

LHC

Tilman Plehn

Motivation

Data

Jets

Simulation

Inference

# Machine Learning for the LHC

Tilman Plehn

Universität Heidelberg

Stuttgart 7/2022



# Modern LHC physics

Motivation

Data

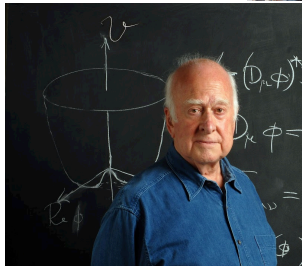
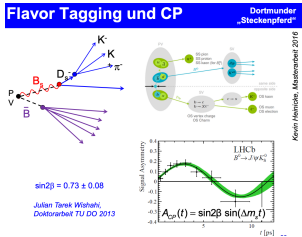
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## Classic motivation

- dark matter
- baryogenesis
- Higgs VEV



# Modern LHC physics

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## LHC physics

- fundamental questions
- huge data set
- complete uncertainty control
- first-principle precision simulations



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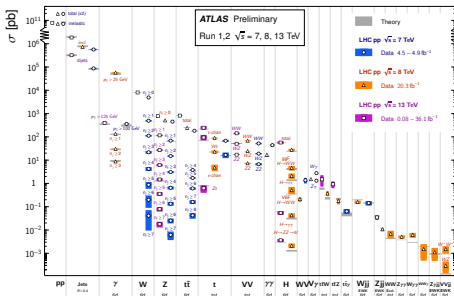
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## Traditional methods

- discover in rates
- unveil little black holes
- find supersymmetry
- travel extra dimensions
- measure couplings



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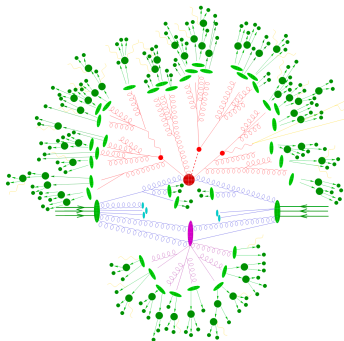
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- start with Lagrangian
- calculate scattering using QFT
- simulate events
- simulate detectors

→ LHC events in virtual worlds



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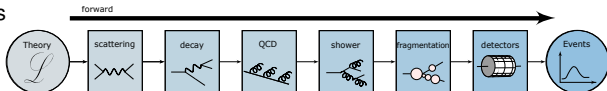
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## New physics searches

- compare simulations and data
  - analyze data systematically
  - understand LHC dataset [SM or BSM]
  - publish useable results
- With a little help from data science...



# LHC data

## Data from ATLAS & CMS

- protons on protons at  $E \approx 13000 \times m_p \rightarrow$  relativistic kinematics
- crossing every 25 ns, 40 MHz, 1.6 MB per event  $\rightarrow$  1 PB/s
- frequency vs size

$$\frac{10 \text{ m}}{3 \times 10^8 \text{ m/s}} \approx 3 \times 10^{-8} \text{ s} = 30 \text{ ns}$$

$\rightarrow$  Big and fast data



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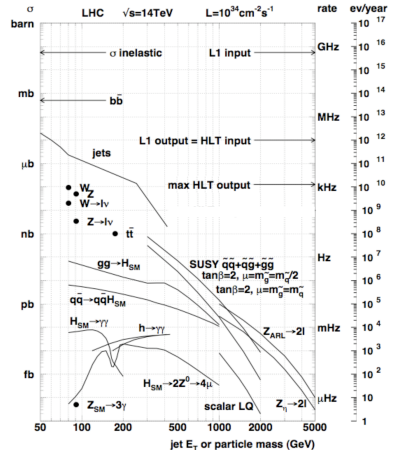
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## Triggering

- $10^{-6}$  suppression physics-loss-less
- L1 hardware 40 MHz  $\rightarrow$  100 kHz
- L2/HL software  $\rightarrow$  3 kHz
- L3 software  $\rightarrow$  200 Hz, 320 MB/s





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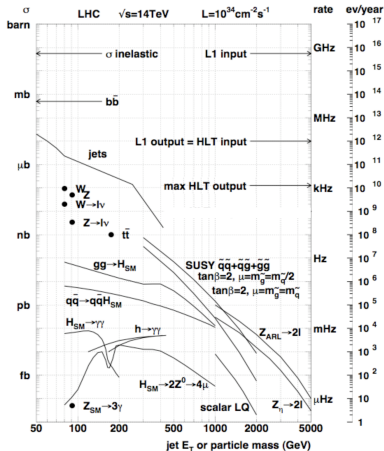
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## Strategies

- classic trigger cuts
- probabilistic prescale trigger
- downsized data scouting



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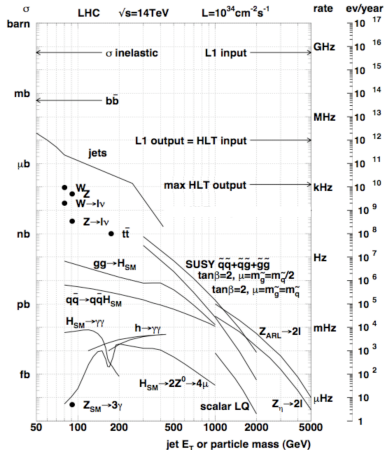
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## ML-questions

- identification of interesting events?
- identification **unexpected events?**
- data compression for analyses?



