}{2} \partial_\mu (\phi^\dagger \phi) \partial^\mu (\phi^\dagger \phi) = \frac{1}{2} (\tilde{H} + v)^2 \partial_\mu \tilde{H} \partial^\mu \tilde{H}$$

proper normalization of combined kinetic term [LSZ]

$$\mathcal{L}_{\text{kin}} = \frac{1}{2} \partial_\mu \tilde{H} \partial^\mu \tilde{H} \left(1 + \frac{f_{\phi,2} v^2}{\Lambda^2} \right) \stackrel{!}{=} \frac{1}{2} \partial_\mu H \partial^\mu H \quad \Leftrightarrow \quad H = \tilde{H} \sqrt{1 + \frac{f_{\phi,2} v^2}{\Lambda^2}}$$

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second operator, minimum condition giving v

$$v^2 = -\frac{\mu^2}{\lambda} - \frac{f_{\phi,3} \mu^4}{4 \lambda^3 \Lambda^2}$$

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both operators contributing to Higgs mass

$$\begin{aligned} \mathcal{L}_{\text{mass}} &= -\frac{\mu^2}{2} \tilde{H}^2 - \frac{3}{2} \lambda v^2 \tilde{H}^2 - \frac{f_{\phi,3}}{\Lambda^2} \frac{15}{24} v^4 \tilde{H}^2 \stackrel{!}{=} -\frac{m_H^2}{2} H^2 \\ &\Leftrightarrow m_H^2 = 2\lambda v^2 \left(1 - \frac{f_{\phi,2} v^2}{\Lambda^2} + \frac{f_{\phi,3} v^2}{2\Lambda^2 \lambda} \right) \end{aligned}$$

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$$\begin{aligned} \mathcal{L}_{\text{self}} = & -\frac{m_H^2}{2v} \left[\left(1 - \frac{f_{\phi,2} v^2}{2\Lambda^2} + \frac{2f_{\phi,3} v^4}{3\Lambda^2 m_H^2} \right) H^3 - \frac{2f_{\phi,2} v^2}{\Lambda^2 m_H^2} H \partial_\mu H \partial^\mu H \right] \\ & - \frac{m_H^2}{8v^2} \left[\left(1 - \frac{f_{\phi,2} v^2}{\Lambda^2} + \frac{4f_{\phi,3} v^4}{\Lambda^2 m_H^2} \right) H^4 - \frac{4f_{\phi,2} v^2}{\Lambda^2 m_H^2} H^2 \partial_\mu H \partial^\mu H \right] \end{aligned}$$

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alternatively, strong multi-Higgs interactions

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⇒ operators and distributions linked to poor UV behavior

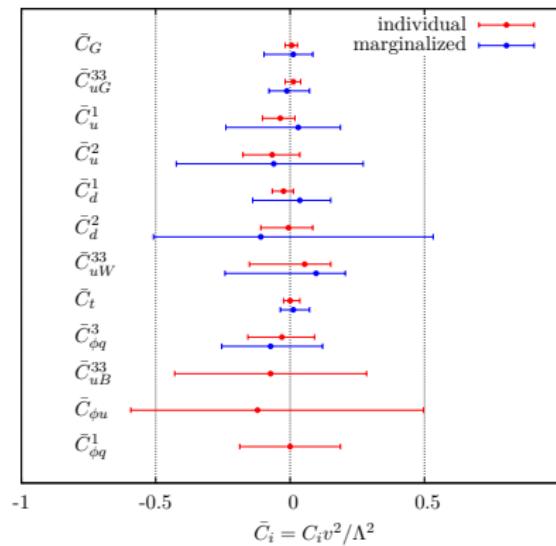
D6 top operators

Same for tops [TopFitter: Buckley, Englert, Ferrando, Miller, Moore, Russell, White]

