

Tilman Plehn

CP vs \hat{T}

SMEFT

Information

WBF

ZH

Decays

Comparison

Higgs CP Through Information Geometry

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Freiburg, July 2018

CP symmetry tests

CP vs \hat{T}

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Comparison

C and P and T and \hat{T}

- transformations on state with spin/momentum [review: Valencia]

$$C |\phi(p, s)\rangle = |\phi^*(p, s)\rangle \quad P |\phi(p, s)\rangle = \eta_\phi |\phi(-p, s)\rangle \quad T |\phi(p, s)\rangle = \langle \phi(-p, -s)|$$

- transformation of complex scalar

$$C\phi(t, \vec{x})C^{-1} = \eta_C \phi^*(t, \vec{x}) \quad P\phi(t, \vec{x})P^{-1} = \eta_P \phi(t, -\vec{x}) \quad T\phi(t, \vec{x})T^{-1} = \phi(-t, \vec{x})$$

- CPT symmetry generally assumed, T proxy for CP
- naive time reversal \hat{T} avoiding initial \leftrightarrow final state

$$\hat{T} |\phi(p, s)\rangle = |\phi(-p, -s)\rangle$$

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- CPT symmetry generally assumed, T proxy for CP
- naive time reversal \hat{T} avoiding initial \leftrightarrow final state

$$\hat{T} |\phi(p, s)\rangle = |\phi(-p, -s)\rangle$$

- genuine U-odd is what we want [$U = C, P, \hat{T}$, Atwood, Bar-Shalom, Eilam]

$$\langle O \rangle_{\mathcal{L}=U\mathcal{L}U^{-1}} = 0 .$$

- U-odd is what we usually use

$$O(U|i\rangle \rightarrow U|f\rangle) \stackrel{\text{odd}}{=} -O(|i\rangle \rightarrow |f\rangle) \xrightarrow{p(|i\rangle)=p(U|i\rangle)} \langle O \rangle_{\mathcal{L}=U\mathcal{L}U^{-1}} = 0 .$$

- finite genuine \hat{T} -odd measurement means CP-violating theory, provided

- 1- phase space \hat{T} -symmetric
- 2- initial state distribution invariant under \hat{T}
- 3- no re-scattering

\Rightarrow use \hat{T} as proxy to CP

Naive time reversal at LHC

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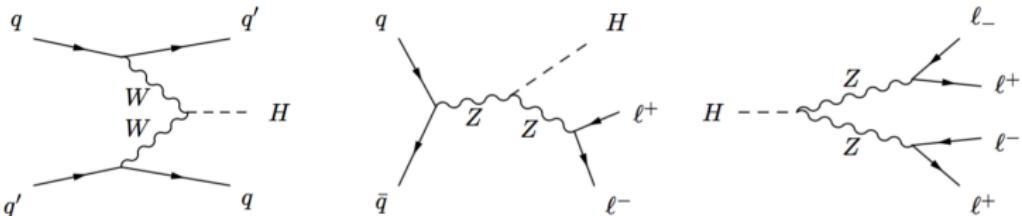
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Comparison

Processes with same vertex/information



- four 4-momenta defining 10 observables [Han, Li,...]
 - four scalar products giving external masses
 - four C -even, P -even, and \hat{T} -even scalar products
 - two C -odd, P -even, and \hat{T} -even scalar products
 - plus** C -even, P -odd, and \hat{T} -odd observable from Levi-Civita-tensor
 - 1– CP -odd and \hat{T} -odd
 - for $q\bar{q} \rightarrow ZH$ also genuine CP -odd and genuine \hat{T} -odd
 - non-zero expectation value means CP violation
 - 2– CP -odd and \hat{T} -even
 - for $q\bar{q} \rightarrow ZH$ also genuine CP -odd
 - if theory CP -violating, CP -expectation value should be non-zero
 - without re-scattering, \hat{T} -expectation value zero
 - with re-scattering or complex phase $\langle O \rangle$ matches symmetry
- ⇒ **focus on one CP -odd and \hat{T} -odd observable**

CP in SMEFT Lagrangian

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Dimension-6 Lagrangian

- for EFT fit choose symmetries first
- CP-conserving couplings [defining Higgs properties]

$$\mathcal{O}_B = i \frac{g}{2} (D^\mu \phi^\dagger)(D^\nu \phi) B_{\mu\nu} \quad \mathcal{O}_W = i \frac{g}{2} (D^\mu \phi)^\dagger \sigma^k (D^\nu \phi) W_{\mu\nu}^k$$

$$\mathcal{O}_{BB} = -\frac{g'^2}{4} (\phi^\dagger \phi) B_{\mu\nu} B^{\mu\nu} \quad \mathcal{O}_{WW} = -\frac{g^2}{4} (\phi^\dagger \phi) W_{\mu\nu}^k W^{\mu\nu k}$$

$$\mathcal{O}_{\phi,2} = \frac{1}{2} \partial^\mu (\phi^\dagger \phi) \partial_\mu (\phi^\dagger \phi)$$