# Jets

## Partons as QCD jets

- most interactions  $q\bar{q}, gg \rightarrow q\bar{q}, gg$

$$\sigma_{pp \rightarrow jj} \times \mathcal{L} \approx 10^8 \text{ fb} \times \frac{80}{\text{fb}} \approx 10^{10} \text{ events}$$

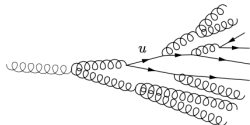
- quarks/gluon visible as jets  
splittings described by QCD  
hadronization and hadron decays in jets

- jets as decay products

$$67\% W \rightarrow jj \quad 70\% Z \rightarrow jj \quad 60\% H \rightarrow jj \quad 67\% t \rightarrow jjj \quad 60\% \tau \rightarrow j \dots$$

- new physics in 'dark jets'
- typical process  $pp \rightarrow t\bar{t}H + \text{jets} \rightarrow bjj \bar{b}jj b\bar{b} + \text{jets}$

→ **Everywhere in LHC physics**



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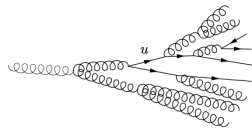
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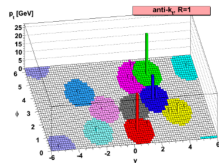
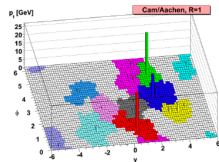
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## Dealing with jets

- 50-200 constituents per jet  
40 pile-up events on top
- calorimeter + tracking = particle-flow
- jet algorithms returning parton 4-momentum
- sub-jet physics new for LHC



# Jets

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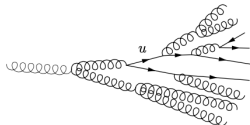
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## ML-questions

- fast particle/parton identification?
- data denoising against jet radiation and pileup?
- combination of calorimeter and tracking resolution?
- combination of low-level and high-level observables?



# ML-tagging: nothing is ever new

## LHC visionaries

- 1991: NN-based quark-gluon tagger

### USING NEURAL NETWORKS TO IDENTIFY JETS

Leif LÖNNBLAD\*, Carsten PETERSON\*\* and Thorsteinn RÖGNVALDSSON\*\*\*

*Department of Theoretical Physics, University of Lund, Sölvegatan 14A, S-22362 Lund, Sweden*

Received 29 June 1990

A neural network method for identifying the ancestor of a hadron jet is presented. The idea is to find an efficient mapping between certain observed hadronic kinematical variables and the quark-gluon identity. This is done with a neuron expansion in terms of a network of sigmoidal functions using a gradient descent procedure, where the errors are back-propagated through the network. With this method we are able to separate gluon from quark jets originating from Monte Carlo generated  $e^+e^-$  events with  $\sim 85\%$  accuracy. The result is independent of the MC model used. This approach for isolating the gluon jet is then used to study the so-called string effect.

In addition, heavy quarks (b and c) in  $e^+e^-$  reactions can be identified on the 50% level by just observing the hadrons. In particular we are able to separate b-quarks with an efficiency and purity, which is comparable with what is expected from vertex detectors. We also speculate on how the neural network method can be used to disentangle different hadronization schemes by compressing the dimensionality of the state space of hadrons.



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- 1994: jet-algorithm  $W$ /top-tagger

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### Searches for new particles using cone and cluster jet algorithms: a comparative study

Michael H. Seymour

*Department of Theoretical Physics, University of Lund, Sölvegatan 14A, S-22362 Lund, Sweden*

Received 18 June 1993; in revised form 16 September 1993

**Abstract.** We discuss the reconstruction of the hadronic decays of heavy particles using jet algorithms. The ability to reconstruct the mass of the decaying particle is compared between a traditional cone-type algorithm and a recently proposed cluster-type algorithm. The specific examples considered are the semileptonic decays of a heavy Higgs boson at  $\sqrt{s}=16$  TeV, and of top quark-antiquark pairs at  $\sqrt{s}=1.8$  TeV. We find that the cluster algorithm offers considerable advantages in the former case, and a slight advantage in the latter. We briefly discuss the effects of calorimeter energy resolution, and show that a typical resolution dilutes these advantages, but does not remove them entirely.

except that the invariant mass of a pair is replaced by the transverse momentum of the softer particle relative to the other.

More recently, this algorithm was extended to collisions with incoming hadrons [5], and a longitudinally-invariant  $k_t$ -clustering algorithm for hadron-hadron collisions was proposed [6]. This algorithm has been compared with the more commonly used cone algorithm from the viewpoints of a parton-shower Monte Carlo program [6, 7], and a fixed-order matrix-element calculation [8], and advantages of the cluster algorithm were reported in both cases. This paper is concerned with a comparison between the algorithms for the task of reconstructing the hadronic decays of heavy particles, which was also studied in a preliminary way in [9].