- single, pair-wise, and associated top production [plus decays]
- including anomalous A_{FB} from Tevatron
- 4-quark, Yang-Mills, electroweak operators

$$\mathcal{O}_{qq} = \bar{q}\gamma_\mu q \bar{t}\gamma^\mu t \quad \mathcal{O}_G = f_{ABC} G_\mu^{A\nu} G_\nu^{B\lambda} G_\lambda^{C\mu} \quad \mathcal{O}_{\phi G} = \phi^\dagger \phi G_{\mu\nu}^a G^{a\mu\nu} \dots$$

- profile likelihoods and individual limits
- \Rightarrow generic D6 reach ~ 500 GeV $[C = 1]$



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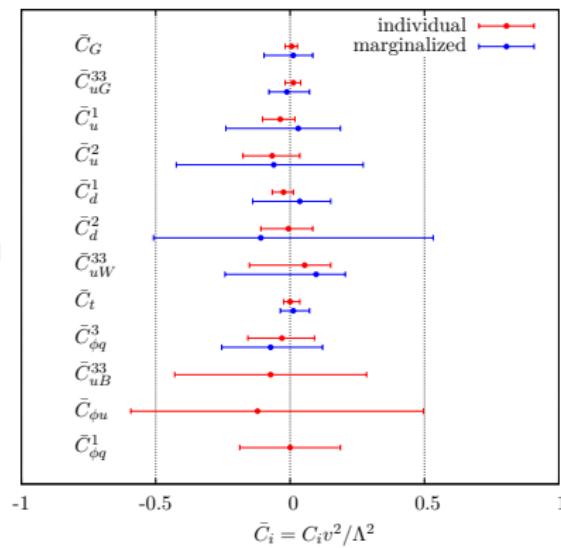
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For theorists: in terms of models

- axigluon: $M_A > 1.4$ TeV [$t\bar{t}$ resonance]
 - SM-like W' : $M_{W'} > 1.2$ TeV [t -channel,...]
- \Rightarrow models less sensitive to correlations



D6 dark matter operators

Combining direct, indirect, collider results for WIMPs [Tait et al]

- choose dark matter candidate [Majorana/Dirac fermion, scalar, dark photon]
- consider D6 scattering process $\chi\chi \rightarrow \text{SM SM}$
- relic density from annihilation [$m_\chi / T \sim 30$]
- indirect detection even later
- direct detection non-relativistic [$E \sim 10 \text{ MeV}$]
- LHC tricky: single scale $m_\chi \ll m_{\text{mediator}}$?
- example: scalar dark matter

Label	Coefficient	Operator	$\sigma_{\text{SI}} \langle \sigma \text{ann} v \rangle$
Real scalar			
R1	$\lambda_1 \sim 1/(2M^2)$	$m_q \chi^2 \bar{q} q$	✓ s-wave
R2	$\lambda_2 \sim 1/(2M^2)$	$i m_q \chi^2 \bar{q} \gamma^5 q$	s-wave
R3	$\lambda_3 \sim \alpha_s/(4M^2) \chi^2 G_{\mu\nu} G^{\mu\nu}$		✓ s-wave
R4	$\lambda_4 \sim \alpha_s/(4M^2) i \chi^2 G_{\mu\nu} \tilde{G}^{\mu\nu}$		s-wave
Complex scalar			
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C3	$\lambda_3 \sim 1/(M^2)$	$\chi^\dagger \partial_\mu \chi \bar{q} \gamma^\mu q$	✓ p-wave
C4	$\lambda_4 \sim 1/(M^2)$	$\chi^\dagger \partial_\mu \chi \bar{q} \gamma^\mu \gamma^5 q$	p-wave
C5	$\lambda_5 \sim \alpha_s/(8M^2) \chi^\dagger \chi G_{\mu\nu} G^{\mu\nu}$		✓ s-wave
C6	$\lambda_6 \sim \alpha_s/(8M^2) i \chi^\dagger \chi G_{\mu\nu} \tilde{G}^{\mu\nu}$		s-wave