- CP-violating couplings [defining CP-violation]

$$\mathcal{O}_{B\bar{B}} = -\frac{g'^2}{4} (\phi^\dagger \phi) \tilde{B}_{\mu\nu} B^{\mu\nu} \quad \mathcal{O}_{W\bar{W}} = -\frac{g^2}{4} (\phi^\dagger \phi) \tilde{W}_{\mu\nu}^k W^{\mu\nu k}$$

- dimension six means non-SM momentum dependence
- link to loop-induced coupling: $\text{Tr}(\gamma_\mu \gamma_\nu \gamma_\rho \gamma_\sigma \gamma_5) \rightarrow \epsilon_{\mu\nu\rho\sigma}$
- complex phases [re-scattering]
 - typical from absorptive diagrams with light particles
 - not foreseen in EFT approach
 - mimic with complex CP-conserving Wilson coefficients, lacking better idea
- ⇒ all established and known

Information geometry

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Comparison

Quantify what is there to learn

- covariance matrix [measurement error in model space \mathbf{g}]

$$C_{ij}(\mathbf{g}) \equiv E [(\hat{g}_i - \bar{g}_i)(\hat{g}_j - \bar{g}_j) | \mathbf{g}]$$

- Fisher information [sensitivity in model space]

$$I_{ij}(\mathbf{g}) \equiv -E \left[\frac{\partial^2 \log f(\mathbf{x}|\mathbf{g})}{\partial g_i \partial g_j} \middle| \mathbf{g} \right]$$

- phase space distribution [phase space \mathbf{x} , additive]

$$I_{ij} = \frac{L}{\sigma} \partial \sigma g_i \partial \sigma g_j - L \sigma E \left[\frac{\partial^2 \log f^{(1)}(\mathbf{x}|\mathbf{g})}{\partial g_i \partial g_j} \right]$$

- Cramér-Rao bound defining lowest possible covariance

$$C_{ij}(\mathbf{g}) \geq (I^{-1})_{ij}(\mathbf{g})$$

- model-space distance [probability to measure \mathbf{g}_b with true \mathbf{g}_a]

$$d(\mathbf{g}_b; \mathbf{g}_a) = \sqrt{(\mathbf{g}_a - \mathbf{g}_b)_i I_{ij}(\mathbf{g}_a) (\mathbf{g}_a - \mathbf{g}_b)_j}$$

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- 1– parametrization-invariant ellipses of constant distance/reach in model space
 - 2– diagonalize I_{ij} , define model-space eigenvectors
 - 3– compute information in distributions or phase space regions
- ⇒ best tool to understand modern analyses

WBF production

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Testing CP in WBF

- four external 4-momenta \rightarrow 10 scalar products

four external masses [zero]

four C -even, P -even, \hat{T} -even [not interesting]

two C -odd, P -even, \hat{T} -even [not possible]

one C -even, P -odd, \hat{T} -odd [yeah]

$$\epsilon_{\mu\nu\rho\sigma} k_1^\mu k_2^\nu q_1^\rho q_2^\sigma \rightarrow O \equiv \epsilon_{\mu\nu\rho\sigma} k_1^\mu k_2^\nu q_1^\rho q_2^\sigma \text{ sign}[(k_1 - k_2) \cdot (q_1 - q_2)]$$

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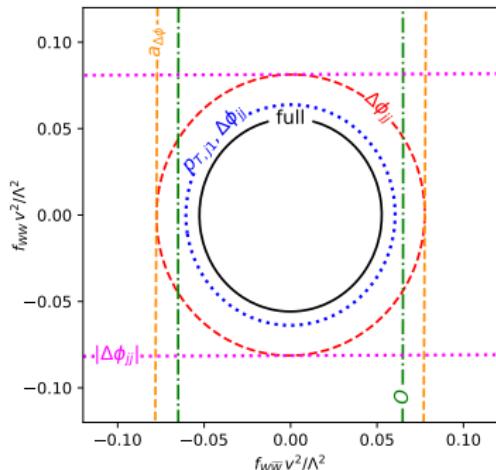
- azimuthal angle difference [lab frame]

$$O = 2E_- (\vec{q}_- \times \vec{q}_+) \cdot \vec{k}_+ \rightarrow \sin \Delta\phi_{jj}$$

- CP asymmetry

$$a_{\Delta\phi_{jj}} \equiv \frac{d\sigma(\Delta\phi_{jj}) - d\sigma(-\Delta\phi_{jj})}{d\sigma(\Delta\phi_{jj}) + d\sigma(-\Delta\phi_{jj})}$$

- separating dimension-6 effects
- check with imaginary Wilson coefficients
- ⇒ testing CP , but assuming no re-scattering



