

Examining nonextensive statistics in relativistic heavy-ion collisions

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Simon, A., & Wolschin, G. (2018). Examining nonextensive statistics in relativistic heavy-ion collisions. *Physical Review C*, 97(4), 044913.

Setting and Problem

- Model stopping of (proton - antiproton) @ heavy-ion collisions (RHIC, PbPb)
- Use relativistic diffusion model

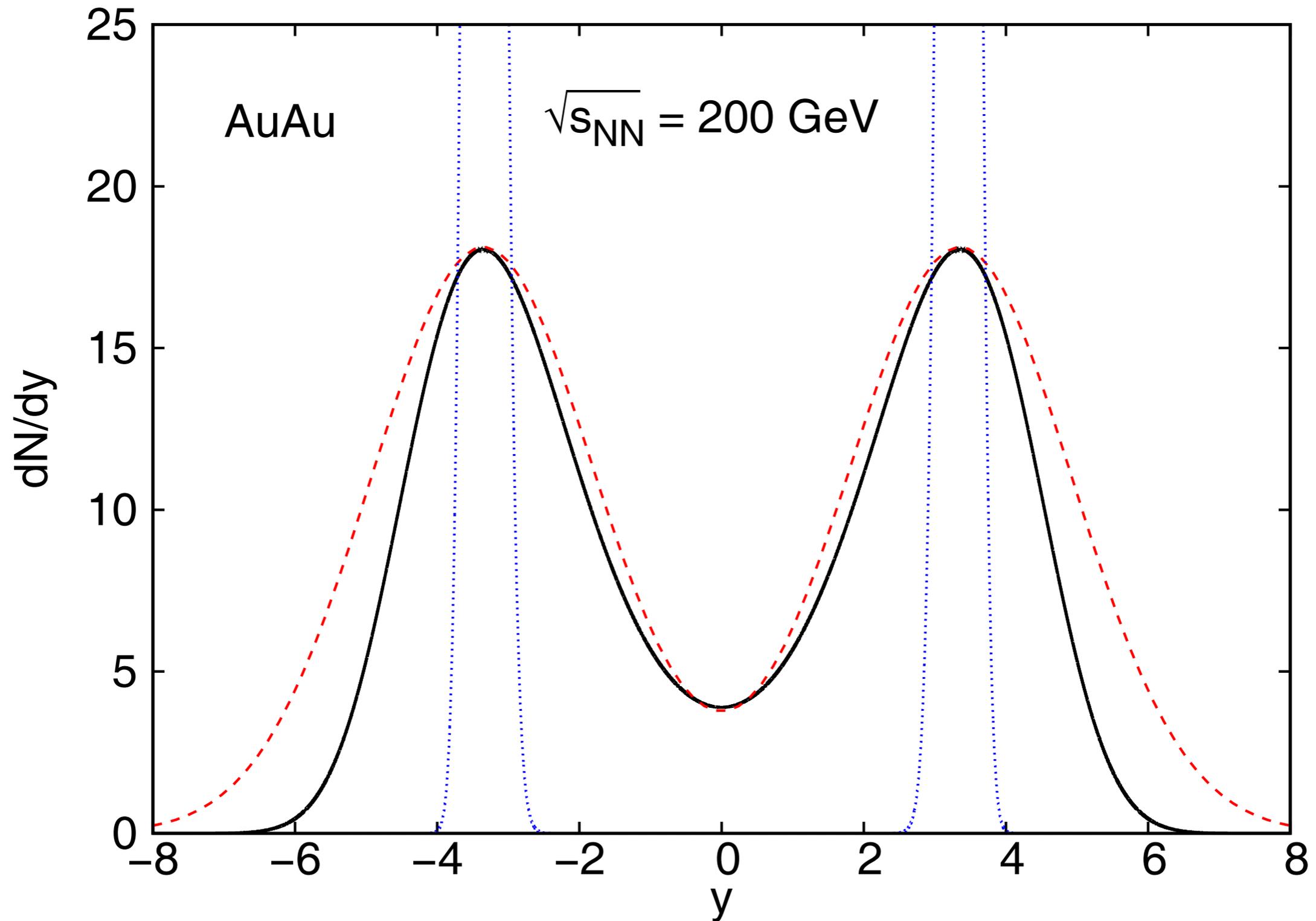
$$\frac{\partial}{\partial t} W(y, t) = - \frac{\partial}{\partial y} [J(y, t) W(y, t)] + \frac{\partial^2}{\partial y^2} [D(y, t) W(y, t)]$$