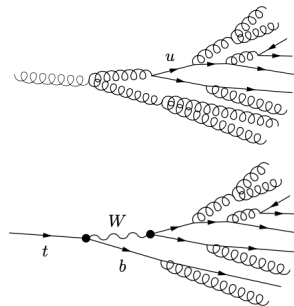
The only as-yet unobserved particles of the minimal Standard Model are the top quark and Higgs boson. The search for, and study of, these particles are among the most important goals of current and planned hadron-



# QCD jet representation

## Jet constituents

- historically  
only hard parton 4-momentum interesting  
parton content from 'tagging'  
QCD tests from theory observables

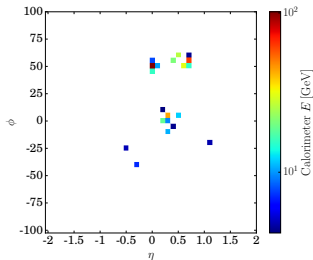




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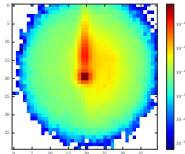
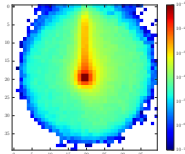
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- ML-excitement phase [since 2015]
  - data-driven jet analyses
  - include as much data as possible
  - avoid intermediate high-level variables
  - calorimeter output as image



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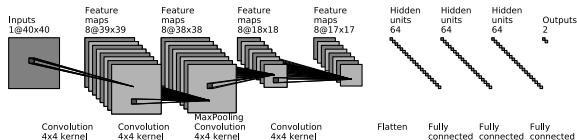
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## Convolutional network

- image recognition standard ML task
- top tagging on 2D jet images
- $40 \times 40$  bins with calorimeter resolution



# Physics representation

## Networks with 4-vector input

- sparsely filled picture: graph CNN
- physics objects from calorimeter and tracker
- distance measure known from e&m

## ML-jet algorithm — combination layer

- input 4-vectors

$$(k_{\mu,i}) = \begin{pmatrix} k_{0,1} & k_{0,2} & \cdots & k_{0,N} \\ k_{1,1} & k_{1,2} & \cdots & k_{1,N} \\ k_{2,1} & k_{2,2} & \cdots & k_{2,N} \\ k_{3,1} & k_{3,2} & \cdots & k_{3,N} \end{pmatrix}$$

- combining them

$$k_{\mu,i} \xrightarrow{\text{CoLa}} \tilde{k}_{\mu,j} = k_{\mu,i} C_{ij} \quad C = \begin{pmatrix} 1 & 0 & \cdots & 0 & C_{1,N+2} & \cdots & C_{1,M} \\ 0 & 1 & & \vdots & C_{2,N+2} & \cdots & C_{2,M} \\ \vdots & \vdots & \ddots & 0 & \vdots & & \vdots \\ 0 & 0 & \cdots & 1 & C_{N,N+2} & \cdots & C_{N,M} \end{pmatrix}$$



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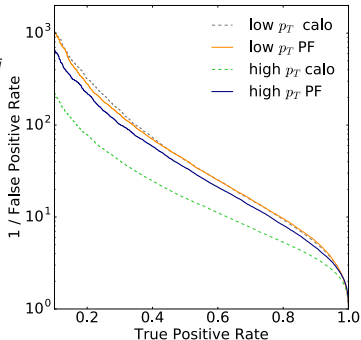
## Invariants — Lorentz layer

- DNN on Lorentz scalars

$$\tilde{k}_j \xrightarrow{\text{LoLa}} \hat{k}_j = \begin{pmatrix} m^2(\tilde{k}_j) \\ p_T(\tilde{k}_j) \\ \vdots \end{pmatrix}$$

→ Learn Minkowski metric

$$g = \text{diag}(0.99 \pm 0.02, \\ -1.01 \pm 0.01, -1.01 \pm 0.02, -0.99 \pm 0.02)$$



# Meet the professionals

## A brief history of achievement

- 2014/15: first jet image papers
- 2017: first (working) ML top tagger
- ML4Jets 2017: What architecture works best?
- ML4Jets 2018: Lots of architectures work

→ Jet classification understood and done

SciPost Physics

Submission

### The Machine Learning Landscape of Top Taggers

G. Kasieczka (ed)<sup>1</sup>, T. Plehn (ed)<sup>2</sup>, A. Butter<sup>2</sup>, K. Cranmer<sup>3</sup>, D. Debnath<sup>4</sup>,  
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**1** Institut für Experimentalphysik, Universität Hamburg, Germany

**2** Institut für Theoretische Physik, Universität Heidelberg, Germany

**3** Center for Cosmology and Particle Physics and Center for Data Science, NYU, USA

**4** NHECT, Dept. of Physics and Astronomy, Rutgers, The State University of NJ, USA

**5** Theoretical Particle Physics and Cosmology, King's College London, United Kingdom

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**13** LPTHE, CNRS & Sorbonne Université, Paris, France

**14** III. Physics Institute A, RWTH Aachen University, Germany

gregor.kasieczka@uni-hamburg.de

plehn@uni-heidelberg.de

April 12, 2019

### Abstract

Based on the established task of identifying boosted, hadronically decaying top quarks, we compare a wide range of modern machine learning approaches. We find that they are extremely powerful and great fun.

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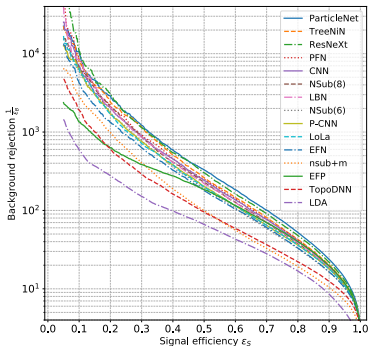
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## Path to LHC reality

- application in analyses?
- beyond top and QCD jets?
- uncertainties?
- resilience in experimental reality?
- beyond fully supervised learning?
- from jets to events?
- analyses only ML will allow us to do?