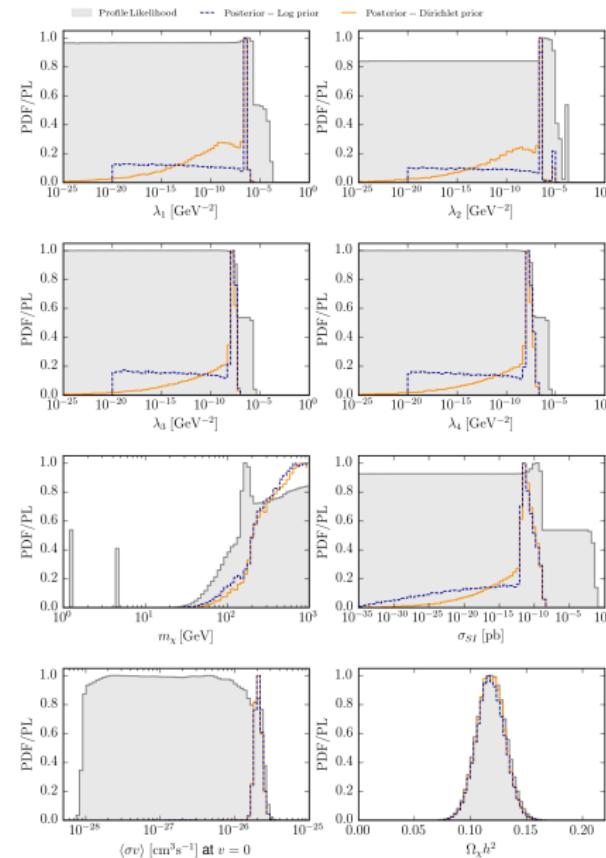
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Relic density plus Hooperon [Liem, Bertone, Calore, Ruiz de Austri, Tait, Trotta, Weniger]

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- profile likelihood
- flat prior on $\log \lambda_i$ [prior $1/\lambda_i$]
- Dirichlet prior preferring similar-sized Wilson coefficients



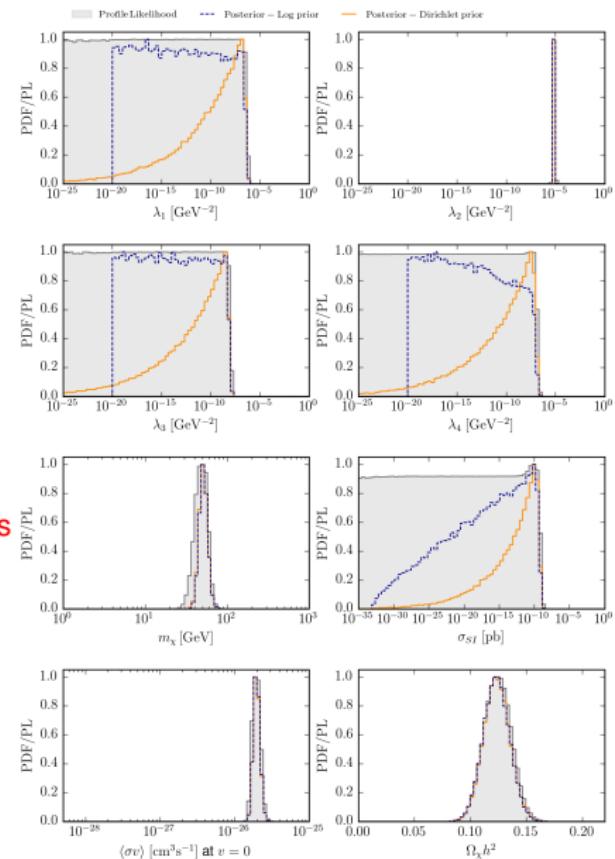
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- profile likelihood
- flat prior on $\log \lambda_i$ [prior $1/\lambda_i$]
- Dirichlet prior preferring similar-sized Wilson coefficients
- Fermi: GCE plus dwarf galaxies
- ⇒ with data, the method hardly matters



Higgs couplings

Higgs EFT

Top EFT

DM EFT

Higgs models

(Simplified) scalar/gauge extensions

Higgs singlet/doublet extensions [Higgs portal]

- one or more new (pseudo-) scalars
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Scalar top partners, non-Higgs [simplified supersymmetry]