= drift **= diffusion**

- Know drift coefficient $J(y,t)$ by:
 - interaction time and peak position: amplitude
 - stationary solution: functional form
- Determine diffusion by fluctuation-dissipation theorem

$$J(y) = -\frac{m_{\perp} D}{T} \sinh(y) \quad \begin{array}{|c|} \hline \text{---} \\ \hline \text{---} \\ \hline \end{array}$$

$$J(y) = \frac{y_{\text{eq}} - y}{\tau_y} \quad \begin{array}{|c|} \hline \text{---} \\ \hline \end{array}$$



Account for collective expansion

- ‘Artificially’ through stronger diffusion coefficient, or
- Change underlying equation: non-linear Fokker-Planck equation

$$\frac{\partial}{\partial t} W(y, t)^\mu = -\frac{\partial}{\partial y} [J(y, t) W(y, t)^\mu] + \frac{\partial^2}{\partial y^2} [D(y, t) W(y, t)^\nu]$$

Non-extensive entropy

Boltzmann-Gibbs entropy

$$S_{BG} = -k \sum_i p_i \ln p_i$$

$$S_B(\mathbf{A} + \mathbf{B}) = S_B(\mathbf{A}) + S_B(\mathbf{B})$$

q entropy

$$S_q = \frac{k}{q-1} \left(1 - \sum_i p_i^q \right)$$

$$S_q(\mathbf{A} + \mathbf{B}) = S_q(\mathbf{A}) + S_q(\mathbf{B}) \\ + (1 - q) S_q(\mathbf{A}) S_q(\mathbf{B})$$

$$S_q = \left\langle \ln_q \frac{1}{p_i} \right\rangle$$

Connection to NLFPE?

- Maximize entropy $S_q[p] = \frac{1 - \int du [p(u)]^q}{q - 1}$

$$p_q(y, t) = \frac{\{1 - \beta(t)(1 - q)[y - y_m(t)]^2\}^{1/(1-q)}}{Z_q(t)}$$

- Insert as ansatz into NLFPE and chose linear drift
- Leads to: $q = 1 + \mu - \nu$