[etc](#)



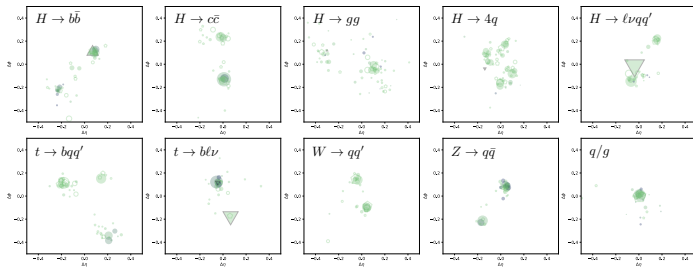


# THE NEED FOR A LARGE DATASET

- **JetClass**: a new large-scale public jet dataset
  - 100M jets for training: ~ two orders of magnitude larger than existing public datasets
  - 10 classes: several unexplored scenarios, e.g.,  $H \rightarrow WW^* \rightarrow 4q$ ,  $H \rightarrow WW^* \rightarrow \ell\nu qq$ , etc.
  - comprehensive information per particle: kinematics, particle ID, track displacement

H. Qu, C. Li, S. Qian,  
arXiv:2202.03772,  
[https://github.com/jet-universe/  
particle-transformer/](https://github.com/jet-universe/particle-transformer/)

Simulated w/ MadGraph +  
Pythia + Delphes



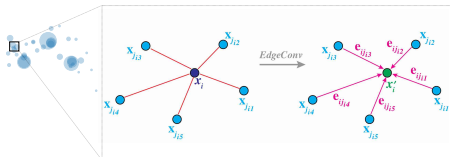
Jet Tagging in the Era of Deep Learning - June 9, 2022 - Huilin Qu (CMS)



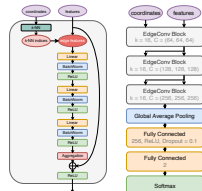
# PARTICLENET

- ParticleNet: jet tagging via particle clouds
  - treating a jet as an **unordered set of particles**, distributed in the  $\eta - \phi$  space
  - graph neural network architecture**, adapted from Dynamic Graph CNN [arXiv:1801.07829]
    - treating a point cloud as a graph: each point is a vertex
      - for each point, a local patch is defined by finding its  $k$ -nearest neighbors
    - designing a permutation-invariant "convolution" function
      - define "edge feature" for each center-neighbor pair:  $e_{ij} = \text{he}(x_i, x_j)$
      - aggregate the edge features in a symmetric way:  $x'_i = \text{mean}_j e_{ij}$

Jet Tagging in the Era of Deep Learning - June 9, 2022 - Huilin Qu (CMS)



ParticleNet architecture



cf. P.T. Komiske, E.M. Metodiev, J. Thaler: *JHEP* **01** (2019) 121;  
V.Mkuni and F.Canello, *Eur. Phys. J. Plus* **135**, 463 (2020); *Mach. Learn. Sci. Tech.* **2** (2021) 3, 035027

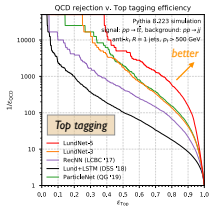


# LUNDNET: PERFORMANCE

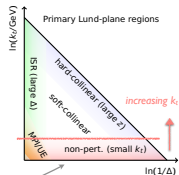
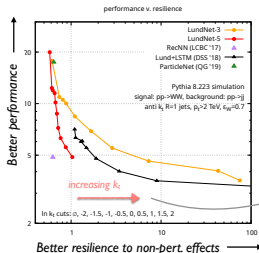
F. Dreyer and H. Qu,  
JHEP 03 (2021) 052

- LundNet achieves very high performance at significant lower computational cost than ParticleNet
  - due to fewer number of neighbors in a binary tree & static graph structure
- Moreover, LundNet provides a systematic way to control the robustness of the tagger
  - the non-perturbative region can be effectively rejected by applying a  $k_T$  cut on the Lund plane

Jet Tagging in the Era of Deep Learning - June 9, 2022 - Huilin Qu (CMS)



	Number of parameters	Training time [ms/sample/epoch]	Inference time [ms/sample]
LundNet	30k	0.472	0.117
ParticleNet	30k	3.488	1.036
Lund+LSTM	67k	0.424	0.131



\* Resilience assessed by applying the model trained on hadron-level samples to parton-level samples and compare the difference



# PARTICLE TRANSFORMER

H. Qu, C. Li, S. Qian,  
arXiv:2202.03772,  
[https://github.com/jet-universe/  
particle-transformer/](https://github.com/jet-universe/particle-transformer/)

- **Attention mechanism and Transformers:** the new state-of-the-art architecture in ML

- Large Language Models: BERT, GPT-3, ...
- Computer Vision: ViT, Swin-T, ...
- AlphaFold2 for protein structure prediction

- **Particle Transformer (ParT)**

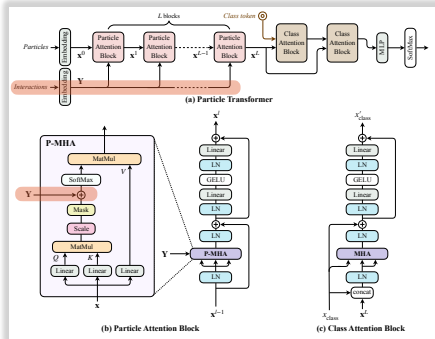
- Transformer-based architecture for jet tagging
- injecting physics-inspired pairwise features to "bias" the dot-product self-attention

$$\text{P-MHA}(Q, K, V) = \text{SoftMax}(QK^T / \sqrt{d_k} + \mathbf{Y})V,$$

"Interaction" features

$$\begin{aligned} \Delta &= \sqrt{(y_a - y_b)^2 + (\phi_a - \phi_b)^2}, \\ k_T &= \min(p_{T,a}, p_{T,b})\Delta, \\ z &= \min(p_{T,a}, p_{T,b}) / (p_{T,a} + p_{T,b}), \\ m^2 &= (E_a + E_b)^2 - \|\mathbf{p}_a + \mathbf{p}_b\|^2, \end{aligned}$$

and more...