- Lagrangian with scalar top partner, singlet plus doublet

$$\begin{aligned} \mathcal{L} \supset & (D_\mu \tilde{Q})^\dagger (D^\mu \tilde{Q}) + (D_\mu \tilde{t}_R)^* (D^\mu \tilde{t}_R) - \tilde{Q}^\dagger M^2 \tilde{Q} - M^2 \tilde{t}_R^* \tilde{t}_R \\ & - \kappa_{LL} (\phi \cdot \tilde{Q})^\dagger (\phi \cdot \tilde{Q}) - \kappa_{RR} (\tilde{t}_R^* \tilde{t}_R) (\phi^\dagger \phi) - \left[\kappa_{LR} M \tilde{t}_R^* (\phi \cdot \tilde{Q}) + \text{h.c.} \right] \end{aligned}$$

- contribution through loops all over Higgs-gauge sector

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- contribution through loops all over Higgs-gauge sector

Triplet gauge extension [whatever that becomes in the UV]

- additional vector triplet field V_μ

$$\begin{aligned} \mathcal{L} \supset & -\frac{1}{4} \tilde{V}_{\mu\nu}^a \tilde{V}^{\mu\nu a} + \frac{M_V^2}{2} \tilde{V}_\mu^a \tilde{V}^{\mu a} + i \frac{g_V}{2} c_H \tilde{V}_\mu^a \left[\phi^\dagger \sigma^a \overleftrightarrow{D}^\mu \phi \right] + \frac{g_w^2}{2g_V} \tilde{V}_\mu^a \sum_{\text{fermions}} c_F \bar{F}_L \gamma^\mu \sigma^a F_L \\ & + \frac{g_V}{2} c_{VVV} \epsilon_{abc} \tilde{V}_\mu^a \tilde{V}_\nu^b D^{[\mu} \tilde{V}^{\nu]}{}^c + g_V^2 c_{VWH} \tilde{V}_\mu^a \tilde{V}^{\mu a} (\phi^\dagger \phi) - \frac{g_w}{2} c_{VWW} \epsilon_{abc} W^{\mu\nu} \tilde{V}_\mu^b \tilde{V}_\nu^c \end{aligned}$$

- new states, mixing with W^\pm and Z
weak gauge coupling to W, Z mass eigenstates

Higgs D6 breakdown

D6-Lagrangian breakdown [Brehmer, Freitas, Lopez-Val, TP]

- phenomenology: does D6 capture all model features at LHC?
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 - LHC simulations: D6-Lagrangian vs full model
production: WBF, VH, HH
decays: $H \rightarrow \gamma\gamma, 4\ell$
 - check where differences appear at 13 TeV
kinematic distributions like $p_{T,j}$ or m_{VH} ?
resonance peaks of new states?