Final equation

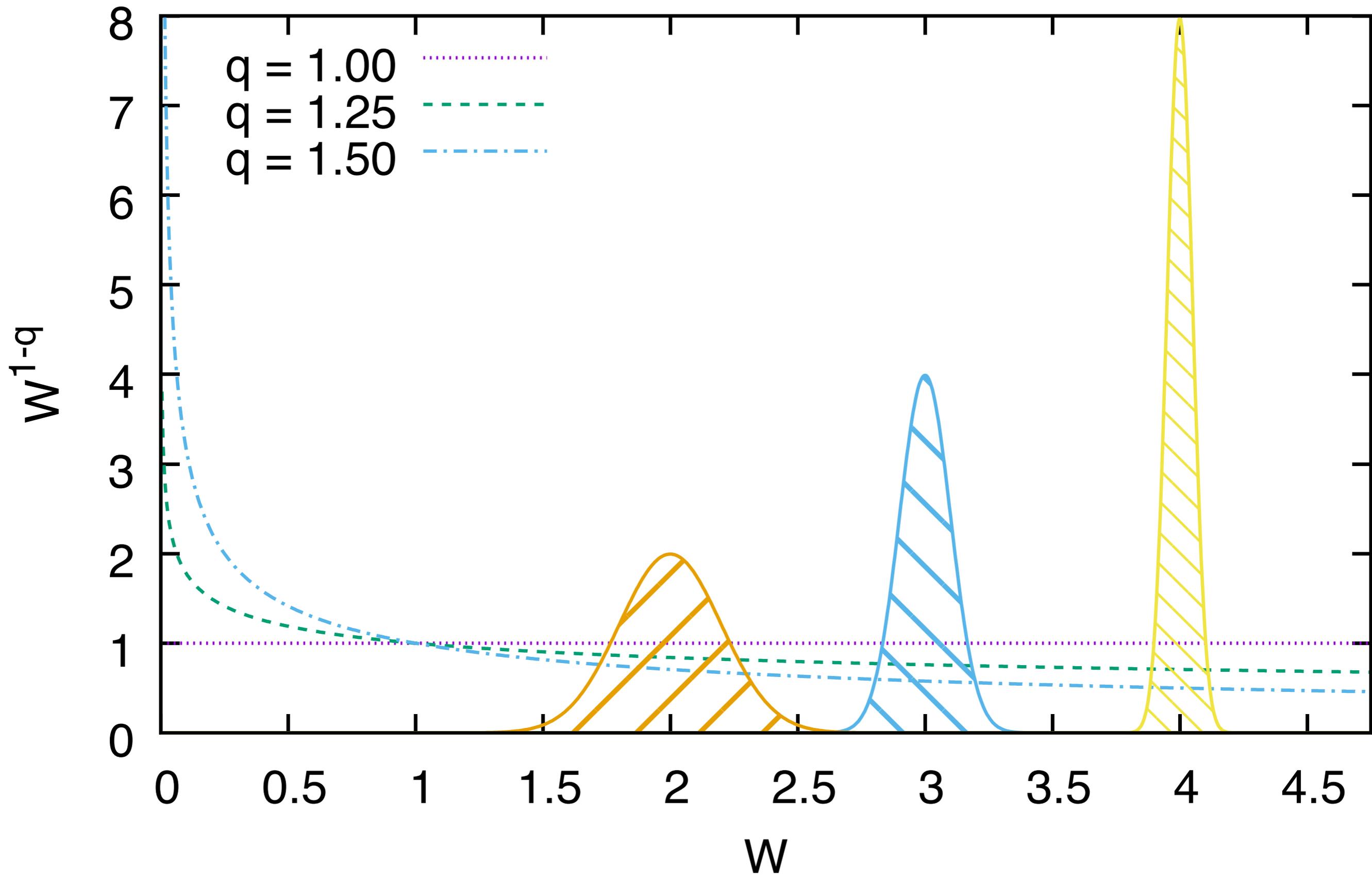
- Rescale time and get

$$\frac{\partial f}{\partial \tau} = \frac{\partial}{\partial y} [\sinh(y) f(y, t)] + \gamma \frac{\partial^2}{\partial y^2} [f(y, t)^{2-q}], \quad \gamma = T/m_{\perp}$$