# Simulation

## Event generation

- start from Lagrangian

$$\mathcal{L} = \sum_q \bar{\psi}_q (i\gamma^\mu \partial_\mu - m - gG_\mu) \psi_q - \frac{1}{4} G_{\mu\nu} G^{\mu\nu} + \dots - \mu^2 |\phi|^2 - \lambda |\phi|^4$$

- total rate for proton-proton collisions

$$\sigma_{\text{tot}} = \int_0^1 dx_1 \int_0^1 dx_2 \sum_{\text{partons } ij} f_i(x_1) f_j(x_2) \hat{\sigma}_{ij}(x_1 x_2 E^2)$$

- simulation factorized by energy
  - Monte Carlo generation, LO or NLO in QCD
  - production process particle decays  
QCD jet radiation  
QCD showering  
fragmentation/hadronization
- **Theory task**  
Pythia, Madgraph, Sherpa, Herwig

---

### Machine Learning and LHC Event Generation

Anja Butter<sup>1,2</sup>, Tilman Plehn<sup>1</sup>, Steffen Schumann<sup>3</sup> (Editors),  
 Simon Badger<sup>4</sup>, Sascha Caron<sup>5, 6</sup>, Kyle Cranmer<sup>7,8</sup>, Francesco Armando Di Bello<sup>9</sup>,  
 Etienne Dreyer<sup>10</sup>, Stefano Forte<sup>11</sup>, Sanmay Ganguly<sup>12</sup>, Dorival Goncalves<sup>13</sup>, Eilan Gross<sup>10</sup>,  
 Theo Heemel<sup>1</sup>, Gudrun Heinrich<sup>14</sup>, Lukas Heinrich<sup>14</sup>, Alexander Held<sup>15</sup>, Stefan H $\ddot{o}$ che<sup>17</sup>,  
 Jessica N. Howard<sup>18</sup>, Phillip Ilten<sup>19</sup>, Joshua Isaacson<sup>17</sup>, Timo Janßen<sup>2</sup>, Stephen Jones<sup>20</sup>,  
 Marumi Kado<sup>9,21</sup>, Michael Kogan<sup>22</sup>, Gregor Knieczka<sup>23</sup>, Felix Kling<sup>24</sup>, Sabine Kram<sup>25</sup>,  
 Claudius Krause<sup>26</sup>, Frank Krauss<sup>26</sup>, Kevin Kröninger<sup>27</sup>, Rahul Kumar Barman<sup>13</sup>,  
 Michel Luchmann<sup>1</sup>, Vitaly Magerya<sup>14</sup>, Daniel Maître<sup>28</sup>, Bogdan Malaescu<sup>2</sup>,  
 Fabio Maltoni<sup>28,29</sup>, Till Marini<sup>30</sup>, Olivier Mattelaer<sup>28</sup>, Benjamin Nachman<sup>31,32</sup>,  
 Sebastian Pitz<sup>1</sup>, Juan Rojo<sup>33,34</sup>, Matthew Schwartz<sup>25</sup>, David Shih<sup>25</sup>, Frank Siegert<sup>30</sup>,  
 Roy Stegeman<sup>11</sup>, Bob Stienen<sup>1</sup>, Jesse Thaler<sup>27</sup>, Rob Verheyen<sup>28</sup>, Daniel Whiteson<sup>14</sup>,  
 Ramon Winterhalder<sup>28</sup>, and Jure Zupan<sup>19</sup>

### Abstract

First-principle simulations are at the heart of the high-energy physics research program. They link the vast data output of multi-purpose detectors with fundamental theory predictions and interpretation. This review illustrates a wide range of applications of modern machine learning to event generation and simulation-based inference, including conceptual developments driven by the specific requirements of particle physics. New ideas and tools developed at the interface of particle physics and machine learning will improve the speed and precision of forward simulations, handle the complexity of collision data, and enhance inference as an inverse simulation problem.



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---



# Simulation

## Detector simulation

- process-independent response function  
millions of output channels
- full MC simulation Geant4  
lepton/pion/photon shower in matter  
built from detector plans
- fast detector simulation  
Gaussian approximation of response



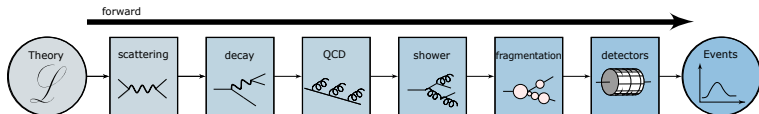
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Gaussian approximation of response

## ML-questions

- fast and precise surrogates for individual steps?
- full phase space coverage?
- full feature mapping?
- variable-dimensional and high-dimensional phase spaces?
- improved data- and theory-driven models?