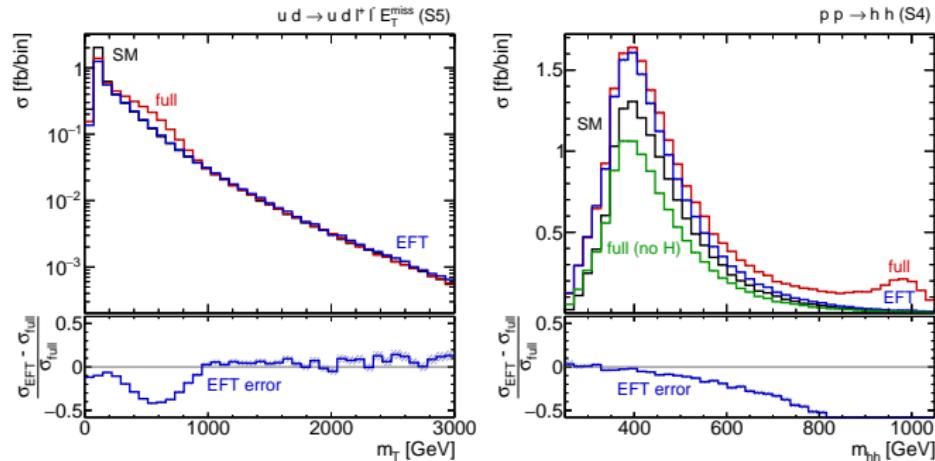
Higgs D6 breakdown

Higgs portal

- testable benchmarks for LHC

Singlet				EFT			EFT (v -improved)	
m_H	$\sin \alpha$	v_S/v	$\Delta_x^{\text{singlet}}$	Λ	\bar{c}_H	Δ_x^{EFT}	\bar{c}_H	Δ_x^{EFT}
500	0.2	10	-0.020	491	0.036	-0.018	0.040	-0.020
350	0.3	10	-0.046	336	0.073	-0.037	0.092	-0.046
200	0.4	10	-0.083	190	0.061	-0.031	0.167	-0.083
1000	0.4	10	-0.083	918	0.183	-0.092	0.167	-0.092
500	0.6	10	-0.200	407	0.461	-0.231	0.400	-0.200

- effects in WBF and hh



Higgs D6 breakdown

Higgs portal

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2HDM

- testable benchmarks for LHC

Type	$\tan \beta$	2HDM					EFT		
		α/π	m_{12}	m_{H^0}	m_{A^0}	m_{H^\pm}	$ \Lambda [\text{GeV}]$	\bar{c}_u	$\bar{c}_{d,\ell}$
I	1.5	-0.086	45	230	300	350	100	-0.744	-0.744
II	15	-0.023	116	449	450	457	448	0.000	0.065
II	10	0.032	157	500	500	500	99	0.465	-46.5
I	20	0	45	200	500	500	142	0.003	0.003

Higgs D6 breakdown

Higgs portal

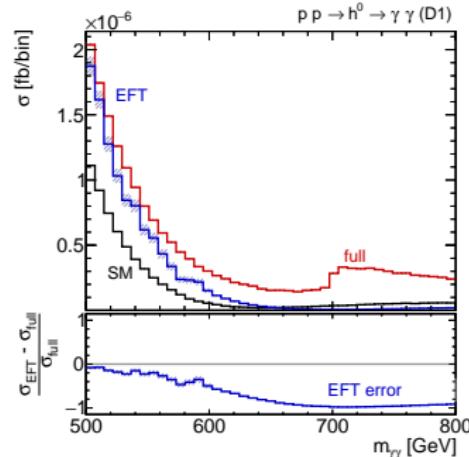
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- effects in $H \rightarrow \gamma\gamma$



Higgs D6 breakdown

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Top partners

- testable benchmarks for LHC

Scalar top-partner model						EFT		
M	κ_{LL}	κ_{RR}	κ_{LR}	$m_{\tilde{t}_1}$	$m_{\tilde{t}_2}$	\bar{c}_H	\bar{c}_W	\bar{c}_{HW}
500	-1.16	2.85	0.147	500	580	$6.22 \cdot 10^{-3}$	$-3.11 \cdot 10^{-7}$	$3.99 \cdot 10^{-7}$
350	-3.16	-2.82	0.017	173	200	$4.30 \cdot 10^{-3}$	$-2.55 \cdot 10^{-4}$	$2.55 \cdot 10^{-4}$
500	-7.51	-7.17	0.012	173	200	$1.66 \cdot 10^{-2}$	$-2.97 \cdot 10^{-4}$	$2.97 \cdot 10^{-4}$

Higgs D6 breakdown

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2HDM

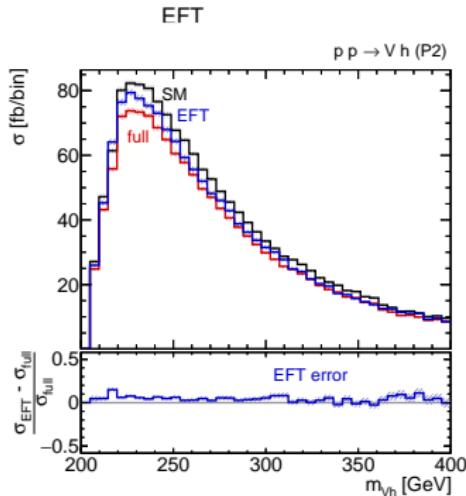
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Higgs couplings