- Non-linear diffusion coefficient

$$\frac{\partial^2}{\partial y^2} [D W^{2-q}] = \frac{\partial^2}{\partial y^2} [(D W^{1-q}) W]$$

- What changes?



FEM

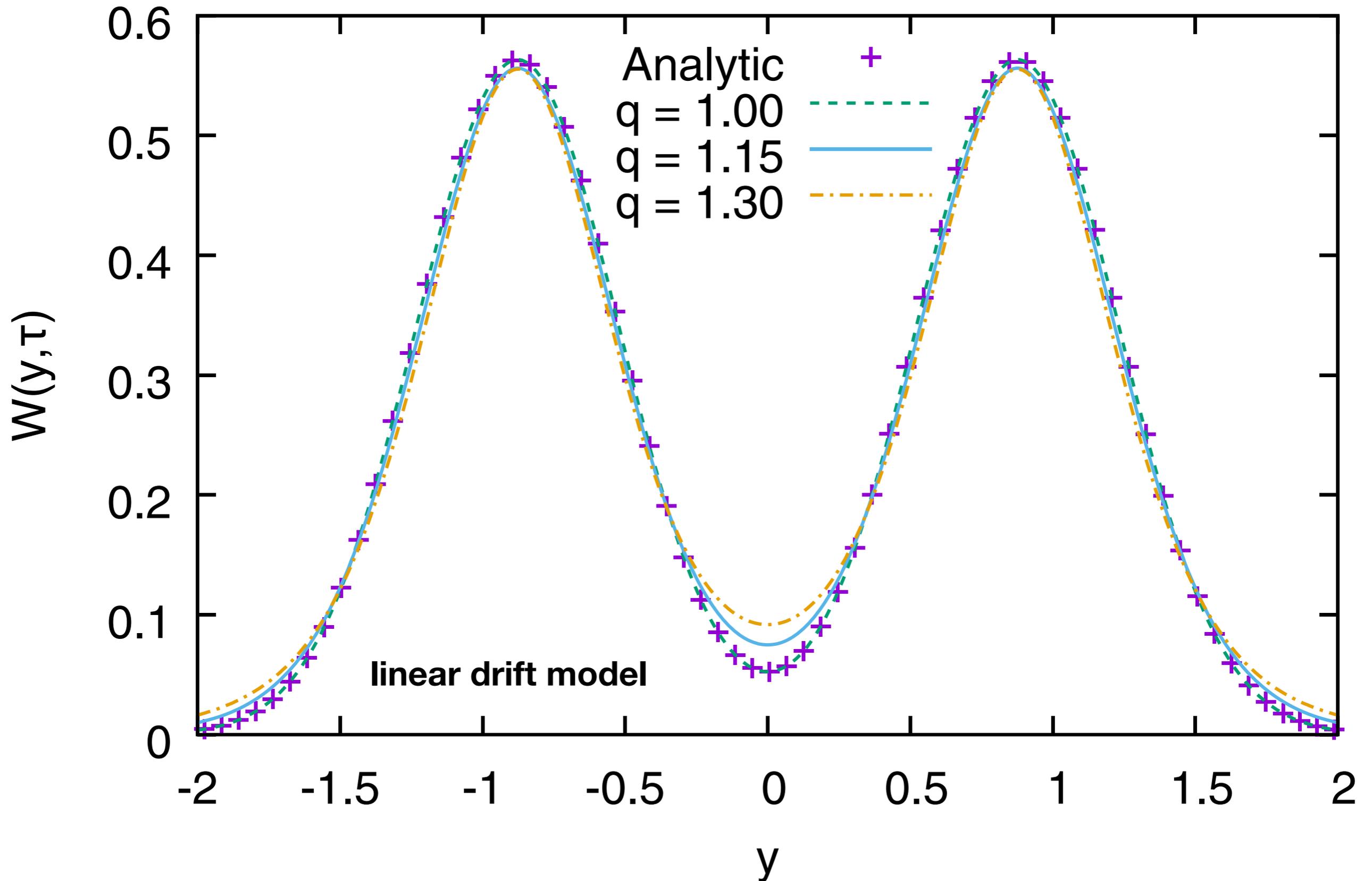
- Weak formulation (integral equation)