## Likelihood-based inference

## Unlabeled likelihood ratio [CWoLa]

- Neyman-Pearson lemma: LR optimal discriminator
- likelihood ratio for event samples

$$\text{LR}(x) = \frac{p(x|H_{S+B})}{p(x|H_B)} = \frac{\text{Pois}(n|s+b) \prod_{j=1}^n f_{S+B}(x_j)}{\text{Pois}(n|b) \prod_{j=1}^n f_B(x_j)} = e^{-s} \left( \frac{s+b}{b} \right)^n \frac{\prod_j f_{S+B}(x_j)}{\prod_j f_B(x_j)}$$

- additive log-likelihood ratio

$$\text{LLR}(x) = -s + \sum_j \log \left( 1 + \frac{s f_S(x_j)}{b f_B(x_j)} \right)$$

- LLR from simulation and/or classifier



## Likelihood-based inference

## Unlabeled likelihood ratio [CWoLa]

- Neyman-Pearson lemma: LR optimal discriminator
- problem no signal and background samples to train on  
instead samples  $p_j$  with signal fractions  $f_j$  and background fractions  $1 - f_j$
- phase space densities

$$\begin{pmatrix} p_1(x) \\ p_2(x) \end{pmatrix} = \begin{pmatrix} f_1 & 1 - f_1 \\ f_2 & 1 - f_2 \end{pmatrix} \begin{pmatrix} p_S(x) \\ p_B(x) \end{pmatrix}$$

$$\Leftrightarrow \begin{pmatrix} p_S(x) \\ p_B(x) \end{pmatrix} = \frac{1}{f_1 - f_2} \begin{pmatrix} 1 - f_2 & f_1 - 1 \\ -f_2 & f_1 \end{pmatrix} \begin{pmatrix} p_1(x) \\ p_2(x) \end{pmatrix}$$

- goal: train classifier to extract

$$\frac{p_S(x)}{p_B(x)} = \frac{(1 - f_2)p_1(x) + (f_1 - 1)p_2(x)}{-f_2 p_1(x) + f_1 p_2(x)}$$



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- trick: train classifier for

$$\begin{aligned} \frac{p_1(x)}{p_2(x)} &= \frac{f_1 p_S(x) + (1 - f_1)p_B(x)}{f_2 p_S(x) + (1 - f_2)p_B(x)} = \frac{f_1 \frac{p_S(x)}{p_B(x)} + 1 - f_1}{f_2 \frac{p_S(x)}{p_B(x)} + 1 - f_2} \\ \frac{d}{d(p_S/p_B)} \frac{p_1(x)}{p_2(x)} &= \frac{f_1 \left[ f_2 \frac{p_S(x)}{p_B(x)} + 1 - f_2 \right] - f_2 \left[ f_1 \frac{p_S(x)}{p_B(x)} + 1 - f_1 \right]}{\left[ f_2 \frac{p_S(x)}{p_B(x)} + 1 - f_2 \right]^2} = \frac{f_1 - f_2}{\left[ f_2 \frac{p_S(x)}{p_B(x)} + 1 - f_2 \right]^2} \end{aligned}$$

→ Apply mixed instead of pure classifier



## Likelihood-based inference

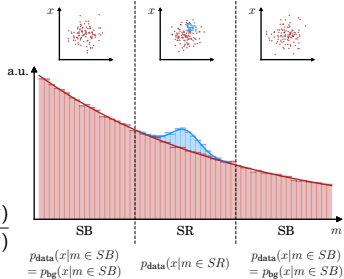
## Improved bump hunts [CWoLa, Anode, Cathode]

- bump hunt in  $m$   
orthogonal information in  $x$

## 1. CWoLa on SB and SR samples

$$\frac{x \sim p_{\text{data}}(x|m \in SR)}{x \sim p_{\text{data}}(x|m \in SB)} \xrightarrow{\text{class}} \frac{p_{S+B}(x)}{p_B(x)} \rightarrow \frac{p_S(x)}{p_B(x)}$$

- but problem with correlations in  $m$  and  $x$



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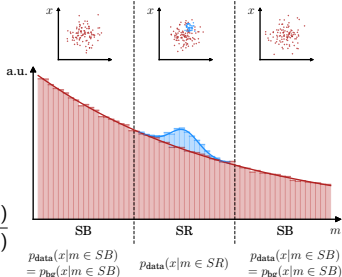
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## 2. density estimation through normalizing flow

$$p_{\text{model}}(x|m \in \text{SB}) \xrightarrow{\text{interpol}} p_{\text{model}}(x|m \in \text{SR})$$

- computable LR in signal regions

$$\text{LR}(x) = \frac{p_{\text{data}}(x|m \in \text{SR})}{p_{\text{model}}(x|m \in \text{SR})} \sim \frac{p_{\text{S+B}}(x)}{p_{\text{B}}(x)}$$



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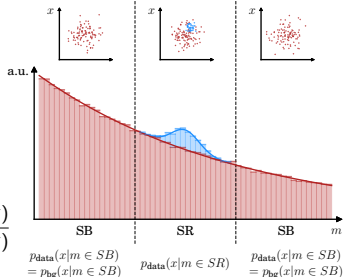
## 3. background generation using normalizing flow

$$p_{\text{model}}(x|m \in \text{SB}) \xrightarrow{\text{sample}} x \sim p_{\text{model}}(x|m \in \text{SR})$$

- classifier on event samples

$$\frac{x \sim p_{\text{model}}(x|m \in \text{SR})}{x \sim p_{\text{model}}(x|m \in \text{SB})} \xrightarrow{\text{class}} \frac{p_{\text{S+B}}(x)}{p_{\text{B}}(x)}$$

→ Guess which works best?