Higgs EFT

Top EFT

DM EFT

Higgs models

Higgs D6 breakdown

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Top partners

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Vector triplet [Brehmer, Biekötter, Krämer, TP]

- testable benchmarks for LHC

Triplet model						EFT			
M_V	g_V	c_H	c_F	c_{VVHH}	m_ξ	\bar{c}_W	\bar{c}_H	\bar{c}_6	\bar{c}_f
591	3.0	-0.47	-5.0	2.0	1200	-0.044	0.000	0.000	0.000
946	3.0	-0.47	-5.0	1.0	1200	-0.017	0.000	0.000	0.000
941	3.0	-0.28	3.0	1.0	1200	0.006	0.075	0.100	0.025
1246	3.0	-0.50	3.0	-0.2	1200	0.006	0.103	0.138	0.034
846	1.0	-0.56	-1.32	0.08	849	-0.007	-0.020	-0.027	-0.007

Higgs D6 breakdown

Higgs portal

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2HDM

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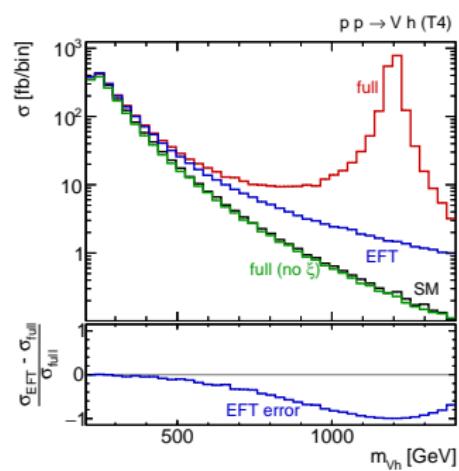
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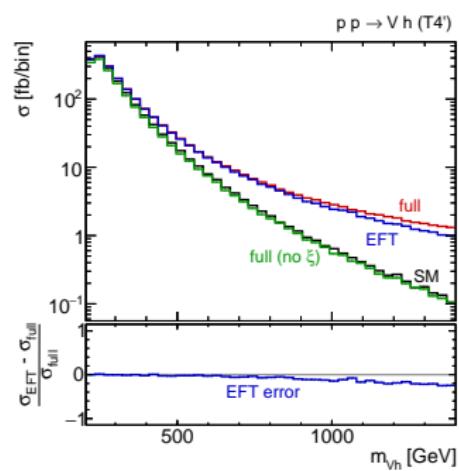
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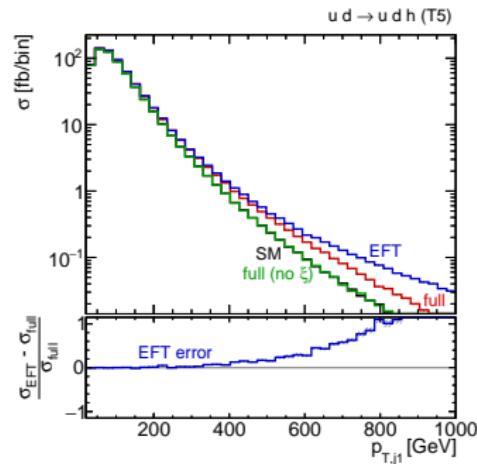
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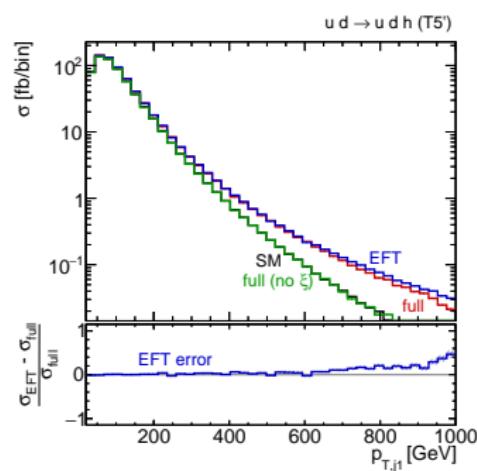
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Higgs D6 breakdown