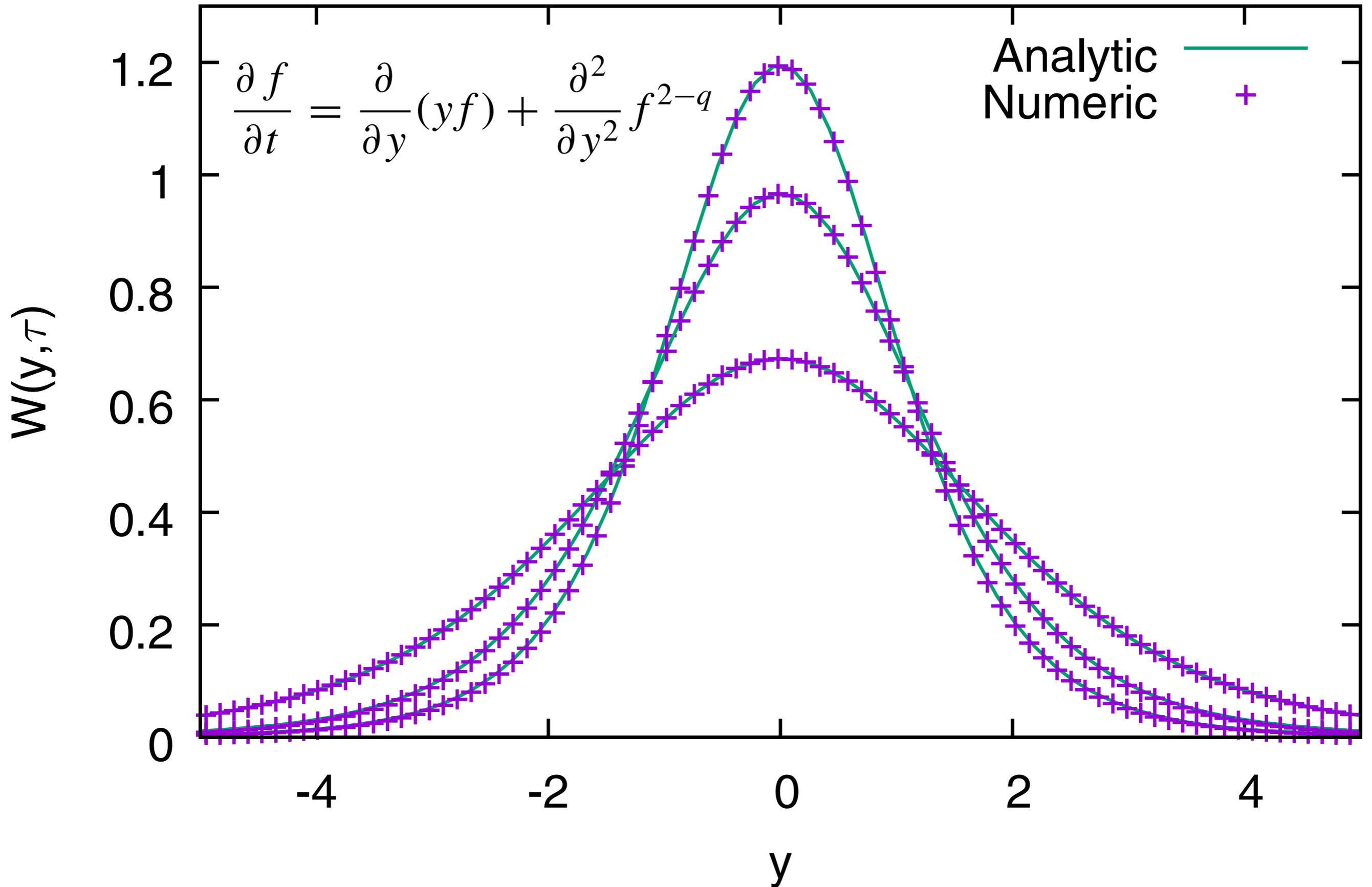
$$\int_{\Omega} dy \left\{ g(y) \frac{\partial}{\partial y} \left[\sinh(y) f(y,t) + \gamma \frac{\partial}{\partial y} f(y,t)^{2-q} \right] \right\}$$
$$\left[g(y) \left\{ \sinh(y) f(y,t) + \gamma \frac{\partial}{\partial y} f(y,t)^{2-q} \right\} \right] \Big|_{\partial\Omega}$$
$$- \int_{\Omega} dy \left\{ \frac{\partial g}{\partial y} \left[\sinh(y) f(y,t) + \gamma \frac{\partial}{\partial y} f(y,t)^{2-q} \right] \right\}$$