## Likelihood-based inference

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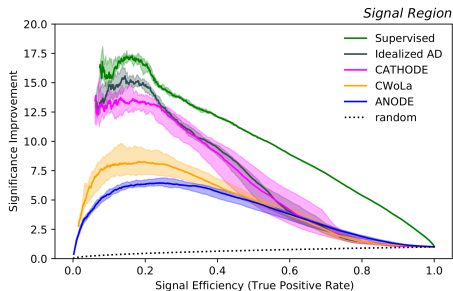
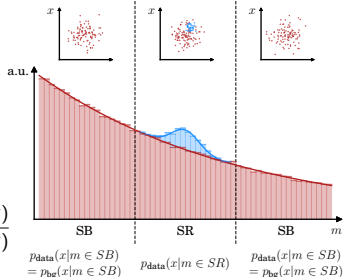
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# ML-Parton densities

## Dirty LHC secret

- proton-proton collisions from parton-parton predictions  $[x = E_{\text{parton}}/E_{\text{proton}}]$

$$\sigma_{\text{tot}} = \int_0^1 dx_1 \int_0^1 dx_2 \sum_{\text{partons } ij} f_i(x_1) f_j(x_2) \hat{\sigma}_{ij}(x_1 x_2 E^2)$$

- DGLAP equation, including factorization scale  $\mu$

$$\frac{df_i(x, \mu)}{d \log \mu^2} = \sum_{\text{partons } j} \int_x^1 \frac{dz}{z} \frac{\alpha_s}{2\pi} P_{i \leftarrow j}(z) f_j\left(\frac{x}{z}, \mu\right) = \frac{\alpha_s}{2\pi} \sum_j (P_{i \leftarrow j} \otimes f_j)(x, \mu)$$

- historic parametrization

$$f_i(x, \mu_0) = a_0 x^{a_1} (1-x)^{a_2} e^{a_3 x + a_4 x^2}$$

→ WTF... → lattice gauge theory?





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## Non-parametric network fit

- parametrizations not useful
- bias problematic

→ NNPDF

### Neural Network Parametrization of Deep-Inelastic Structure Functions

Stefano Forte<sup>a</sup>, Luis Garrido<sup>b</sup>, José I. Latorre<sup>b</sup> and Andrea Piccione<sup>c</sup>

<sup>a</sup>INFN, Sezione di Roma Tre  
Via della Vasca Navale 84, I-00146 Rome, Italy

<sup>b</sup>Departament d'Estructura i Constituents de la Matèria, Universitat de Barcelona,  
Diagonal 647, E-08028 Barcelona, Spain

<sup>c</sup>INFN sezione di Genova and Dipartimento di Fisica, Università di Genova,  
via Dodecaneso 33, I-16146 Genova, Italy

#### Abstract

We construct a parametrization of deep-inelastic structure functions which retains information on experimental errors and correlations, and which does not introduce any theoretical bias while interpolating between existing data points. We generate a Monte Carlo sample of pseudo-data configurations and we train an ensemble of neural networks on them. This effectively provides us with a probability measure in the space of structure functions, within the whole kinematic region where data are available. This measure can then be used to determine the value of the structure function, its error, point-to-point correlations and generally the value and uncertainty of any function of the structure function itself. We apply this technique to the determination of the structure function  $F_2$  of the proton and deuteron, and a precision determination of the isospin combination  $F_2^p - F_2^d$ . We discuss in detail these results, check their stability and accuracy, and make them available in various formats for applications.



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→ NNPDF 6 Summary

We have presented a determination of the probability density in the space of structure functions for the structure function  $F_2$  for proton, deuteron and nonsinglet, as determined from experimental data of the NMC and BCDMS collaborations. Our results, for each of the three structure functions, take the form of a set of 1000 neural nets, each of which gives a determination of  $F_2$  for given  $x$  and  $Q^2$ . The distribution of these functions is a Monte Carlo sampling of the probability density. This Monte Carlo sampling has been obtained by first, producing a sampling of the space of data points based on the available experimental information through a set of Monte Carlo replicas of the original data, and then, training each neural net to one of these replicas.

In practice, all functions are given by a FORTRAN routine which reproduces a feed-forward neural network (described in Section 3) entirely determined by a set of 47 real parameters. Each function is then specified by the set of values for these parameters. Our results are available at the web page <http://sophia.ecm.ub.es/f2neural/>. The full set of FORTRAN routines and parameters can be downloaded from this page. On-line plotting and computation facilities for

### Neural Network Parametrization of Deep-Inelastic Structure Functions

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Diagonal 647, E-08028 Barcelona, Spain

<sup>c</sup>INFN sezione di Genova and Dipartimento di Fisica, Università di Genova,  
via Dodecaneso 33, I-16146 Genova, Italy

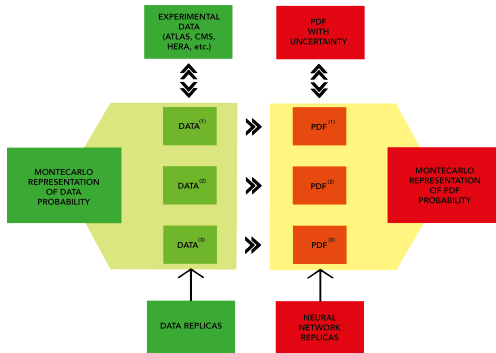
#### Abstract

We construct a parametrization of deep-inelastic structure functions which retains information on experimental errors and correlations, and which does not introduce any theoretical bias while interpolating between existing data points. We generate a Monte Carlo sample of pseudo-data configurations and we train an ensemble of neural networks on them. This effectively provides us with a probability measure in the space of structure functions, within the whole kinematic region where data are available. This measure can then be used to determine the value of the structure function, its error, point-to-point correlations and generally the value and uncertainty of any function of the structure function itself. We apply this technique to the determination of the structure function  $F_2$  of the proton and deuteron, and a precision determination of the iscriptip combination  $F_2^D - F_2^P$ . We discuss in detail these results, check their stability and accuracy, and make them available in various formats for applications.