Reasons for D6-breakdown in Higgs sector at LHC

Model	Process	EFT failure		
		resonance	kinematics	matching
singlet	on-shell $h \rightarrow 4\ell$, WBF, Vh , ...			X
	off-shell WBF, ...		(x)	X
2HDM	hh	X	X	X
	on-shell $h \rightarrow 4\ell$, WBF, Vh , ...			X
top partner	off-shell $H \rightarrow \gamma\gamma$, ...		(x)	X
	hh	X	X	X
vector triplet	WBF, Vh			X
	WBF		(x)	X
	Vh	X	(x)	X

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	off-shell $H \rightarrow \gamma\gamma$, ...		(x)	X
	hh	X	X	X
top partner	WBF, Vh			X
vector triplet	WBF		(x)	X
	Vh	X	(x)	X

Lessons from Higgs sector

- start with D6 description [data-driven era of particle physics]
 - EFT expansion in E/Λ known to be dodgy
 - test D6 in comparison to (simplified) models
 - all relevant effect at tree level
 - resonance peaks the key feature
- ⇒ D6 limitations not from matter-of-principle arguments

Questions

Questions waiting to be answered

- is it really the Standard Model Higgs? [No]
- is there WIMP dark matter? [Yes]
- is there TeV-scale physics beyond the Standard Model? [Yes]
- are EFT analyses boring? [Yes]
- will we stop EFT analyses once we find new states [Definitely]
⇒ **welcome to a data-driven era!**

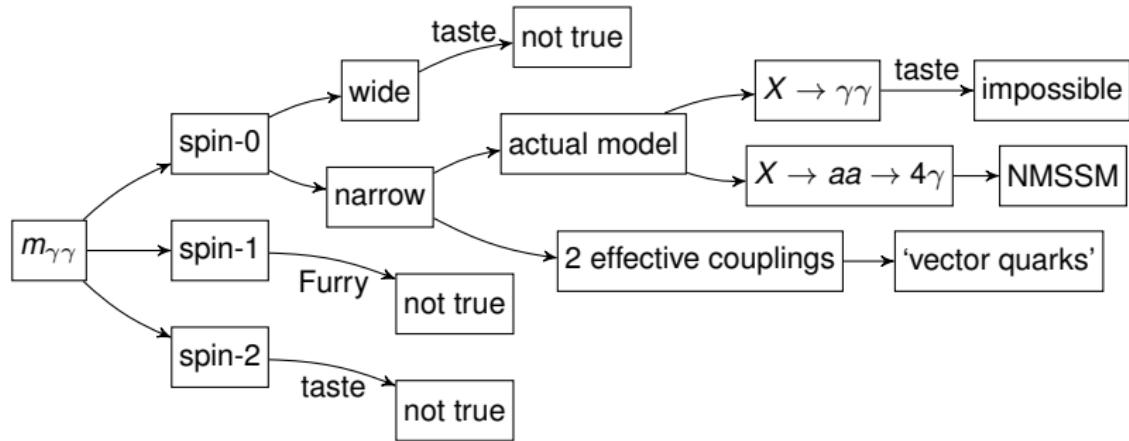
Lectures on LHC Physics and dark matter updated under www.thphys.uni-heidelberg.de/~plehn/

Much of this work was funded by the BMBF Theorie-Verbund which is ideal for relevant LHC work



Bundesministerium
für Bildung
und Forschung

750 GeV — the finger to particle theory



Effective theories in action?

- key question: another Higgs scalar?
 - dimension-5 operators $XG^{\mu\nu}G_{\mu\nu}$ and $XA^{\mu\nu}A_{\mu\nu}$ → avoid di-jet constraints
 - gauge invariant $XB^{\mu\nu}B_{\mu\nu}$ and $XW^{\mu\nu}W_{\mu\nu}$ → avoid VV constraints
 - no clear link to other data
- ⇒ everyone writing models papers!