- Time integration by Backward-Euler

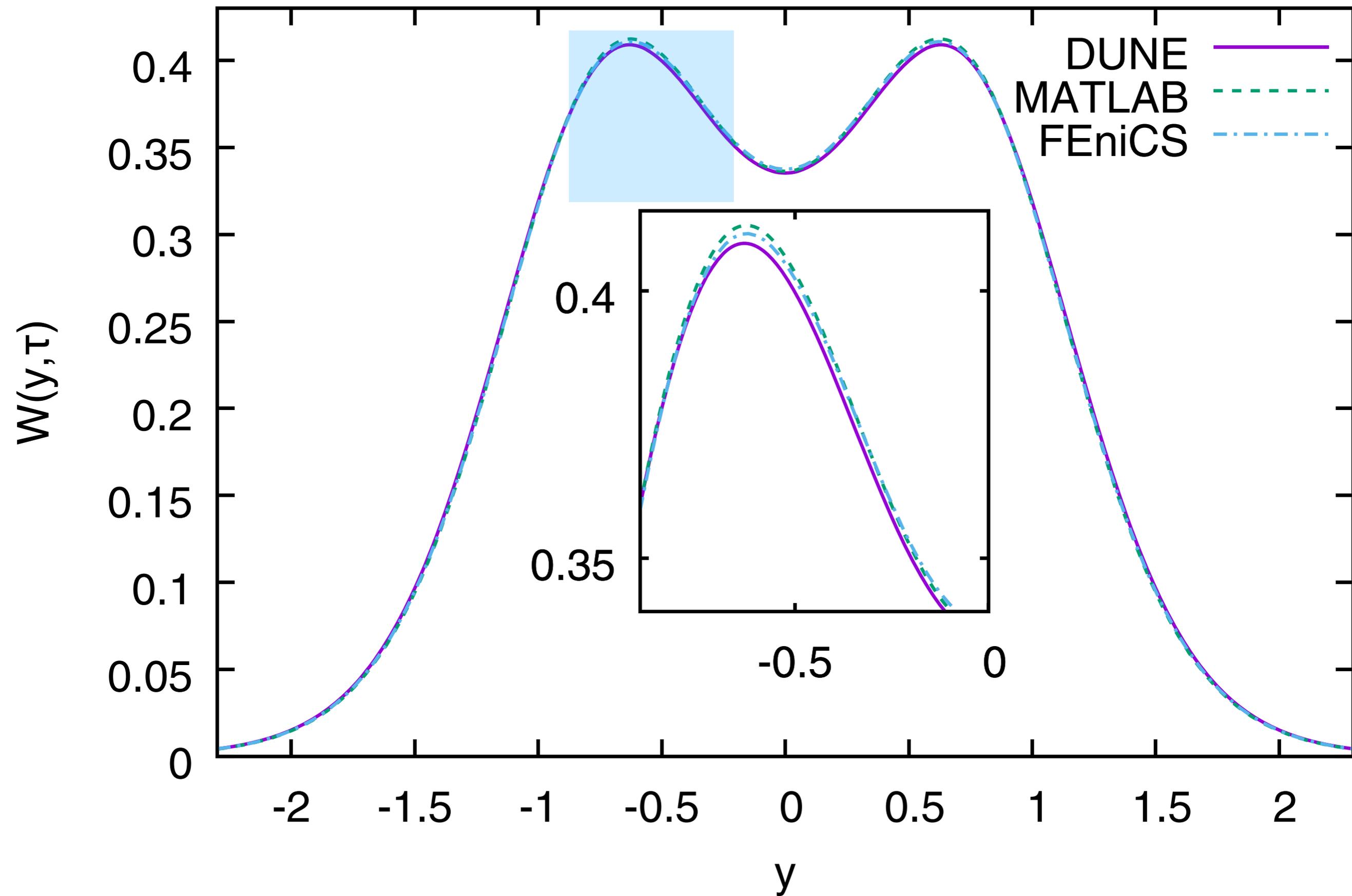
$$\frac{\partial f(t_n)}{\partial t} = \frac{f(t_n) - f(t_{n-1})}{\Delta t} + O(\|\Delta t^2\|)$$