## THE FUNCTIONAL MONTE CARLO

REPLICA SAMPLE OF FUNCTIONS  $\Leftrightarrow$  PROBABILITY DENSITY IN FUNCTION SPACE  
 KNOWLEDGE OF LIKELIHOOD SHAPE (FUNCTIONAL FORM) NOT NECESSARY



FINAL PDF SET:  $f_i^{(a)}(x, \mu)$ ;

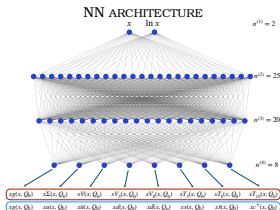
$i = \text{up, antiup, down, antidown, strange, antistrange, charm, gluon}; j = 1, 2, \dots, N_{\text{rep}}$



## THE ML METHODOLOGY

HYPEROPTIMIZED PARAMETERS

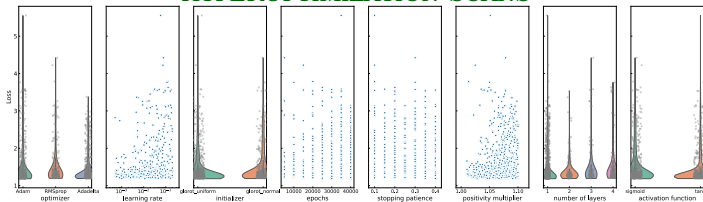
Parameter	NNPDF4.0	$L$ as in Eq. (3.21)	Flavour basis Eq. (3.2)
Architecture	25-20-8	70-50-8	7-26-27-8
Activation function	hyperbolic tangent	hyperbolic tangent	sigmoid
Initializer	glorot_normal	glorot_uniform	glorot_normal
Optimizer	Nadan	Adadelata	Nadan
Clipnorm	$6.0 \times 10^{-6}$	$5.2 \times 10^{-2}$	$2.3 \times 10^{-5}$
Learning rate	$2.6 \times 10^{-3}$	$2.5 \times 10^{-1}$	$2.6 \times 10^{-3}$
Maximum # epochs	$17 \times 10^3$	$45 \times 10^3$	$45 \times 10^3$
Stopping patience	10% of max epochs	12% of max epochs	16% of max epochs
Initial positivity $\Lambda^{(100)}$	185	106	2
Initial integrability $\Lambda^{(101)}$	10	10	10



- HYPEROPT ADAPTS TO EXTERNAL CHOICES (E.G. PARAMETRIZATION BASIS)
- SIMILAR RESULTS CAN BE OBTAINED WITH RATHER DIFFERENT SETTINGS
- $\sim 800$  FREE PARAMETERS



## FITTING THE METHODOLOGY HYPEROPTIMIZATION SCANS



### HYPEROPT PARAMETERS

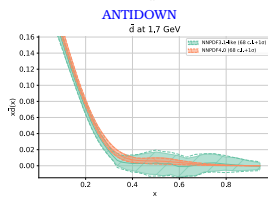
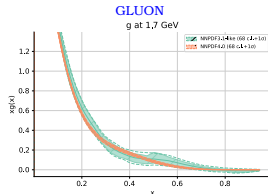
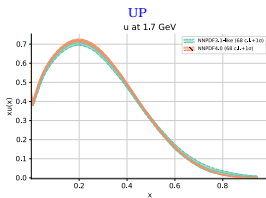
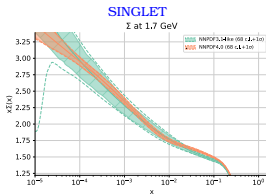
NEURAL NETWORK	FIT OPTIONS
NUMBER OF LAYERS (*)	OPTIMIZER (*)
SIZE OF EACH LAYER	INITIAL LEARNING RATE (*)
DROPOUT	MAXIMUM NUMBER OF EPOCHS (*)
ACTIVATION FUNCTIONS (*)	STOPPING PATIENCE (*)
INITIALIZATION FUNCTIONS (*)	POSITIVITY MULTIPLIER (*)

- **SCAN** PARAMETER SPACE
- **OPTIMIZE** FIGURE OF MERIT: **VALIDATION**  $\chi^2$
- **BAYESIAN** UPDATING



## NNPDF4.0 vs. NNPDF3.1

- FULL BACKWARD COMPATIBILITY
- SUBSTANTIAL REDUCTION IN UNCERTAINTY



# ML-LHC introduction

## Summary

- particle physics has questions
- LHC is big and fast data
- data needs regression and classification
- knowledge comes through theory and simulation
- stochastic data and uncertainty craziness

## Outlook

1. introduction (done)
2. uncertainties and Bayesian networks [TP]
3. generation and inversion [AB]
4. tutorial/hands-on fun [AB]
5. favorite cool ideas [AB,TP]