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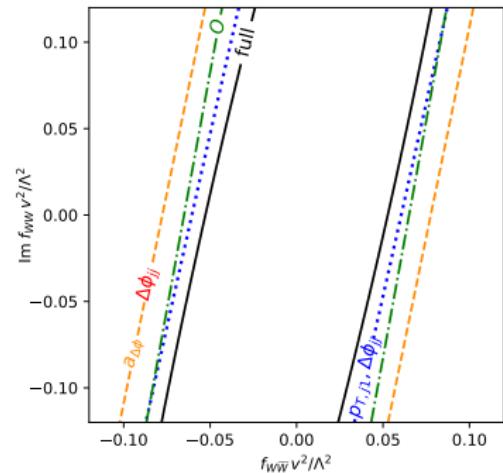
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Testing CP in ZH production

- same 10 scalar products as for WBF
- CP -odd and \hat{T} -odd angle as for WBF

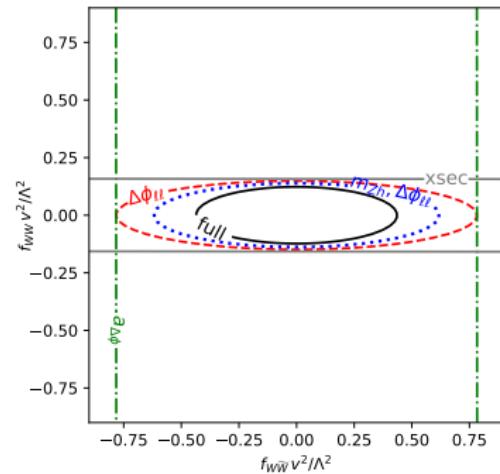
$$O_1 = \epsilon_{\mu\nu\rho\sigma} k_1^\mu k_2^\nu q_{\ell+}^\rho q_{\ell-}^\sigma \text{ sign}((k_1 - k_2) \cdot (q_1 - q_2)) \rightarrow \sin \Delta\phi_{\ell\ell}$$

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- separating dimension-6 effects

\Rightarrow testing CP without assumptions [to leading order]



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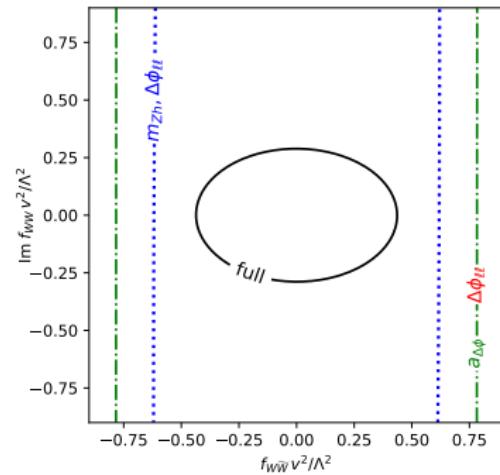
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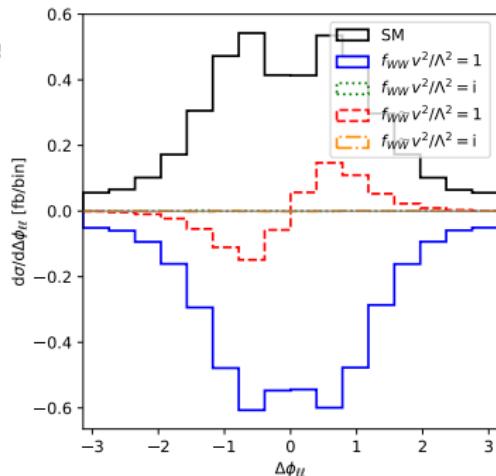
\Rightarrow testing CP without assumptions [to leading order]

- CP-odd and \hat{T} -even, requiring second phase

$$O_2 \rightarrow \Delta E_{\ell\ell}$$

$$O_3 \rightarrow \Delta p_{T,\ell\ell}$$

\Rightarrow interesting subsequent test



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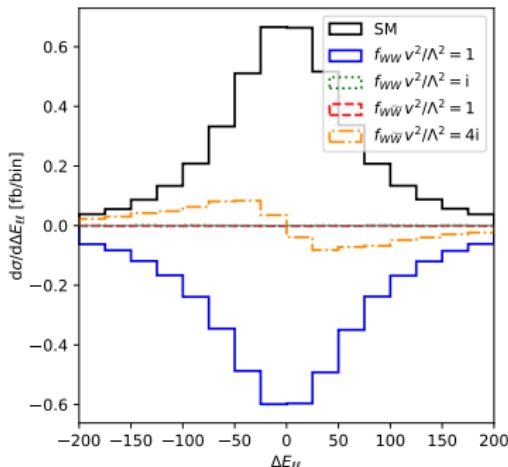
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Higgs decays

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Testing *CP* in $H \rightarrow 4\ell$ decays

- again same 10 scalar products
 - momentum flow limited by m_H
 - reach for *CP*-even operators shit [1612.05261]
 - even making slide is waste of time
- ⇒ what's the point...

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