Numerical checks

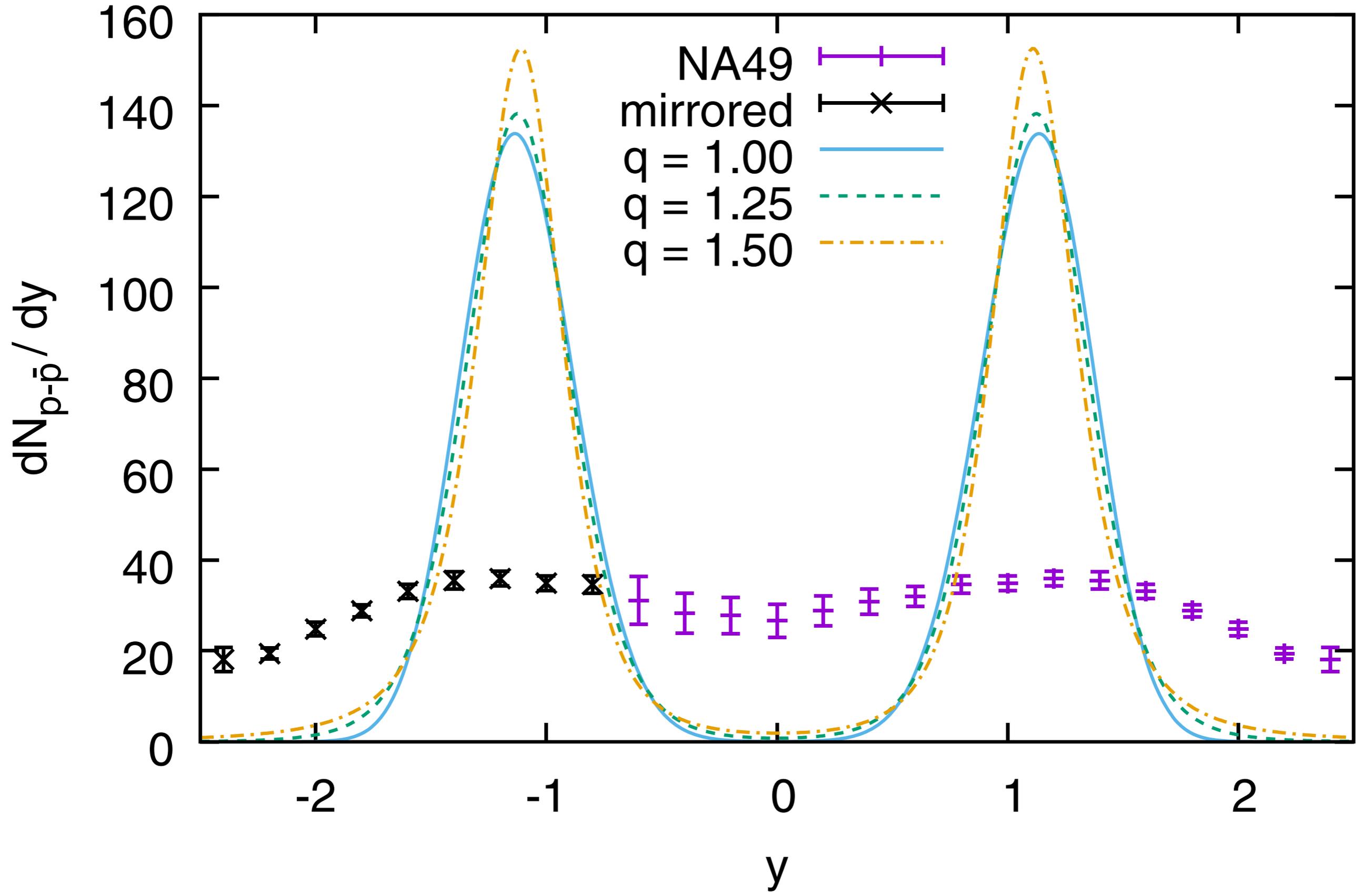




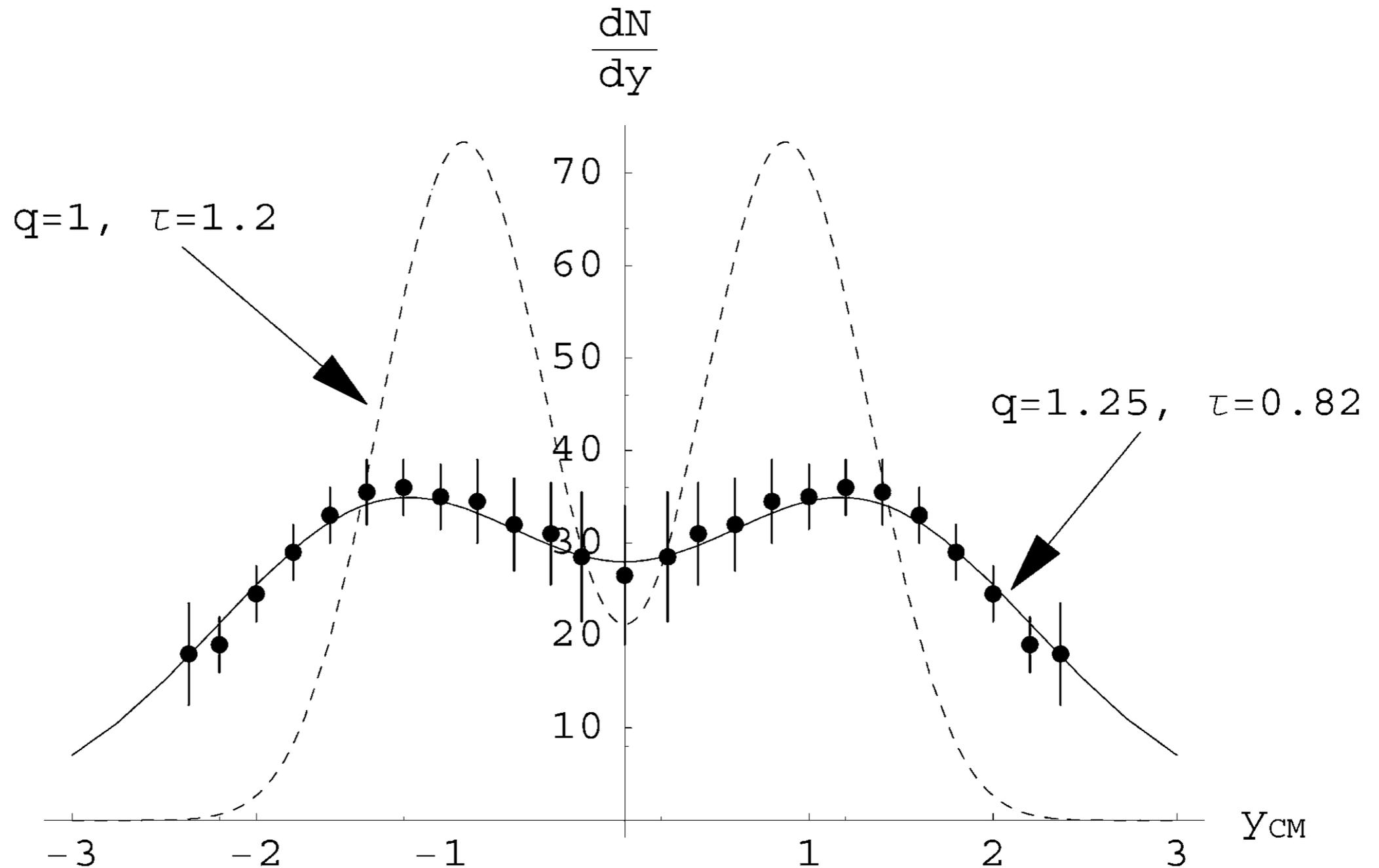
Method comparison

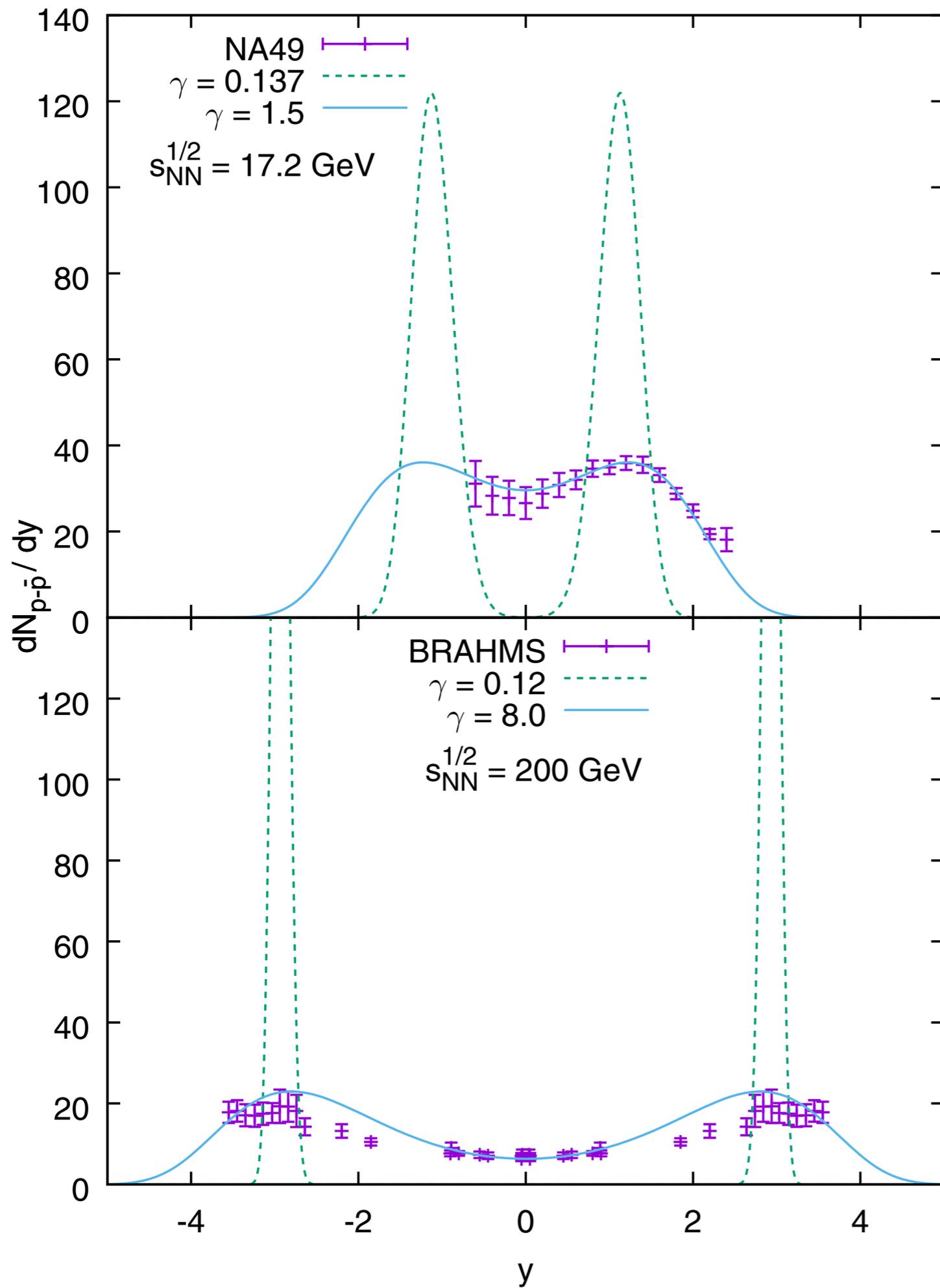


Result



≠ Lavagno's result





- Return to linear diffusion
- Adjust diffusion coefficient manually

Summary

- Measured rapidity spectra broader than anticipated
- Nonextensive statistics supposed to produce collective expansion
- Numerically solved the resulting nonlinear Fokker-Planck equation
- Effect is too weak to account for experimental data